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How to Resolve the Proton Radius Puzzle?

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Introduction: The proton radius puzzle

Form Factors

- Matrix element of EM current between nucleon states give rise to two form factors ($q = p_f - p_i$)

$$\langle N(p_f) | \sum_q e_q \bar{q} \gamma^\mu q | N(p_i) \rangle = \bar{u}(p_f) \left[\gamma^\mu F_1(q^2) + \frac{i\sigma_{\mu\nu}}{2m} F_2(q^2) q^\nu \right] u(p_i)$$

- Sachs electric and magnetic form factors

$$G_E(q^2) = F_1(q^2) + \frac{q^2}{4m_p^2} F_2(q^2) \quad G_M(q^2) = F_1(q^2) + F_2(q^2)$$

$$G_E^p(0) = 1$$

$$G_M^p(0) = \mu_p \approx 2.793$$

- The slope of G_E^p

$$\langle r^2 \rangle_E^p = 6 \left. \frac{dG_E^p}{dq^2} \right|_{q^2=0}$$

determines the charge radius $r_E^p \equiv \sqrt{\langle r^2 \rangle_E^p}$

- The proton *magnetic* radius

$$\langle r^2 \rangle_M^p = \frac{6}{G_M^p(0)} \left. \frac{dG_M^p(q^2)}{dq^2} \right|_{q^2=0}$$

Charge radius from atomic physics

$$\langle p(p_f) | \sum_q e_q \bar{q} \gamma^\mu q | p(p_i) \rangle = \bar{u}(p_f) \left[\gamma^\mu F_1^P(q^2) + \frac{i\sigma_{\mu\nu}}{2m} F_2^P(q^2) q^\nu \right] u(p_i)$$

- For a point particle amplitude for $p + \ell \rightarrow p + \ell$

$$\mathcal{M} \propto \frac{1}{q^2} \Rightarrow U(r) = -\frac{Z\alpha}{r}$$

- Including q^2 corrections from proton structure

$$\mathcal{M} \propto \frac{1}{q^2} q^2 = 1 \Rightarrow U(r) = \frac{4\pi Z\alpha}{6} \delta^3(r) (r_E^p)^2$$

- Proton structure corrections $\left(m_r = m_\ell m_p / (m_\ell + m_p) \approx m_\ell \right)$

$$\Delta E_{r_E^p} = \frac{2(Z\alpha)^4}{3n^3} m_r^3 (r_E^p)^2 \delta_{\ell 0}$$

- Muonic hydrogen can give the best measurement of r_E^p !

Charge radius from atomic physics



- Lamb shift in muonic hydrogen [Pohl et al. Nature **466**, 213 (2010)]
 $r_E^p = 0.84184(67) \text{ fm}$
more recently $r_E^p = 0.84087(39) \text{ fm}$ [Antognini et al. Science **339**, 417 (2013)]
- CODATA value [Mohr et al. RMP **80**, 633 (2008)]
 $r_E^p = 0.87680(690) \text{ fm}$
more recently $r_E^p = 0.87750(510) \text{ fm}$ [Mohr et al. RMP **84**, 1527 (2012)]
extracted mainly from (electronic) hydrogen
- **5σ discrepancy!**
- This is the proton radius puzzle

What could be the reason for the discrepancy?

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 - 2) Hadronic Uncertainty? (Part 2 of this talk)

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 - 3) New Physics?

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 - 1) Problem with the electronic extraction? (Part 1 of this talk)
 - 2) Hadronic Uncertainty? (Part 2 of this talk)
 - 3) New Physics?
- Disclaimer: I will focus on published work I am involved in
- For a review of other approaches see
[Carlson, Prog. Part. Nucl. Phys. **82**, 59 (2015)]

Part 1: Proton radii from scattering

What does the PDG say?

K. Nakamura et al. (Particle Data Group), J. Phys. G 37, 075021 (2010)

p CHARGE RADIUS

This is the rms charge radius, $\sqrt{\langle r^2 \rangle}$.

