



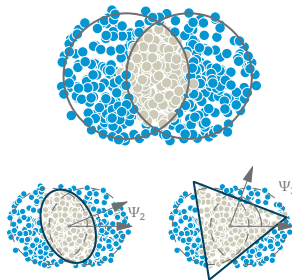
# Heavy-flavor $R_{AA}$ and event-by-event $v_n$ correlations in heavy ion collisions

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Jorge Noronha and Alexandre A. P. Suaide

Quark Matter 2017

- Nuclear modification factor;
- Collective flow;
  - Higher order coefficients:  $v_3$  for heavy flavor!
- Multiparticle cumulants:
  - Unambiguous measurement;
  - Account for event-by-event flow fluctuations;



- Event-plane method is obsolete:

Luzum M. and Ollitrault J.Y. PRC **87**, 044907 (2013).

- Ambiguous results;
  - $v_n\{2\}$  and  $v_n\{4\}$  already have all information contained in  $v_n\{EP\}$ .
- Using Q-Cumulants method define  $v_n\{m\}$ :

Bilandzic A. *et al*, PRC **83** 044913 (2011), PRC **89** 064904 (2014)

$$\bullet \quad v_n\{2\}(p_T) = \frac{d_n\{2\}(p_T)}{(c_n\{2\})^{1/2}}$$

$$\bullet \quad v_n\{6\}(p_T) = \frac{d_n\{6\}(p_T)}{[4(c_n\{6\})^5]^{1/6}}$$

$$\bullet \quad v_n\{4\}(p_T) = \frac{-d_n\{4\}(p_T)}{(-c_n\{4\})^{3/4}}$$

$$\bullet \quad v_n\{8\}(p_T) = \frac{-d_n\{8\}(p_T)}{[33(-c_n\{8\})^7]^{1/8}}$$

- Reference  $m$ -particle cumulants:

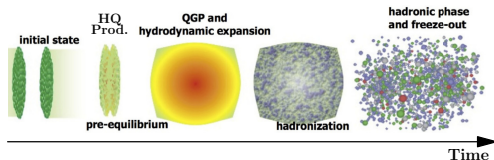
$$\bullet \quad c_n\{2\} = \langle\langle 2 \rangle\rangle$$

$$\bullet \quad c_n\{4\} = \langle\langle 4 \rangle\rangle - 2 \langle\langle 2 \rangle\rangle^2$$

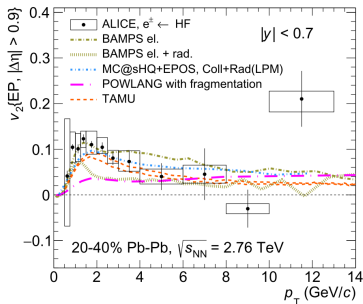
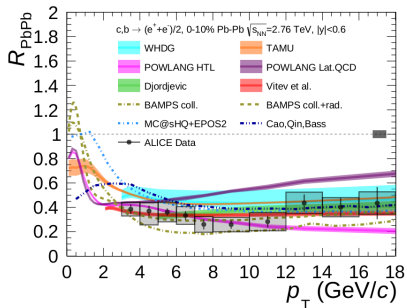
$$\bullet \quad c_n\{6\} = \langle\langle 6 \rangle\rangle - 9 \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle + 12 \langle\langle 2 \rangle\rangle^3$$

$$\bullet \quad c_n\{8\} = \langle\langle 8 \rangle\rangle - 16 \langle\langle 6 \rangle\rangle \langle\langle 2 \rangle\rangle - 18 \langle\langle 4 \rangle\rangle^2 + 144 \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle^2 - 144 \langle\langle 2 \rangle\rangle^4$$

- Correct wide centrality bins with multiplicity weighting.



- Hard in general to get both observables right;
- Try new observables in order to distinguish models.



- What happens with the higher flow harmonics?
- How does the collision energy affects these observables?
- Is there collectivity in the heavy flavor sector?
- How do fluctuations in an event-by-event approach affect these observables?

