



# Boosted $W/Z$ tagging at ATLAS

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On behalf of the ATLAS experiment

BOOST 2016

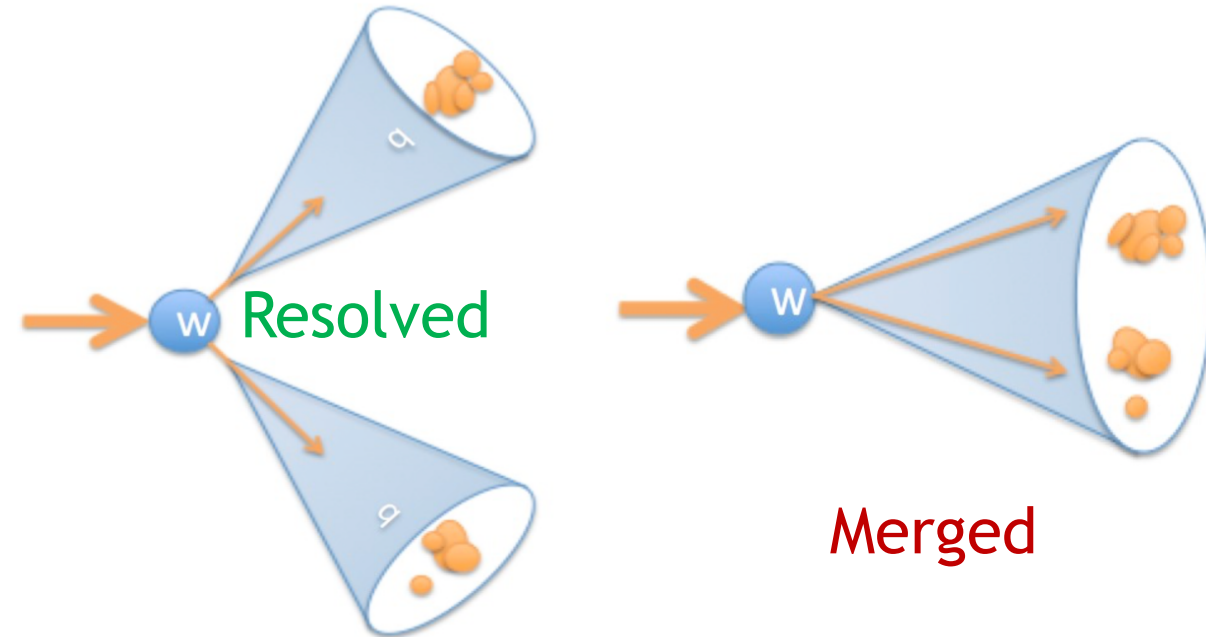
18<sup>th</sup> - 22<sup>nd</sup> July, 2016

- Hadron collider experiments deal with a large fraction of events from strongly produced jets - QCD multi-jet production
- XS of QCD multi-jet is much higher than W/Z signals  $\rightarrow$  leptonic decays are studied to reduce QCD multi-jet backgrounds
- Bosons similar to W/Z are predicted in BSM theories which don't couple to leptons

When boson momentum  $\gg$  mass, spatial proximity of decay products can help distinguish between QCD and W/Z-jets

- Many substructure algorithms have been studied in Run1

Tradeoff between using relatively pure leptonic decays of bosons and high-branching-ratio hadronic decays



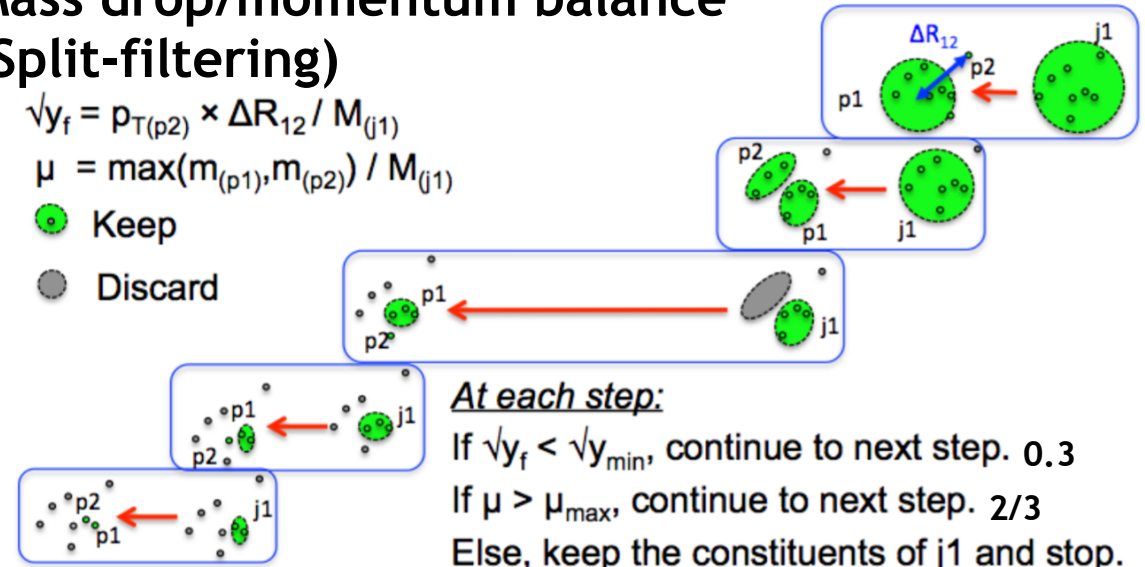
- Remove soft component targeting Pileup (PU) and Underlying event (UE)
- Better signal-background separation
- Improves resolution of signal jet mass peak

