

Characteristic signals of right-handed currents at future colliders

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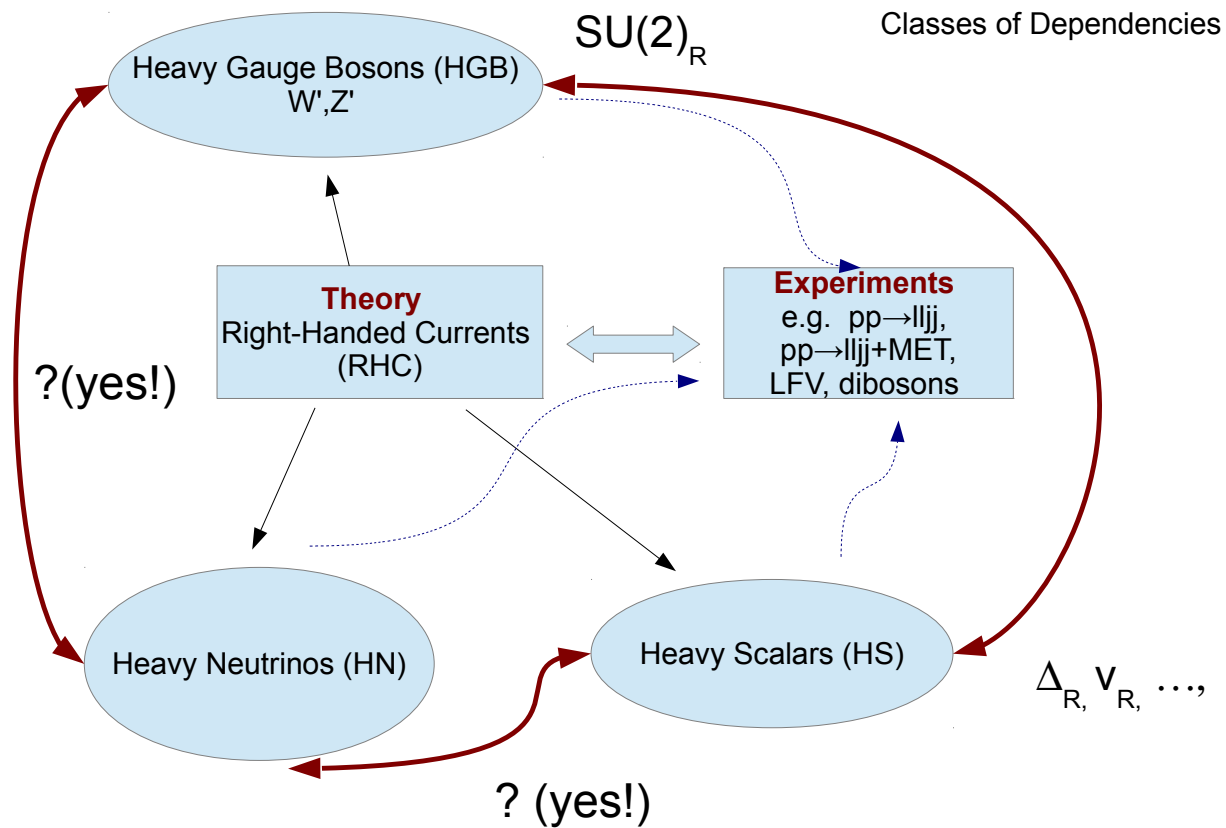
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in collaboration with:

G. Bambhaniya, J. Chakraborty, J. Gluza, M. Kordiaczyńska, T. Srivastava, R. Szafron

FCC Week 2016, Roma, 11-15 April 2016

RHC includes plenty of connected issues



Studies including FCC-hh:

