

Cosmological Collider Physics

COSMO 2015

Warsaw

Juan Maldacena

IAS

Based on: N. Arkani-Hamed and JM, [arXiv: 1503.0804](https://arxiv.org/abs/1503.0804)

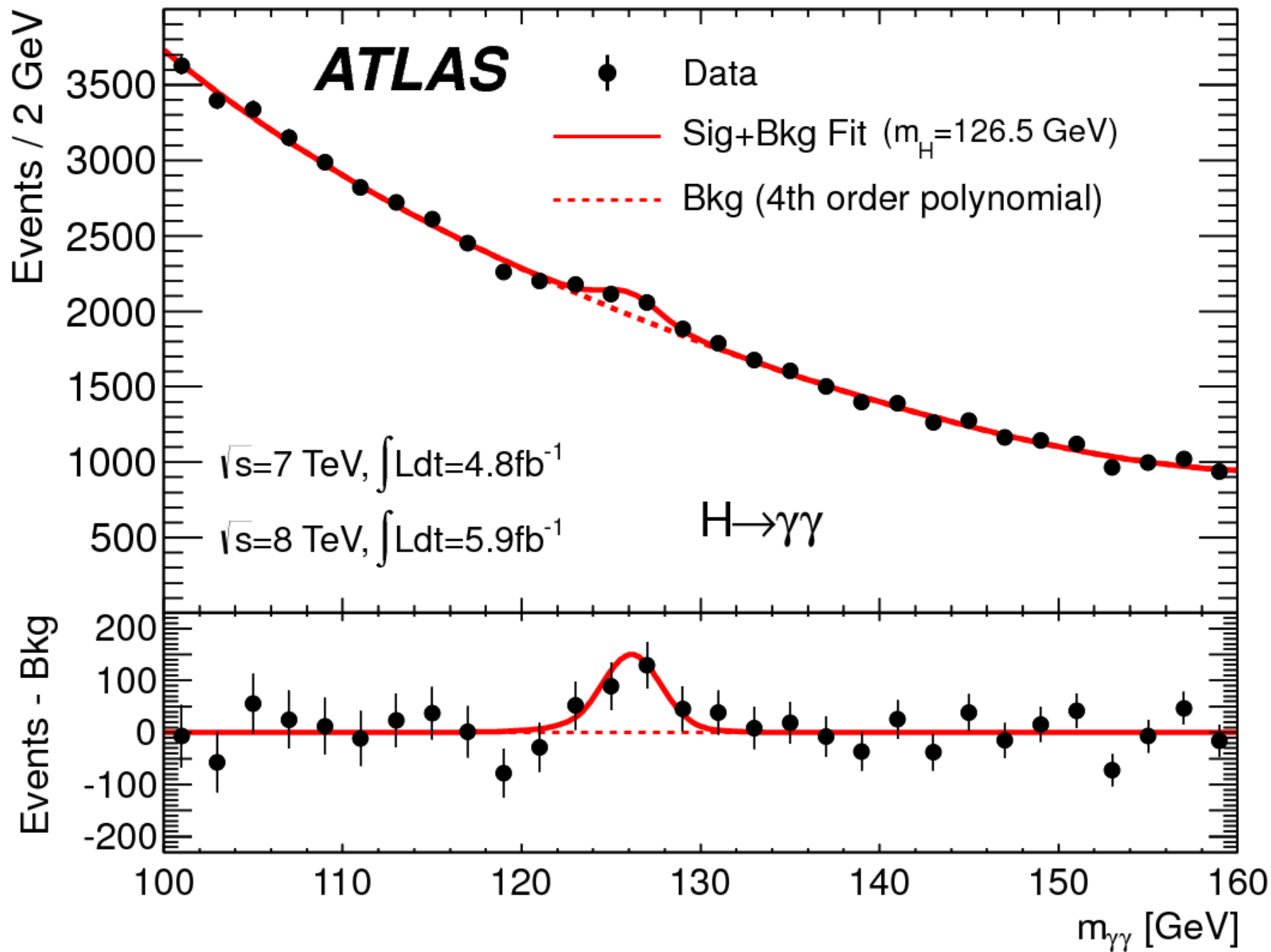
- According to inflationary theory cosmological perturbations have a quantum mechanical origin.
- They were created during inflation

$$T \sim \frac{H}{2\pi} , \quad 3M_{pl}^2 H^2 = V$$

- Relevant modes have energies of order H .
- Hubble scale could be as high as 10^{14} Gev !

- New particles and their interactions during inflation → leave a small imprint on the perturbations.
- We need to do the “collider physics”, i.e. go from the signatures to the basic interactions.
- How do we recognize new particles, measure their masses and spins ?

In flat space



In Cosmology

Non-Gaussianities

- Study non-gaussianities in the cosmological correlators.