## Mass drop/momentum balance (Split-filtering)

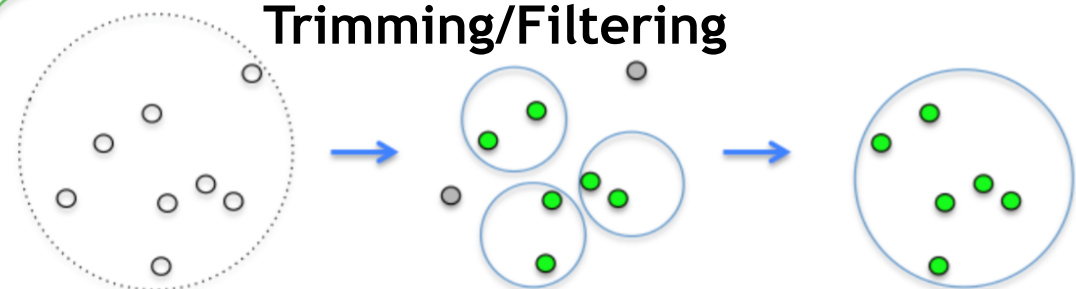
$$\sqrt{y_f} = p_{T(p2)} \times \Delta R_{12} / M_{(j1)}$$

$$\mu = \max(m_{(p1)}, m_{(p2)}) / M_{(j1)}$$

- Keep
- Discard

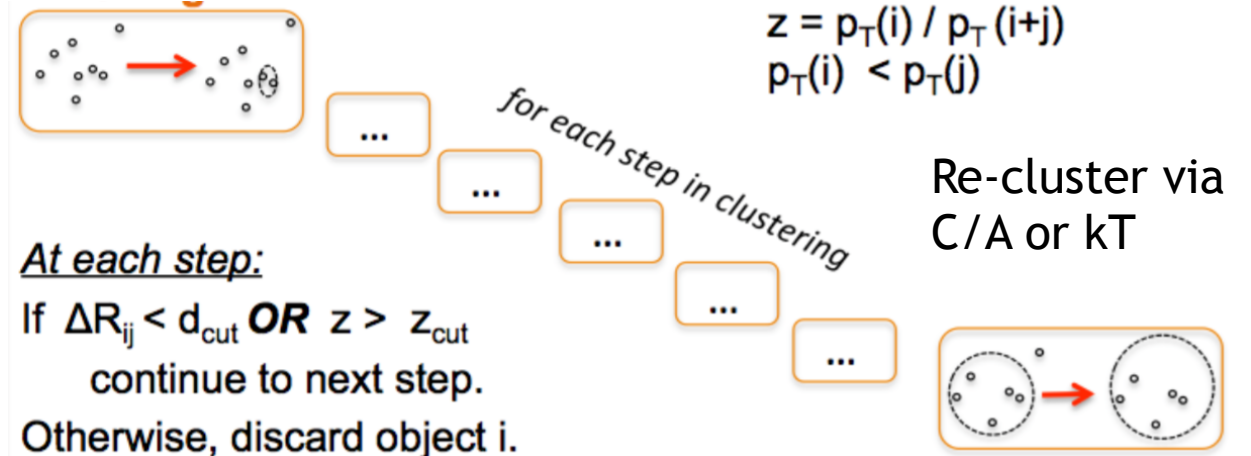


## Trimming/Filtering



**Type 1 (Trimming)** : If  $p_T(\text{subjet } i) / p_T(\text{jet}) < f_{cut}$  : discard subjet.  
**Type 2** : If  $N_{subjets} \leq N_{min}$  : discard jet.  
 Resulting jet is sum of subjets.

