

Constraints on hybrid metric-Palatini models from background evolution

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Characteristics of hybrid metric-Palatini gravity

It is a type of $f(R)$ theory.

These can be approached from two directions:

- Metric variation.
- Connection independent from the metric.

In hybrid metric-Palatini theories

The Lagrangian is a function of the Ricci scalars of both the connection and the metric.

The action describing the hybrid metric-Palatini gravity is given by

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} [R + f(\mathcal{R})] + S_m, \quad (1)$$

where $\mathcal{R} = g^{\mu\nu} \mathcal{R}_{\mu\nu}$ is the Palatini curvature.

The modified Friedmann equations are

$$3H^2 = \frac{1}{1 + f_{\mathcal{R}}} \left[\kappa^2 \rho - 3H\dot{f}_{\mathcal{R}} - \frac{3\dot{f}_{\mathcal{R}}^2}{4f_{\mathcal{R}}} + \frac{\mathcal{R}F(\mathcal{R}) - f(\mathcal{R})}{2} \right], \quad (2)$$

and,

$$2\dot{H} = \frac{1}{1 + f_{\mathcal{R}}} \left[-\kappa^2 \rho + p + H\dot{f}_{\mathcal{R}} - \ddot{f}_{\mathcal{R}} + \frac{3\dot{f}_{\mathcal{R}}^2}{2f_{\mathcal{R}}} \right], \quad (3)$$

Identifying $\phi = F(\mathcal{R})$ the effective Klein-Gordon equation for the scalar field (by T. Harko, T. S. Koivisto et al, Phys. Rev. D 85, 084016, 2012)

$$\ddot{\phi} + 3H\dot{\phi} - \frac{\dot{\phi}^2}{2\phi} + \frac{\phi}{3} \left[2V - 1 + \phi \frac{dV}{d\phi} \right] = \frac{\kappa^2 \phi}{3} T, \quad (4)$$

Recalling $\phi \equiv F(\mathcal{R}) \equiv f_{\mathcal{R}}$ Eq.4 may be rewritten as

$$\ddot{\mathcal{R}} = -\frac{1}{f_{\mathcal{R}\mathcal{R}}} \left[\dot{\mathcal{R}}^2 \left[f_{\mathcal{R}\mathcal{R}\mathcal{R}} - \frac{f_{\mathcal{R}\mathcal{R}}^2}{2f_{\mathcal{R}}} \right] + 3H\dot{\mathcal{R}}f_{\mathcal{R}\mathcal{R}} + \frac{f_{\mathcal{R}}}{3} [\mathcal{R} [f_{\mathcal{R}} - 1] - 2f(\mathcal{R})] - \kappa^2 \frac{f_{\mathcal{R}}}{3} T \right] \quad (5)$$

The effective potential is

$$\frac{dV_{\text{eff}}}{d\mathcal{R}} = \frac{f_{\mathcal{R}}}{3f_{\mathcal{R}\mathcal{R}}} (\mathcal{R} [f_{\mathcal{R}} - 1] - 2f(\mathcal{R})), \quad (6)$$

meaning $V(\mathcal{R})$ will be given by the indefinite integral

$$V(\mathcal{R}) = \int^{\mathcal{R}} \frac{f_{\mathcal{R}}}{3f_{\mathcal{R}\mathcal{R}}} [\mathcal{R} (f_{\mathcal{R}} - 1) - 2f(\mathcal{R})] d\mathcal{R}. \quad (7)$$

