

V+jet at NNLO

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QCD, EW and Tools at 100 TeV

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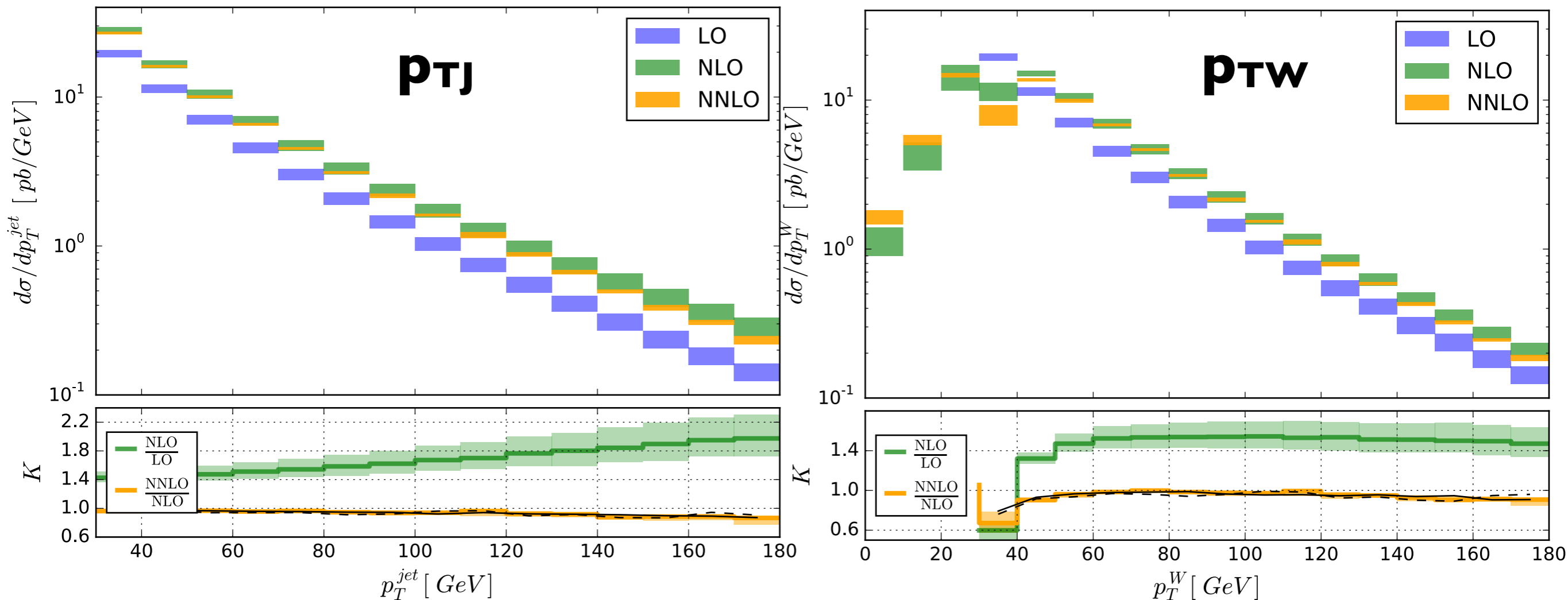


Outline

- Review of Wj phenomenology at 8 TeV
- Wj phenomenology at 100 TeV
- All $Wj@NNLO$ results based on the following calculation:
1504.02131 (Boughezal, Focke, Liu, FP)
- Comparison of exact NNLO with LoopSim (**preliminary**)
(in collaboration with Daniel Maître)

Phenomenology at 8 TeV

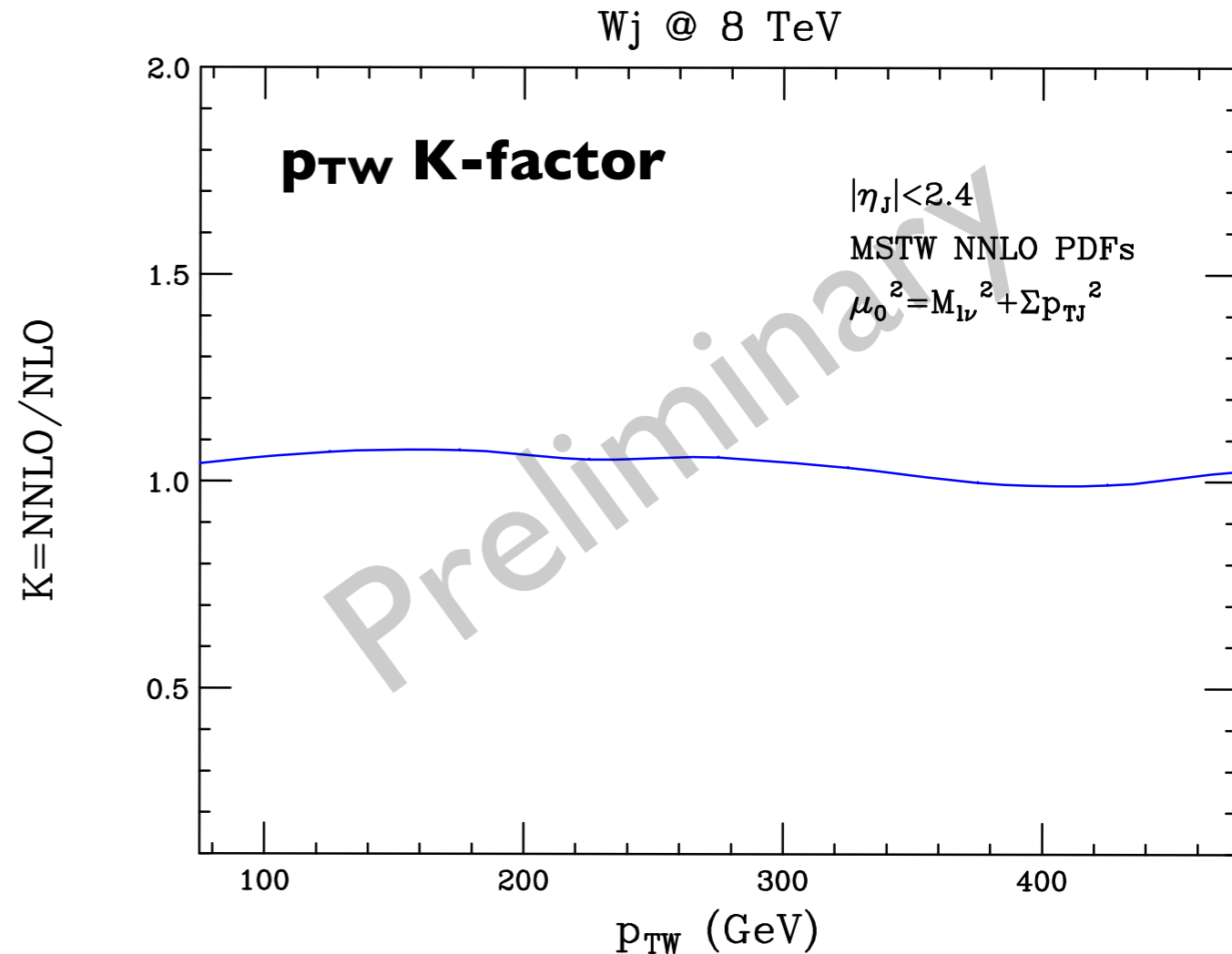
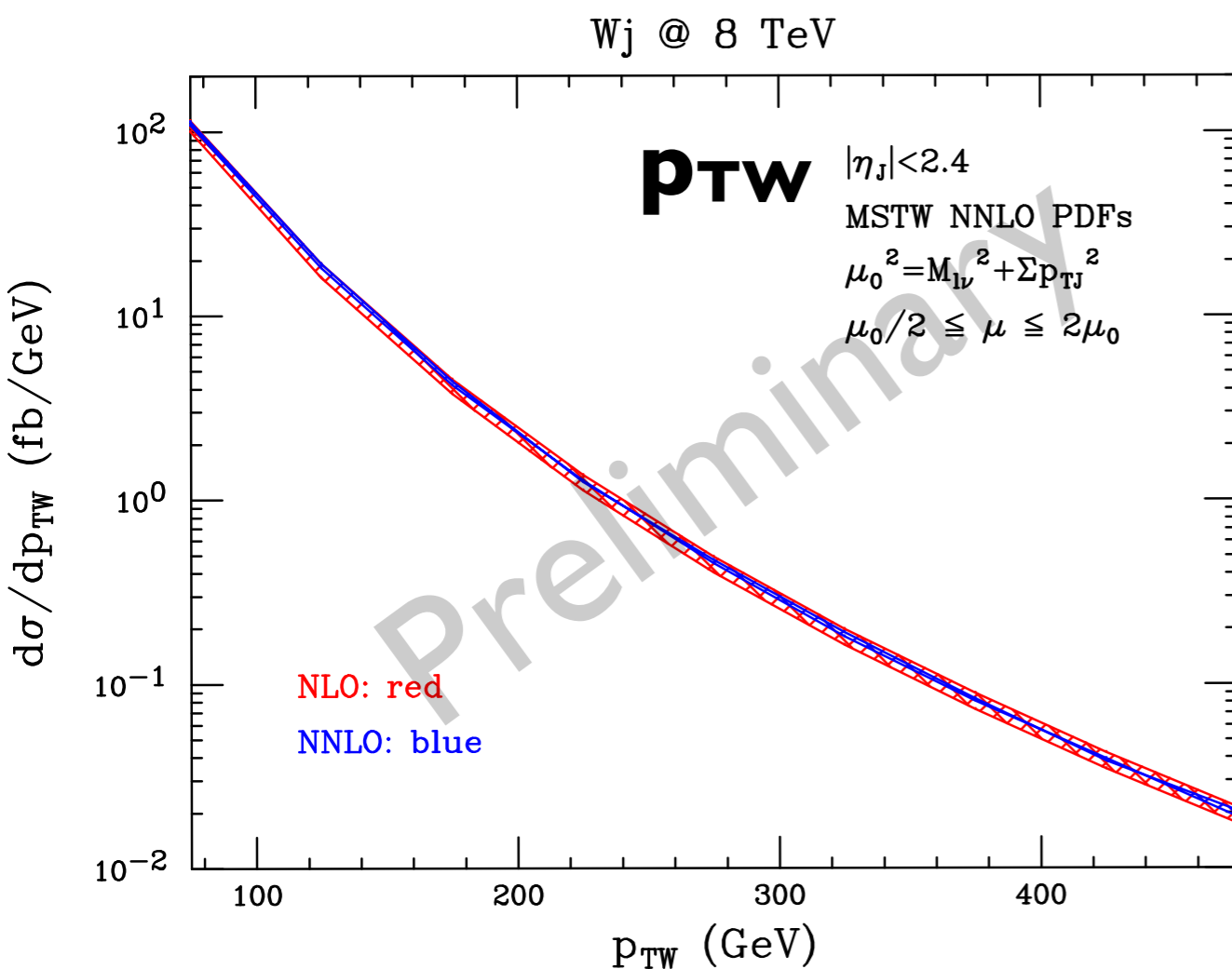
- Results are for CT14 PDFs, central scale $\mu=M_W$, estimate error with a factor-of-2 variation around this central scale



