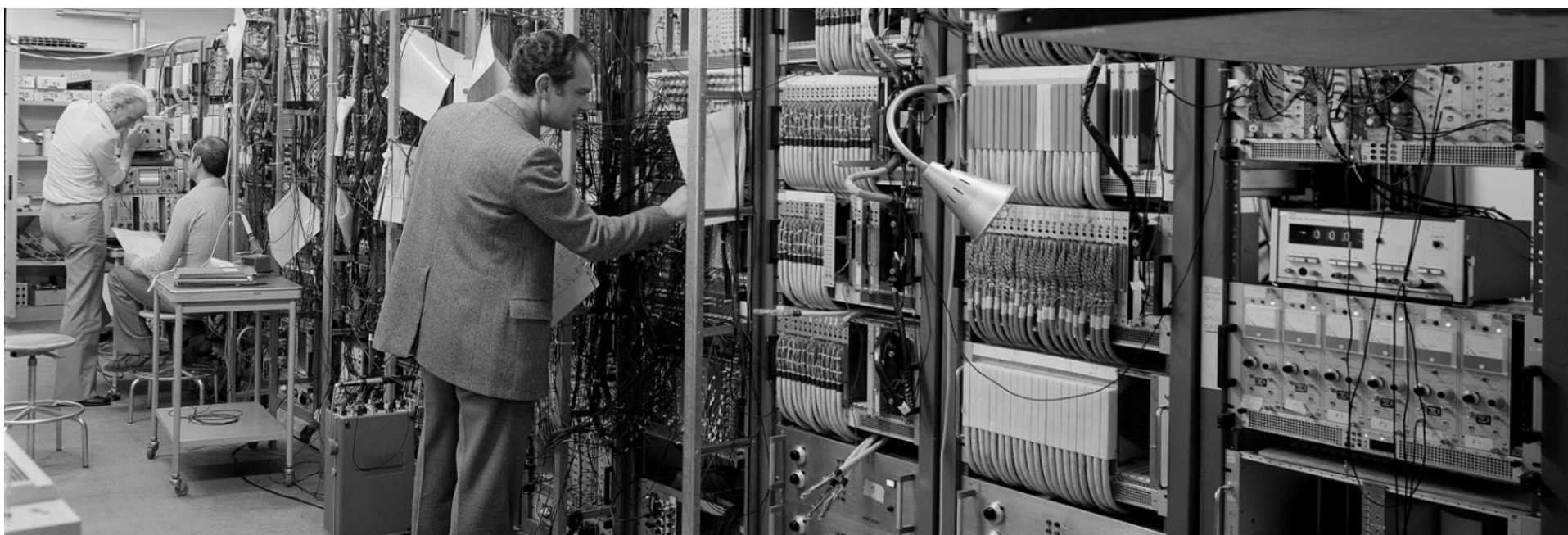


# DAQ HW

## Hands-on Approach



**ISOTDAQ 2016: 7th International School of Trigger and Data Acquisition**

Rehovot, 25 Jan 2016

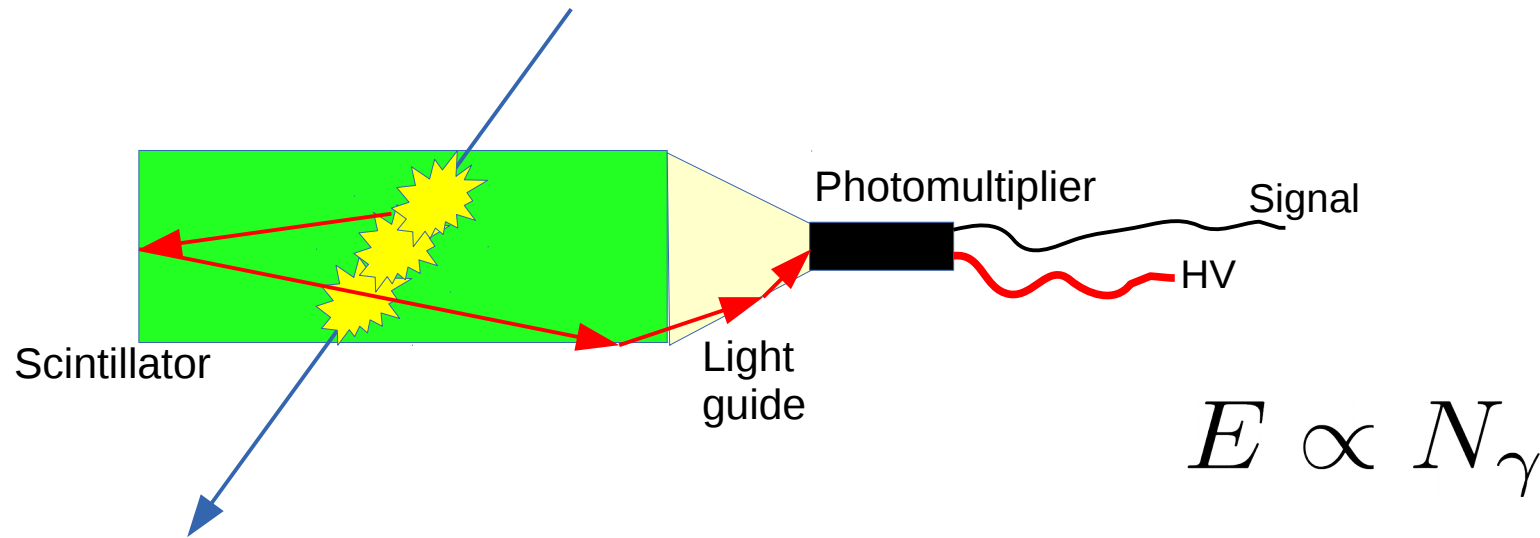
Andrea.Negri@pv.infn.it

© Wainer Vandelli & Sergio Ballestrero

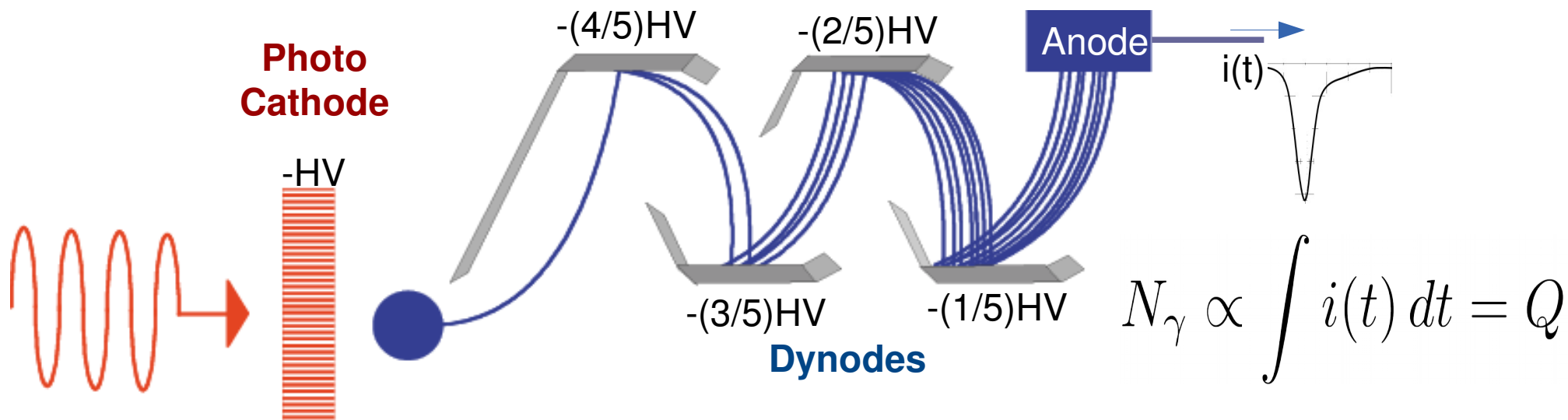
- This wants to be a hands-on approach to the basic DAQ hardware
  - We will discuss two different experiments, requiring different techniques and components
  - We also have some good real data to discuss
  - You will see, we are talking about real life here
- Acknowledgements
  - © Wainer Vandelli (CERN/PH-ATD)
  - © Sergio Ballestrero (Univ. Johannesburg & CERN)
    - Material and ideas have been taken from CERN Summer Student lectures of P.Farthouat, C.Joram and O.Ullaland

- Introduction
- Measure energy deposition
  - Scintillator setup
  - Photomultiplier
  - Analog-to-Digital conversion
  - Charge-to-Digital conversion
  - QDC in real life
- Measure position
  - Wire chamber setup
  - Time-to-Digital conversion
  - TDC in real life
- Corollary





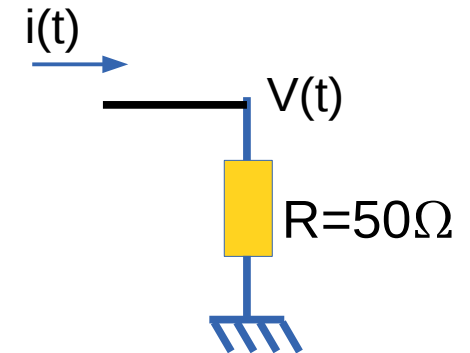
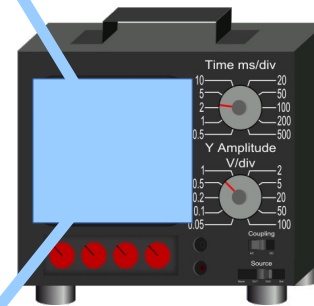
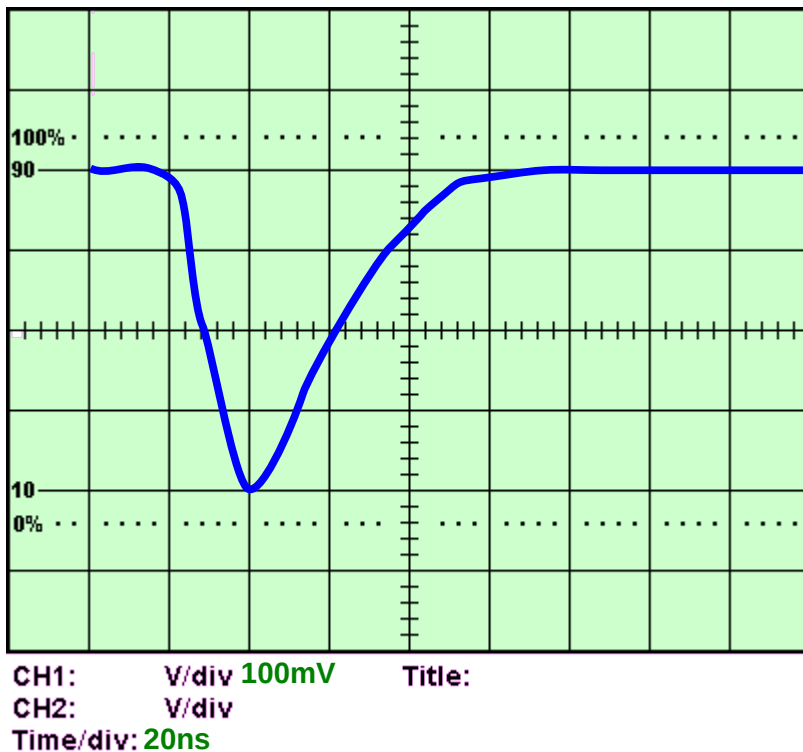
- Measure energy deposited by a particle traversing a medium
- The medium (detector) is a **scintillator**
  - Molecules, excited by the passing particle, relax emitting light
  - The amount of light is proportional to the deposited energy
- The light is then
  - collected, using dedicated optical means (**light guide**)
  - fed into a photo-detector: **photomultiplier**



- **Photo cathode:** photon to electron conversion via photo-electric effect
  - typical quantum efficiency  $\approx 1\text{-}10\%$  (max  $30\%$ ), depends on material and wavelength
- **Dynodes:** electrodes that amplify number of electrons thanks to secondary emission
  - typical overall gain  $\approx 10^6$
- **Dark current:** noise
  - current flowing in PMT without light

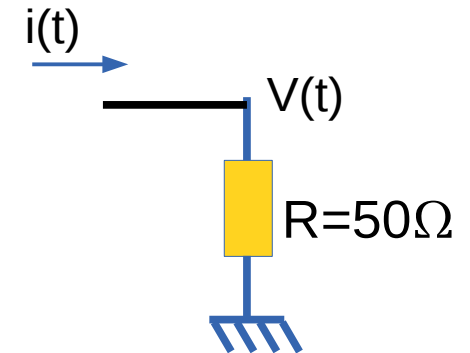
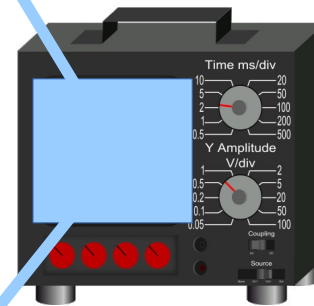
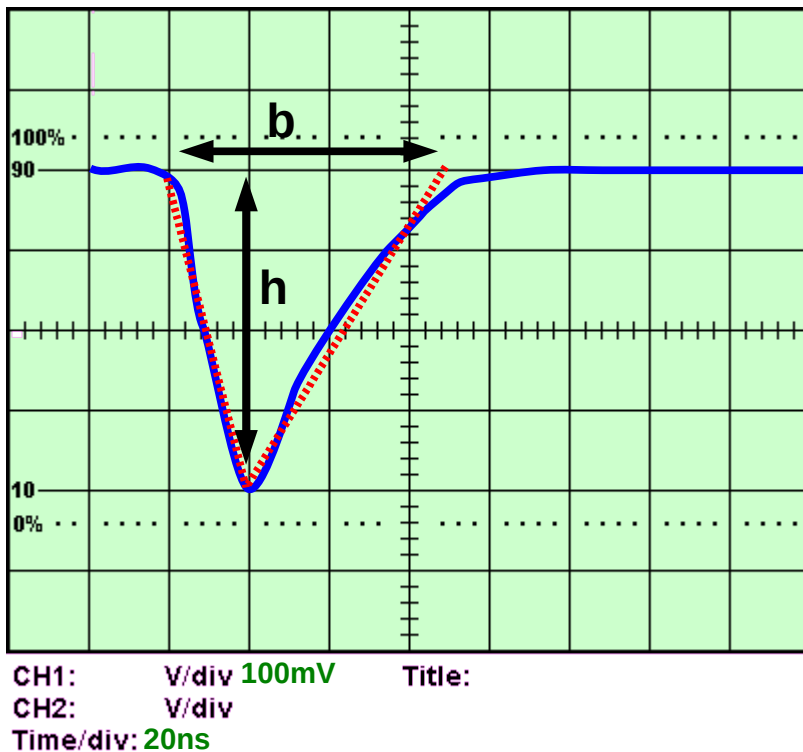


- Approximate Q measurement using oscilloscope
  - Linear approximation of a exponential decay



$$Q = \int i(t) dt = \frac{1}{R} \int V(t) dt$$

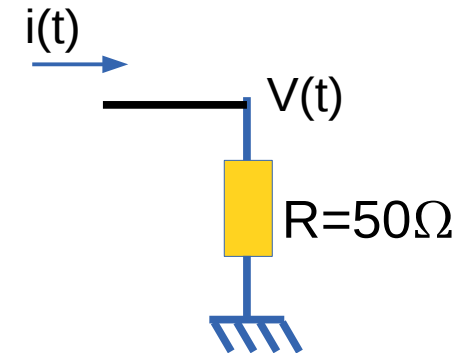
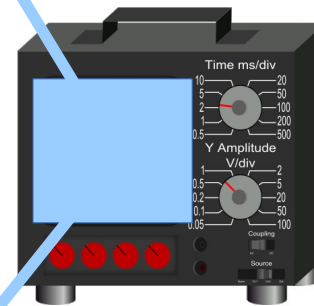
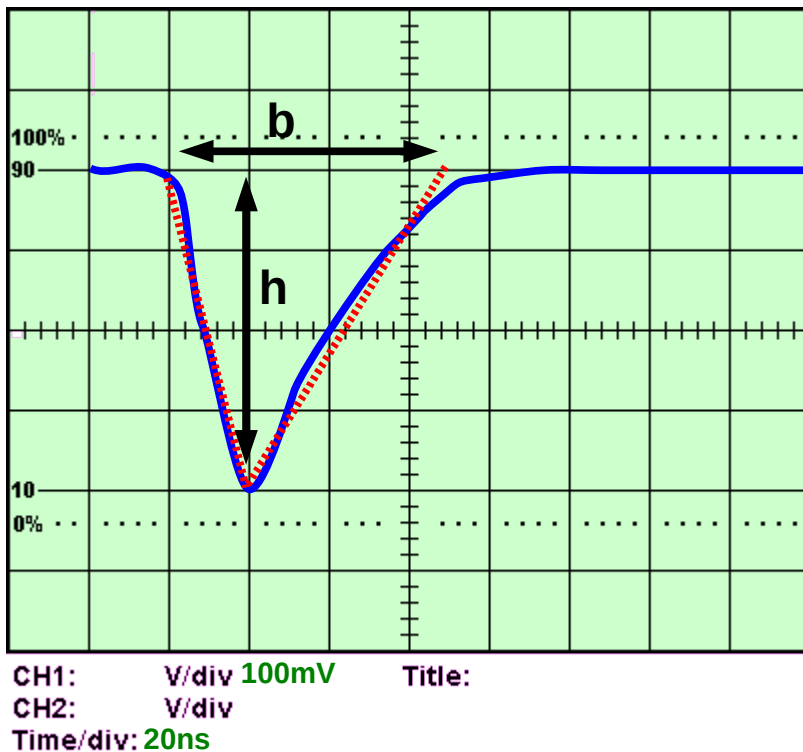
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$$Q = \int i(t) dt = \frac{1}{R} \int V(t) dt$$

$$Q \approx \frac{1}{R} \frac{bh}{2} = \frac{1}{50\Omega} \frac{(3.5 \cdot (20\text{ns}))(4 \cdot (100\text{mV}))}{2} = 280\text{pC}$$

