

New Angles on Groomed Observables

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Based on upcoming works with Lina Necib and Jesse Thaler
and with Andrew Larkoski and Duff Neill

Motivation

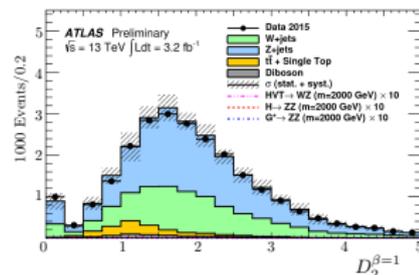
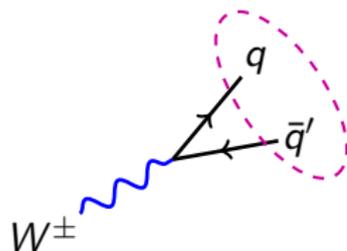
- A major focus of jet substructure has been on observables for boosted object tagging:

$$\tau_{2,1}^{(\beta)}, \tau_{3,2}^{(\beta)}, C_2^{(\beta)}, D_2^{(\beta)}, \dots$$

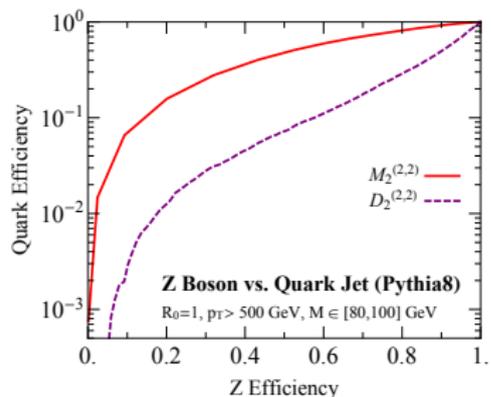
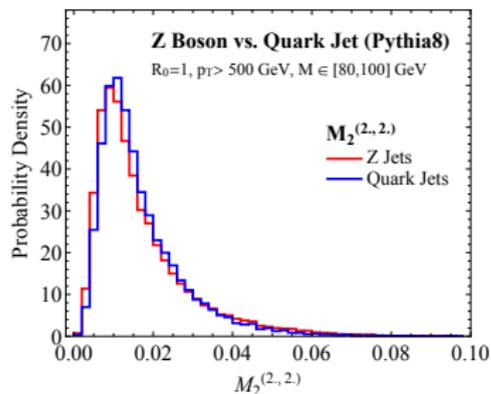
- Large push towards groomed observables

$$\begin{aligned}\tau_{2,1}^{(\beta)} &\rightarrow \tau_{2,1}^{(\beta),\text{groomed}} \\ \tau_{3,2}^{(\beta)} &\rightarrow \tau_{3,2}^{(\beta),\text{groomed}} \\ D_2^{(\beta)} &\rightarrow D_2^{(\beta),\text{groomed}}\end{aligned}$$

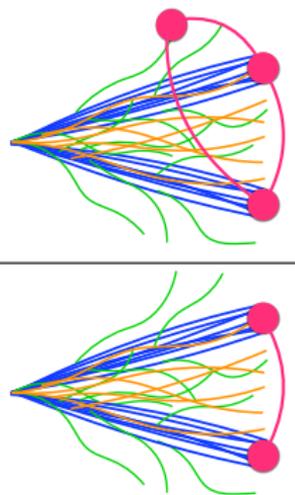
- For ungroomed substructure, many improvements came from more detailed theoretical understanding, and explicit calculations.
 - Recoil free observables, ECFs, WTA axes, $D_2^{(\beta)}$, ...



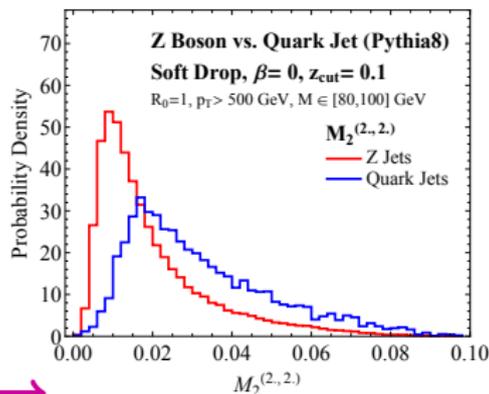
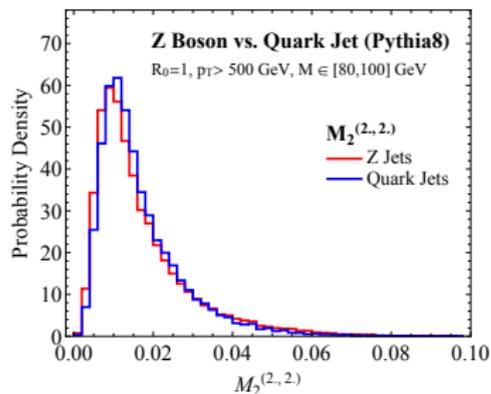
An Interesting Example: M_2



$$M_2^{(\beta)} = \frac{1 e_3^{(\beta)}}{e_2^{(\beta)}} =$$

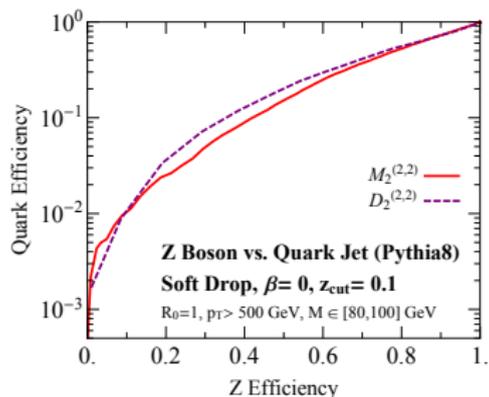
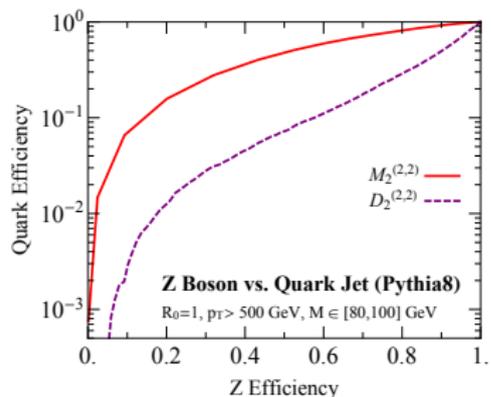