## Pruning

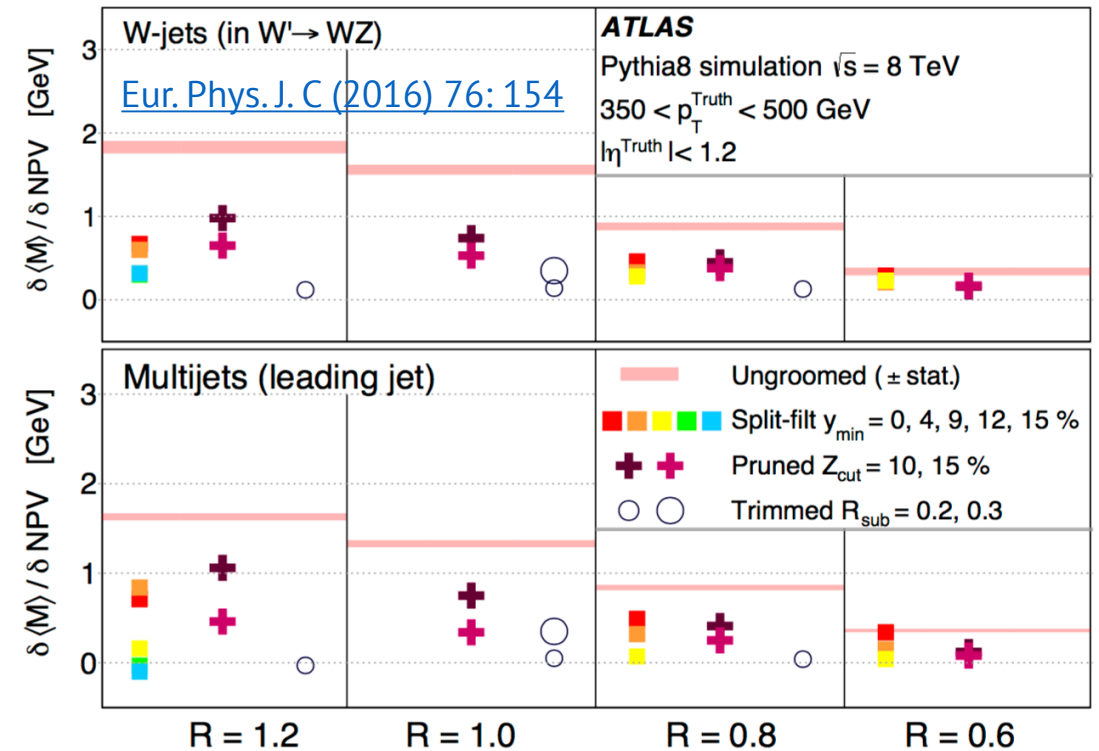




Grooming techniques mentioned on previous slide have been studied in great detail in Run 1 and the best combinations have been studied in Run 2

- Trimming : various large R jet values, different  $R_{\text{subject}}$  and  $f_{\text{cut}}$  values.
- Pruning : various large R jet values, different re-clustering algorithms,  $Z_{\text{cut}}$  and  $R_{\text{cut}}$  values.
- Split-filtering: various large R jet values,  $R_{\text{subject}}$ , mass drop fraction and  $y_{\text{cut}}$  values.