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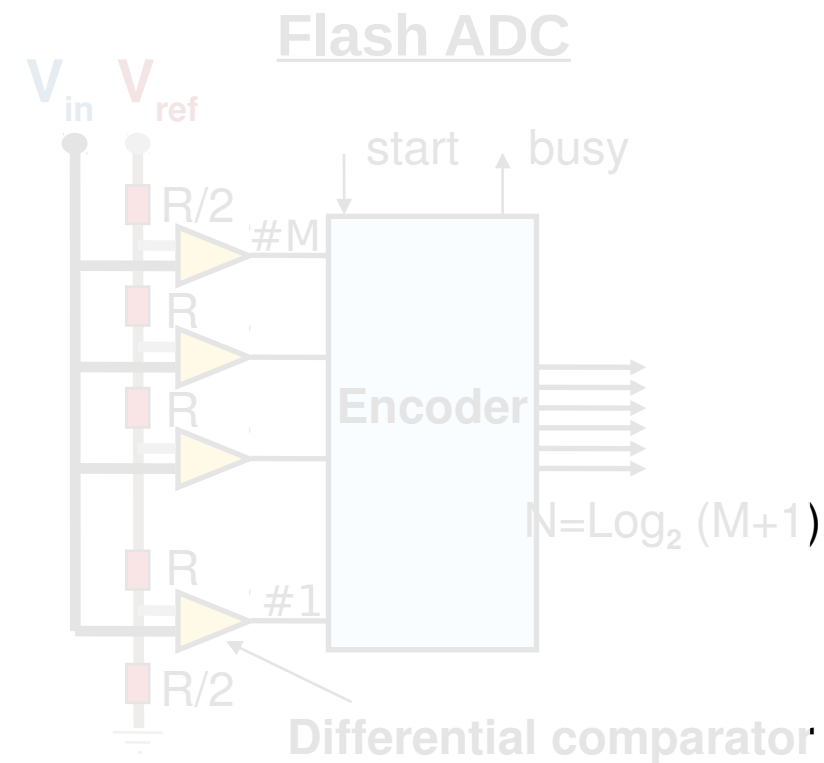
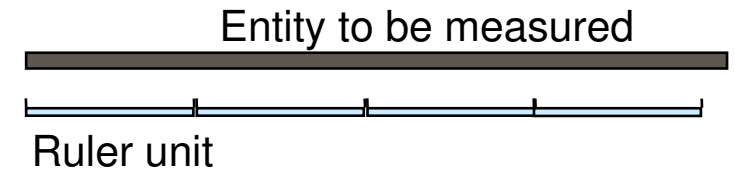
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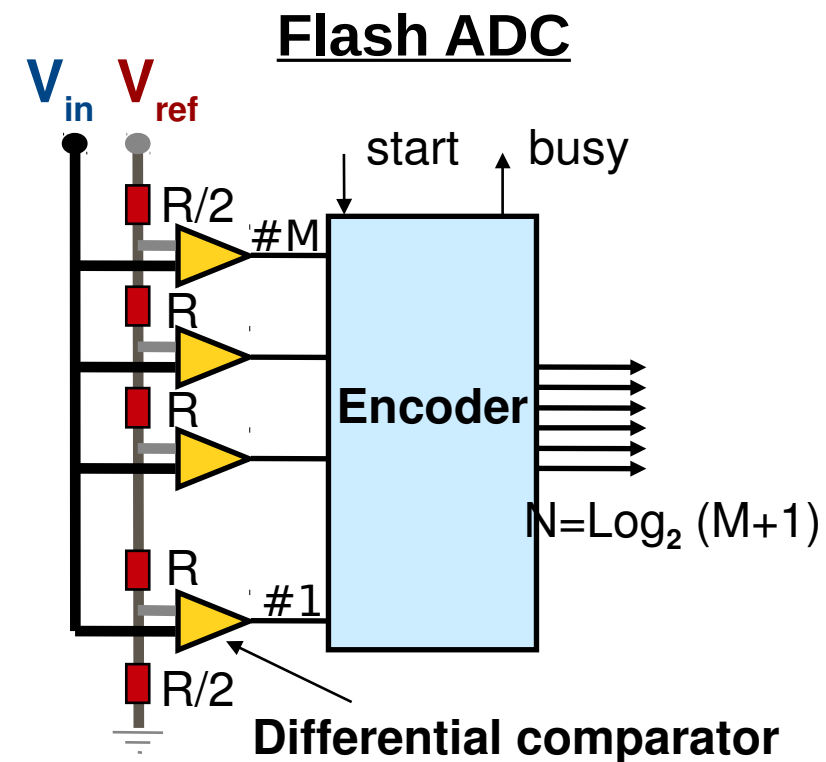
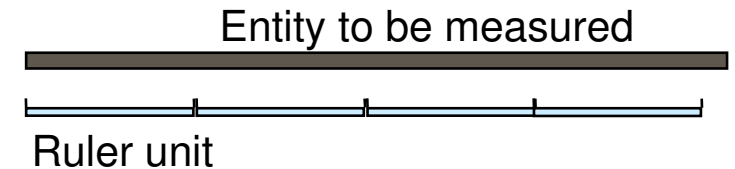
- Approximate Q measurement using oscilloscope
  - Linear approximation of a exponential decay
- Easy, but
  - Deadtime 5 min,  $\sim 3 \times 10^{-3}$ /Hz (if you are good)
  - Necessary to encode data into some sort of electronic format by hand
- Wouldn't be much more convenient to have a direct electronic measurement?
  - It could save the data in some digital format and fill a histogram on-line. Wouldn't be cool?
- N.B.: the oscilloscope method is still fundamental
  - it allows for the **validation** of your DAQ
  - yes, you should never thrust it a priori!



- Digitization — Stefan' talk and lab 8
  - Encoding an analog value into a binary representation
  - By comparing entity with a ruler
- Flash ADC simplest and fastest implementation
  - M comparisons in parallel
  - Input voltage  $V_{in}$  compared with M fractions of a reference voltage
    - $(1/2) V_{ref} / M \rightarrow (M-1/2) V_{ref} / M$
    - E.g.: M=3
  - Result is encoded into a compact binary form of N bits
    - $N = \log_2 (M+1)$



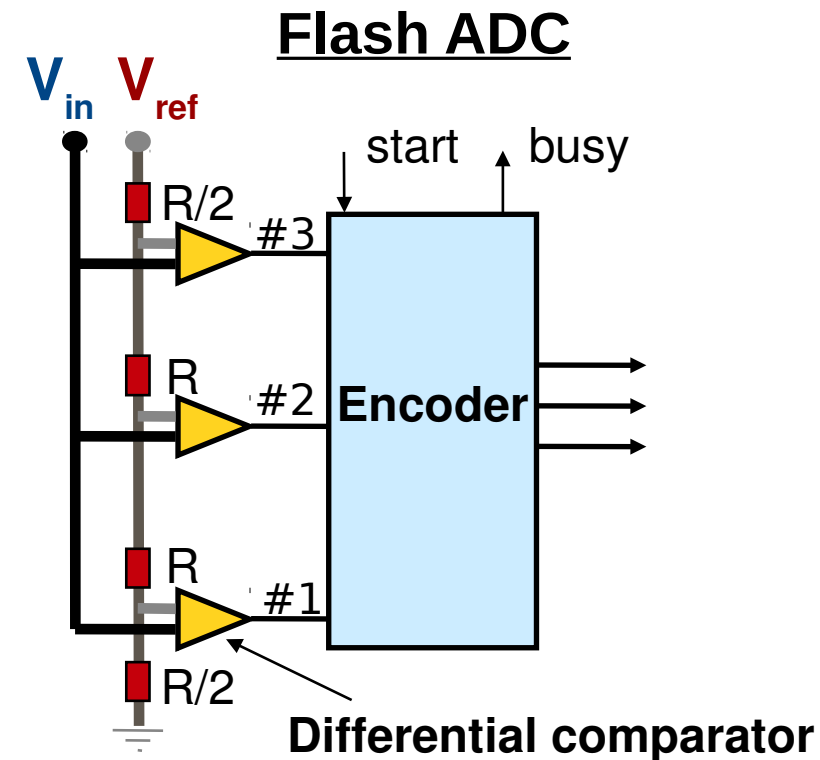
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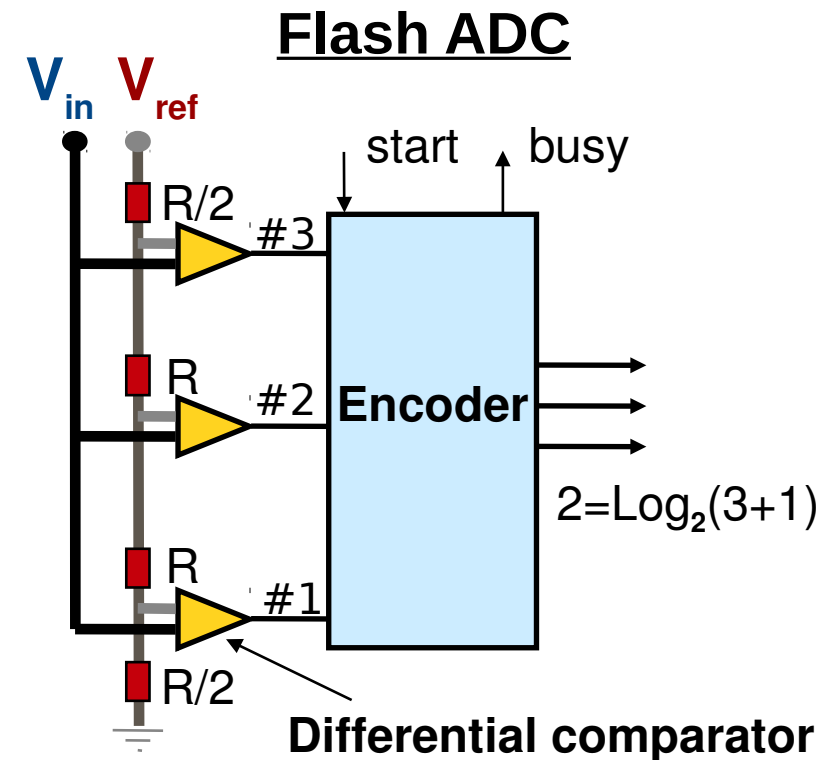
$x = V_{in}/V_{ref}$	Comparison results	Encoded form
$x < 1/6$	000	00
$1/6 \leq x < 3/6$	001	01
$3/6 \leq x < 5/6$	011	10
$5/6 \leq x$	111	11

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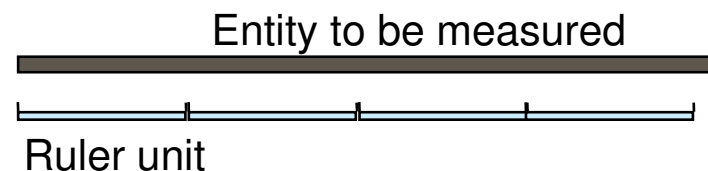


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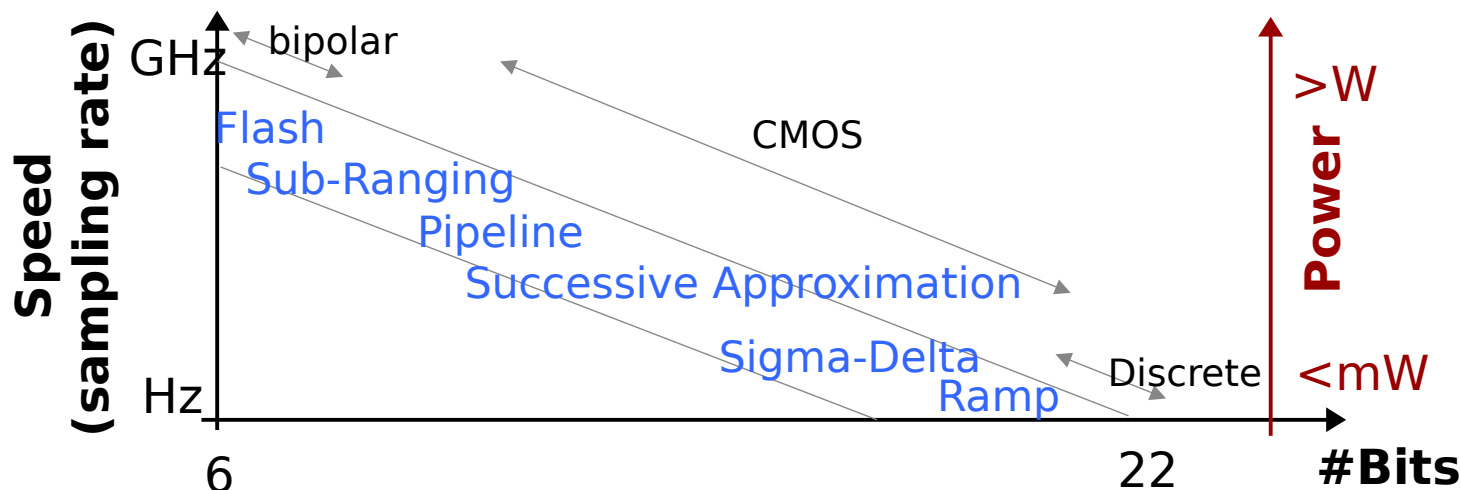
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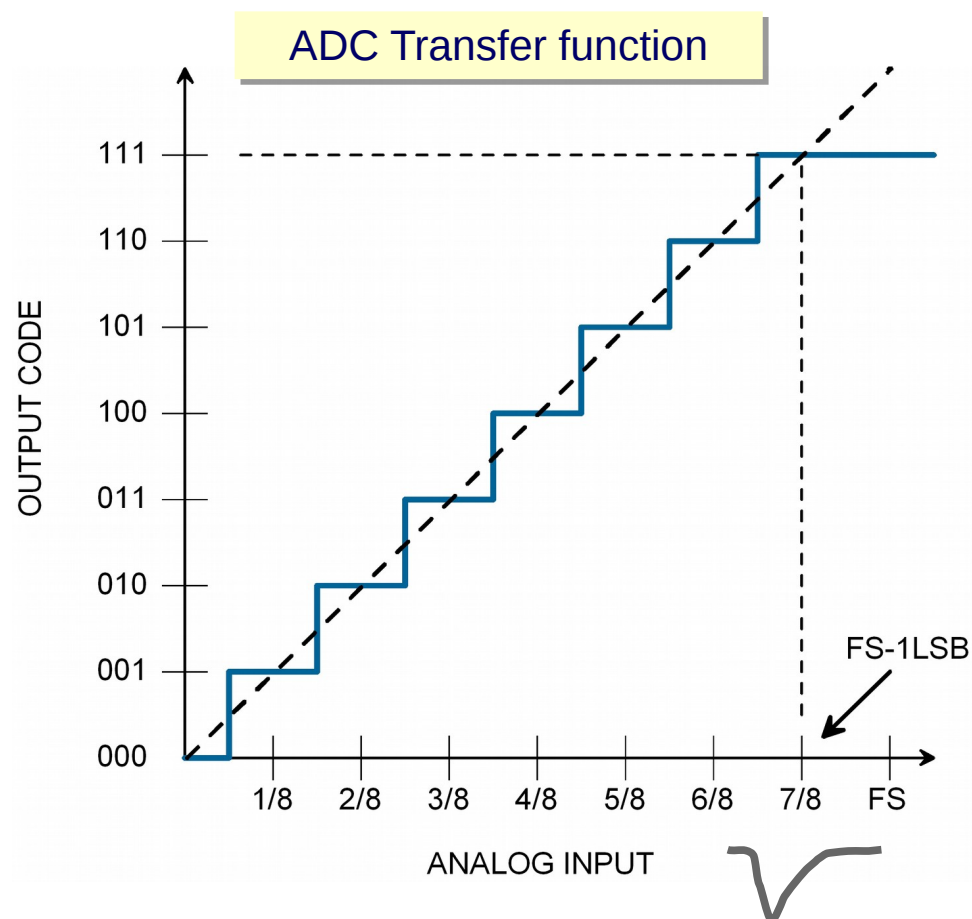
- Resolution (LSB), the ruler unit:  $V_{\max}/N$ 
  - e.g.: 1V and 8bit ( $N=256$ )  $\rightarrow$  LSB = 3.9 mV
- Quantization error:  $\pm \text{LSB}/2$
- Dynamic range: ratio largest /smallest value (in  $\log_2$ )
  - $N$  for linear ADC
  - $>N$  for non-linear ADC  
(Constant relative resolution on the valid input range)
- Many different ADC technique exists
  - mostly because of the trade-off between speed and resolution



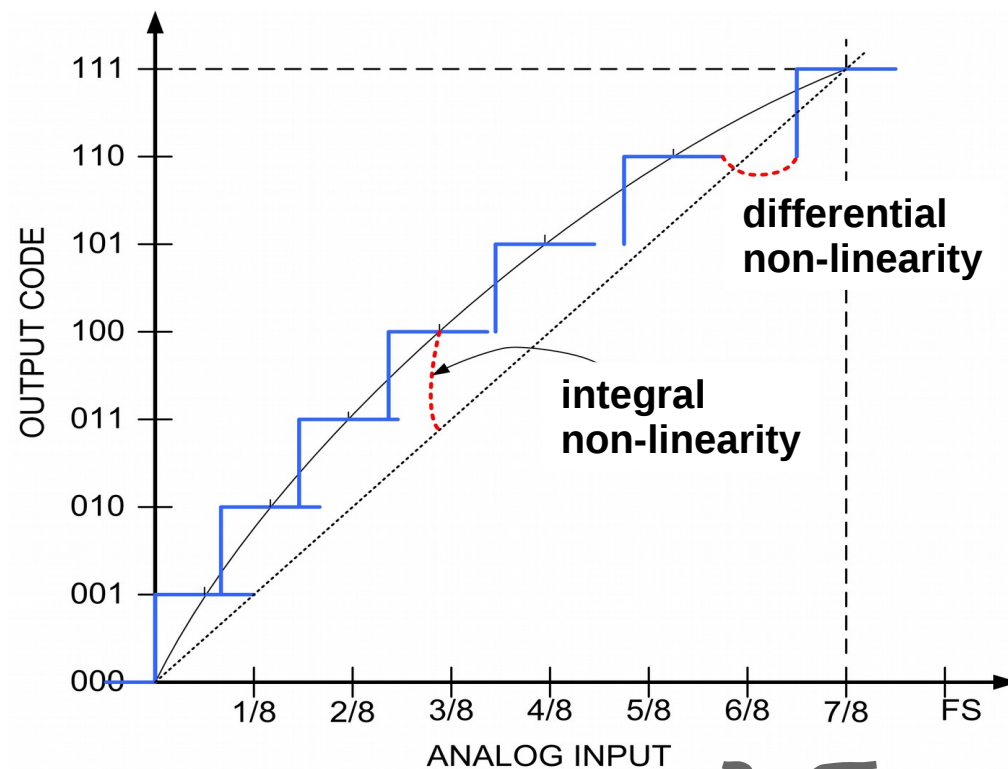
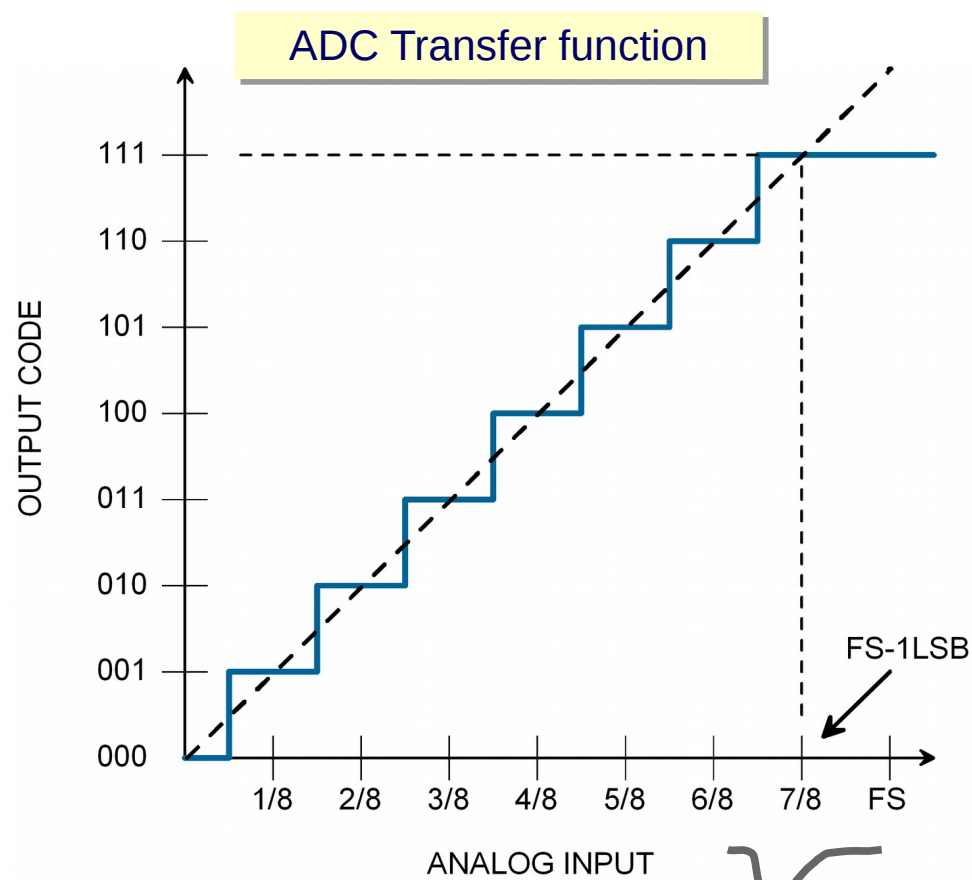
Lab 4



