

The squeezed limit of the bispectrum for multi-field inflation

Zac Kenton

Based on arxiv:1507.08629 with David Mulryne

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COSMO-15, Warsaw, Poland, 08/09/2015

Take-home message

- ▶ Standard δN calculates bispectrum in the equilateral limit
- ▶ We find the bispectrum in the squeezed limit

Outline

Primordial Curvature Perturbation ζ

Squeezed Limit of the Bispectrum

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Definition of ζ

The δN Formula

Correlation Functions of ζ

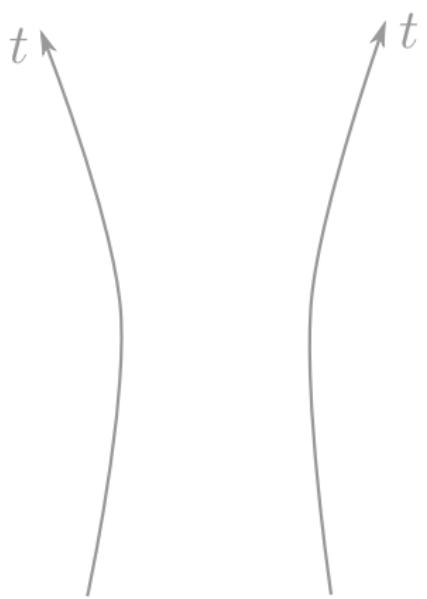
Squeezed Limit of the Bispectrum

Squeezed Limit

Single Field

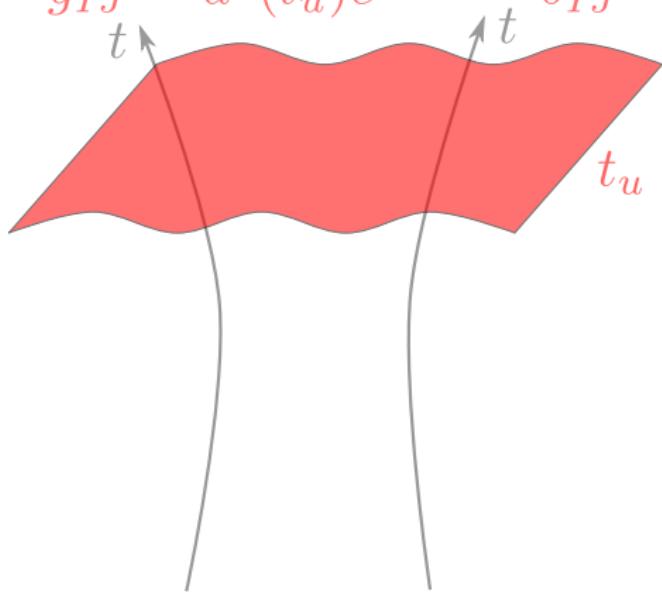
Multiple Fields

Observational Prospects

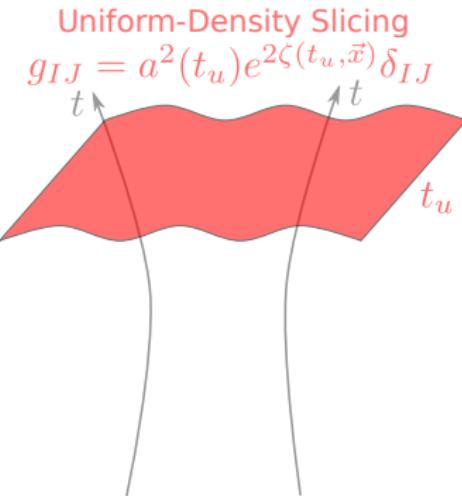


Uniform-Density Slicing

$$g_{IJ} = a^2(t_u) e^{2\zeta(t_u, \vec{x})} \delta_{IJ}$$



Primordial Curvature Perturbation, ζ



- ▶ This defines the ‘Primordial Curvature Perturbation’, ζ
- ▶ ζ important for observations of CMB and LSS
- ▶ ζ originates from inflation