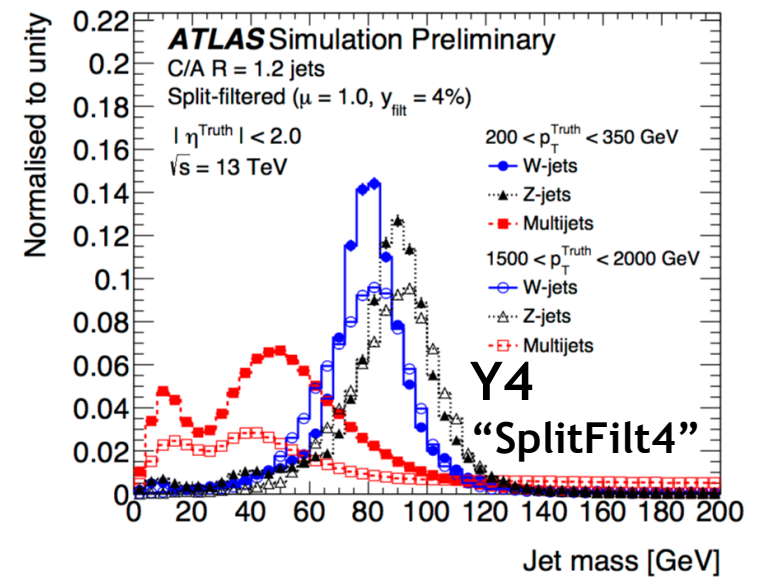
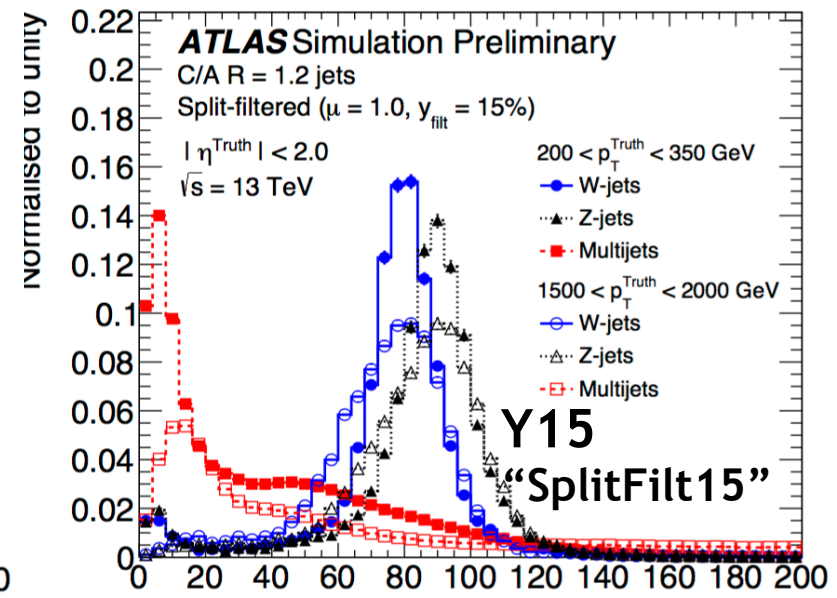
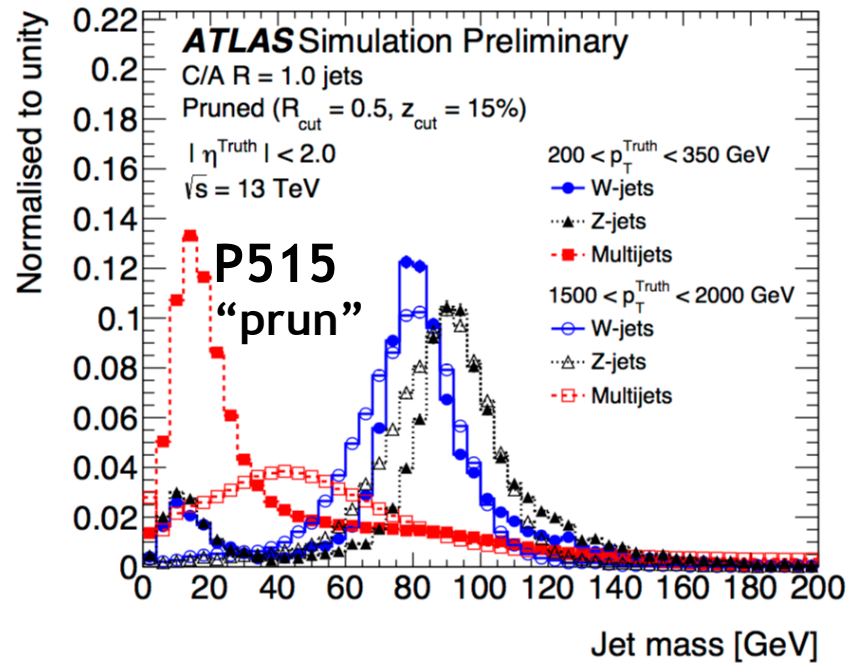
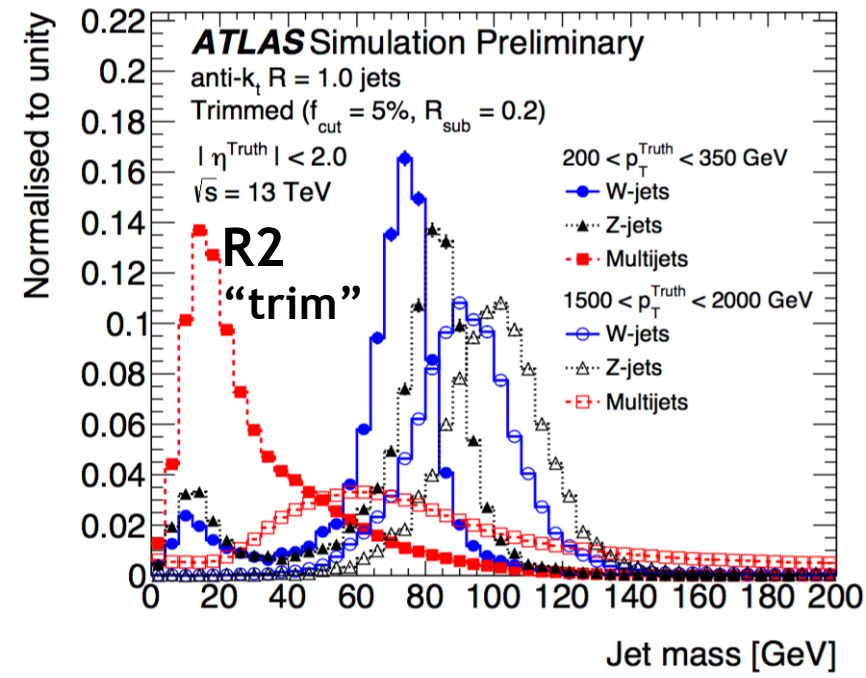
- **R2 (“trim”)**: AntiKt, R=1.0 trimmed,  $R_{\text{subject}} = 0.2$ ,  $f_{\text{cut}} = 5\%$ : **best BG rejection + PU stability**
- **P515 (“prun”)**: C/A, R=1.0 pruned  $Z_{\text{cut}} = 15\%$ ,  $R_{\text{cut}} = 1/2$   
*particularly good for W/Z-tagging at Pt ~500 GeV*
- **Y15 (“SplitFilt15”)**: C/A, R=1.2 split-filtered  $R_{\text{subject}} = 0.3$ ,  $y_{\text{cut}} > 15\%$
- **Y4 (“SplitFilt4”)**: C/A, R=1.2 split-filtered  $R_{\text{subject}} = 0.3$ ,  $y_{\text{cut}} > 4\%$