- NLO corrections large: 40% ranging to 80% for $p_{TJ} < 180$ GeV with a dependence on p_{TJ} , 40% for $p_{TW} < 180$ GeV
- NNLO corrections are $\sim 1\%$ with little dependence on kinematics

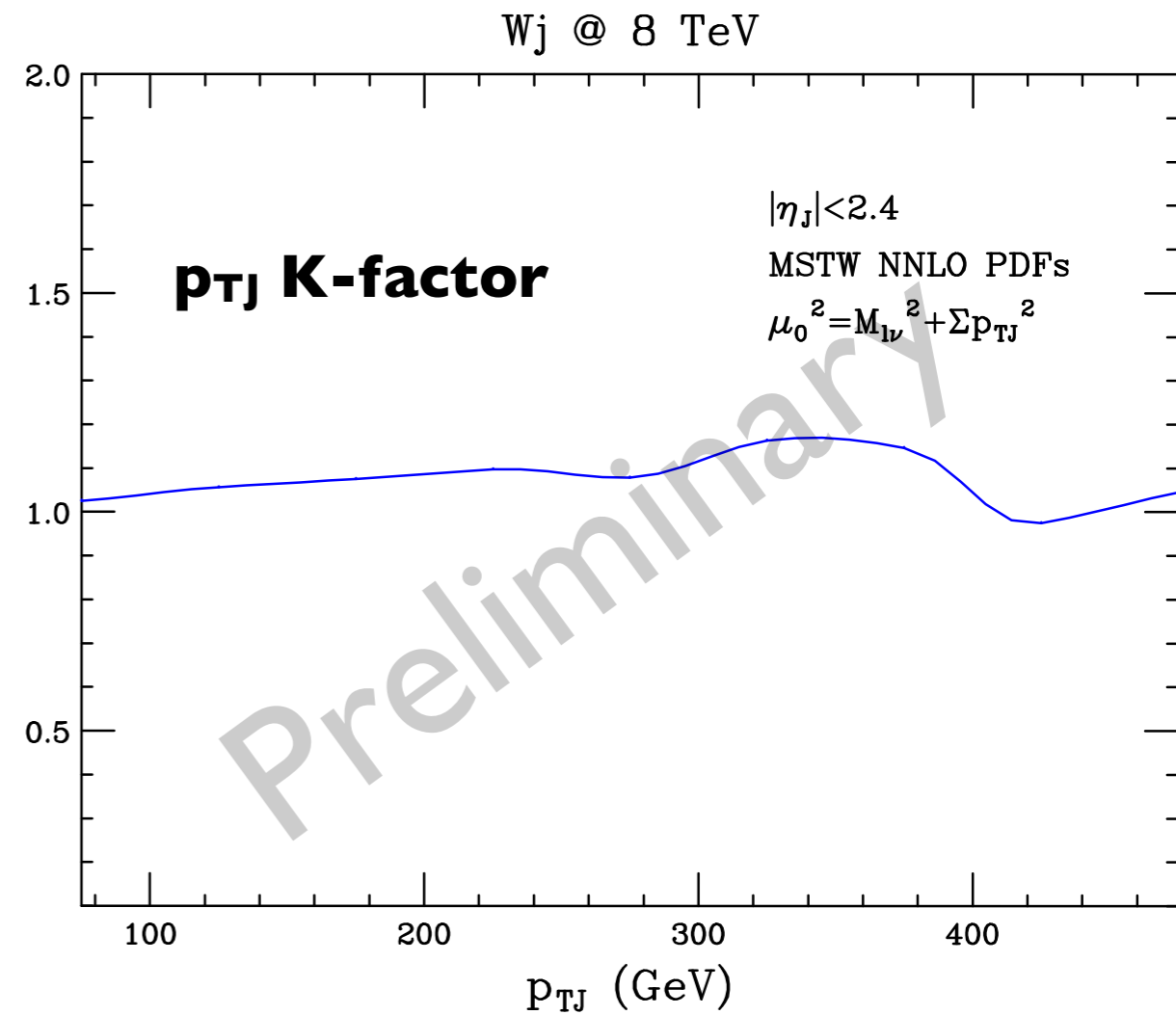
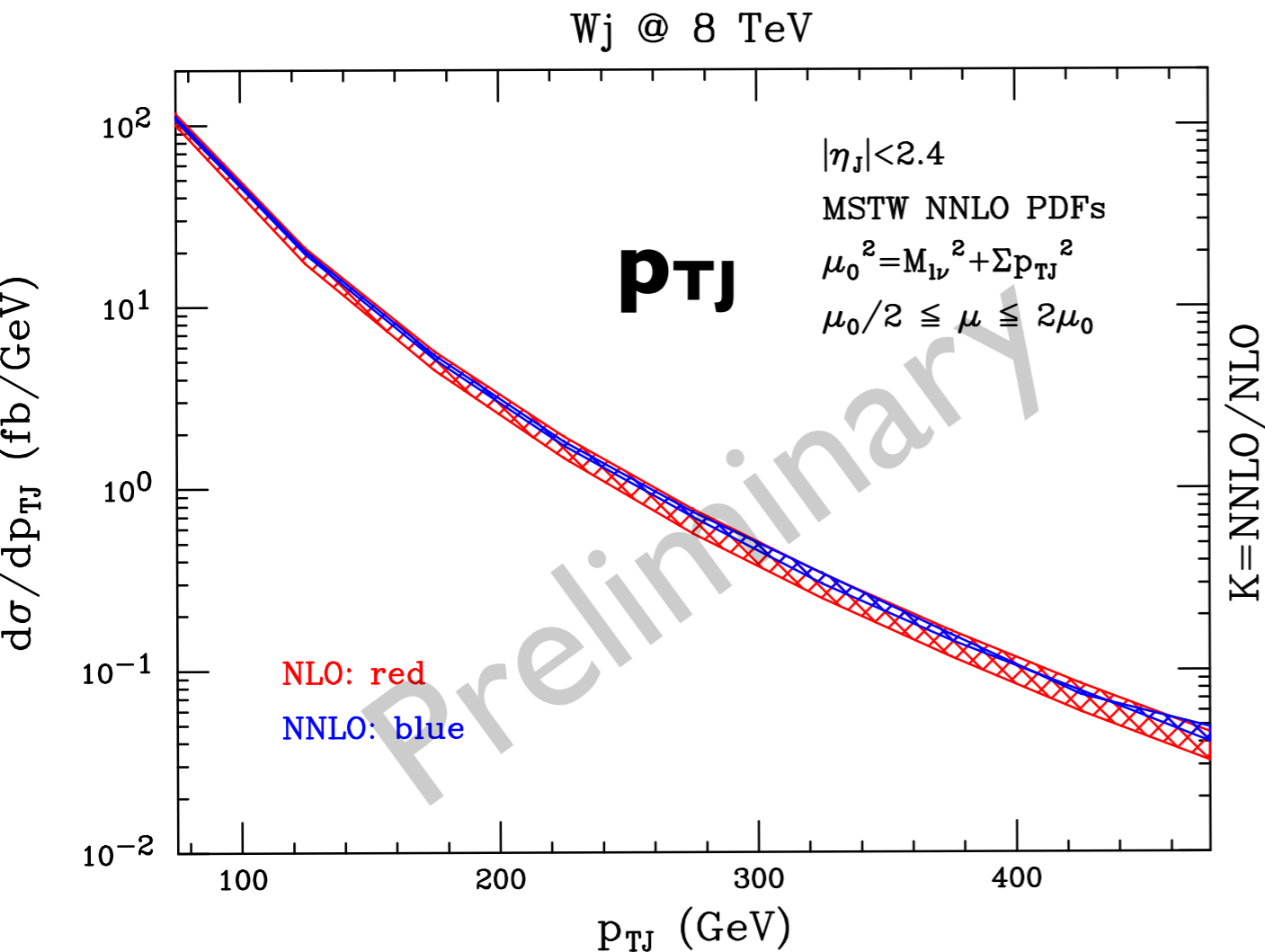
Phenomenology at 8 TeV: p_{TW}

