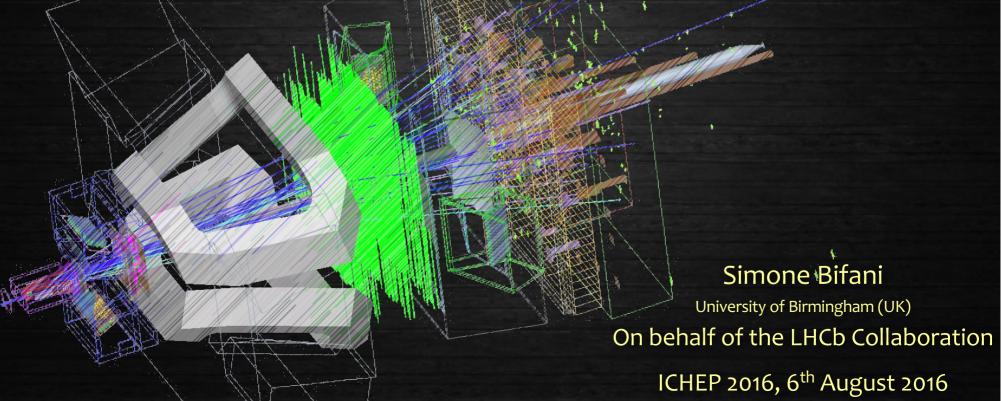




Rare b Decays and Tests of LFU at LHCb





A Forward Spectrometer



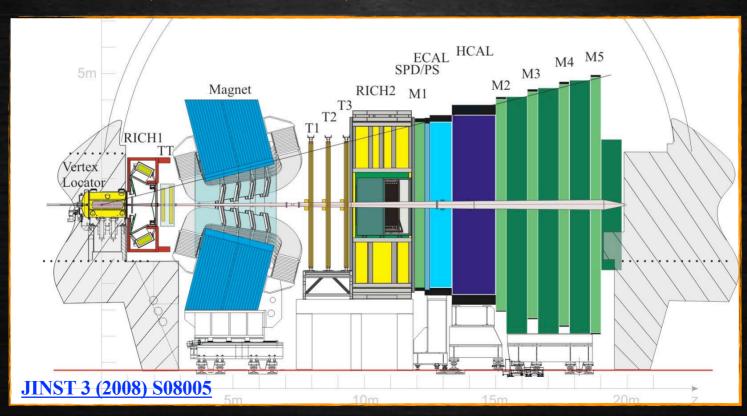
> Optimized for beauty and charm physics at large pseudorapidity ($2<\eta<5$)

» Trigger: ~90% efficient for di-muon channels, ~30% for all-hadronic

» Tracking: σ_p/p 0.4%–0.6% (p from 5 to 100 GeV), σ_{IP} < 20 μm

» Vertexing: σ_{τ} ~45 fs

» PID: 97% μ ID for 1–3% $\pi \rightarrow \mu$ misID

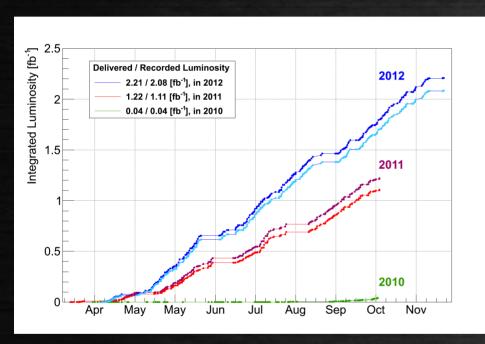


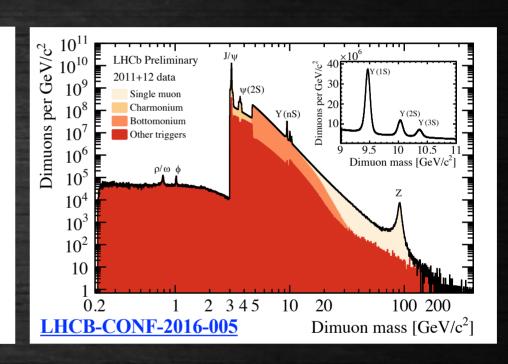


Datasets



> Analyses presented today based on the Run 1 dataset





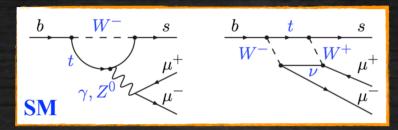
> Due to luminosity levelling, same running conditions throughout fills