An Interesting Example: M_2

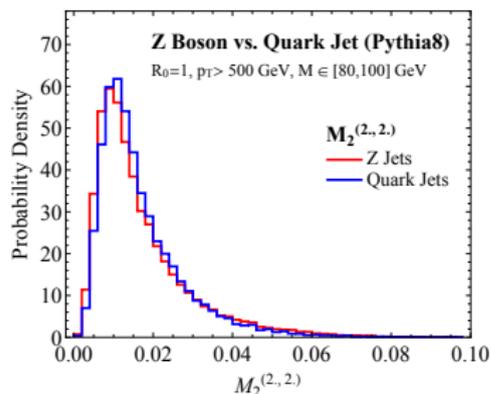


Soft Drop

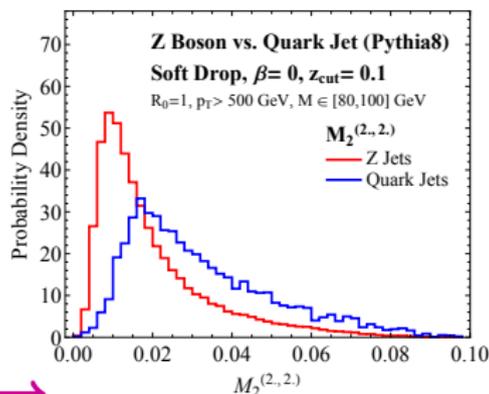
Grooming



An Interesting Example: M_2



Soft Drop
Grooming

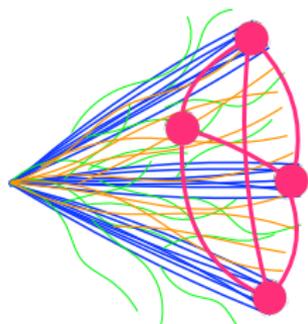


- Grooming has highly non-trivial impact on radiation pattern!
- Are the observables we are currently using optimal?
- It is essential to understand properties of groomed multi-prong observables from first principles QCD.

For progress for groomed masses, see: [Dasgupta, Fregoso, Marzani, Salam], Andrew Larkoski's Talk Later Today

Outline

- New Observables for Groomed Jet Substructure
 - Generalizing the energy correlation functions
 - The N_i series of observables and M_2
- Analytic Boosted Boson Discrimination with Grooming
 - Factorization for groomed multi-prong observables
 - Universality of groomed D_2 distribution



Collinear Soft Subjet

$$e_3^{(\alpha)} \ll (e_2^{(\alpha)})^3$$

$J_{n_{\alpha j}}$

$S_{n_a n_b n_c}^+$

jet axis

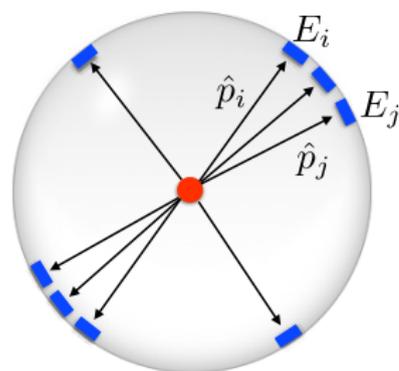
J_n

R

New Observables for Groomed Jet Substructure

Back to Basics

- What is an observable?



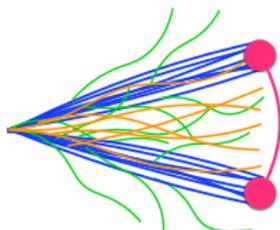
$$F_N(P) = \sum E_{i_1} \cdots E_{i_N} f_N(\hat{p}_{i_1}, \cdots, \hat{p}_{i_N})$$

- Linear in the energies by IRC safety.
- f_N is symmetric, and $f_N \rightarrow 0$ if $\hat{p}_i || \hat{p}_j$
- Known that from this one can reconstruct any IRC safe observable in the QFT (i.e. the energy momentum tensor). [Tkachov]
- Is this useful for jet substructure?

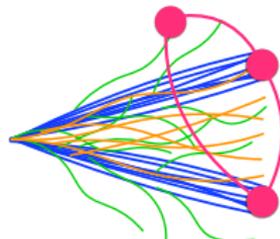
Example: Energy Correlation Functions

- Energy Correlation Functions: $f_N(\hat{p}_{i_1}, \dots, \hat{p}_{i_N}) = \prod R_{ij}^\beta$
[Larkoski, Salam, Thaler]

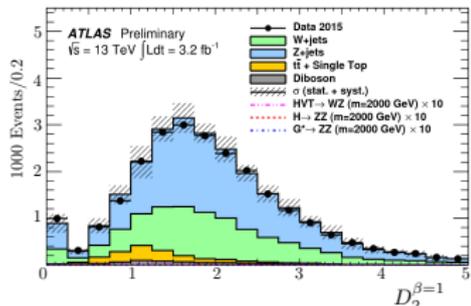
$$e_2^{(\beta)} = \frac{1}{p_{TJ}^2} \sum_{i < j \in J} p_{Ti} p_{Tj} (R_{ij})^\beta$$



$$e_3^{(\beta)} = \frac{1}{p_{TJ}^3} \sum_{i < j < k \in J} p_{Ti} p_{Tj} p_{Tk} (R_{ij} R_{jk} R_{ik})^\beta$$



- Powerful discriminant $D_2^{(\beta)} = \frac{e_3^{(\beta)}}{(e_2^{(\beta)})^3}$
[Larkoski, Moult, Neill]
- Rigid structure.



New Angles on the Energy Correlation Functions

- To tailor jet substructure observables, need more freedom.

$$F_N(P) = \sum E_{i_1} \cdots E_{i_N} f_N(\hat{p}_{i_1}, \cdots, \hat{p}_{i_N})$$

- Energy dependence of general IRC safe observable is fixed.
- Angular dependence, $f_N(\hat{p}_{i_1}, \cdots, \hat{p}_{i_N})$, more flexible.
- A convenient set of functions for identifying hierarchical structures is

$$f_N(\hat{p}_{i_1}, \cdots, \hat{p}_{i_N}) = \min \left(\prod_{s,t}^i R_{st}^\beta \right)$$

General Energy Correlation Functions

$$i e_j^{(\beta)} = \frac{1}{P_{TJ}^j} \sum_{1 \leq n_1 < \cdots < n_j \leq n} P_{Tn_1} P_{Tn_2} \cdots P_{Tn_j} \min \left(\prod_{s,t}^i R_{st}^\beta \right)$$

Generalized Energy Correlation Functions

General Energy Correlation Functions

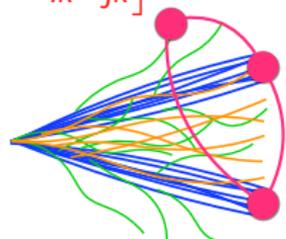
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- Example: Three different ways to probe three particle correlations.

$$1 e_3^{(\beta)} = \frac{1}{p_{TJ}^3} \sum_{1 \leq i < j < k \leq n_j} p_{Ti} p_{Tj} p_{Tk} \min \left[R_{ij}^\beta, R_{ik}^\beta, R_{jk}^\beta \right],$$

$$2 e_3^{(\beta)} = \frac{1}{p_{TJ}^3} \sum_{1 \leq i < j < k \leq n_j} p_{Ti} p_{Tj} p_{Tk} \min \left[R_{ij}^\beta R_{ik}^\beta, R_{ij}^\beta R_{jk}^\beta, R_{ik}^\beta R_{jk}^\beta \right],$$

$$3 e_3^{(\beta)} = \frac{1}{p_{TJ}^3} \sum_{1 \leq i < j < k \leq n_j} p_{Ti} p_{Tj} p_{Tk} R_{ij}^\beta R_{ik}^\beta R_{jk}^\beta = e_3^{(\beta)}$$



- Flexible basis for substructure observables.

Applications

- Observables for particular applications constructed using “power counting”.