- Can be produced by self interactions of the inflaton.

$$\int (\nabla\phi)^2 + \lambda(\nabla\phi)^4$$

- There are interesting patterns produced by new particles.

$$\int (\nabla\phi)^2 + (\nabla\sigma)^2 + m^2\sigma^2 + \lambda\sigma(\nabla\phi)^2$$

Analogy

QCD

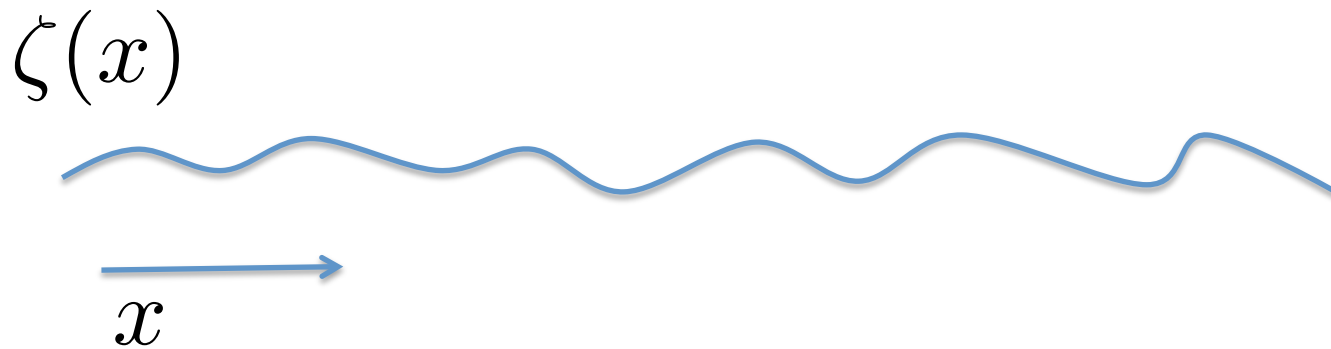
Cosmology

Hadrons	=	galaxies
Hadronization	=	Structure formation
Energy correlators	=	correlators of primordial density fluctuations.
Weak coupling at high energies	=	weak coupling during inflation
Approximate scale invariance	=	approximate scale invariance of wavefunction = approximate de-Sitter invariance.
OPE of energy correlators	=	squeezed limits of primordial correlators.
Time \rightarrow scale	=	time \rightarrow scale

Both are controlled by
(slightly broken) conformal symmetry

Basic Observable

Primordial Curvature Perturbations



$$\langle \zeta(x_1) \cdots \zeta(x_n) \rangle$$

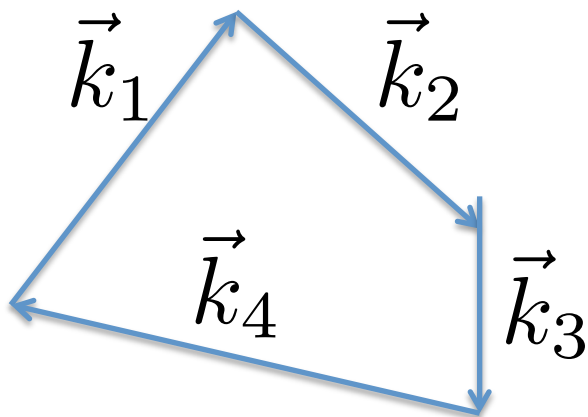
Kinematics

$$\langle \zeta(x_1) \zeta(x_2) \cdots \zeta(x_n) \rangle \rightarrow \langle \zeta(k_1) \cdots \zeta(k_n) \rangle$$

Fourier transform \rightarrow set of momenta

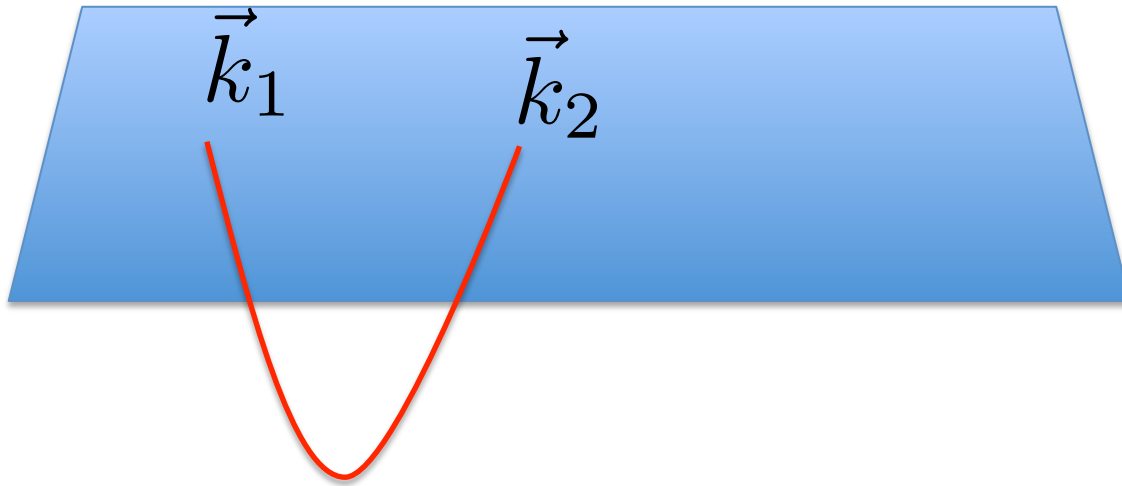
Statistical homogeneity of the universe \rightarrow Momentum conservation

This is similar to amplitudes. But no “energy conservation”.