VALUE (fm)	DOCUMENT ID	TECN	COMMENT
0.8768 ± 0.0069	MOHR	08	RVUE 2006 CODATA value
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.897 ± 0.018	BLUNDEN	05	SICK 03 + 2 γ correction
0.8750 ± 0.0068	MOHR	05	RVUE 2002 CODATA value
0.895 ± 0.010 ± 0.013	SICK	03	$ep \rightarrow ep$ reanalysis
0.830 ± 0.040 ± 0.040	²⁴ ESCHRICH	01	$ep \rightarrow ep$
0.883 ± 0.014	MELNIKOV	00	1S Lamb Shift in H
0.880 ± 0.015	ROSENFELDR.	00	ep + Coul. corrections
0.847 ± 0.008	MERGELL	96	ep + disp. relations

Citation: K. Nakamura et al. (Particle Data Group), JPG 37, 075021 (2010) (URL: <http://pdg.lbl.gov>)

0.877 ± 0.024	WONG	94	reanalysis of Mainz ep data
0.865 ± 0.020	MCCORD	91	$ep \rightarrow ep$
0.862 ± 0.012	SIMON	80	$ep \rightarrow ep$
0.880 ± 0.030	BORKOWSKI	74	$ep \rightarrow ep$
0.810 ± 0.020	AKIMOV	72	$ep \rightarrow ep$
0.800 ± 0.025	FREREJACQ...	66	$ep \rightarrow ep$ (CH ₂ tgt.)
0.805 ± 0.011	HAND	63	$ep \rightarrow ep$

²⁴ESCHRICH 01 actually gives $\langle r^2 \rangle = (0.69 \pm 0.06 \pm 0.06) \text{ fm}^2$.

Form Factors: What we don't know

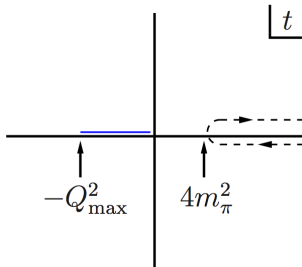
- The form factors are non-perturbative objects.
- **Nobody** knows the *exact* functional form of G_E^p and G_M^p
- They don't have to have a dipole/polynomial/spline or any other functional form
- Including such models can bias your extraction of r_E^p and r_M^p

Form Factors: What we do know

- Analytic properties of $G_E^p(t)$ and $G_M^p(t)$ are known
- They are analytic outside a cut $t \in [4m_\pi^2, \infty]$

[Federbush, Goldberger, Treiman, Phys. Rev. **112**, 642 (1958)]

- $e - p$ scattering data is in $t < 0$ region



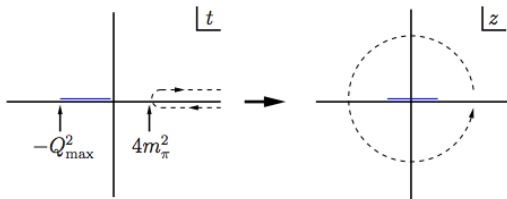
z expansion

- z expansion:

We can map the domain of analyticity onto the unit circle

$$z(t, t_{\text{cut}}, t_0) = \frac{\sqrt{t_{\text{cut}} - t} - \sqrt{t_{\text{cut}} - t_0}}{\sqrt{t_{\text{cut}} - t} + \sqrt{t_{\text{cut}} - t_0}}$$

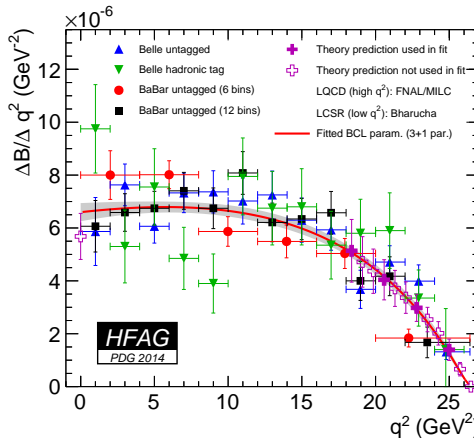
where $t_{\text{cut}} = 4m_\pi^2$, $z(t_0, t_{\text{cut}}, t_0) = 0$