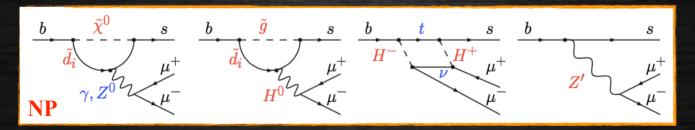
Why Rare b Decays?



> b→sll decays proceed via FCNC transitions that only occur at loop order (and beyond) in SM



> New particles can contribute to loop or tree level enhancing/suppressing decay rates, introducing new sources of CP violation or modifying the angular distribution of the final-state particles



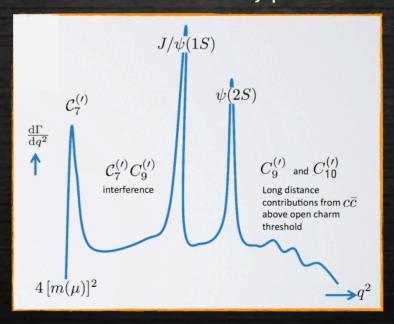
- > Goal
 - » Make precise measurements of rare FCNC decays as precision tests of the SM
 - » Make null tests of the SM, e.g. look for LFV or LNV decays that are essentially forbidden in the SM



Shopping List



- > Differential branching fractions of $B^0 \rightarrow K^{(*)0} \mu \mu$, $B^+ \rightarrow K^{(*)+} \mu \mu$, $B_s \rightarrow \Phi \mu \mu$, $B^+ \rightarrow \pi^+ \mu \mu$ and $\Lambda_b \rightarrow \Lambda \mu \mu$ decays
 - » Large hadronic uncertainties in theory predictions
- > **Angular analyses** of $B \rightarrow K^{(*)} \mu \mu$, $B_s \rightarrow \phi \mu \mu$, $B^o \rightarrow K^{*o}$ ee and $\Lambda_b \rightarrow \Lambda \mu \mu$
 - » Define observables with small theory uncertainties
- > Test of Lepton Flavour Universality in B+→K+ll
 - » Cancellation of hadronic uncertainties in theory predictions

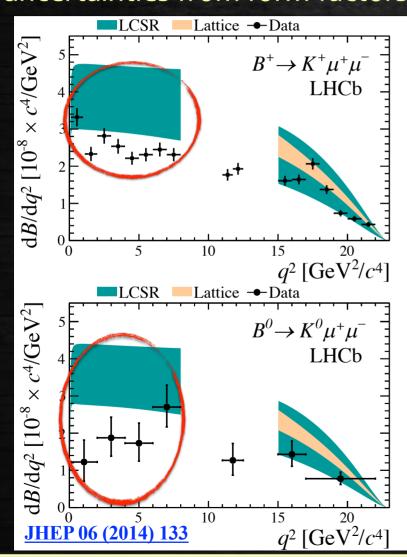


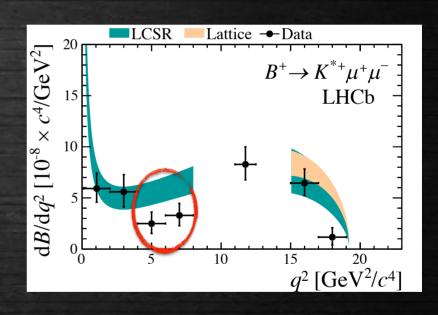


Differential Branching Fractions



> Results **consistently lower than SM predictions** despite large theory uncertainties from form-factors



