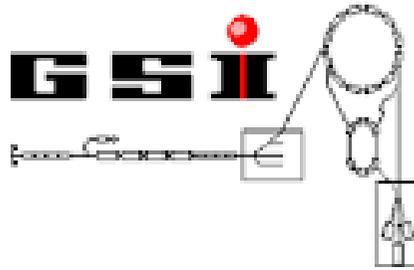




FIAS Frankfurt Institute  
for Advanced Studies



HIC  
for FAIR  
Helmholtz International Center

GOETHE  
UNIVERSITÄT  
FRANKFURT AM MAIN

# Tomography of the Quark-Gluon-Plasma by Charm Quarks

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In collaboration with **Taesoo Song**, Hamza Berrehrah, Daniel Cabrera, Juan Torres-Rincon, Laura Tolos, Wolfgang Cassing, Jörg Aichelin and Pol-Bernard Gossiaux

32nd Winter Workshop on Nuclear Dynamics  
Guadeloupe, France, 28 February - 5 March 2016

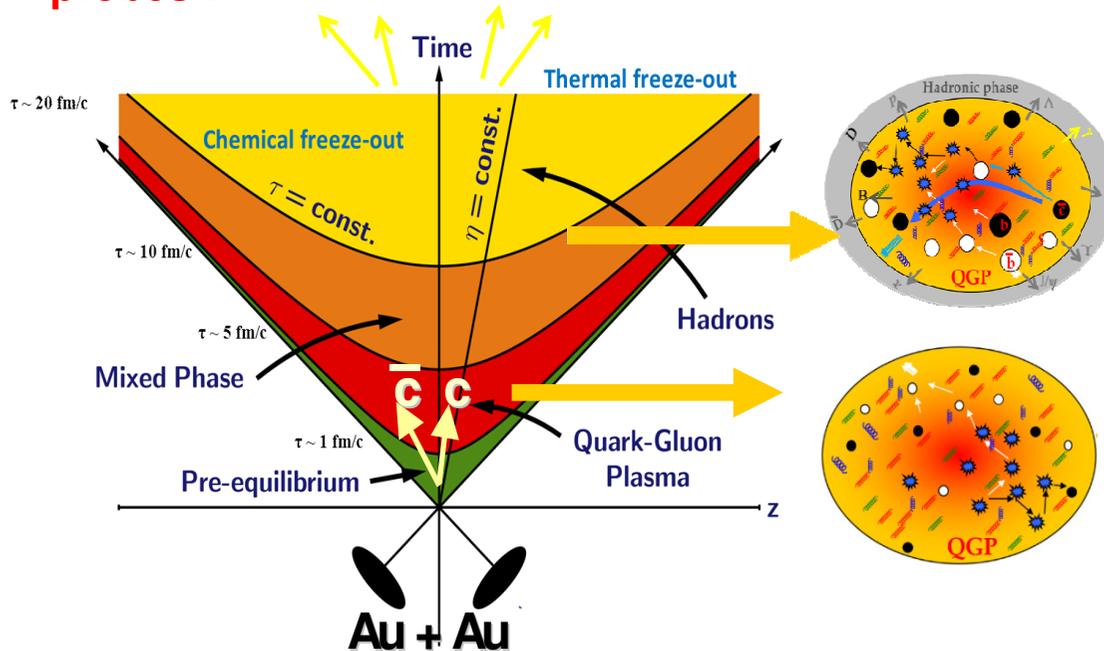
DAAD

DFG Deutsche  
Forschungsgemeinschaft



# Motivation

□ study of the properties of hot and dense nuclear and partonic matter by **'charm probes'**:



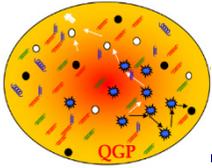
The **advantages** of the 'charm probes':

□ dominantly produced in the very early stages of the reactions in **initial binary collisions** with large energy-momentum transfer

□ initial charm production is well described by **pQCD** – FONLL

□ scattering cross sections are small (compared to the light quarks) → **not in an equilibrium** with the surrounding matter

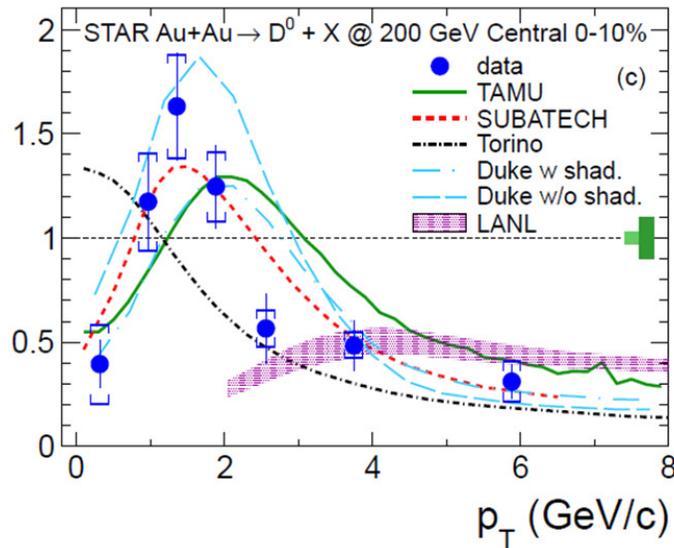
→ Hope to use 'charm probes' for an **early tomography of the QGP**



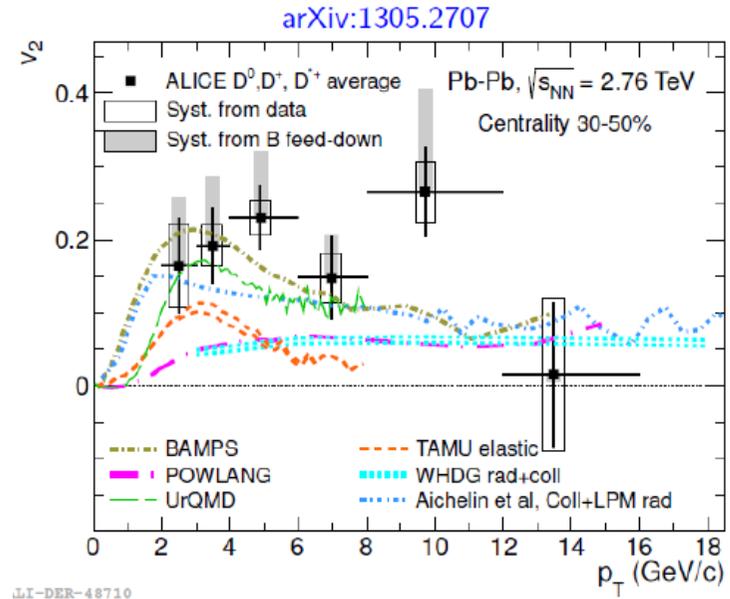
# Charm: experimental signals

## 1. Nuclear modification factor:

$$R_{AA}(p_T) \equiv \frac{dN_D^{Au+Au}/dp_T}{N_{\text{binary}}^{Au+Au} \times dN_D^{p+p}/dp_T}$$



## 2. Elliptic flow $v_2$ :



□ What is the origin for the “energy loss” of charm at large  $p_T$ ?

Collisional energy loss (elastic scattering  $Q+q \rightarrow Q+q$ )  
vs radiative (gluon bremsstrahlung  $Q+q \rightarrow Q+q+g$ ) ?

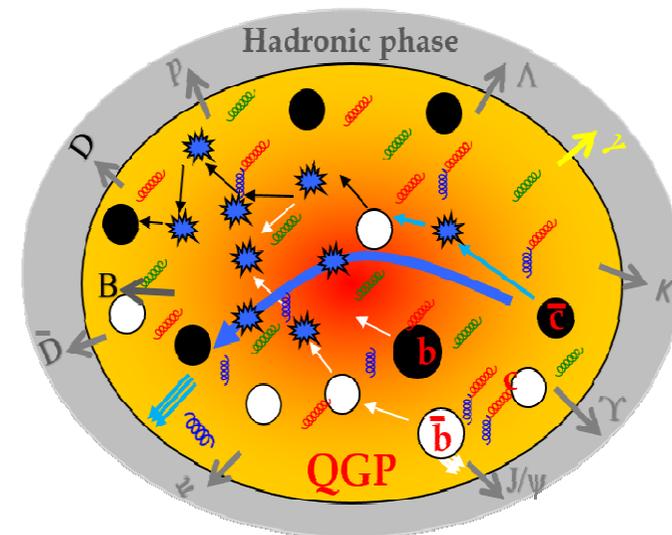
→ **Challenge for theory:** simultaneous description of  $R_{AA}$  and  $v_2$  !

# Dynamics of charm quarks in A+A

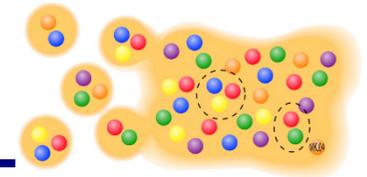
1. **Production** of charm quarks in initial binary collisions
2. **Interactions in the QGP:**
  - elastic scattering  $Q+q \rightarrow Q+q$  → collisional energy loss
  - gluon bremsstrahlung  $Q+q \rightarrow Q+q+g$  → radiative energy loss
3. **Hadronization:** c/cbar quarks → D(D\*)-mesons:  
coalescence vs fragmentation
4. **Hadronic interactions:**  
D+baryons; D+mesons

**The goal:** to model the dynamics of charm quarks/mesons in all phases on a **microscopic basis**

**The tool:** PHSD approach



# From SIS to LHC: from hadrons to partons



**The goal:** to study of the phase transition from hadronic to partonic matter and properties of the Quark-Gluon-Plasma on a **microscopic level**

→ need a **consistent non-equilibrium transport model**

- ❑ with explicit **parton-parton interactions** (i.e. between quarks and gluons)
- ❑ explicit **phase transition** from hadronic to partonic degrees of freedom
- ❑ **IQCD EoS** for partonic phase (‘cross over’ at  $\mu_q=0$ )

❑ **Transport theory for strongly interacting systems:** off-shell Kadanoff-Baym equations for the Green-functions  $S_h^<(x,p)$  in phase-space representation for the **partonic** and **hadronic phase**



**Parton-Hadron-String-Dynamics (PHSD)**

**QGP phase described by**

**Dynamical QuasiParticle Model  
(DQPM)**

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919;  
NPA831 (2009) 215;  
W. Cassing, EPJ ST 168 (2009) 3