First assume exact
de Sitter symmetry

Leading effect



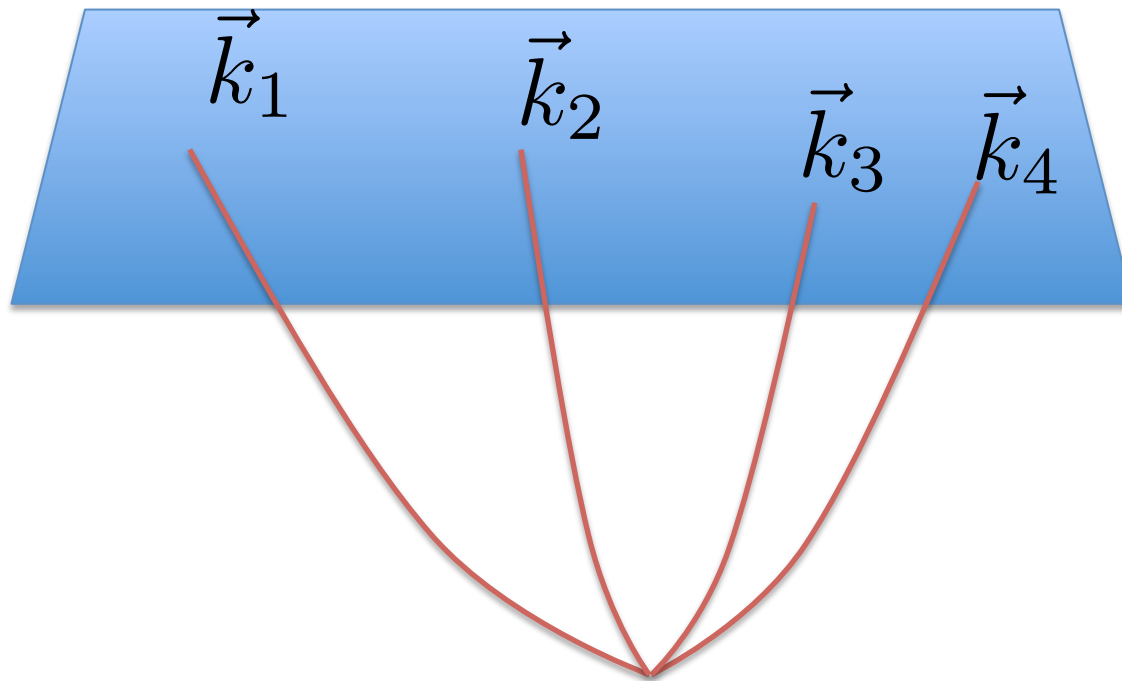
Two point function

$$\langle \phi_{\vec{k}_1} \phi_{\vec{k}_2} \rangle = \frac{H^2}{k_1^3} \delta^3(\vec{k}_1 + \vec{k}_2)$$

$$\langle \phi(0) \phi(x) \rangle \sim H^2 \log |x| + \dots$$

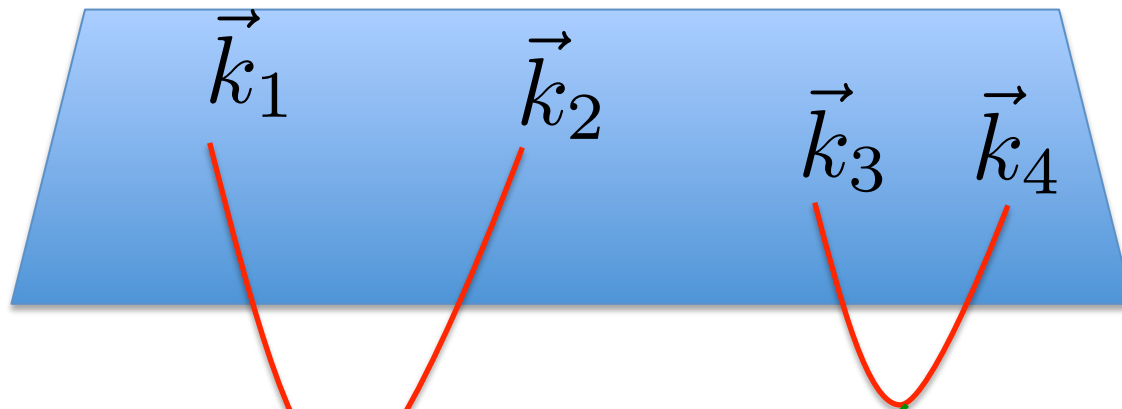
Four point function

- From self interactions



Four point function

- From a new massive particle existed during inflation, m of order H .



Density fluctuations \rightarrow
Inflaton \sim massless scalar
field in de-Sitter.

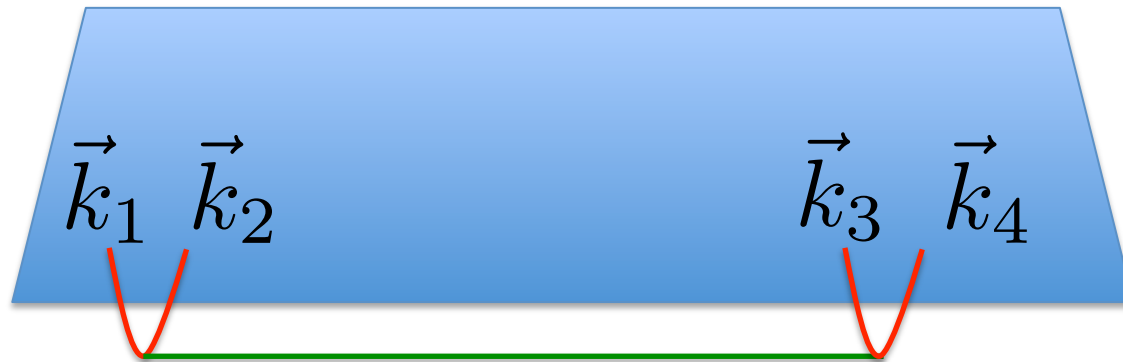
Massive particle

Another field,
Beyond the inflaton

$$\vec{k}_I = \vec{k}_1 + \vec{k}_2$$

Squeezed limit

$$k_I = |k_I| \ll |\vec{k}_i| = k_i, \quad i = 1, 2, 3, 4$$



Chen, Wang, (quasi single field),
 Noumi, Yamaguchi, Yokoyama,
 Assassi, Baumann, Green,
 Senatore, Silverstein,
 Zaldarriaga, Suyama, Yamaguchi,...

