# EXTRA DIMENSION VERSUS SUPERSYMMETRY AT THE LHC

#### S. Nandi

Oklahoma State University & Oklahoma Center for High Energy Physics

"[Work done in collaboration with Kirtiman Ghosh and Durmus Karabacak published in JHEP, and K. Ghosh, D. Karabacak and S. Nandi (work in progress)]

Talk presented at ICHEP 2016, in Chicago, IL, August 3-10, 2016."

### Overview

- Introduction
  - Universal Extra Dimension
- 2 Model: nmUED
  - Constraints on the parameters
- 3 Phenomenology
  - Higgs Phenomenology
  - Collider Phenomenology at the LHC
- 4 Conclusion

### minimal Universal Extra Dimension

In minimal UED  $\Rightarrow$  mUED, 5D space-time is  $M^4 \times S_1/Z_2$ 

• There are KK-excitations of every SM particles

$$\Longrightarrow$$

- Can only be pair produced
- Gives Dark Matter

$$m_n = \frac{n}{R}$$

Radiative corrections

breaks mass
degeneracy.

 $\implies$  Spectra depends on the cut-off scale  $\Lambda$  and  $R^{-1}$ 

- Λ can not be too large, otherwise gauge couplings blow off.
- Spectra of KK-particles is pretty degenerate at each level

With only level-1 KK excitations, similar to compressed SUSY



However, most SUSY breaking scenarios do not give such a compressed spectra

- Pair production of SUSY particles and their decay gives high p<sub>T</sub> jets plus missing transverse energy.
  - Classic signals of SUSY
  - Will distinguish SUSY from UED

However, LHC Higgs data do not agree with mUED at the  $1\sigma$  level

Need to go beyond mUED  $\Longrightarrow$  non-minimal UED  $\Longrightarrow$  nmUED

### Model: nmUED

$$S = S_{bulk} + S_{BLKT}$$

$$\begin{split} \mathcal{S}_{bulk} & = \int \mathrm{d}^4 \times \int_{-L}^L \mathrm{d}y \bigg[ \sum_{\mathcal{A}}^{G,W,B} - \frac{1}{4} \mathcal{A}_{MN} \mathcal{A}^{MN} + \sum_{\Psi}^{Q,U,D,L,E} i \bar{\Psi} \overleftrightarrow{D}_M \Gamma^M \Psi - M_{\Psi} \bar{\Psi} \Psi \bigg] \\ \mathcal{S}_{bdry} & = \int \mathrm{d}^4 \times \int_{-L}^L \mathrm{d}y \bigg( \sum_{\mathcal{A}}^{G,W,B} - \frac{r_{\mathcal{A}}}{4} \mathcal{A}_{\mu\nu} \mathcal{A}^{\mu\nu} + \sum_{\Psi = Q,L} i r_{\Psi} \bar{\Psi}_L D_{\mu} \gamma^{\mu} \Psi_L \\ & + \sum_{\Psi = U,D,E} i r_{\Psi} \bar{\Psi}_R D_{\mu} \gamma^{\mu} \Psi_R \bigg) \times [\delta(y - L) + \delta(y + L)]. \end{split}$$

#### Parameters:

- Coefficients of the BLKT terms
  - $\bullet$   $r_A \Longrightarrow r_g, r_W, r_B$
  - $r_{\psi} \Longrightarrow r_{U}, r_{D}, r_{E}$
- Also involve  $M_{\psi} = \mu \theta(y) \Longrightarrow 5D$  fermion bulk mass.

Masses of the KK-gauge bosons and fermions are determined by solving transcendental equations



# Constraints on the parameters

- $\frac{r_{\Psi}}{L} > \frac{\exp^{-2\mu L} 1}{2\mu L}$  to avoid ghosts and/or tachyons in the fermion sector.
- $r_A/L > -1$  to avoid ghosts and/or tachyons in the gauge sector.
- The bounds on the parameters are also obtained from the low-energy observables. [T. Flacke, K. Kong and S. C. Park, JHEP 1305, 111 (2013)]
  - KK-parity conserving interactions,  $\mathcal{L}_{002n}$
  - $Z^{(2)}$  contribute to 4-fermi interactions
    - $r_A > 0.5L$  for  $\mu L = -0.1$  and fixed  $R^{-1}$
  - However, for  $0 > \mu L > -0.03$  and  $R^{-1} \approx 850$  GeV,  $g_{200}$  will be small and  $M_{Z^{2n}}$  will be heavy enough to escape this bound
- EW precision test and the collider searches are insensitive to small values of  $\mu$ , for example  $\mu L = -0.02$