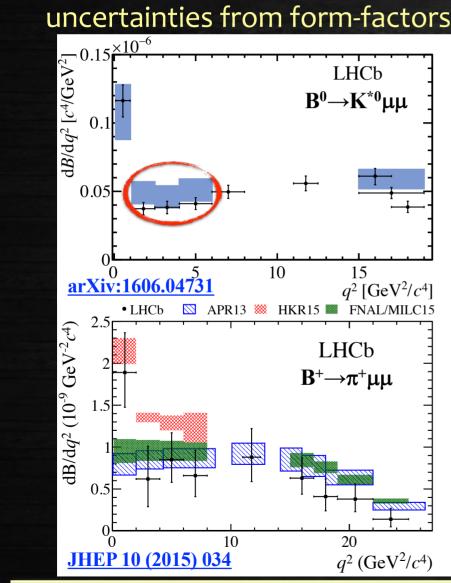


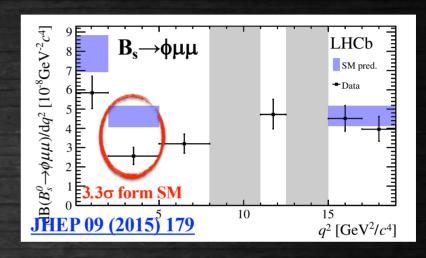


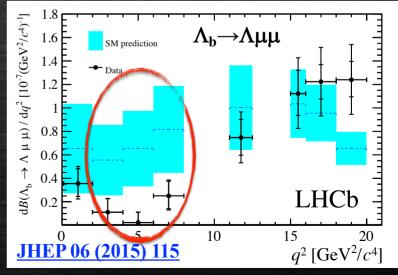
Differential Branching Fractions



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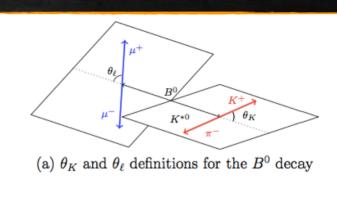


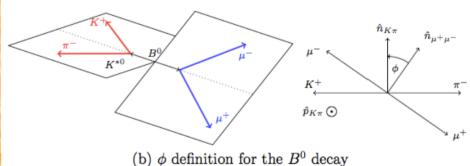


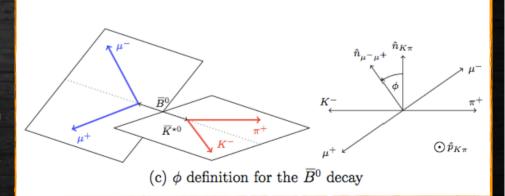




- > Four-body final states
- > System described by three angles (helicity basis) and the di-lepton invariant mass squared, q²
- Complex angular distribution that provides many observables sensitive to different types of NP
- Each observable depends on different Wilson coefficients (underlying short-distance physics) and form-factors (hadronic matrix elements)



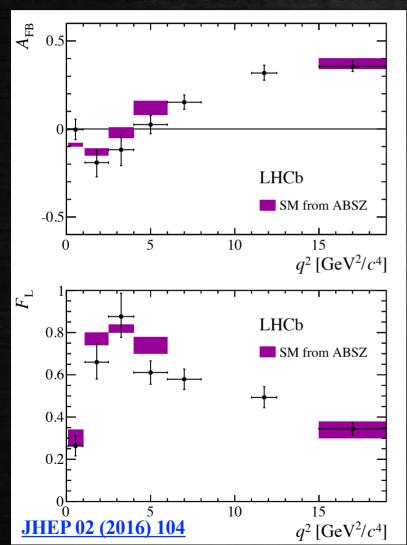




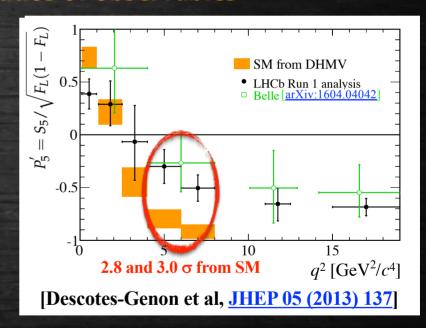




> First **full angular analysis** of $B^0 \rightarrow K^{*0} \mu \mu$: full set of CP-averaged angular terms and correlations as well as full set of CP-asymmetries