- How do the perturbative corrections change when an extended kinematic range is studied? Switch to a dynamical scale $\mu^2 = M_{l\nu}^2 + \sum p_{Tj}^2$



- NNLO corrections to the p_{TW} distribution remain small, ~few percent, over the entire studied range

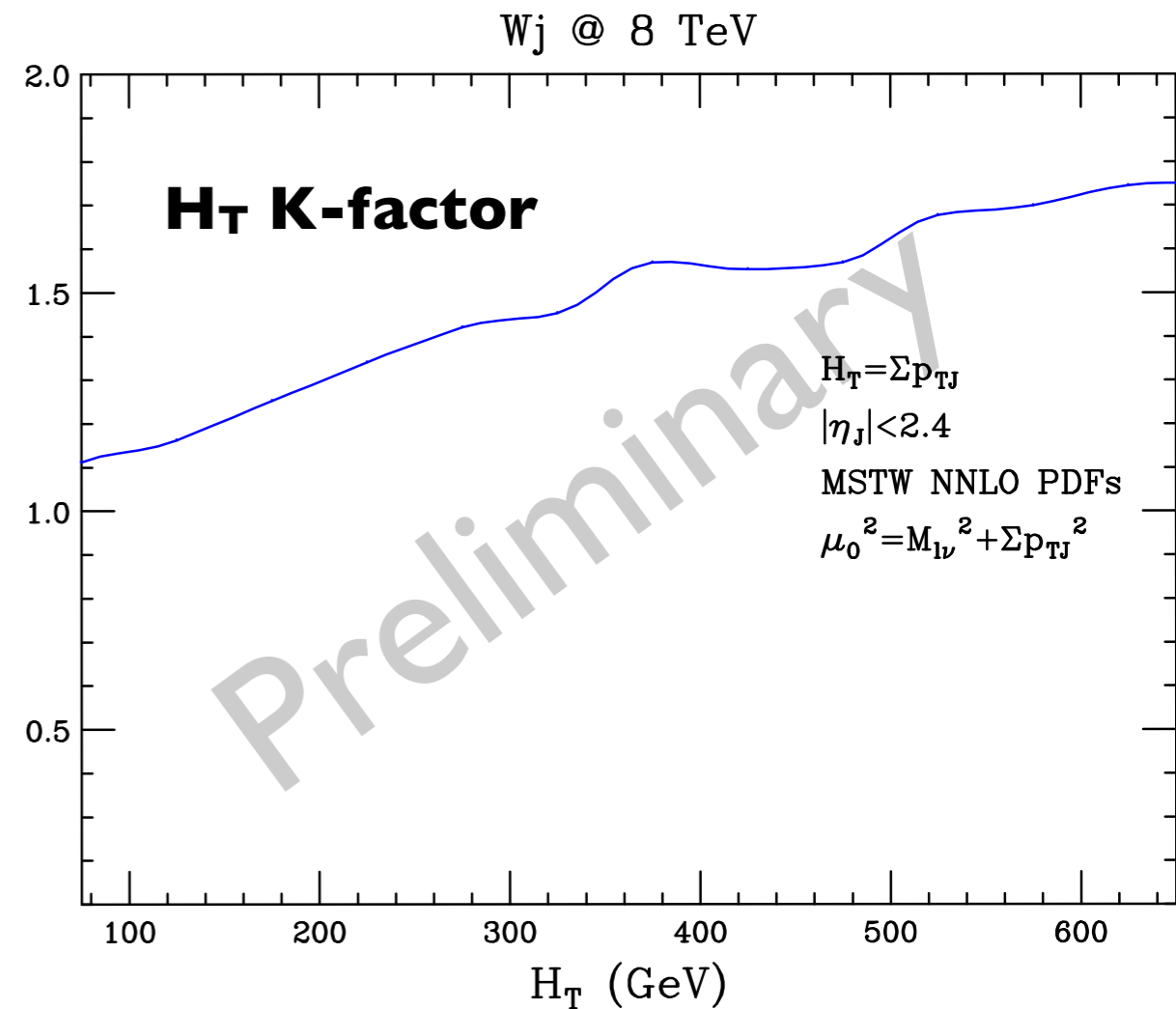
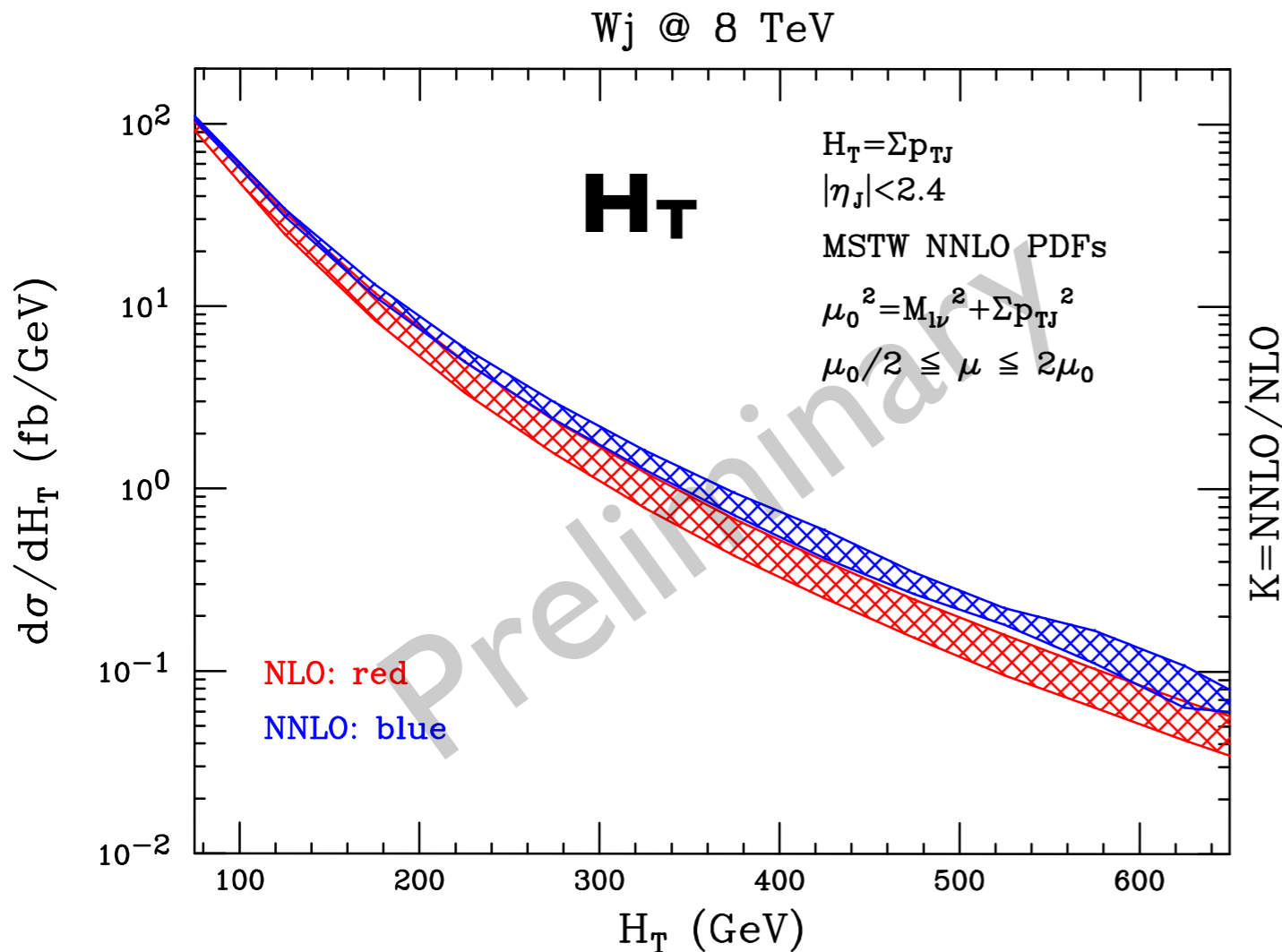
Phenomenology at 8 TeV: p_{TJ}



