

SOTTO L'ALTO PATRONATO DEL PRESIDENTE DELLA REPUBBLICA
UNDER THE HIGH PATRONAGE OF THE PRESIDENT OF THE ITALIAN REPUBLIC

Prospects for the measurement of the Higgs Potential @ FCC

Seung J. Lee

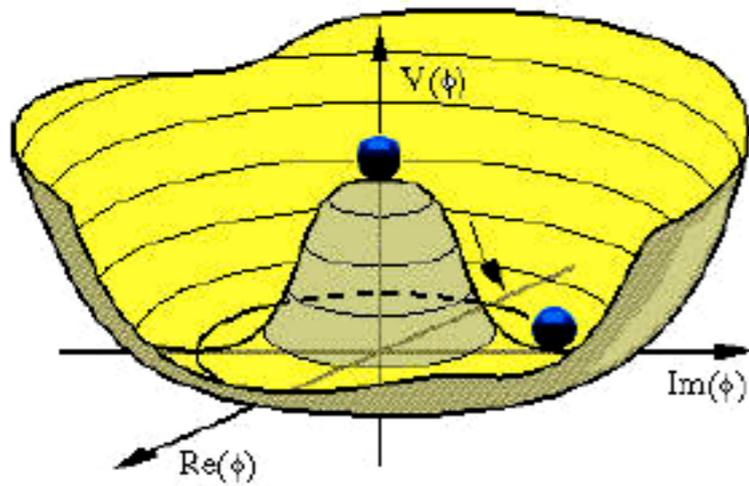


Based on collaboration with Benjamin Fuks, Jeong Han Kim (arXiv:1510.07697)

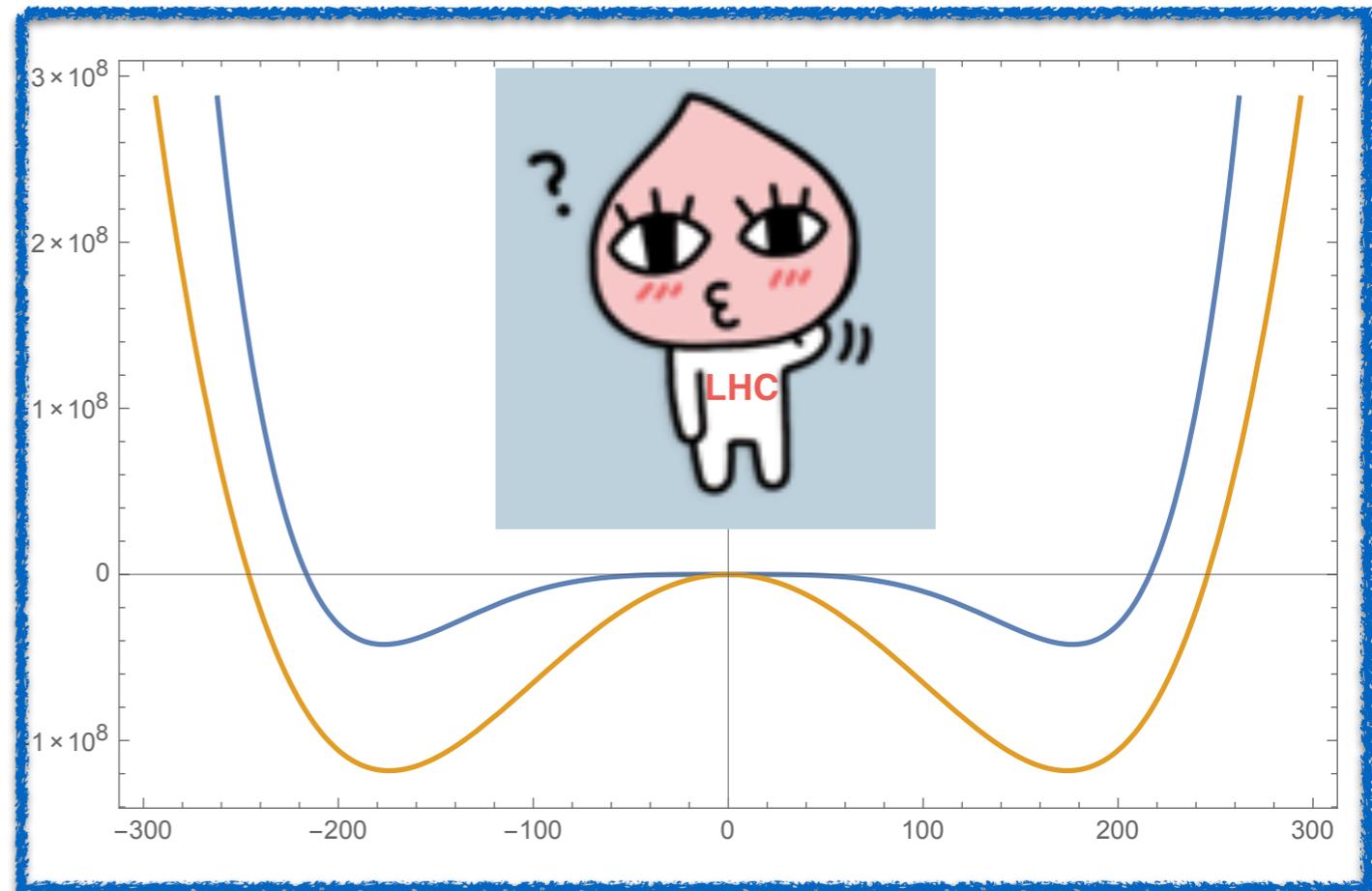
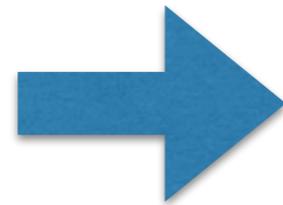
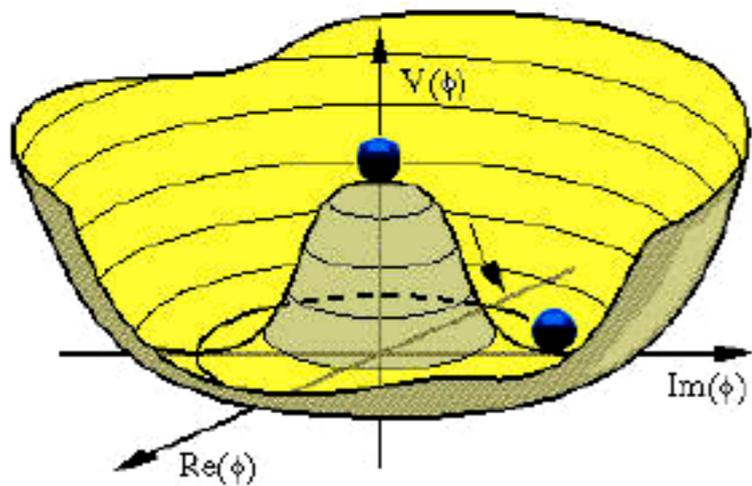
Higgs Chapter of the report on Physics at 100 TeV (R. Contino et al.)

And also with B. Bellazzini, C. Csaki, J. Hubisz, J. Serra, J. Terning (arXiv:1511.08218)

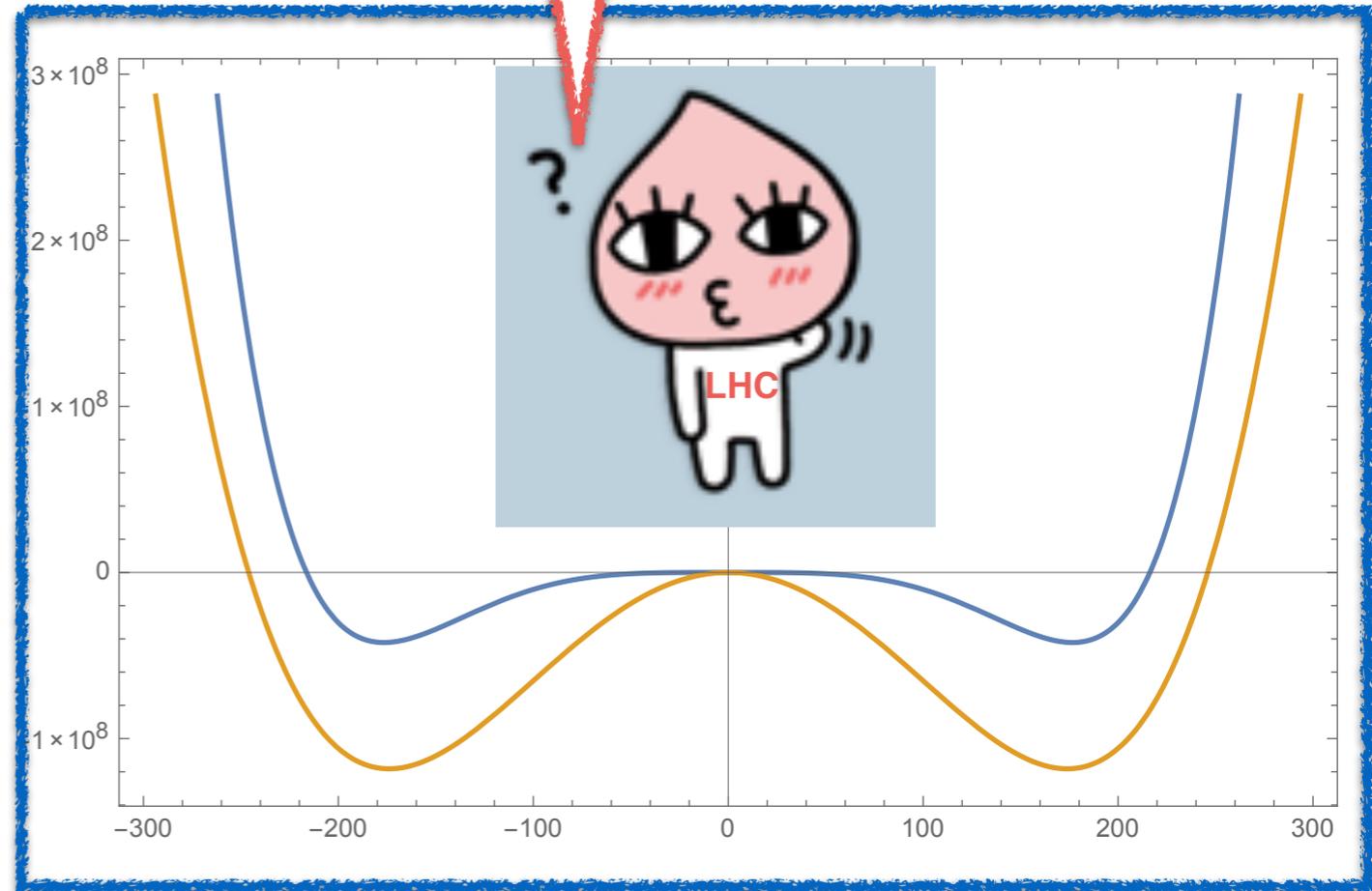
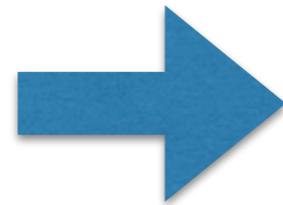
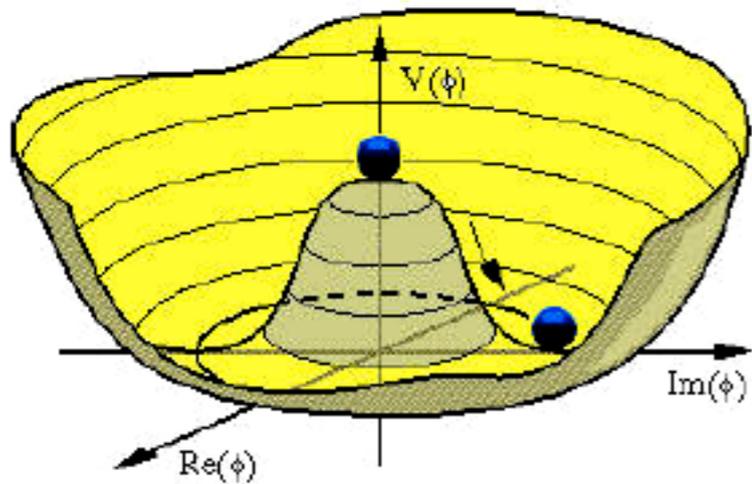
How well do we understand the EWSB, i.e. Higgs Potential?



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$$v^2 = \frac{4\lambda}{3c_6}, \quad \frac{m_h^2}{v^2} = 2\lambda \quad \xrightarrow{\text{experimentally}} \quad \lambda \approx 0.13$$

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This is just one simple example that shows that the form of the Higgs potential is completely undetermined

Higgs Potential: toy model

- ◆ However, FCC can potentially distinguish between the two possibilities just presented.
- ◆ In particular, via double Higgs production, the triple Higgs coupling can be probed: i.e.