# Phenomenology

We discuss the implications of nmUED in the context of the Higgs data and multijets plus  $E_T$  searches at the 8 TeV LHC

- Parameters of the model:
  - $\bullet$   $r_{\psi}$ ,  $\Psi = Q$ , U, D, L, E
  - $\bullet$   $r_A$ , A = G, W, B
  - $\bullet$   $\mu \Longrightarrow \mathsf{Bulk}$  fermion mass term
- Take universal boundary parameters for all quarks and leptons  $\implies r_F$
- ullet For gauge sector, we choose  $r_g 
  eq r_W = r_B$
- Also choose  $\mu L = -0.02$

# Mass spectra for level-1 KK-excitations

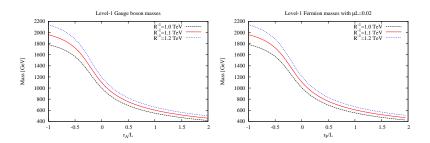


Figure : Level-1 gauge boson (left panel) and fermion (right panel) as a function of r/L for three different values of  $R^{-1}$ . For the level-1 fermion masses in the right panel, we consider  $\mu = -0.02L$ .

- Masses of KK-excitations are very sensitive to the BLKT parameters 

  Large splitting even at the tree level
- Both level-1 KK-fermions and gauge bosons masses increase if we decrease r/L.



- $m_H = 125 \text{ GeV} \Longrightarrow \lambda = 0.129 \text{ in the SM}$
- $\lambda \to 0$  at  $Q = 10^{11}$  GeV  $\Longrightarrow$  Vacuum instability
- $\bullet$  For mUED,  $\lambda$  evolves much faster because of KK-excitations
  - $\lambda \rightarrow 0$  at 4-6  $R^{-1}$
  - For subsequent analysis, we take  $\Lambda = 5R^{-1}$
- Higgs production measured in different channels at the LHC

$$\mu_i = \frac{(\sigma \times BR)_i}{(\sigma \times BR)_i^{SM}}$$

- $H \rightarrow gg$ : Only KK-tower of top contributes
- $H \rightarrow \gamma \gamma$ : KK-towers of top and W both contribute

#### [G. Belanger, B. Dumont, U. Ellwanger, J. F. Gunion and S. Kraml, PRD 88, 075008 (2013)]

$\gamma\gamma$ decay	/ channel	VV decay channel			
$\hat{\mu}^{ggF}$	$\hat{\mu}^{ ext{VBF}}$	$\hat{\mu}^{ggF}$	$\hat{\mu}^{ ext{VBF}}$		
0.98±0.28	1.72±0.59	$0.91 \pm 0.16$	1.01±0.49		

Table: Combined best-fit Higgs signal strengths for different Higgs production and decay modes.

Results for mUED: has only 2 parameters  $\Longrightarrow \Lambda$  and  $R^{-1}$ 

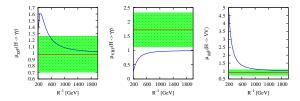


Figure : Higgs signal strengths relative to the SM expectations:  $\mu_{ggF}^{H\to\gamma\gamma}$  (left panel),  $\mu_{VBF}^{H\to\gamma\gamma}$  (middle panel) and  $\mu_{ggF}^{H\to\gamma\nu}$  (right panel), in the context of mUED scenario as a function of  $R^{-1}$ . The combined best-fit values (from Table 1) of the abovementioned Higgs signal strengths are also presented. Can not fit data at  $1\sigma$  level (see middle panel)