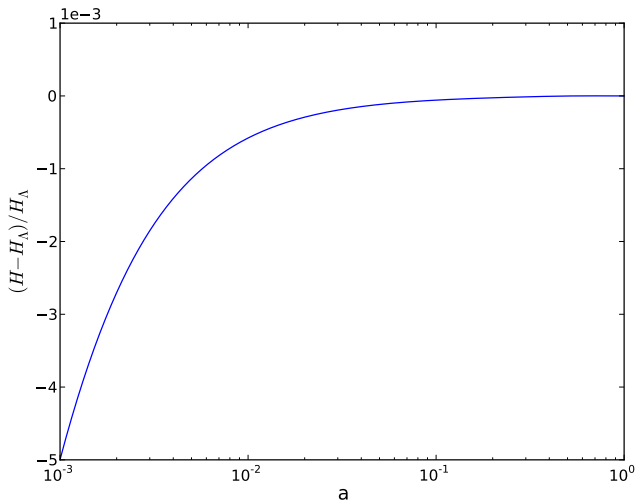
Exponential model

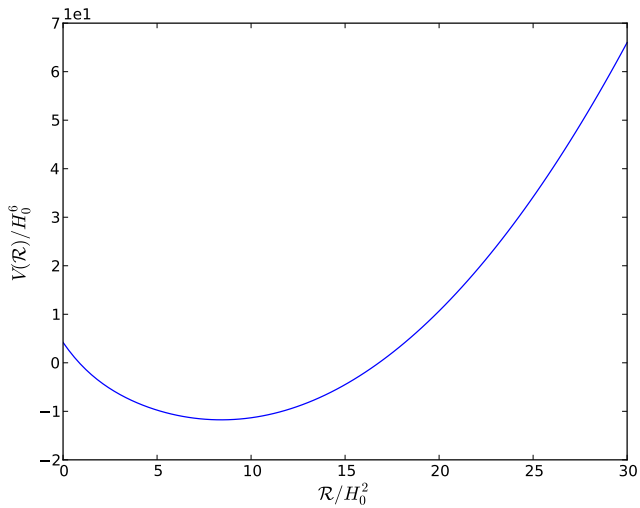
It is defined by the exponential function

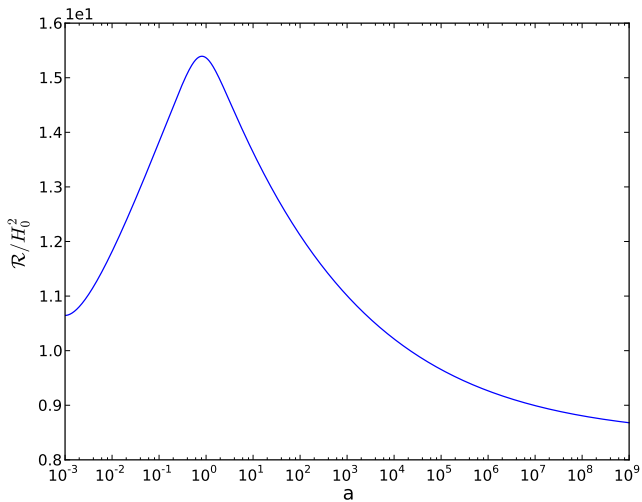
$$f(\mathcal{R}) = \Lambda_\star \left[1 + e^{-\mathcal{R}/\mathcal{R}_\star} \right], \quad (8)$$

with effective potential of the form

$$V(\mathcal{R}) = \frac{\mathcal{R}_\star}{3} \left[2\Lambda_\star \mathcal{R} + \frac{1}{2} \mathcal{R}^2 - 2\Lambda_\star \mathcal{R}_\star e^{-\mathcal{R}/\mathcal{R}_\star} - e^{-\mathcal{R}/\mathcal{R}_\star} (\Lambda_\star \mathcal{R} + \Lambda_\star \mathcal{R}_\star) \right]. \quad (9)$$