[Larkoski, Moult, Neill], Larkoski Boost 2014, Moult Boost 2015

⇒ Improving boosted top tagging:

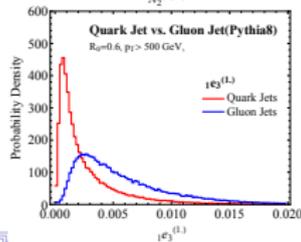
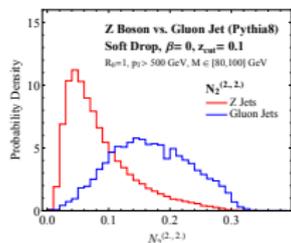
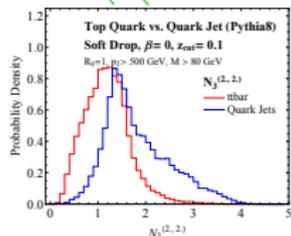
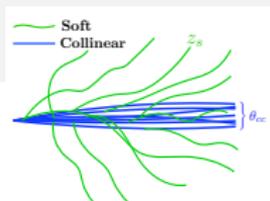
$$N_3^{(\beta)} = \frac{2e_4^{(\beta)}}{(1e_3^{(\beta)})^2}$$

⇒ New observables for $W/Z/H$ tagging:

$$M_2^{(\beta)} = \frac{1e_3^{(\beta)}}{e_2^{(\beta)}}, \quad N_2^{(\beta)} = \frac{2e_3^{(\beta)}}{(e_2^{(\beta)})^2}, \quad D_2^{(\alpha,\beta)} = \frac{3e_3^{(\alpha)}}{(e_2^{(\beta)})^{3\alpha/\beta}}$$

⇒ Improving Quark vs. Gluon discrimination:

$$1e_3^{(\beta)}$$



fjcontrib: Energy correlator tag 1.2.0-alpha0

Boosted Top Tagging with N_3

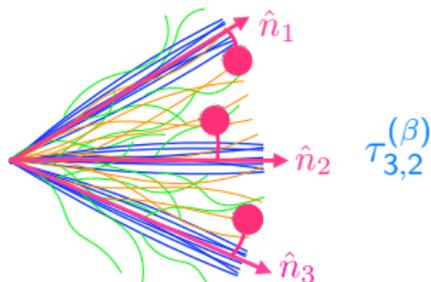
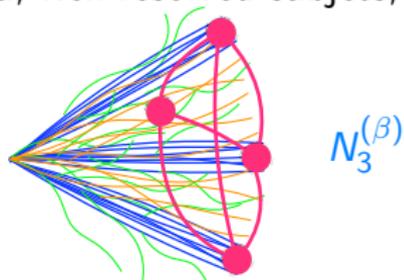
- Appropriate observable for boosted top tagging

$$N_3^{(\beta)} = \frac{2e_4^{(\beta)}}{(1e_3^{(\beta)})^2}$$

- It can be shown that in the limit of groomed, well resolved subjects, four point correlations factorize as

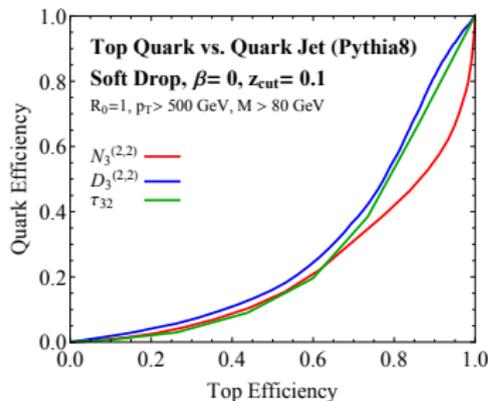
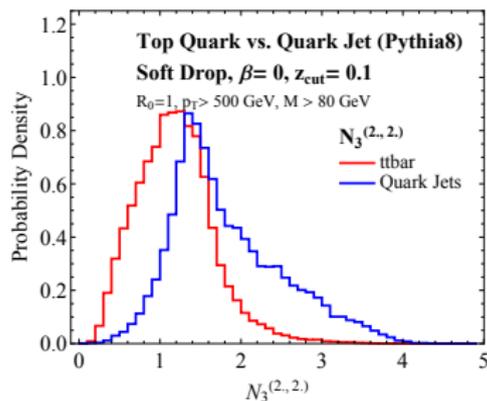
$$N_3^{(\beta)} \sim \frac{\tau_3^{(\beta)} \cdot \tau_2^{(\beta)}}{(\tau_2^{(\beta)})^2} = \frac{\tau_3^{(\beta)}}{\tau_2^{(\beta)}}$$

- N_3 is defined as a sum of four particle correlations.
- No axes, reclustering into subjects, etc.
- Smooth transition from resolved to unresolved regime.



Boosted Top Tagging with N_3

- $\tau_N^{(\beta)}$ behave well in the limit of well resolved substructure, $\tau_N^{(\beta)} \ll 1$.
- Can have pathological or undesired behavior in the $\tau_3^{(\beta)} / \tau_2^{(\beta)} \rightarrow 1$ limit. Precisely where used for high efficiency taggers!
- N_3 well behaved at high efficiency. Leads to significantly improved performance.



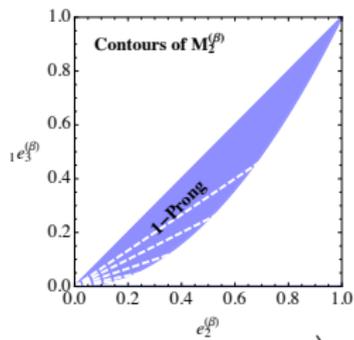
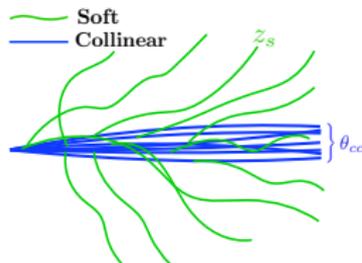
New Observables for Two-Prong Tagging

- Also allows for new two prong observables.
- Interestingly, powerful observables can be constructed from each of the $i e_3^{(\beta)}$, $i = 1, 2, 3$.
- Each probes different angular correlations
 \implies Very different behavior under grooming.
- Power counting uniquely fixes the structure of the observables:

$$M_2^{(\beta)} = \frac{1 e_3^{(\beta)}}{e_2^{(\beta)}}, \quad N_2^{(\beta)} = \frac{2 e_3^{(\beta)}}{(e_2^{(\beta)})^2}, \quad D_2^{(\alpha, \beta)} = \frac{3 e_3^{(\alpha)}}{(e_2^{(\beta)})^{3\alpha/\beta}}$$

The Power of Power Counting

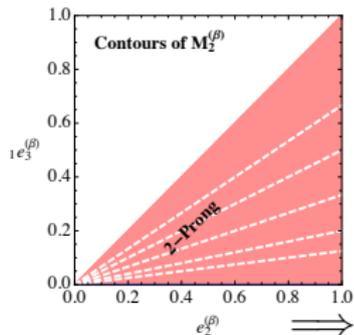
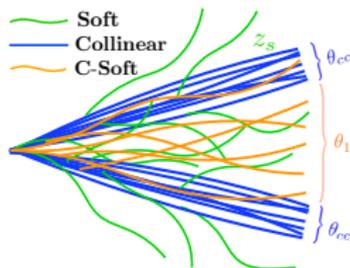
- Power counting allows us to understand $M_2^{(\beta)} = \frac{1e_3^{(\beta)}}{e_2^{(\beta)}}$:



$$e_2^{(\beta)} \sim \theta_{cc}^\beta + z_s,$$

$$1e_3^{(\beta)} \sim \theta_{cc}^\beta + z_s^2$$

$$\Rightarrow \text{1-prong: } (e_2^{(\beta)})^2 \lesssim 1e_3^{(\beta)} \lesssim (e_2^{(\beta)})^2$$