# In addition to $R^{-1}$ and $\Lambda$ , we have several additional parameters coming from the BLK terms

- $\bullet$   $r_{\psi}$ ,  $\Psi = Q$ , U, D, L, E
- $\bullet$   $r_A$ , A = G, W, B
- $\mu \Longrightarrow$  Bulk fermion mass term
- We choose cut-off scale  $\Lambda = 5R^{-1}$
- Universal r<sub>F</sub>
- But, for gauge sector, we choose  $r_g \neq r_W = r_B$
- Also choose  $\mu L = -0.02$

We have scanned the parameter space  $(r_W, r_F)$  to fit the Higgs data in different channels for 2 values of  $R^{-1}$  (1.2 and 1.3 TeV)

## Continued.

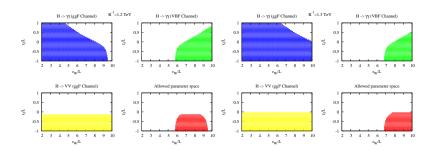


Figure : Scattered points in  $r_W/L$ - $r_F/L$  plane which are consistent with the combined best-fit results of  $\mu_{ggF}^{H\to\gamma\gamma}$ ,  $\mu_{VBF}^{H\to\gamma\gamma}$ ,  $\mu_{ggF}^{H\to\gamma\gamma}$  and all three together for  $R^{-1}=1.2$  TeV (left panel) and  $R^{-1}=1.3$  TeV (right panel).

All combined best-fit Higgs data is consistent with nmUED for large  $r_W$  and negative  $r_F$ 



# Collider Phenomenology at the LHC

- Multijets at high  $p_T$  plus large  $\not \! E_T$  signal
- Use the LHC data for SUSY search and if we can reproduce the exact limits for nmUED using the allowed parameter space
- We use 8 TeV LHC data with  $L=20.3~{\rm fb}^{-1}$
- SUSY limits: For  $m_{\tilde{q}}=m_{\tilde{g}}<1.7$  TeV is excluded from jets +  $\not\!\!E_T$  channel.
- Two questions:
  - Can we reproduce any SUSY signals from the level-1 KK particles using the allowed parameter space of nmUED?
  - What limits we can put on  $q^{(1)}$  and  $g^{(1)}$  masses in nmUED ?

# nmUED Benchmark point for multijets analysis

Benchmark Point (BP)							
$R^{-1}$	ΛR	$\mu$ L	$r_g/L$	r <sub>F</sub> /L	$r_W/L$		
1.2 TeV	5	-0.02	-0.05	-0.42	7.4		
Masses in GeV							
$m_{Q^{(1)}}$	$m_{L^{(1)}}$	$m_{G^{(1)}}$	$m_{W^{(1)\pm}}$	$m_{Z^{(1)}}$	$m_{\gamma^{(1)}}$		
1800	1800	1265	275	275	260		

- Produce  $q^1q^1$ ,  $g^1g^1$ ,  $q^1g^1$
- Decay these using the spectra for the benchmark point
- Apply the same cuts as ATLAS multijets SUSY searches

# Results: ATLAS cuts

Cuts	A (2-			-jets)	C (4	jets <u>)</u>	D		E (6-jets)	_
	L	М	M		M	ı	(5-jets)	L	M	1
$\not\models_T > [GeV]$	160									
$ ho_T^{j_1} > [GeV]$	130									
$p_T^{j_2} > [GeV]$	60									
$p_T^{j_3} > [GeV]$	-		60		60		60	60		
$p_T^{j_4} > [GeV]$	-		-		60		60	60		
$p_T^{j_5} > [GeV]$	-		-		-		60	60		
$p_T^{j_6} > [GeV]$	-							60		
$\Delta \phi(j_i, \vec{E}_T)_{min} >$	0.4 {i=1,2,3 if $p_T^{j3} > 40 \text{ GeV}$ }			0.4 {i=1,2,3}, 0.2 $p_T^{j_i} > 40 \text{ GeV}$						
$\not\!\!E_T/M_{eff}(N_j) >$	0.2	-	0.3	0.4	0.25	0.25	0.2	0.15	0.2	0.25
m <sub>eff</sub> (incl.) [TeV]	1.0	1.6	1.8	2.2	1.2	2.2	1.6	1.0	1.2	1.5
$\sigma_{BSM}$ [fb]	66.07	2.52	0.73	0.33	4.00	0.12	0.77	4.55	1.41	0.41