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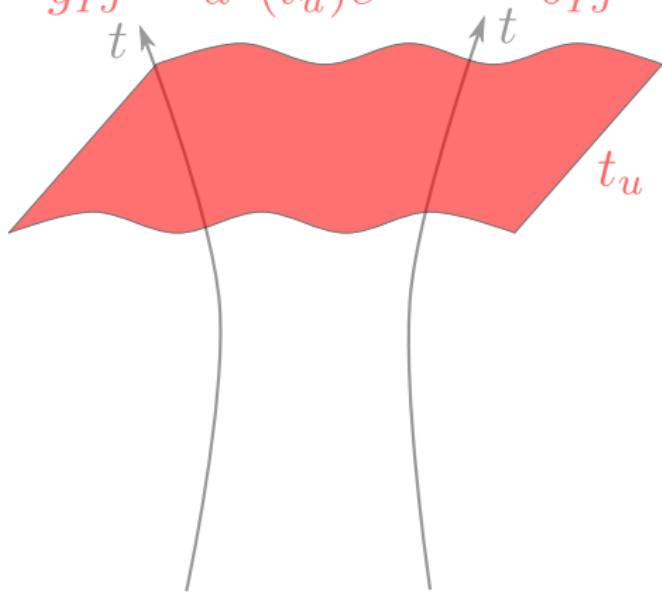
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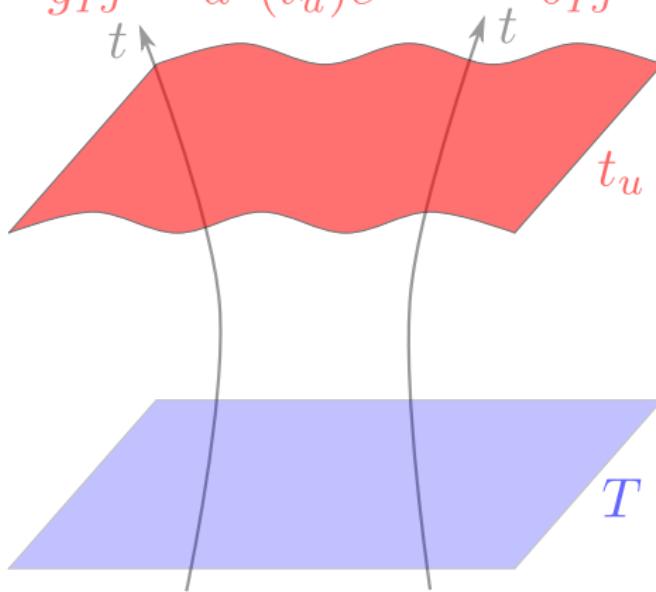
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$$g_{IJ} = a^2(T) \delta_{IJ}$$

Flat Slicing

δN For Inflation

- ▶ Inflation with multiple scalar fields $\phi_i(T, \vec{x}) = \phi_i^{(T)} + \delta\phi_i^{(T)}(\vec{x})$.

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$$\zeta(t_u, \vec{x}) = N_i^{(T)} \delta\phi_i^{(T)}(\vec{x}) + \frac{1}{2} N_{ij}^{(T)} \delta\phi_i^{(T)}(\vec{x}) \delta\phi_j^{(T)}(\vec{x}) + \dots$$

where $N_i^{(T)} \equiv \frac{\partial N_0(t_u, T)}{\partial \phi_i^{(T)}}$ with $N_0(t_u, T) \equiv \int_T^{t_u} H(t) dt$

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- So ζ at t_u can be expressed in perturbations of inflationary fields at early time T , multiplied by background quantities

[Sasaki, Stewart '96; Lyth, Rodriguez '05; ...]