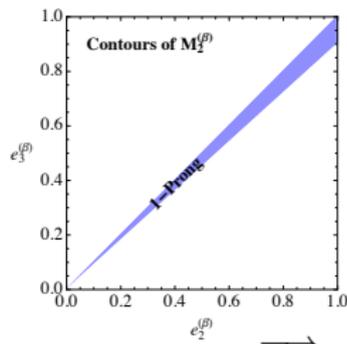
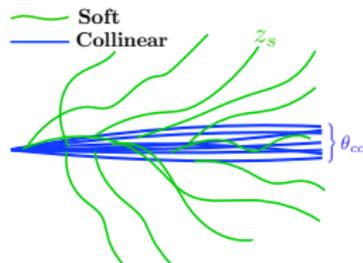
$$e_2^{(\beta)} \sim \theta_{12}^\beta,$$

$$1e_3^{(\beta)} \sim \theta_{12}^\beta z_s + \theta_{cc}^\beta + \theta_{12}^\beta z_{cs}$$

$$\Rightarrow \text{2-prong: } 0 < 1e_3^{(\beta)} \ll e_2^{(\beta)}$$

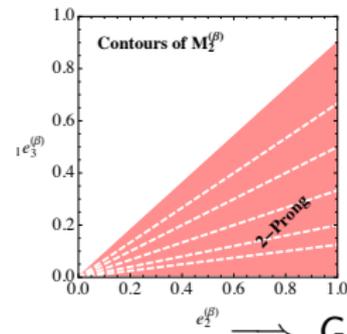
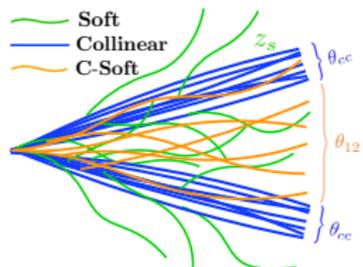
The Power of Power Counting

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$$e_2^{(\beta)} \sim \theta_{cc}^\beta + z_s,$$

$$1e_3^{(\beta)} \sim \theta_{cc}^\beta + z_s$$



$$e_2^{(\beta)} \sim \theta_{12}^\beta,$$

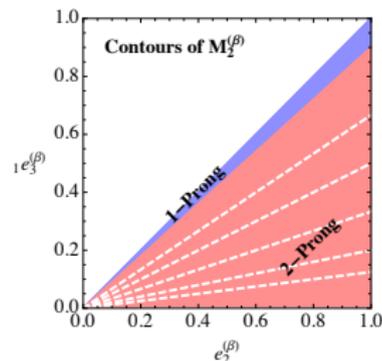
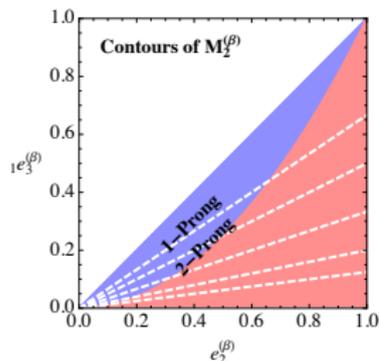
$$1e_3^{(\beta)} \sim \theta_{12}^\beta z_s + \theta_{cc}^\beta + \theta_{12}^\beta z_s$$

$e_2^{(\beta)} \Rightarrow$ Groomed 1-prong: $e_2^{(\beta)} \sim 1e_3^{(\beta)} !!$

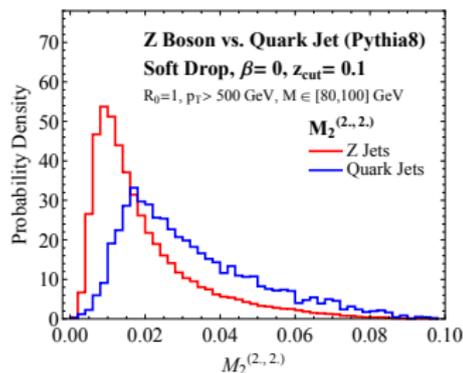
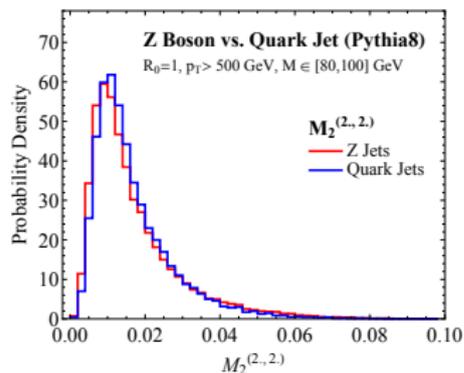
$e_2^{(\beta)} \Rightarrow$ Groomed 2-prong: $0 < 1e_3^{(\beta)} \ll e_2^{(\beta)}$

Taking Advantage of Grooming

- Grooming “cleans” phase space, allowing for discrimination.

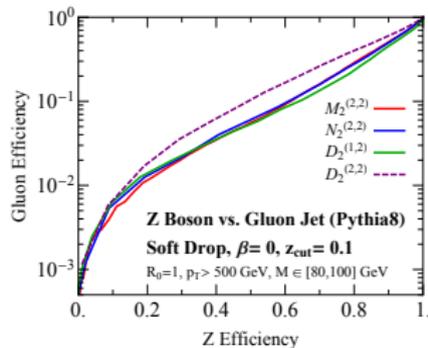


Soft Drop
Grooming

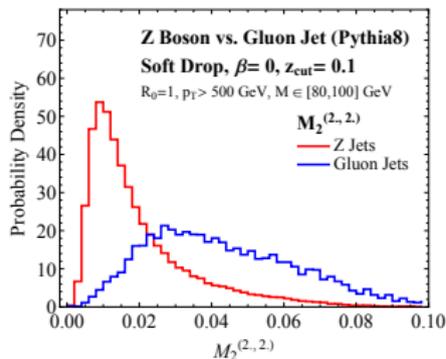


Taking Advantage of Grooming

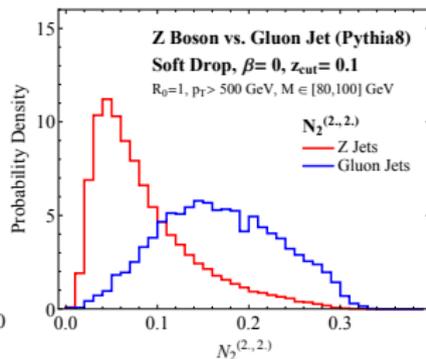
- Similar power counting used to study other observables.
- Three interesting observables for two prong discrimination on (un)groomed jets.
- Each outperforms $D_2^{(2,2)}$.
- Would be interesting to study correlations.