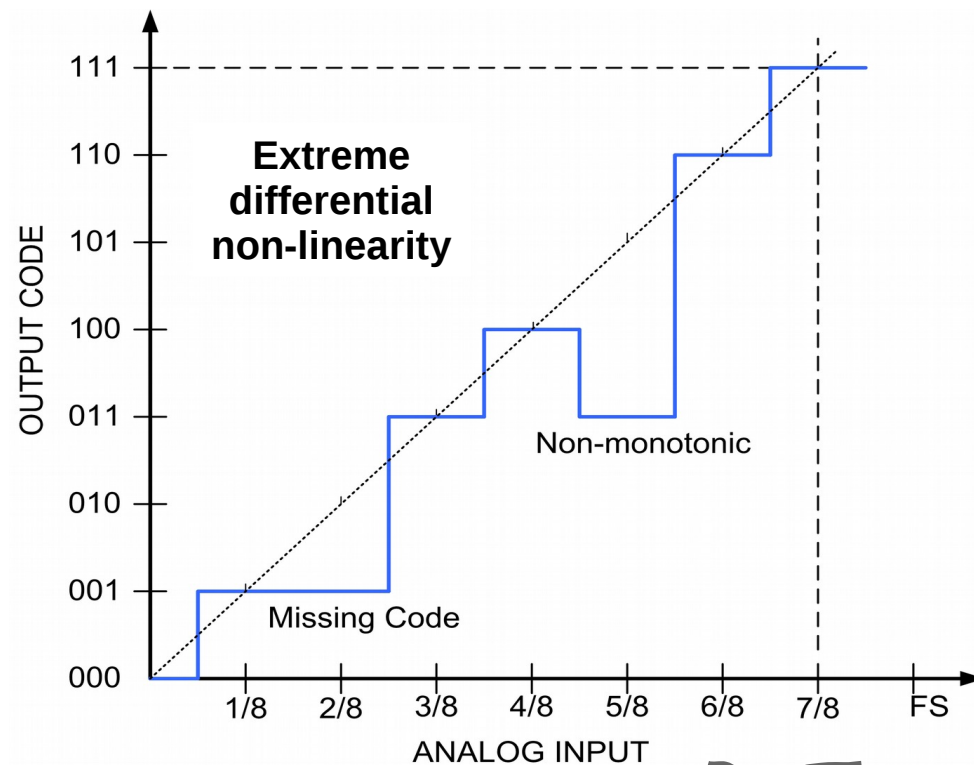
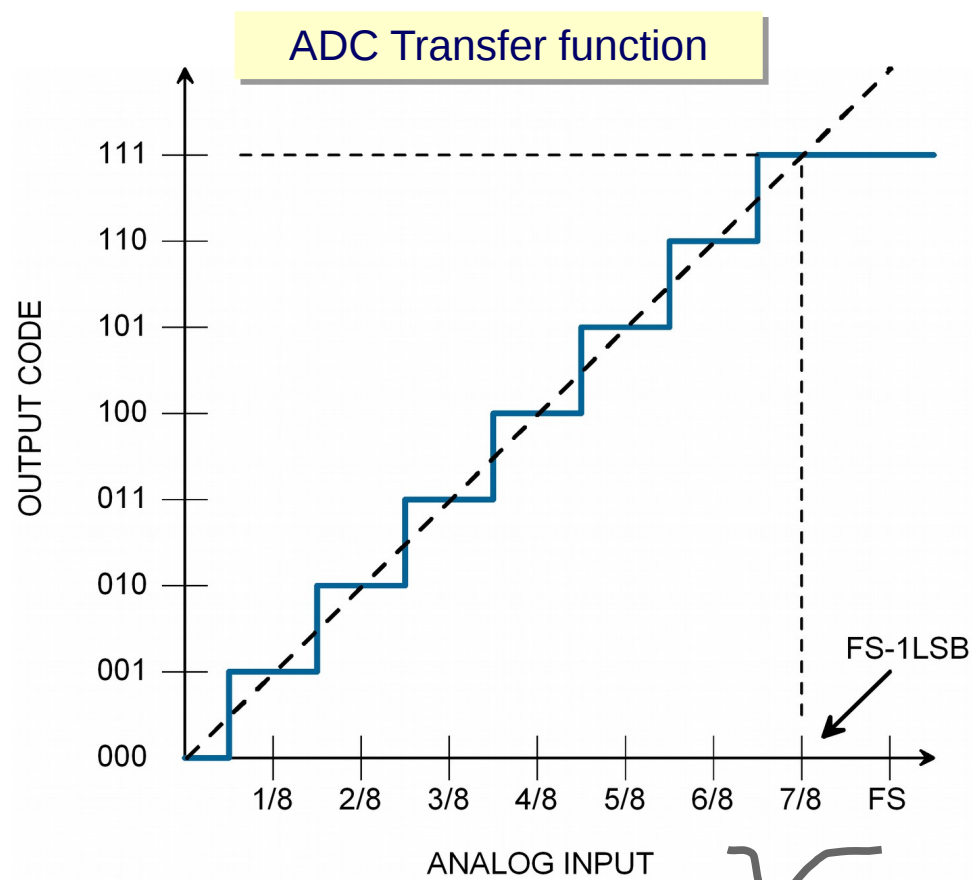
- ADC transfer function
  - Output code vs analog input



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  - Output code vs analog input

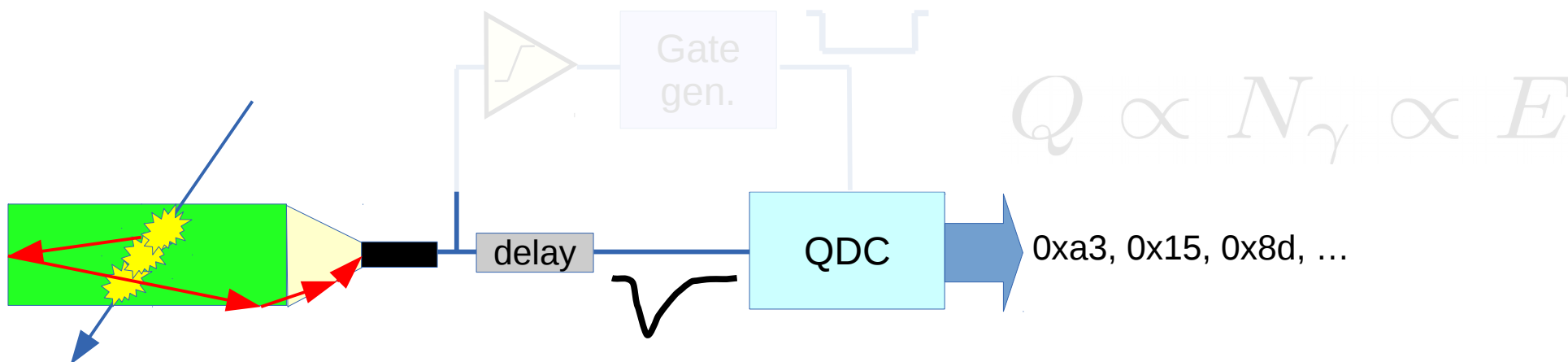


- ADC transfer function
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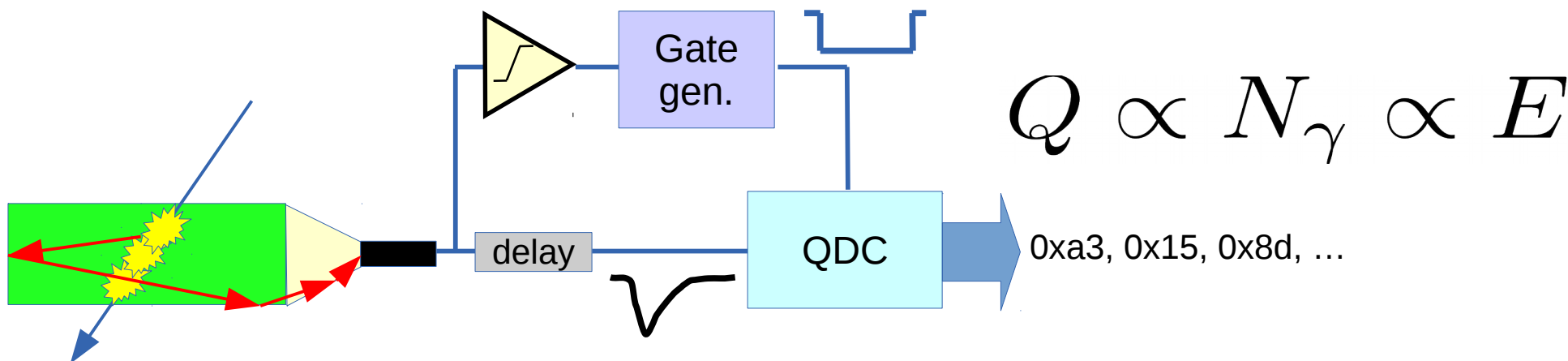
- ADC converts a voltage into a digital representation
  - However, in our experiment, we have a current and we are interested in the total charge
- We need a **QDC** (Charge to Digital Converter)
  - Essentially an integration step followed by an ADC
  - Integration requires limits → gate

$$I = \int_a^b f(x) dx$$



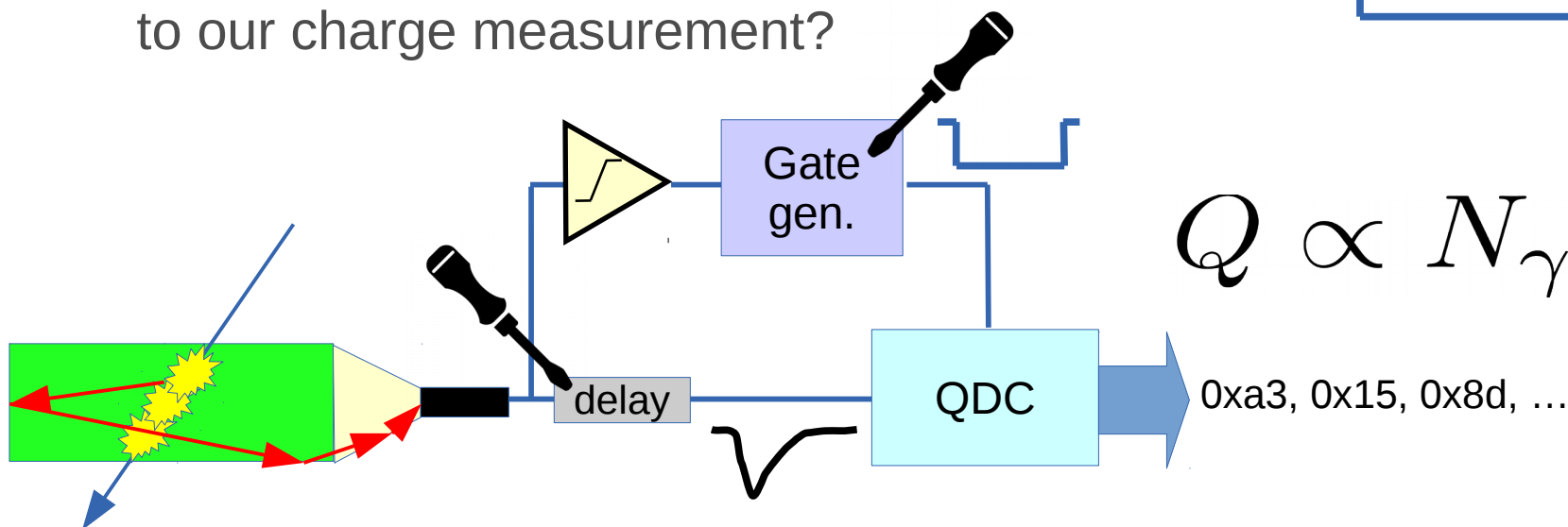
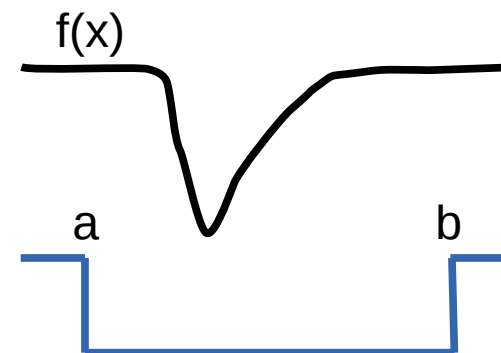
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$$I = \int_a^b f(x) dx$$



- Relative timing between signal and gate is important
  - Delay tuning
- Gate should be **large enough** to contain the full pulse and to accommodate for the jitter
  - Fluctuations are always with us!
- Gate should **not** be **too large**
  - Increases the noise level
  - By the way, which is the noise contribution to our charge measurement?

Labs 2, 3, 4



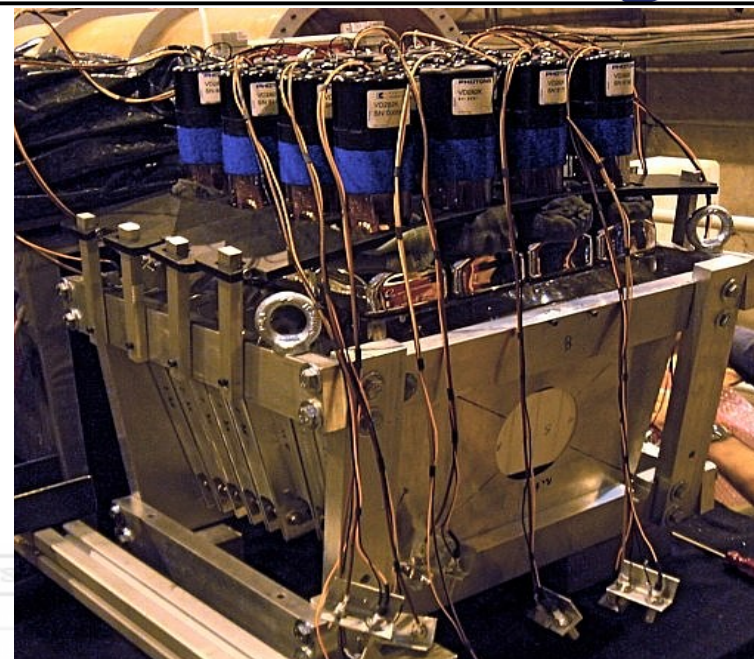
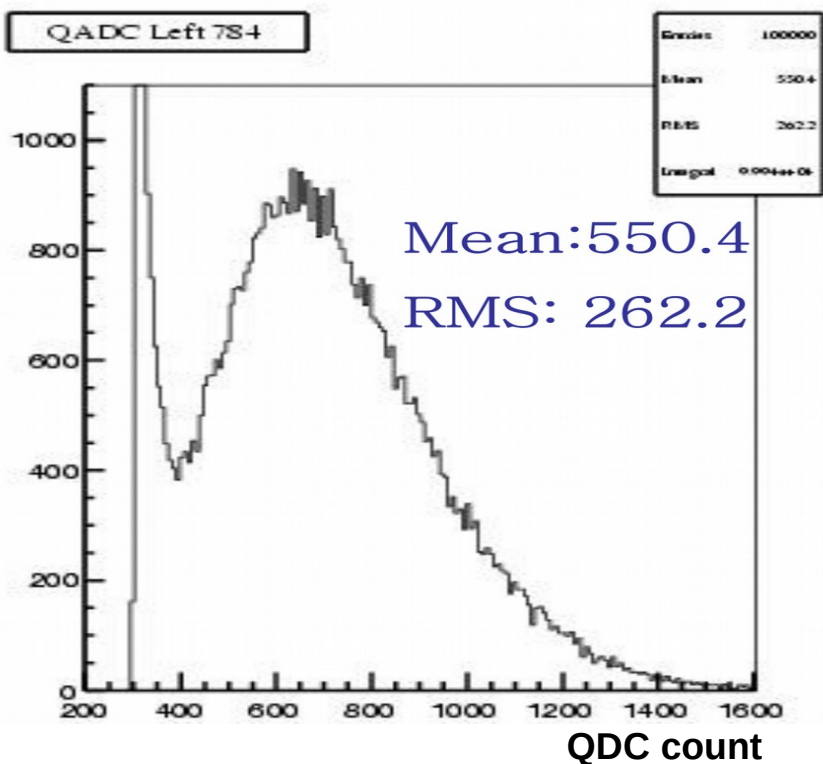
$$Q \propto N_{\gamma} \propto E$$

# Example: QDC spectrum

- QDC spectra from data taken during a test beam @CERN (calorimetry R&D)

$$Q \propto N_{\gamma} \propto E$$

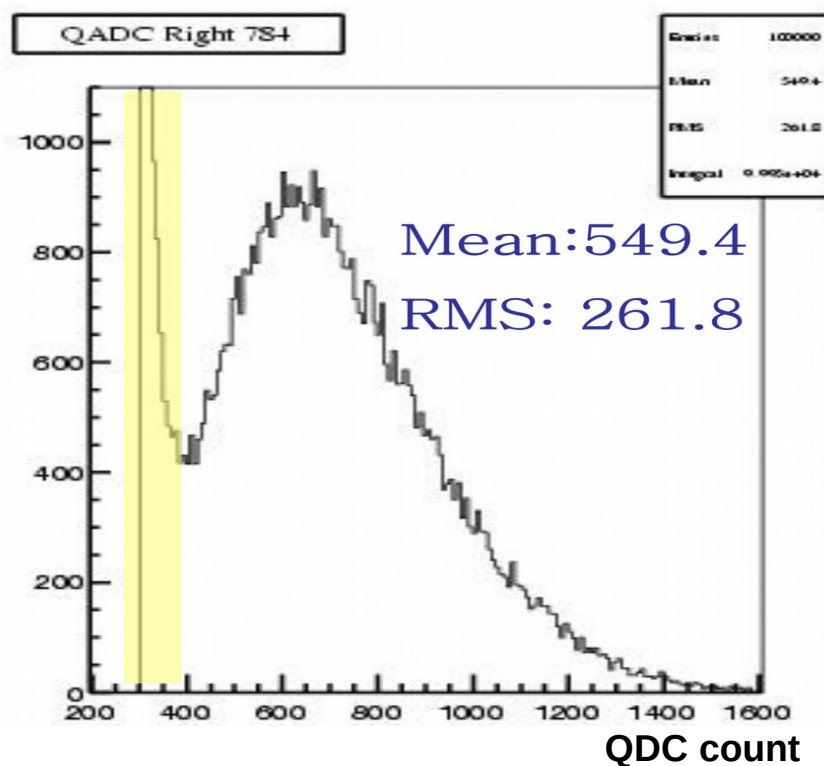
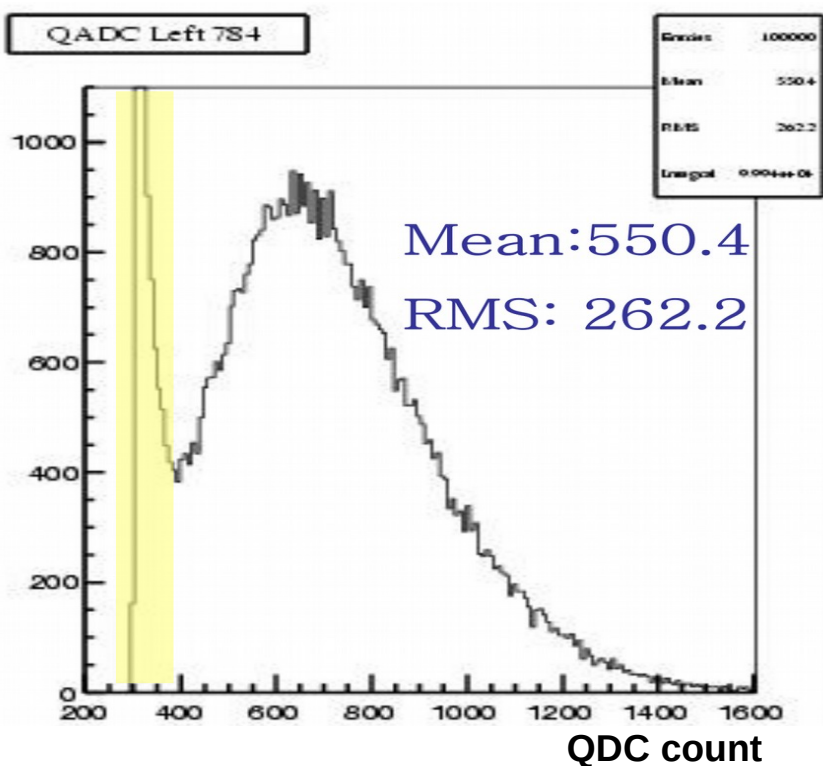
- But, what is the 1<sup>st</sup> peak?