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Bispectrum of ζ from δN

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$k_1 \leq k_2 \leq k_3$. Want bispectrum $\langle \zeta_{\vec{k}_1} \zeta_{\vec{k}_2} \zeta_{\vec{k}_3} \rangle$ in terms of

$$\langle \delta\phi_{i,\vec{k}_1}^{(T)} \delta\phi_{j,\vec{k}_2}^{(T)} \rangle \sim \Sigma_{ij}^{(T)}(k_1)$$

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- ▶ Equilateral limit $k_1 \approx k_2 \approx k_3 \approx k_*$
 - ▶ all modes exit at about the same time, in which case
$$\Sigma_{ij}^{(3)}(k_A) \approx \Sigma_{ij}^{(*)}(k_*) \Big|_{k_* \mapsto k_A}$$
$$\alpha_{ijk}^{(3)}(k_1, k_2, k_3) \approx \alpha_{ijk}^{(*)}(k_1, k_2, k_3) \Big|_{k_1 \approx k_2 \approx k_3}$$
have been calculated
[Stewart, Lyth '93; Nakamura, Stewart '96; Seery, Lidsey '05; ...]

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[Stewart, Lyth '93; Nakamura, Stewart '96; Seery, Lidsey '05;...]
- ▶ Squeezed limit $k_1 \ll k_2, k_3$
 - ▶ We need $\Sigma_{ij}^{(3)}(k_1)$ and $\alpha_{ijk}^{(3)}(k_1, k_2, k_3)$. This is our work.

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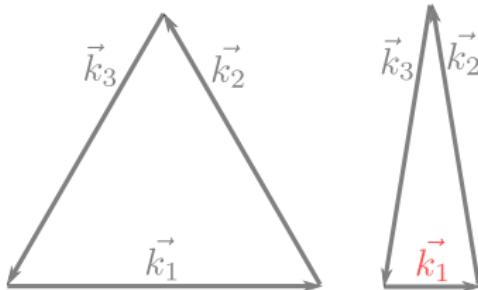
Single Field

Multiple Fields

Observational Prospects

Squeezed Limit

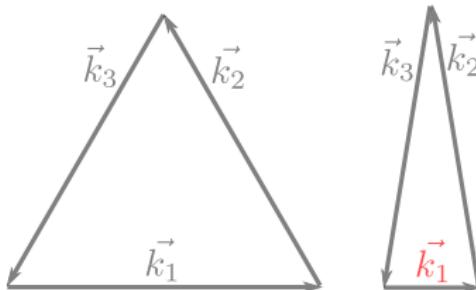
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- ▶ Bispectrum: Momentum conservation \Rightarrow closed triangle



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Squeezed Limit

- ▶ Soft limit: $\vec{k}_1 \ll \text{other } k_a$
- ▶ Bispectrum: Momentum conservation \Rightarrow closed triangle



- ▶ Soft limit is a squeezed triangle \equiv ‘squeezed limit’
- ▶ Observationally good: CMB, LSS, halo bias, μ -distortion
- ▶ Theoretically simple
- ▶ Sensitive to number of light fields during inflation!

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Single Field Case

$$\lim_{k_1 \ll k_3} \langle \zeta_{\vec{k}_1} \zeta_{\vec{k}_2} \zeta_{\vec{k}_3} \rangle = -(n_s - 1) \langle \zeta_{\vec{k}_1} \zeta_{-\vec{k}_1} \rangle \langle \zeta_{\vec{k}_2} \zeta_{\vec{k}_3} \rangle$$
$$n_s - 1 = -0.032 \pm 0.006$$

- ▶ In squeezed limit f_{NL} is small
- ▶ Result independent of slow roll!

[Maldacena '03; Creminelli, Zaldarriaga '04; Cheung et al. '08;...]