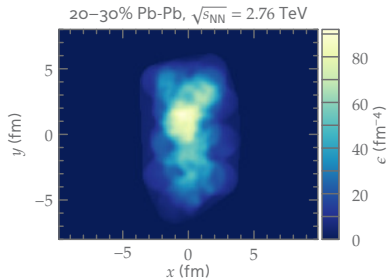
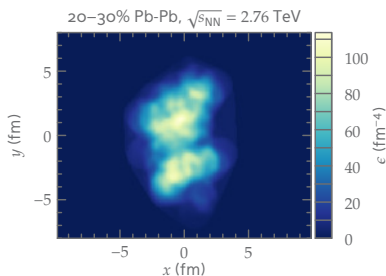
Prado *et al*, arXiv:1611.02962 [nucl-th] (2016)

- Monte Carlo simulation (C++ with ROOT and Pythia8);
- Modules:
  - Energy loss model;
  - Hadronization;
  - Meson decay.
- Different models can be freely plugged in/out;
- Possible to run the simulation on top of any hydrodynamics code;
- Heavy quarks (bottom and charm) are probes:
  - Sampled at the beginning of simulation;
  - Evolved with the medium;
  - We currently neglect any effect of the probes on the medium.
- Currently no coalescence is implemented (future work).

Noronha-Hostler *et al*, PRL **116**, 252301 (2016), PRC **90**, 034907 (2014), PRC **88** 044916 (2013)

- MCKLN initial conditions + v-USPhydro;
- 2D profiles:
  - Energy density;
  - Temperature;
  - Transverse velocity.
- Heavy quarks initial state:
  - Position: number of binary collisions;
  - Momentum: FONLL;
  - Direction: random.

Cacciari *et al*, JHEP **1210**, 137 (2012)



Prado *et al*, arXiv:1611.02962 [nucl-th] (2016)

# Energy Loss and Hadronization

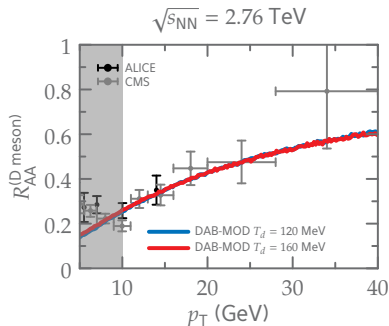
- Energy loss parametrization:  $\frac{dE}{dx}(T, v) = \alpha \Gamma_{\text{flow}}$ .  
Das *et al*, PLB **747** 260–264 (2015)
  - Simple model; can be replaced anytime for more realistic ones;
- Flow factor:  $\Gamma_{\text{flow}} = \gamma [1 - v \cos(\varphi_{\text{quark}} - \varphi_{\text{flow}})]$ .  
Armesto *et al*, PRC **72** 064910 (2005); Noronha-Hostler *et al*, PRL **116**, 252301 (2016)
  - Couples differences in path length in the medium to energy loss experienced by partons.
- Free parameter  $\alpha$  fit:
  - 1 Fit  $\alpha_{\text{charm}}$  using D meson  $R_{AA}$  data;
  - 2 Fit  $\alpha_{\text{bottom}}$  using B meson or electron  $R_{AA}$  data;



# Nuclear modification factor

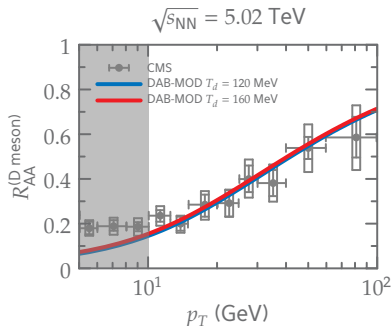
- Collision energies: 2.76 TeV and 5.02 TeV;
- After fitting  $\alpha$ , temperature has little influence on  $R_{AA}$  results.

## D meson $R_{AA}$ ; 0–10% Pb-Pb



JHEP 03, 081 (2016) (ALICE)

CMS-PAS-HIN-15-005 (CMS)

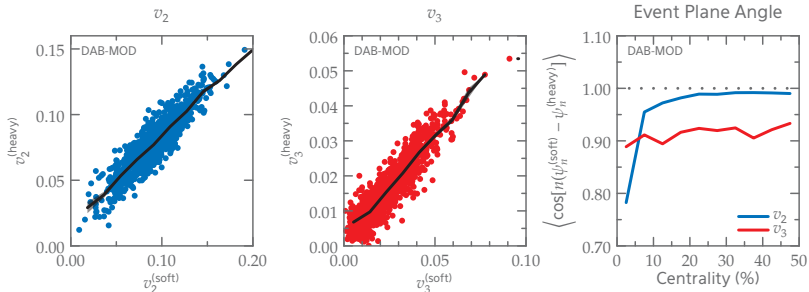


CMS-PAS-HIN-16-001 (2016)

\*Gray area: where coalescence may be important.