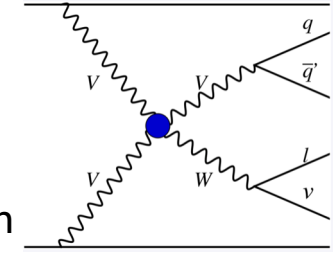
reduced pileup dependence after grooming: Each value of  $\delta \langle M \rangle / \delta \text{NPV}$  is the slope of a *straight line* fit of  $\langle M \rangle$  versus NPV



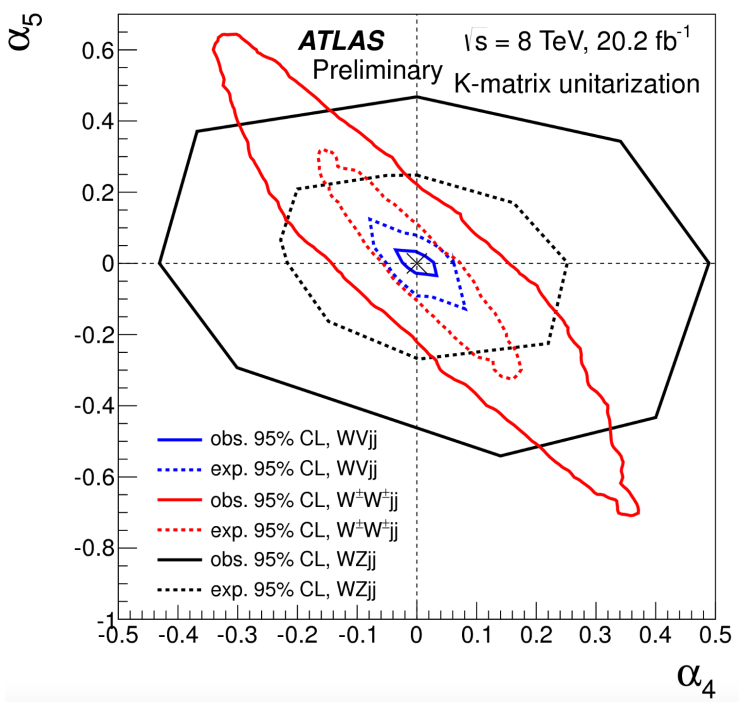
# Jet mass after grooming



- No jet-level calibrations applied. Jet mass distributions for  $200 < Pt < 350$  GeV and  $1500 < Pt < 2000$  GeV.
- R2 and p515: multi jet shape has strong dependence on Pt
- Y15 and Y4: more stable multijet shape across Pt bins
- W-jet and Z-jet mass shapes similar across Pt bins for Y4 and Y15