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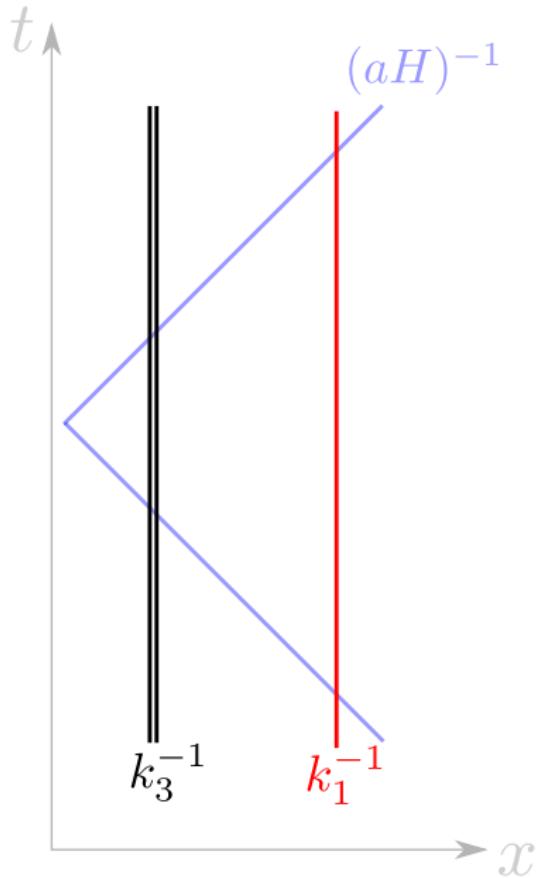
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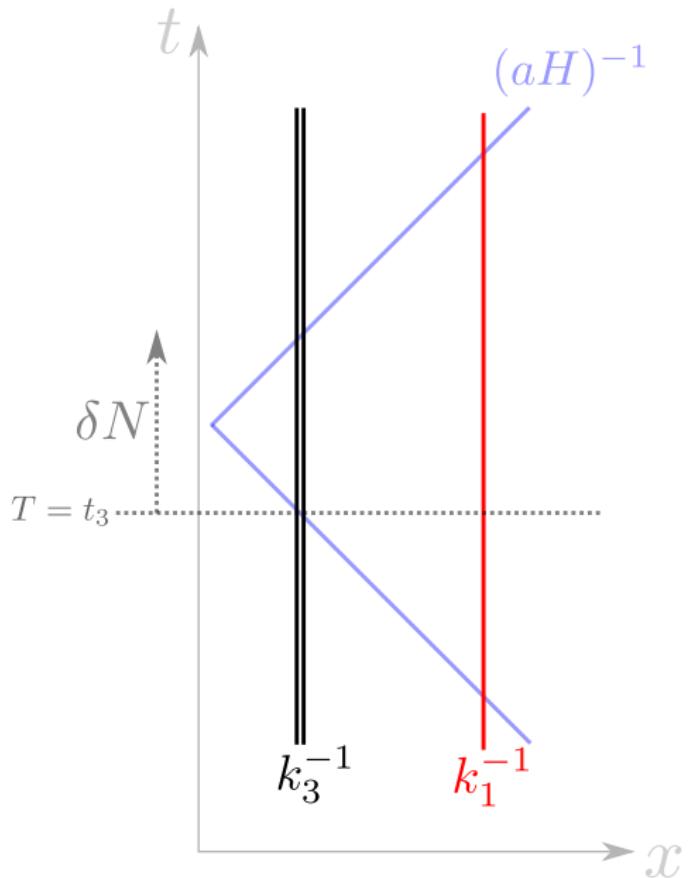
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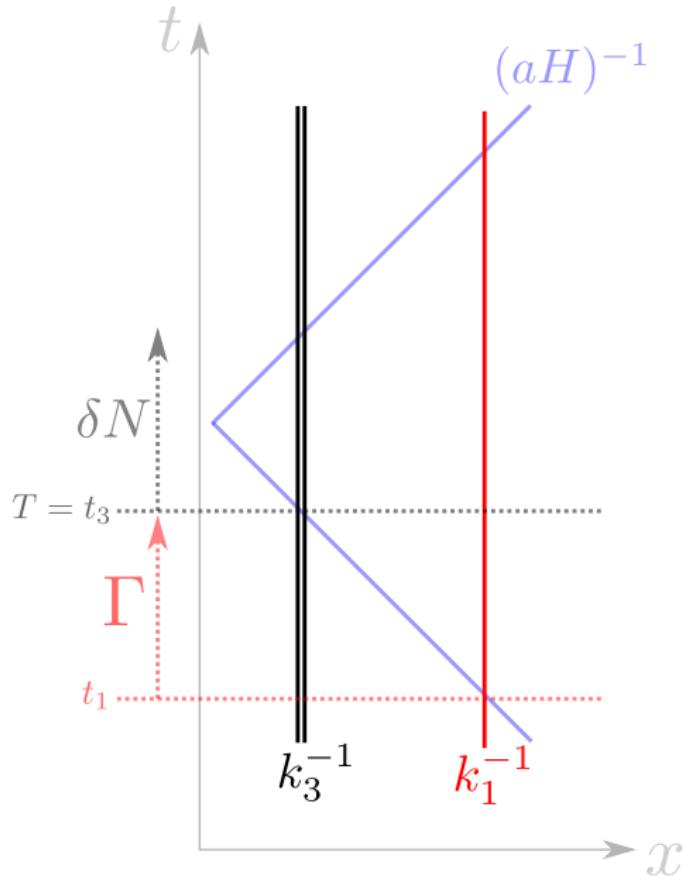
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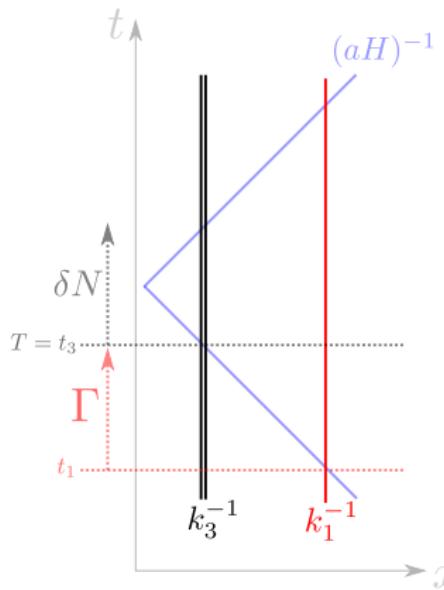