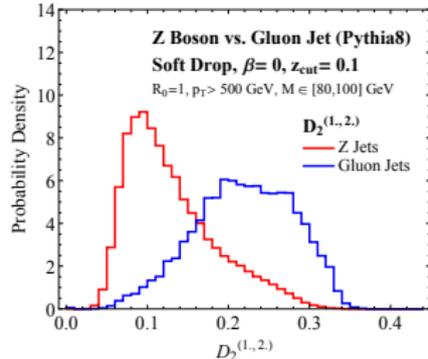
$$M_2^{(\beta)}$$



$$N_2^{(\beta)}$$



$$D_2^{(1,2)}$$



Analytic Boosted Boson Discrimination with Grooming

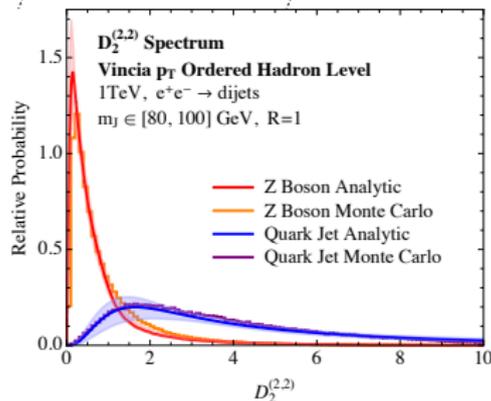
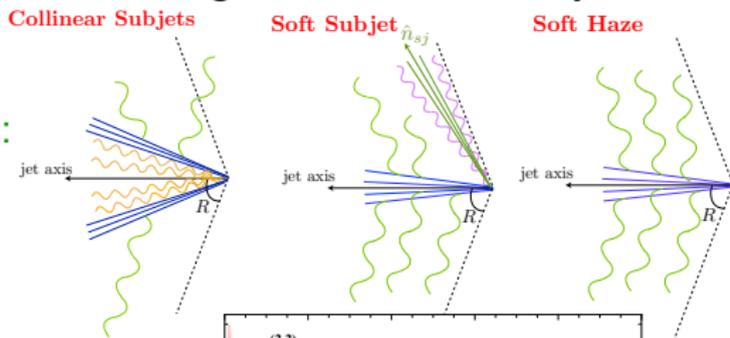
Analytic Boosted Boson Discrimination

[Larkoski, Moulton, Neill]

- Calculation of $D_2 = \frac{e_3^{(\beta)}}{(e_2^{(\beta)})^3}$ in e^+e^- using effective field theory.

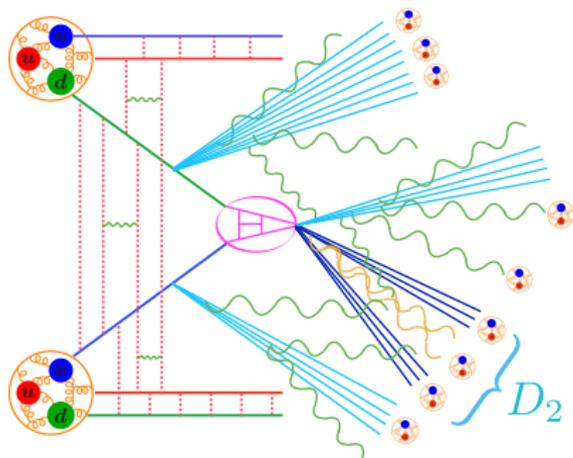
EFTs for 2-prong Substructure:

- All orders factorization theorems ✓
- Treatment of non-perturbative physics ✓
- Resummation to first non-trivial order (NLL) ✓
- Treatment of $D_2 < 1$ and $D_2 > 1$ ✓



Analytic Boosted Boson Discrimination

- Difficulties in extending to pp :



- Global color correlations
- Hadronization corrections
- Pile-Up
- Underlying event*

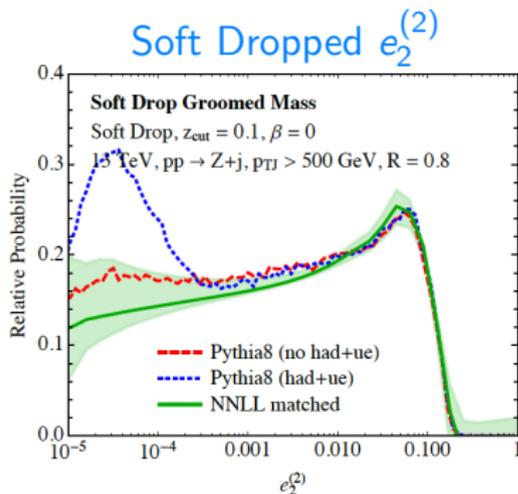
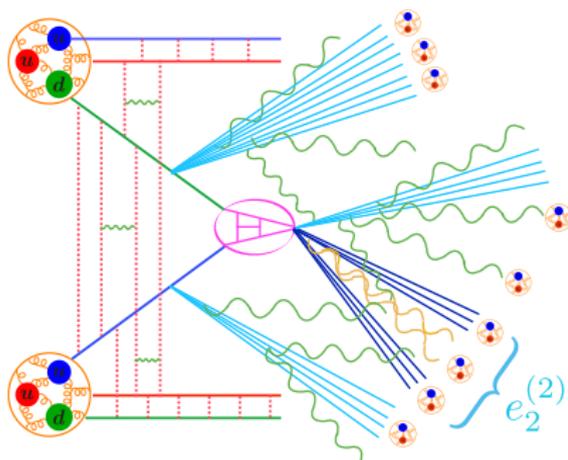
- All complications associated with soft radiation.
⇒ Groomers can make calculations simpler and more universal.

Factorization for Soft Dropped D_2

[Larkoski, Marzani, Soyez, Thaler]

- Focus on Soft Drop groomer. Involves a single additional scale z_{cut} .
- Consider first Soft Dropped $e_2^{(2)} \sim \frac{m_J^2}{p_{TJ}^2}$.

[See Andrew Larkoski's Talk Later Today]

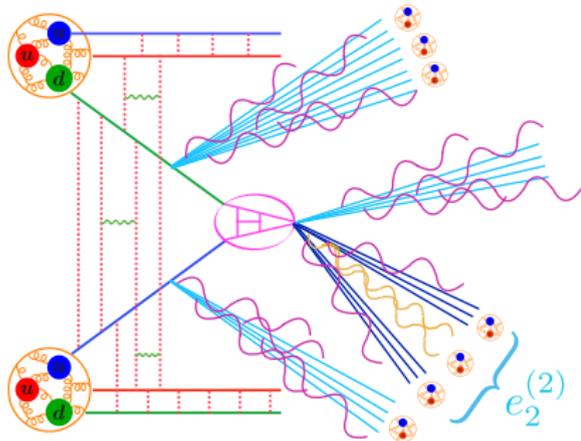


[Frye, Larkoski, Schwartz, Yan]

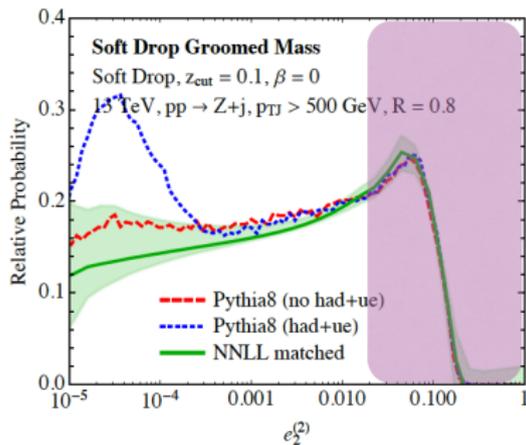
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- IF $e_2^{(2)} \ll z_{\text{cut}}$, soft radiation cannot contribute. Factorization formula in terms of collinear physics.