- Expand $G_{E,M}^p$ in a Taylor series in z : $G_{E,M}^p(q^2) = \sum_{k=0}^{\infty} a_k z(q^2)^k$

z expansion

- For meson form factors, z expansion **is** the method
- E.g. $|V_{ub}|$ from exclusive $B \rightarrow \pi \ell \bar{\nu}$



[Heavy Flavor Averaging Group, arXiv:1412.7515]

Extracting r_E^p using the z expansion

- First use of the z expansion to extract r_E^p
[Hill, GP PRD **82** 113005 (2010)]

- Proton: $Q^2 < 0.5 \text{ GeV}^2$

$$r_E^p = 0.870 \pm 0.023 \pm 0.012 \text{ fm}$$

- Proton and neutron data

$$r_E^p = 0.880^{+0.017}_{-0.020} \pm 0.007 \text{ fm}$$

- Proton, neutron and $\pi\pi$ data

$$r_E^p = 0.871 \pm 0.009 \pm 0.002 \pm 0.002 \text{ fm}$$

- Lamb shift in muonic hydrogen [Pohl et al. Nature **466**, 213 (2010)]

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Extracting r_M^p using the z expansion

- z expansion study

[Zachary Epstein, GP, Joydeep Roy PRD **90**, 074027 (2014)]

- Proton data : $r_M^p = 0.91_{-0.06}^{+0.03} \pm 0.02$ fm
- Proton and neutron data: $r_M^p = 0.87_{-0.05}^{+0.04} \pm 0.01$ fm
- Proton, neutron and $\pi\pi$ data: $r_M^p = 0.87 \pm 0.02$ fm

- PDG 2014:

- $r_M^p = 0.777 \pm 0.017$ fm [Bernauer et al. PRL **105**, 242001 (2010)]
- $r_M^p = 0.876 \pm 0.019$ fm [Borisyuk NPA **843**, 59 (2010)]
- $r_M^p = 0.854 \pm 0.005$ fm [Belushkin et al. PRC **75**, 035202 (2007)]

- Other non-PDG values:

- $r_M^p = 0.855 \pm 0.035$ fm [Sick Prog.Part.Nucl.Phys. **55**, 440 (2005)]
- $r_M^p = 0.86_{-0.03}^{+0.02}$ fm [Lorenz et al. EPJA **48**, 151 (2012)]
- $r_M^p = 0.78 \pm 0.08$ fm [Karshenboim PRD **90** 053013 (2014) 5]

Latest z expansion fit

- Most recent study using the z expansion

[Gabriel Lee, Arrington, Hill, PRD **92**, 013013 (2015)]

Analyze the “Mainz” data set

[Bernauer et al. PRL **105**, 242001 (2010)]

and world data (excluding Mainz)

- World data

[Lee, Arrington, Hill '15]

$$r_E^p = 0.918 \pm 0.024 \text{ fm}$$

[Hill, GP '10]

$$r_E^p = 0.870 \pm 0.023 \pm 0.012 \text{ fm}$$

[Lee, Arrington, Hill '15]

$$r_M^p = 0.913 \pm 0.037 \text{ fm}$$

[Epstein, GP, Roy '14]

$$r_M^p = 0.910^{+0.030}_{-0.060} \pm 0.020 \text{ fm}$$

- Mainz data

$$r_E^p = 0.895 \pm 0.020 \text{ fm} \quad r_M^p = 0.773 \pm 0.038 \text{ fm}$$

Part 2: Connecting muon-proton scattering and muonic hydrogen

The bottom line

- Scattering:
 - World $e - p$ data [Lee, Arrington, Hill '15]
 $r_E^p = 0.918 \pm 0.024$ fm
 - Mainz $e - p$ data [Lee, Arrington, Hill '15]
 $r_E^p = 0.895 \pm 0.020$ fm
 - Proton, neutron and π data [Hill, GP '10]
 $r_E^p = 0.871 \pm 0.009 \pm 0.002 \pm 0.002$ fm
- Muonic hydrogen
 - [Pohl et al. Nature **466**, 213 (2010)]
 $r_E^p = 0.84184(67)$ fm
 - [Antognini et al. Science **339**, 417 (2013)]
 $r_E^p = 0.84087(39)$ fm
- The bottom line:
using z expansion scattering disfavors muonic hydrogen
- Is there a problem with muonic hydrogen *theory*?