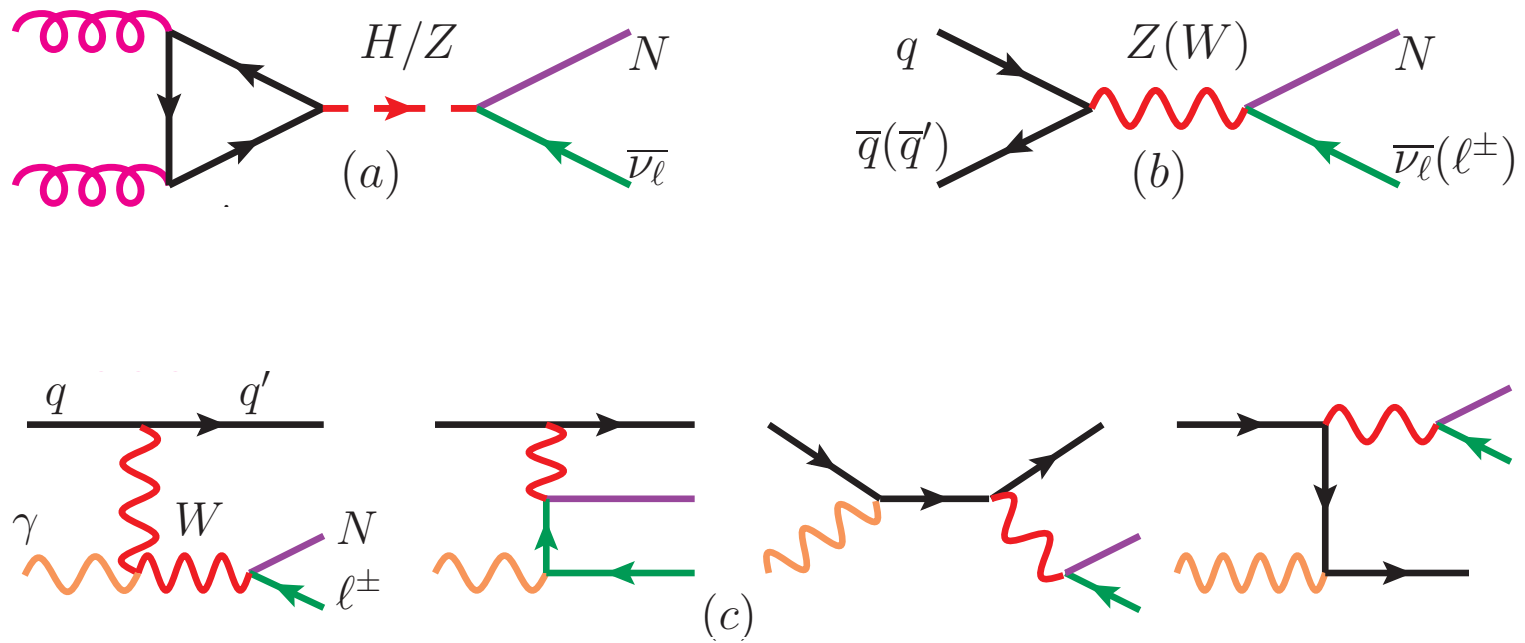
- ❖ G. Bambhaniya, J. Chakraborty, J. Gluza, T. Jeliński and R. Szafron, “Search for doubly charged Higgs bosons through VBF at the LHC and beyond ,” [arXiv:1504.03999](#), PRD 92 (2015) 1, 015016
 - ❖ G. Bambhaniya, J. Chakraborty, J. Gluza, T. Jeliński and M. Kordiaczyńska, “Lowest limits on the doubly charged Higgs boson masses in the minimal left-right symmetric model,” [arXiv:1408.0774](#), PRD 90 (2014) 9, 095003
 - ❖ J .Gluza, T. Jeliński, “Heavy neutrinos and the $pp \rightarrow lljj$ CMS data,” [arXiv:1504.05568](#), PLB 748 (2015) 125
 - ❖ J. Gluza, T. Jeliński and R. Szafron, “Lepton Number Violation and ‘Diracness’ of massive neutrinos composed of Majorana states,” [arXiv:1604.01388](#)
 - ❖ FCC-hh (pre-CDR) Report: ” Probing BSM Higgs Sectors at the FCC-hh” , ” Beyond the Standard Model physics at 100 TeV”
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Heavy neutrinos interactions