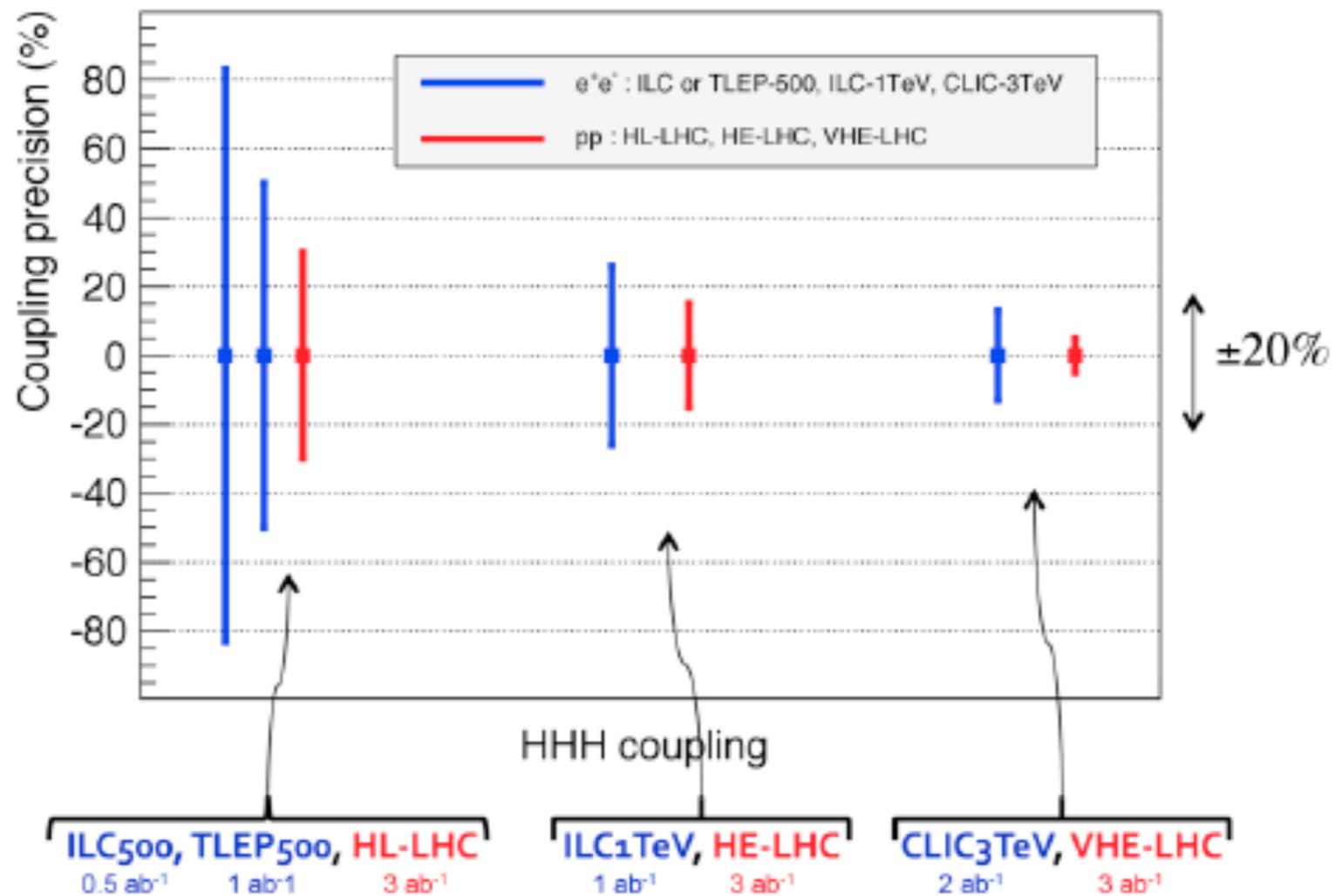
$$\lambda_{hhh} = 3 \frac{m_h^2}{v} \quad \longleftrightarrow \quad \tilde{\lambda}_{hhh} = 7 \frac{m_h^2}{v}$$

Higgs Potential: toy model

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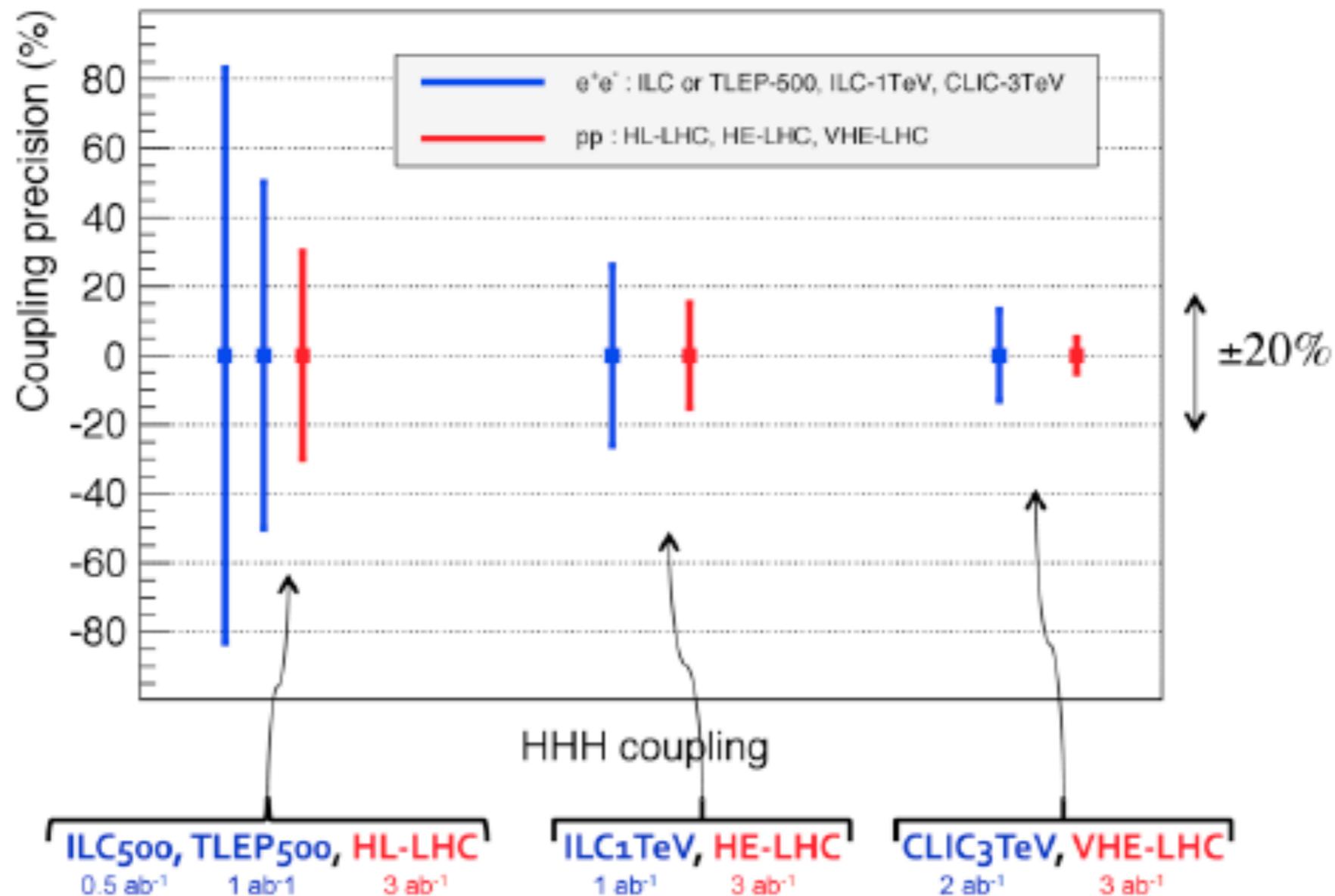
◆ In particular, Higgs couplings

λ_{hhh}



ple

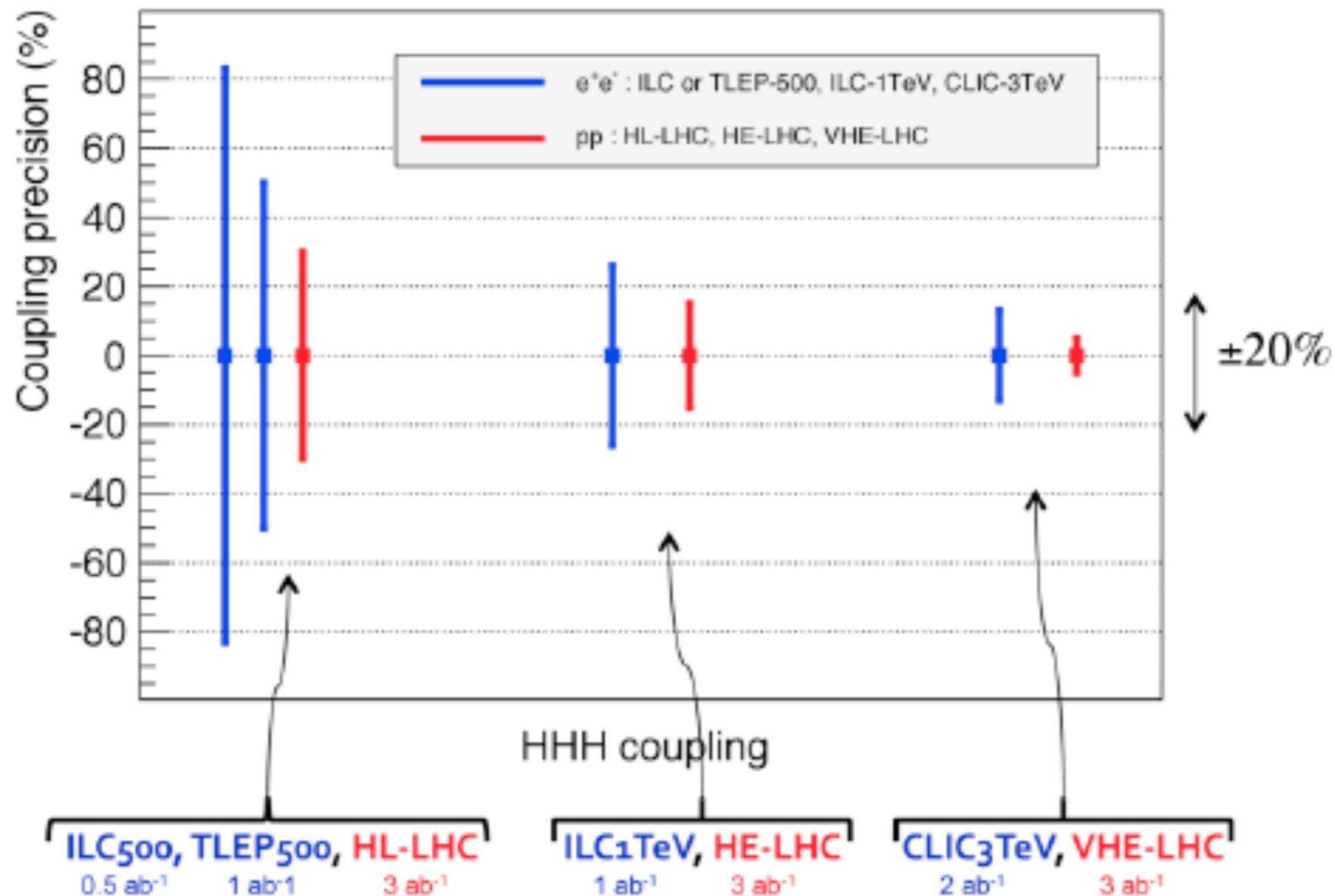
Higgs Potential



- ◆ we need future colliders to understand the underlying dynamics of the EWSB

Higgs Potential

see Roberto Contino's talk!



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Probing the EWSB mechanism

with Benjamin Fuks, Jeong Han Kim

- ◆ Establishing the SM nature of the electroweak symmetry breaking mechanism
 - ♣ Finding a Higgs is the first step ✓
 - ♣ Deriving the form of the scalar potential is necessary ✗
- ◆ Establishing the SM nature of the Higgs boson ?
 - ♣ Measurements of the Yukawa interaction strengths (i.e., the fermion masses)

◆ The scalar potential

$$V_h = \frac{m_h^2}{2} h^2 + (1 + \kappa_3) \lambda_{hhh}^{\text{SM}} v h^3 + \frac{1}{4} (1 + \kappa_4) \lambda_{hhhh}^{\text{SM}} h^4 \quad \text{with} \quad \lambda_{hhh}^{\text{SM}} = \lambda_{hhhh}^{\text{SM}} = \frac{m_h^2}{2v^2}$$

Measured

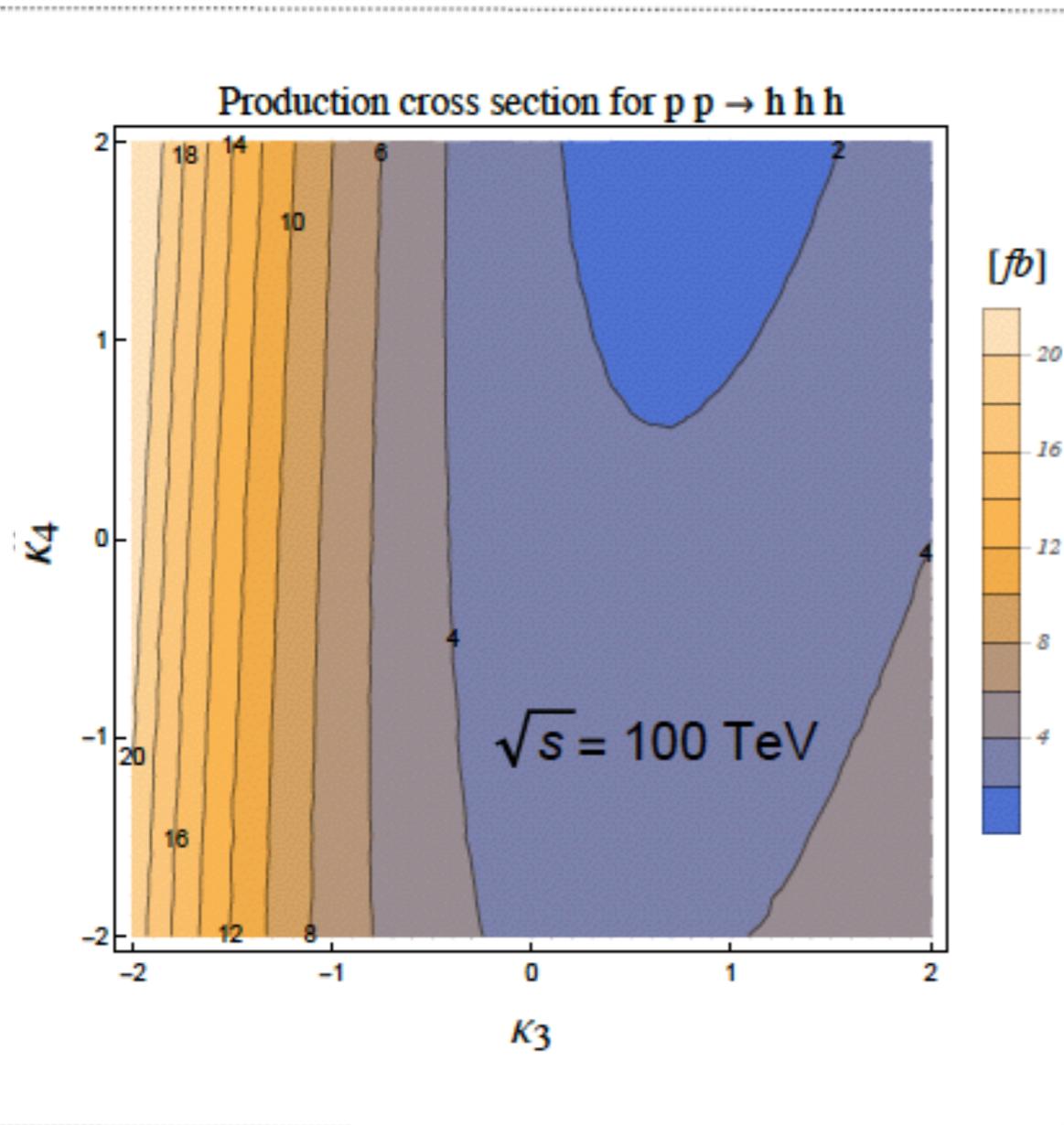
is $\kappa_3 = 0$?
Di-Higgs production
at the LHC can help

is $\kappa_4 = 0$?
No way for the LHC
(hhh cross section of ~ 0.01 fb)