Multifield Case



Multifield Case

- Multifield: perturbations evolve
- Relate later perturbation at t_3 to earlier one at t_1



$$\delta\phi_i^{(3)}(\mathbf{x}) = \Gamma_{ij}\delta\phi_j^{(1)}(\mathbf{x}) + \dots$$

$$\Gamma_{ij} \equiv \frac{\partial \phi_i^{(3)}}{\partial \phi_j^{(1)}}$$

[ZK & Mulryne '15; Yokoyama et al. '07;...]

- Previous work has approximations to this

$$\Gamma_{ij} \approx \delta_{ij} - (\text{small}),_{ij} \log\left(\frac{k_3}{k_1}\right) + \dots$$

but not valid for large $\log(k_3/k_1)$

[Byrnes et al. '10, Dias, Ribeiro, Seery '13;...]

Multifield Case

- Squeezed limit:

$$\lim_{k_1 \ll k_3, k_2} \langle \zeta_{\vec{k}_1} \zeta_{\vec{k}_2} \zeta_{\vec{k}_3} \rangle \sim N_i^{(3)} N_j^{(3)} N_k^{(3)} \alpha_{ijk}^{(3)}(k_1, k_2, k_3) + 2N_i^{(3)} N_{jk}^{(3)} N_l^{(3)} \Sigma_{ij}^{(3)}(k_1) \Sigma_{kl}^{(3)}(k_3)$$

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- First find $\Sigma_{ij}^{(3)}(k_1)$:
- Γ -evolution gives

$$\Sigma_{ij}^{(3)}(k_1) = \Gamma_{il} \Gamma_{jm} \Sigma_{lm}^{(1)}(k_1) = \Gamma_{il} \Gamma_{jl} \frac{H^{(1)2}}{2k_1^3}$$