$$\begin{pmatrix} \nu_L \\ N_R \end{pmatrix} = \begin{pmatrix} U_{3 \times 3} & V_{3 \times n} \\ X_{n \times 3} & K_R^\dagger_{n \times n} \end{pmatrix} \begin{pmatrix} \nu \\ N \end{pmatrix}$$

$$\begin{aligned} \mathcal{L} = & -\frac{g_L}{\sqrt{2}} W_{1\mu}^+ \bar{\nu} U^\dagger \gamma^\mu P_L \ell^- - \frac{g_L}{\sqrt{2}} W_{1\mu}^+ \bar{N} V^\dagger \gamma^\mu P_L \ell^- \\ & -\frac{g_R}{\sqrt{2}} W_{2\mu}^+ \bar{\nu} X^\dagger \gamma^\mu P_R \ell^- - \frac{g_R}{\sqrt{2}} W_{2\mu}^+ \bar{N} K_R \gamma^\mu P_R \ell^- \\ & + \dots \end{aligned}$$

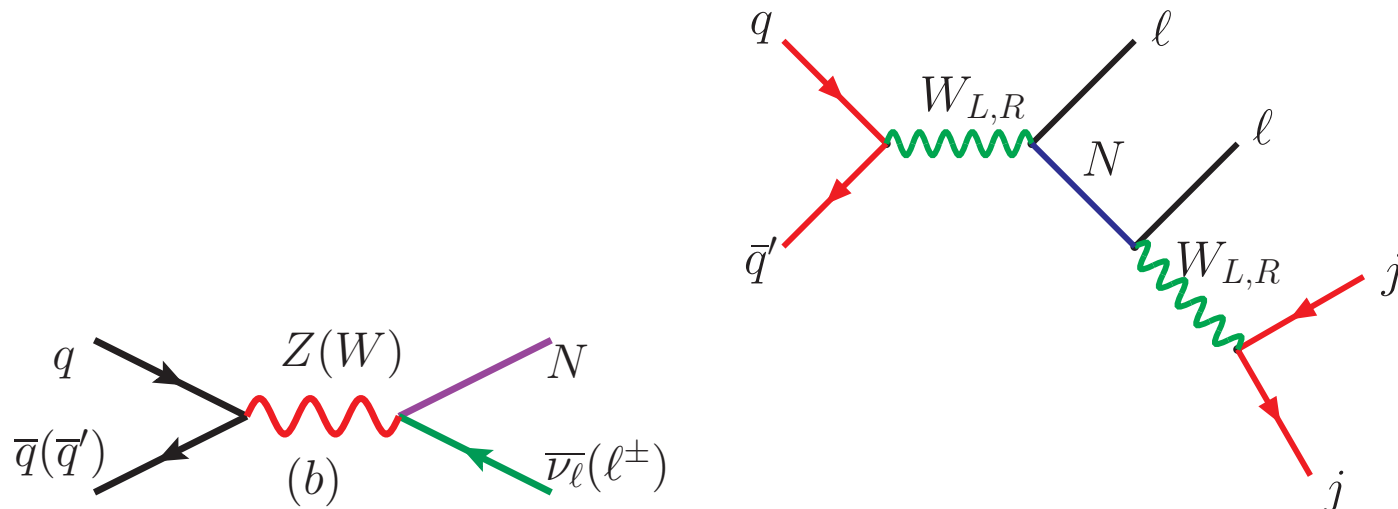
Heavy neutrinos production



[Degrande et al., arXiv:1602.06957]

ongoing studies: estimation of the SM background and analysis of possible decay channels