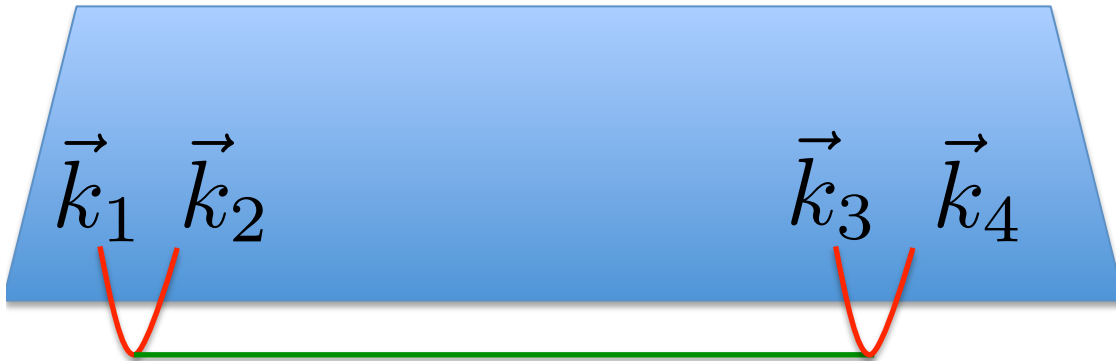
$$\frac{\langle 4pt \rangle}{\langle 2pt \rangle^2} \propto \frac{1}{k_I^3} \left[\left(\frac{k_I^2}{k_1 k_3} \right)^{\frac{3}{2} + i\mu} e^{i\delta} + c.c. \right]$$

$$\mu = \sqrt{m^2/H^2 - 9/4}$$

This non-trivial power of $k_I \rightarrow$ signature of a new physical particle.
 Not obtained from self interactions.

$$\frac{\langle 4pt \rangle}{\langle 2pt \rangle^2} \sim e^{-\pi\mu} \frac{1}{k_I^3} \left[\left(\frac{k_I^2}{k_1 k_3} \right)^{\frac{3}{2} + i\mu} e^{i\delta} + c.c. \right]$$

$$\mu = \sqrt{m^2/H^2 - 9/4}$$



We see clear oscillations a function of the log of the ratio of scales

Boltzman suppression $e^{-\pi\mu}$ vs. $1/(\mu)^k$

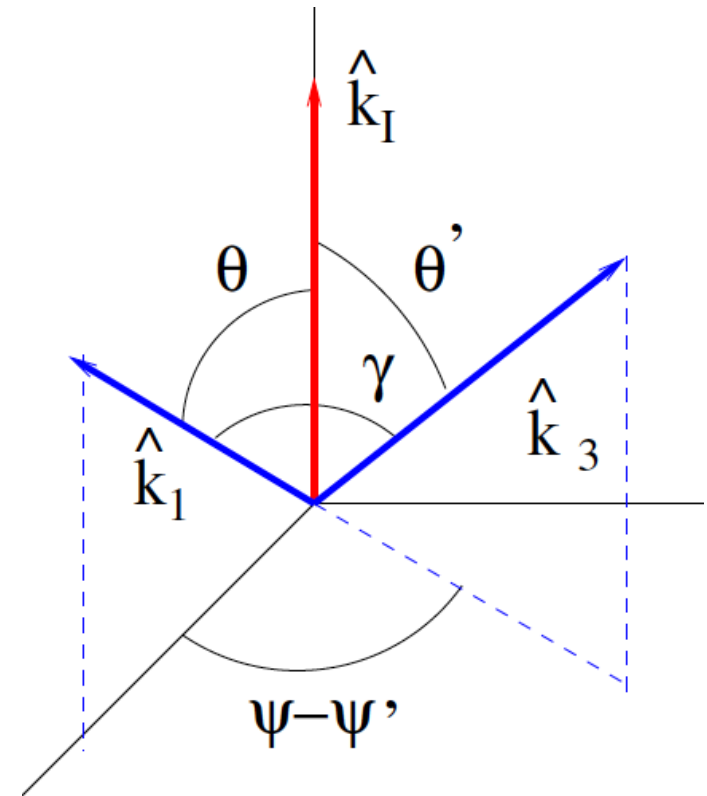
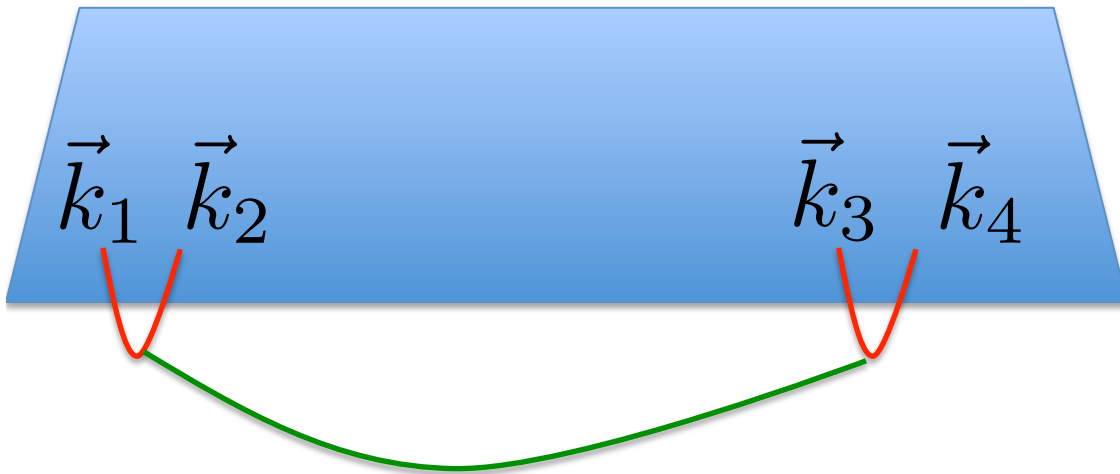
Interference effect: $\Psi_{\text{nopair}} + e^{-\pi\mu} \Psi_{\text{pair}}$

Phase is a function of the mass.