Idea: probing triple-Higgs
production at the FCC

Sensitivity to both
 κ parameters

Triple Higgs production at the FCC



◆ Triple Higgs production total rate

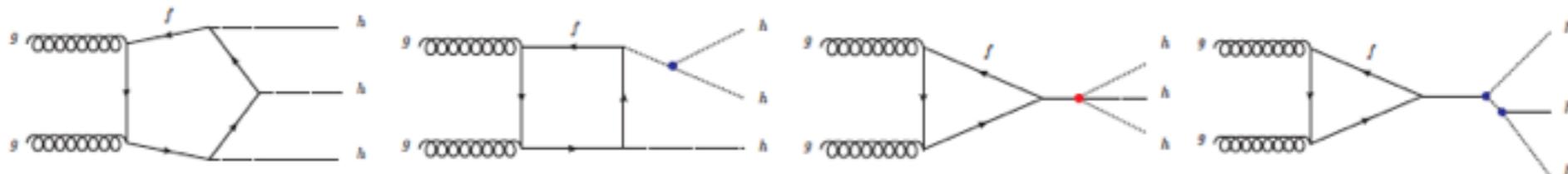
- ❖ Small in the Standard Model ($\sim 3\text{-}4 \text{ fb}$)
- ❖ **Could be much larger with new physics**

◆ Dependence on the κ parameters

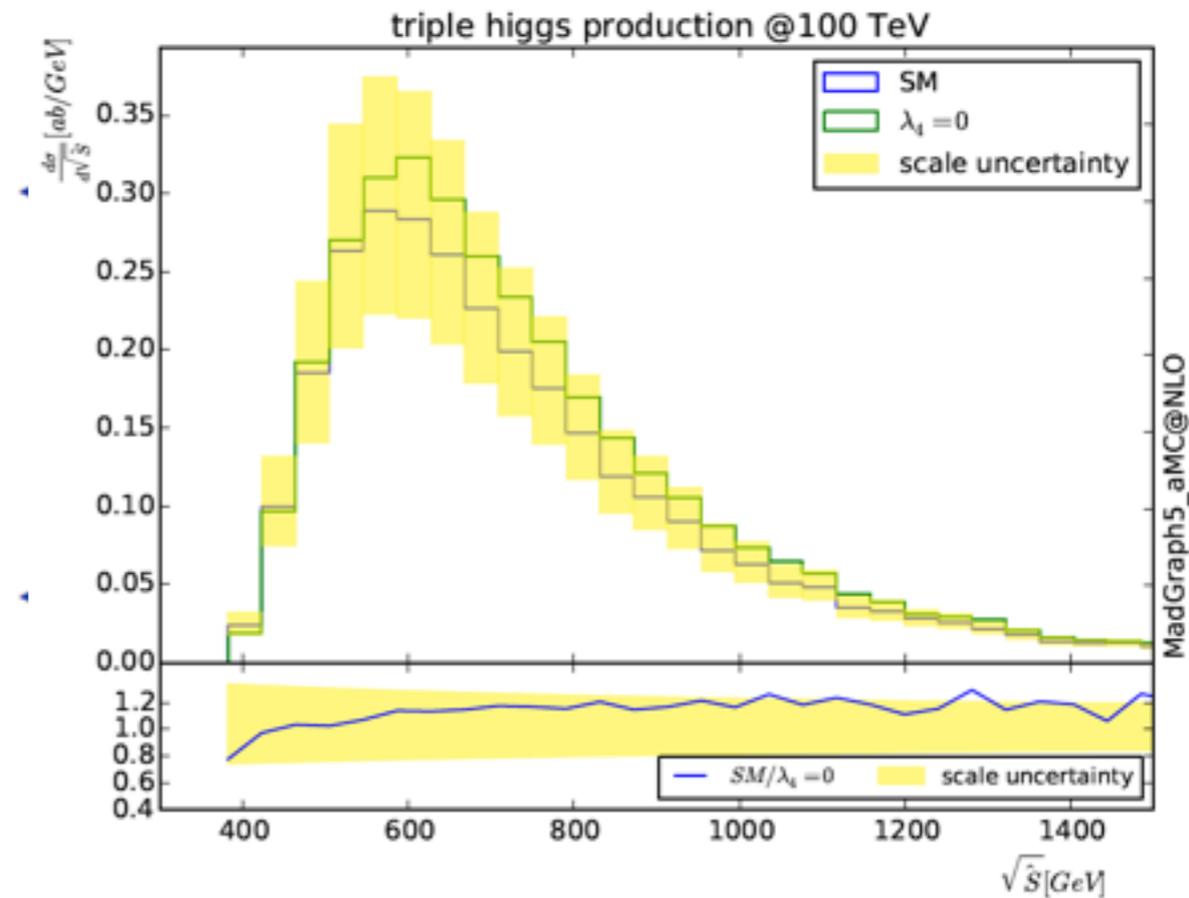
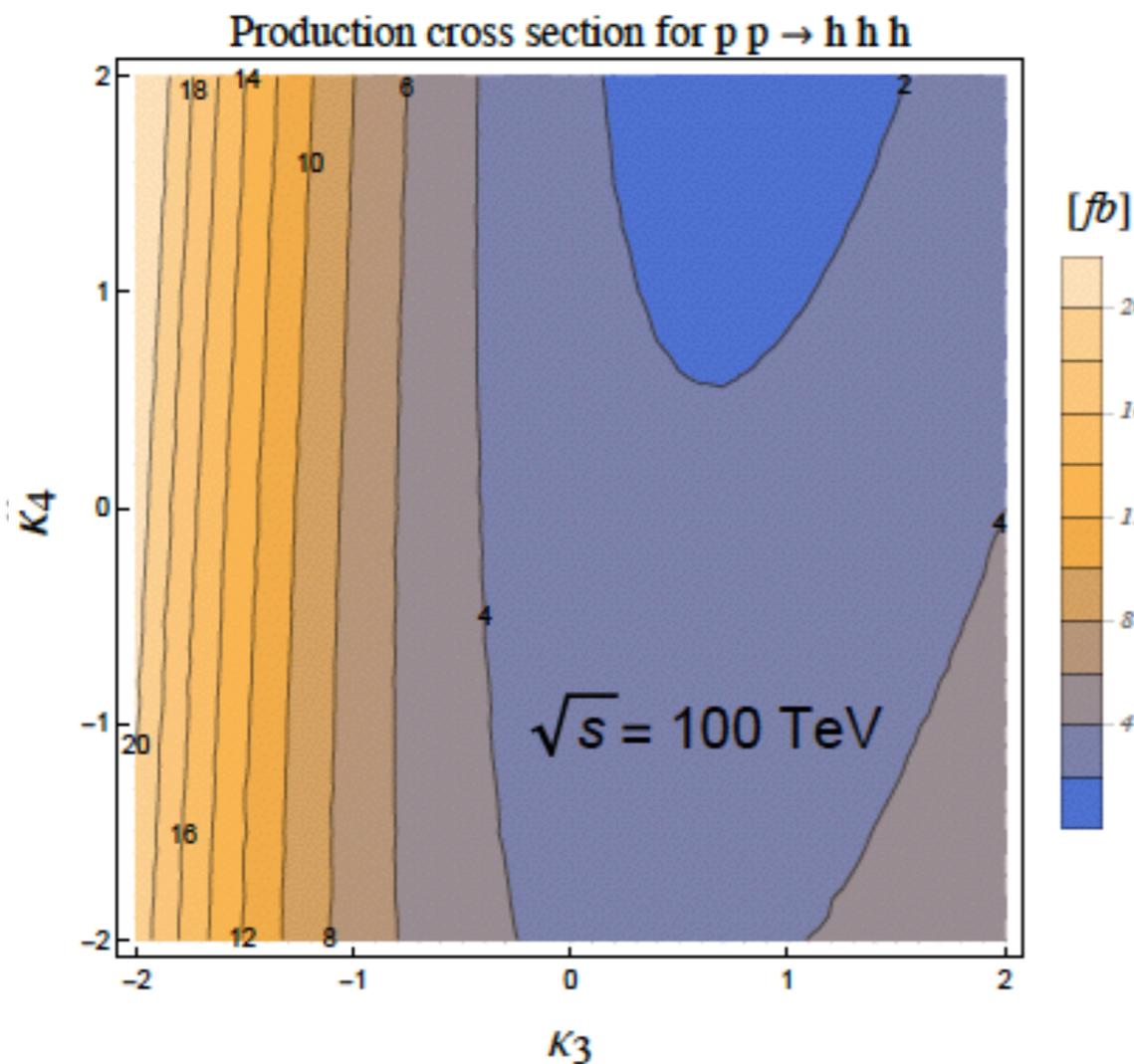
- ❖ κ_3 : very strong
 - Can reach 20 fb for large and negative values
- ❖ κ_4 : milder (but not negligible)

Could it be constrained?

$$(1 + \kappa_3) \lambda_{hhh}^{\text{SM}} v h^3 + \frac{1}{4} (1 + \kappa_4) \lambda_{hhhh}^{\text{SM}} h^4$$



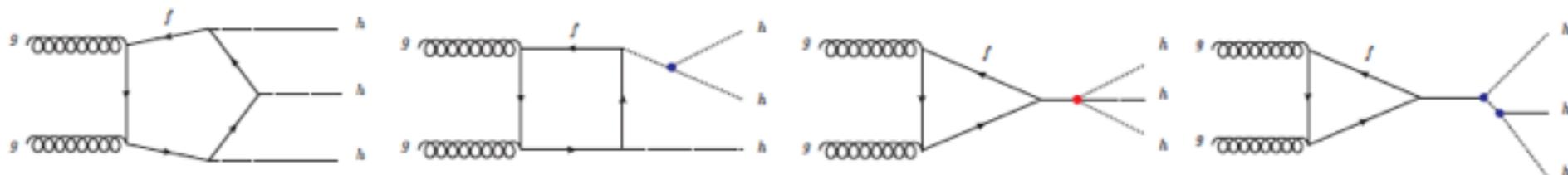
Triple Higgs production at the FCC



values

Could it be constrained?

$$(1 + \kappa_3) \lambda_{hhh}^{\text{SM}} v h^3 + \frac{1}{4} (1 + \kappa_4) \lambda_{hhhh}^{\text{SM}} h^4$$



Considered triple Higgs signals

◆ Decay modes with the larger branching ratios

- ❖ 4 b-jets and a 2 W's (22%): boosted techniques?
 - see the di-Higgs case at the LHC
- ❖ 6 b-jets (19.5%): boosted techniques or angular information?
 - see the di-Higgs case at the LHC

Necessitates FCC
detector details



Study of
other channels

◆ Decay in the $4b + 2\gamma$ channel

- ❖ Clean, low background (cf. the diphoton)
- ❖ Small branching fraction (0.232%)
- ❖ **Only studied channel so far** [Papaefstathiou & Sakurai] [Chen, Yan, Zhao, Zhong & Zhao]

◆ Decay in the $4b + 2\tau$ channel

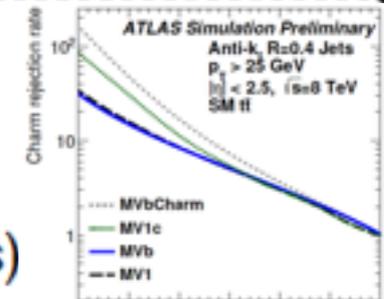
- ❖ Large branching fraction (6.46%)
- ❖ **Not studied yet**, but the $2b+2\tau$ channel is interesting in the di-Higgs case at the LHC

Why not at the FCC?

The $4b + 2\gamma$ channel: generalities

◆ Simulation details

- ❖ Parton-level study + smearing of the four-momenta (*à la* ATLAS)
- ❖ **b-tagging performance**: two LHC-inspired working points
 - ★ Very efficient (70%) with larger fake rates (18%, for c-jets, 1% for lighter jets)
 - ★ Less efficient (60%) with smaller fake rates (1.8%, for c-jets, 0.1% for lighter jets)



How good should the b-tagging be to observe a triple-Higgs signal

◆ Selection strategy for 20 ab^{-1}

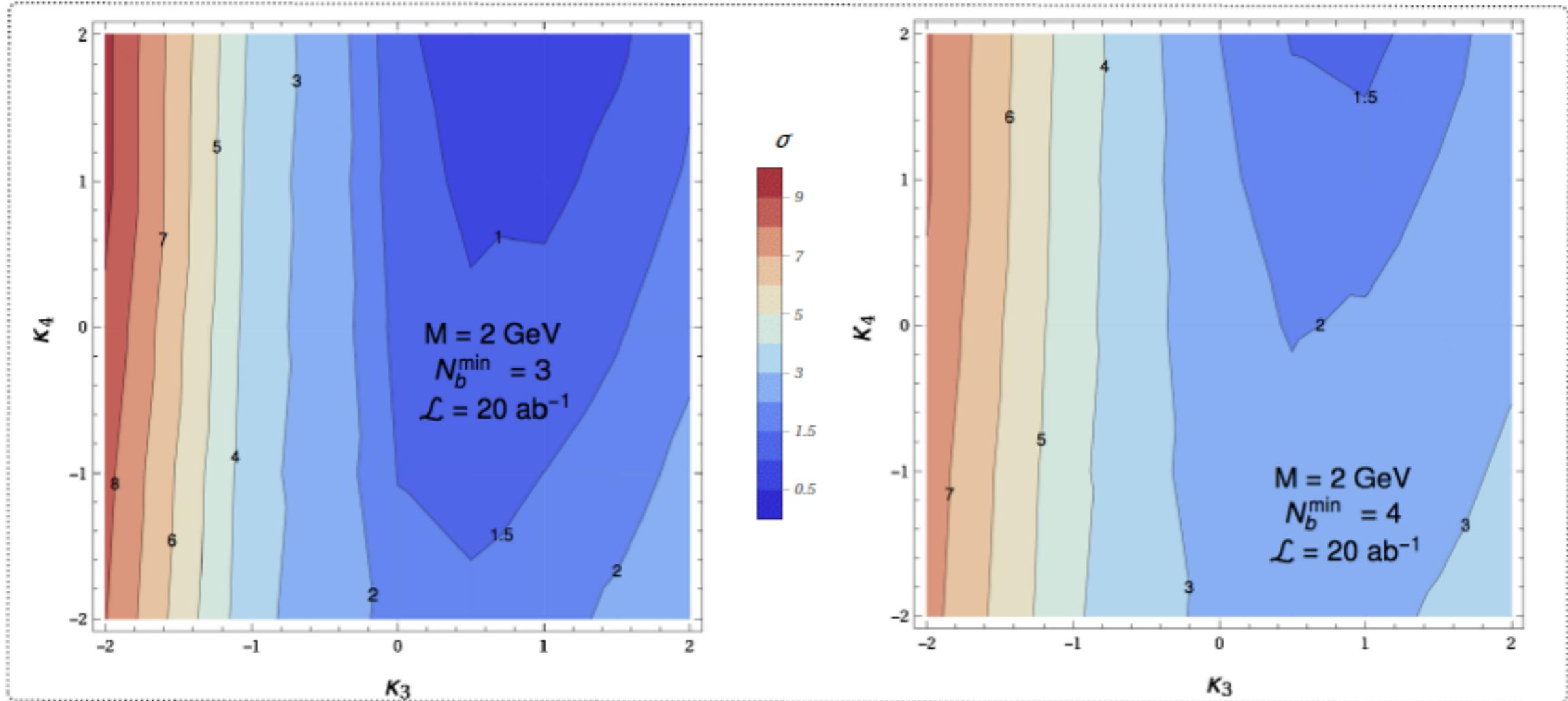
- ❖ Four jets (with an invariant mass smaller than 600 GeV), 2 photons
- ❖ Two dijet systems compatible with a Higgs ($m_{jj} \in [105, 140]$ GeV)
- ❖ The diphoton system compatible with a Higgs ($m_{\gamma\gamma} \in [125-M, 125+M]$ GeV)

What is the best M-value?