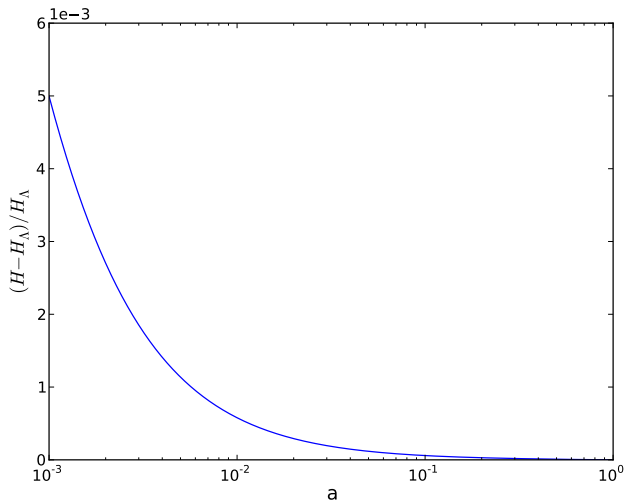
Quadratic model

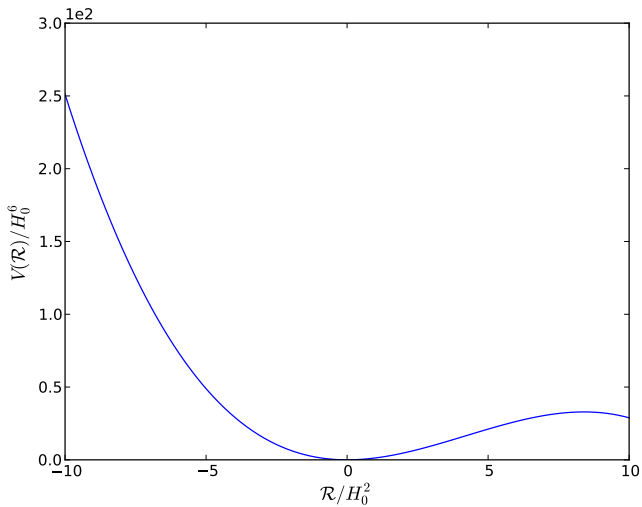
It is defined by the quadratic function

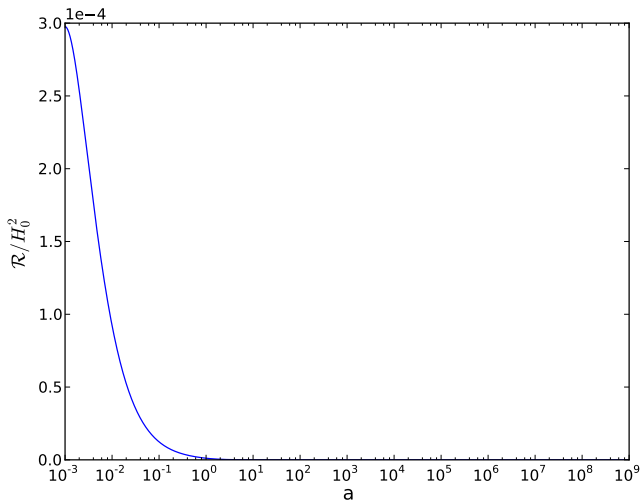
$$f(\mathcal{R}) = \Lambda_\star \left[1 + \mathcal{R}^2 / \mathcal{R}_\star^2 \right], \quad (10)$$

with the associated potential

$$V(\mathcal{R}) = -\frac{1}{3}\Lambda_\star \mathcal{R}^2 - \frac{1}{9}\mathcal{R}^3, \quad (11)$$







General priors are

- Ω_m : [0.01, 0.99].
- Ω_b : [0.001, 0.080].
- H_0 : [40.0, 100.0].