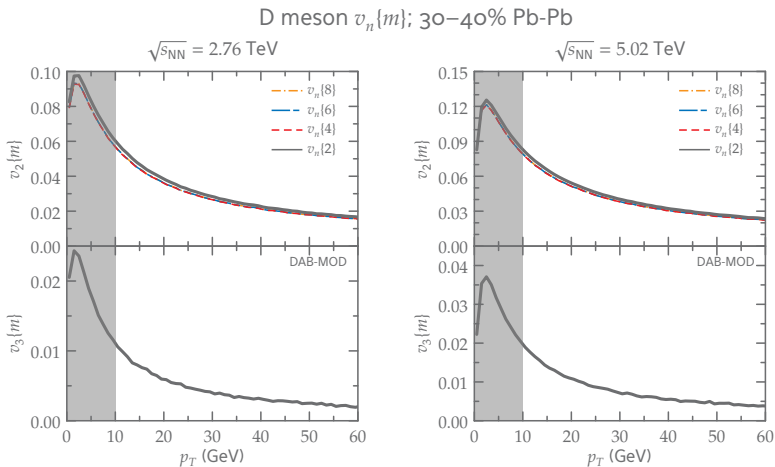
- 2-particle correlation: 
$$v_n\langle 2 \rangle(p_T) = \frac{\langle v_n^{(\text{heavy})}(p_T) v_n^{(\text{soft})} \cos [n(\psi_n^{(\text{heavy})}(p_T) - \psi_n^{(\text{soft})})] \rangle}{\sqrt{\langle (v_n^{(\text{soft})})^2 \rangle}}$$
  - This formula maximizes  $v_n$ .
  - Heavy and soft sectors have the medium as the same underlying cause for flow;

D meson; 30–40% Pb-Pb,  $\sqrt{s_{\text{NN}}} = 5.02$  TeV



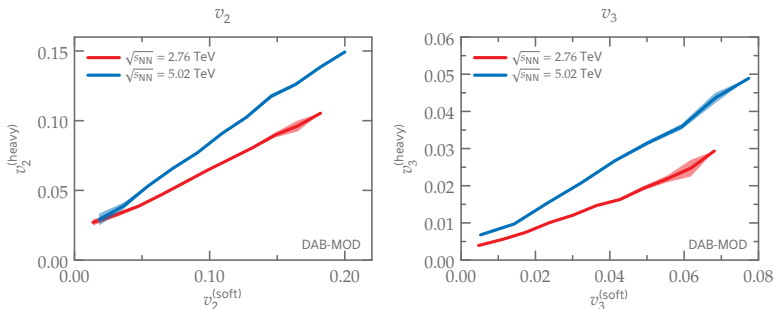
Noronha-Hostler *et al*, PRL **116**, 252301 (2016) Nahrgang *et al*, PRC **91**, 014904 (2015)

- 2, 4, 6 and 8-particle correlations for  $v_2$  and  $v_3\{2\}$ .



ATLAS Coll., PRC 92, 034903 (2015)

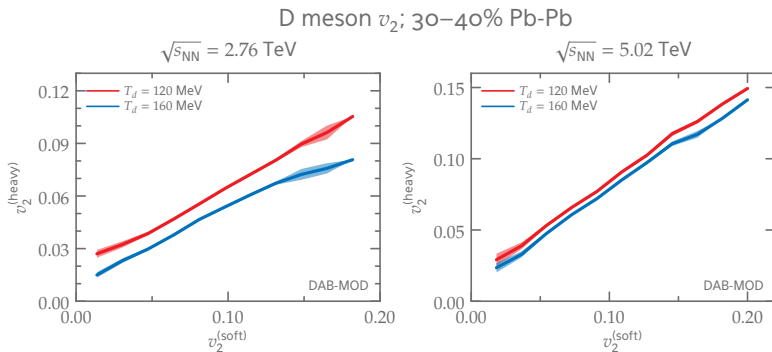
- Heavy-soft correlation is affected by the collision energy;
- Only possible to compute in event-by-event analysis;