A. Peshier, W. Cassing, PRL 94 (2005) 172301;  
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

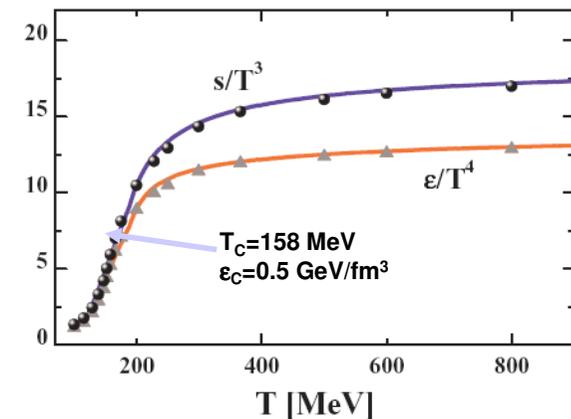
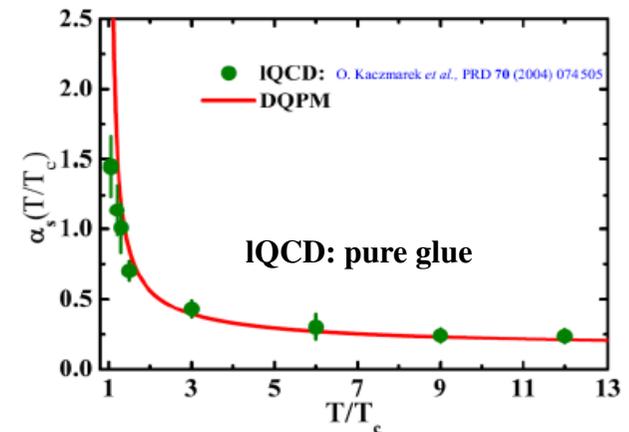
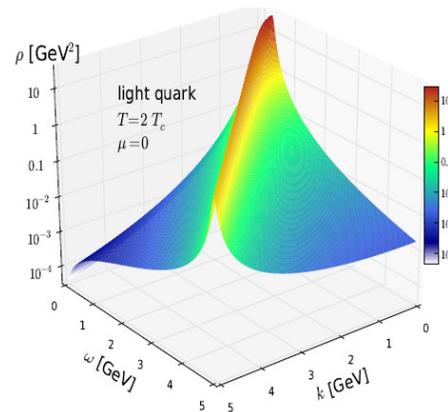
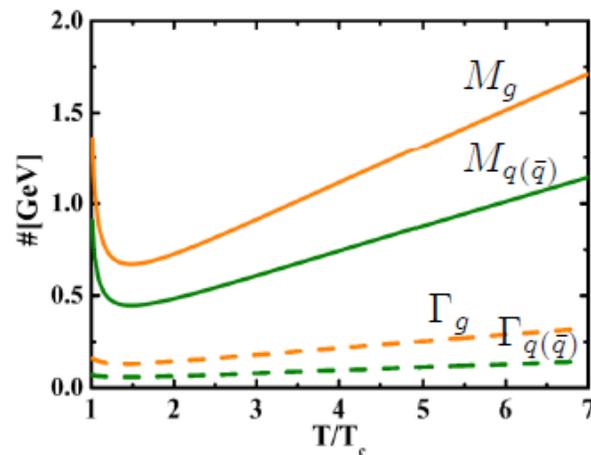
# The Dynamical QuasiParticle Model (DQPM)

- Basic idea: **interacting quasi-particles: massive quarks and gluons (g, q, q<sub>bar</sub>)** with **Lorentzian spectral functions** :

$$\rho_i(\omega, T) = \frac{4\omega\Gamma_i(T)}{(\omega^2 - \vec{p}^2 - M_i^2(T))^2 + 4\omega^2\Gamma_i^2(T)} \quad (i = q, \bar{q}, g)$$

- fit to lattice (IQCD) results (e.g. entropy density) with 3 parameters

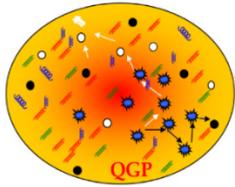
→ **Quasi-particle properties:**  
large width and mass for gluons and quarks



- DQPM provides mean-fields (1PI) for gluons and quarks as well as effective 2-body interactions (2PI)
- DQPM gives transition rates for the formation of hadrons → PHSD

DQPM: Peshier, Cassing, PRL 94 (2005) 172301; Cassing, NPA 791 (2007) 365; NPA 793 (2007)

# Basic idea of PHSD



## QGP in equilibrium

**Dynamical QuasiParticle Model (DQPM):**

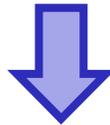
Quasiparticle properties:  
 ,resummed' self-energies, propagators  
 → Calculation of cross sections



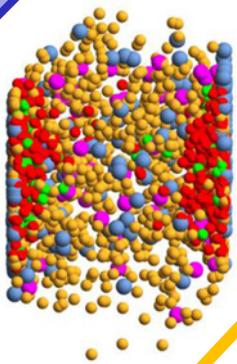
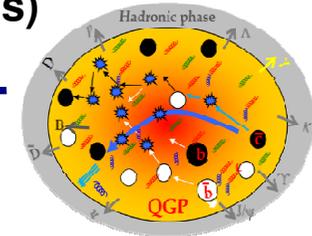
**IQCD**

Controlled by IQCD!  
 Calculation of transport coefficients  $\eta, \zeta, \sigma_0, \dots$

DQPM: consider all the **effects of the nonperturbative nature** of the strongly interacting quark-gluon plasma (**sQGP**) constituents (vs. pQCD models)



## QGP out-of equilibrium ↔ HIC



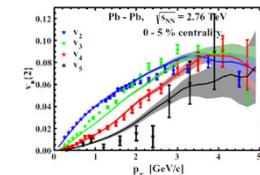
**Parton-Hadron-String-Dynamics (PHSD)**

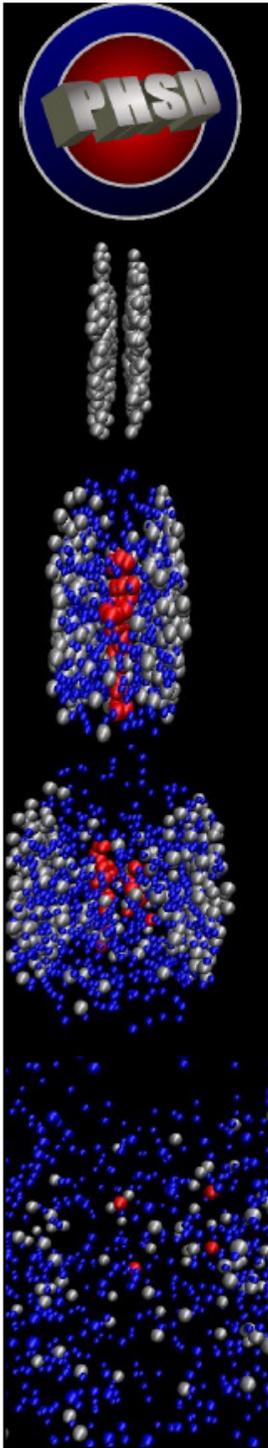
Controlled by

**experimental data**



Partonic interactions → DQPM  
 hadronic interactions → hadron physics  
 \* In-medium hadronic interactions → many-body physics: G-matrix

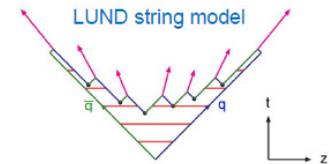




# Parton-Hadron-String-Dynamics (PHSD)

## Initial A+A collisions – HSD:

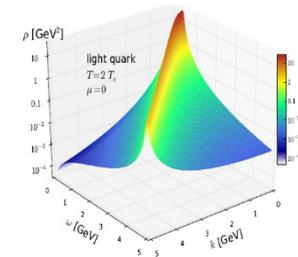
$N+N \rightarrow$  string formation  $\rightarrow$  decay to pre-hadrons



## Formation of QGP stage if $\epsilon > \epsilon_{\text{critical}}$ :

dissolution of pre-hadrons  $\rightarrow$  (DQPM)  $\rightarrow$

$\rightarrow$  massive quarks/gluons + mean-field potential  $U_q$



## Partonic stage – QGP:

based on the Dynamical Quasi-Particle Model (DQPM)

### (quasi-) elastic collisions:

$$q + q \rightarrow q + q \quad g + q \rightarrow g + q$$

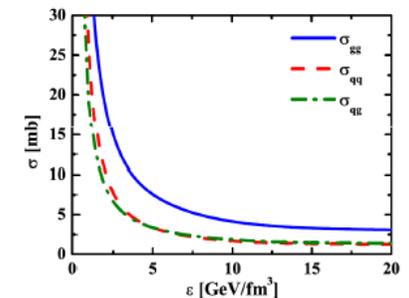
$$q + \bar{q} \rightarrow q + \bar{q} \quad g + \bar{q} \rightarrow g + \bar{q}$$

$$\bar{q} + \bar{q} \rightarrow \bar{q} + \bar{q} \quad g + g \rightarrow g + g$$

### inelastic collisions:

$$q + \bar{q} \rightarrow g \quad q + \bar{q} \rightarrow g + g$$

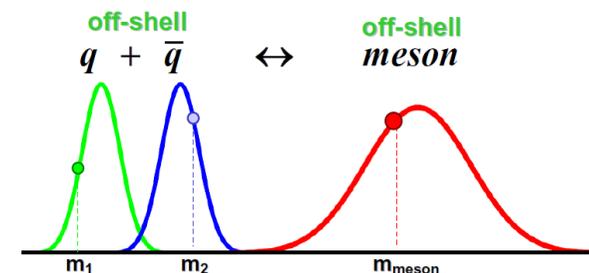
$$g \rightarrow q + \bar{q} \quad g \rightarrow g + g$$



## Hadronization (based on DQPM):

$$g \rightarrow q + \bar{q}, \quad q + \bar{q} \leftrightarrow \text{meson (or 'string')}$$

$$q + q + q \leftrightarrow \text{baryon (or 'string')}$$

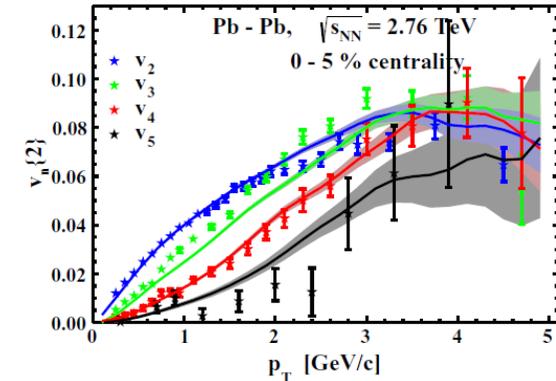
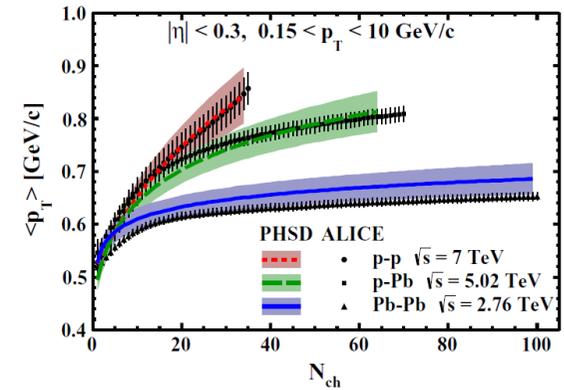
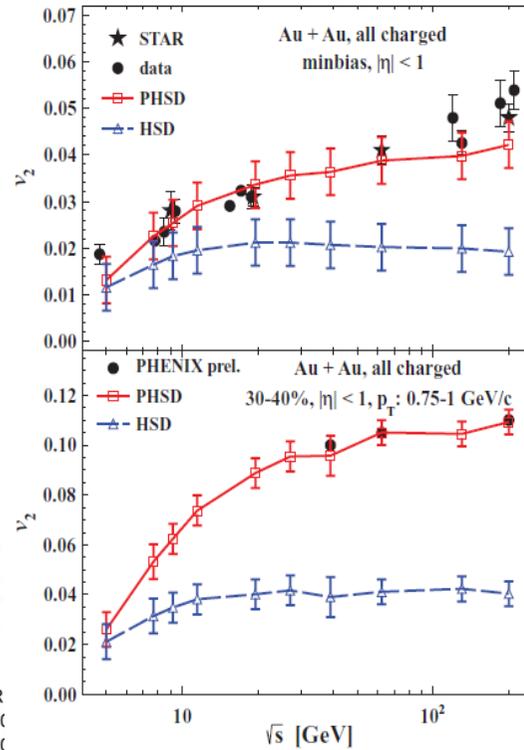
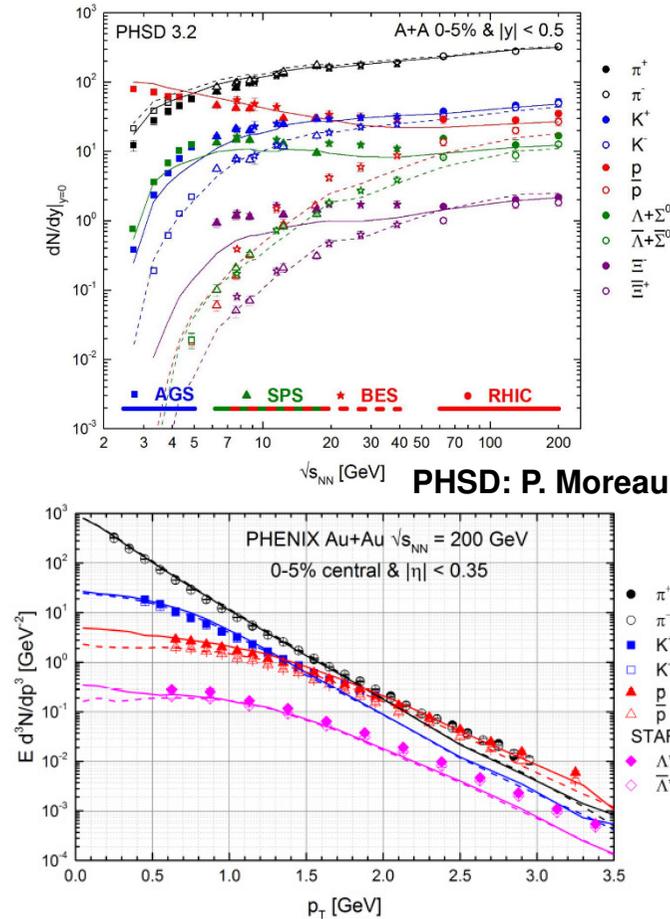