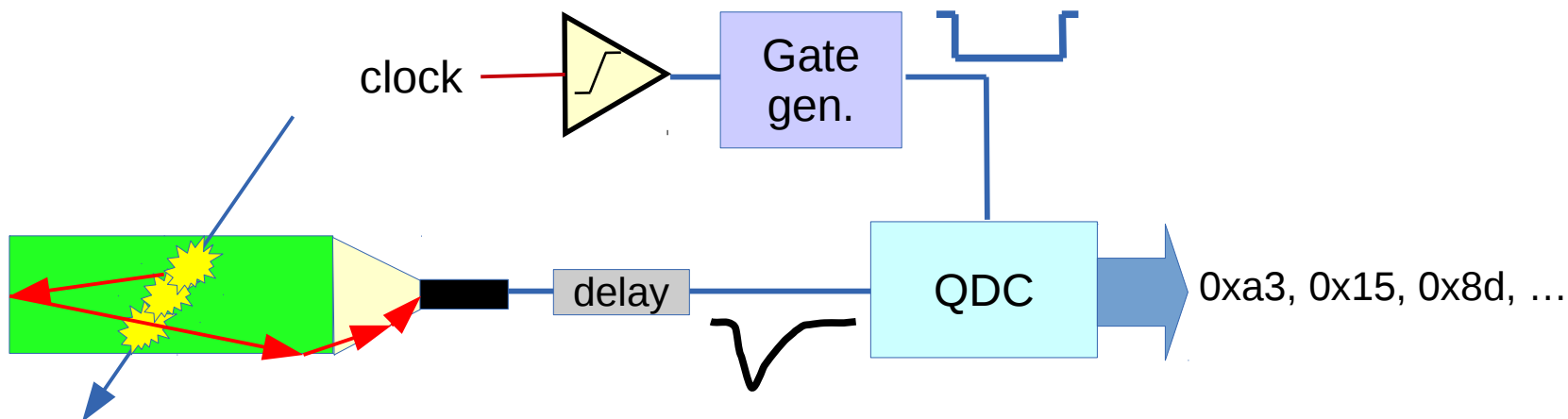
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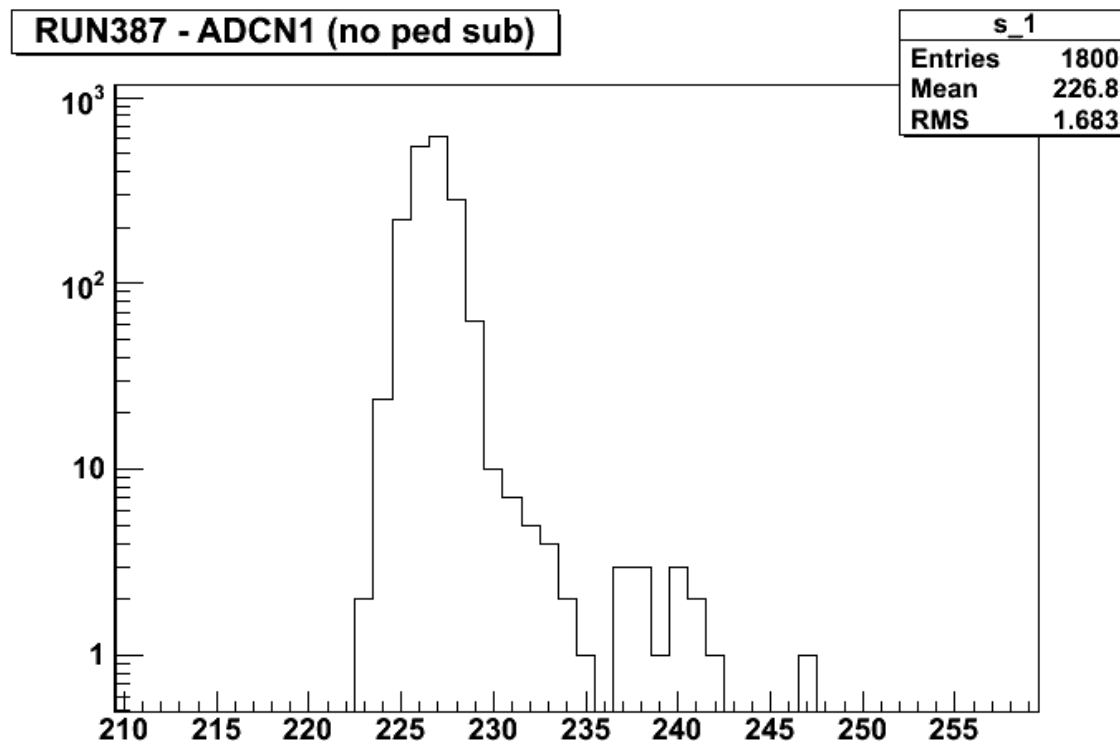
- But, what is the 1<sup>st</sup> peak? How can we estimate it?



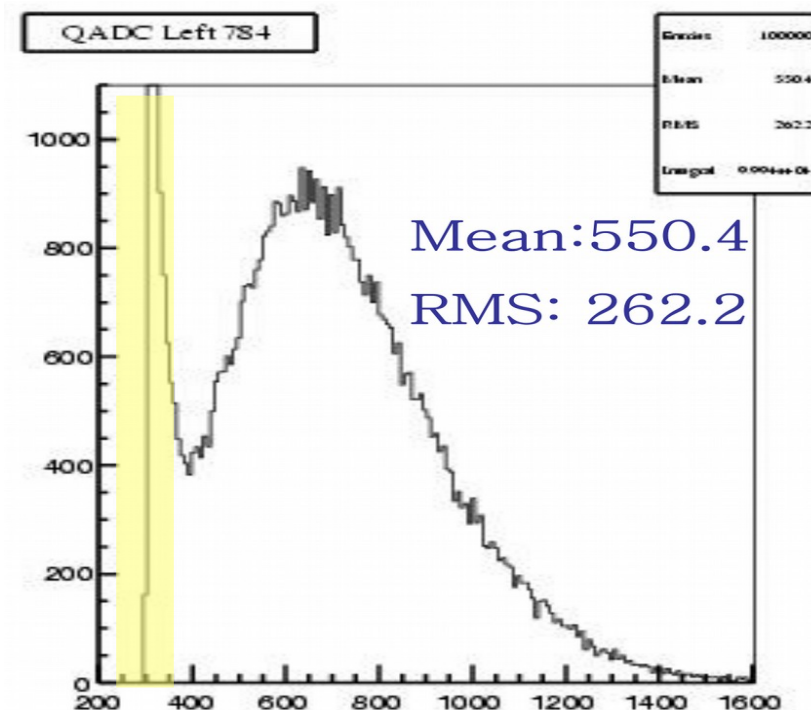
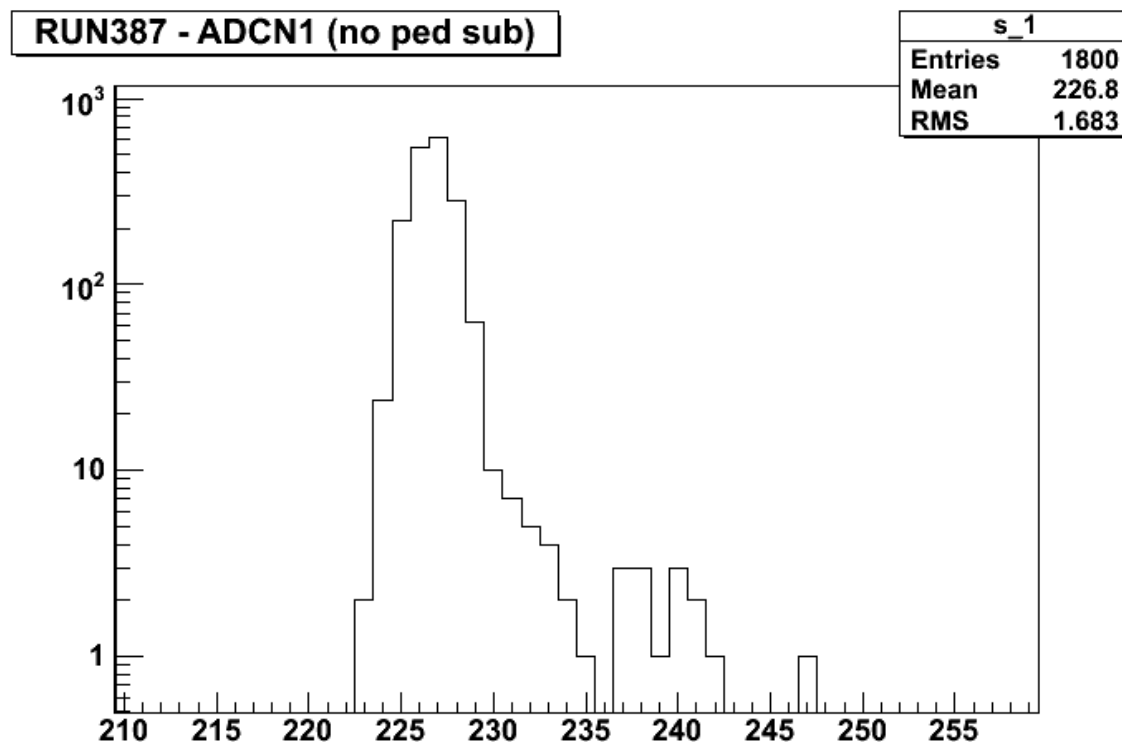
- The **pedestal** can be measured with an out-of-phase trigger
  - PMT dark current, thermal noise, ...
  - The same noise enters our physics measurements and contributes with an offset to the distribution
- The result of a pedestal measurement has to be subtracted from our charge measurements



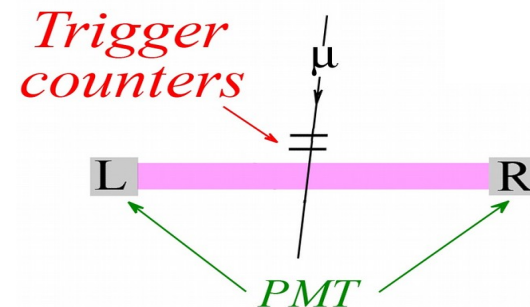
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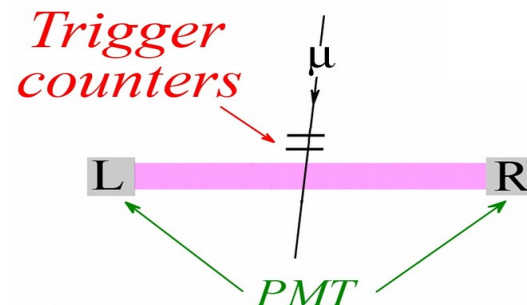
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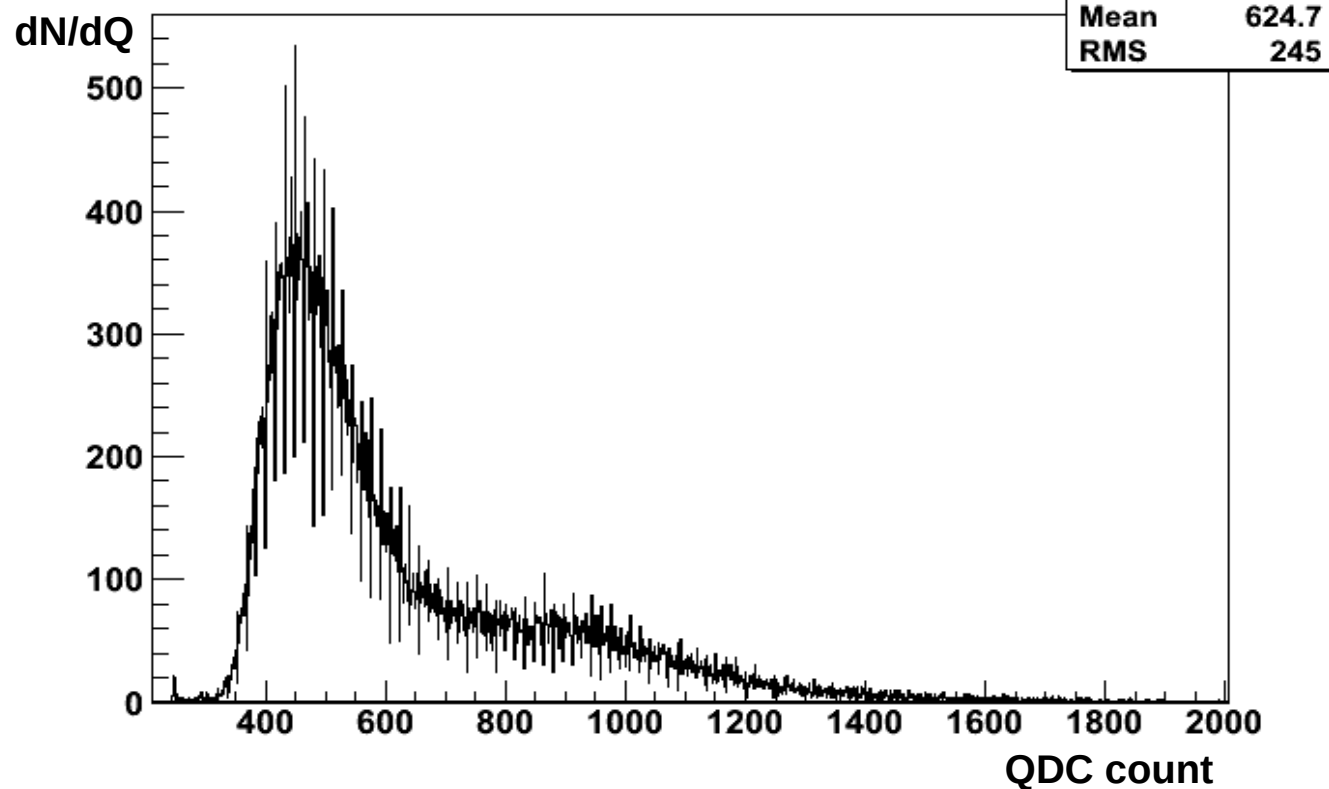
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  - Real data from a test beam @CERN



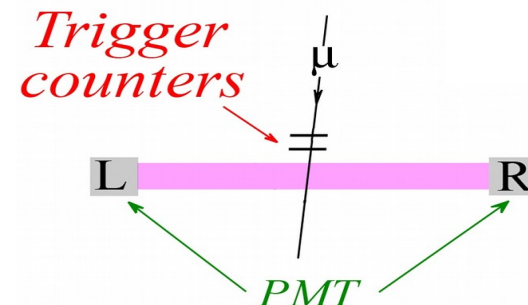
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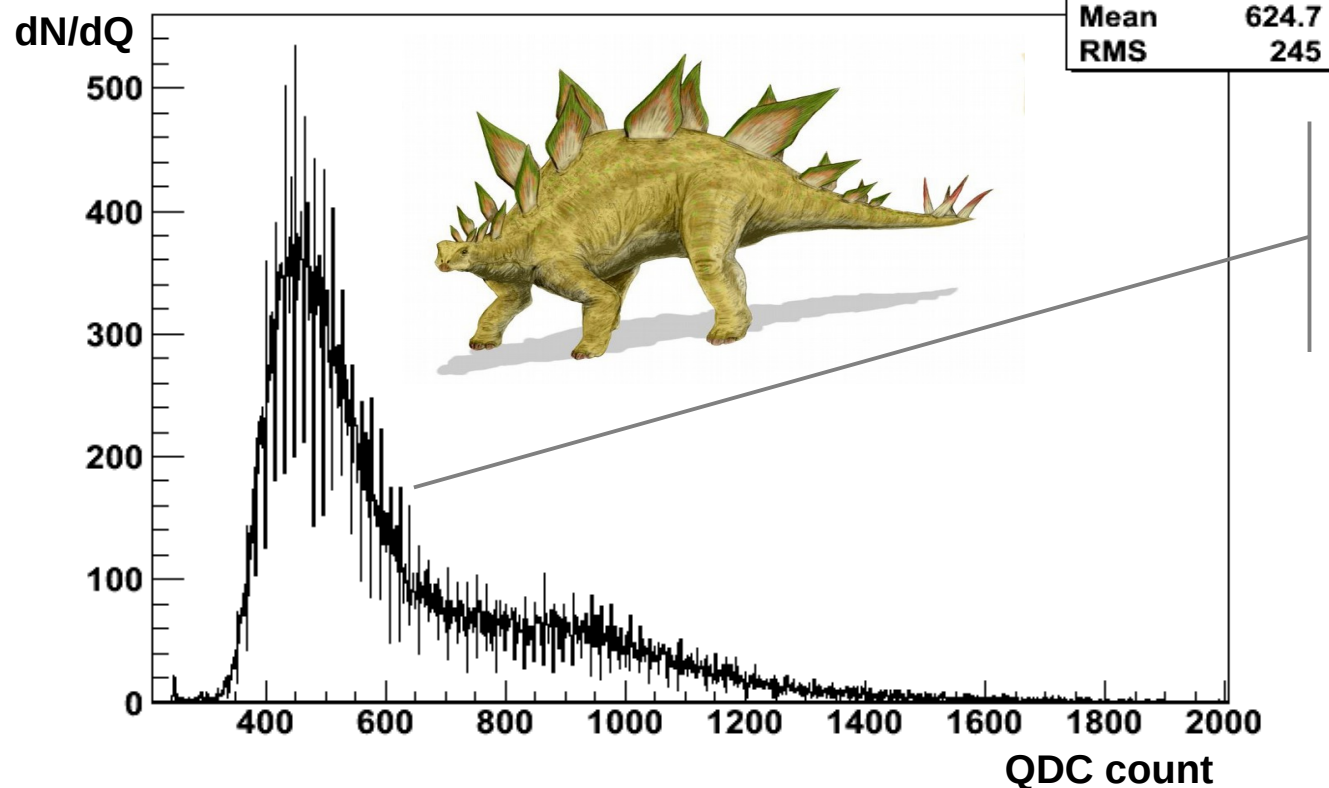
$\pi$ -beam charge-distribution for one PMT