Interesting test of the quantum nature of fluctuations.

Spin

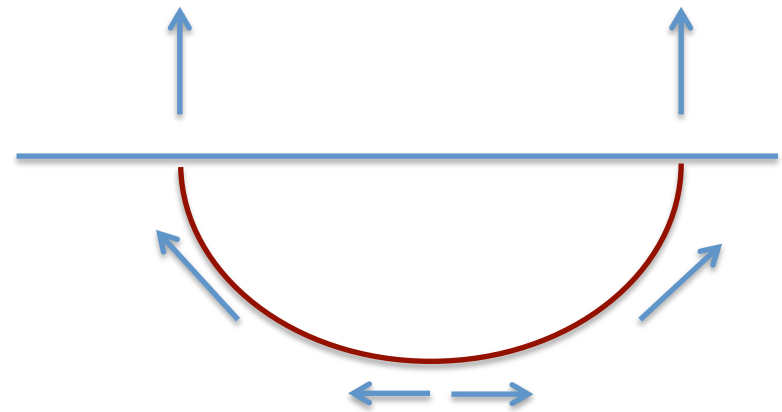
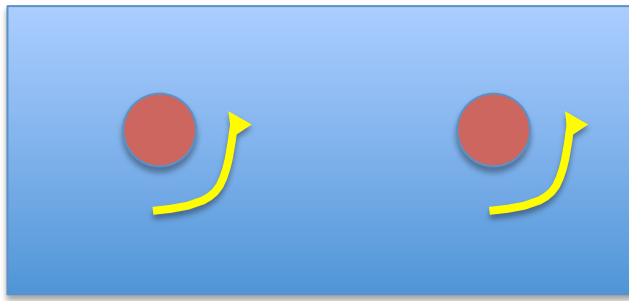
$$\langle 4pt \rangle \sim F(\gamma, \theta, \theta')$$



Further evidence of quantum mechanics ! \rightarrow
View it as a measurement of the correlated
spins of pair of produced particles.

There is a constraint on their masses.

Spin



$$\langle \epsilon_1^s \cdot O \epsilon_2^s \cdot O \rangle \sim \frac{[\epsilon_1 \cdot \epsilon_2 - 2(\epsilon_1 \cdot \hat{x})(\epsilon_2 \cdot \hat{x})]^s}{|x|^{2\Delta}}$$

Overall size estimate

$$\frac{\langle 4pt \rangle}{\langle 2pt \rangle^2} \sim \frac{H^2}{M_{pl}^2} \frac{e^{-\pi\mu}}{k_I^3} \left[\left(\frac{k_I^2}{k_1 k_3} \right)^{\frac{3}{2} + i\mu} e^{i\delta} + c.c. \right]$$

Overall size is small. $\lambda \sim 1/M_{pl}$

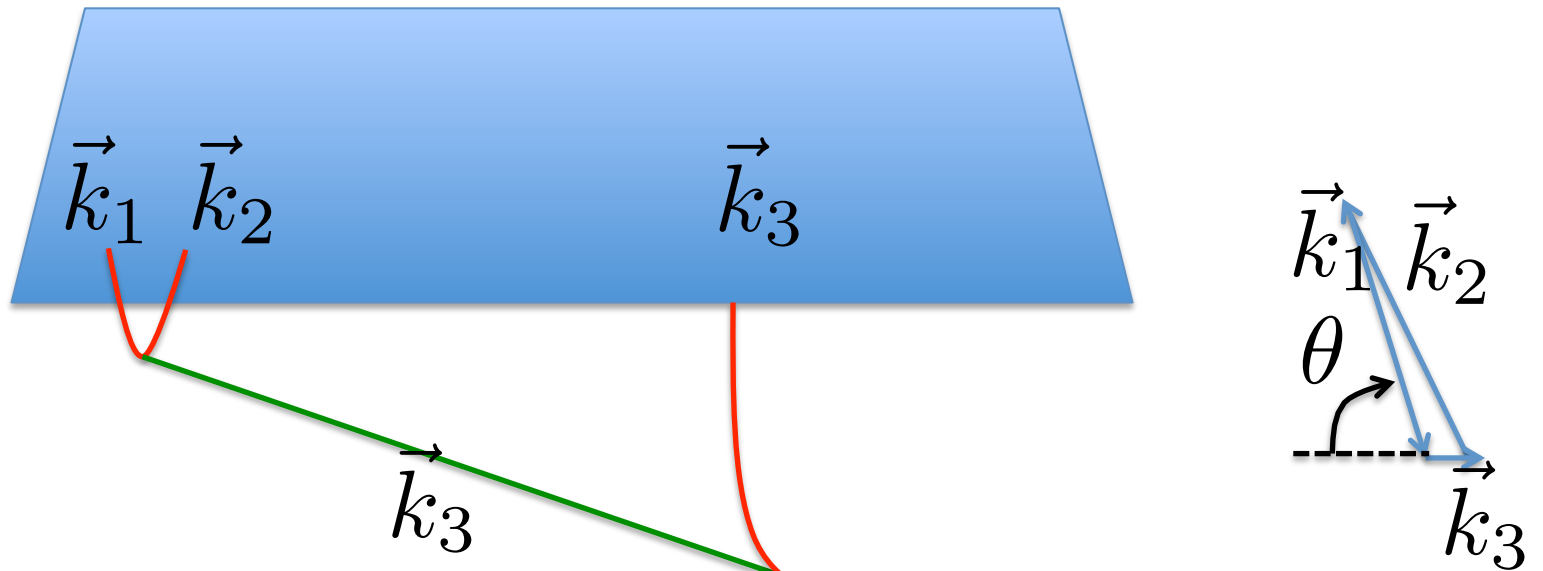
One factor of H/M from each interaction.