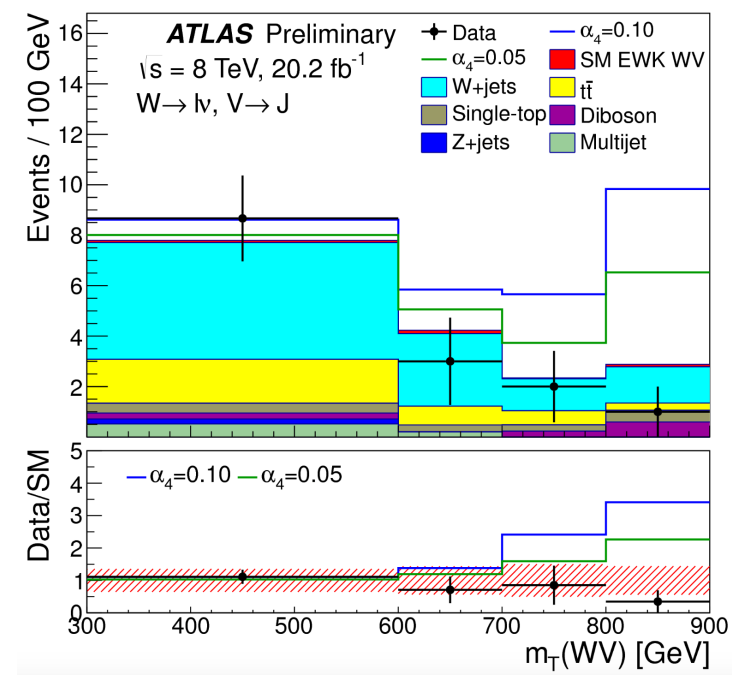


- Vector Boson scattering (VBS) with the W decaying hadronically and Z boson leptonically
  - Key probe of EW Symmetry Breaking
  - Sensitive to anomalous quartic gauge couplings: is the coupling as predicted by SM?
  - Previous searches for aQGCs in VBS have focused on leptonic decays → smaller branching fraction
- Mass-drop filtered jets are used, for multiple large-R jets, the one with mass closest to the nominal W mass is used.



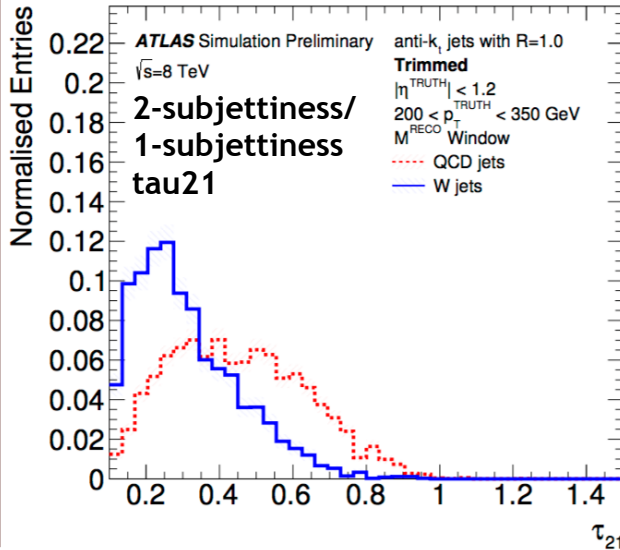
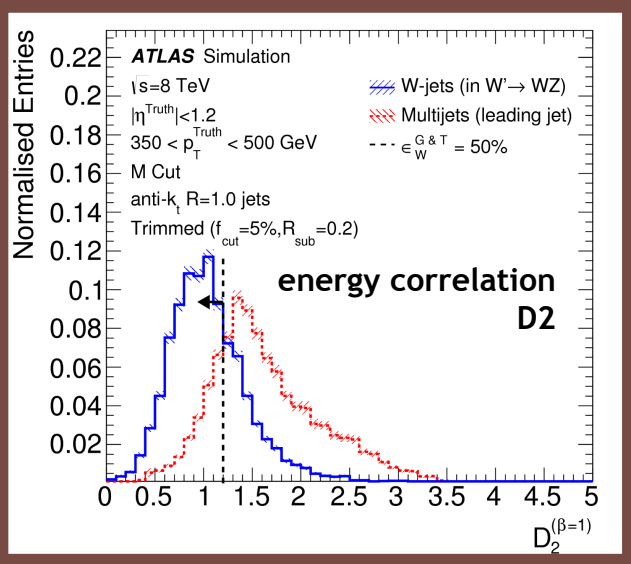
2D limits on aQGC parameters for current result (blue) and previous results (red).

- merged regime
  - aQGC limits are 40% better compared to only using resolved events.
  - most of the aQGC sensitivity comes from the highest- $m_T$  ( $W V$ ) bins, where the merged category is powerful.
- More stringent limits obtained than the previous constraints on these parameters in searches for VBS in the  $W_{\pm}W_{\pm}$  and  $WZ$  leptonic decay channels