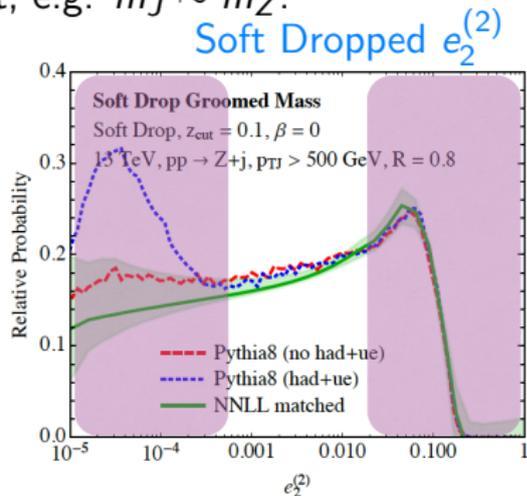
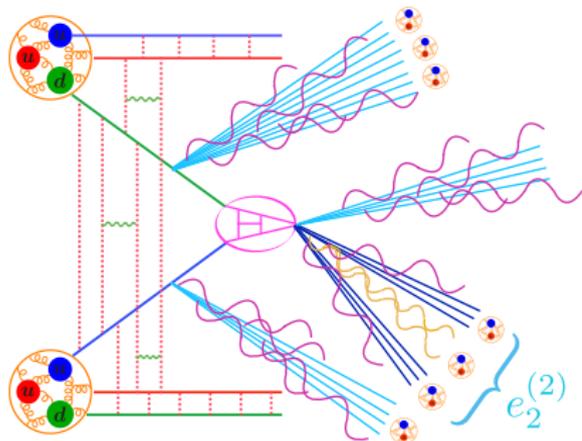
Soft Dropped $e_2^{(2)}$



Factorization for Soft Dropped D_2

[Larkoski, Marzani, Soyez, Thaler]

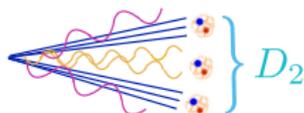
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- IF $e_2^{(2)} \ll z_{\text{cut}}$, soft radiation cannot contribute. Factorization formula in terms of collinear physics.
- For substructure, apply a mass cut, e.g. $m_J \sim m_Z$.



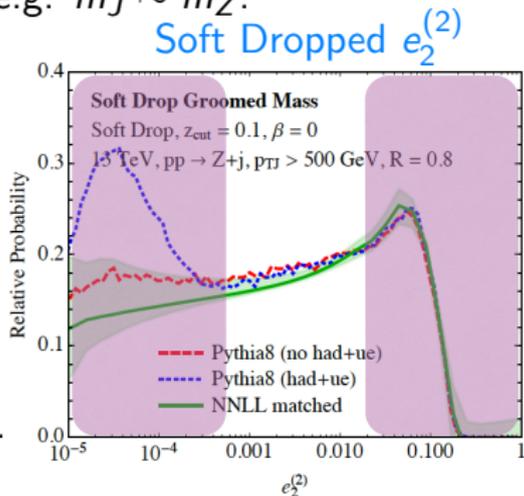
Factorization for Soft Dropped D_2

[Larkoski, Marzani, Soyez, Thaler]

- Focus on Soft Drop groomer. Involves a single additional scale z_{cut} .
- Consider first Soft Dropped $e_2^{(2)} \sim \frac{m_J^2}{p_{TJ}^2}$. [See Andrew Larkoski's Talk Later Today]
- IF $e_2^{(2)} \ll z_{\text{cut}}$, soft radiation cannot contribute. Factorization formula in terms of collinear physics.
- For substructure, apply a mass cut, e.g. $m_J \sim m_Z$.
- Jet decouples!
Can now measure D_2 .

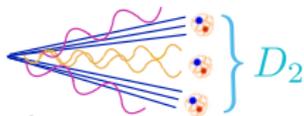


- For $z_{\text{cut}} = 0.1$, need $p_{TJ} \gtrsim 500$ GeV.



Factorization for Soft Dropped D_2 : Universality

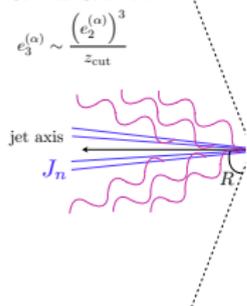
- With a mass cut jet decouples. Can measure its substructure.
 \implies Soft dropped D_2 is probably process independent!



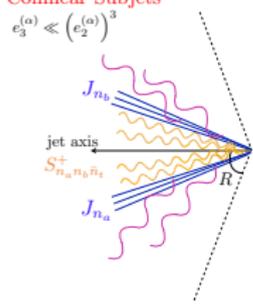
- Possible dependence only on jet properties p_{TJ} , m_J , flavor...
- All orders (partonic) factorization formula can be proven using EFT.
- Allows for resummation using renormalization group.

EFTs for Groomed 2-prong Substructure:

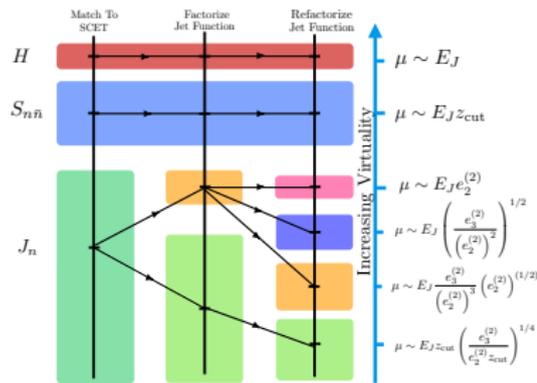
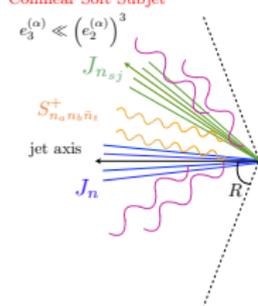
Collinear Soft Haze