## Hadronic phase: hadron-hadron interactions – off-shell HSD



# Non-equilibrium dynamics: description of A+A with PHSD

**Important:** to be conclusive on charm observables, the **light quark dynamics** must be well under control!



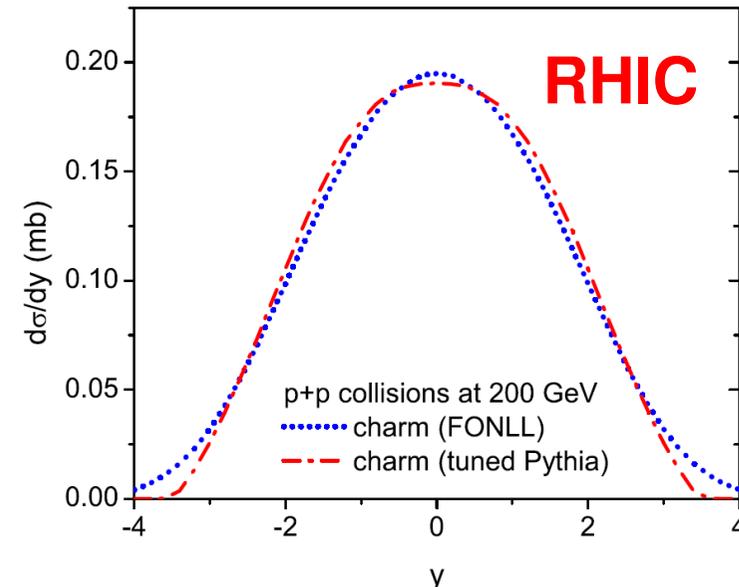
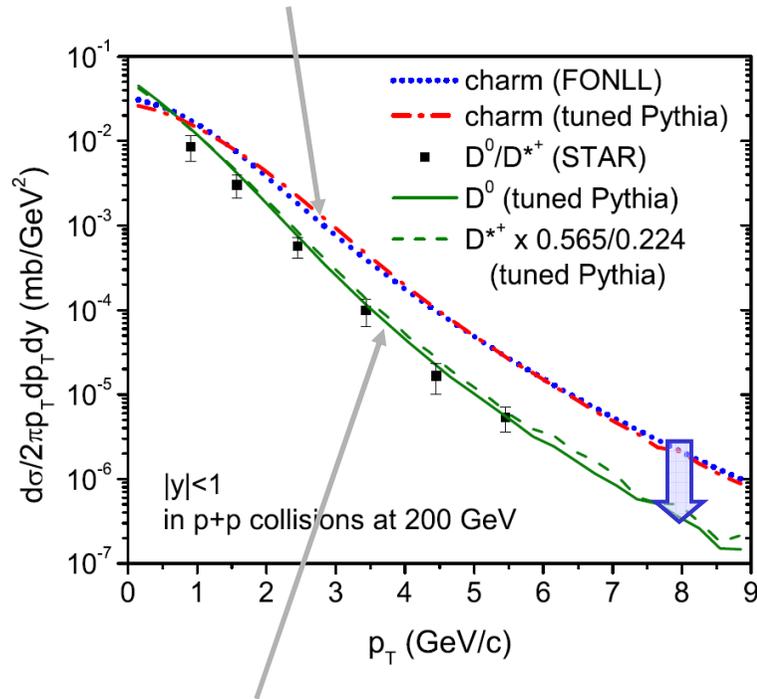
V. Konchakovski et al.,  
PRC 85 (2012) 011902; JPG42 (2015) 055106

**PHSD** provides a **good description of 'bulk' observables** ( $y$ -,  $p_T$ -distributions, flow coefficients  $v_n$ , ...) from SPS to LHC



# Charm quark/hadrons production in p+p collisions

1) **Momentum distribution of charm quark:** Use **tuned** PYTHIA event generator to reproduce **FONLL** (fixed-order next-to-leading log) results (R. Vogt et al.)



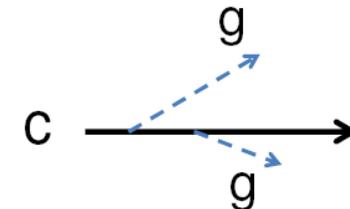
T. Song et al., PRC 92 (2015) 014910, arXiv:1503.03039

2) **Charm hadron production by c-quark fragmentation:**

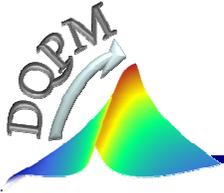
- D<sup>0</sup> 20 %
- D<sup>+</sup> 17.4 %
- D<sup>\*0</sup> 21.3 %
- D<sup>\*+</sup> 22.4 %
- Ds<sup>+</sup> 8 %
- Λ<sub>c</sub> 9.4 %

$$D_Q^H(z) \sim \frac{1}{z[1 - 1/z - \epsilon_Q/(1-z)]^2}$$

From C. Peterson et al., PRD27 (1983) 105

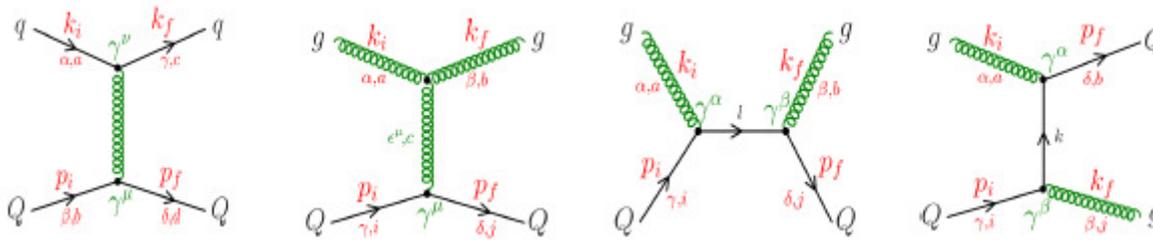


- $z$  : momentum fraction of hadron H in heavy quark Q
- $\epsilon_Q=0.01$  for Q=charm

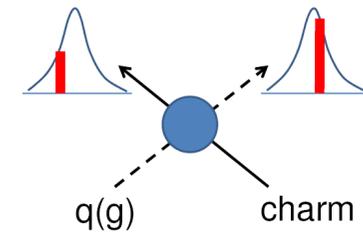


# Charm quark scattering in the QGP (DQPM)

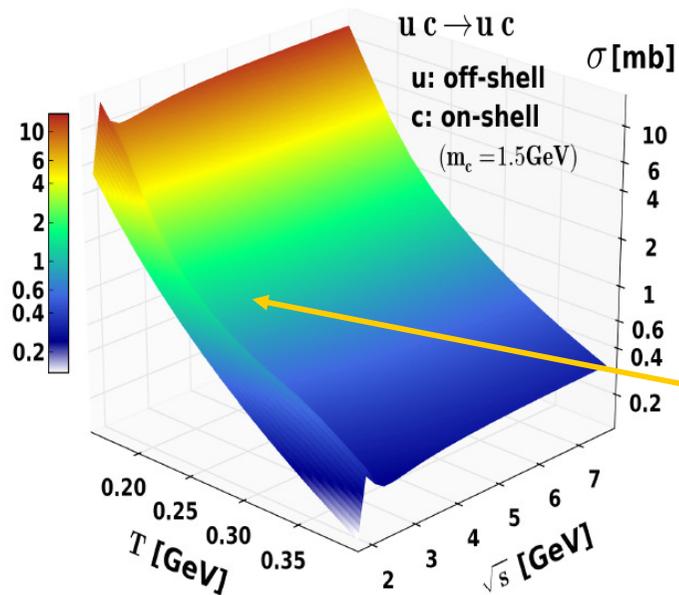
- Elastic scattering with off-shell massive partons  $Q+q(g) \rightarrow Q+q(g)$



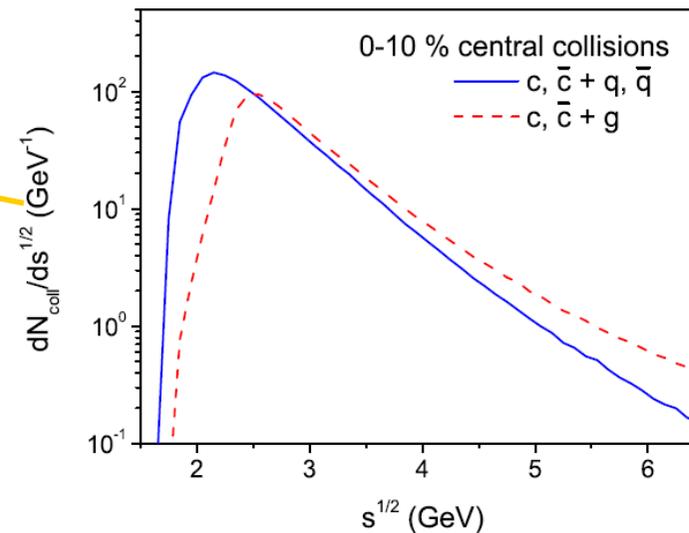
Non-perturbative QGP!