Can we find a bigger effect ?

Three point functions

- Consider instead the inflationary background.
- Now, we have a time dependent background

$$\phi(t)$$

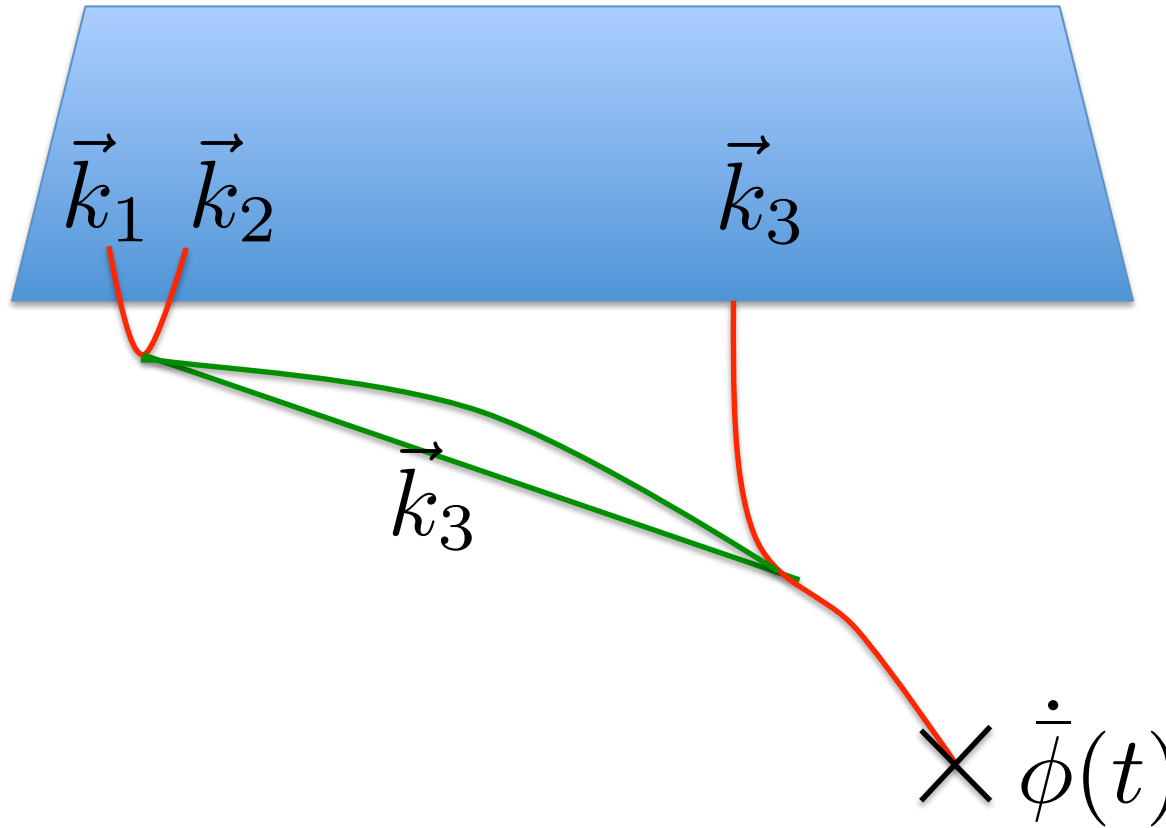


Only one very small coupling H/M_{pl}

$$\times \dot{\phi}(t) \vec{k}_4 = 0$$

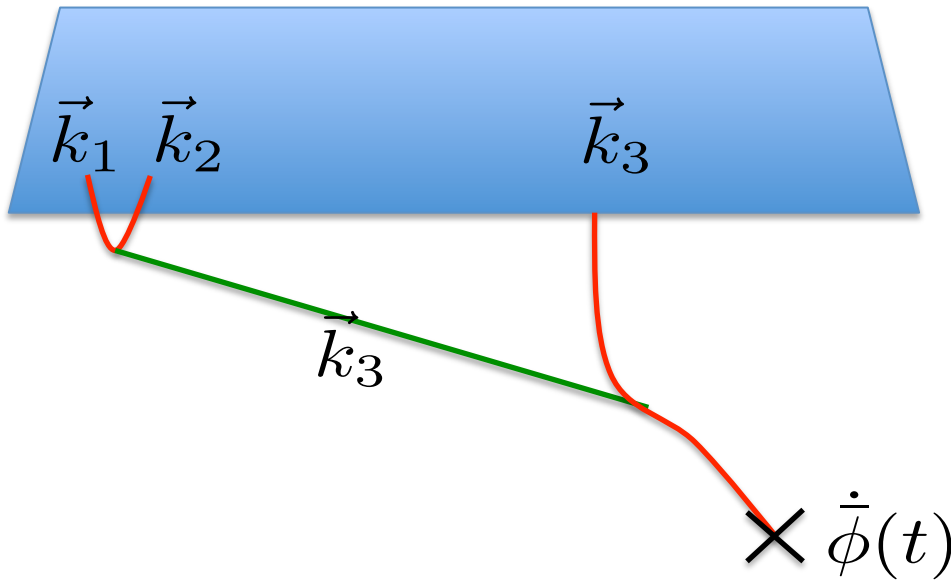
$$\langle 3pt \rangle \propto \overbrace{\frac{\dot{\phi}}{k_1^3 k_3^3}}^{\text{Usual}} e^{-\pi\mu} \left[\left(\frac{k_3}{k_1} \right)^{\frac{3}{2} + i\mu} e^{i\delta} + c.c. \right] P_s(\cos \theta)$$

