

Direct CP violation in $K \rightarrow \pi\pi$ decays and supersymmetry

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$$A_0 \equiv A(K^0 \rightarrow (\pi\pi)_{I=0}) \quad \text{and} \quad A_2 \equiv A(K^0 \rightarrow (\pi\pi)_{I=2}),$$

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Direct CP violation (from decay amplitude):

$$\epsilon'_K \simeq \frac{\epsilon_K}{\sqrt{2}} \left[\frac{\langle (\pi\pi)_{I=2} | K_L \rangle}{\langle (\pi\pi)_{I=0} | K_L \rangle} - \frac{\langle (\pi\pi)_{I=2} | K_S \rangle}{\langle (\pi\pi)_{I=0} | K_S \rangle} \right] = (16.6 \pm 2.3) \cdot 10^{-4} \cdot \epsilon_K$$

discovered in **1999**

Experimentally well-known:

$$\begin{aligned}\operatorname{Re}A_0 &= (3.3201 \pm 0.0018) \times 10^{-7} \text{ GeV}, \\ \operatorname{Re}A_2 &= (1.4787 \pm 0.0031) \times 10^{-8} \text{ GeV}.\end{aligned}$$



assumes PDG convention for CKM elements

Master equation for ϵ'_K :

$$\frac{\epsilon'_K}{\epsilon_K} = \frac{\omega_+}{\sqrt{2}|\epsilon_K^{\text{exp}}|\text{Re}A_0^{\text{exp}}} \left\{ \frac{\text{Im}A_2}{\omega_+} - \left(1 - \hat{\Omega}_{\text{eff}}\right) \text{Im}A_0 \right\}.$$

Here:

$$\omega_+ \simeq \frac{\text{Re}A_2}{\text{Re}A_0} = (4.53 \pm 0.02) \cdot 10^{-2}$$

ϵ_K quantifies indirect CP violation in $K-\bar{K}$ mixing,
 $\hat{\Omega}_{\text{eff}} = (14.8 \pm 8.0) \cdot 10^{-2}$ quantifies isospin breaking.

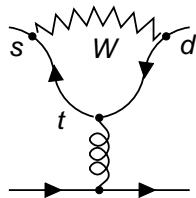
Important theoretical ingredients: $\text{Im}A_0$ and $\text{Im}A_2$, calculated from the effective $|\Delta S| = 1$ hamiltonian describing $s \rightarrow dq\bar{q}$ decays.

The enhanced sensitivity to $\Delta I = 3/2$ transitions (such as electroweak penguins and boxes) is a **special feature** of ϵ'_K .

$\text{Im}A_0$ is dominated by gluon penguins:

Operator: $Q_6 = \bar{s}_L^j \gamma_\mu d_L^k \sum_q \bar{q}_R^k \gamma^\mu q_R^j$

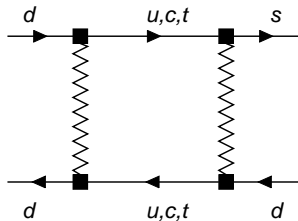
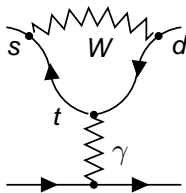
Matrix element: $\langle (\pi\pi)_{I=0} | Q_6 | K^0 \rangle$



$\text{Im}A_2$ is dominated by photon penguin and box diagrams:

Operator: $Q_8 = \frac{3}{2} \bar{s}_L^j \gamma_\mu d_L^k \sum_q e_q \bar{q}_R^k \gamma^\mu q_R^j$

Matrix element: $\langle (\pi\pi)_{I=2} | Q_8 | K^0 \rangle$



$$\frac{\epsilon'_K}{\epsilon_K} = (16.6 \pm 2.3) \times 10^{-4} \quad (\text{experiments: NA62, KTeV})$$

$$\frac{\epsilon'_K}{\epsilon_K} = (1.0 \pm 4.7 \pm 1.5 \pm 0.6) \times 10^{-4} \quad (\text{SM-NLO})$$