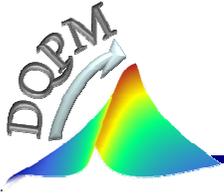


- Elastic cross section  $uc \rightarrow uc$



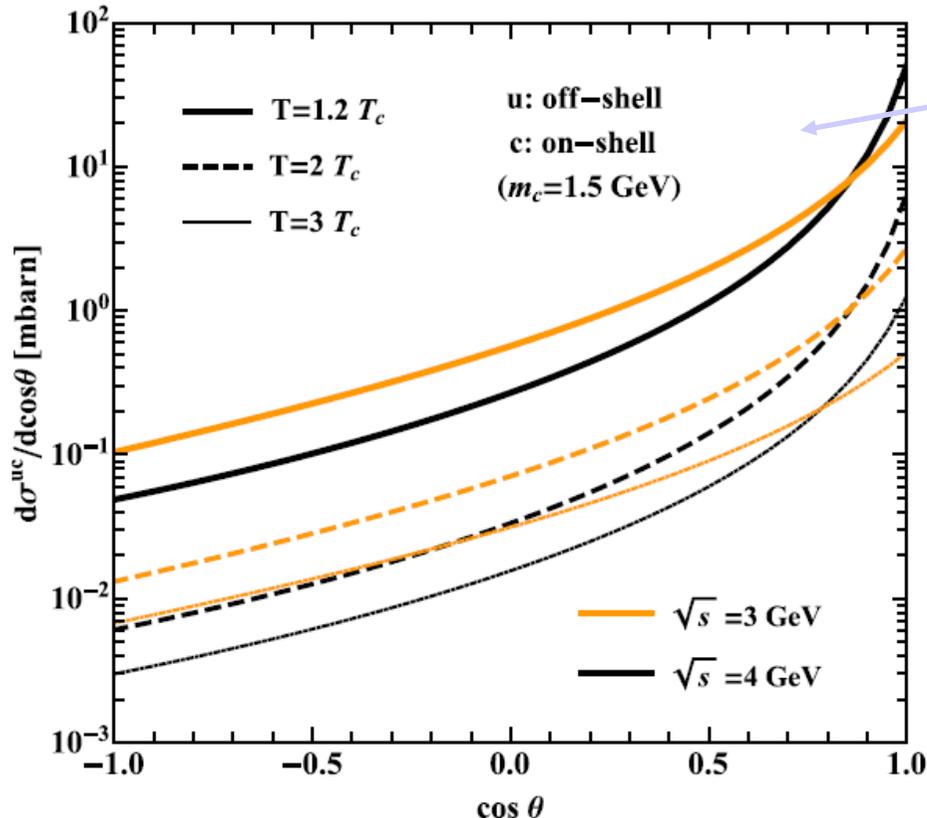
- Distributions of  $Q+q, Q+g$  collisions vs  $s^{1/2}$  in Au+Au, 10% central





# Charm quark scattering in the QGP

□ **Differential elastic cross section** for  $uc \rightarrow uc$  for  $s^{1/2}=3$  and 4 GeV at  $1.2T_c$ ,  $2T_c$  and  $3T_c$



□ **DQPM - anisotropic angular distribution**

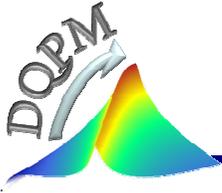
Note: pQCD - strongly forward peaked  
 → Differences between DQPM and pQCD :  
 less forward peaked angular distribution  
 leads to **more efficient momentum transfer**

□  $N(cc) \sim 19$  pairs,  
 $N(Q+q) \sim 130$ ,  $N(Q+g) \sim 85$  collisions

→ each charm quark makes  
 $\sim 6$  elastic collisions

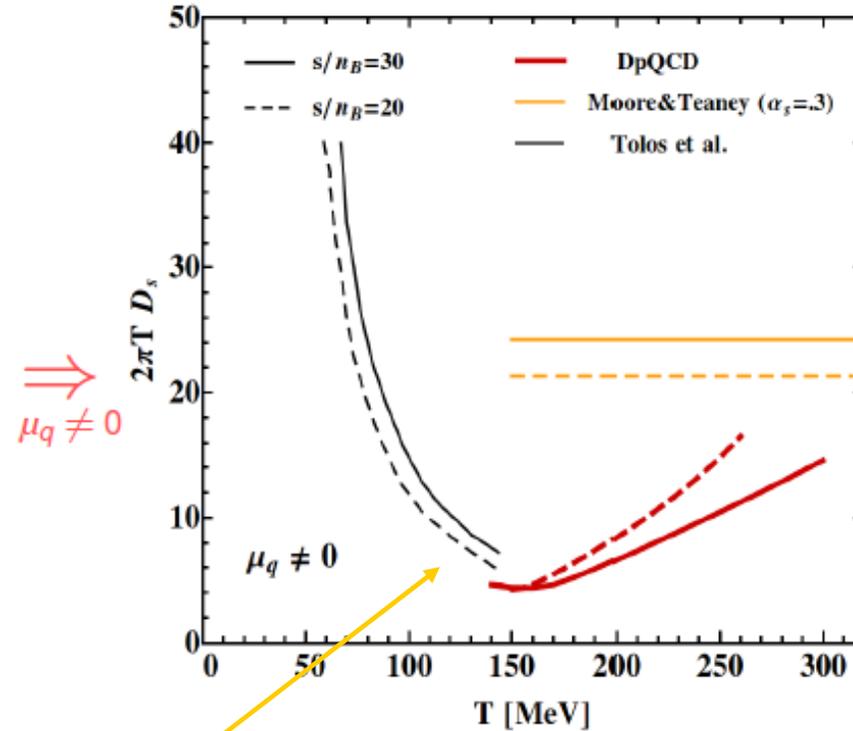
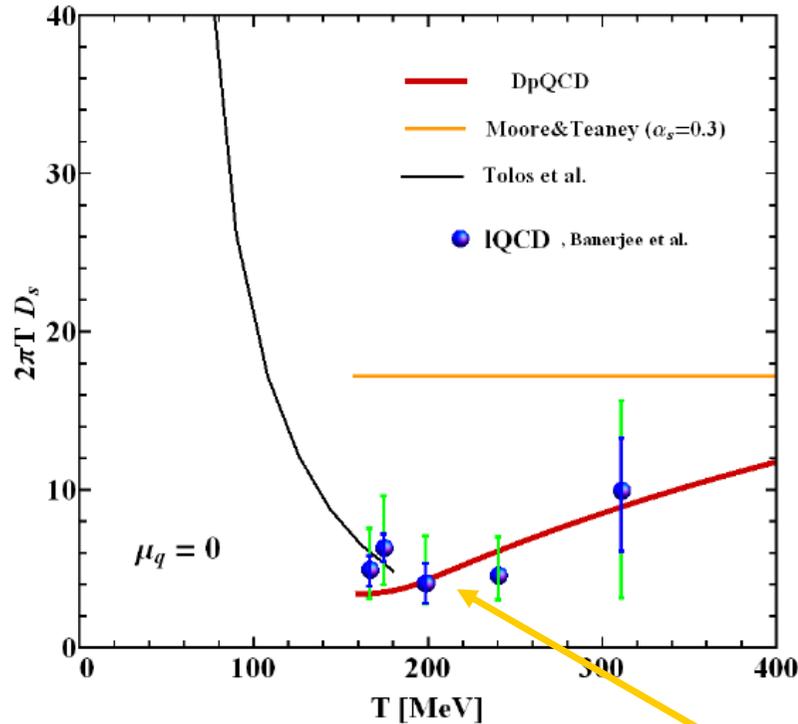
→ **Smaller number** (compared to pQCD)  
 of elastic scatterings with massive  
 partons leads to a **large energy loss**

**! Note:** radiative energy loss is NOT included yet in PHSD (work in progress);  
 expected to be **small** due to the large gluon mass in the DQPM



# Charm spatial diffusion coefficient $D_s$ in the hot medium

- $D_s$  for heavy quarks as a function of  $T$  for  $\mu_q=0$  and finite  $\mu_q$  assuming adiabatic trajectories (constant entropy per net baryon  $s/n_B$ ) for the expansion



□  $T < T_c$  : hadronic  $D_s$

→ Continuous transition at  $T_c$ !

L. Tolos, J. M. Torres-Rincon, PRD 88 (2013) 074019  
 V. Ozvenchuk et al., PRC90 (2014) 054909

H. Berrehrah et al, PRC 90 (2014) 051901, arXiv:1406.5322



# Hadronization of charm quarks in A+A

□ PHSD: if the local energy density  $\varepsilon < \varepsilon_c \rightarrow$  hadronization of charm quarks to hadrons:

## 1. Dynamical coalescence model

$$\varepsilon_c = 0.5 \text{ GeV/fm}^3$$

Probability for charm quark/antiquark and the light quark/antiquark to form a meson:

$$f(\rho, \mathbf{k}_\rho) = \frac{8g_M}{6^2} \exp \left[ -\frac{\rho^2}{\delta^2} - \mathbf{k}_\rho^2 \delta^2 \right]$$

where  $\rho = \frac{1}{\sqrt{2}}(\mathbf{r}_1 - \mathbf{r}_2)$ ,  $\mathbf{k}_\rho = \sqrt{2} \frac{m_2 \mathbf{k}_1 - m_1 \mathbf{k}_2}{m_1 + m_2}$

Degeneracy factor:

$g_M = 1$  for D  
 $g_M = 3$  for D\*

D\*:

D\*<sub>0</sub>(2400)<sup>0</sup>  
 D\*<sub>1</sub>(2420)<sup>0</sup>  
 D\*<sub>2</sub>(2460)<sup>0±</sup>

Width  $\delta \leftarrow$  from root-mean-square radius of meson:

$$\langle r^2 \rangle = \frac{3}{2} \frac{m_1^2 + m_2^2}{(m_1 + m_2)^2} \delta^2$$

### Hadronization scenarios :

- 1: only fragmentation
- 2: coalescence with  $\langle r \rangle = 0.5 \text{ fm}$  + fragmentation
- 3: coalescence with  $\langle r \rangle = 0.9 \text{ fm}$  + fragmentation

## 2. Fragmentation (as in pp)

- if NOT hadronized by coalescence

Note: large  $\langle r \rangle$  used also in Refs:

S. Cao, G. Y. Qin and S. A. Bass, PRC 88, 044907 (2013).

Y. Oh, C. M. Ko, S. H. Lee and S. Yasui, PRC 79, 044905 (2009)



# Hadronization of charm quarks in A+A : LHC

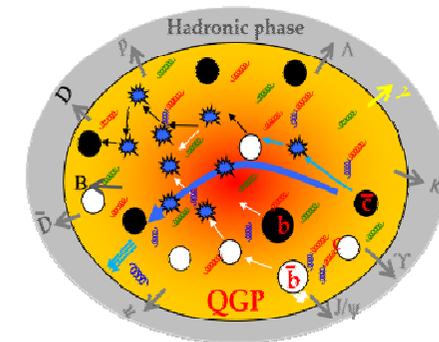
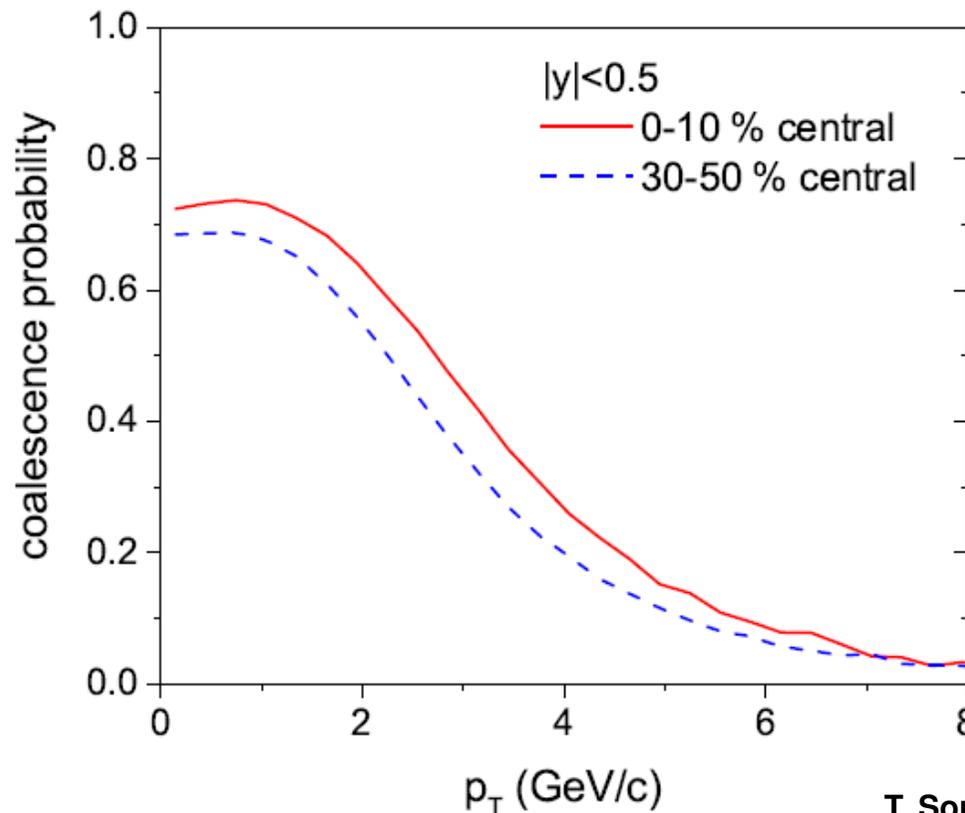
## Dynamical Hadronization scenarios :

4. coalescence with  $\langle r \rangle = 0.9$  fm... + fragmentation

$0.4 < \epsilon < 0.75$  GeV/fm<sup>3</sup>

$\epsilon < 0.4$  GeV/fm<sup>3</sup>

## Coalescence probability in Au+Au



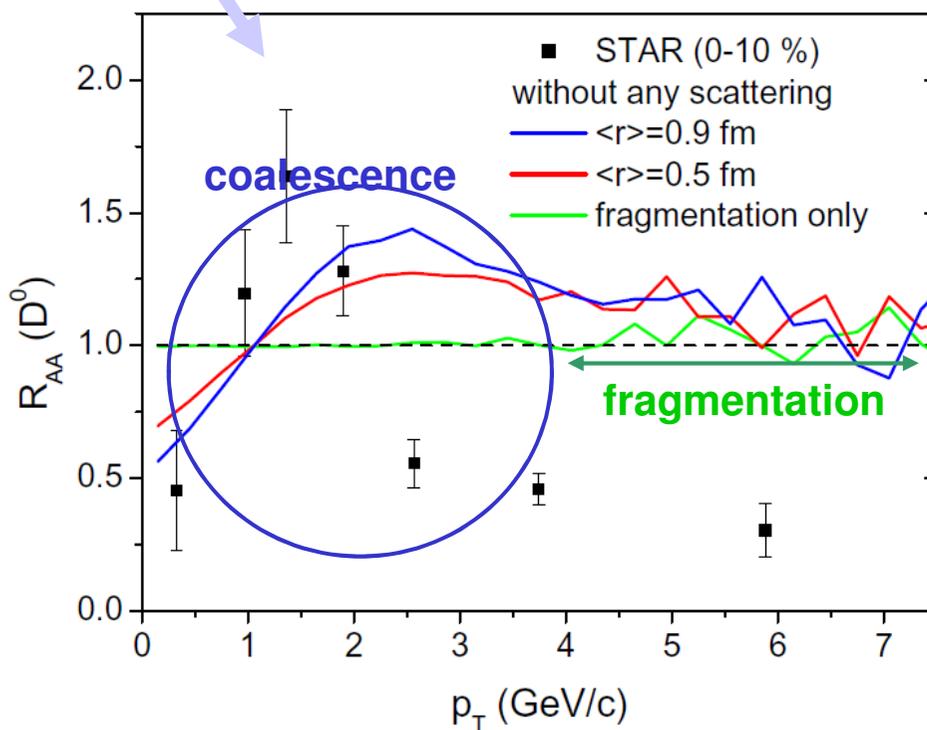
T. Song et al., PRC (2016), arXiv:1512.0089



# $R_{AA}$ at RHIC - coalescence vs fragmentation

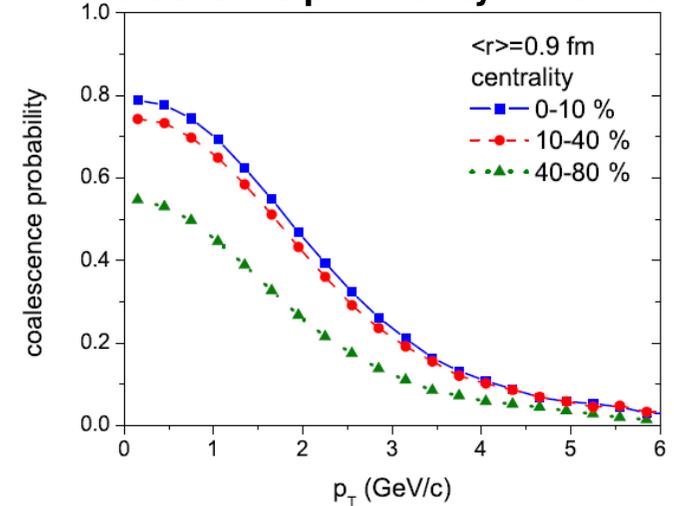
## 1. Influence of hadronization scenarios: coalescence vs fragmentation

**! Model study: without any rescattering (partonic and hadronic)**



$$R_{AA}(p_T) \equiv \frac{dN_D^{Au+Au}/dp_T}{N_{binary}^{Au+Au} \times dN_D^{P+P}/dp_T}$$

Coalescence probability in Au+Au



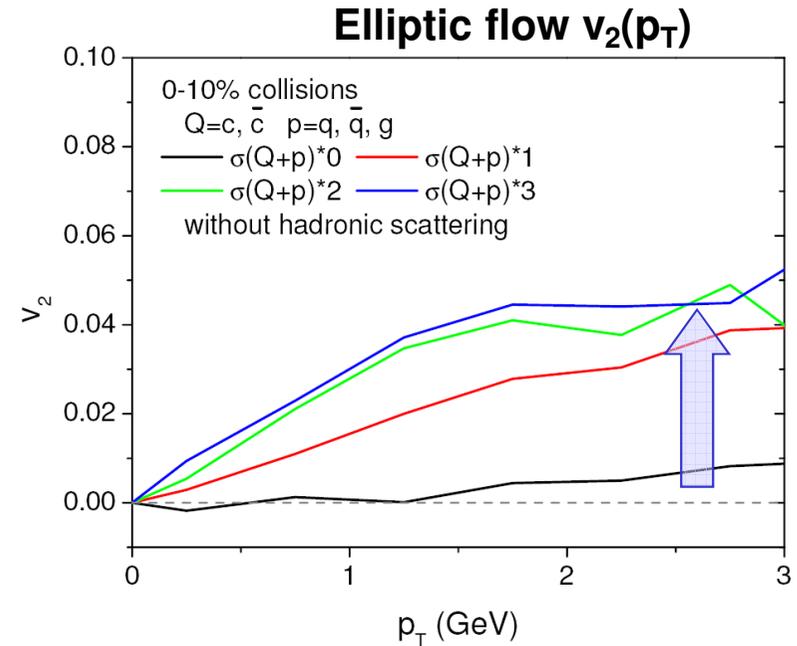
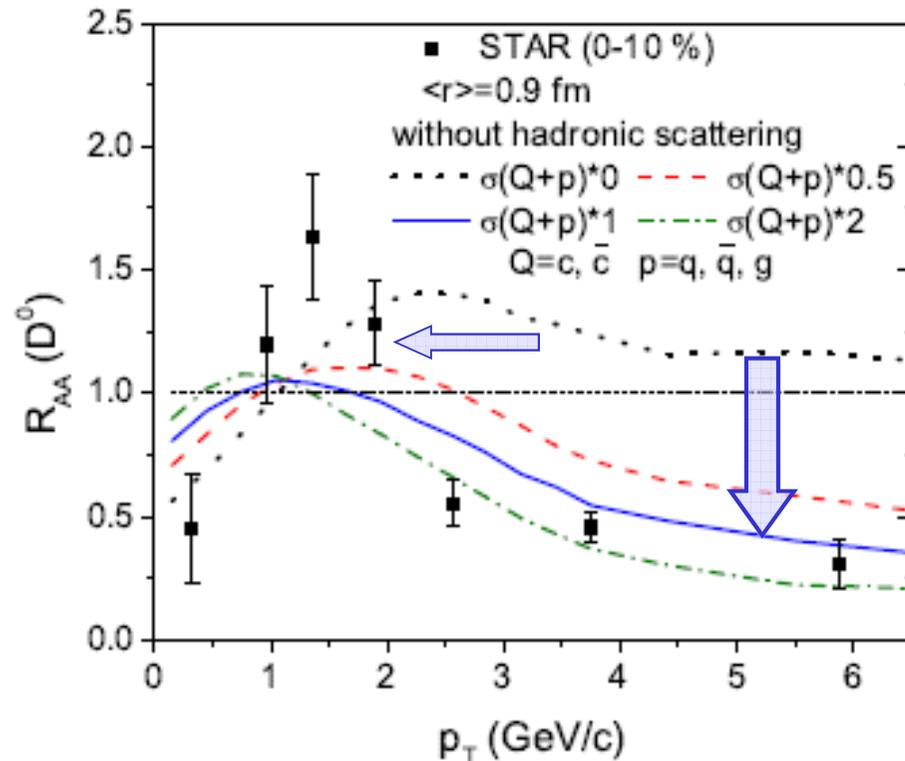
- Expect: no scattering:  $R_{AA}=1$
- Hadronization by **fragmentation** only (as in pp)  $\rightarrow R_{AA}=1$
- **Coalescence** (not in pp!) shifts  $R_{AA}$  to larger  $p_T \rightarrow$  'nuclear matter' effect
- The **height of the  $R_{AA}$  peak** depends on the balance: coalescence vs. fragmentation



# $R_{AA}$ at RHIC: partonic scattering

## 2. Influence of partonic rescattering

! Model study: by scaling of parton cross sections:  $\sigma(Q+q(g))^*\alpha$  by  $\alpha=0, 0.5, 1, 2$  (without hadronic rescattering)



### Elastic partonic rescattering

- moves  $R_{AA}$  to lower  $p_T$  and suppresses large  $p_T$
- increases  $v_2$

Central Au+Au at  $s^{1/2} = 200$  GeV :