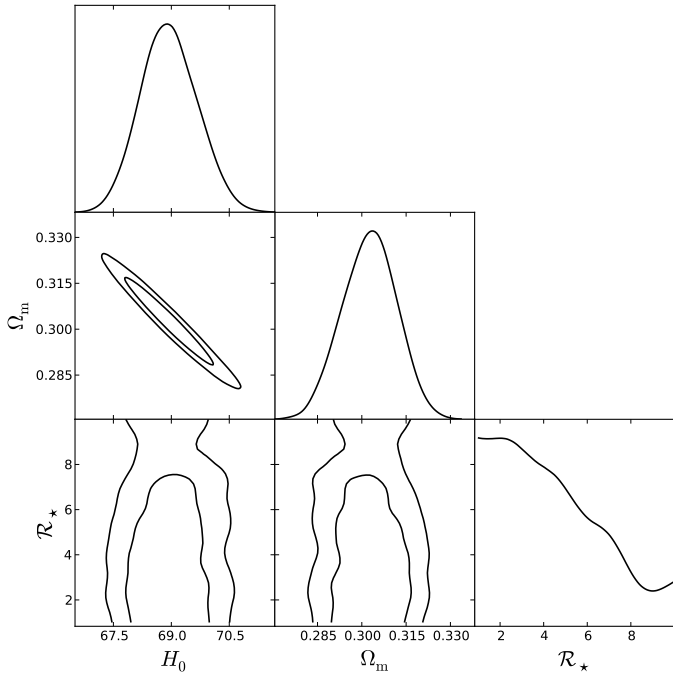
Exponential priors

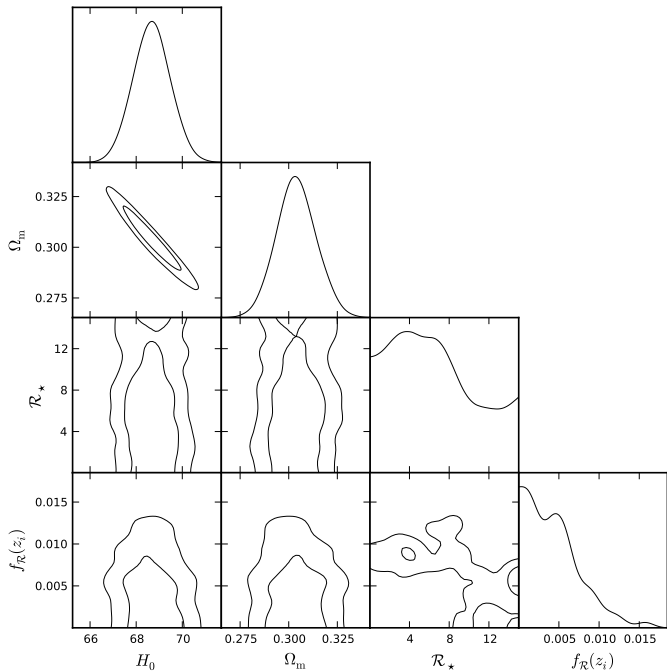
$f_{\mathcal{R}}$ at $z_i = 10^{-4}$, \mathcal{R}_* from [1.0, 10.0].

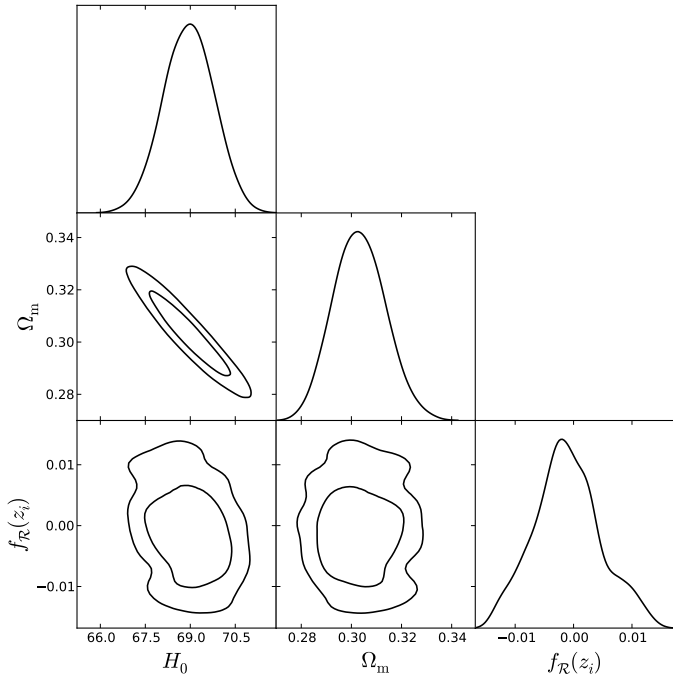
$f_{\mathcal{R}}(z_i)$ from $[1 \times 10^{-6}, 0.1]$ and \mathcal{R}_* from [0.01, 15.0].

Quadratic priors

\mathcal{R}_* is fix and $f_{\mathcal{R}}$ at z_i from $[-0.1, 0.1]$.







Conclusions

- Small deviation from Λ CDM in both scenarios.
- \mathcal{R} tends to the minimum at late times.
- Constraints on parameters after imposing observations available.
- Interesting enough theory to add perturbative analysis.

Thank you for listening.