Collinear Subjets



Collinear Soft Subjet

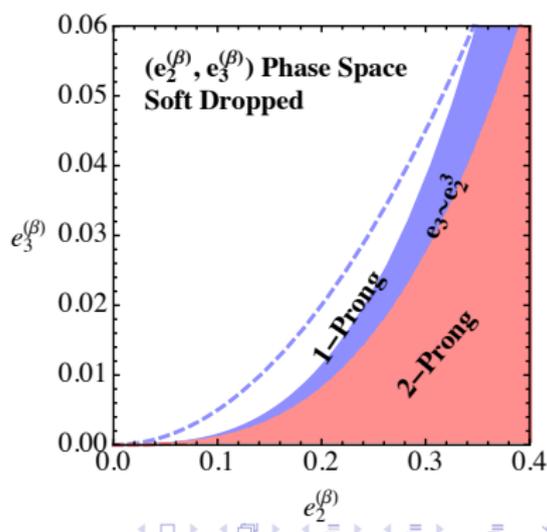
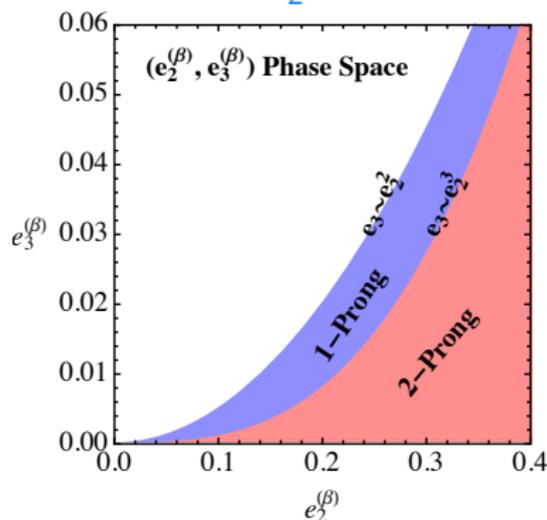


Stability

- Soft Drop modifies the $D_2 = \frac{e_3^{(\beta)}}{(e_2^{(\beta)})^3}$ phase space.

$$D_2^{\max} \sim \frac{1}{2e_2^{(\alpha)}} \sim \frac{1}{2} \frac{p_{TJ}^2}{m_J^2}$$

$$D_2^{\max, \text{softdrop}} \sim \frac{1}{2Z_{\text{cut}}}$$

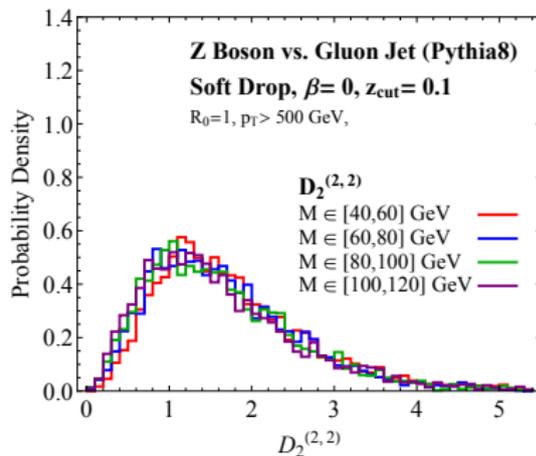
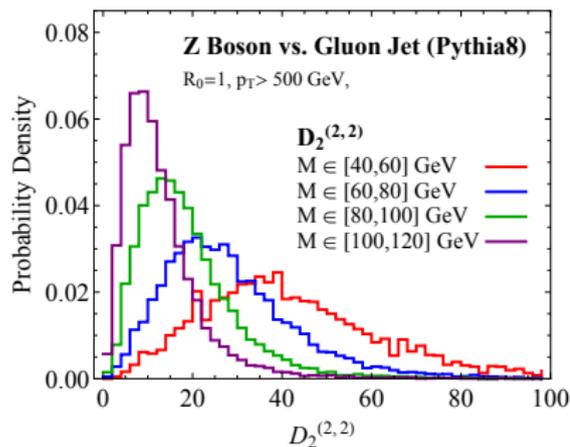


Stability

- Soft dropped D_2 distribution largely independent of
 - jet mass

$$D_2^{\max} \sim \frac{1}{2} \frac{p_{TJ}^2}{m_j^2}$$

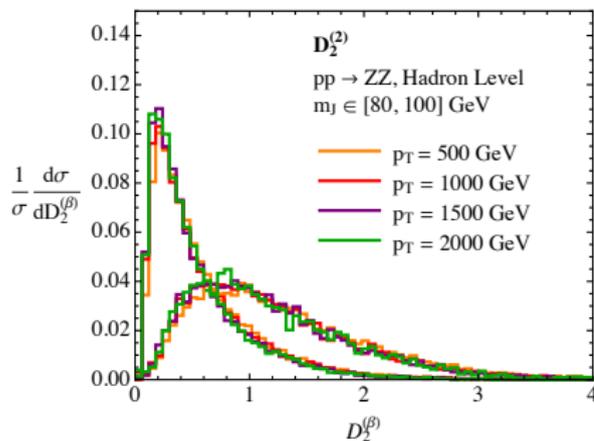
$$D_2^{\max, \text{softdrop}} \sim \frac{1}{2z_{\text{cut}}}$$



Stability

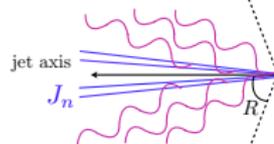
- Soft dropped D_2 distribution largely independent of
 - jet mass
 - jet p_T

$$D_2^{\text{max,softdrop}} \sim \frac{1}{2z_{\text{cut}}}$$



Collinear Soft Haze

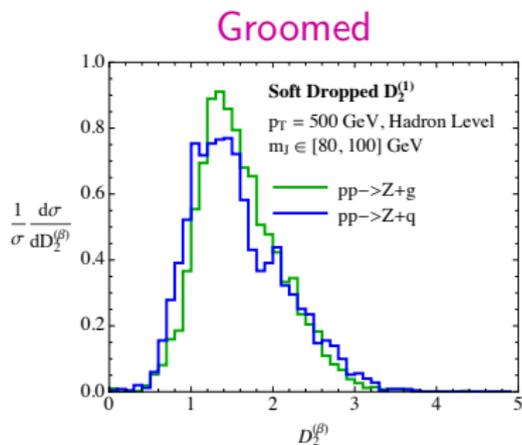
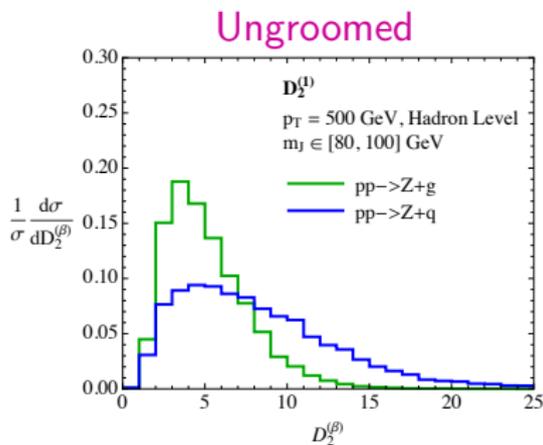
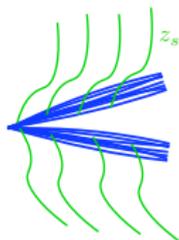
$$e_3^{(\alpha)} \sim \frac{(e_2^{(\alpha)})^3}{z_{\text{cut}}}$$



- Extremely stable discriminant!
- Expansion about small values insufficient to describe endpoint.