- ❖ At least N_b^{\min} b-tagged jets

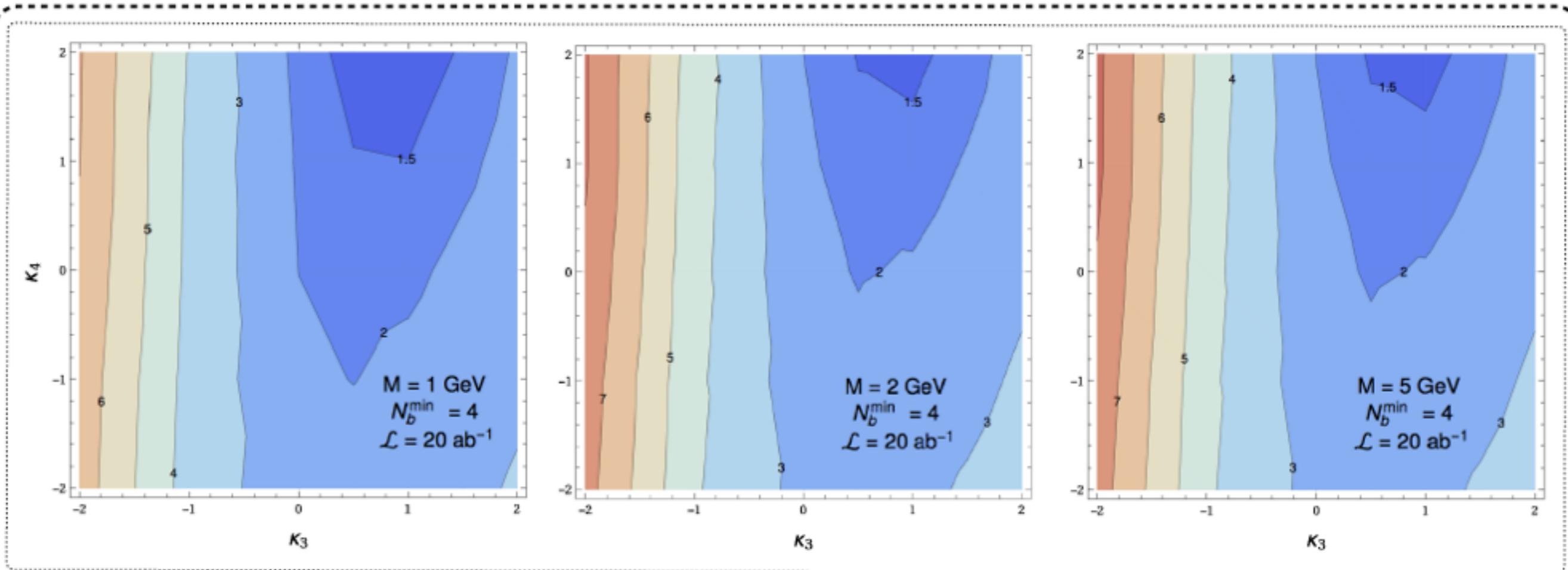
What is the best choice?

The $4b + 2\gamma$ channel: b-tagging considerations



- ◆ A low fake rate is primordial for a good sensitivity (1.8%/0.1% for a 60% efficiency)
 - ♣ Requiring at least 4 b-jets gives slightly better results (the background efficiency drops faster than the signal one)
 - ♣ **Poorer results for a fake rate of 18%/1% for a 70% efficiency**
 - ★ Better signal acceptance
 - ★ Much worse background contamination due to the fakes

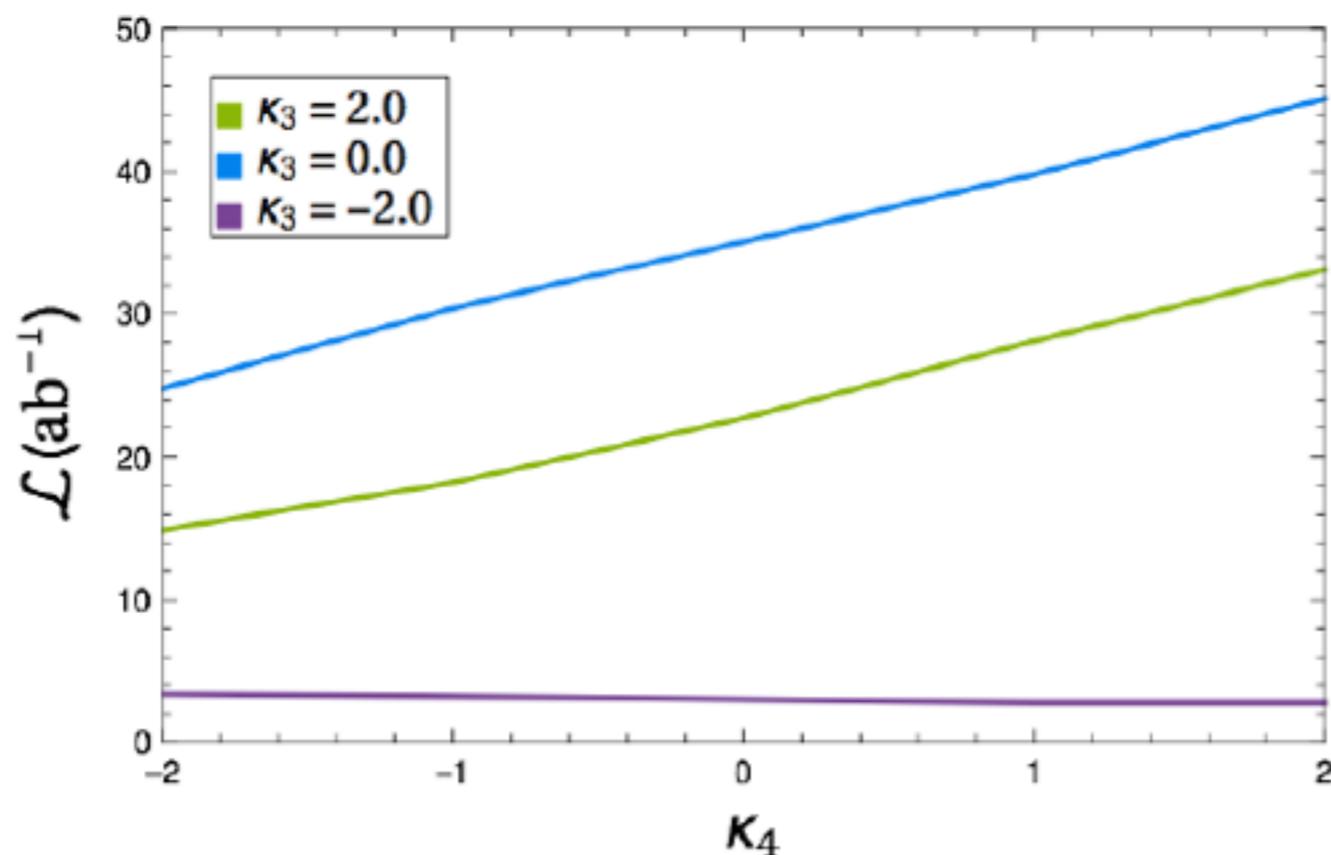
The $4b + 2\gamma$ channel: diphoton mass resolution



- ◆ Photons with a p_T greater than 20 GeV are very well reconstructed ($\sigma/E \sim 0.1/\sqrt{E}$)
- ♣ A loss of signal efficiency implies to maintain M not too small
- ♣ A too large M implies a more important background contamination
 - ★ However, mild effects on the sensitivity

$M = 2$ GeV gives the best results

The $4b + 2\gamma$ channel: luminosity goals for 3σ



- ✦ At least 4 b-tags
- ✦ $M = 2$ GeV
- ✦ b-tagging: 60% / 1.8% / 0.1%

- ◆ Large and negative $\kappa_3 \Leftrightarrow$ huge cross section \Leftrightarrow BSM hints reachable with low luminosity
- ◆ Other κ_3 values \Leftrightarrow more luminosity is required
 - ✦ The scanned region of the parameter space cannot be entirely covered

The $4b + 2\tau$ channel: generalities

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- ❖ **b-tagging performance**: two LHC-inspired working points
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 - ★ Less efficient (60%) with smaller fake rates (1.8%, for c-jets, 0.1% for lighter jets)
- ❖ **tau-tagging performance**: two LHC-inspired working points
 - ★ Very efficient (80%) with a small fake rate (0.1%)
 - ★ More conservative (50%) with a larger fake rates (1%,)

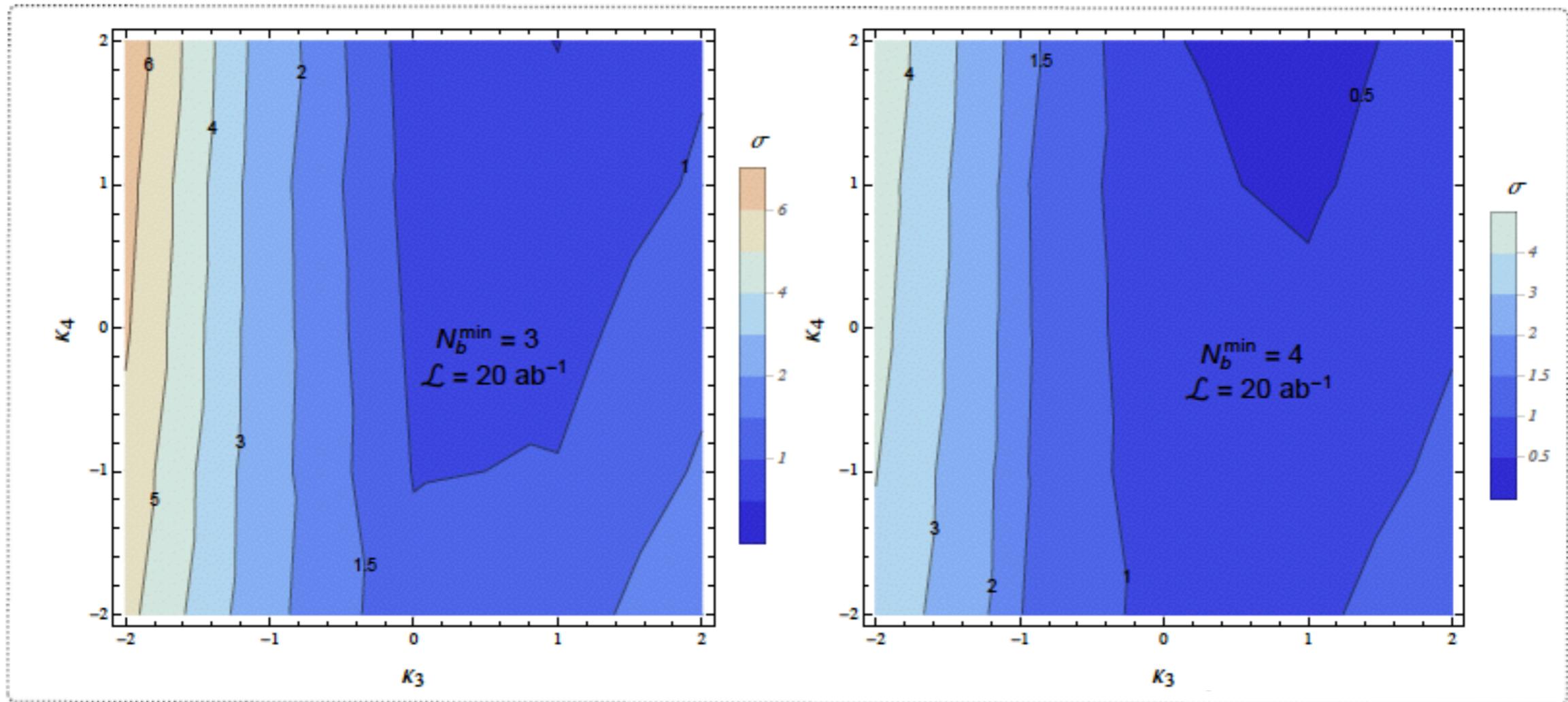
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- ❖ At least N_b^{min} b-tagged jets

What is the best choice?

The $4b + 2\tau$ channel: b- and tau-tagging



- ◆ A good b-tagging efficiency is essential (70% for an 18% / 1% fake rate)
 - ♣ Requiring at least 3 b-jets gives slightly better results (better signal efficiency)
 - ♣ **Smaller fake rate and efficiency: the signal efficiency drops faster than the background one**
- ◆ A very efficient tau-tagger is essential (80% for a 0.1% fake rate)
 - ♣ All sensitivity is lost for a more conservative choice of (50% / 1%)

Summary

- ◆ We considered triple Higgs production at the FCC

- ❖ Decay in the $4b + 2\gamma$ and $4b + 2\tau$ channels
- ❖ Effects of the b-tagging and tau-tagging performances on the sensitivity
- ❖ Effect of the diphoton mass resolution

- ◆ The diphoton channel yields the best sensitivity

- ❖ Controlling the b-tagging fake rate is mandatory
- ❖ The diphoton mass requirement cannot be too tight
- ❖ A good fraction of the parameter space can be covered (but not the SM case)
 - ★ A small part even with a few ab^{-1}

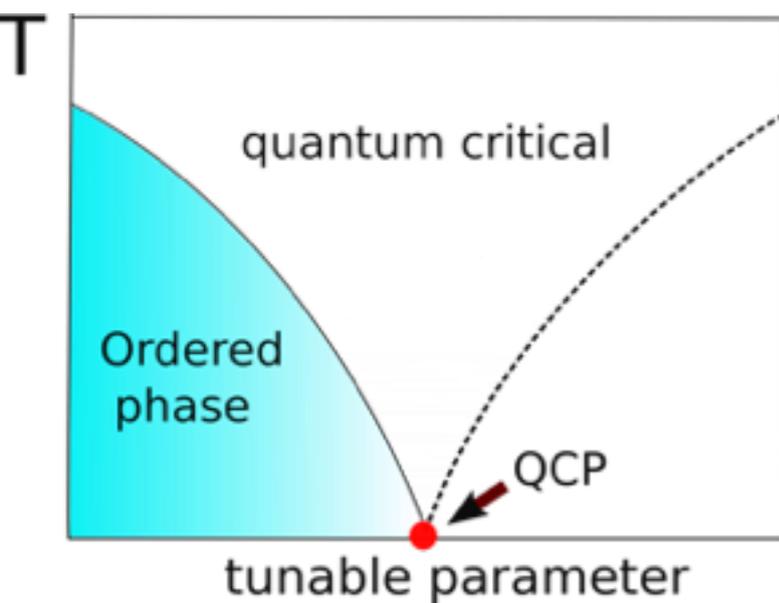
- ◆ The ditau channel could be complementary