Dilution



Loops \rightarrow
 give rise to a faster decay

$$\left(\frac{k_3}{k_1} \right)^{3+2i\mu}$$



Story: Particle is created by long wave mode k_3 . It then decays. We see interference between decay products and the original unperturbed state.

A striking evidence of quantum mechanics.

Phase of oscillation is calculable!.

Cosmological double slit experiment

Finding massive particles

- Collider \rightarrow peaks in the invariant mass distribution.
- Cosmology \rightarrow peaks in the Fourier transform of the cosmological correlator as a function of
$$\ell = \log(k_{short}/k_{long})$$
- Spin \rightarrow angular dependence.

How difficult is it to detect ?

Compare it with the standard 3pt function

- The standard 3 point function can be viewed as exchanging a graviton.

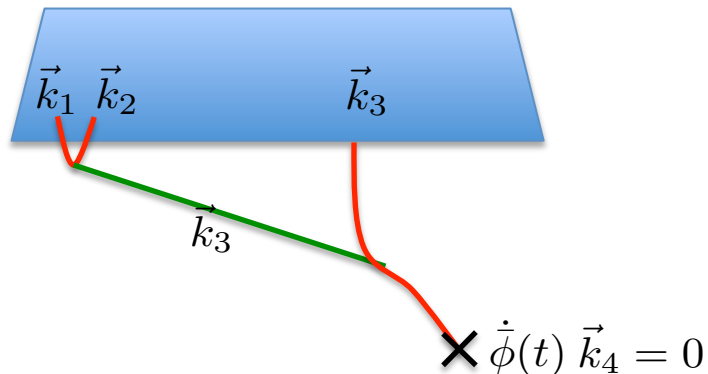
Kundu, Shukla, Trivedi

Planck:

$$|f_{NL}^{\text{experimental}}| \lesssim 5 ,$$

$$f_{NL}^{\text{standard}} \sim (n_s - 1)$$

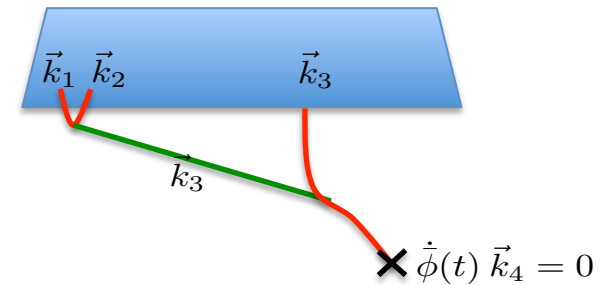
JM



How difficult is it to detect ?

- This one has extra factors of

$$(\lambda M_{pl})^2 e^{-\pi\mu} \left(\frac{k_3}{k_1} \right)^{3/2+i\mu}$$



- The last two suppress the signal. So the number of modes has to grow like the square of the above factor.
- The interactions could be larger than gravitational !

De Sitter isometries and conformal symmetry

$$ds^2 = \frac{-d\eta^2 + dx^2}{\eta^2}$$

$$\langle \phi(\eta_1, \vec{x}_1) \cdots \phi(\eta_n, \vec{x}_n) \rangle$$

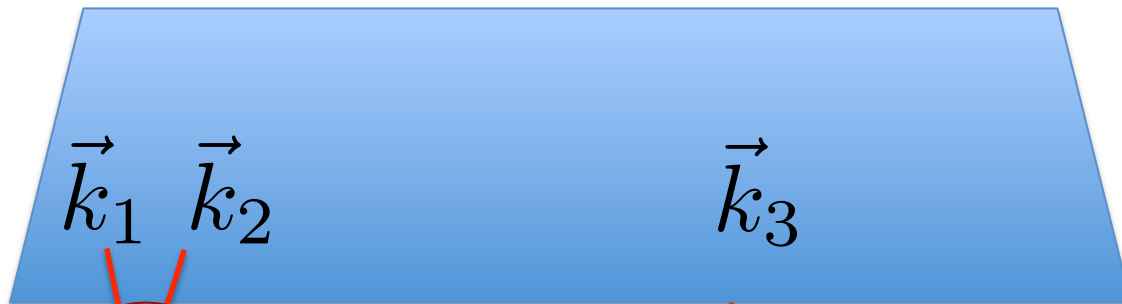
Invariant under de-Sitter isometries.

At late times, de-Sitter isometries act on x as conformal symmetries.

At late times we can often expand $\phi \sim \sum_i \eta^{\Delta_i} O_i(\vec{x})$

Strominger, Witten

3d operator of conformal dimension Δ_i



\vec{k}_1 \vec{k}_2

\vec{k}_3

\vec{k}_3

Decompose in terms of

$$\sum_i \eta^{\Delta_i} O_i(\vec{x})$$

$\times \dot{\phi}(t)$

Leads to

$$\sum_i \left(\frac{k_3}{k_1} \right)^{\Delta_i} C_i$$

$$\sum_i \left(\frac{k_3}{k_1} \right)^{\Delta_i} c_i$$

Powers that appear: Dimensions of 3d Operators \rightarrow
energies of quasinormal modes in the de-Sitter static patch.