D meson  $v_n$ ; 30–40% Pb-Pb

...see also Dobrin's talk: Tue 2pm

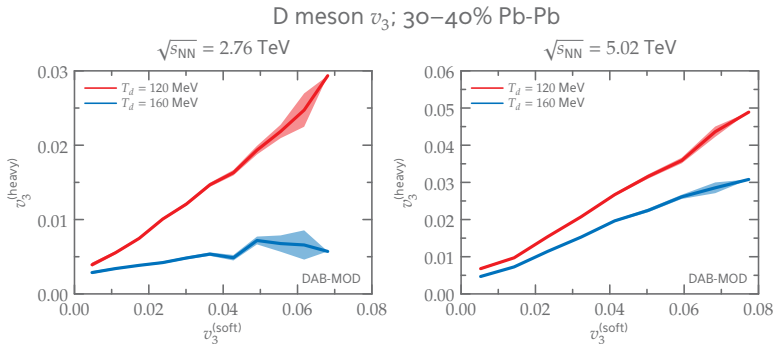
ATLAS Coll., PRC 92, 034903 (2015)

- Heavy-soft correlation is also affected by  $T_d$ ;
- With higher  $T_d$  quarks lifetime within the plasma is shortened leading to smaller  $v_n$ .



ATLAS Coll., PRC 92, 034903 (2015)

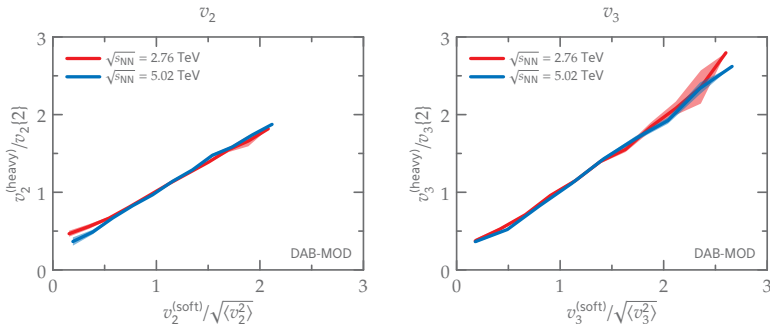
- $v_3$  is heavily affected by the change of temperature!



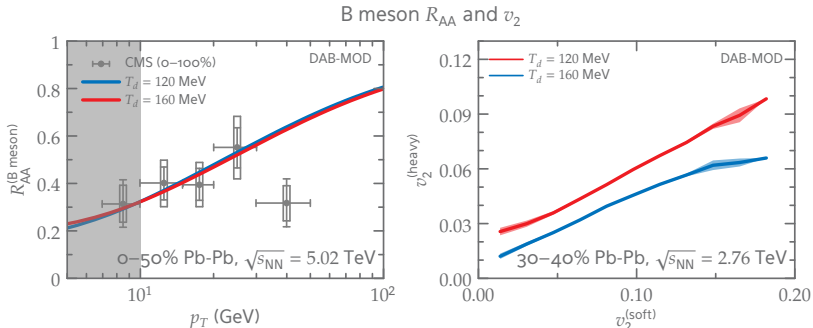
- By normalizing the axes all the correlation curves superpose;
- This is a direct consequence of the linear response between  $v_2$  and  $\varepsilon_2$  both in the heavy and the soft sector.

Prado *et al*, arXiv:1611.02962 [nucl-th] (2016); Noronha-Hostler *et al*, arXiv:1609.05171 (2016)

### D meson $v_n$ ; 30–40% Pb-Pb



- Preliminary results for B meson.