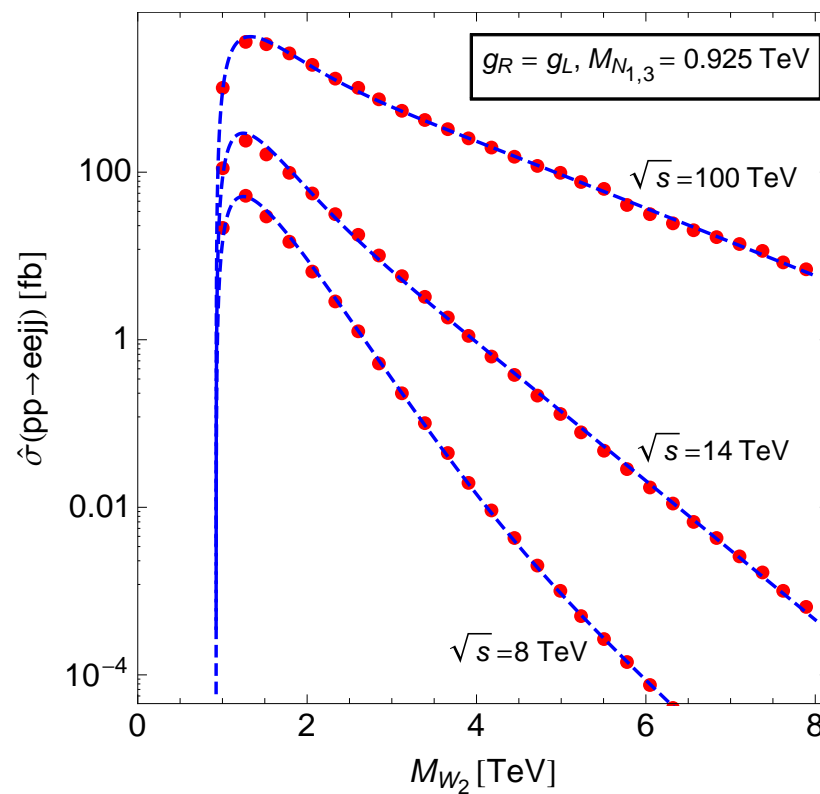
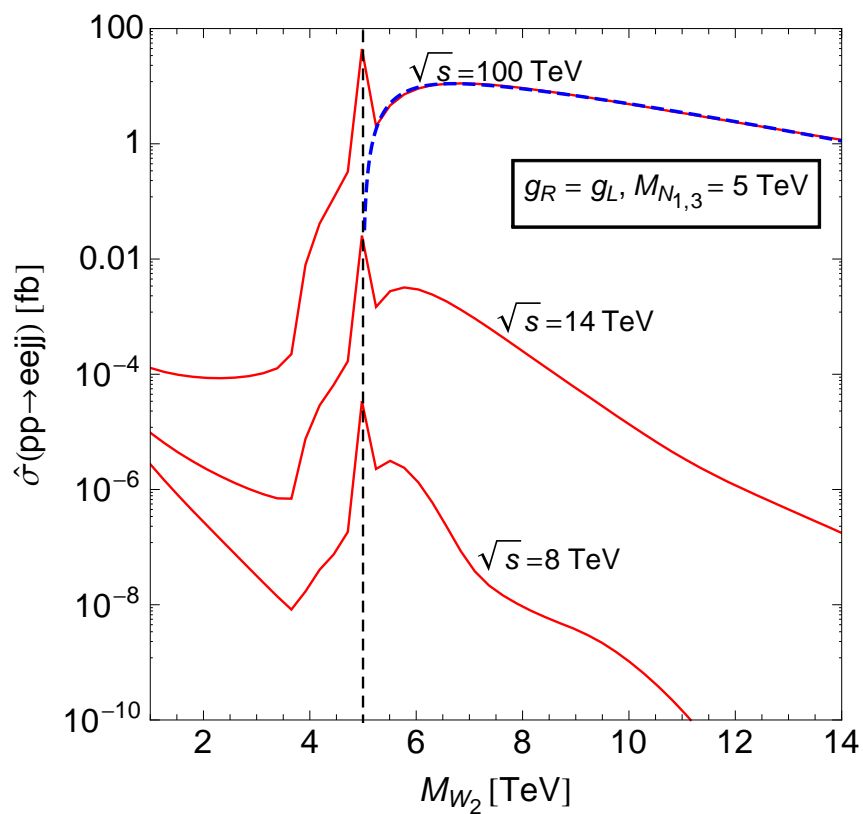
Drell-Yan production & $N \rightarrow lj$