Muonic hydrogen theory

- Is there a problem with muonic hydrogen *theory*?

- Potentially yes!

[Hill, GP PRL **107** 160402 (2011)]

- Muonic hydrogen measures ΔE and translates it to r_E^p

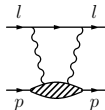
- [Pohl et al. Nature **466**, 213 (2010)]

$$\Delta E = 209.9779(49) - 5.2262(r_E^p)^2 + 0.0347(r_E^p)^3 \text{ meV}$$

- [Antognini et al. Science **339**, 417 (2013)]

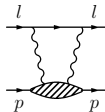
$$\Delta E = 206.0336(15) - 5.2275(10)(r_E^p)^2 + 0.0332(20) \text{ meV}$$

- In both cases apart from r_E^p need two-photon exchange



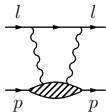
Two photon exchange

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Two photon exchange

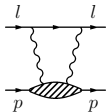
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$$\begin{aligned} & \frac{1}{2} \sum_s i \int d^4x e^{iq \cdot x} \langle \mathbf{k}, s | T \{ J_{\text{e.m.}}^\mu(x) J_{\text{e.m.}}^\nu(0) \} | \mathbf{k}, s \rangle \\ &= \left(-g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) W_1 + \left(k^\mu - \frac{k \cdot q q^\mu}{q^2} \right) \left(k^\nu - \frac{k \cdot q q^\nu}{q^2} \right) W_2 \end{aligned}$$

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- Dispersion relations ($\nu = 2k \cdot q$, $Q^2 = -q^2$)

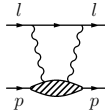
$$W_1(\nu, Q^2) = W_1(0, Q^2) + \frac{\nu^2}{\pi} \int_{\nu_{\text{cut}}(Q^2)^2}^{\infty} d\nu'^2 \frac{\text{Im} W_1(\nu', Q^2)}{\nu'^2(\nu'^2 - \nu^2)}$$

$$W_2(\nu, Q^2) = \frac{1}{\pi} \int_{\nu_{\text{cut}}(Q^2)^2}^{\infty} d\nu'^2 \frac{\text{Im} W_2(\nu', Q^2)}{\nu'^2 - \nu^2}$$

- W_1 requires subtraction...

Two photon exchange

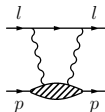
- In both cases apart from r_E^p we have two-photon exchange



- Imaginary part of TPE related to data:
form factors, structure functions

Two photon exchange

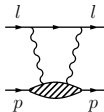
- In both cases apart from r_E^p we have two-photon exchange



- Imaginary part of TPE related to data:
form factors, structure functions
- Cannot reproduce it from its imaginary part:
Dispersion relation requires subtraction
- Need poorly constrained non-perturbative function $W_1(0, Q^2)$
Can calculate it
 - In small Q^2 limit using NRQED [Hill, GP, PRL **107** 160402 (2011)]
 - In large Q^2 limit using OPE [J. C. Collins, NPB **149**, 90 (1979)]

Two photon exchange

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 - In small Q^2 limit using NRQED [Hill, GP, PRL **107** 160402 (2011)]
 - In large Q^2 limit using OPE [J. C. Collins, NPB **149**, 90 (1979)]
- Introduces hard to quantify hadronic uncertainty

Variety of approaches

- Considering far-reaching implications of the puzzle we should explore a variety of approaches:

[Nevado, Pineda, PRC **77**, 035202 (2008)]

[Peset, Pineda, EPJA **51**, 32 (2015)]

[Peset, Pineda, NPB **887**, 69 (2014)]

[Carlson, Vanderhaeghen, PRA **84**, 020102 (2011)]