- NNLO corrections are again small for the p_{TJ} distribution over the entire studied range; slight rise for higher p_{TJ}

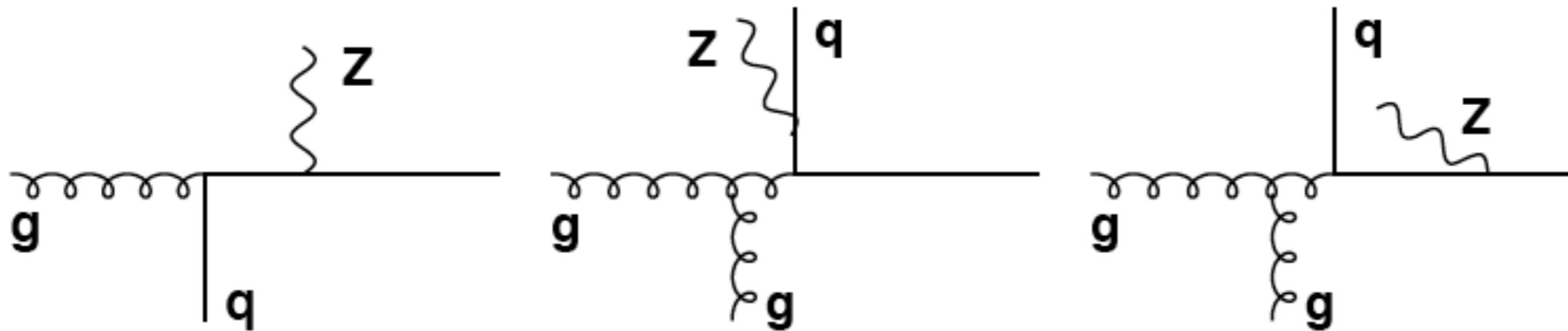
Phenomenology at 8 TeV: H_T

- Define $H_T = \sum p_{Tj}$. Larger corrections to this distribution; K-factor rises above 1.75 for $H_T > 600$ GeV



Giant K-factors and LoopSim

- The origin of these large corrections to the H_T distribution have been discussed in the literature (for example, see Salam, Rubin, Sapeta 1006.2144)