- ❖ golden channel for Majorana heavy neutrinos \rightarrow SS and OS content of the dilepton signal, LFV
 - ❖ no \cancel{E}_T , relatively easy to separate from the SM background
 - ❖ heavy neutrino masses and mixings are crucial for interference effects \rightarrow more careful treatment of $pp \rightarrow eejj$ is needed; rough analyses can lead to overconstraining parameter space.
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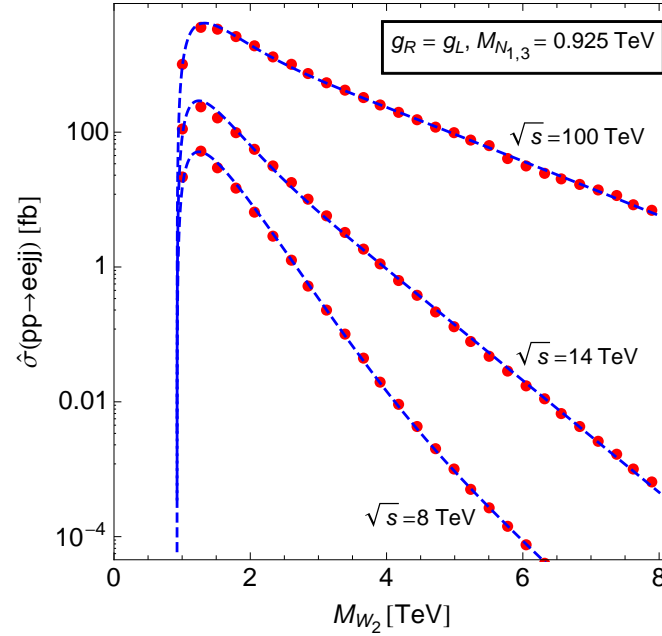
Total cross section $\sigma(pp \rightarrow eejj)$

MADGRAPH5 2.3.3 & NN23LO1 PDFs



NWA valid only for $M_{N_{1,3}} < M_{W_2} < \sqrt{s}$

Fitting dependence on M_{W_2}



$$\text{Naive fit: } \sigma(pp \rightarrow eejj) = \frac{F_W(x_1)}{1 + \frac{1}{18} \sum_b F_W(x_b)} a(e^{-b\mu} + ce^{-d\mu}), \quad \mu = M_{W_2}/1 \text{ TeV}$$

$$(8 \text{ TeV}) \quad a = 0.18 \times 10^5 \text{ fb}, \quad b = 3.62, \quad c = 0.002, \quad \text{and} \quad d = 2.17,$$

$$(14 \text{ TeV}) \quad a = 1.32 \times 10^5 \text{ fb}, \quad b = 3.97, \quad c = 0.016, \quad \text{and} \quad d = 1.92,$$