[Birse, McGovern, EPJA **48**, 120 (2012)]

[Miller PLB **718**, 1078 (2013)]

[Gorchtein, Llanes-Estrada, Szczepaniak PRA **87**, 052501 (2013)]

[Alarcon, Lensky, Pascalutsa, EPJC **74**, 2852 (2014)]

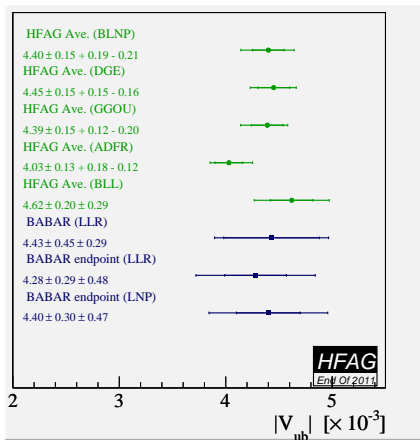
[Tomalak, Vanderhaeghen, PRD **90**, 013006 (2014)]

[Tomalak, Vanderhaeghen, EPJC **76**, 125 (2016)]

...

Variety of approaches

- Considering far-reaching implications of the puzzle we should explore a variety of approaches
- What you want to have:



- In the following: $W_i \rightarrow \text{NRQED}$

Why NRQED?

- Solving Schrödinger equation for $1/r$ potential

$$E_n = -\frac{1}{2}m_r c^2 \alpha^2 \frac{1}{n^2}$$

where $m_r = m_\mu m_p / (m_\mu + m_p) \approx m_\mu$

- Muon momentum in muonic hydrogen $p \sim m_\mu c \alpha \sim 1 \text{ MeV}$

- Muon is non-relativistic

Can use Non Relativistic QED (NRQED)

[Caswell, Lepage PLB **167**, 437 (1986); Kinoshita Nio PRD **53**, 4909 (1996); Manohar PRD **56**, 230 (1997)]

- For an introduction to NRQED see [GP MPLA **30**, 1550128 (2015)]

$$\begin{aligned}
 \mathcal{L}_p = & \psi_p^\dagger \left\{ iD_t + \frac{\mathbf{D}^2}{2m_p} + \frac{\mathbf{D}^4}{8m_p^3} + c_F e \frac{\boldsymbol{\sigma} \cdot \mathbf{B}}{2m_p} + c_D e \frac{[\partial \cdot \mathbf{E}]}{8m_p^2} \right. \\
 & + i c_S e \frac{\boldsymbol{\sigma} \cdot (\mathbf{D} \times \mathbf{E} - \mathbf{E} \times \mathbf{D})}{8m_p^2} + c_{W1} e \frac{\{\mathbf{D}^2, \boldsymbol{\sigma} \cdot \mathbf{B}\}}{8m_p^3} \\
 & - c_{W2} e \frac{D^i \boldsymbol{\sigma} \cdot \mathbf{B} D^i}{4m_p^3} + c_{p'p} e \frac{\boldsymbol{\sigma} \cdot \mathbf{D} \mathbf{B} \cdot \mathbf{D} + \mathbf{D} \cdot \mathbf{B} \boldsymbol{\sigma} \cdot \mathbf{D}}{8m_p^3} \\
 & \left. + i c_M e \frac{\{\mathbf{D}^i, [\partial \times \mathbf{B}]^i\}}{8m_p^3} + c_{A1} e^2 \frac{\mathbf{B}^2 - \mathbf{E}^2}{8m_p^3} - c_{A2} e^2 \frac{\mathbf{E}^2}{16m_p^3} + \dots \right\} \psi_p
 \end{aligned}$$

- The $1/m_p^4$ calculated in
[Hill, Lee, GP, Mikhail P. Solon, PRD **87** 053017 (2013)]
- Need also

$$\mathcal{L}_{\text{contact}} = d_1 \frac{\psi_p^\dagger \boldsymbol{\sigma} \psi_p \cdot \psi_\ell^\dagger \boldsymbol{\sigma} \psi_\ell}{m_\ell m_p} + d_2 \frac{\psi_p^\dagger \psi_p \psi_\ell^\dagger \psi_\ell}{m_\ell m_p}$$