Table : Cuts used by the ATLAS collaboration to define the signal regions.  $\Delta \phi(jet, \vec{\xi}_T)$  is the azimuthal separations between  $\vec{E}_T'$  and the reconstructed jets.  $m_{eff}(N_j)$  is defined to be the scalar sum of the transverse momenta of the leading N jets together with  $\vec{\xi}_T$ . However, for  $m_{eff}^{incl.}$ , the sum goes over all jets with  $\rho_T > 40$  GeV. Last column corresponds to the 95% C.L. observed upper limits on the non-SM contributions  $\sigma_{BSM}$ .

# Results: Cut-flow table.

			LIED	
Process	Supersyn	nmUED		
	$\tilde{g}\tilde{g}$ one	$g^{(1)}g^{(1)}$ one-step		
Point	$m_{\tilde{g}} = 12$	$m_{g(1)} = 1265 \text{ GeV}$		
	$m_{\tilde{\chi}_1^{\pm}} = 8$	865 GeV	$m_{\tilde{W}(1)\pm} = 865 \text{ GeV}$	
	$m_{\tilde{\chi}_1^0} = 4$	$m_{ ilde{\gamma}(1)} = 465 \; { m GeV}$		
Cuts		n %		
(E-tight)	ATLAS	Our Simulation	Our Simulation	
	Appendix-C of [?]			
0-lepton	63.5	66.1	57.3	
<i>E/<sub>T</sub></i> > 160 GeV	55.6	57.6	54.7	
$p_T^{j_1} > 130 \text{ GeV}$	55.6	57.5	54.7	
$p_T^{j_2} > 60 \text{ GeV}$	55.6	57.5	54.6	
$p_T^{j_3} > 60 \text{ GeV}$	55.4	57.3	51.8	
$p_T^{j_4} > 60 \text{ GeV}$	53.4	55.2	41.3	
$p_T^{j_5} > 60 \text{ GeV}$	46.3	47.1	27.4	
$ ho_T^{j_6} > 60 \; { m GeV}$	31.7	31.1	15.0	
$\Delta \phi(j_i, E/T), i = 1, 2, 3$	26.5	26.1	12.2	
$\Delta\phi(j, E_T), p_T^j > 40 \text{ GeV}$	21.3	21.6	9.7	
$E_T/m_{eff}(N_i) > 0.25$	12.0	12.7	4.7	
$m_{\rm eff}({\it incl.}) > 1.5~{ m TeV}$	7.9	8.3	4.5	

- Our simulation agrees very well with the ATLAS simulations



### Results

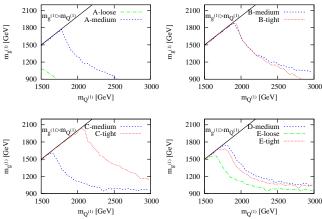


Figure : The exclusion limits on  $m_{Q^1}$ - $m_{g^1}$  plane from 8 TeV 20.3 inverse femtobarn integrated luminosity ATLAS data for different ATLAS defined signal regions. We have assumed fixed mass for the level-1 electroweak KK gauge bosons ( $m_{Q^1}^{(1)\pm} = m_{Z(1)} = 275$  GeV and  $m_{W^1} \gamma^{(1)} = 260$  GeV).

For  $m_{q^1} = m_{g^1}$ , the limit is 2.1 TeV.



### Conclusion

- Though mUED signals are very different from the SUSY signals at the LHC, nmUED signals are not.
- With switable choice of BLK terms, we can reproduce any multijets + ∉<sub>T</sub> signal given by SUSY in nmUED.
- If enhancement in  $H \to \gamma \gamma$  persists in LHC Run 2, it is allowed in nmUED, but not in mUED.
- nmUED with the assumption  $m_{q^1}=m_{g^1}$ , 8 TeV LHC limit is  $\sim 2.1$  TeV.
- Since nmUED can reproduce any SUSY signals, the production of level-2 KK excitations will be the key to distinguish between extra-dimension and supersymmetry if any signal is seen.