$N(cc) \sim 19$  pairs,

$N(Q+q) \sim 130$  collisions

$N(Q+g) \sim 85$  collisions

→ each charm quark makes  
~ 6 elastic collisions



# Modelling of D-meson scattering in the hadronic gas

## 1. D-meson scattering with mesons

**Model:** effective chiral Lagrangian approach with heavy-quark spin symmetry

L. M. Abreu, D. Cabrera, F. J. Llanes-Estrada, J. M. Torres-Rincon, *Annals Phys.* 326, 2737 (2011)

Interaction of  $D=(D^0, D^+, D^+_s)$  and  $D^*=(D^{*0}, D^{*+}, D^{*+}_s)$  with octet  $(\pi, K, Kbar, \eta)$  :

$$\begin{aligned} \mathcal{L}_{LO} = & \langle \nabla^\mu D \nabla_\mu D^\dagger \rangle - m_D^2 \langle DD^\dagger \rangle - \langle \nabla^\mu D^{*\nu} \nabla_\mu D_\nu^{*\dagger} \rangle \\ & + m_D^2 \langle D^{*\mu} D_\mu^{*\dagger} \rangle + ig \langle D^{*\mu} u_\mu D^\dagger - D u^\mu D_\mu^{*\dagger} \rangle \\ & + \frac{g}{2m_D} \langle D^*_\mu u_\alpha \nabla_\beta D_\nu^{*\dagger} - \nabla_\beta D^*_\mu u_\alpha D_\nu^{*\dagger} \rangle \epsilon^{\mu\nu\alpha\beta} \end{aligned}$$

with

$$u_\mu = i (u^\dagger \partial_\mu u - u \partial_\mu u^\dagger)$$

$$U = u^2 = \exp\left(\frac{\sqrt{2}i\Phi}{f}\right) \quad \Phi = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & \pi^+ & K^+ \\ \pi^- & -\frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & K^0 \\ K^- & \bar{K}^0 & -\frac{2}{\sqrt{6}}\eta \end{pmatrix}$$

## 2. D-meson scattering with baryons

**Model:** G-matrix approach: interactions of  $D=(D^0, D^+, D^+_s)$  and  $D^*=(D^{*0}, D^{*+}, D^{*+}_s)$  with nucleon octet  $J^P=1/2^+$  and Delta decuplet  $J^P=3/2^+$

C. Garcia-Recio, J. Nieves, O. Romanets, L. L. Salcedo, L. Tolos, *Phys. Rev. D* 87, 074034 (2013)

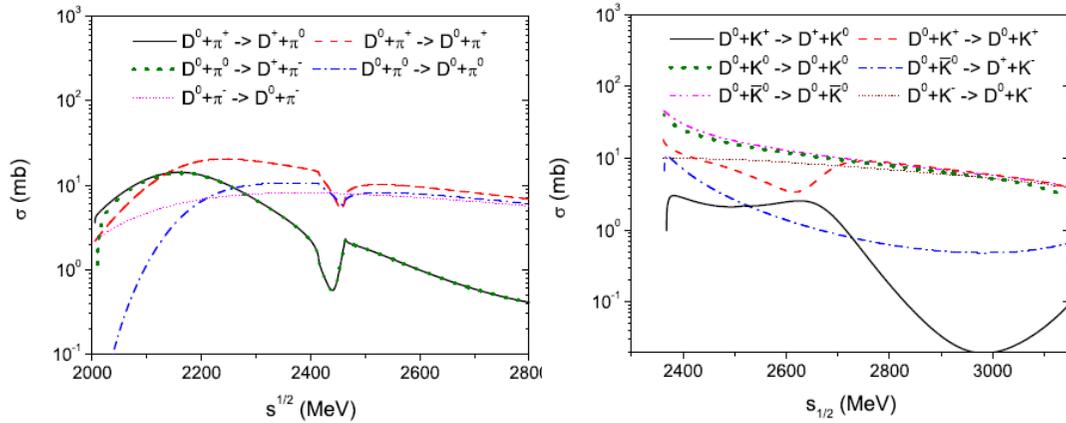
Unitarized scattering amplitude  $\rightarrow$  from solution of coupled-channel Bethe-Salpeter equations:

$$T = T + VGT$$



# D-meson scattering in the hadron gas

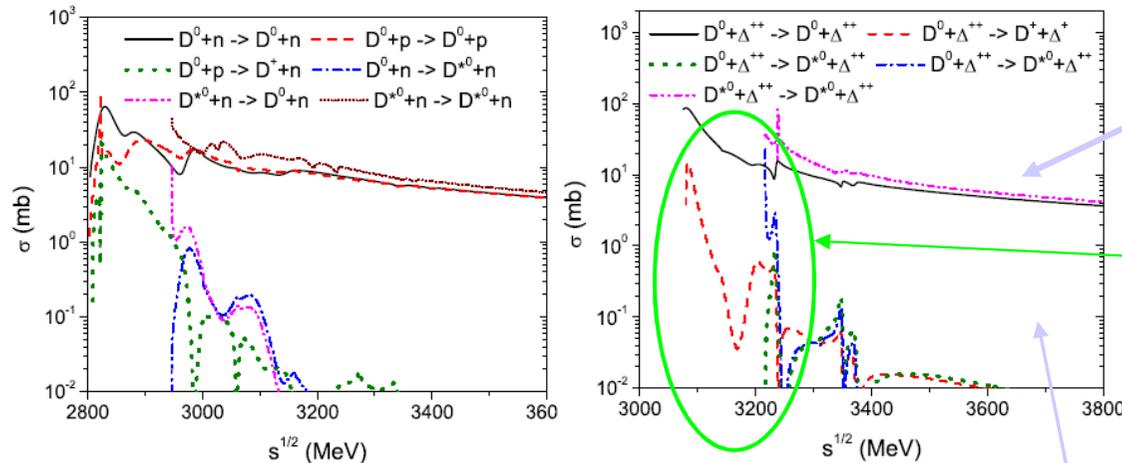
## 1. D-meson scattering with mesons



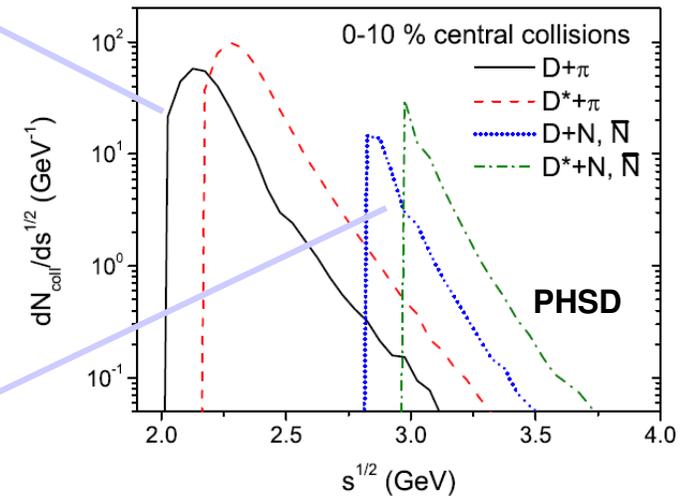
1a) cross sections with  $m = \rho, \omega, \phi, K^*, \dots$  taken as

$$\sigma(D, D^* + m) = 10 \text{ mb}$$

## 2. D-meson scattering with baryons



Distribution  $dN/ds^{1/2}$



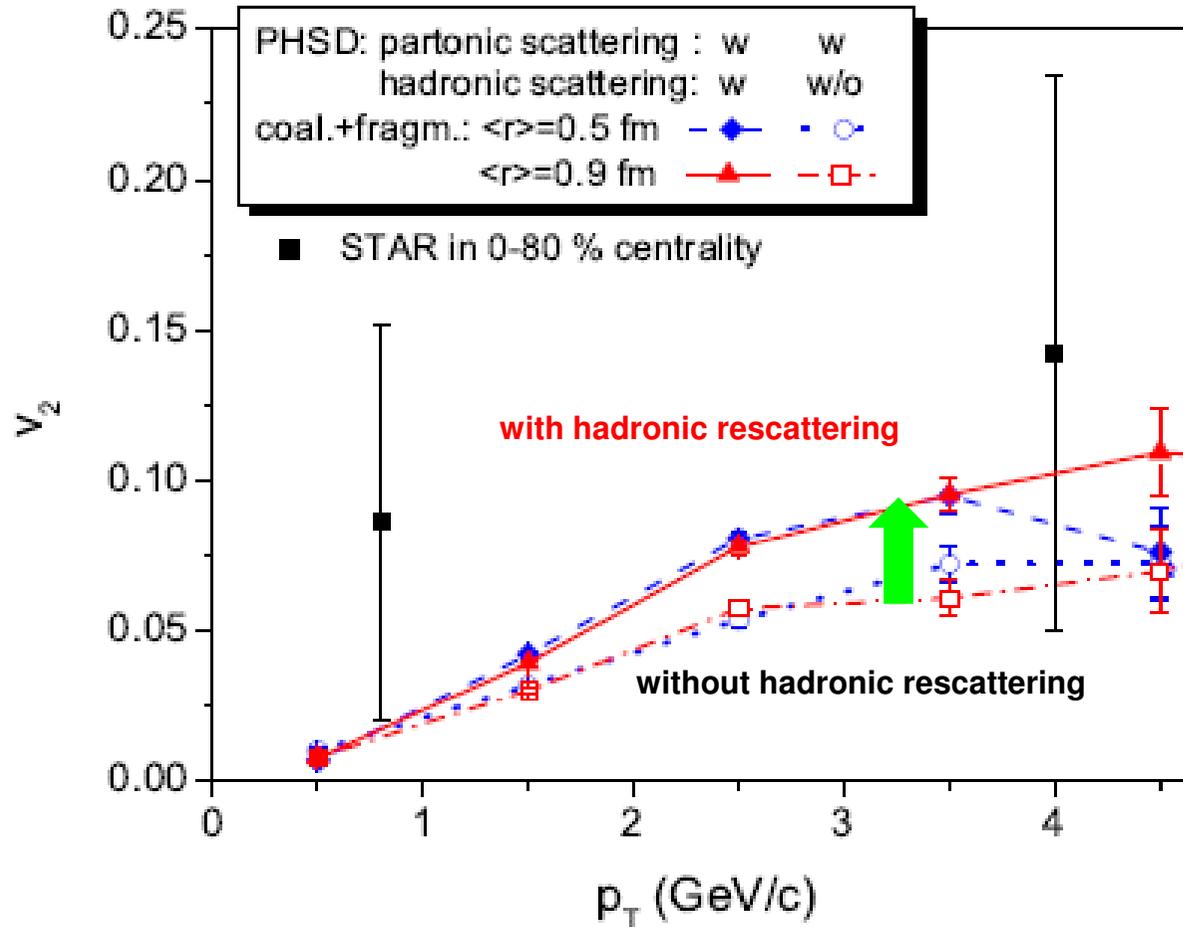
→ Strong **isospin dependence** and complicated structure (due to the resonance coupling) of  $D+m$ ,  $D+B$  cross sections!