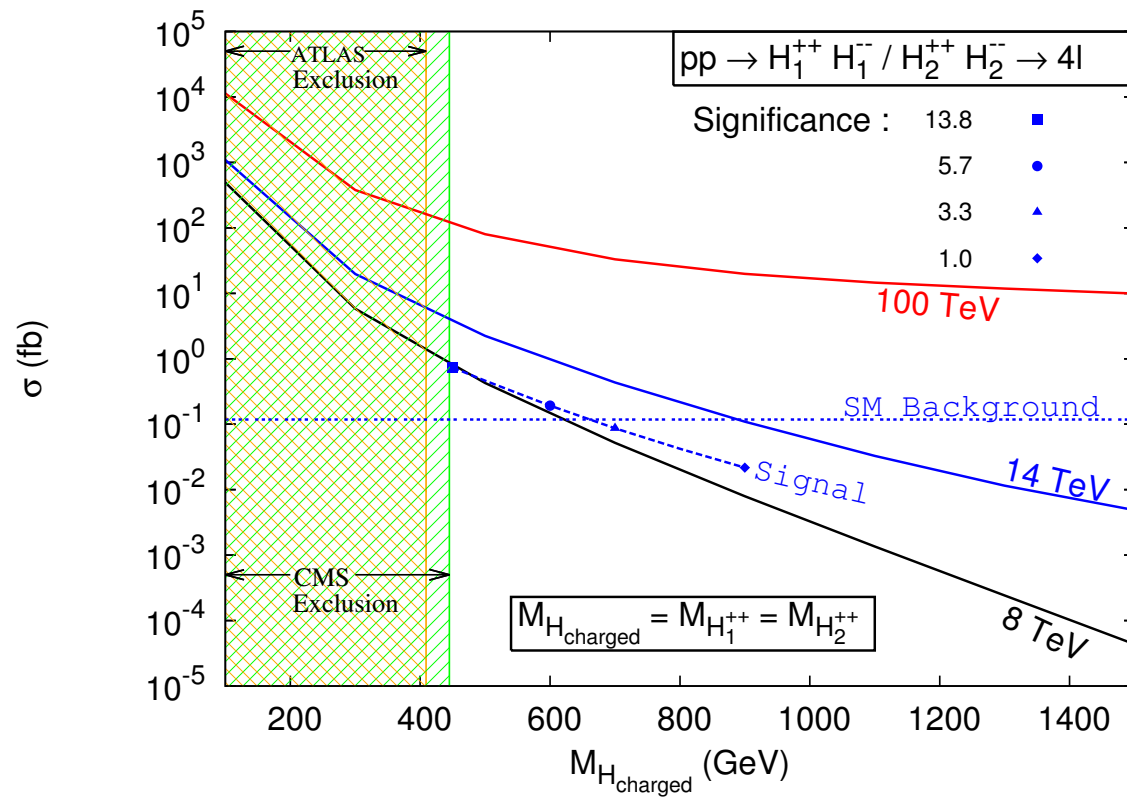
$$(100 \text{ TeV}) \quad a = 5.40 \times 10^5 \text{ fb}, \quad b = 3.04, \quad c = 0.020, \quad \text{and} \quad d = 0.94.$$

Primary production	Secondary production	Signal
I. $H_1^+ H_1^-$	$l^+ l^- \nu_L \nu_L$	$l^+ l^- \oplus MET$
-	$l^+ l^- \nu_R \nu_R$	depends on ν_R decay modes
-	$l^+ l^- \nu_L \nu_R$	depends on ν_R decay modes
II. $H_2^+ H_2^-$	$l^+ l^- \nu_L \nu_L$	$l^+ l^- \oplus MET$
-	$l^+ l^- \nu_R \nu_R$	depends on ν_R decay modes
-	$l^+ l^- \nu_L \nu_R$	depends on ν_R decay modes
III. $H_1^{++} H_1^{--}$	-	$l^+ l^+ l^- l^-$
-	$H_1^+ H_1^+ H_1^- H_1^-$	See I
-	$H_1^\pm H_1^\pm H_2^\mp H_2^\mp$	See I & II
-	$H_2^+ H_2^+ H_2^- H_2^-$	See II
-	$W_i^+ W_i^+ W_j^- W_j^-$	depends on W 's decay modes
IV. $H_2^{++} H_2^{--}$	-	$l^+ l^+ l^- l^-$
-	$H_2^+ H_2^+ H_2^- H_2^-$	See II
-	$H_1^\pm H_1^\pm H_2^\mp H_2^\mp$	See I & II
-	$H_1^+ H_1^+ H_1^- H_1^-$	See I
-	$W_i^+ W_i^+ W_j^- W_j^-$	depends on W 's decay modes
V. $H_1^{\pm\pm} H_1^\mp$	-	$l^\pm l^\pm l^\mp \nu_L$
VI. $H_2^{\pm\pm} H_2^\mp$	-	$l^\pm l^\pm l^\mp \nu_L$
VII. $H_1^\pm Z_i, H_1^\pm W_i$	-	See I & Z_i, W_i decay modes
VIII. $H_2^\pm Z_i, H_1^\pm W_i$	-	See II & Z_i, W_i decay modes
IX. $H_1^\pm A$	-	See I
X. $H_2^\pm A$	-	See II