- ❖ An excellent tau-tagger is required

Post Higgs question: The Quantum Critical Higgs

with B. Bellazzini, C. Csaki, J. Hubisz, J. Serra, J. Terning (arXiv:1511.08218)

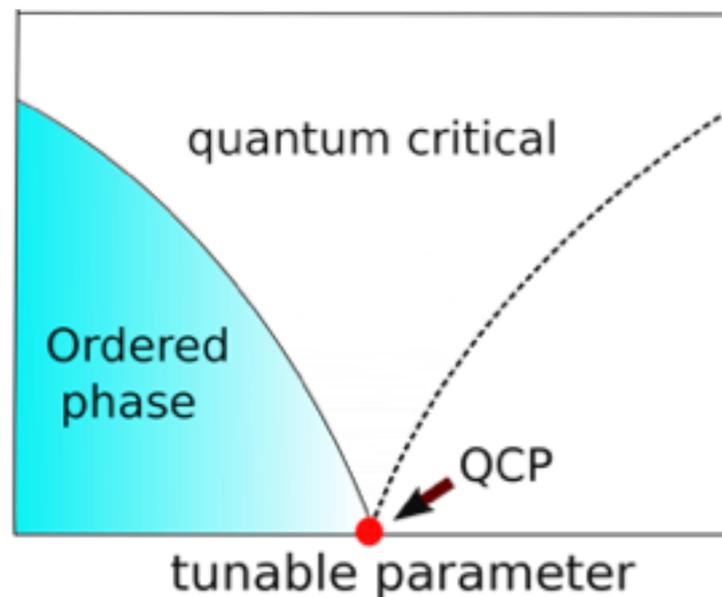
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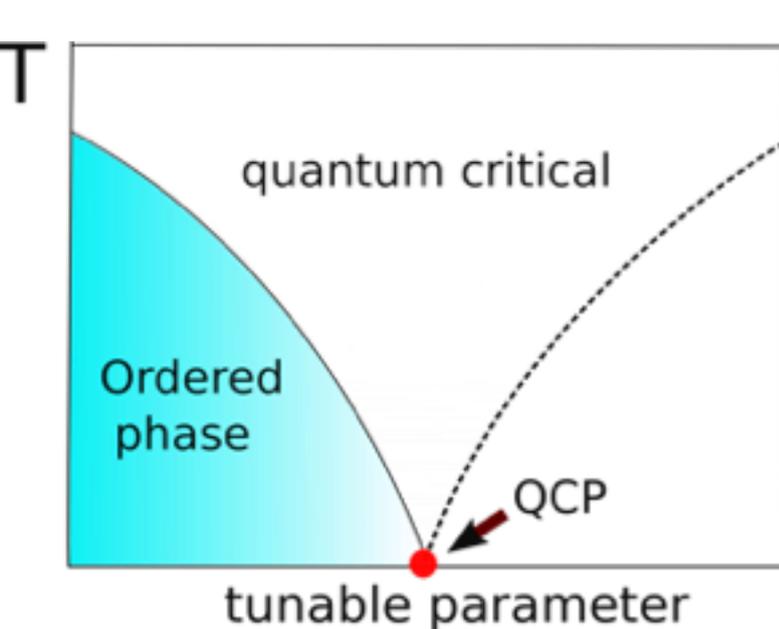


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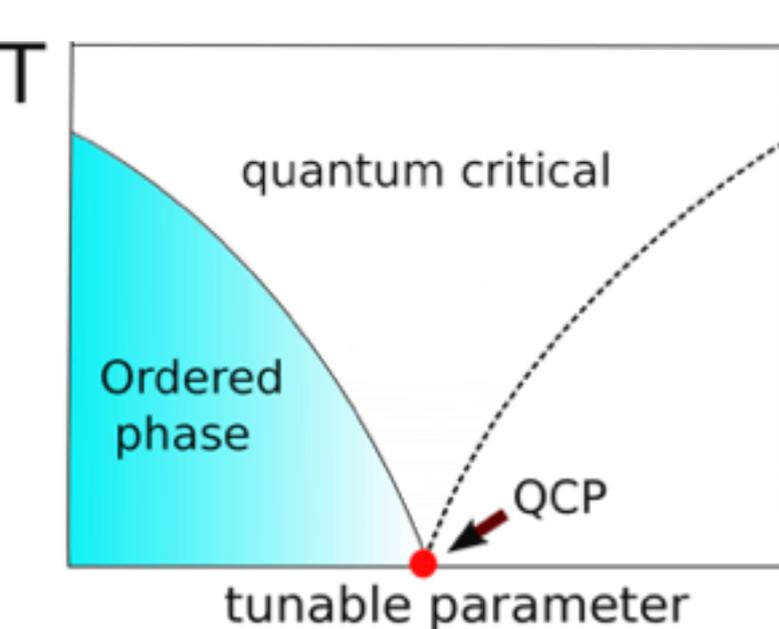


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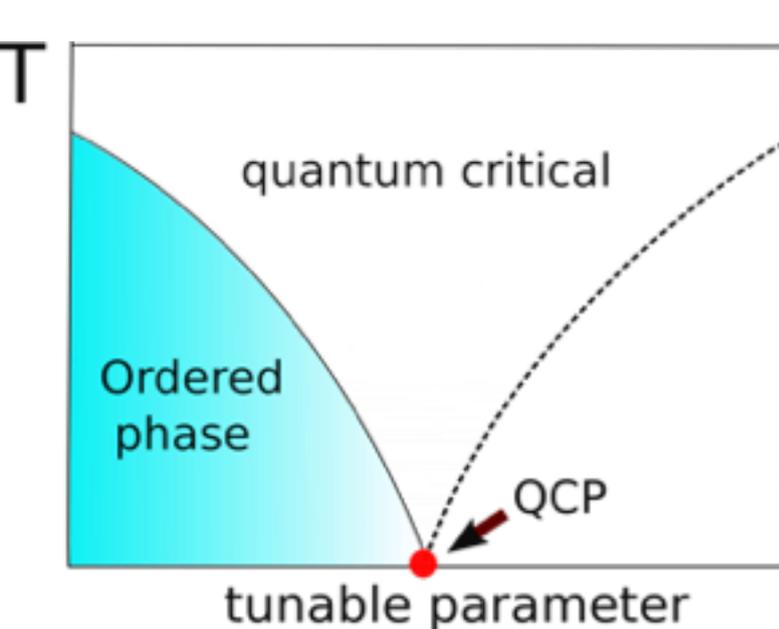


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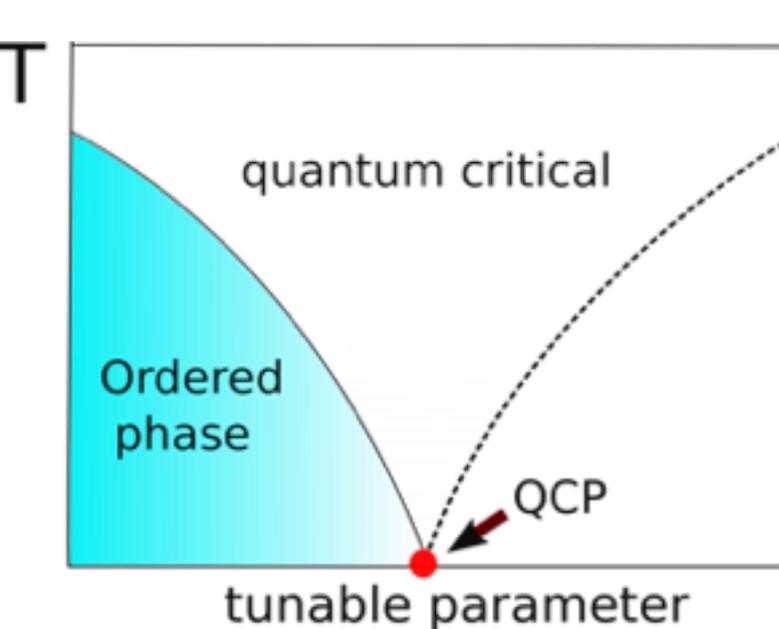
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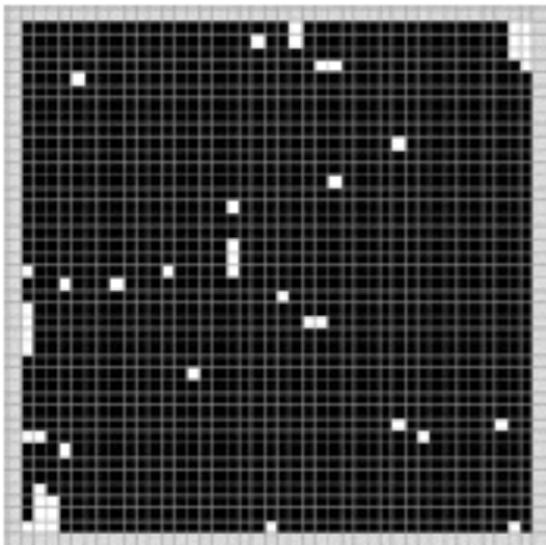
- ◆ and if so, whether it is more interesting than mean-field theory like the Standard Model of particle physics.

Ising Model

$$H = -J \sum s(x)s(x+n)$$

$$s(x) = \pm 1$$

Low T



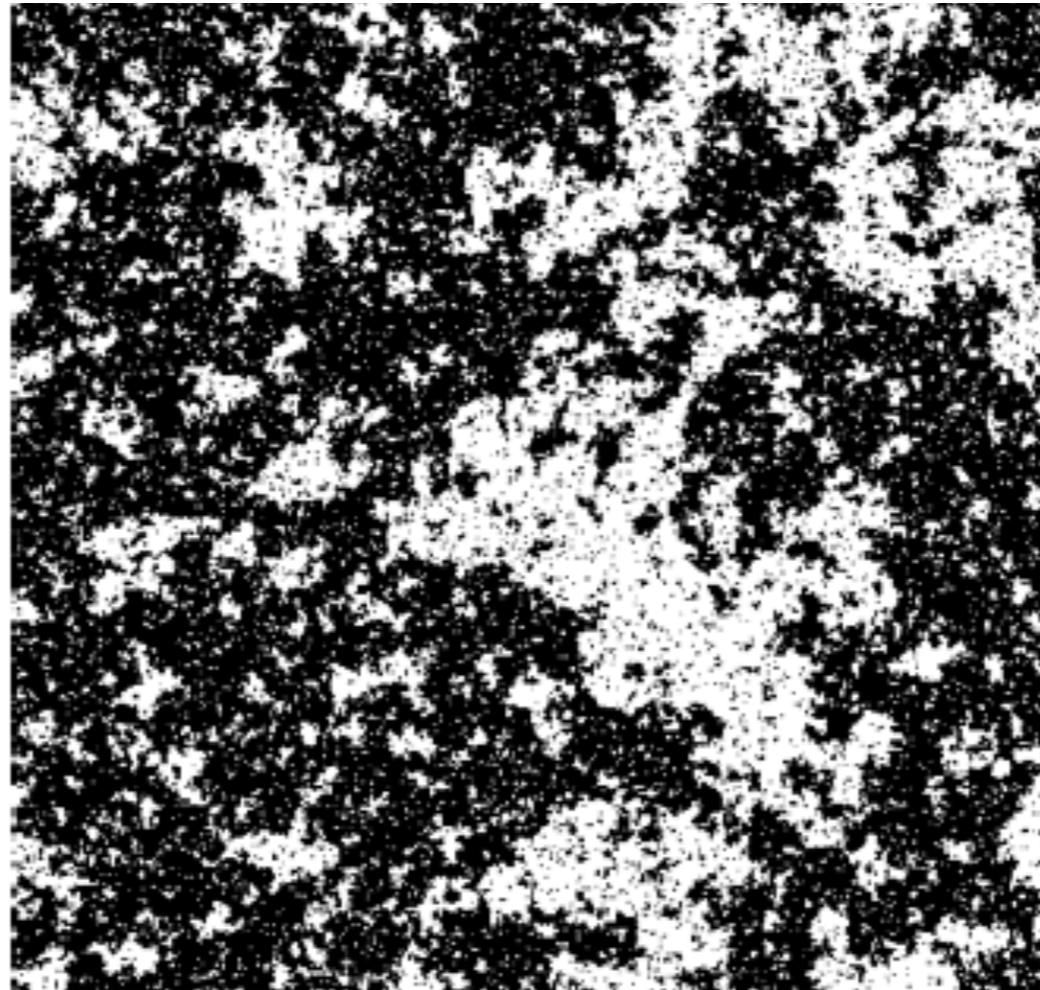
High T

T_c

$$\langle s(0)s(x) \rangle = e^{-|x|/\xi}$$

at $T=T_c$ $\xi \rightarrow \infty$

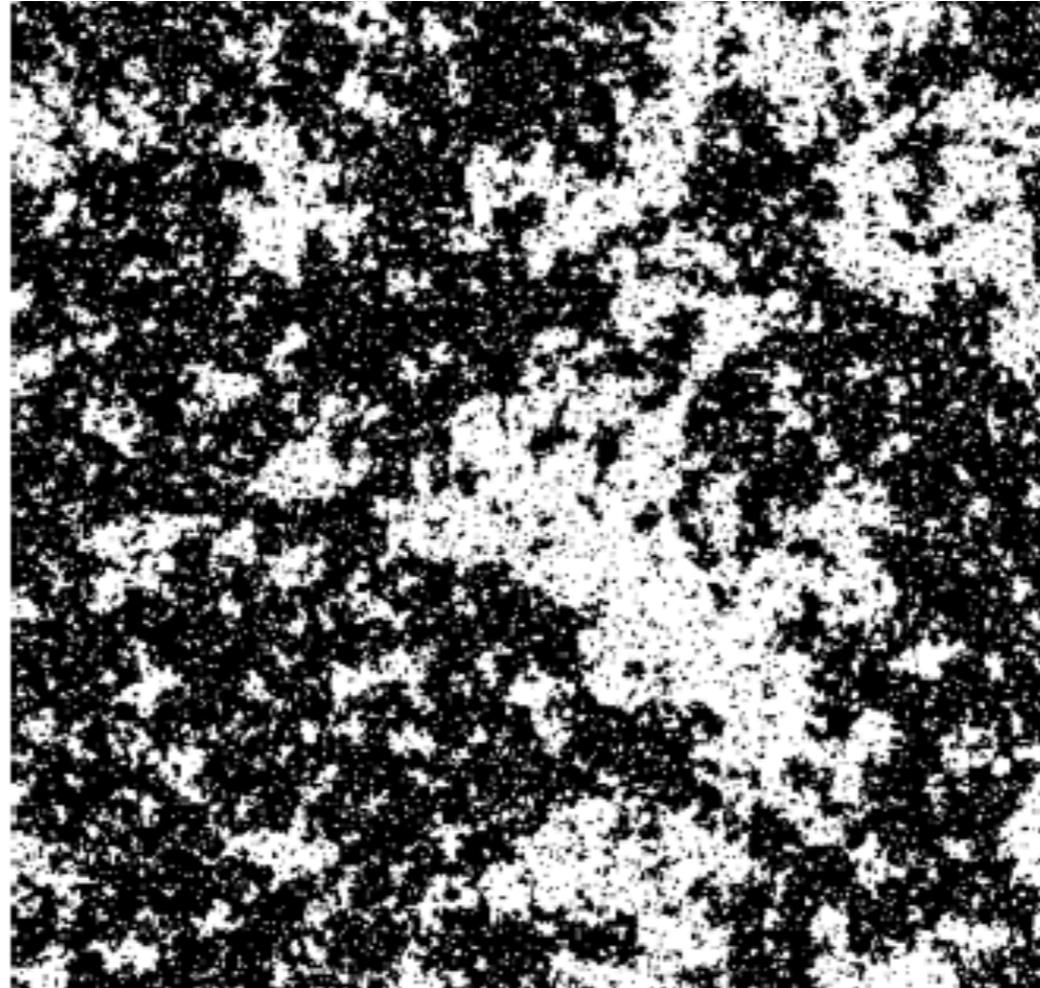
Critical Ising Model is Scale Invariant



<http://bit.ly/2Dcrit>

$$\text{at } T=T_c \quad \langle s(0)s(x) \rangle \propto \frac{1}{|x|^{2\Delta-1}}$$