Can be complex!

Sensitive to the spectrum of masses in the theory.

**Powers in the squeezed limit
=
Quasinormal mode spectrum**

The squeezed region of the correlator, $k_3 \ll k_1$, k_2 is not where the largest non-gaussian signal lies.

But it is the region containing direct information about the spectrum of the theory.

$$e^{-\pi\mu} \quad \text{vs.} \quad 1/(\mu)^k$$

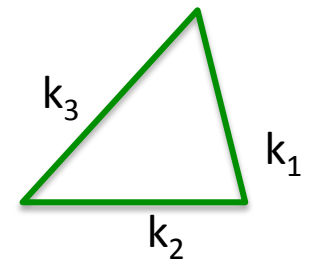
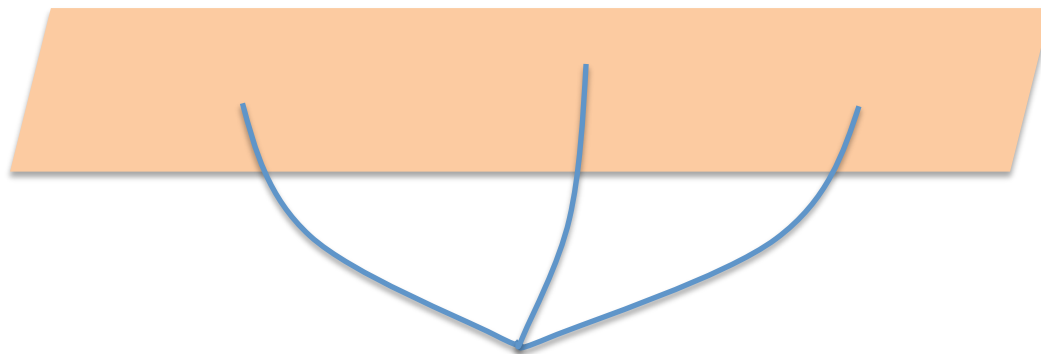
squeezed Leading effect

Very stringy inflation?

- Usual picture: Strings \rightarrow 10d \rightarrow KK theory \rightarrow inflation.
- Another possibility:
$$l_s \lesssim 1/H = R = \text{Hubble radius}$$
- Observations: higher spin massive particles!
- New structures in graviton three point functions.
- I do not know of a concrete stringy model...
- Could be evidence for string theory.

Higher spin ($S > 2$) weakly coupled
particles \rightarrow String theory

Graviton 3pt function

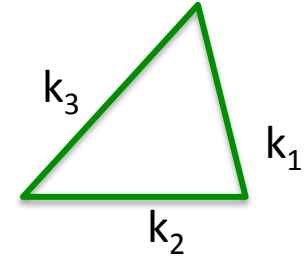
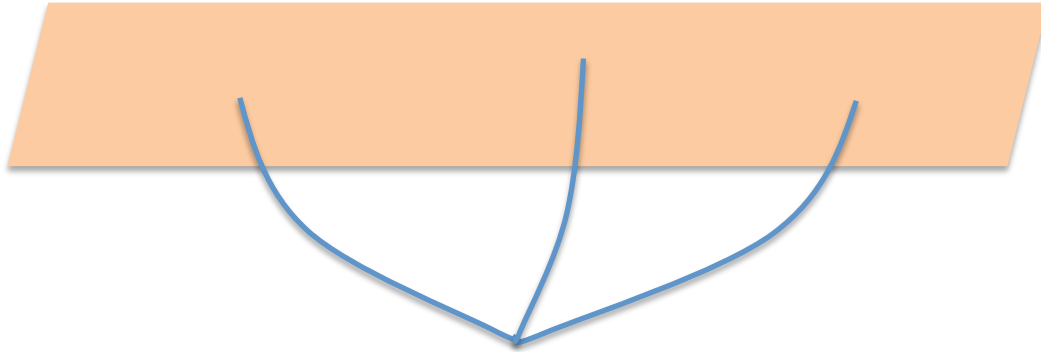


$$\frac{\langle hhh \rangle}{\langle hh \rangle^{3/2}} \sim \frac{H}{M_{pl}} \left[F_E(\text{shape}) + \alpha_4^2 H^4 F_2(\text{shape}) + o(\epsilon) \right]$$

JM, Pimentel

Overall small coupling

This is allowed by the approximate scale and conformal invariance of inflation



$$\frac{\langle hhh \rangle}{\langle hh \rangle^{3/2}} \sim \frac{H}{M_{pl}} \left[F_E(\text{shape}) + \alpha_4^2 H^4 F_2(\text{shape}) + o(\epsilon) \right]$$

JM, Pimentel

$$S = \frac{1}{G_N} \int R + \alpha_4^2 R^3 + \dots$$

If this is observed + causality of the de-Sitter theory \rightarrow massive higher spin states

This is only power suppressed in $l_s H$.

Camanho, Edelstein, J.M., Zhiboedov

Long string creation \rightarrow suppressed exponentially as $e^{-\frac{2}{(l_s H)^2}}$

Conclusions

- Non gaussianities in cosmological correlators have very interesting information.
- Squeezed limit directly probes the spectrum of the theory during inflation.
- Mass and spin information.
- Very interesting evidence of the quantum nature of the perturbations.
- Could be observable with futuristic experiments... (e.g. 21 cm tomography). (After seeing other non-gaussian signals.)