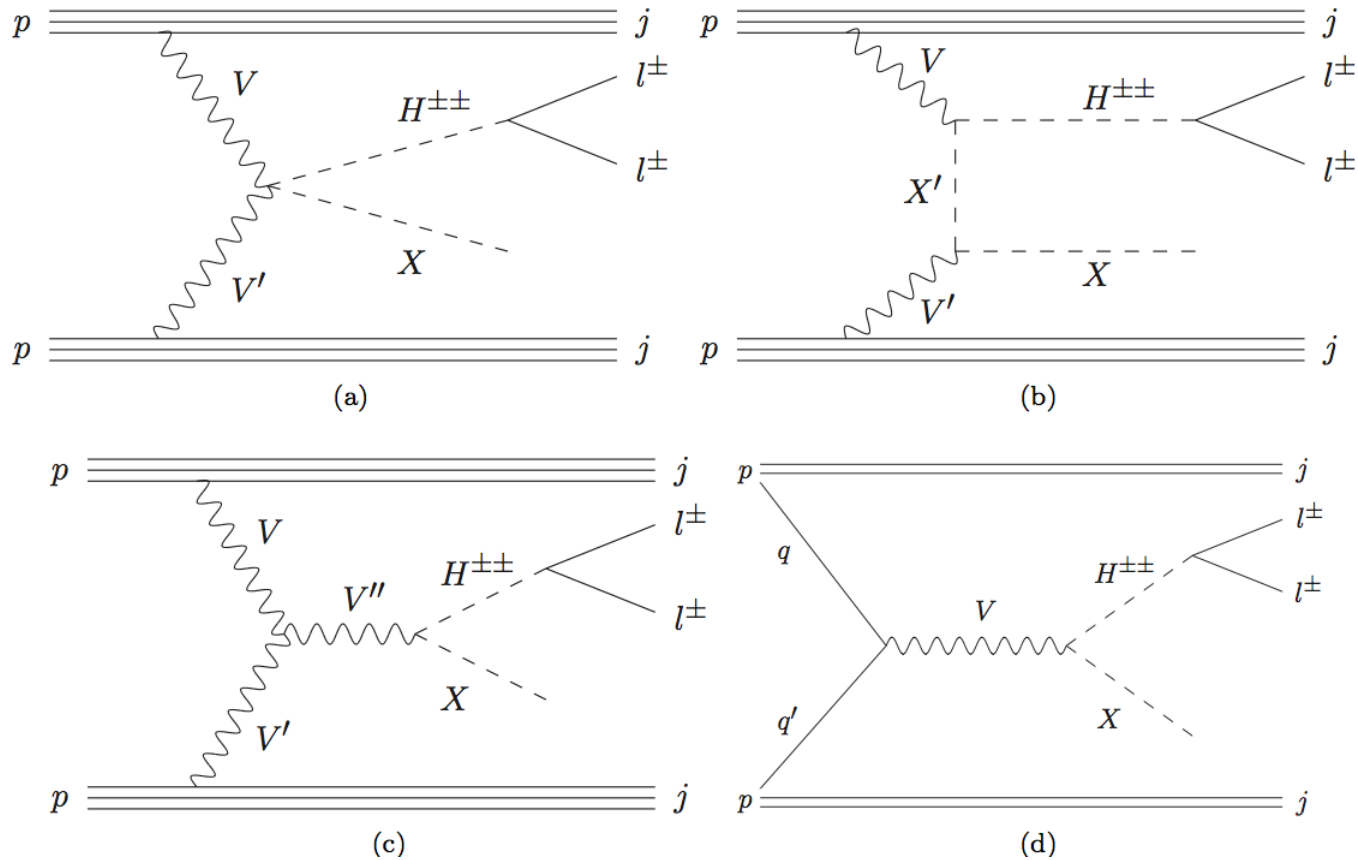
Doubly charged Higgses production (Drell-Yan)

$$m_{H_{1,2}^{\pm\pm}} = 600 \text{ GeV} :$$

$$\sigma(pp \rightarrow H_{1,2}^{++} H_{1,2}^{--} \rightarrow l_i^+ l_i^+ l_j^- l_j^-) = 0.144(0.9498) \text{ fb for } \sqrt{s} = 8(14) \text{ TeV}.$$