Multifield Case: Find α

- ▶ Next find $\alpha_{ijk}^{(3)}(k_1, k_2, k_3)$
- ▶ Equilateral case known [Seery, Lidsey '05]
- ▶ Squeezed case [ZK & Mulryne '15]

$$\lim_{k_1 \ll k_3, k_2} \alpha_{ijk}^{(3)}(k_1, k_2, k_3) \approx \Sigma_{im}^{(3)}(k_1) \frac{\partial}{\partial \phi_m^{(3)}} \Sigma_{jk}^{(3)}(k_3).$$

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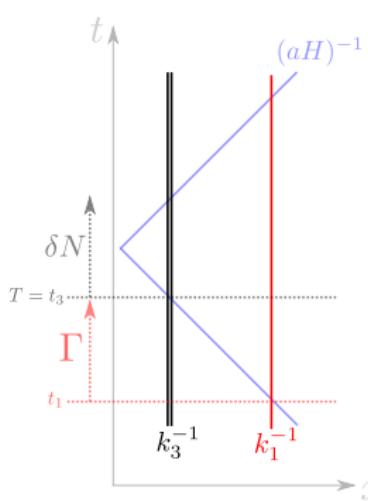
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- ▶ Our proof uses soft limit technique
- ▶ Similar results in [Allen et al. '05; Li & Wang '08]

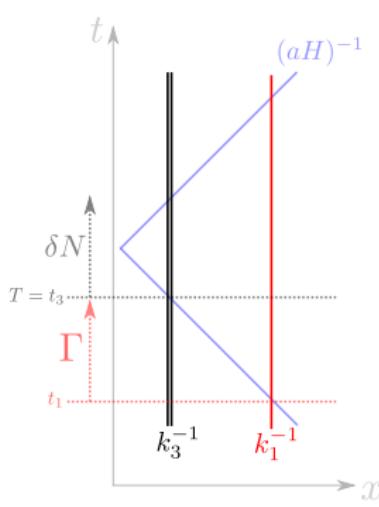
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Put these into:

$$\begin{aligned} \lim_{k_1 \ll k_3, k_2} \langle \zeta_{\vec{k}_1} \zeta_{\vec{k}_2} \zeta_{\vec{k}_3} \rangle &\sim N_i^{(3)} N_j^{(3)} N_k^{(3)} \alpha_{ijk}^{(3)}(k_1, k_2, k_3) \\ &+ 2N_i^{(3)} N_{jk}^{(3)} N_l^{(3)} \Sigma_{ij}^{(3)}(k_1) \Sigma_{kl}^{(3)}(k_3) \end{aligned}$$

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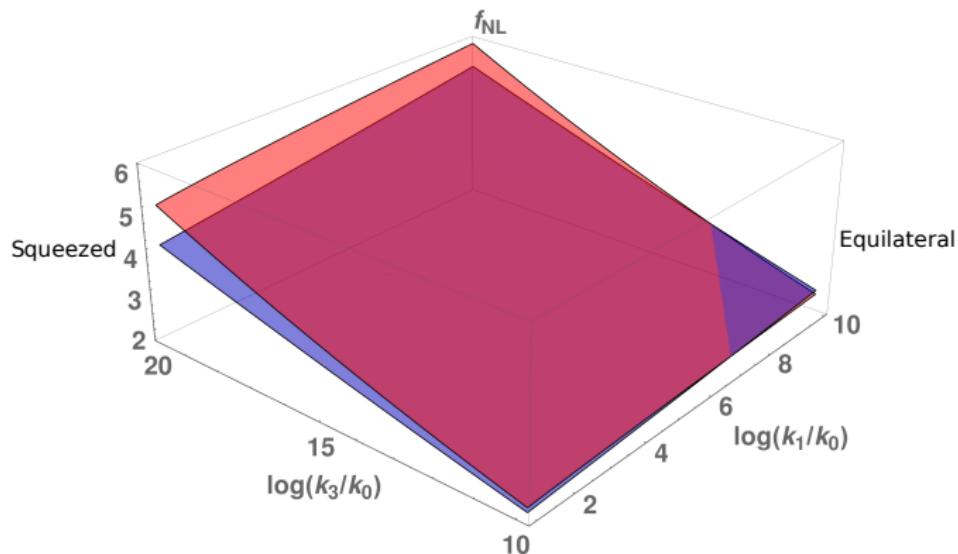
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Example: Interacting Curvaton