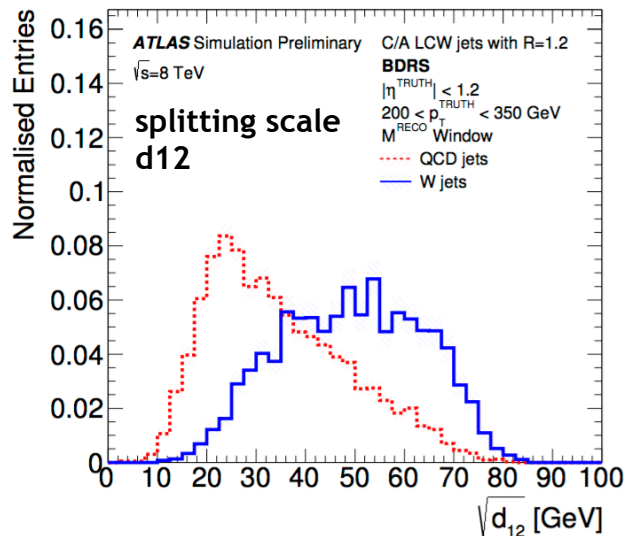
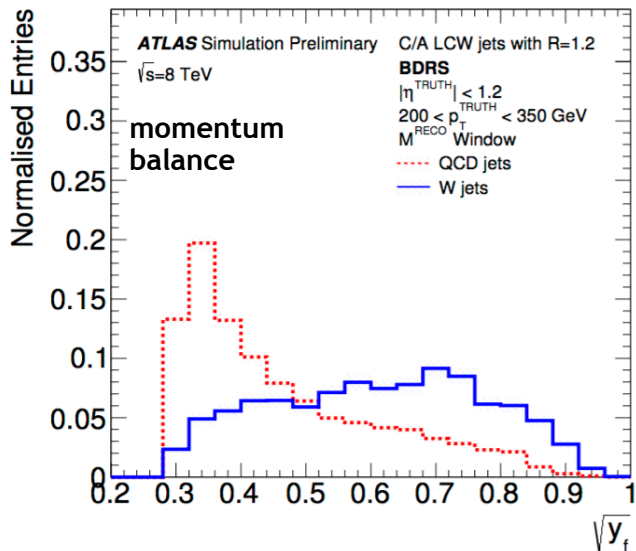