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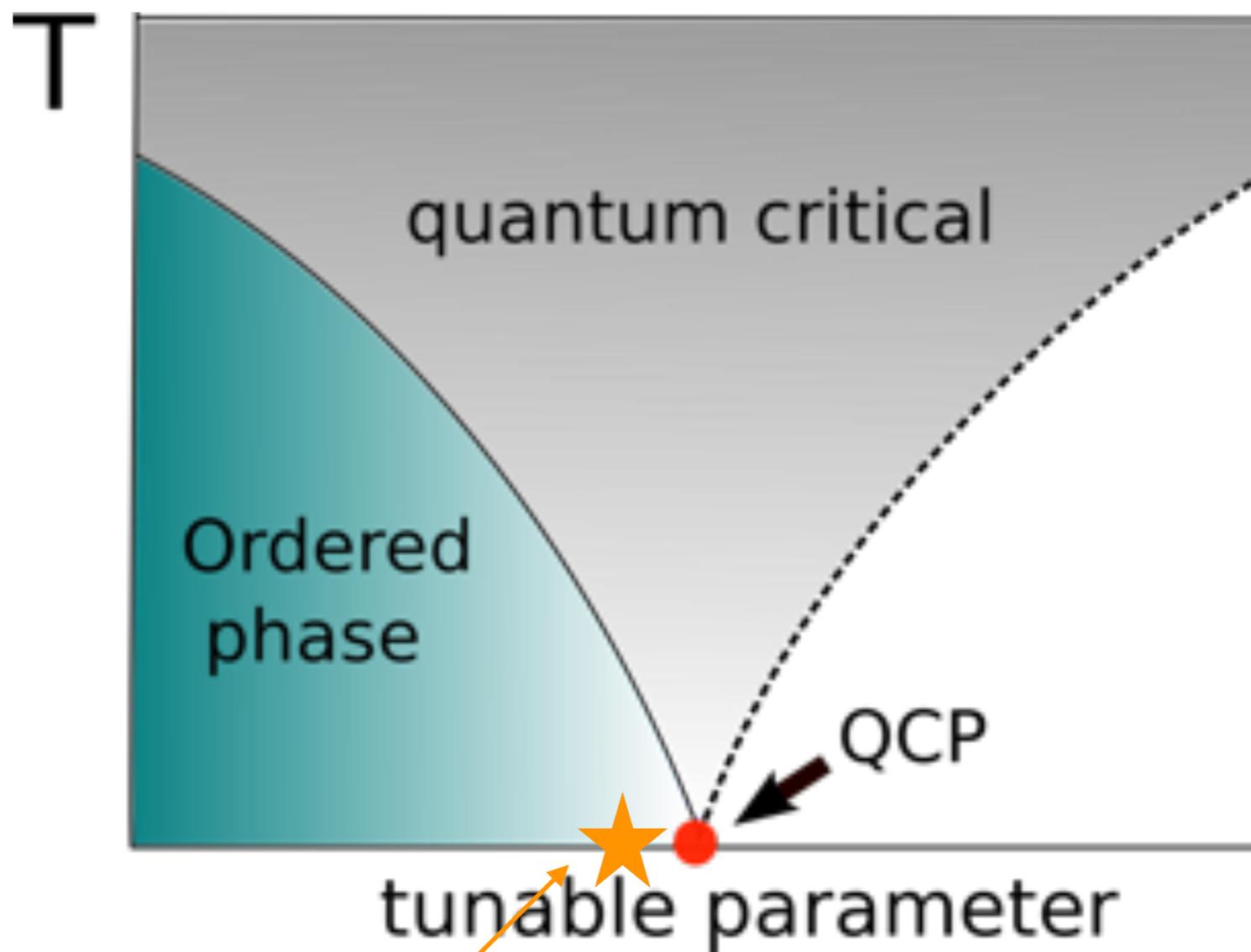


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$$\text{at } T=T_c \quad \langle s(0)s(x) \rangle \propto \frac{1}{|x|^{2\Delta-1}} = \int d^3p \frac{e^{ip \cdot x}}{|p|^{4-2\Delta}}$$

↑
critical exponent

Quantum Phase Transition



We are here

The Quantum Critical Higgs

- ◆ At a QPT the approximate scale invariant theory is characterized by [the scaling dimension \$\Delta\$](#) of the gauge invariant operators, where $1 < \Delta < 2$. In general, the two point function of a scalar with scaling dimension Δ in a CFT is

$$G_{\text{CFT}}(p) = -\frac{i}{(-p^2 + i\epsilon)^{2-\Delta}} \quad \text{SM: } \Delta = 1 + \mathcal{O}(\alpha/4\pi).$$

- ◆ We propose to present a general class of theories describing a Higgs field near a non-mean-field QPT.

- ◆ In such theories, in addition to the pole (Higgs), there can also be a [Higgs continuum](#), representing additional states associated with the dynamics underlying the QPT

$$G_h(p^2) = \frac{i}{p^2 - m_h^2} + \int_{\mu^2}^{\infty} dM^2 \frac{\rho(M^2)}{p^2 - M^2}$$

- ◆ One result of the presence of the continuum will be the appearance of [form factors](#) in couplings of the Higgs to the SM particles.

Form Factors for QCH

- ◆ What would be a general low-energy EFT consistent with a QPT and no new massless particles?

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 - => Higgs pick up a significant anomalous dimension, and there is a large mixing with the continuum

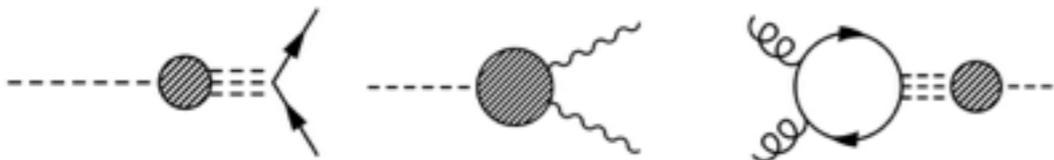
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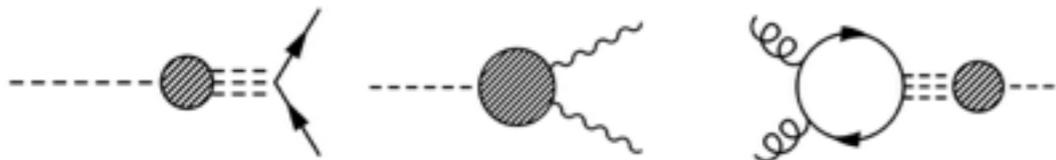
On-shell behavior:
constant form factors



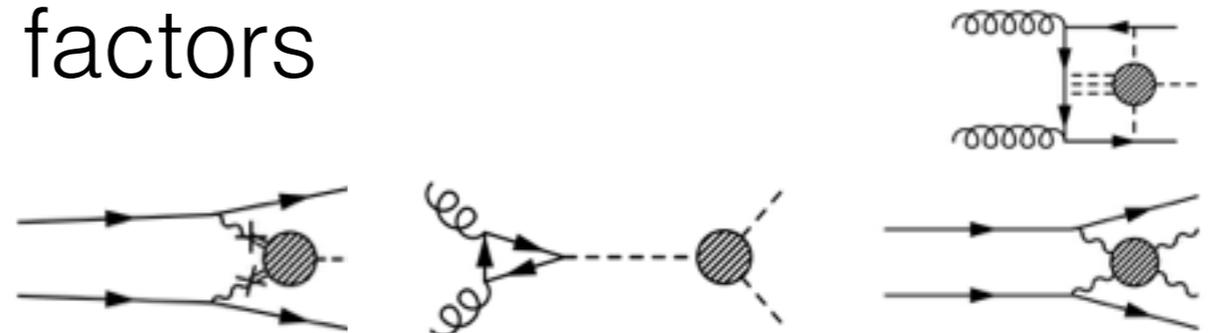
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On-shell behavior:
constant form factors



Off-shell behavior: nontrivial
momentum dependent form
factors



Form Factors for QCH

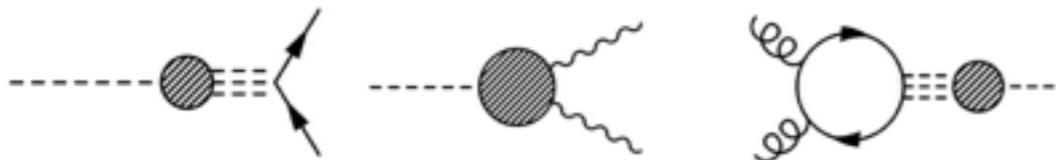
- ◆ What would be a general low-energy EFT consistent with a QPT and no new massless particles?
- ◆ We consider a QPT Higgs scenario where Higgs is (partially) imbedded into a strongly coupled sector (approximately conformal at scale well above the EW scale)

=> Higgs pick up a significant anomalous dimension, and there is a large mixing with the continuum

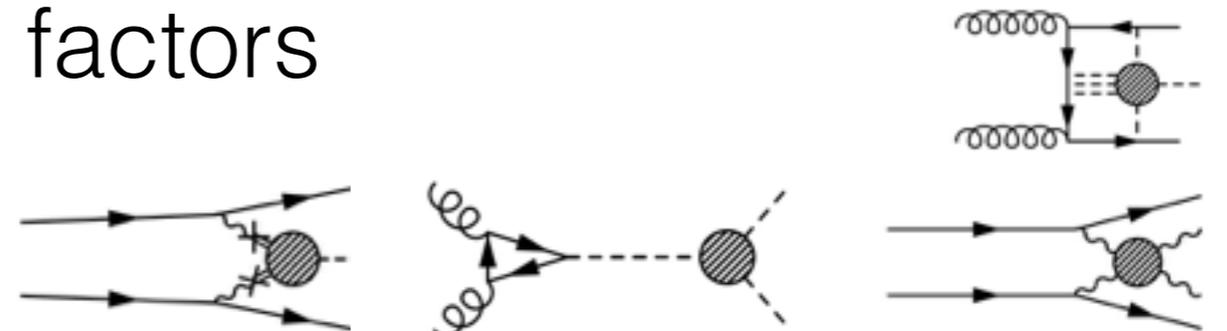
- ◆ The effects of Higgs can be parametrized in terms of form factors in a model independent way

need build a general low-energy EFT!

On-shell behavior:
constant form factors

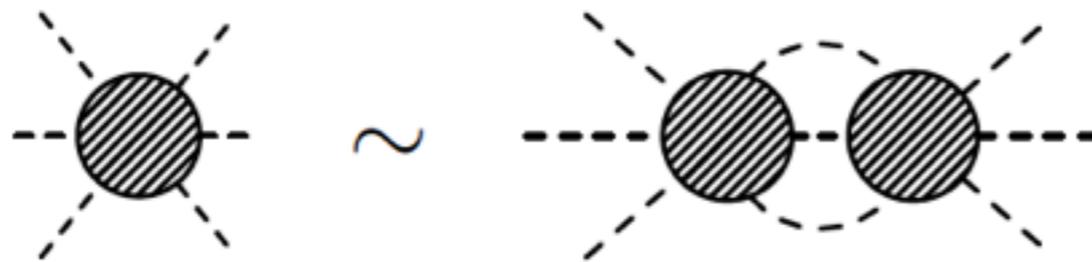


Off-shell behavior: nontrivial momentum dependent form factors



(NDA) Expectations of the size of the form factors

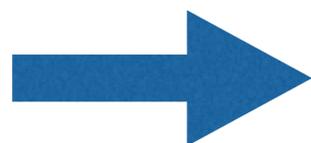
In **NDA**, loop corrections \sim original n-point function. e.g. $n = 6$:



$$\mathcal{L} = \frac{\alpha_n}{\Lambda^{n-4}} \phi^n$$

loop contribution with two insertions of this operator (that contributes to the same n-point) would be one in which each vertex has $n/2$ external lines, and $n/2$ propagators exchanged in $n/2 - 1$ loops.