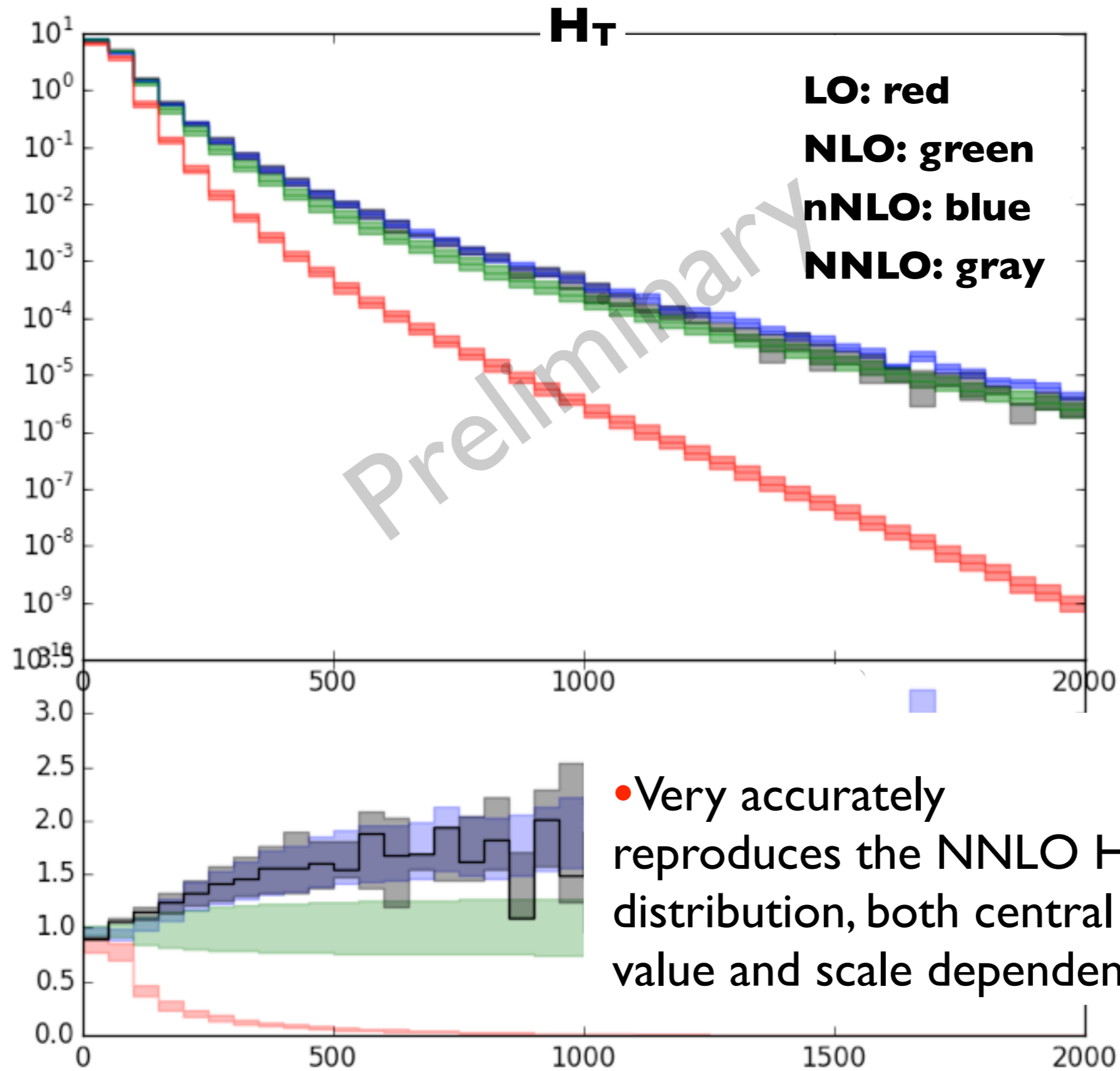


LO contribution to $V+j$; at large p_{Tj} or H_T , V must be a hard emission. Contributes $O(\alpha_S \alpha_{EW})$.

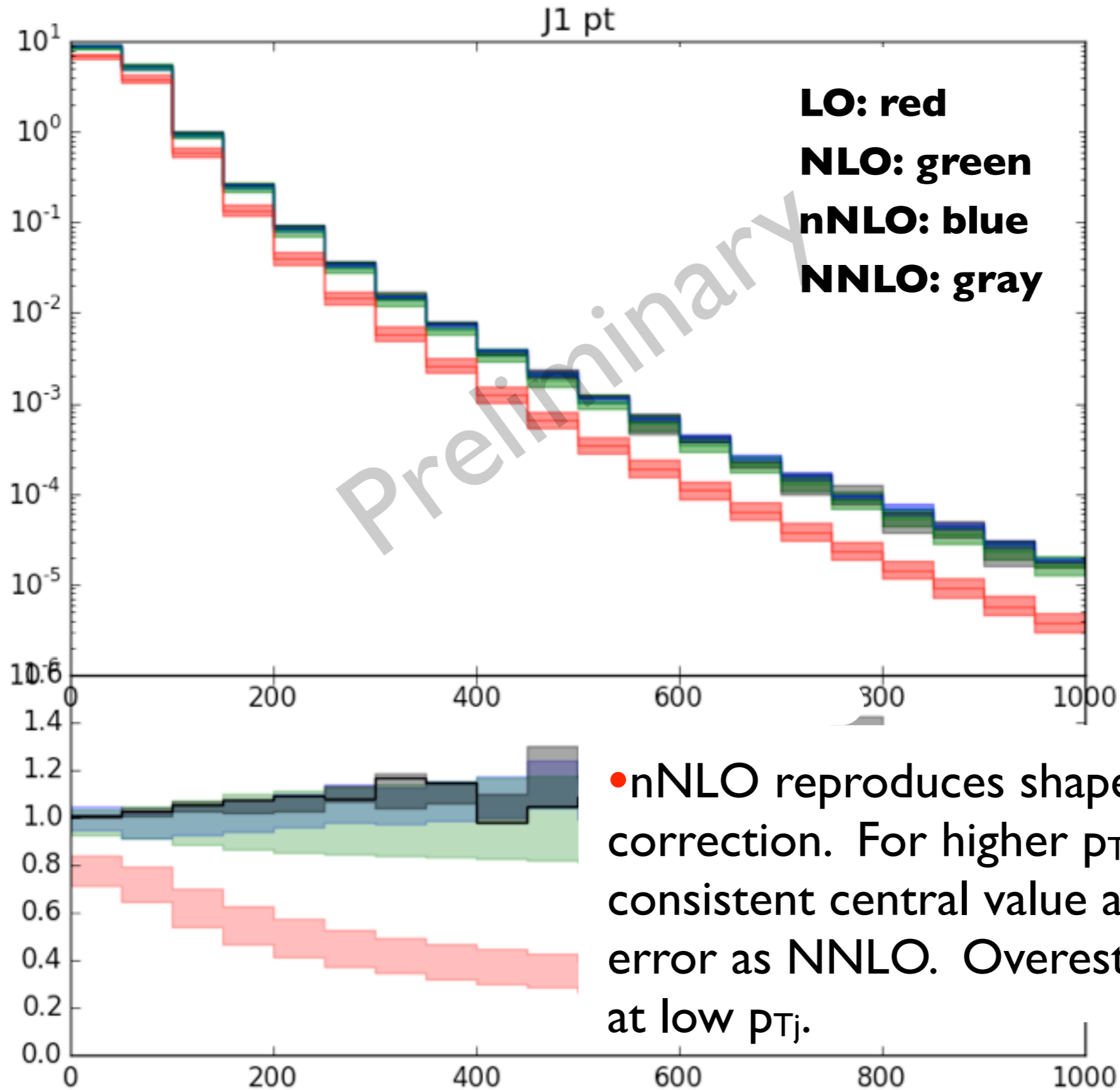
NLO contributions to $V+j$; at large p_{Tj} or H_T , can have soft +collinear V . Contributes $O(\alpha_S^2 \alpha_{EW} \ln^2(Q/M_V))$, with $Q=p_{Tj}$ or H_T , which becomes large at high energies.

- LoopSim (nNLO) is an approximation of the full NNLO calculation that captures these large logarithmic corrections; nNLO has full NLO Vj , and in addition the logarithmically enhanced NNLO corrections