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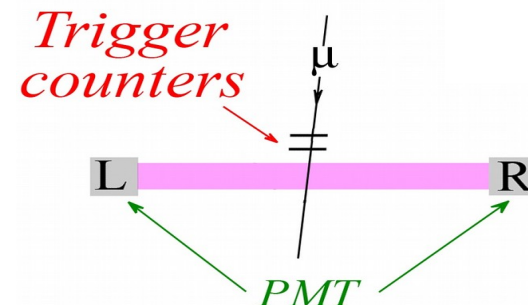
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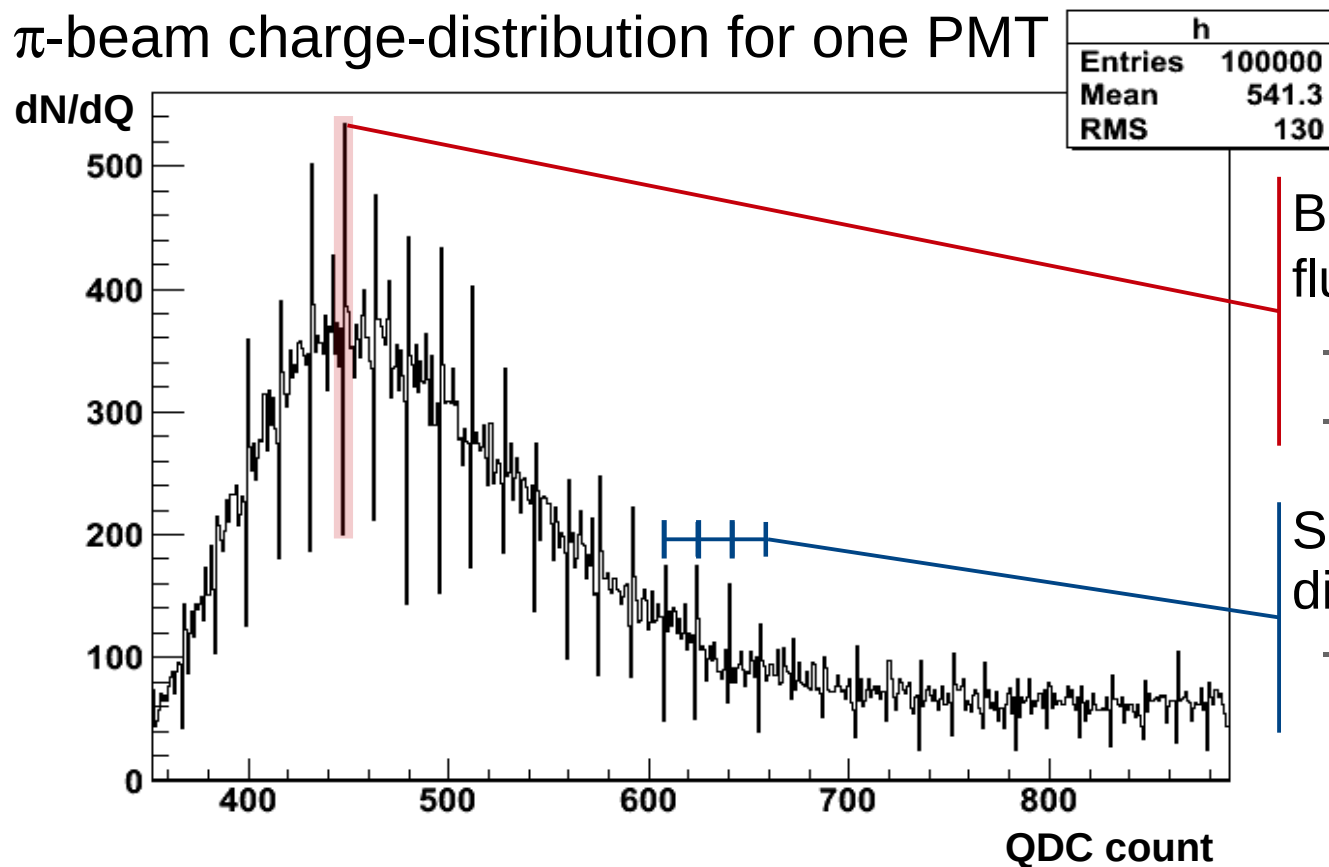
But, what are all those little peaks? Just statistical fluctuations?

Let's zoom in!

- PbWO<sub>4</sub> (scintillating) crystal equipped with two PMTs and exposed to e, $\mu$  and  $\pi$  beams
  - Real data from a test beam @CERN



$\pi$ -beam charge-distribution for one PMT



Bin with N entries can fluctuate with  $\sigma = \sqrt{N}$

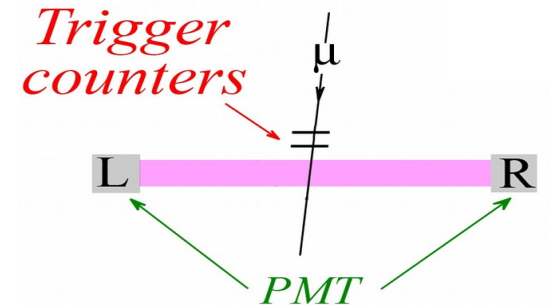
- expected  $\sigma = \sqrt{N} 360 \sim 19$
- observed  $\sim 200$  (**10  $\sigma$** )

Spikes are regularly distributed

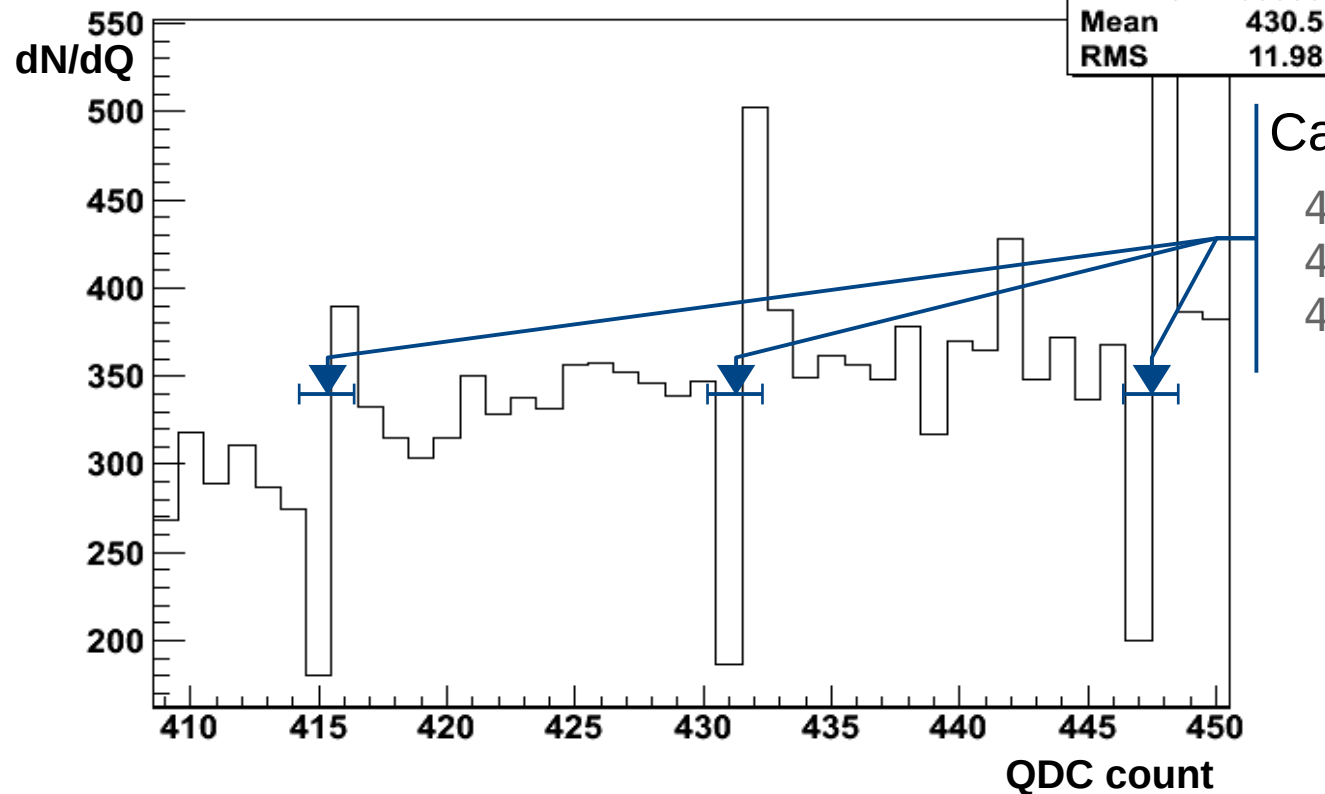
- Some systematic effect must be taking place

**Let's zoom in!**

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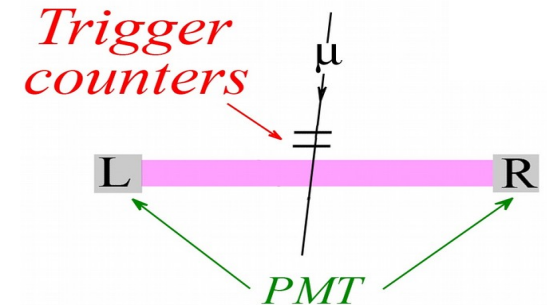
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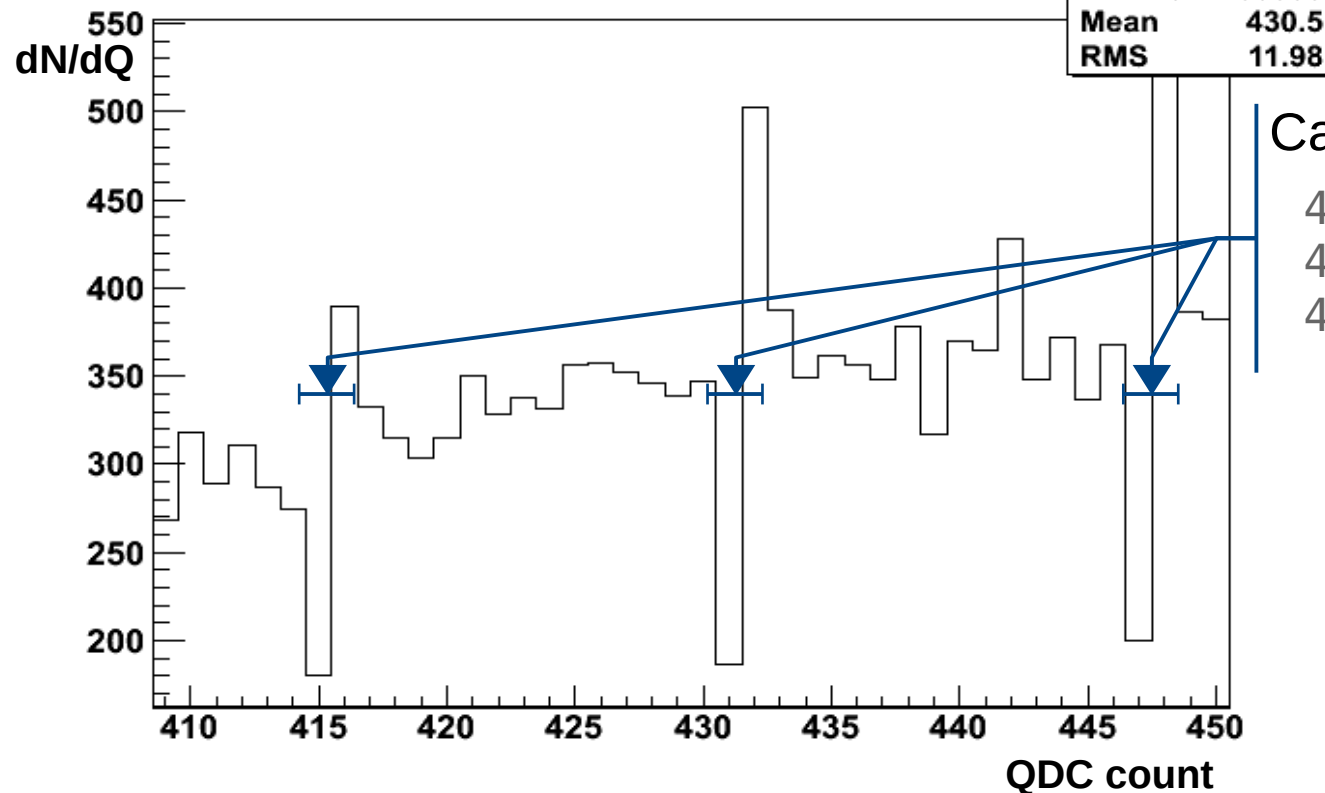
Can you see the effect?

415 & 416  
431 & 432  
447 & 448

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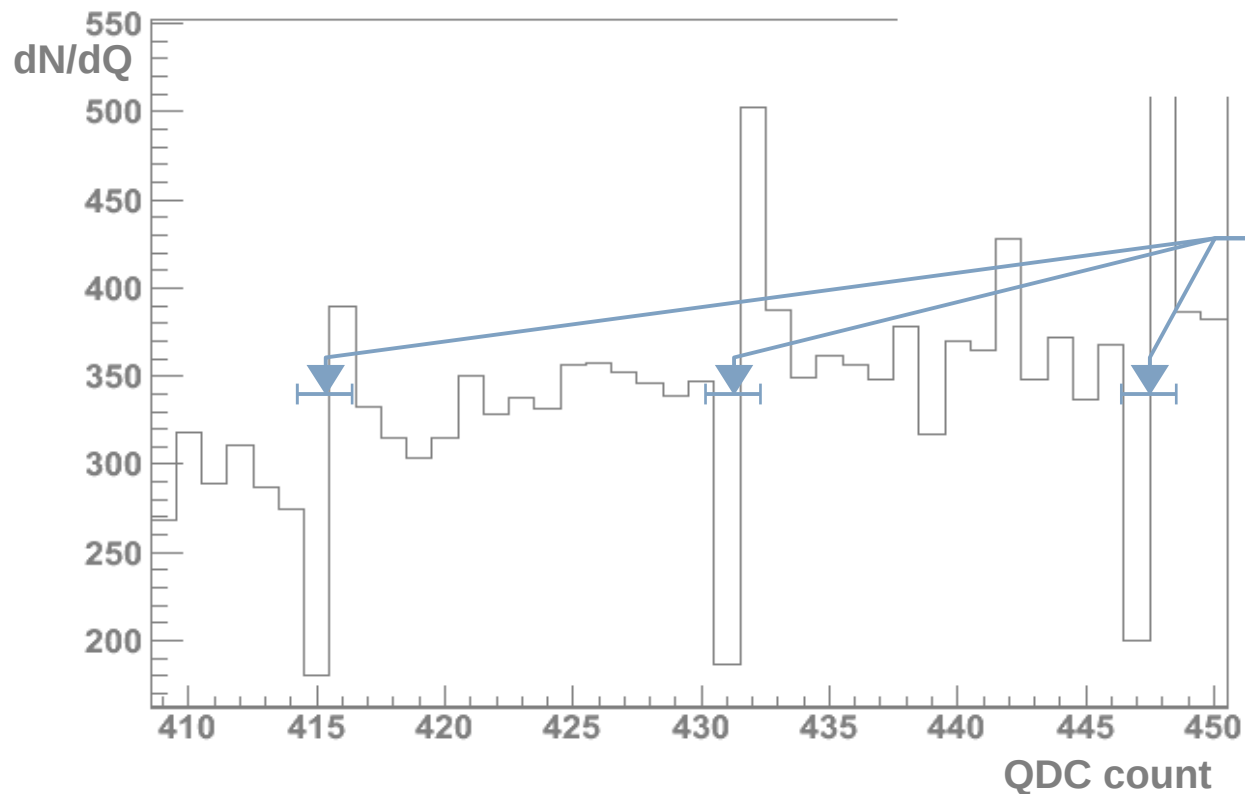


Can you see the effect?

415 & 416  $\rightarrow$  0x19**f** & 0x1A**0**  
 431 & 432  $\rightarrow$  0x1A**f** & 0x1B**0**  
 447 & 448  $\rightarrow$  0x1B**f** & 0x1C**0**

The QDC prefers  
 output of type 0x... $\beta$ **0**  
 in respect of 0x... $\alpha$ **f**  
 where  $\beta = \alpha + 1$

- Which is the simplest way to fix this problem in the data?
  - At which cost?
- Can you understand the module name?
  - Module: 4c6543726f79204c31313832