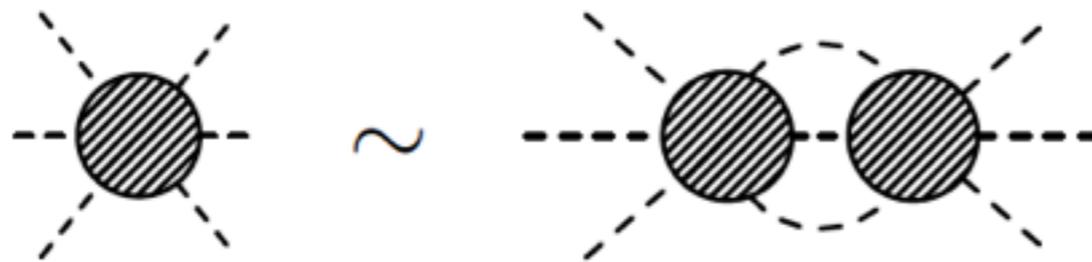
quantum corrections: $\alpha_n \rightarrow \alpha_n \left(1 + \frac{\alpha_n}{(16\pi^2)^{n/2-1}} \right)$



$$\alpha_n \sim (16\pi^2)^{n/2-1}$$

(NDA) Expectations of the size of the form factors

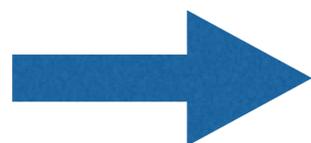
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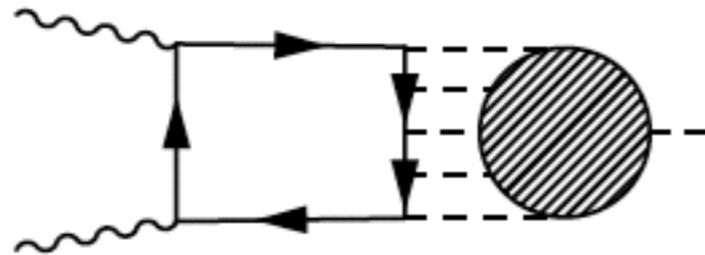
loop contribution with two internal lines (n-point) would be one in which $n/2$ propagators exchanged quantum corrections

e.g. with this, we can see that the n-point contribution to the gluon fusion process is suppressed by insertions of the perturbative coupling of the top-quark to the strongly coupled sector, along with a loop factor that is only partially cancelled by the large coefficient



$$\alpha_n \sim (16\pi^2)^{n/2-1}$$

Off-shell Form Factors for the Quantum Critical Higgs

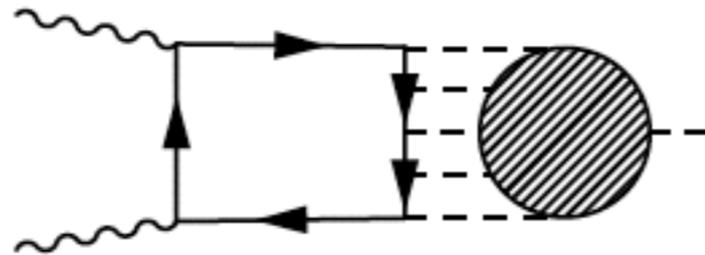


If the shaded region corresponds to the n -point function, there are $n - 1$ insertions of the top Yukawa, and $n - 2$ loops. There are $n - 1$ scalar propagators and $n - 2$ fermionic propagators running in these loops.

estimate for the contribution of the n -point correlator to the $h\bar{t}t$ coupling:

$$g_n^{tth} \sim 4\pi \left(\frac{\lambda_t}{4\pi} \right)^{n-1}$$

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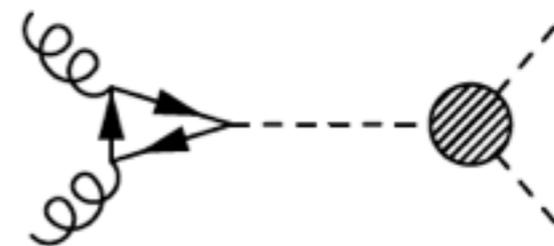


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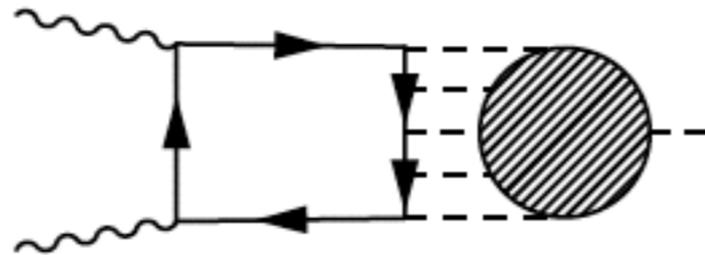
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If the shaded region corresponds to the n -point function, there are $n - 1$ insertions of the top Yukawa, and $n - 2$ loops. There are $n - 1$ scalar propagators and $n - 1$ fermionic propagators running in these loops.

dominant contribution comes from tree diagram

estimate for the contribution to the h production cross-section from the htt coupling:

e.g. double Higgs production through gluon fusion would be dominated by

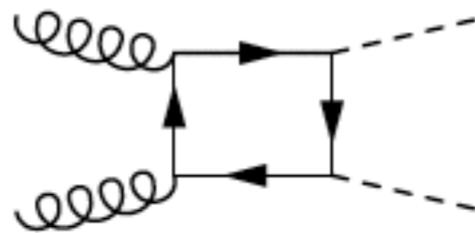


Direct Signal (e.g. Holographic Generalized Free Fields)

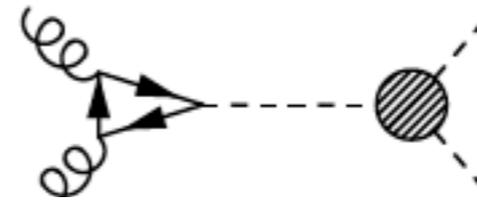
◆ Form factors for trilinear Higgs self coupling

- ❖ SO(4) global symmetry is gauged in the 5D bulk

$$\lambda_5 (H^\dagger H)^2$$

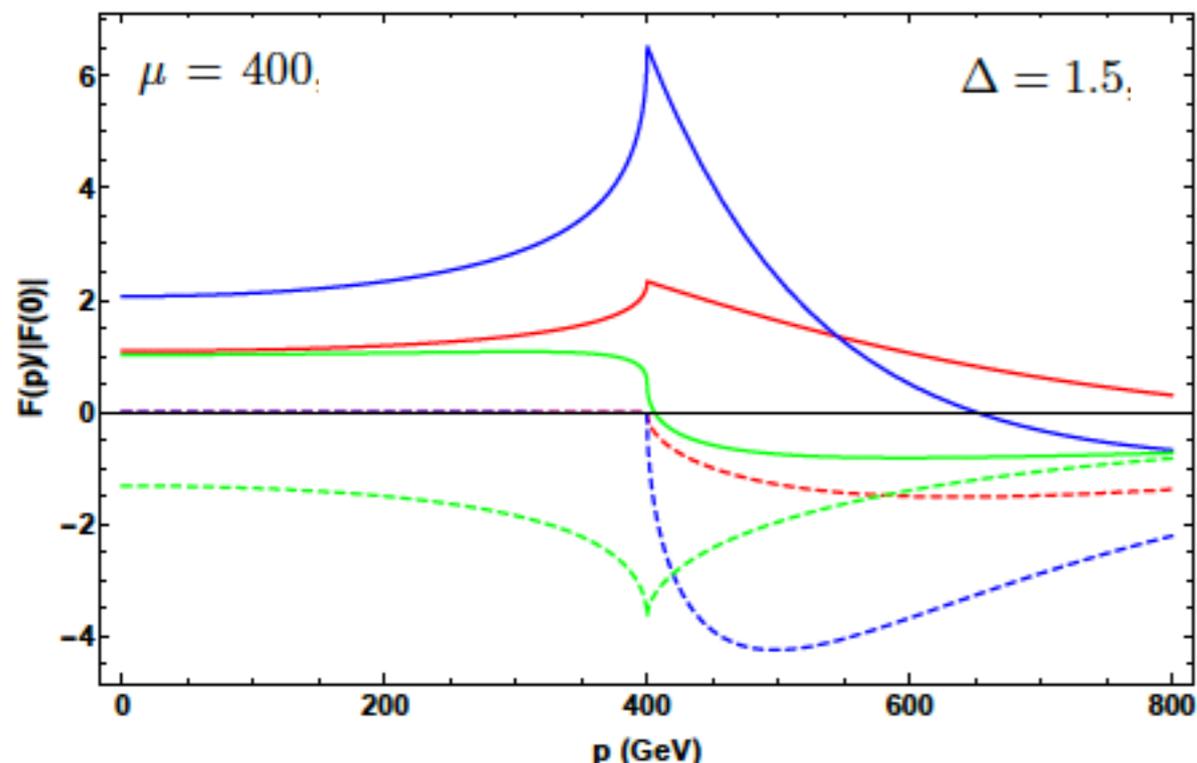


+crossings

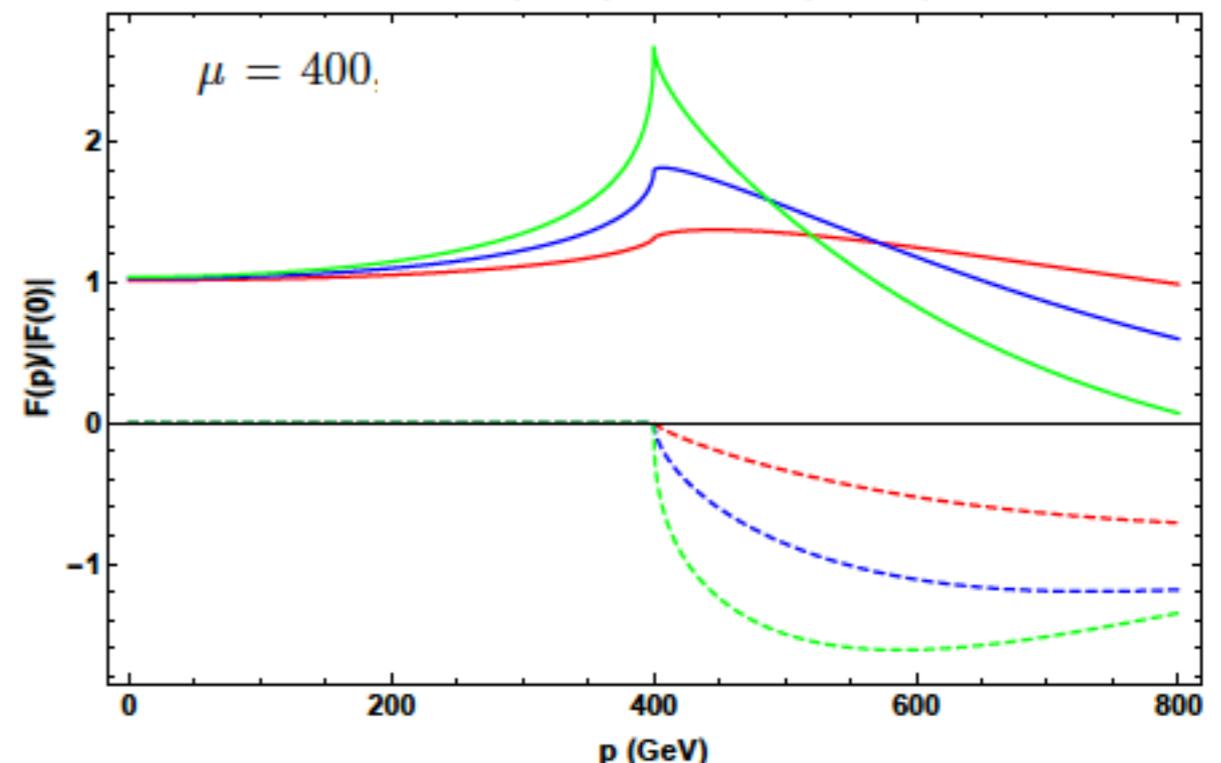


$$F_{hhh} = \frac{\lambda_5}{L^2} \mathcal{V} \int_R^\infty dz \frac{1}{a} \left(\frac{z}{R}\right)^2 \frac{K_{2-\Delta}(\mu z)}{K_{2-\Delta}(\mu R)} \prod_{i=1}^3 \frac{K_{2-\Delta}(\sqrt{\mu^2 - p_i^2} z)}{K_{2-\Delta}(\sqrt{\mu^2 - p_i^2} R)}$$

Higgs momentum: 200 GeV (Red), 400 GeV (Blue), and 600 GeV (Green)



$\Delta = 1.2$ (Red), 1.4 (Blue), and 1.6 (Green).



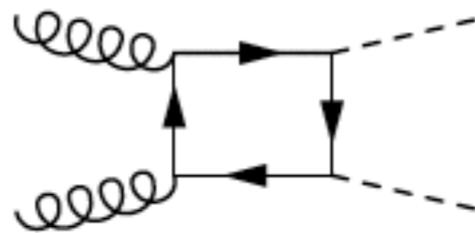
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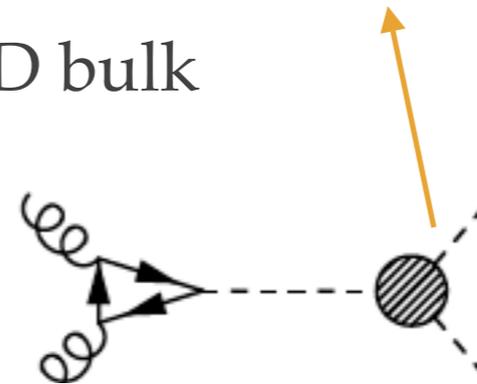
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probe the higher n-point correlators of the CFT.