Can construct form-factor independent ratios of observables

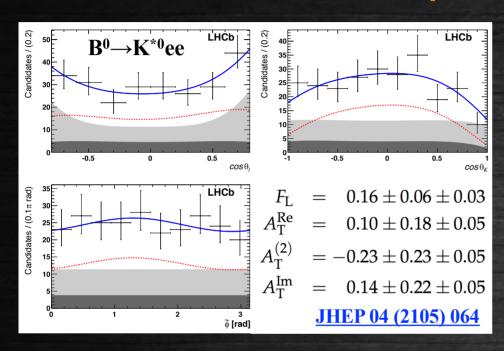


New Belle result consistent with LHCb



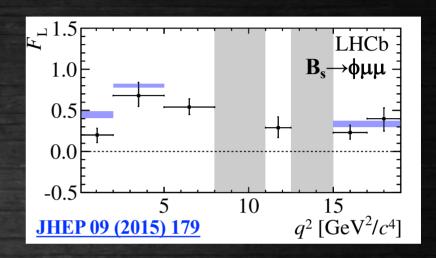


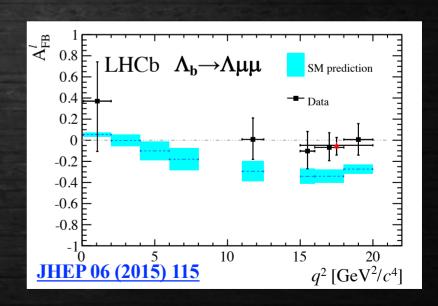
> Results consistent with SM predictions



- > Low-q²: 0.0004–1 GeV²
- > Challenging due to Bremsstrahlung
- > Sensitive to photon polarisation

Λ_b: gives access to different combinations of Wilson coefficients



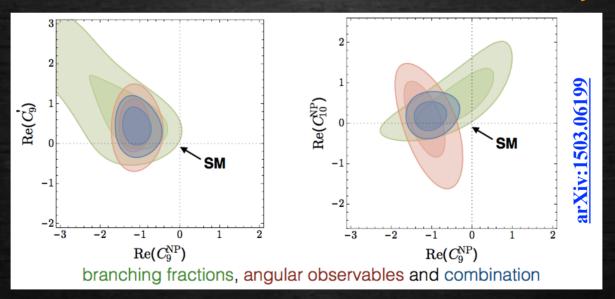




Global Fits



- > Several attempts to interpret results by performing global fits to b→s data (e.g. arXiv:1503.06199, arXiv:1510.04239 and arXiv:1512.07157)
- > Take into account ~80 observables from 6 experiments including $b \rightarrow \mu\mu$, $b \rightarrow sll$ and $b \rightarrow s\gamma$ transitions
- > All global fits require an additional contribution with respect to the SM to accommodate the data, with a preference for NP in C₉ at ~40



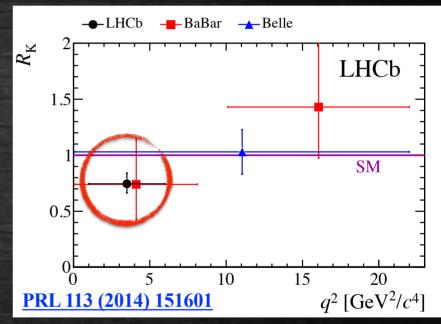
Or is this a problem with our understanding of QCD?
 (e.g. are we correctly estimating the contribution for charm loops?)