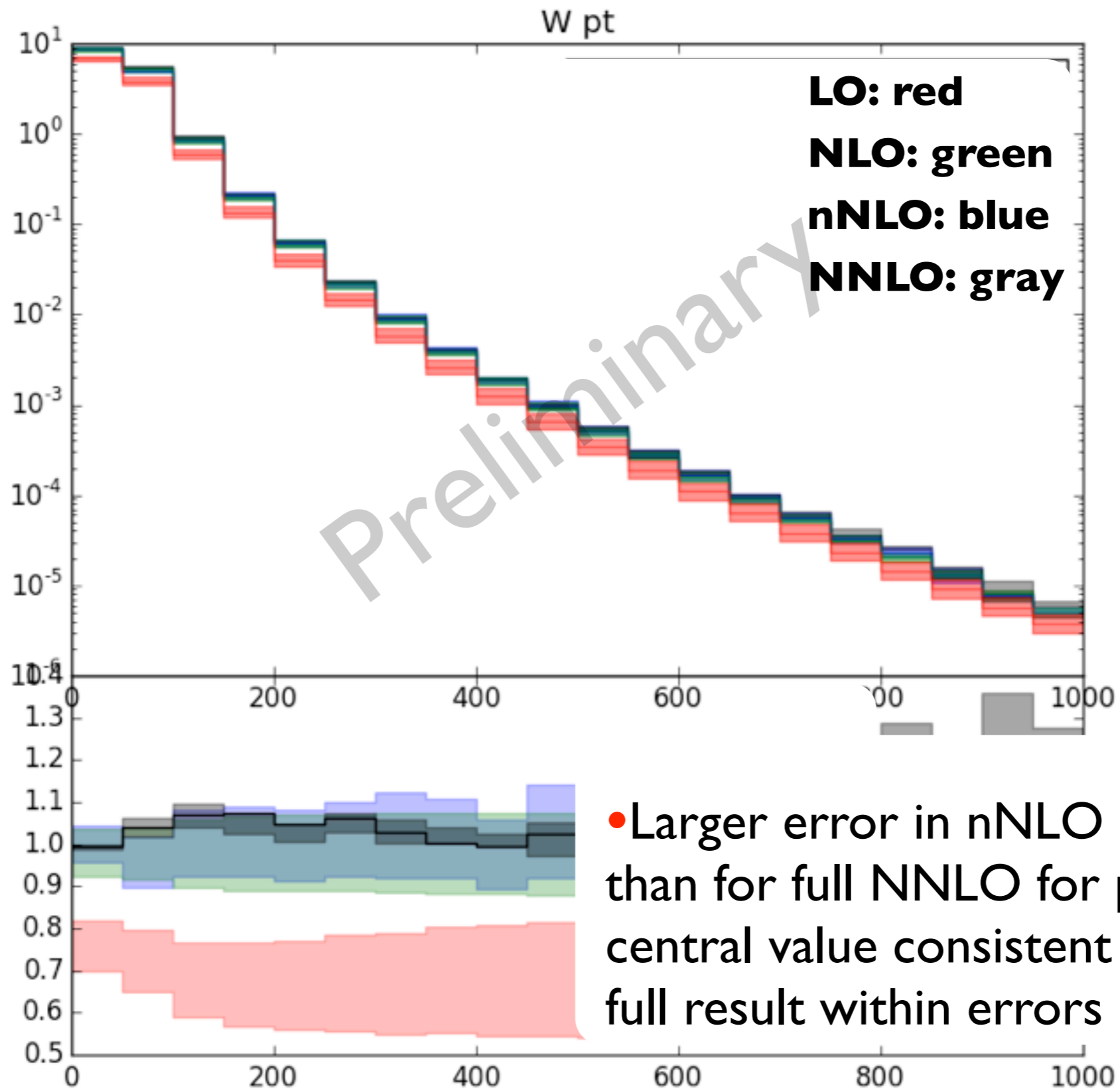
LoopSim at 8 TeV: H_T



LoopSim at 8 TeV: p_{Tj}

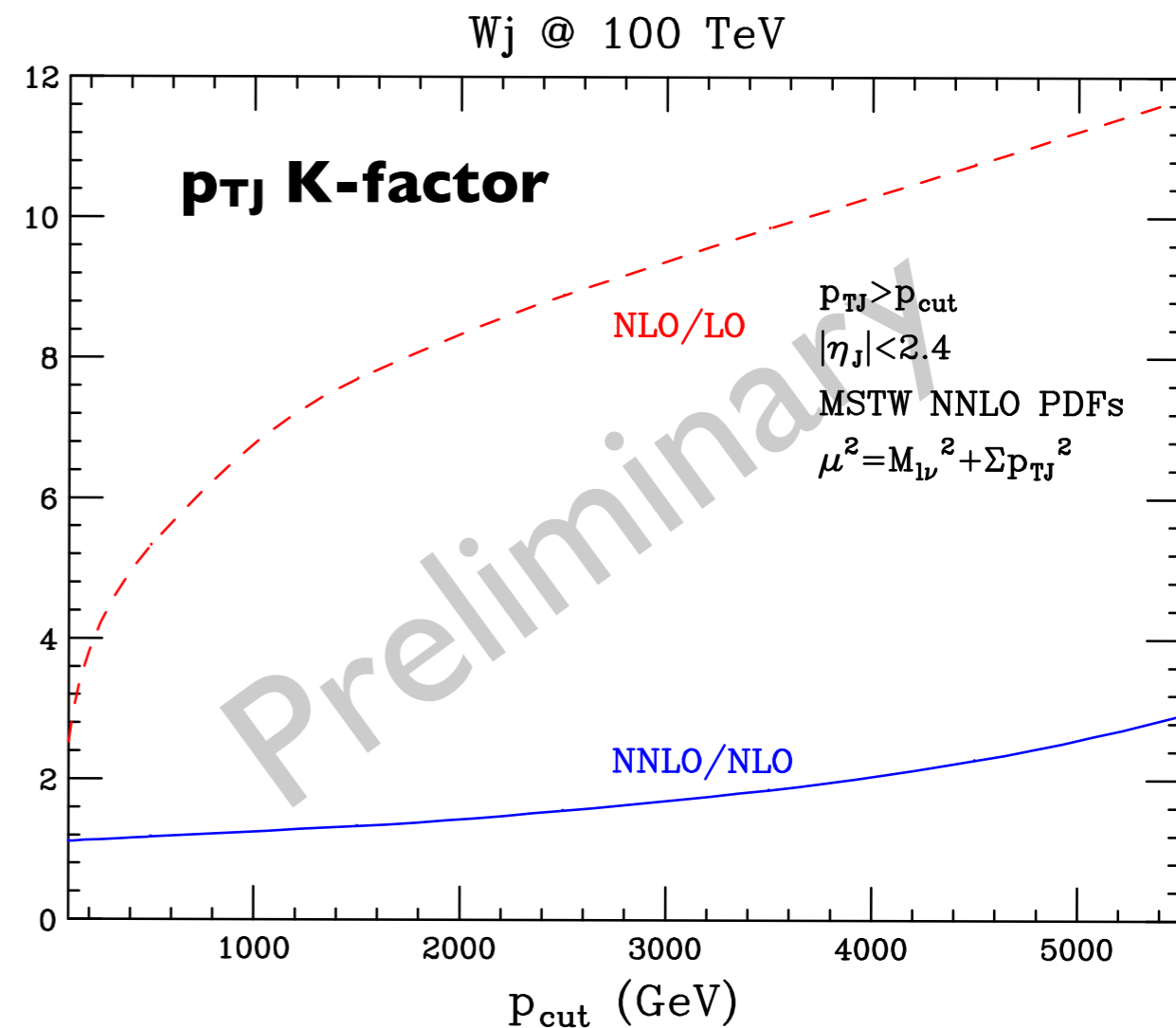
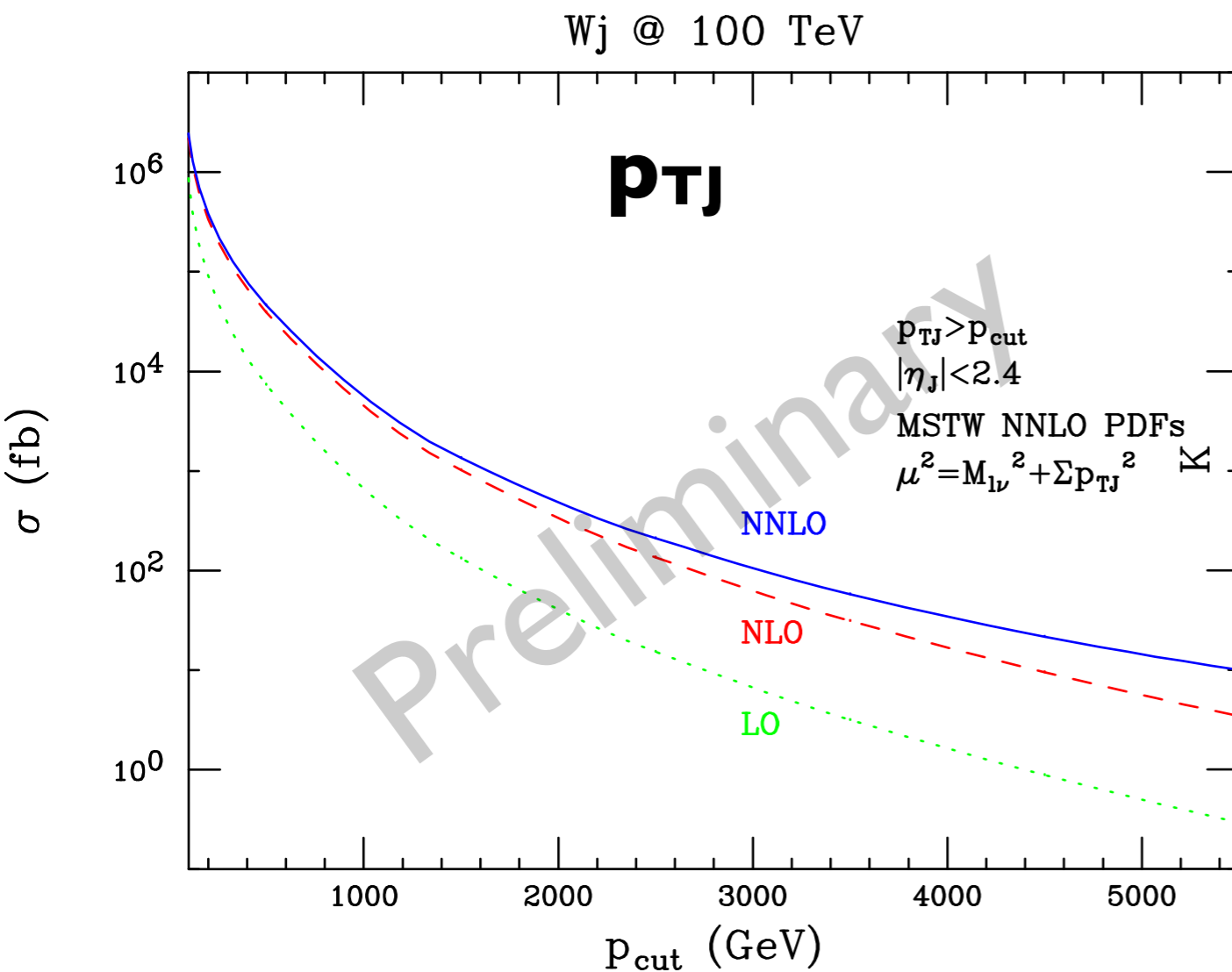