Kitahara, UN, Tremper, 1607.06727

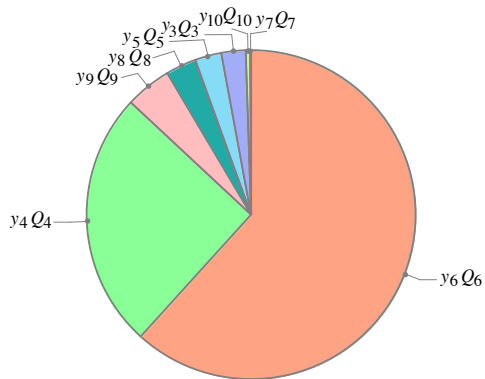
The prediction uses the lattice-QCD results from **RBC-UKQCD**,
Phys. Rev. Lett. **115** 212001 (2015).

Discrepancy with a significance of **2.9 σ** !

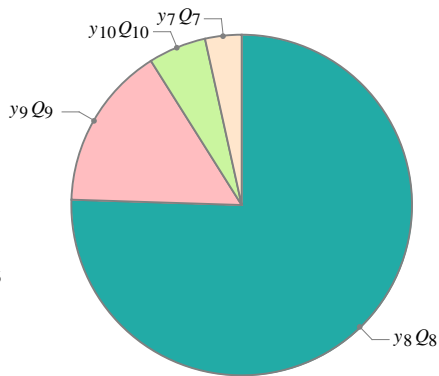
Our prediction uses the methodology of Buras et al. (JHEP 1511 (2015) 202), taking **ReA_{0,2}** from data, and a new formula for the RG evolution, building on Adams, Lee, Phys.Rev. D75 (2007) 074502.

Relative contributions of different operators

$\text{Im}A_0$



$\text{Im}A_2$



Standard Model:

Cabibbo-Kobayashi-Maskawa (CKM) factor:

$$\tau = -\frac{V_{td} V_{ts}^*}{V_{ud} V_{us}^*} \sim (1.5 - 0.6i) \cdot 10^{-3}$$

$$\epsilon_K^{\prime\text{SM}} \propto \text{Im} \frac{\tau}{M_W^2} \quad \text{and} \quad \epsilon_K^{\text{SM}} \propto \text{Im} \frac{\tau^2}{M_W^2}.$$

Generic new physics:

Some flavour-violating parameter: δ

$$\epsilon_K^{\prime\text{NP}} \propto \text{Im} \frac{\delta}{M^2} \quad \text{and} \quad \epsilon_K^{\text{NP}} \propto \text{Im} \frac{\delta^2}{M^2}.$$

with new heavy particle mass $M \gg M_W$.

But data require $|\epsilon_K^{\text{NP}}| \leq |\epsilon_K^{\text{SM}}|$, so that

$$\left| \frac{\epsilon_K^{\text{NP}}}{\epsilon_K^{\text{SM}}} \right| \leq \frac{|\epsilon_K^{\text{NP}} / \epsilon_K^{\text{SM}}|}{|\epsilon_K^{\text{NP}} / \epsilon_K^{\text{SM}}|} = \mathcal{O} \left(\frac{\text{Re } \tau}{\text{Re } \delta} \right).$$

If new physics enters through loops, we need $|\delta| \gg |\tau|$ to compensate for $M \gg M_W$, but then ϵ_K prohibits large effects in ϵ_K' .

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Solutions in the literature:

- tree-level new physics (e.g. Z') with $|\delta| \sim |\tau|$
- fine-tuning of the CP phase to get $\text{Re } \delta \sim 0$
- exploit special features of supersymmetry (this talk).

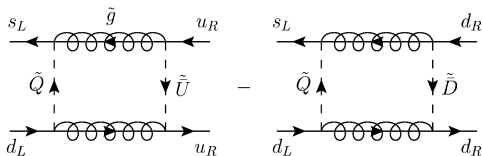
The **MSSM** has a mechanism

- to enhance $\text{Re}A_2$, because it permits strong-isospin violation through splittings between right-handed up-squark and down-squark masses (**Trojan penguins**),

Grossman, Kagan, Neubert 1999.