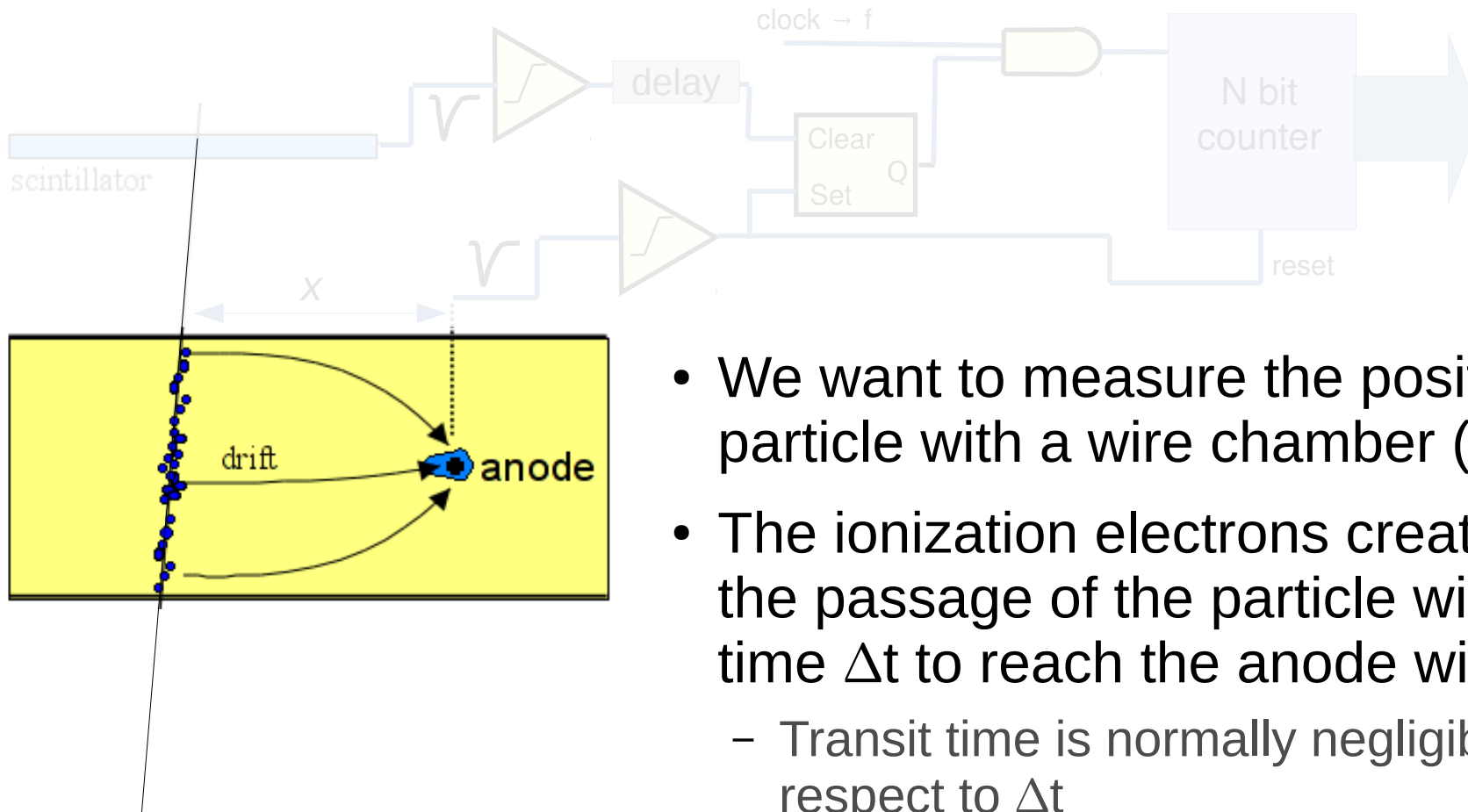
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The QDC prefers  
 output of type 0x...**β0**  
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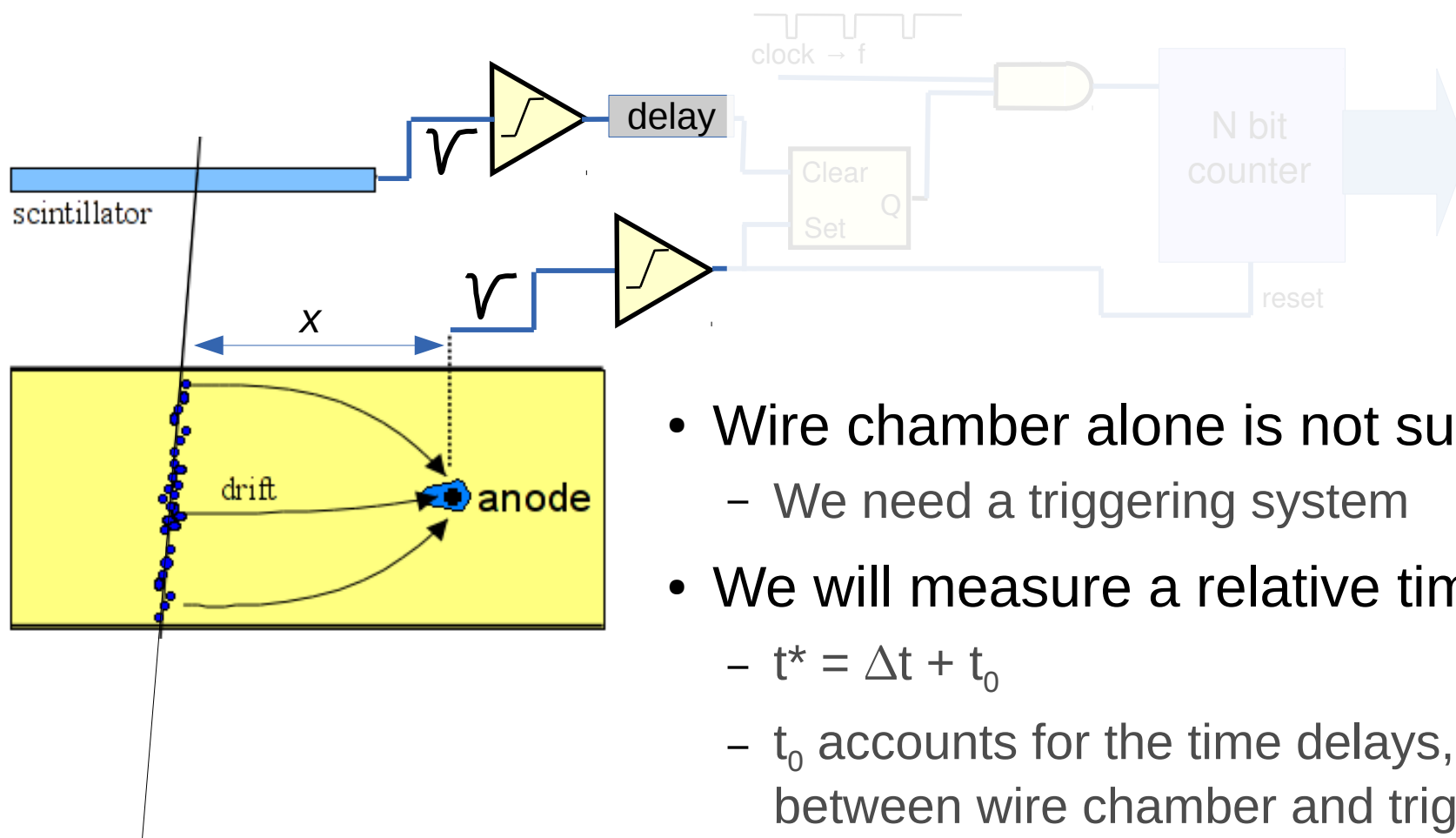
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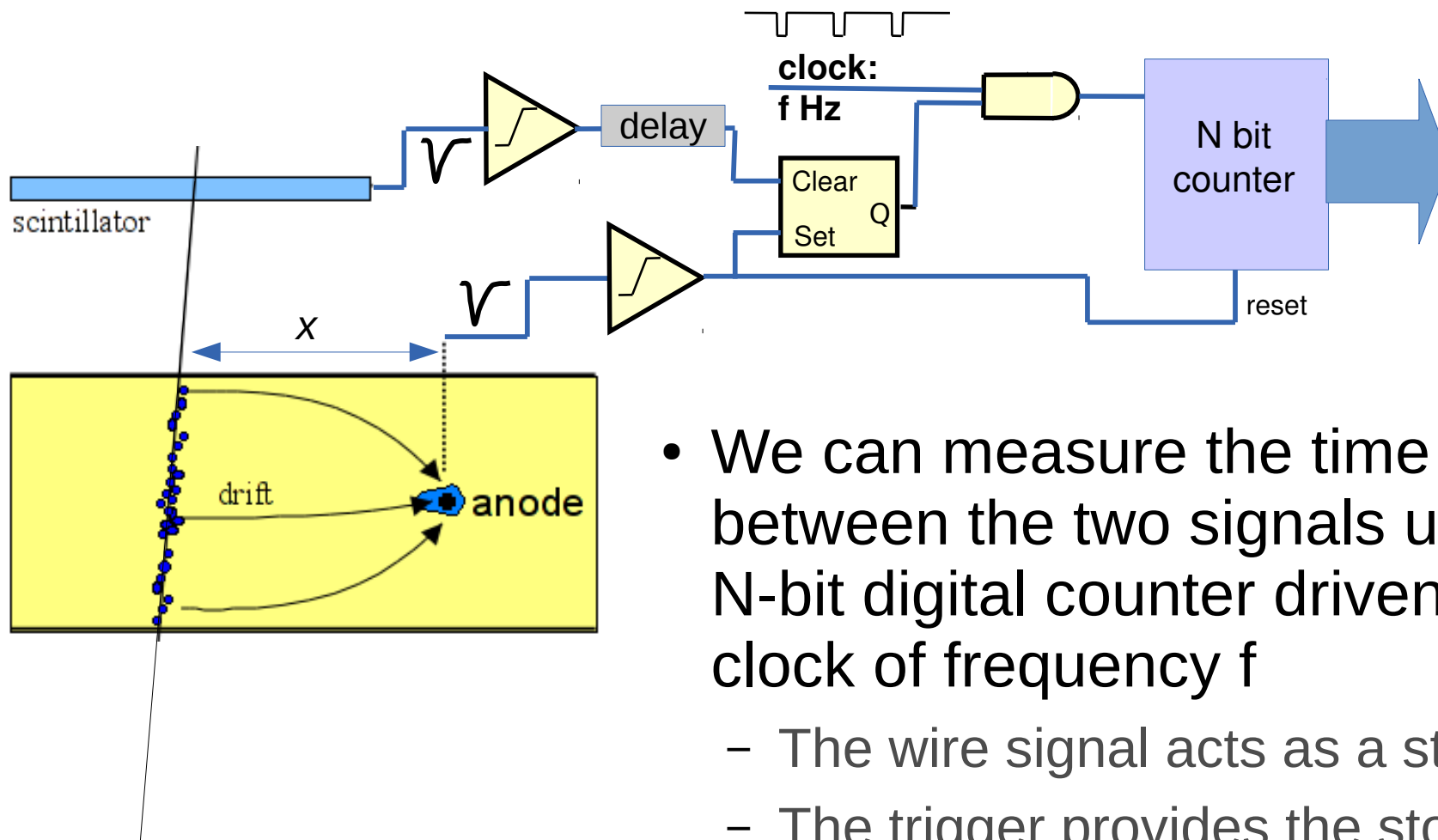
- We want to measure the position of particle with a wire chamber (**drift**)
- The ionization electrons created by the passage of the particle will take a time  $\Delta t$  to reach the anode wire
  - Transit time is normally negligible with respect to  $\Delta t$
  - If we consider a constant drift speed  $v_D$  (e.g.:  $50\mu\text{m/ns}$ ), then position is:

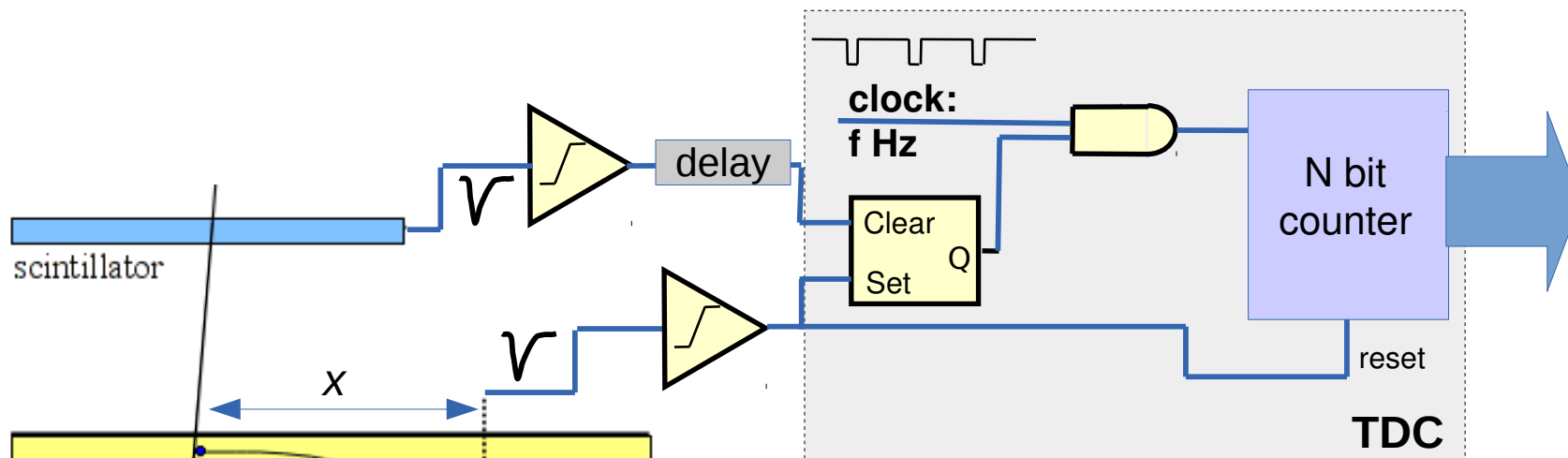
$$x = v_D \cdot \Delta t$$



- Wire chamber alone is not sufficient
  - We need a triggering system
- We will measure a relative time
  - $t^* = \Delta t + t_0$
  - $t_0$  accounts for the time delays, offsets, ... between wire chamber and triggering system
- Assuming a constant drift

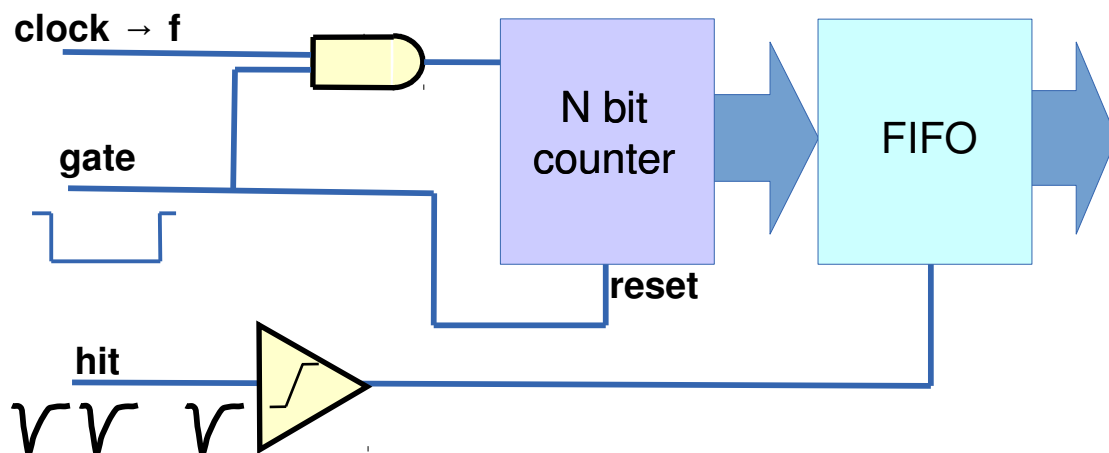
$$x = \alpha t^* + \beta$$



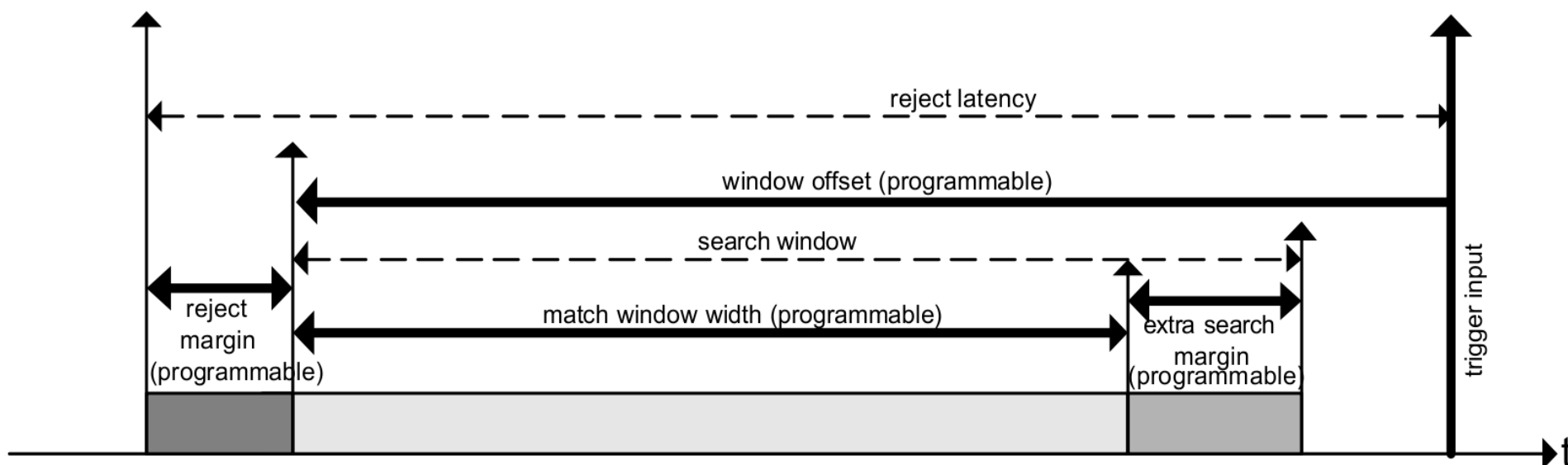


- This device is a **TDC**: Time-to-Digital Converter
  - Resolution:  $1/f$
  - Dynamic range:  $N$
- Single hit TDC
  - if a noise spike comes just before the signal, the measure is lost

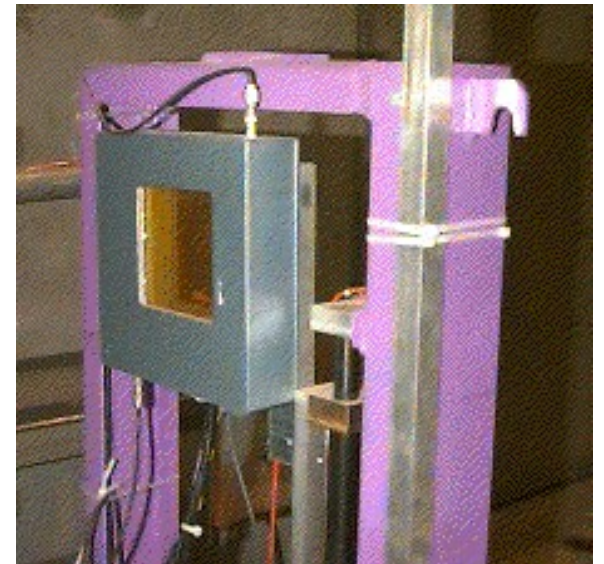
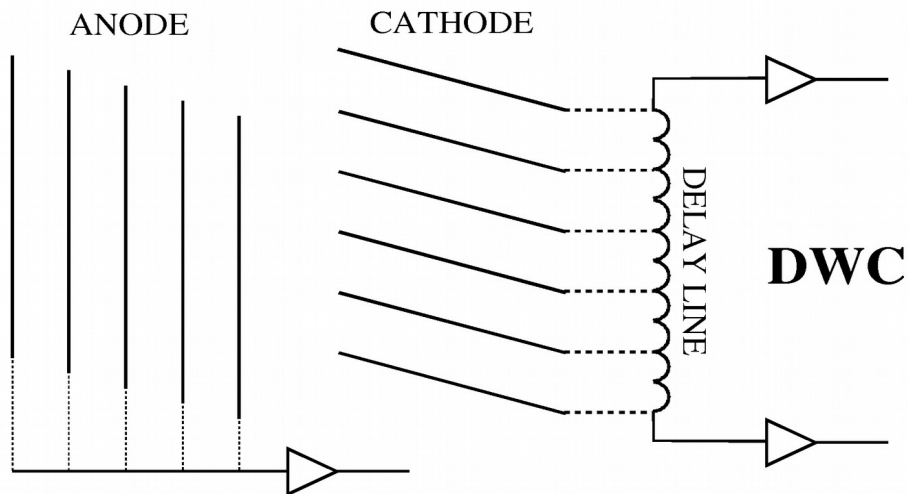
- Gate resets and starts the counter
  - It also provides the measurement period
- Each “hit” (i.e. signal) forces the FIFO to load the current value of the counter, that is the delay after the gate start
  - Common-start configuration
  - In order to distinguish between hits belonging to different gates, some additional logic is needed to tag the data



- Real TDCs provide advanced functionalities for fine-tuning the hit-trigger matching
  - Internal programmable delays
  - Internal generation of programmable gates
  - Programmable rejection frames
  - Usually via a dedicated C library/API