Tests of LFU



- > Ratio of branching fractions of $B^+ \rightarrow K^+ \mu \mu$ to $B^+ \rightarrow K^+ ee$ expected to be unity in the SM (theoretical uncertainty of O(10⁻³))
- > Observation of LFU violations would be a clear sign of NP
- Extremely challenging due to Bremsstrahlung and different trigger / tracking performances between muons and electrons
- > Measured relative to $B^+ \rightarrow K^+ J/\psi(II)$ in order to reduce systematics
- > Observed a 2.6σ deviation from SM



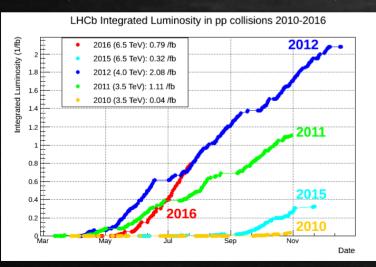
- > Consistent with decay rate if NP does not couple to electrons
- > Pursuing other R-like measurements (e.g see talk by S.Klaver on R(D*))

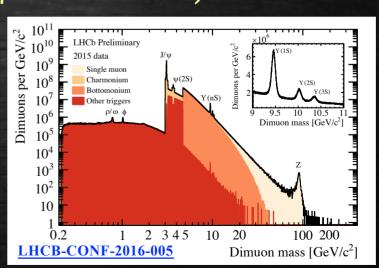


Summary and Outlook



- > Rare b decays place strong constraints on many NP models allowing to probe energy scales higher than direct searches
- > A large number of analyses have been performed using Run 1 data
- > While there is no significant evidence for NP from a single measurement, a clear tension with the SM have been seen in global fits to rare decay observables
- > Rare decays will largely benefit from the increase of energy (cross-section) and collected data (~5 fb⁻¹ expected in LHCb) in Run 2