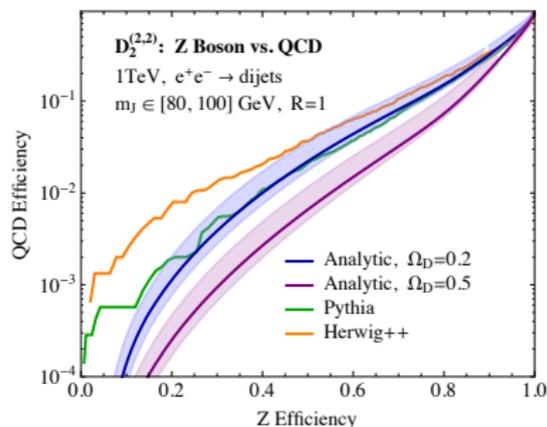
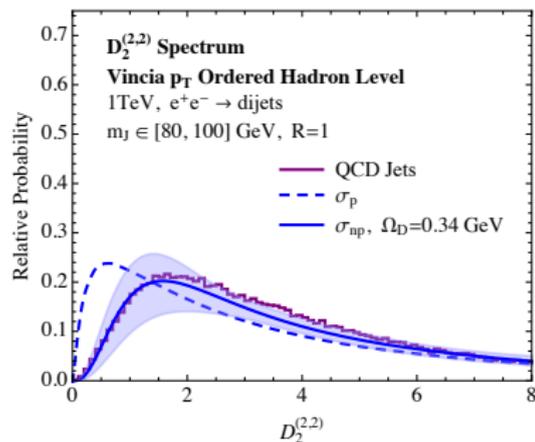
Quarks and Gluons

- Dominant dependence of D_2 on parton flavor carried by soft radiation.
- Enters ungroomed calculation as $C_i \log(z_s)$.
- Replaced in groomed calculation by $C_i \log(z_{\text{cut}})$.
⇒ Doesn't enter normalized distribution!



Non-Perturbative Corrections

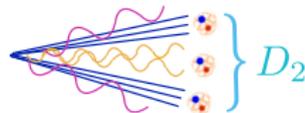
- Observables sensitive to low scales are sensitive to non-perturbative (hadronization) corrections.
- Can introduce non-perturbative parameters, (shift in first moment Ω_D) but leads to uncertainties.
- Discrimination power is particularly sensitive to hadronization.



Non-Perturbative Corrections

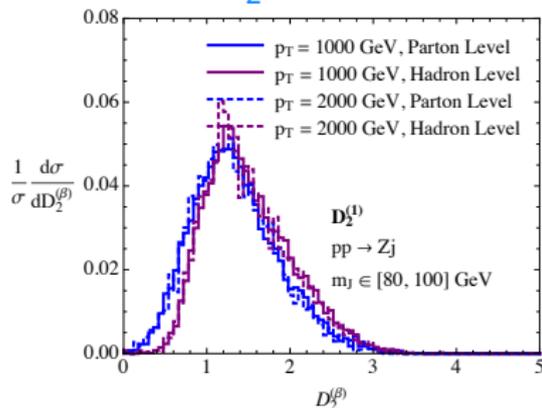
- Factorization is true to all orders, and allows us to address impact of hadronization.
- Can prove that non-perturbative corrections enter at the scale

$$D_2^{(\beta)} \simeq \left(\frac{1}{z_{\text{cut}}} \right)^{2-1/\beta} \left(\frac{\Lambda_{\text{QCD}}}{m_J} \right)$$

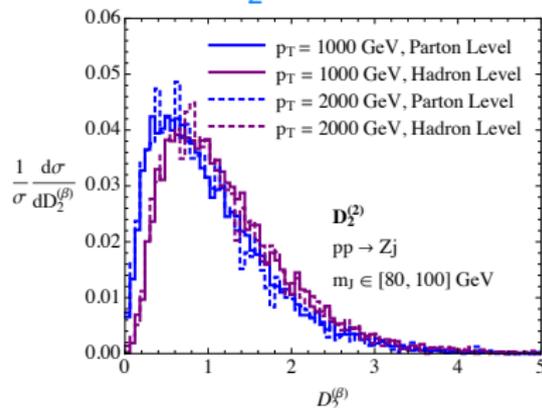


- Reduced hadronization corrections. Independent of Jet Energy !!!!

$$D_2^{(\beta)}, \beta = 1$$



$$D_2^{(\beta)}, \beta = 2$$



Conclusions

- Power counting techniques allow the design of substructure observables with specific properties.
- Introduced a versatile basis of substructure observables $e_j^{(\beta)}$, and the discriminants
 - M_2 , N_2 , D_2 for groomed 2-prong substructure.
 - N_3 for high efficiency top tagging.
- Gaining analytic control over soft dropped multi-prong observables. Have all orders factorization formulae.
- Soft dropped D_2 offers many advantages, that can be proven from first principles QCD.
- Numerical results for soft dropped D_2 at the LHC will be coming soon!

Thanks!

Backup

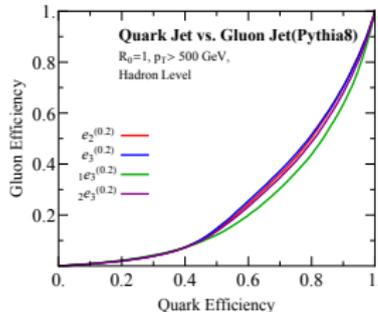
Quark vs. Gluon Discrimination with $1e_3^{(\beta)}$

- The observable $1e_3^{(\beta)}$ probes three particle correlations.
- Use of $1e_3^{(\beta)}$ as compared with $3e_3^{(\beta)}$ essential.

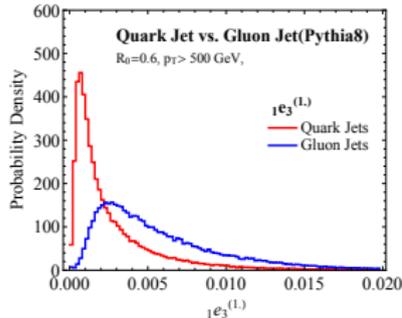
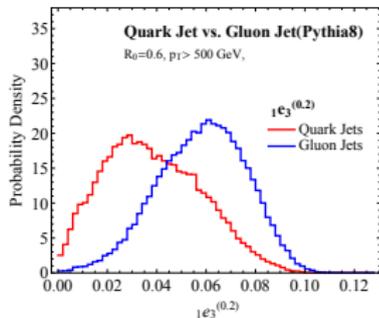
$$e_2^{(\beta)} \sim \theta_{cc}^\beta + z_s$$

$$e_3^{(\beta)} \sim (e_2^{(\beta)})^2$$

$$1e_3^{(\beta)} \sim \theta_{cc}^\beta + z_s^2$$



- Improved stability of discrimination power w.r.t. β .



D_2 with Grooming

- On groomed jets, $D_2^{(\alpha,\beta)}$ can be used with two distinct angular exponents.
- The choice $\alpha = 1$, $\beta = 2$ is good on groomed jets when a mass cut is used, as $\beta = 2$ corresponds to a mass in the denominator.
- The choice $\alpha = 1$, $\beta = 2$ does not work on ungroomed jets.
- Like M_2 , it relies on grooming to clean the phase space

$$D_2^{(\alpha,\beta)} = \frac{3e_3^{(\alpha)}}{(e_2^{(\beta)})^{3\alpha/\beta}}$$

Background:

$$e_3^{(1)} \sim \theta_{cc}^3 + z_s^2 + \theta_{cc} z_s$$

$$e_2^{(2)} \sim \theta_{cc}^2 + \theta_{cc}^2 z_s$$

$$\implies (e_2^{(2)})^2 \lesssim e_3^{(1)} \lesssim (e_2^{(2)})^{3/2}$$

Signal:

$$e_3^{(1)} \lesssim (e_2^{(2)})^{3/2}$$