CMS-PAS-HIN-16-011 (2016)

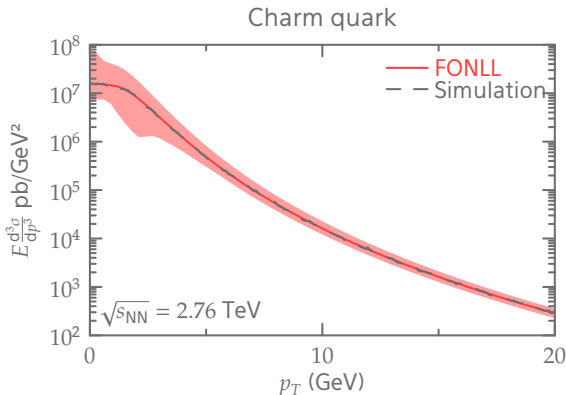
\*Gray area: where coalescence may be important.



- DAB-MOD (D and B mesons MODular code): A new framework to study heavy probes in an event-by-event hydrodynamics expanding QGP;
- Soft-heavy engineering:
  - One can compute the event-by-event correlations between soft and heavy sectors;
  - Leads to a new observable which can provide a novel signature of collectivity in the heavy flavor sector;
  - Needs experimental data for confirmation.
- Future work:
  - Improvement in the low- $p_T$  sector;
  - Correlations between  $v_n$  in the heavy-flavor sector;
  - Check the influence of a  $T$  dependent shear and bulk viscosities;

# Backup

- Quark initial momentum: pQCD (FONLL).
- Quark initial direction ( $\varphi$ ): random;



Cacciari *et al*, JHEP **1210**, 137 (2012)

- Hadronization:

- Heavy quarks hadronize after crossing the  $T_d$  isothermal.
- We use Peterson fragmentation function:

$$D(z) \propto \frac{1}{z \left(1 - \frac{1}{z} - \frac{\varepsilon}{1-z}\right)^2};$$

- The parameter  $\varepsilon$  is fitted to reproduce FONLL calculations;
  - Neglect coalescence: focus on the intermediate regime ( $p_T > 10$  GeV).
- Decays:
    - Performed using Pythia8;
    - Only semi-leptonic channels selected;

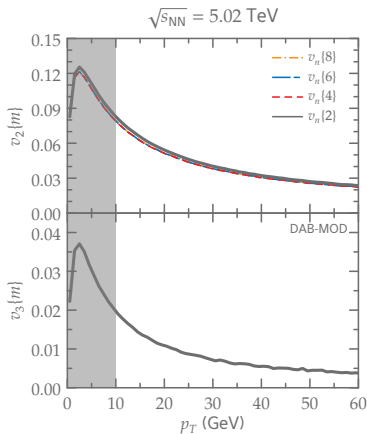
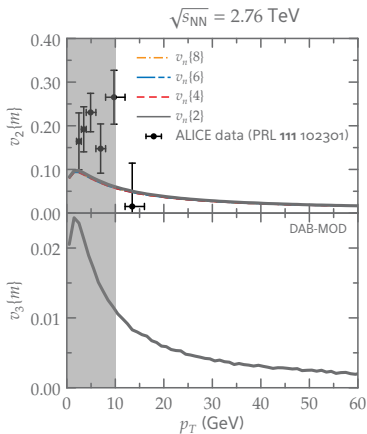
- Viscous event-by-event 2D+1 hydrodynamics:

Noronha-Hostler *et al*, PRL **116**, 252301 (2016), PRC **90**, 034907 (2014), PRC **88** 044916 (2013)

- v-USPhydro code;
  - MCKLN initial conditions;
  - Fast computational time;
  - Well tested code that fits data for soft flow harmonics!
- Freeze-out: Cooper-Frye prescription:
    - Includes viscous corrections;
    - Decoupling temperature  $T_d = 120$  MeV.

- 2, 4, 6 and 8-particle correlations for  $v_2$  and  $v_3\{2\}$ .

D meson  $v_n\{m\}$ ; 30–40% Pb-Pb



\*Gray area: where coalescence may be important.

- By normalizing the axes all the correlation curves superpose;
- This is a direct consequence of the linear response between  $v_2$  and  $\varepsilon_2$  both in the heavy and the soft sector.

Noronha-Hostler *et al*, PRC **93**, 014909 (2016)

D meson  $v_n$ ; 30–40% Pb-Pb  $\sqrt{s_{NN}} = 5.02$  TeV