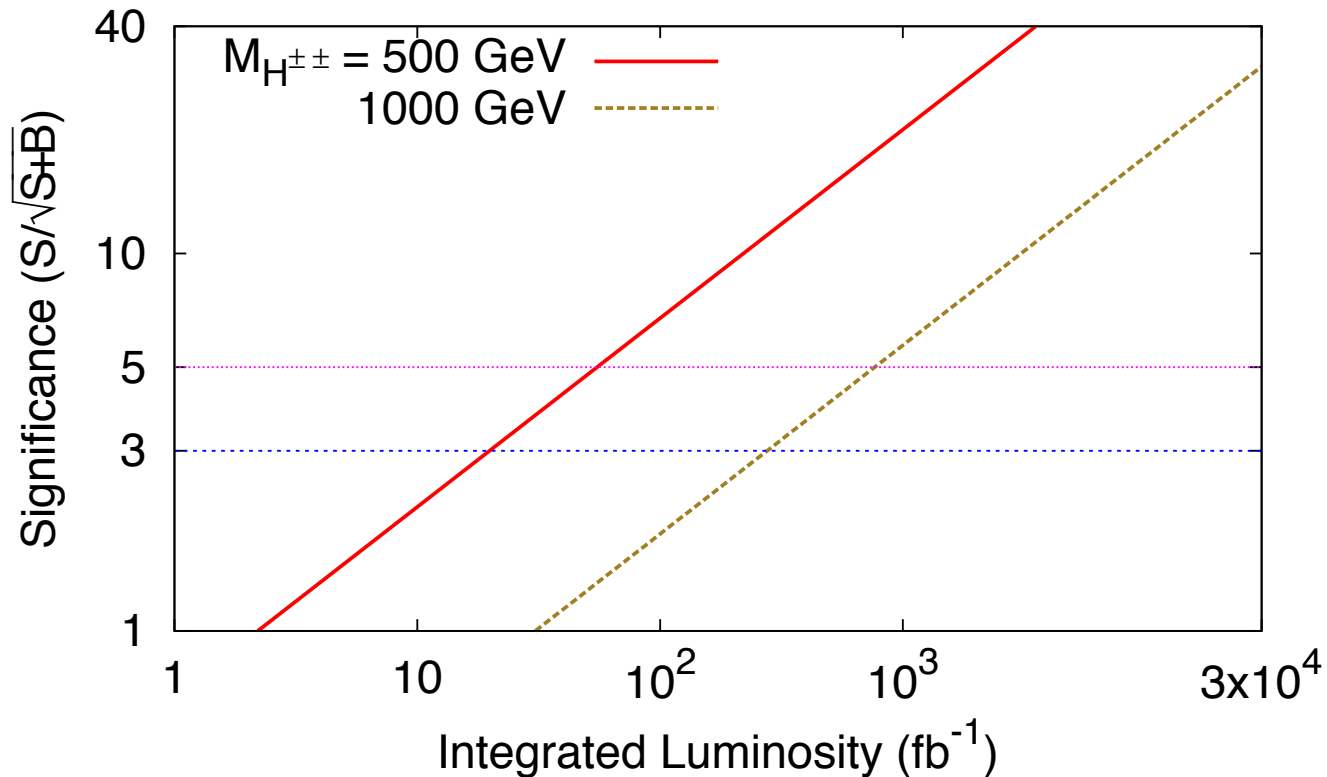
Doubly charged Higgs bosons production (Vector Boson Fusion with 2 jets)



Need for dedicated search channels for tagged forward jets.

Doubly charged Higgs bosons production (Vector Boson Fusion with 2 jets)

$$\sqrt{s} = 100 \text{ TeV}$$



FCC-hh dedicated DELPHES card used which includes the basic isolation and selection criteria.

Summary

- ❖ Discovery of W_2^\pm , $H^{\pm\pm}$ or N_i would be something incredibly new and would define new directions in physics (e.g. issue of supersymmetry)
 - ❖ FCC-hh opens up a very wide range of W_2^\pm , $H^{\pm\pm}$ or N_i masses which can be explored
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