LoopSim at 8 TeV: p_{TW}



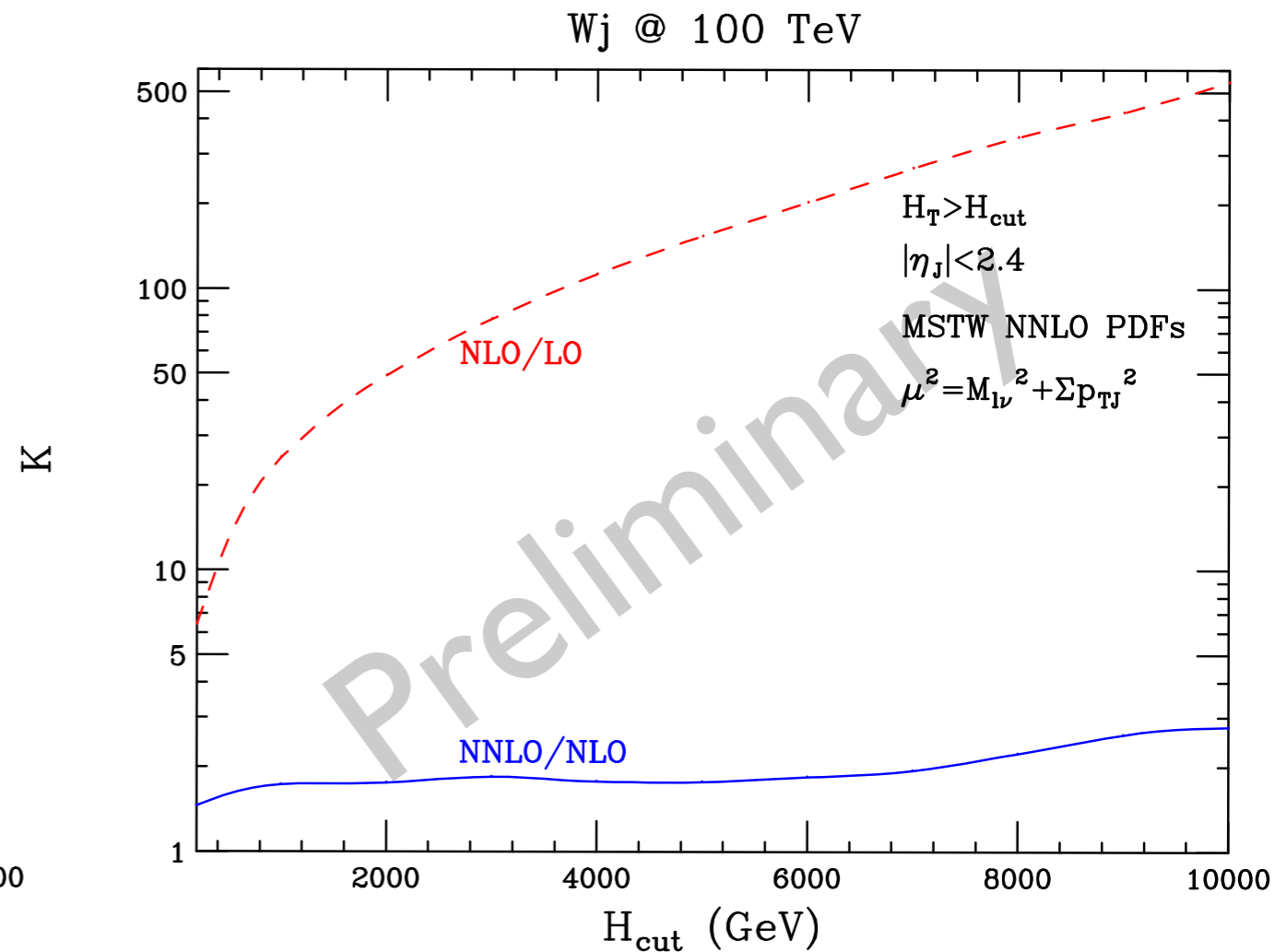
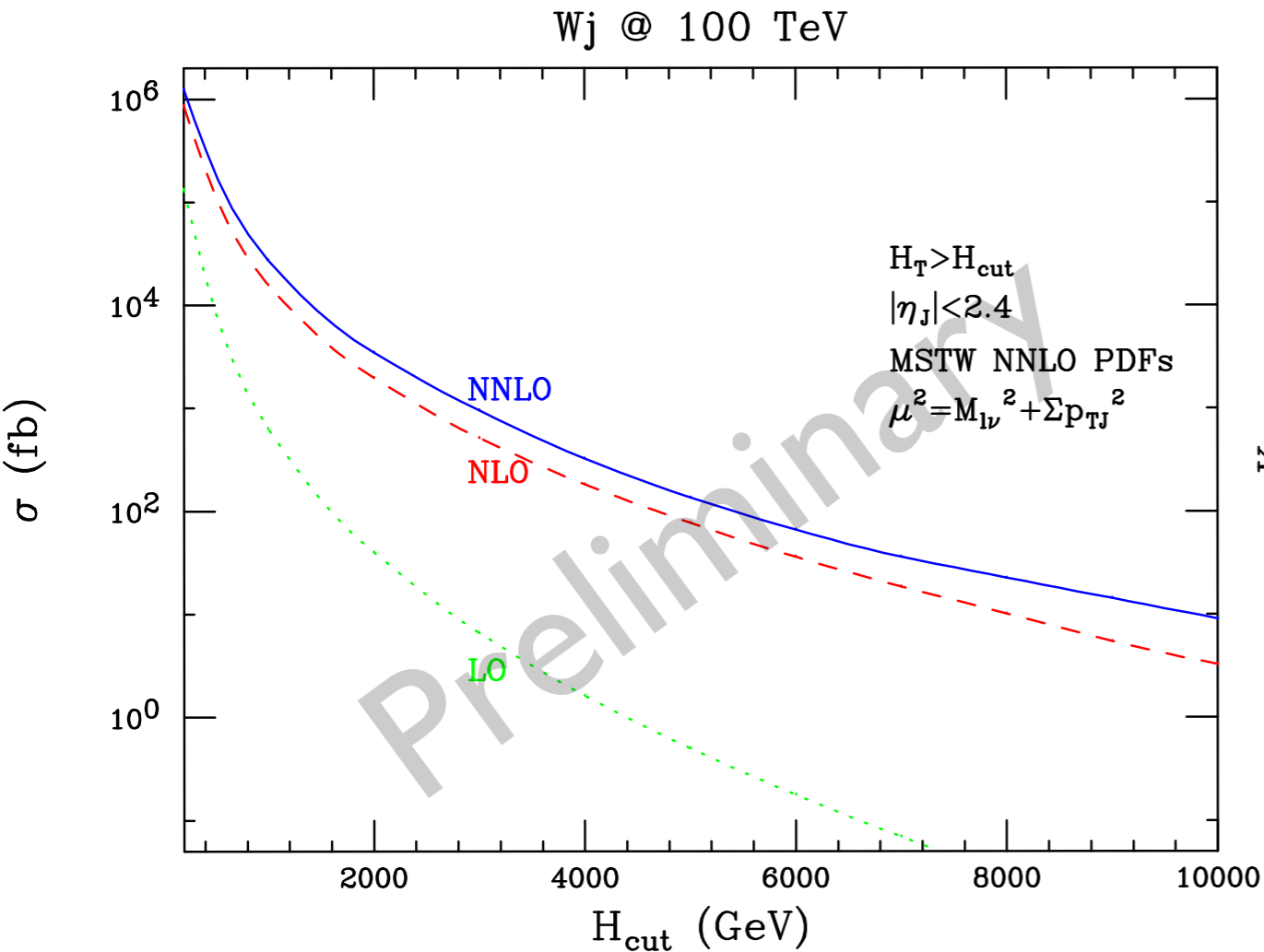
Phenomenology at 100 TeV: p_{TJ}