- ▶ Blue is us
- ▶ Other colours are previous work
- ▶ Reduced bispectrum

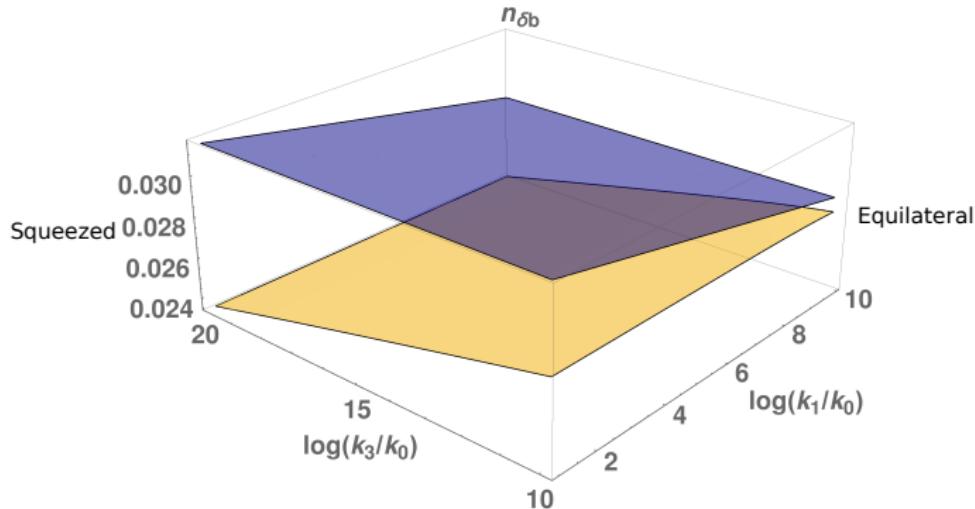
$$f_{NL}(k_1, k_2, k_3) \sim B_\zeta(k_1, k_2, k_3) / P_\zeta(k_1)P_\zeta(k_3)$$



Example: Interacting Curvaton

- ▶ Blue is us
- ▶ Other colours are previous work
- ▶ Spectral index of the halo bias

$$n_{\delta b} \sim \frac{d \log f_{\text{NL}}(k_1, k_2, k_3)}{d \log k_1}$$



Summary

- ▶ Done:
 - ▶ Calculated squeezed limit of bispectrum for multifield inflation
 - ▶ Our approach with Γ -evolution allows a larger squeezing than previous work
 - ▶ Observations: DES, Euclid, μ -distortions will probe large squeezing
 - ▶ Theoretical correction of about 20% for large squeezing

Summary

- ▶ Done:
 - ▶ Calculated squeezed limit of bispectrum for multifield inflation
 - ▶ Our approach with Γ -evolution allows a larger squeezing than previous work
 - ▶ Observations: DES, Euclid, μ -distortions will probe large squeezing
 - ▶ Theoretical correction of about 20% for large squeezing
- ▶ To do:
 - ▶ Higher point correlators, e.g. trispectrum
 - ▶ Multiple soft limits
 - ▶ Soft internal momenta
 - ▶ Understand from symmetries? [Assassi, Baumann, Green '12]
 - ▶ Can α term ever be large? Can't be in equilateral case.
[Lyth, Zaballa '05]

Thanks!