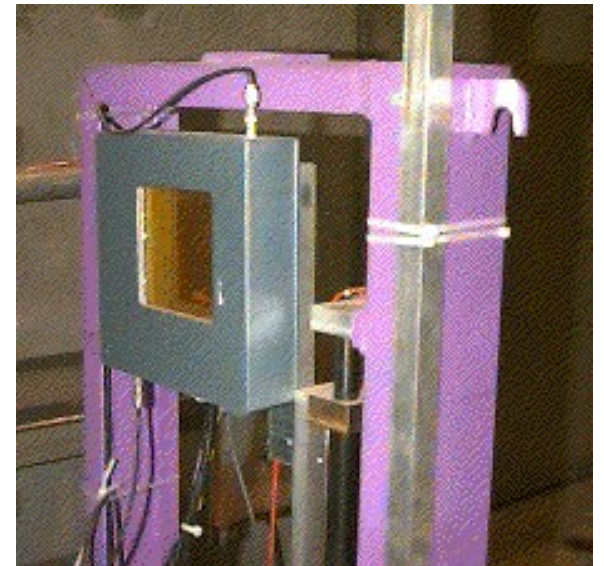
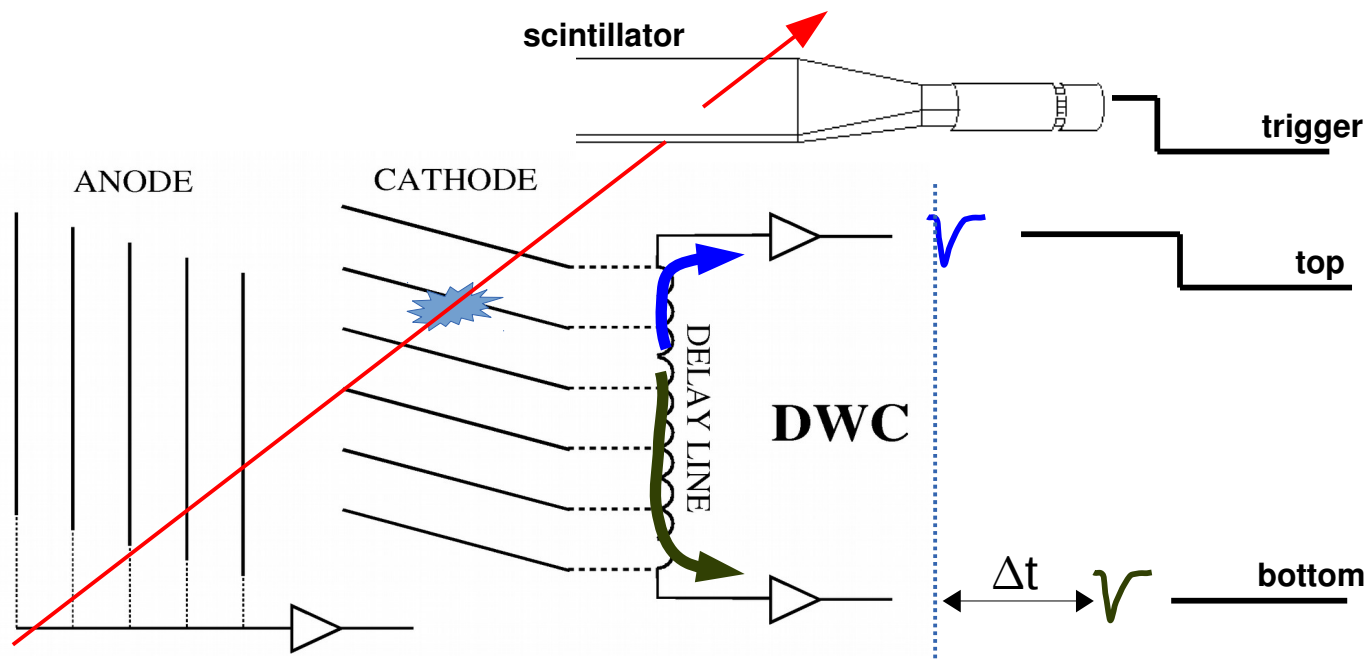


- XDWC: delay wire chambers
  - used on the SPS extracted lines to measure beam profiles
- Two cathode planes provide X and Y positions
  - Measurement based on the delay gained along a delay line

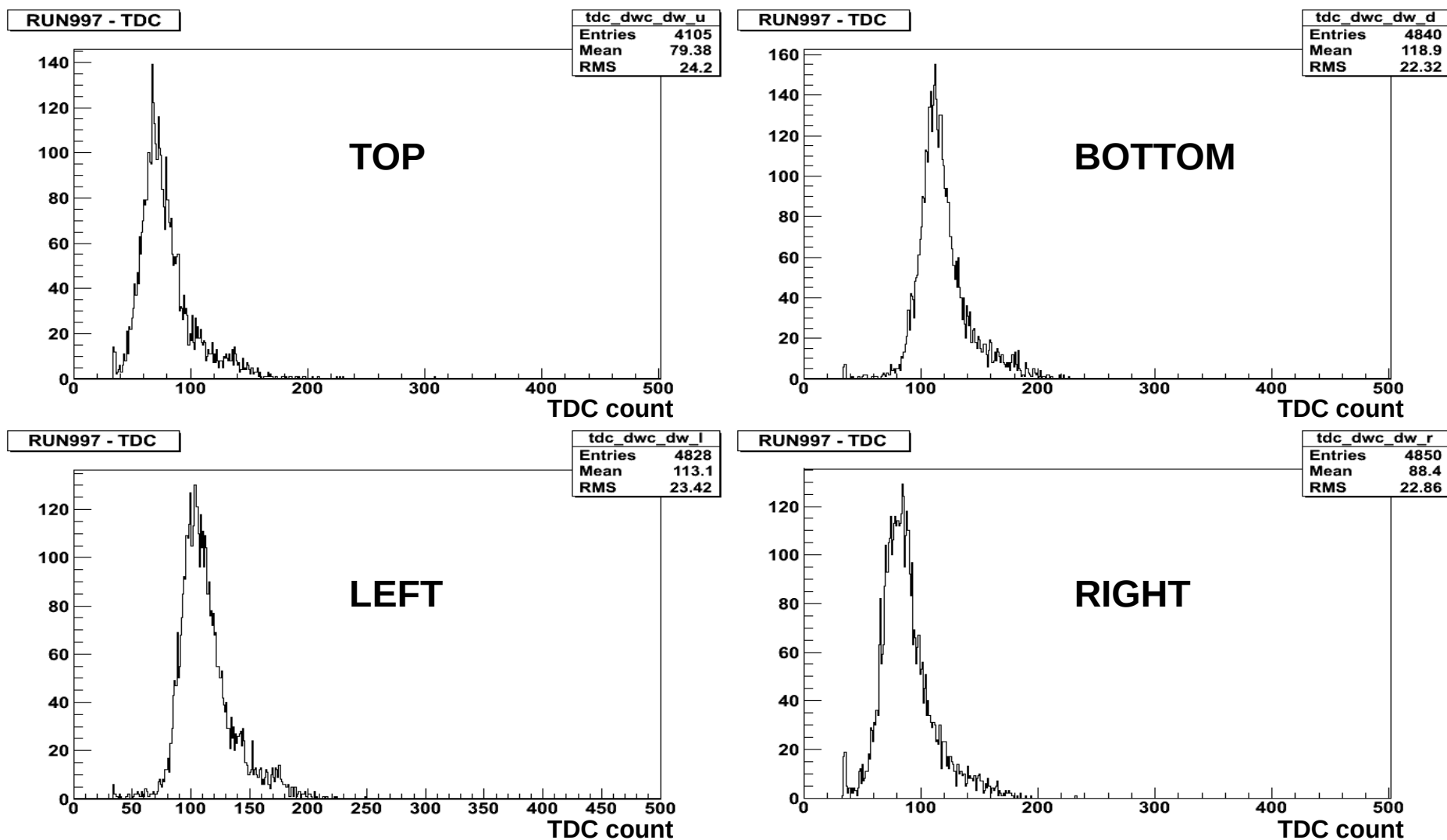


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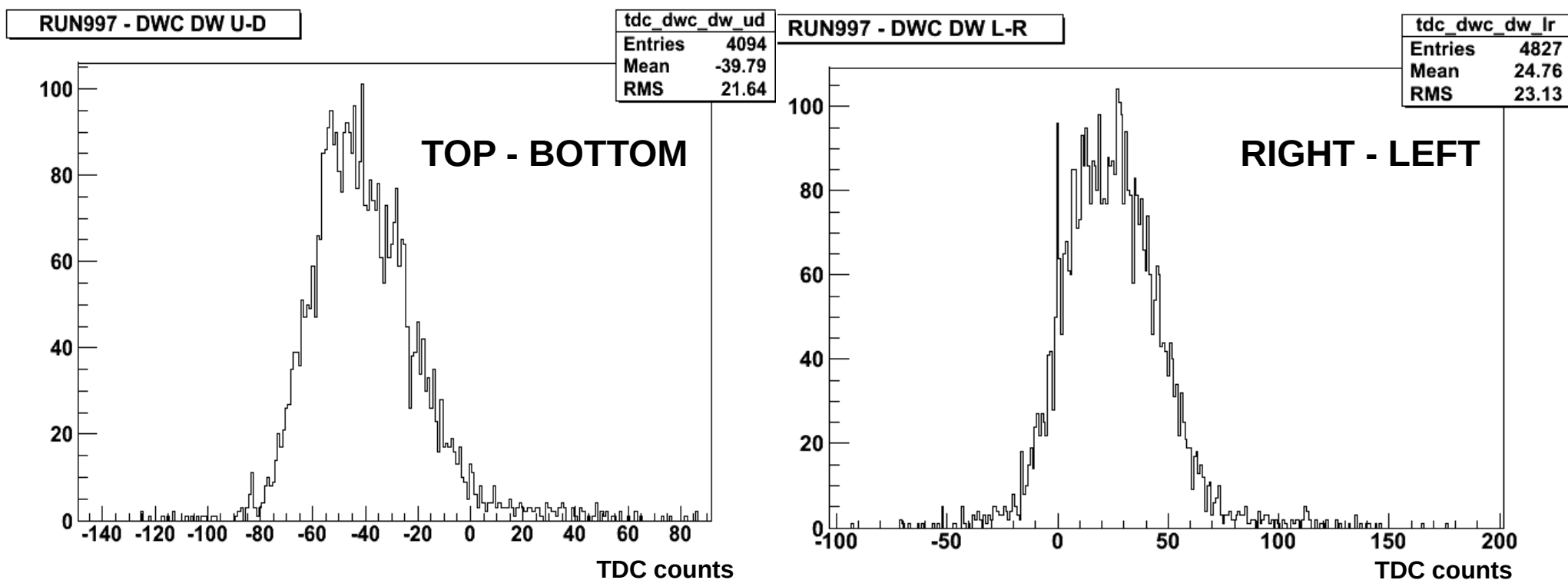
$$y = \alpha \cdot \Delta t + \beta = \alpha \cdot (t_{top} - t_{bottom}) + \beta$$



- Take a run (some thousands events)
  - Individual channel distribution



- Beam sizes are still in TDC counts
  - Not very useful, though
  - How do we convert this into a known scale (e.g. cm)?

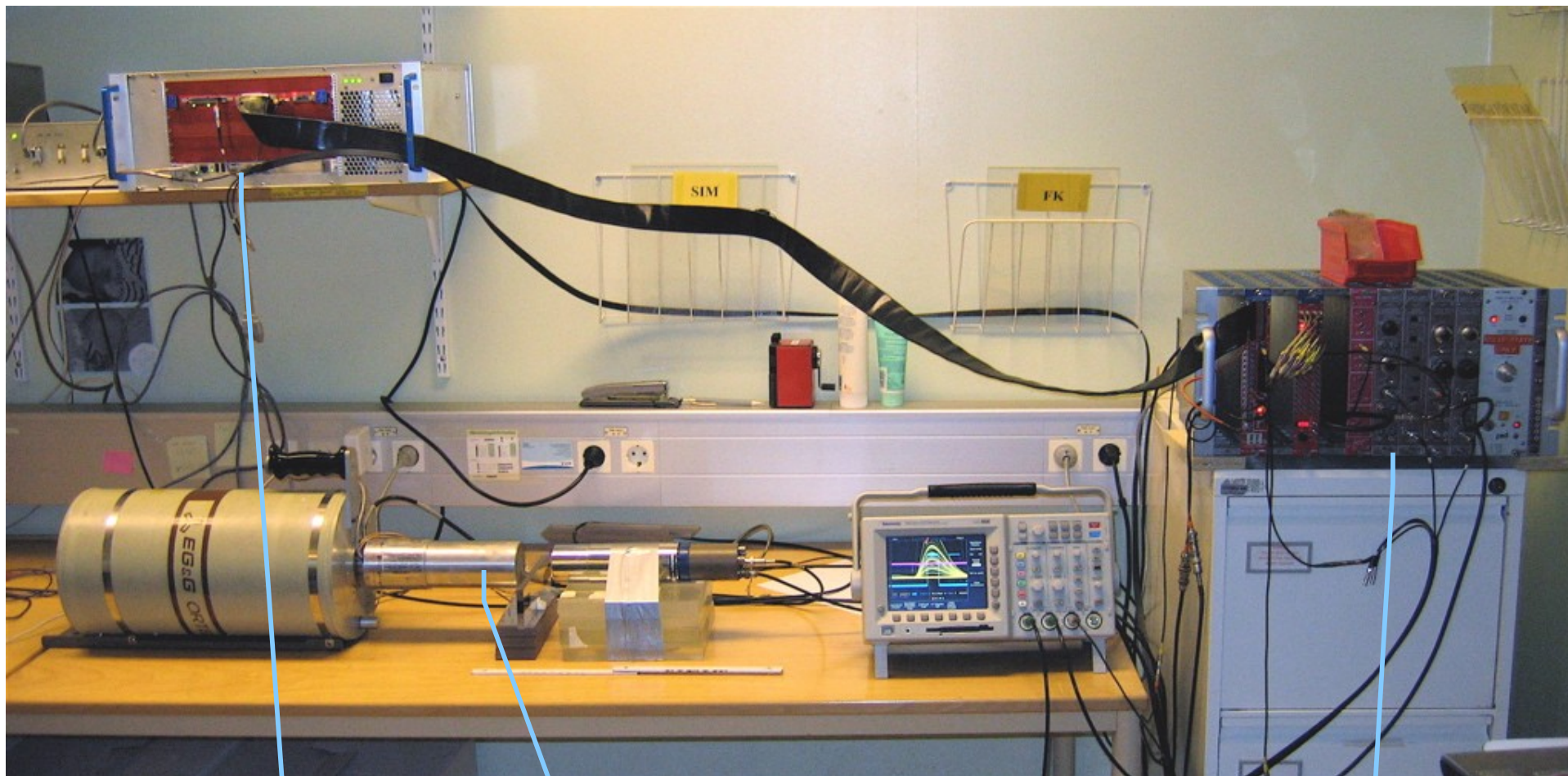


- Introduction
- Measure energy deposition
  - Scintillator setup
  - Photomultiplier
  - Analog-to-Digital conversion
  - Charge-to-Digital conversion
  - QDC in real life
- Measure position
  - Wire chamber setup
  - Time-to-Digital conversion
  - TDC in real life
- Corollary: calibration



- Both experiments provide relative measurements
  - Values obtained via our systems are in some (known) relation with the interesting quantities
    - Scintillator  $Q \propto N_\gamma \propto E$
    - MDWC  $y = \alpha \cdot \Delta t + \beta = \alpha \cdot (t_{top} - t_{bottom}) + \beta$
- Our instruments need to be **calibrated** in order to give us the answer we are looking for
  - We have to determine the parameters that transform the raw data into a physics quantity
  - The parameters normally depend on the experimental setup (e.g. cable length, delay settings, HV settings, ... )
- NB: calibration mechanisms/procedures shall be foreseen in the design of our detector and DAQ

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Crystal HPGe

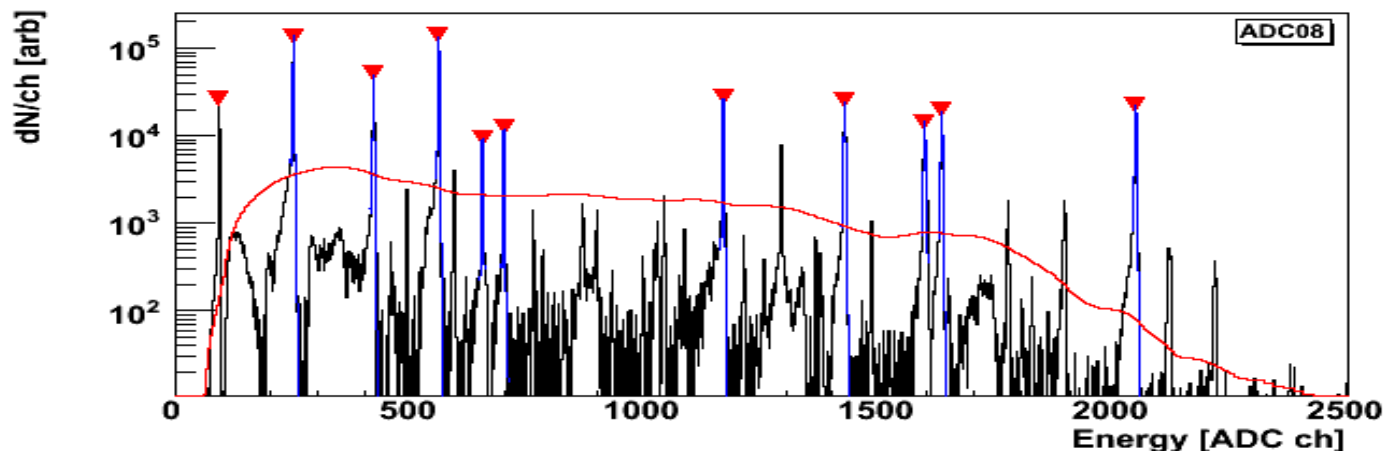
Readout (ADC)

Trigger and front-end

by Sergio Ballestrero

- $^{152}\text{Eu}$  reference source allows for definition of the parameters describing functional relation between ADC count and E
  - Known  $\gamma$  emission lines
- Find the peaks and fit