- Matching

- Operators with one photon coupling:

c_i given by $F_i^{(n)}(0)$

- Operators with only two photon couplings:

c_{A_i} given by forward and backward Compton scattering

- d_i from two-photon amplitude

- From c_i and d_i determine proton structure correction, e.g.

$$\delta E(n, \ell) = \delta_{\ell 0} \frac{m_r^3 (Z\alpha)^3}{\pi n^3} \left(\frac{Z\alpha\pi}{2m_p^2} c_D^{\text{proton}} - \frac{d_2}{m_\ell m_p} \right)$$

- Bottom line: need $c_D \Leftrightarrow r_E^p$ and d_2

d_2 suffers from hadronic uncertainty

What to do?

Hadronic uncertainty d_2

- d_2 suffers from hadronic uncertainty

What to do?

- Improve modeling [Hill, GP *in progress*]

How large can it be?

- But even if d_2 *can* be large it does not follow that it *must* be large

- Experimental test: $\mu - p$ scattering

MUSE (MUon proton Scattering Experiment)

MUSE

- Muonic hydrogen:

Muon momentum $\sim m_\mu c \alpha \sim 1 \text{ MeV}$

Both proton and muon non-relativistic

MUSE

- Muonic hydrogen:

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Both proton and muon non-relativistic

- MUSE:

Muon momentum $\sim m_\mu c \sim 100 \text{ MeV}$

Muon is relativistic, proton is still non-relativistic

MUSE

- Muonic hydrogen:

Muon momentum $\sim m_\mu c \alpha \sim 1 \text{ MeV}$

Both proton and muon non-relativistic

- MUSE:

Muon momentum $\sim m_\mu \sim 100 \text{ MeV}$

Muon is relativistic, proton is still non-relativistic

- QED-NRQED effective theory:

- Use QED for muon

- Use NRQED for proton

$m_\mu/m_p \sim 0.1$ as expansion parameter

- A *new* effective field theory suggested in

[Hill, Lee, GP, Mikhail P. Solon, PRD **87** 053017 (2013)]

QED-NRQED Effective Theory

- Example: TPE at the lowest order in $1/m_p$

[Steven P. Dye, Matthew Gonderinger, GP, PRD **94** 013006 (2016)]

QED-NRQED Effective Theory

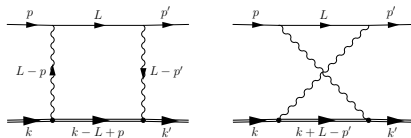
- Example: TPE at the lowest order in $1/m_p$
[Steven P. Dye, Matthew Gonderinger, GP, PRD **94** 013006 (2016)]
- Consider muon-proton scattering $\mu(p) + p(k) \rightarrow \mu(p') + p(k')$
 - At lowest order in $1/m_p$: $p^0 = p'^0 \Rightarrow \delta(p^0 - p'^0)$
 - At the proton rest frame $k = (m_p, \vec{0}) \Rightarrow k^0 = 0$ in NRQED

QED-NRQED Effective Theory

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- NRQED propagator: $\frac{1}{l^0 - \vec{l}^2/2M + i\epsilon}$

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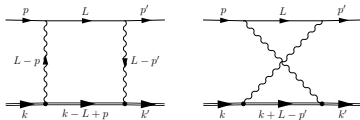


$$\frac{1}{p^0 - L^0 + i\epsilon} + \frac{1}{L^0 - p^0 + i\epsilon} \Rightarrow \delta(L^0 - p^0)$$