$$\lambda_5 (H^\dagger H)^2$$

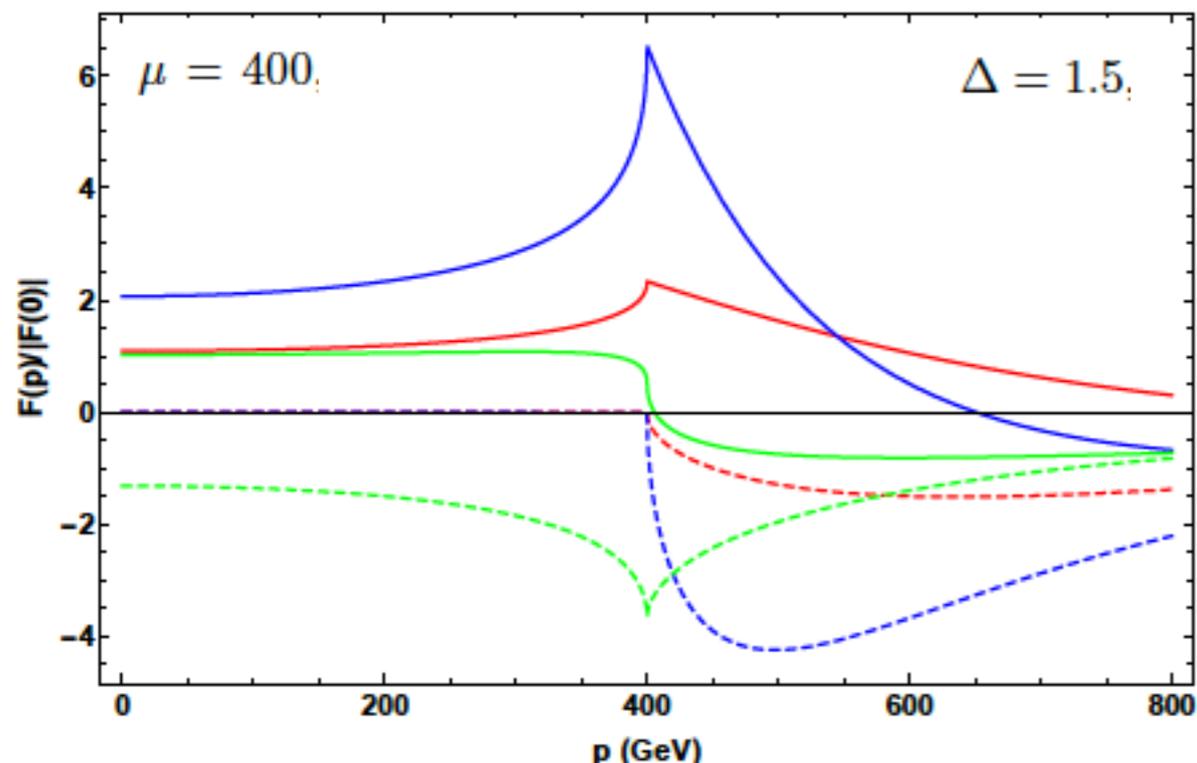


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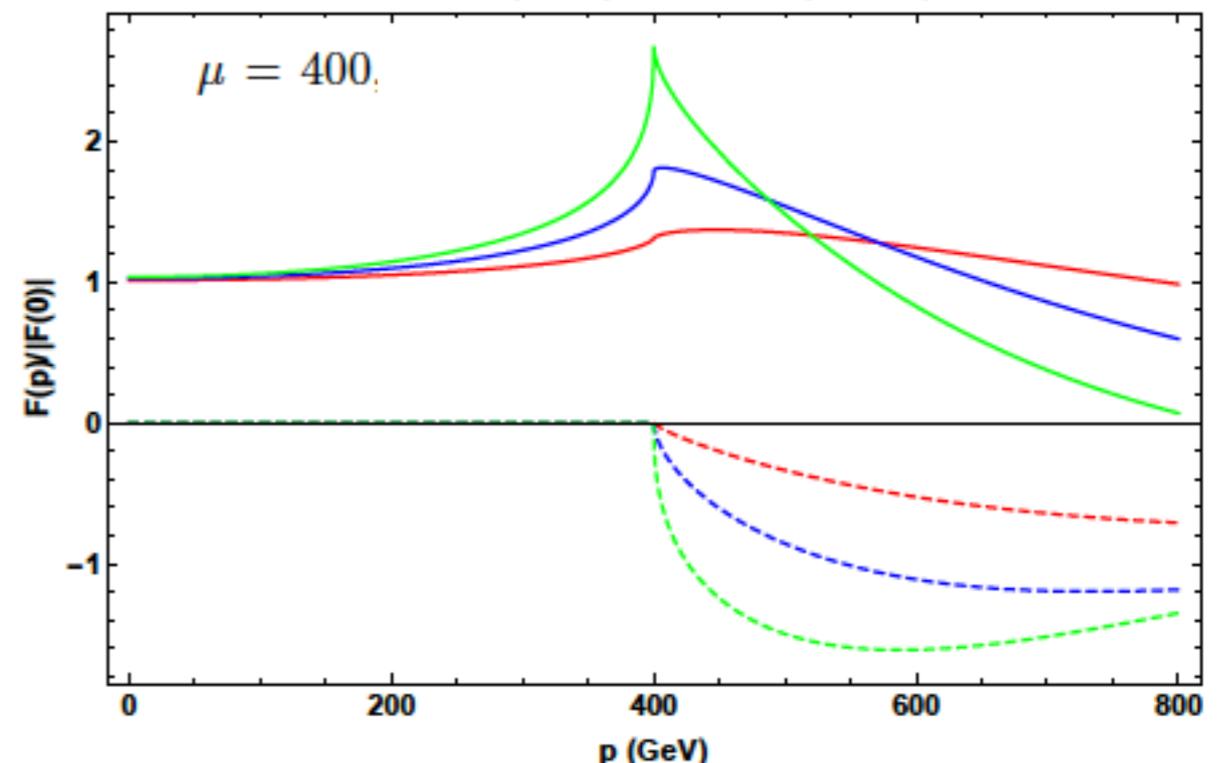


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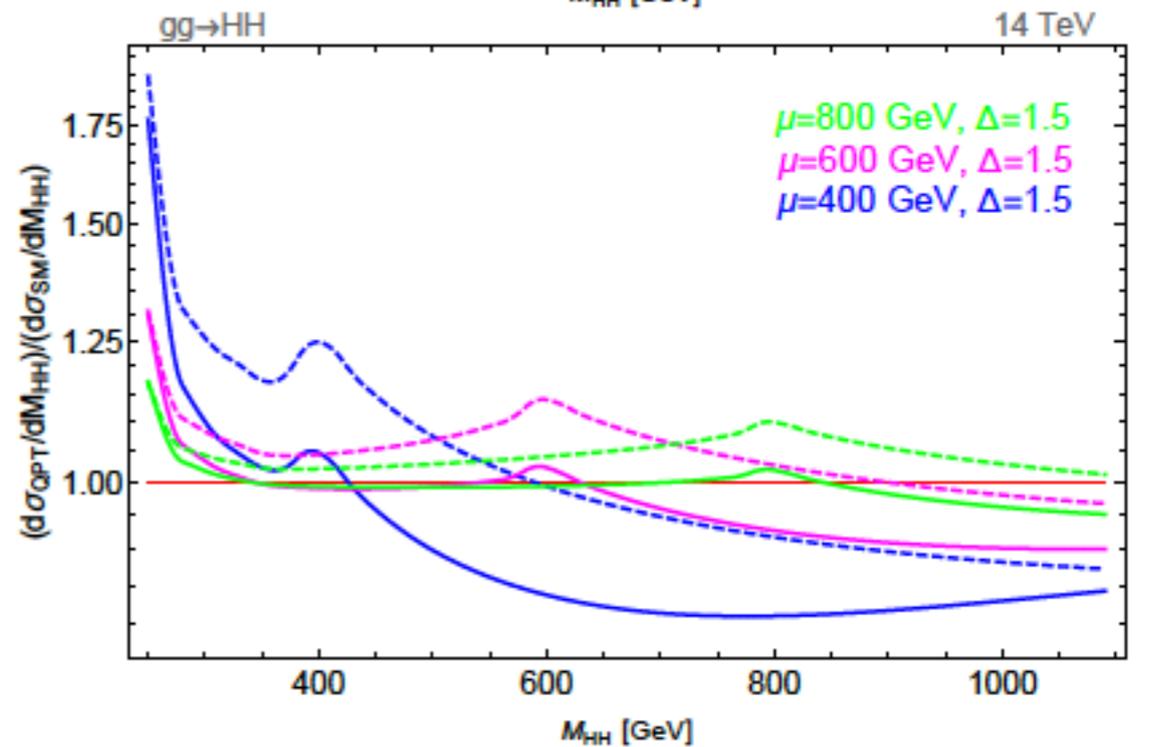
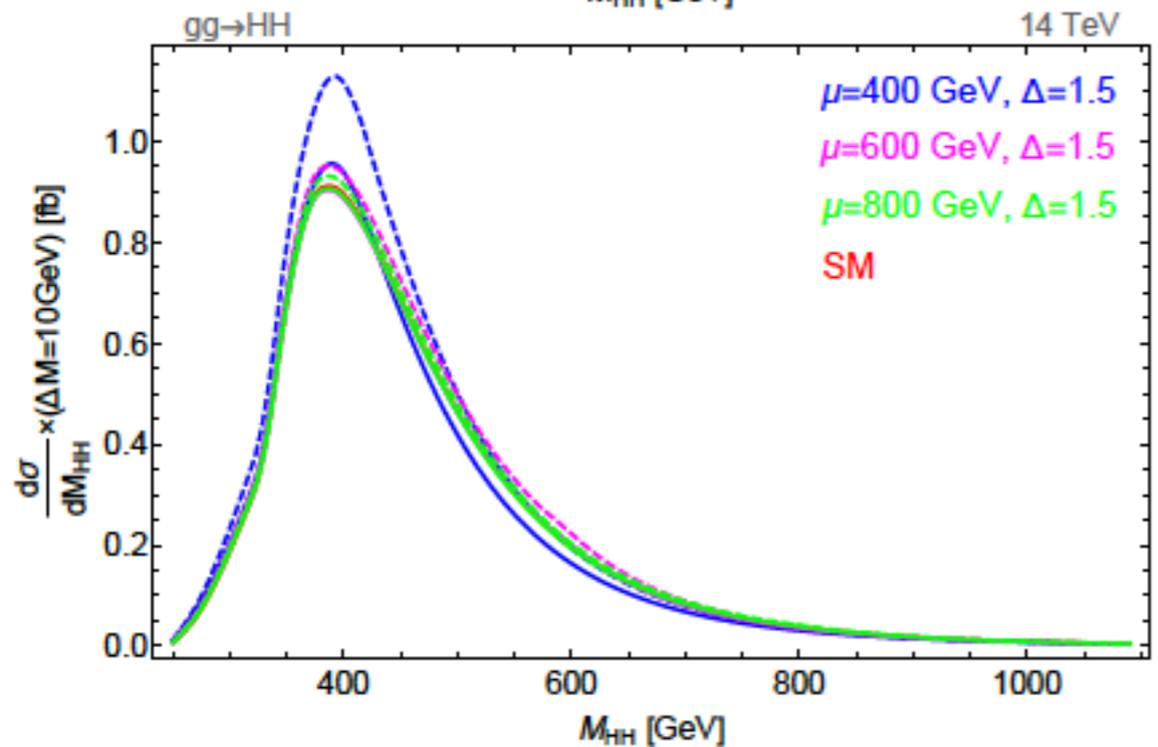
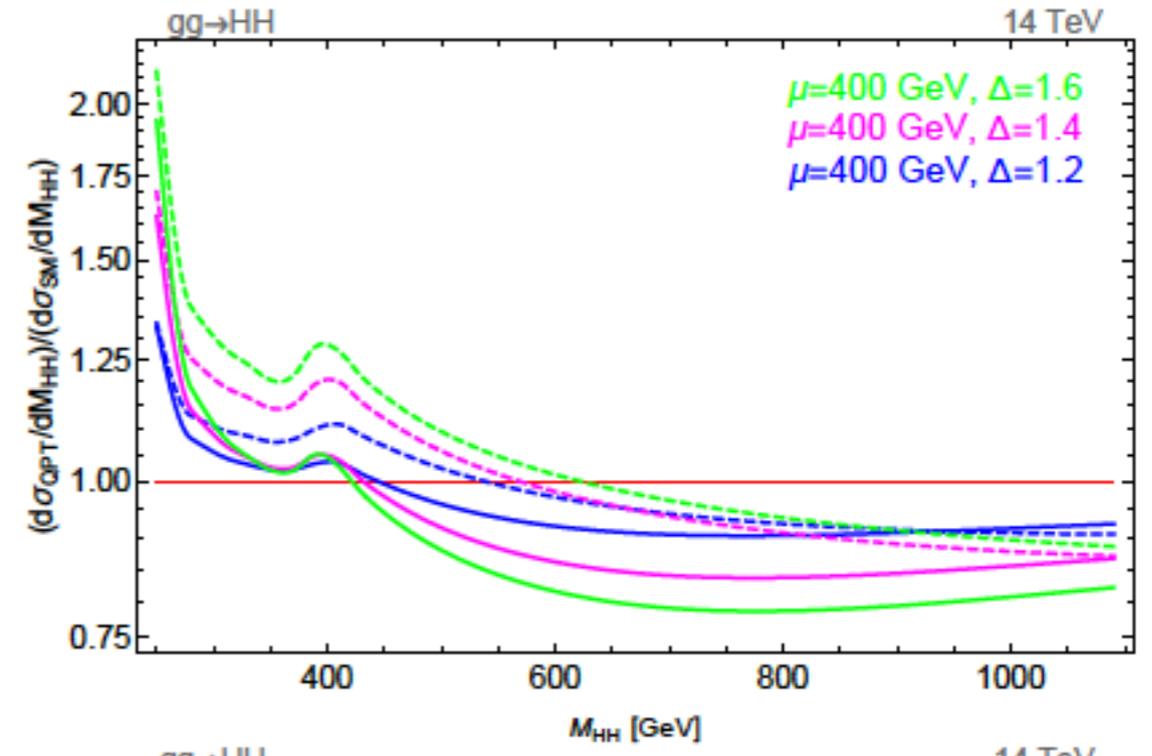
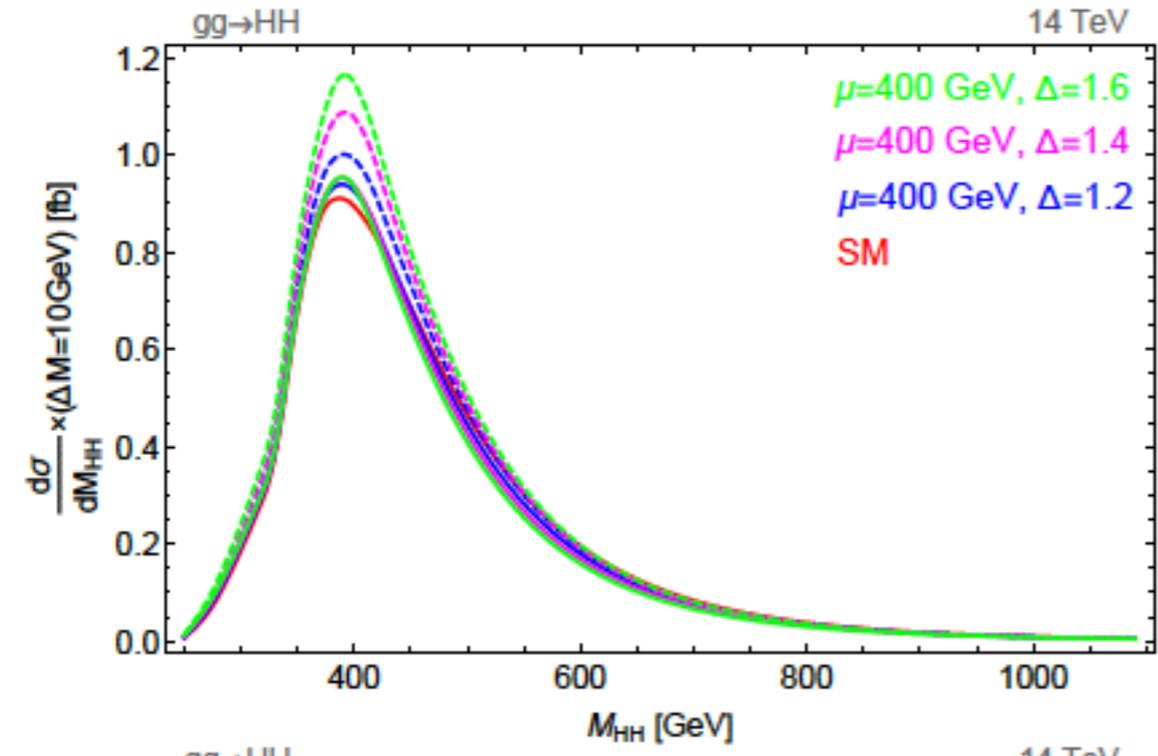
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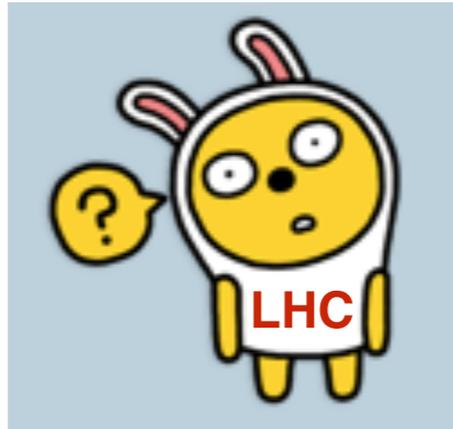
◆ Double Higgs production

dashed lines correspond to the case where only the Higgs two-point function has non-trivial behavior inherited from a sector with strong dynamics.



Summary: Higgs Potential

$$V(\tilde{h}) \stackrel{?}{=} \tilde{v}^{2\Delta} \tilde{m}^{4-2\Delta} \sum_{n=0}^{\infty} \frac{c_n}{n!} \left(\frac{\tilde{h}}{\tilde{v}^{\Delta}} \right)^n$$



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- ◆ We need FCC to really understand the EWSB!