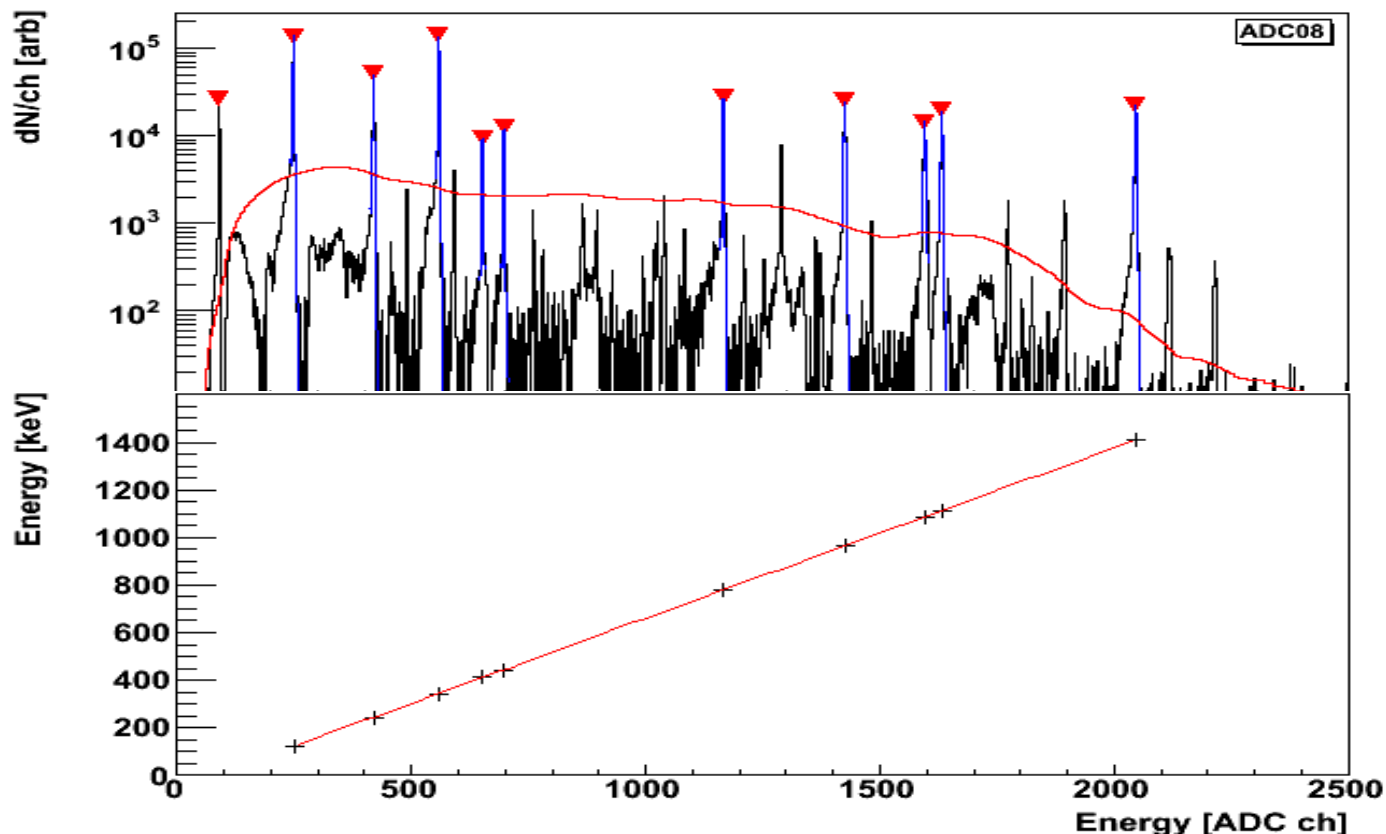
$$Q \propto N_{\gamma} \propto E$$



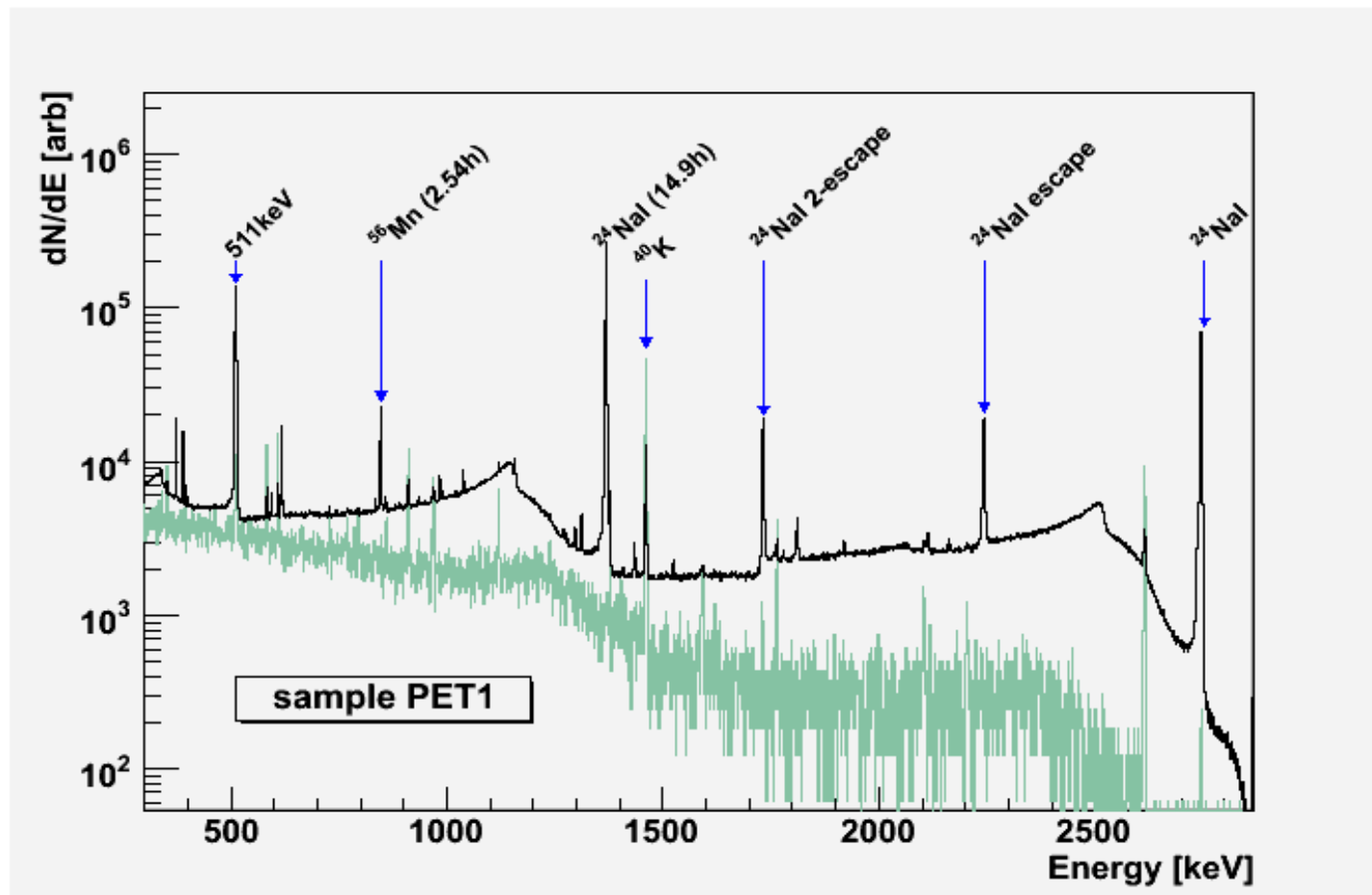
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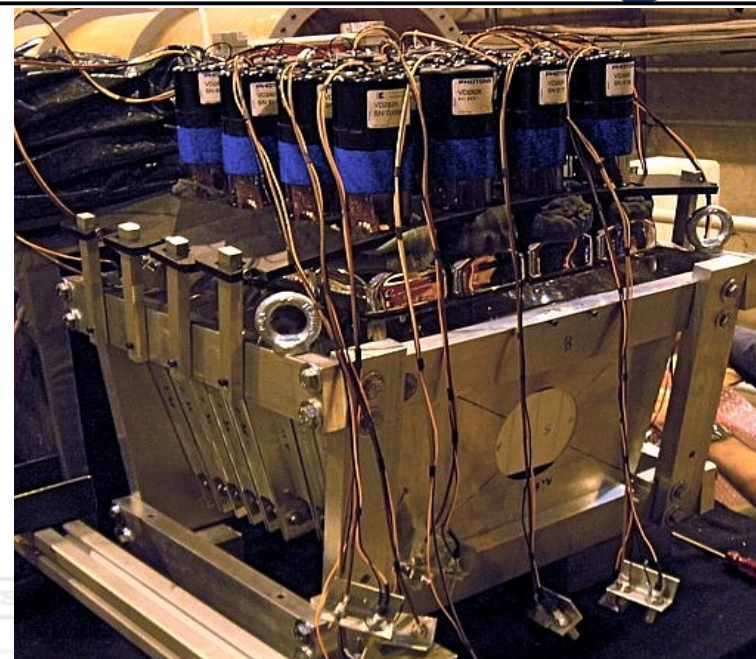
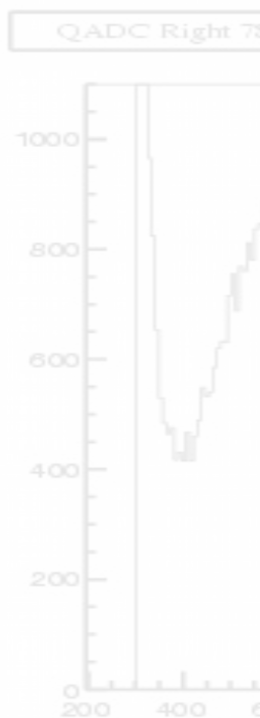
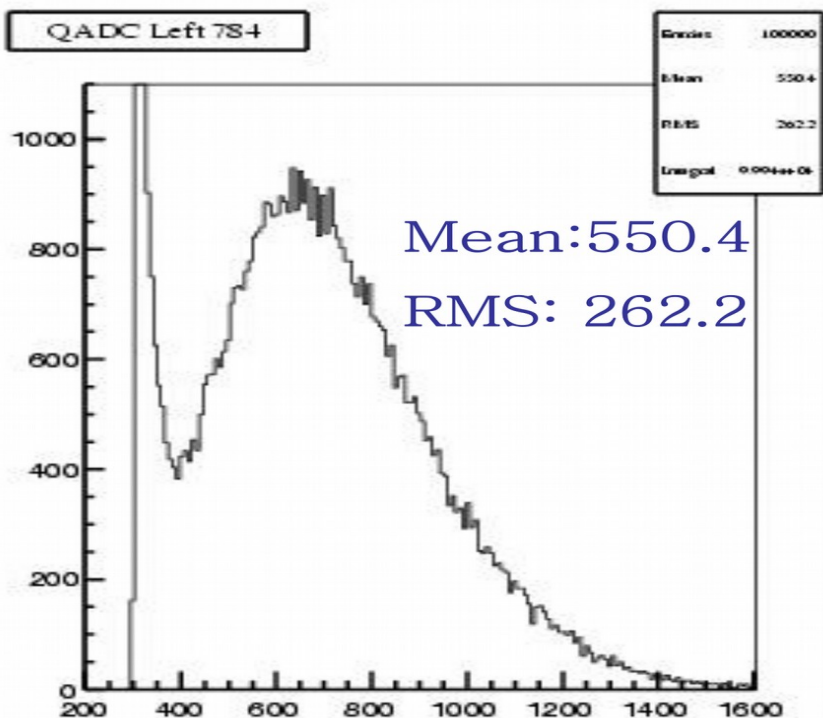


- Calibrate crystal setup can be used to identify isotopes generated in  $\gamma$ -irradiated samples

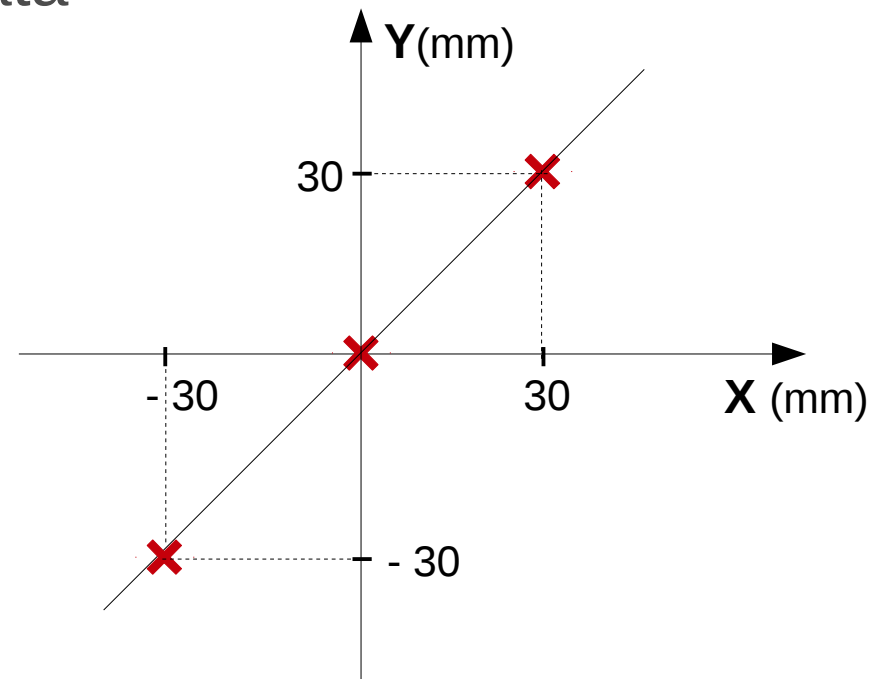


by Sergio Ballestrero

- It is more complicated
  - The beam energy is known
  - But MC is needed to take in account geometry, acceptance, beam profile, etc



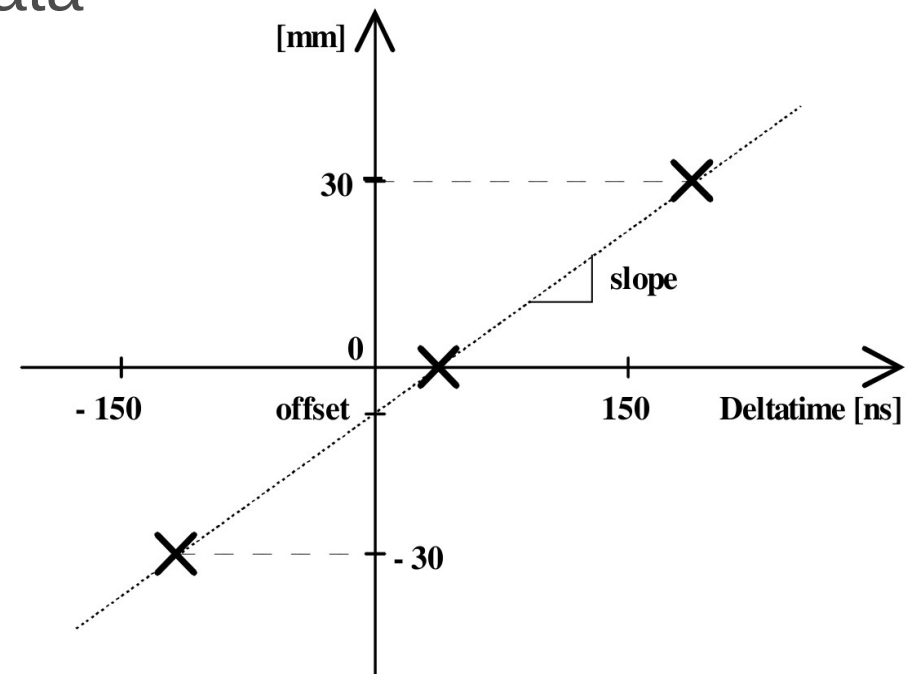
- XDWC chamber have 3 calibration inputs
  - allow for independent calibrations of X and Y axes with only 3 different sets of data
  - Calibration input simulate signals from particles respectively hitting
    - Right-top ( $X=Y=30\text{mm}$ )
    - Center ( $X=Y=0\text{mm}$ )
    - Left-bottom ( $X=Y=-30\text{mm}$ )
  - Interpolating the three points in t-x space, the parameters of the calibration equation can be measured



$$x = \alpha t^* + \beta$$

- Calibration shall be done with final setup and TDC

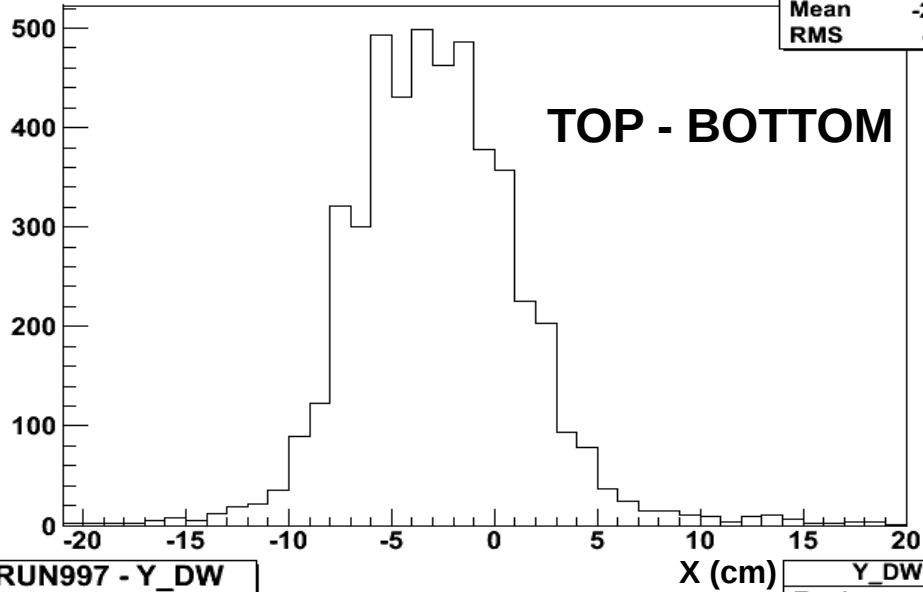
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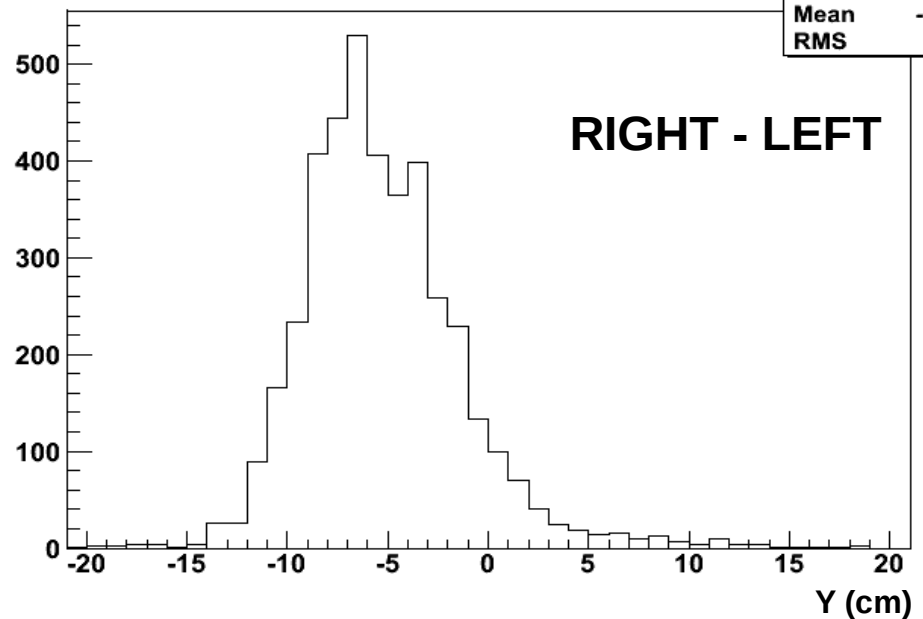
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RUN997 - X\_DW

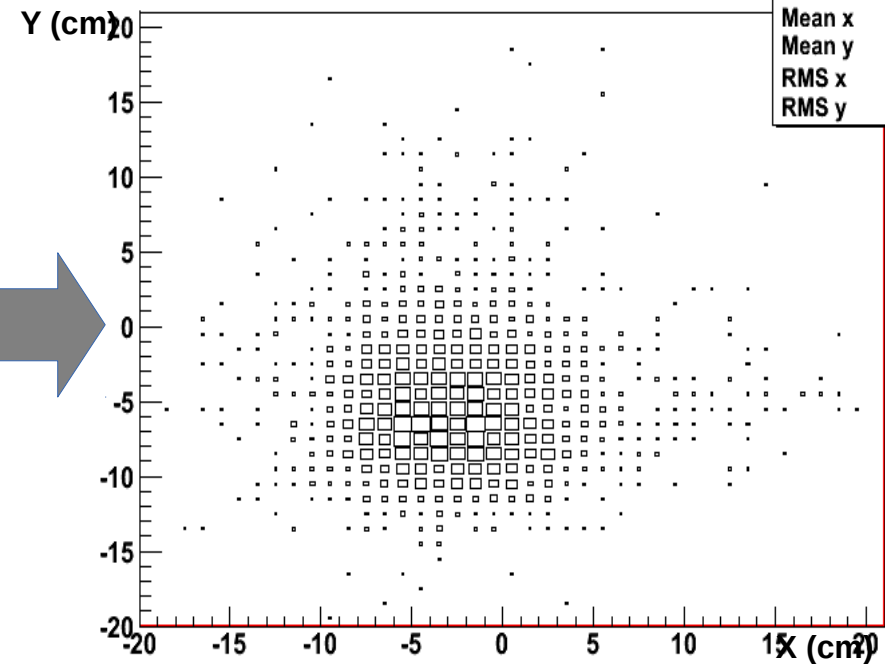


RUN997 - Y\_DW



## Beam profile

RUN997 - Y\_DW vs X\_DW



- Digitization techniques produce data directly manageable by digital systems (e.g. a computer)
  - Greatly simplifies the down-stream data-handling
  - Available on a variety of platforms: VME, ATCA, PCI, USB, ...
  - Root of every modern DAQ system
- Frequently you have to open the “black box” and see where numbers come from
  - Real electronics does not behave as the ideal one
- Trade-offs between speed/precision/cost exist
  - You have to choose the solution that best suits you
- Physics quantities are derived from raw data via calibration
  - Calibration procedures to be foreseen for your detector/DAQ