- In total

$$\delta(p^0 - p'^0) \delta(L^0 - p^0) = \delta(L^0 - p^0) \delta(L^0 - p'^0)$$

QED-NRQED Effective Theory



- The amplitude

$$\begin{aligned}
 i\mathcal{M}(2\pi)^4\delta^4(k+p-k'-p') &= Z^2 e^4 \int \frac{d^4 L}{(2\pi)^4} \frac{1}{(L-p)^2 (L-p')^2} \\
 &\times \bar{u}(p')\gamma^0 \frac{i}{\not{L}-m}\gamma^0 u(p)\chi^\dagger \chi (2\pi)\delta(L^0-p^0) (2\pi)\delta(L^0-p'^0) \\
 &\times (2\pi)^3\delta^4(\vec{p}-\vec{p}'-\vec{k}')
 \end{aligned}$$

- The cross section

$$\frac{d\sigma}{d\Omega} = \frac{Z^2 \alpha^2 4E^2 (1 - v^2 \sin^2 \frac{\theta}{2})}{\vec{q}^4} \left[1 + \frac{Z\alpha\pi v \sin \frac{\theta}{2} (1 - \sin \frac{\theta}{2})}{1 - v^2 \sin^2 \theta} \right]$$

$Z = 1$, $E =$ muon energy, $v = |\vec{p}|/E$, $q = p' - p$, θ scattering angle

QED-NRQED Effective Theory

- QED-NRQED result

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- *Same result* as scattering relativistic lepton off static $1/r$ potential
[Dalitz, Proc. Roy. Soc. Lond. **206**, 509 (1951)]
reproduced in [Itzykson, Zuber, “Quantum Field Theory”]

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- *Same result* as scattering relativistic lepton off static $1/r$ potential [Dalitz, Proc. Roy. Soc. Lond. **206**, 509 (1951)]
reproduced in [Itzykson, Zuber, “Quantum Field Theory”]
- *Same result* as $m_p \rightarrow \infty$ of “point particle proton” QED scattering
(For $m_p \rightarrow \infty$ only proton charge is relevant)

QED-NRQED Effective Theory beyond $m_p \rightarrow \infty$ limit

- QED-NRQED allows to calculate $1/m_p$ corrections
- Example: one photon exchange $\mu + p \rightarrow \mu + p$:
QED-NRQED = $1/m_p$ expansion of form factors
[Steven P. Dye, Matthew Gonderinger, GP, PRD **94** 013006 (2016)]

Connecting muon-proton scattering to muonic hydrogen

- Matching

QED, QCD

$G_{E,M}$, Structure func., $W_1(0, Q^2)$

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QED-NRQED: *MUSE*

$r_E^p, \bar{\mu} \gamma^0 \mu \psi_p^\dagger \psi_p$

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\Downarrow

NRQED-NRQED: *muonic H*

$r_E^p, \psi_\mu^\dagger\psi_\mu\psi_p^\dagger\psi_p$

- Need to match QED-NRQED contact interaction, e.g. $\bar{\mu}\gamma^0\mu\psi_p^\dagger\psi_p$ to NRQED-NRQED contact interaction, e.g. $\psi_\mu^\dagger\psi_\mu\psi_p^\dagger\psi_p$
[Dye, Gonderinger, GP *in progress*]

Connecting muon-proton scattering to muonic hydrogen

- To do list:
 - 1) Relate QED-NRQED contact interactions to NRQED contact interactions and $W_1(0, Q^2)$
 - 2) Calculate $d\sigma(\mu + p \rightarrow \mu + p)$ and asymmetry in terms of r_E^p and d_2
 - 3) *Direct* relation between μ - p scattering and muonic H

Conclusions

- Proton radius puzzle: $> 5\sigma$ discrepancy between
 - r_E^p from muonic hydrogen
 - r_E^p from hydrogen and $e - p$ scattering
- After 6 years the origin is still not clear
 - 1) Is it a problem with the electronic extraction?
 - 2) Is it a hadronic uncertainty?
 - 3) is it new physics?
- Motivates a reevaluation of our understanding of the proton

Conclusions

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Conclusions

- Presented two topics:
 - 1) Extraction of proton radii from scattering:
Use an established tool of the z expansion
Studies disfavor the muonic hydrogen value
 - 2) Direct connection between muon-proton scattering and muonic hydrogen using a new effective field theory: QED-NRQED
- Much more work to do!
- Thank you