→ Hadronic interactions become ineffective for the energy loss of  $D, D^*$  mesons at high transverse momentum (i.e. large  $s^{1/2}$ )





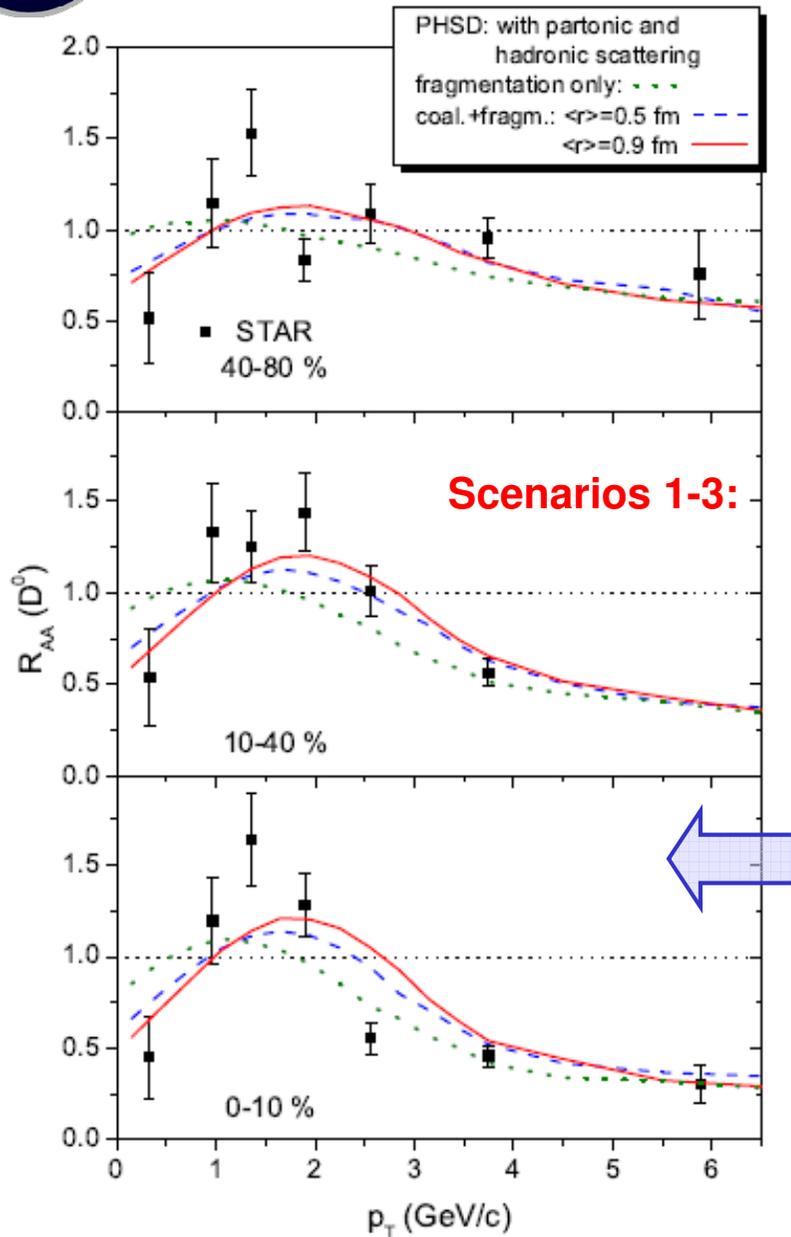
# D-meson elliptic flow $v_2$ at RHIC



- Hadronic rescattering substantially increases  $v_2$  at larger  $p_T$
- $v_2$  is less sensitive to the hadronization scenarios than  $R_{AA}$



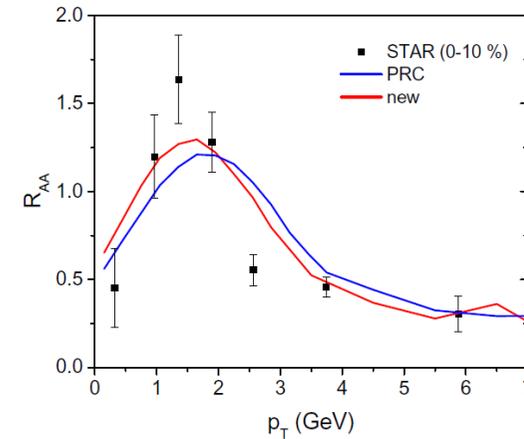
# $R_{AA}$ at RHIC



PHSD results:

with all rescattering (partonic and hadronic)

4: Dynamical Hadronization scenario



□ The height and position of the  $R_{AA}$  peak at low  $p_T$  depends on the hadronization scenario: coalescence/fragmentation!

→ PHSD: the STAR data are better described within scenario „coalescence with  $\langle r \rangle = 0.9$  fm + fragmentation“ and dynamical hadronization scenario



# Shadowing effect

Charm production cross section in **N\*N\*** in HIC:

$$\sigma_{c\bar{c}}^{N^*N^*}(s) = \sum_{i,j} \int dx_1 dx_2 R_i^A(x_1, Q) R_j^A(x_2, Q) \times f_i^N(x_1, Q) f_j^N(x_2, Q) \sigma_{c\bar{c}}^{ij}(x_1 x_2 s, Q).$$

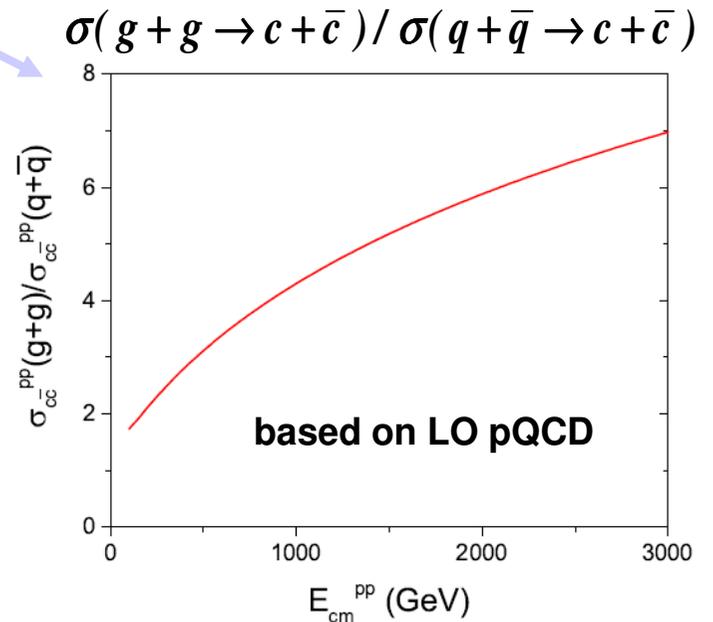
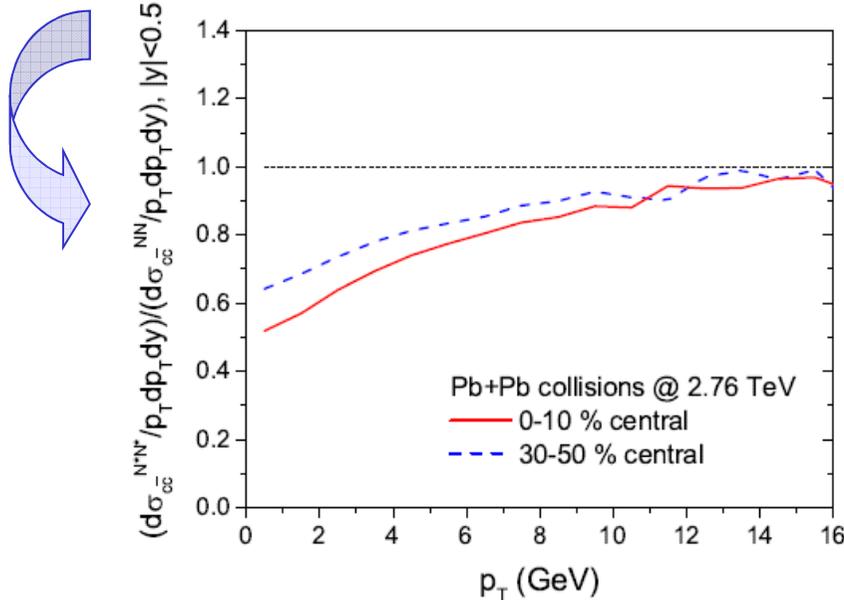
$$x_1 = \frac{M_T}{E_{cm}} e^y, \\ x_2 = \frac{M_T}{E_{cm}} e^{-y},$$

$$\sigma_{c\bar{c}}^{N^*N^*}(s) = \langle R_g^{Pb}(x_1, Q) R_g^{Pb}(x_2, Q) \rangle \sigma_{c\bar{c}}^{NN}(s)$$

$$\text{Scale } Q = (M_T^1 + M_T^2)/2$$

$R_i^A(x_1, Q)$ ,  $R_i^A(x_2, Q)$  for  $i=j=gluon$  are obtained from the EPS09 using that charm production is dominated by **gluon fusion**:

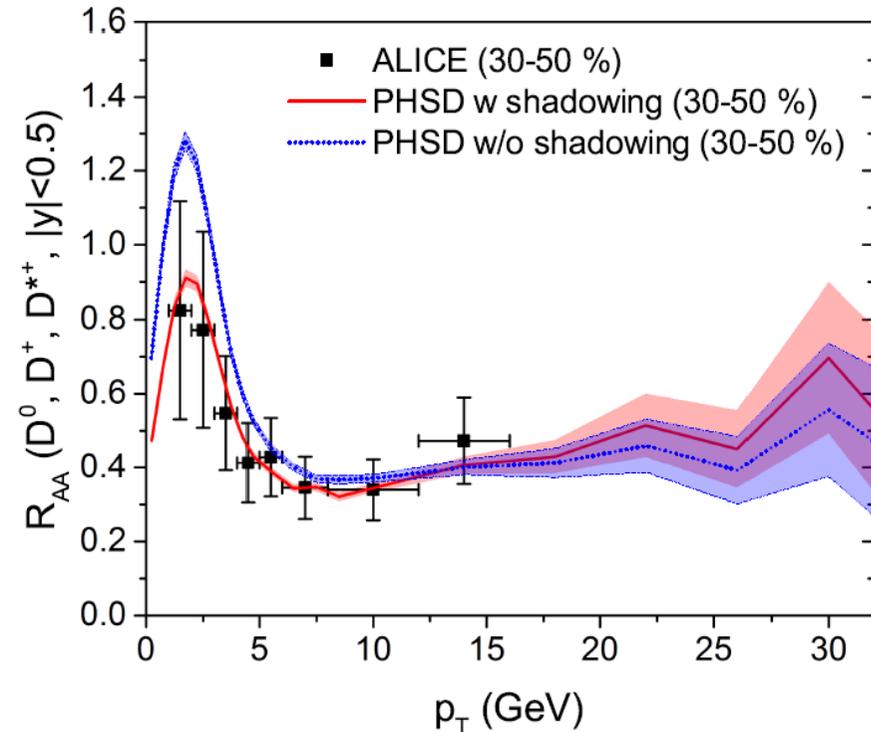
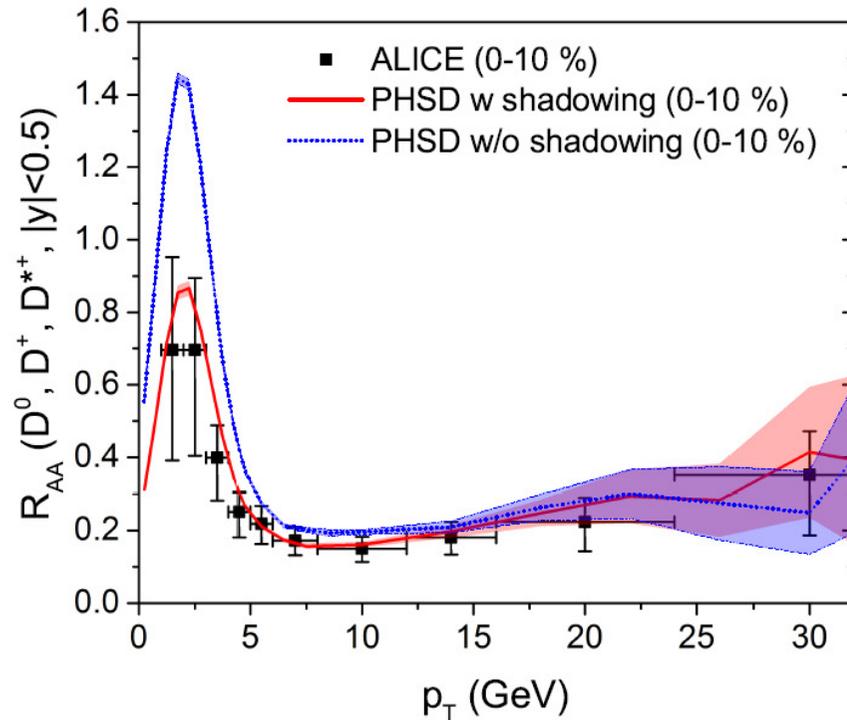
□ The (anti-)shadowing effect depends on the **impact parameter** in HIC:



T. Song et al., PRC (2016), arXiv:1512.0089



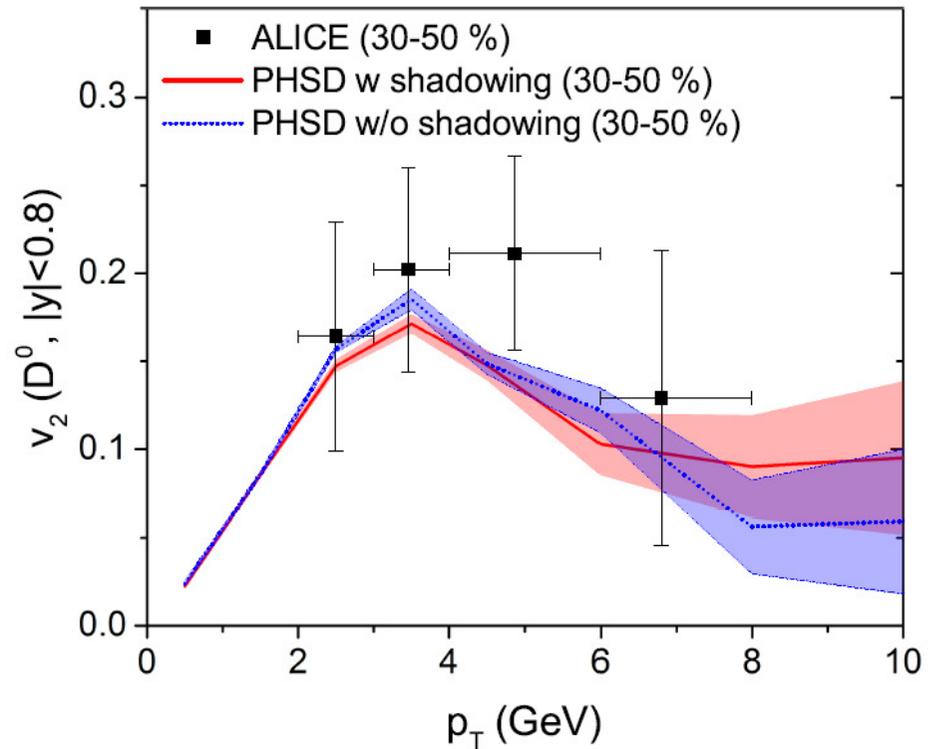
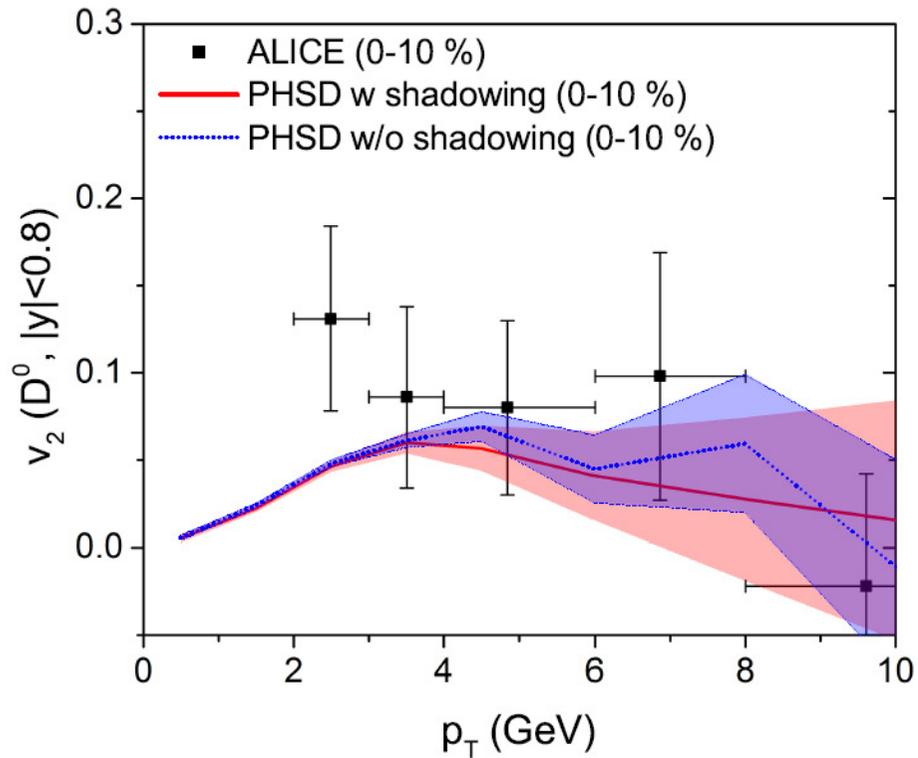
## Charm $R_{AA}$ at LHC



- in PHSD the **energy loss** of D-mesons at high  $p_T$  can be **dominantly attributed to partonic scattering**
- **Shadowing effect** suppresses the low  $p_T$  and slightly enhances the high  $p_T$  part of  $R_{AA}$
- **Hadronic rescattering** moves  $R_{AA}$  peak to higher  $p_T$



# Charm $v_2$ at LHC

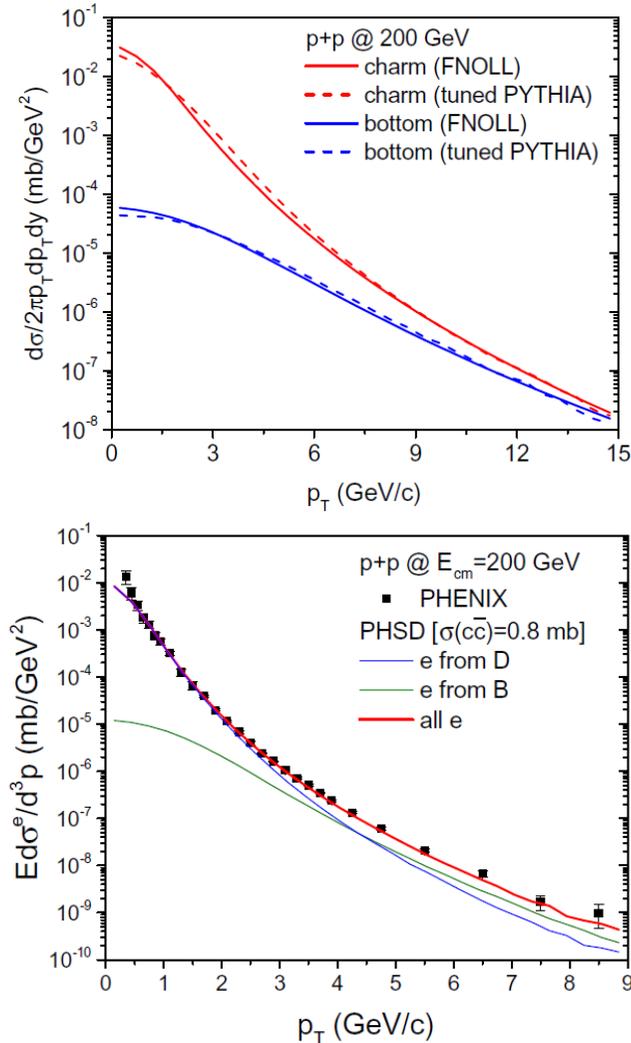


- Shadowing effect has small impact on  $v_2$
- **Hadronic rescattering** increases  $v_2$

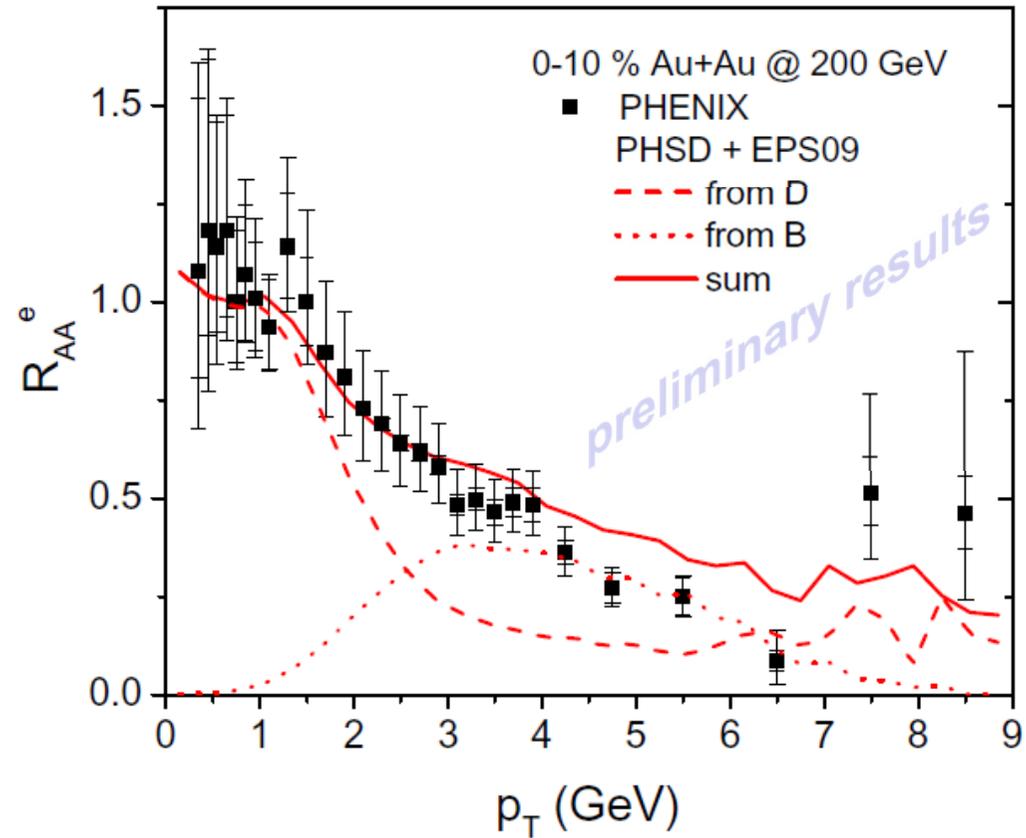


# Beauty contribution

## Beauty production in pp



## $R_{AA}$ from single electrons in Au+Au @ 200 GeV



□ Feed back from beauty contribution  
becoms dominant at  $p_T > 2.5 \text{ GeV}$

T. Song et al., in progress



**Thank you!**