- to suppress the $K-\bar{K}$ mixing amplitude thanks to the Majorana nature of the gluinos, with negative interference of two box diagrams.

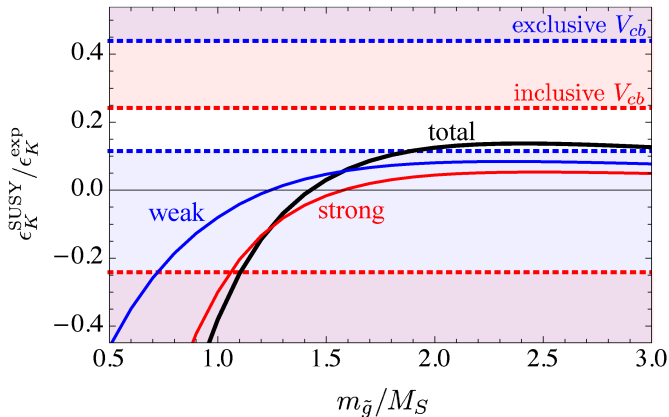
The second feature makes the **MSSM** contribution to $K-\bar{K}$ mixing vanish for $M_{\tilde{g}} \sim 1.5M_{\tilde{q}}$, it stays small for $M_{\tilde{g}} > 1.5M_{\tilde{q}}$.



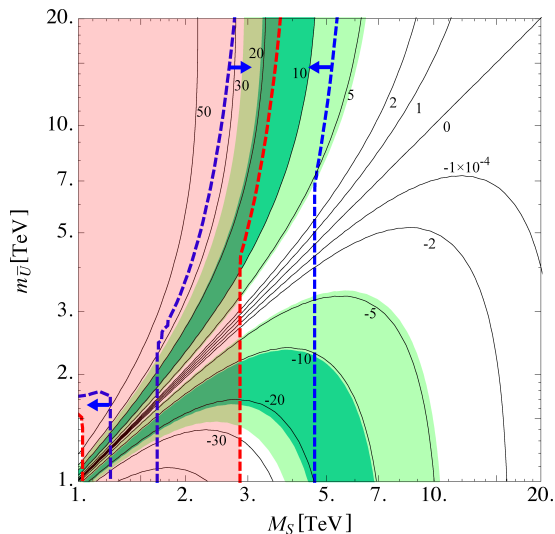
Choose:

Sparticle masses $M_S \sim 10$ TeV, $M_{\tilde{g}} > 1.5M_S$, flavour mixing in down-squark mass matrix only with $\arg \Delta_{sd}^{LL} = \pi/4$.

$M_S = 10$ TeV



Explain ϵ'_K



x-axis: generic sparticle mass, $M_{\tilde{g}} = 1.5M_S$

y-axis: right-handed up-squark mass

red region: excluded by ϵ_K if $|V_{cb}|$ from inclusive decays is correct

blue dashes: delimit allowed region, if $|V_{cb}|$ from exclusive decays is correct

- The new lattice results for the matrix element $\langle (\pi\pi)_{I=0} | Q_6 | K^0 \rangle$ from **RBC-UKQCD** points to a tension between the experimental value of ϵ'_K and the Standard-Model prediction.
- If **new physics** enters through loops, a sizable effect in ϵ'_K requires a new source of flavour violation which is much larger than the **CKM factor** $\text{Im} \frac{V_{td} V_{ts}^*}{V_{ud} V_{us}^*} \sim 6 \cdot 10^{-4}$. But then the effect on ϵ_K will typically be too big.
- In the **MSSM** one can simultaneously enhance ϵ'_K and suppress the new-physics contributions to ϵ_K . This requires flavour mixing among **left-handed squarks**, masses of right-handed **up-type squarks** different from those of the **down-type squarks**, and a **gluino mass** above **1.5** times the mass of the left-handed squarks.

Penguins: Wake-up call for new physics?