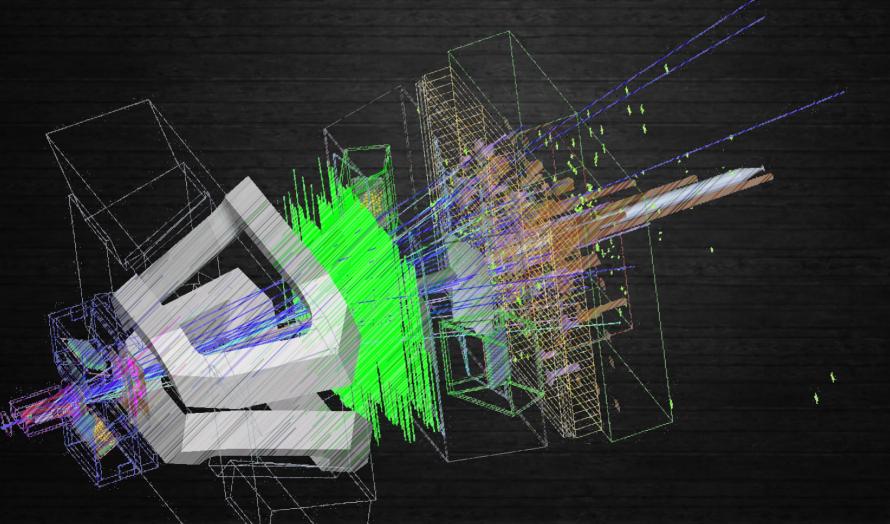








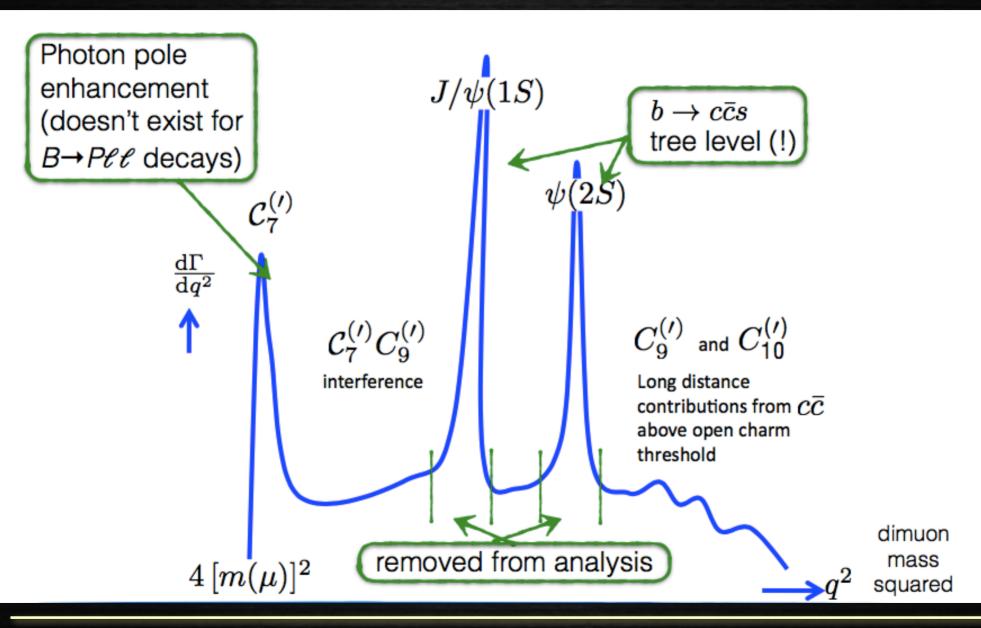
Backup





Di-Lepton Mass

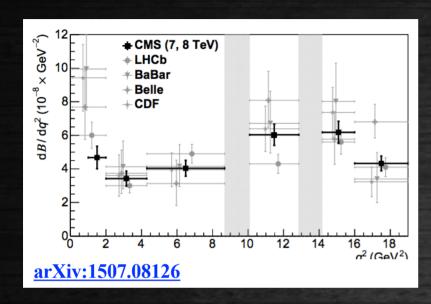


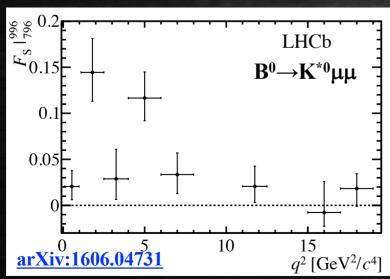


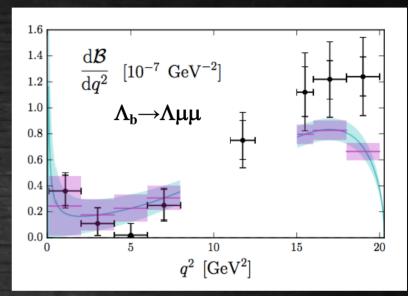


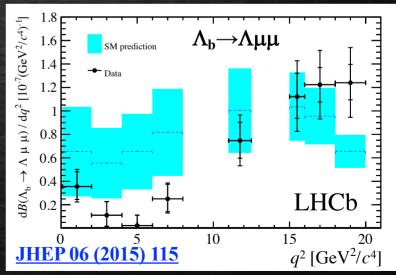
Differential Branching Fractions











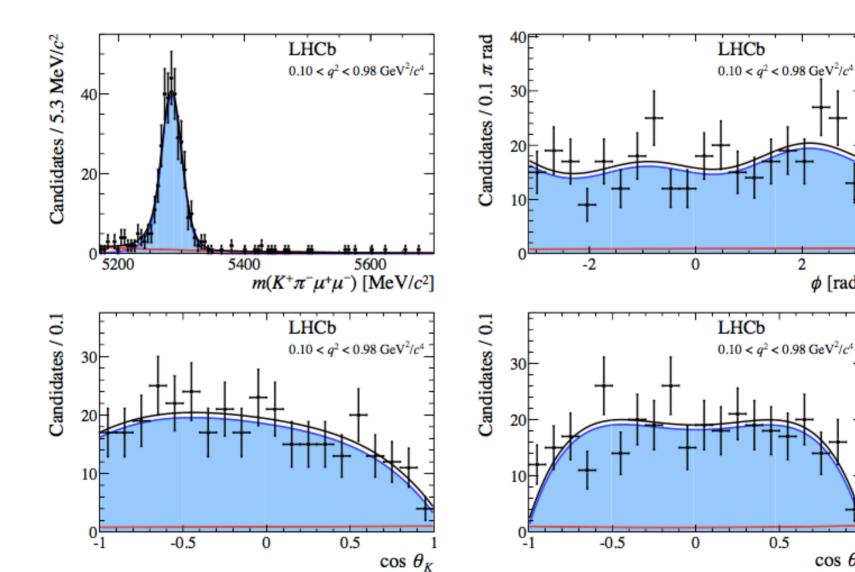




 ϕ [rad]