- Results are for MSTW PDFs, $\mu^2 = M_{1\nu}^2 + \sum p_{TJ}^2$



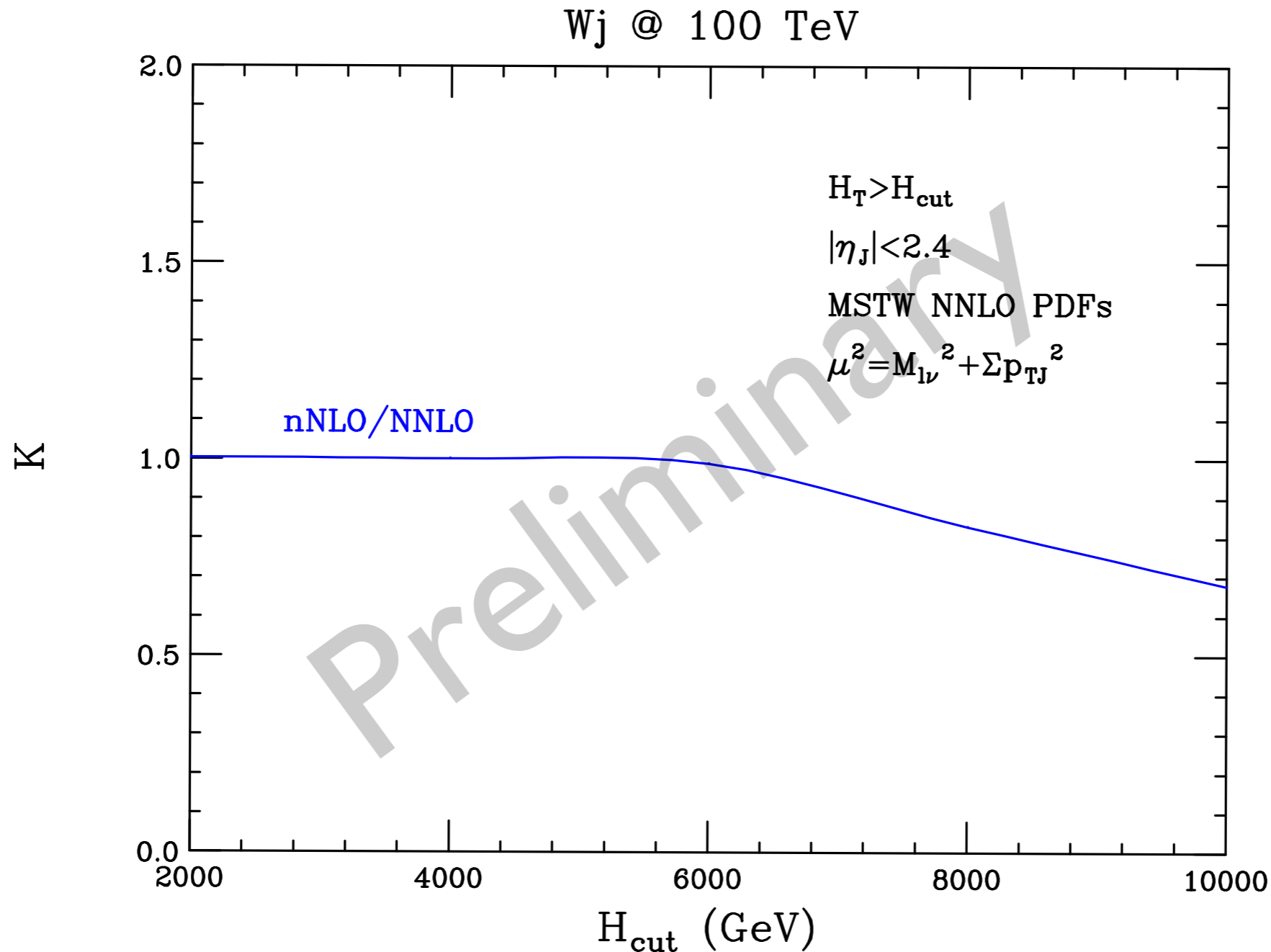
- K-factor at NLO grows above 10 for $p_{TJ} > 4$ TeV
- While NNLO corrections are much milder than NLO corrections, the NNLO/NLO K-factor still grows to 3 for $p_{TJ} > 5$ TeV.

Phenomenology at 100 TeV: H_T



- Enormous K-factor at NLO; grows above 500 for $H_T > 10$ TeV
- NNLO corrections similar to those found for p_{TJ} ; milder than NLO, but still sizable and increasing with H_T . The K-factor grows to $K=3$ for $H_T=10$ TeV.

LoopSim at 100 TeV: H_T



- LoopSim agrees well with NNLO up to 6 TeV, then becomes smaller by 30% at $H_T=10$ TeV; this effect is under investigation

Conclusions

- NNLO corrections to Wj are large and observable-dependent, and should be studied to ensure accurate 100 TeV phenomenology
- NNLO corrections grow with p_{Tj} and H_T , and the K-factor rises to $K=2-3$ in the multi-TeV energy range. While large, these corrections are much smaller than found at NLO.
- The LoopSim approach is found to give a reasonably good approximation to the full NNLO result. For observables dominated by large logarithmic corrections such as the high H_T region, it reproduces the central value and scale variation.
- For observables such as p_{TW} or the low p_{Tj} region, where the logarithmic corrections do not dominate, it gets close to the correct central value and overestimates the scale variation.
- Given the large available phase space for radiation at 100 TeV, and the still-large NNLO corrections at high energies, applications of the LoopSim approach beyond NNLO should be studied