Special thanks to David Mulryne and Steve Thomas

Take-home message

- ▶ Standard δN calculates bispectrum in the equilateral limit
- ▶ We find the bispectrum in the squeezed limit

Multifield Case: Heuristic Proof of α result

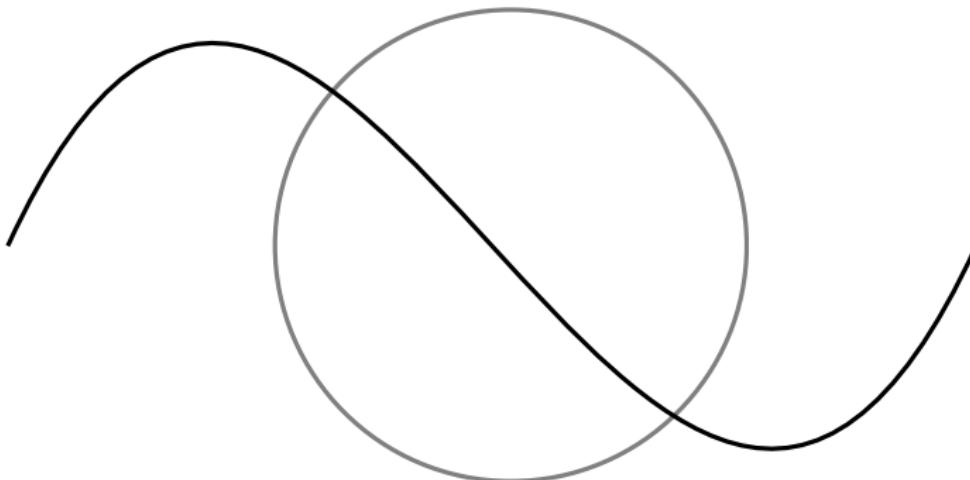
More detailed proof in the paper, in position space.

Let $q \sim k_1 \ll k_2, k_3$

$$\begin{aligned} \lim_{k_1 \ll k_{2,3}} \alpha_{ijk}^{(3)}(k_1, k_2, k_3) &\sim \lim_{k_1 \ll k_{2,3}} \langle \delta\phi_{i,\mathbf{k}_1}^{(3)} \phi_{j,\mathbf{k}_2}^{(3)} \phi_{k,\mathbf{k}_3}^{(3)} \rangle \\ &\sim \langle \delta\phi_{i,\mathbf{k}_1}^{(3)} \langle \phi_{j,\mathbf{k}_2}^{(3)} \phi_{k,\mathbf{k}_3}^{(3)} \rangle_{\delta\phi_{\mathbf{q}}} \rangle \\ &\sim \langle \delta\phi_{i,\mathbf{k}_1}^{(3)} (\langle \phi_{j,\mathbf{k}_2}^{(3)} \phi_{k,\mathbf{k}_3}^{(3)} \rangle + \delta\phi_{m,\mathbf{q}}^{(3)} \frac{\partial}{\partial \phi_m^{(3)}} \langle \phi_{j,\mathbf{k}_2}^{(3)} \phi_{k,\mathbf{k}_3}^{(3)} \rangle + \dots) \rangle \\ &\sim 0 + \langle \delta\phi_{i,\mathbf{k}_1}^{(3)} \delta\phi_{m,\mathbf{q}}^{(3)} \rangle \frac{\partial}{\partial \phi_m^{(3)}} \langle \phi_{j,\mathbf{k}_2}^{(3)} \phi_{k,\mathbf{k}_3}^{(3)} \rangle \\ &\sim \Sigma_{im}^{(3)}(k_1) \frac{\partial}{\partial \phi_m^{(3)}} \Sigma_{jk}^{(3)}(k_3) \end{aligned}$$

Other work this year

- ▶ ZK & Steve Thomas:
 - ▶ D-Brane Potentials in the Warped Resolved Conifold & Natural Inflation
 - ▶ JHEP 1502 (2015) 127. arxiv:1409.1221
- ▶ ZK, David J. Mulryne & Steve Thomas:
 - ▶ Generating Asymmetry with g_{NL}
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