 $\cos \theta_l$

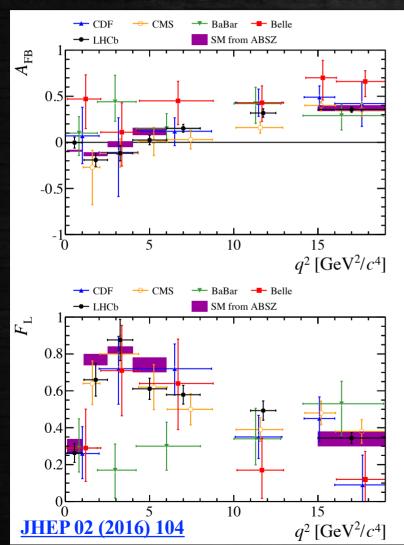
0.5



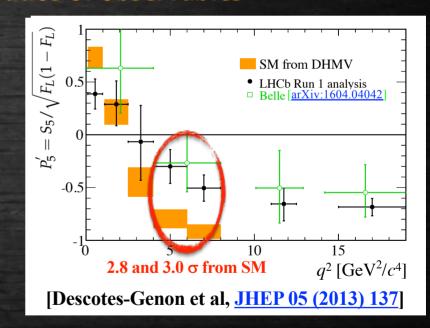




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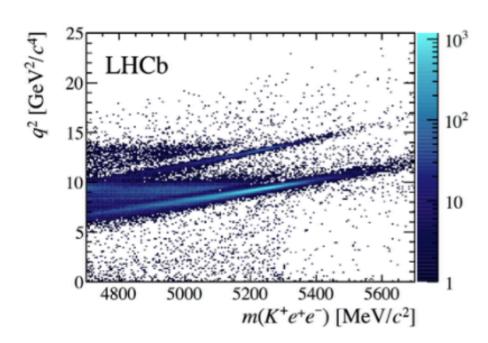
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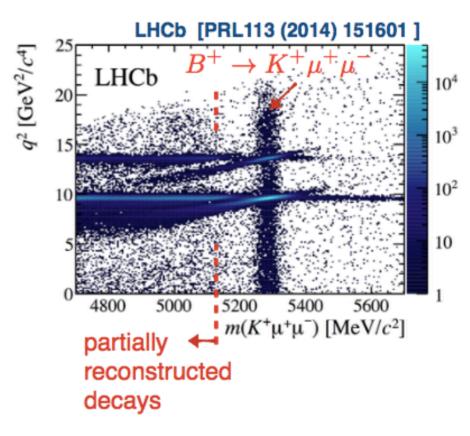


Tests of LFU



 Even after Bremsstrahlung recovery there are significant differences between dielectron and dimuon final states:



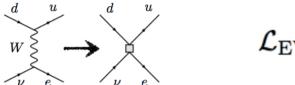




Theoretical Framework



 In the Fermi model of the weak interaction, the full electroweak Lagrangian (which was unknown at the time) is replaced by the low-energy theory (QED) plus a single operator with an effective coupling constant.



$${\cal L}_{
m EW} o {\cal L}_{
m QED} + rac{G_{
m F}}{\sqrt{2}} (\overline{u}d) (ear{
u})$$

Can write a Hamiltonian for the effective theory as

$$\mathcal{H}_{\text{eff}} = -\frac{4 G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{\alpha_e}{4\pi} \sum_{i} C_i(\mu) \mathcal{O}_i(\mu),$$

Wilson coefficient (integrating out scales above μ)

Local operator with different Lorentz structure (vector, axial vector current etc)



Operators



SM operators

photon penguin $\mathcal{O}_7 = rac{m_b}{c} ar{s} \sigma^{\mu u} P_R b F_{\mu u} \,,$ $\mathcal{O}_8 = g_s \frac{m_b}{c^2} \bar{s} \sigma^{\mu\nu} P_R T^a b G^a_{\mu\nu} \,,$ $=ar s\gamma_\mu P_L b\,ar\ell\gamma^\mu\ell\,,$ $|\mathcal{O}_{10}| = \bar{s} \gamma_{\mu} P_L b \, \bar{\ell} \gamma^{\mu} \gamma_5 \ell \quad ,$ vector and axial-vector currents

Beyond SM operators

$$\mathcal{O}_7' = rac{m_b}{e} ar{s} \sigma^{\mu
u} P_L b F_{\mu
u} \,,$$
 $\mathcal{O}_8' = g_s rac{m_b}{e^2} ar{s} \sigma^{\mu
u} P_L T^a b G^a_{\mu
u} \,,$
 $\mathcal{O}_9' = ar{s} \gamma_\mu P_R b \, ar{\ell} \gamma^\mu \ell \,,$
 $\mathcal{O}_{10}' = ar{s} \gamma_\mu P_R b \, ar{\ell} \gamma^\mu \gamma_5 \ell \,.$

right handed currents (suppressed in SM)





Complex angular distribution:

$$\frac{1}{\mathrm{d}(\Gamma + \bar{\Gamma})/\mathrm{d}q^2} \frac{\mathrm{d}^3(\Gamma + \bar{\Gamma})}{\mathrm{d}\bar{\Omega}} \bigg|_{\mathrm{P}} = \frac{9}{32\pi} \Big[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K + F_{\mathrm{L}} \cos^2 \theta_K \Big] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K \right] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_K \right] + \frac{1}{2\pi} \left[\frac{3}{4} (1 - F_{\mathrm{L}}) \sin^2 \theta_$$

fraction of longitudinal polarisation of the K*

$$+\frac{1}{4}(1-F_{\rm L})\sin^2\theta_K\cos2\theta_l$$

 $-F_{\rm L}\cos^2\theta_K\cos 2\theta_l + S_3\sin^2\theta_K\sin^2\theta_l\cos 2\phi$

forward-backward asymmetry of the dilepton system

$$+S_4 \sin 2\theta_K \sin 2\theta_l \cos \phi + S_5 \sin 2\theta_K \sin \theta_l \cos \phi$$

$$+\frac{4}{3}A_{\mathrm{FB}}\sin^2 heta_K\cos heta_l+S_7\sin2 heta_K\sin heta_l\sin\phi$$

$$+S_8\sin 2\theta_K\sin 2\theta_l\sin \phi + S_9\sin^2\theta_K\sin^2\theta_l\sin 2\phi$$

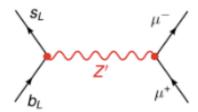
The observables depend on form-factors for the $B \rightarrow K^*$ transition plus the underlying short distance physics (Wilson coefficients).



Interpretation of Global Fits

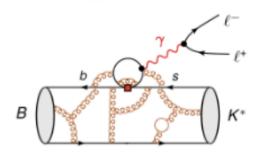


Optimist's view point



Vector-like contribution could come from new tree level contribution from a Z' with a mass of a few TeV

Pessimist's view point



Vector-like contribution could point to a problem with our understanding of QCD, e.g. are we correctly estimating the contribution for charm loops that produce dimuon pairs via a virtual photon.

More work needed from experiment/theory to disentangle the two



Interpretation of Global Fits



 This is the physics we are interested in. (a) ℓ^+ (b) γ $\ell^ \ell^ \delta$ δ δ δ

Short distance part integrates out (as a Wilson coefficient)

 We also get long-distance hadronic contributions.
 Included in the SM but are the predictions correct?