$m_T$  ( $W V$ ) distribution used to extract aQGC limits: in merged regime

## D2 validated with Run 2 data



## signal and background discrimination using substructure variables

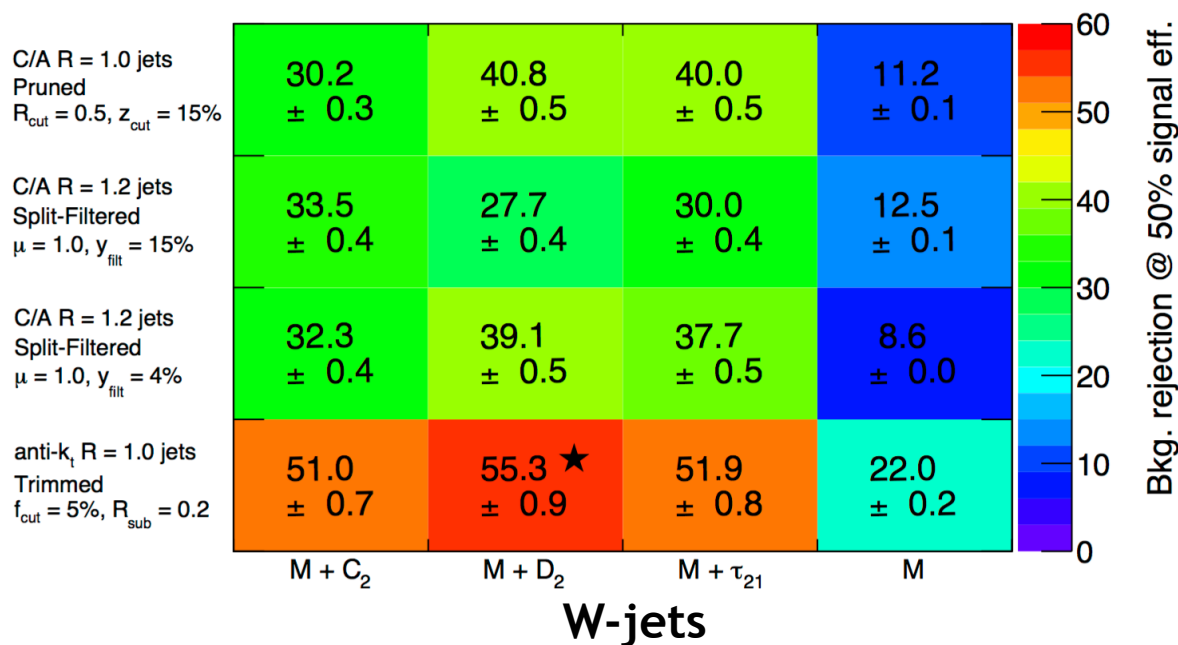


- **D2**: related to ratio of 3- and 2-point correlation functions. [J. High Energ. Phys. \(2014\) 2014: 9.](#)
- **N-subjettiness**: extent to which the substructure of a jet resembles N or fewer subjects. [J. High Energ. Phys. \(2011\) 2011: 15.](#)
- **Momentum balance**: ratio of splitting scale and jet mass of the final clustering step. [Phys. Rev. Lett. 100, 242001](#)
- **Mass-drop**: last step of recombination when two proto jets are combined into one. Fraction of mass carried by the most massive proto jet.
- **Splitting scale**: for a jet re-clustered with the Kt algorithm-Kt distance between two proto jets in the final clustering step. [Phys. Rev. D 65, 096014](#)
- **Jet width**: radial moment in eta, phi space weighted by Pt. [Phys. Rev. D 86, 072006](#)
- **Planar flow**: how uniformly spread out the jet energy is perpendicular to it's axis. [Phys. Rev. D 79, 074017](#)

**ATLAS** Simulation Preliminary

$\sqrt{s} = 13$  TeV ★ = Optimal grooming + tagging combination

$|\eta^{\text{Truth}}| < 2.0, 200 < p_T^{\text{Truth}} < 350$  GeV,  $M^{\text{Reco}}$  Cut W-jets

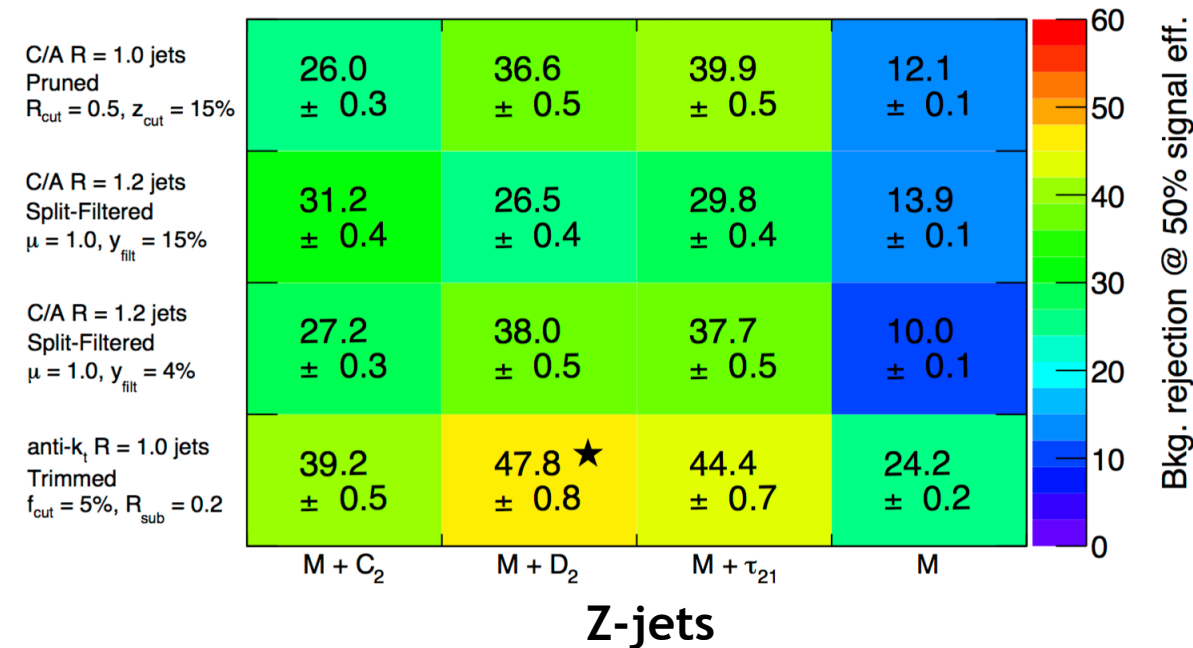


## Low Pt

**ATLAS** Simulation Preliminary

$\sqrt{s} = 13$  TeV ★ = Optimal grooming + tagging combination

$|\eta^{\text{Truth}}| < 2.0, 200 < p_T^{\text{Truth}} < 350$  GeV,  $M^{\text{Reco}}$  Cut Z-jets



- M+C2, M+D2, M+t21: substructure + Mass, M: mass window (tighter) requirement only,  $200 < Pt < 350$  GeV, 50% signal eff.
- Mass only requirement yields higher background rejection for Z-jets than W-jets: Z boson mass is ~10 GeV higher than W boson so better separated than the multijet background.
- Tagger+mass combination yields better background rejection for W-jets than Z-jets - broader Z-jet mass distribution (slide 5) than differences in substructure variable distributions.

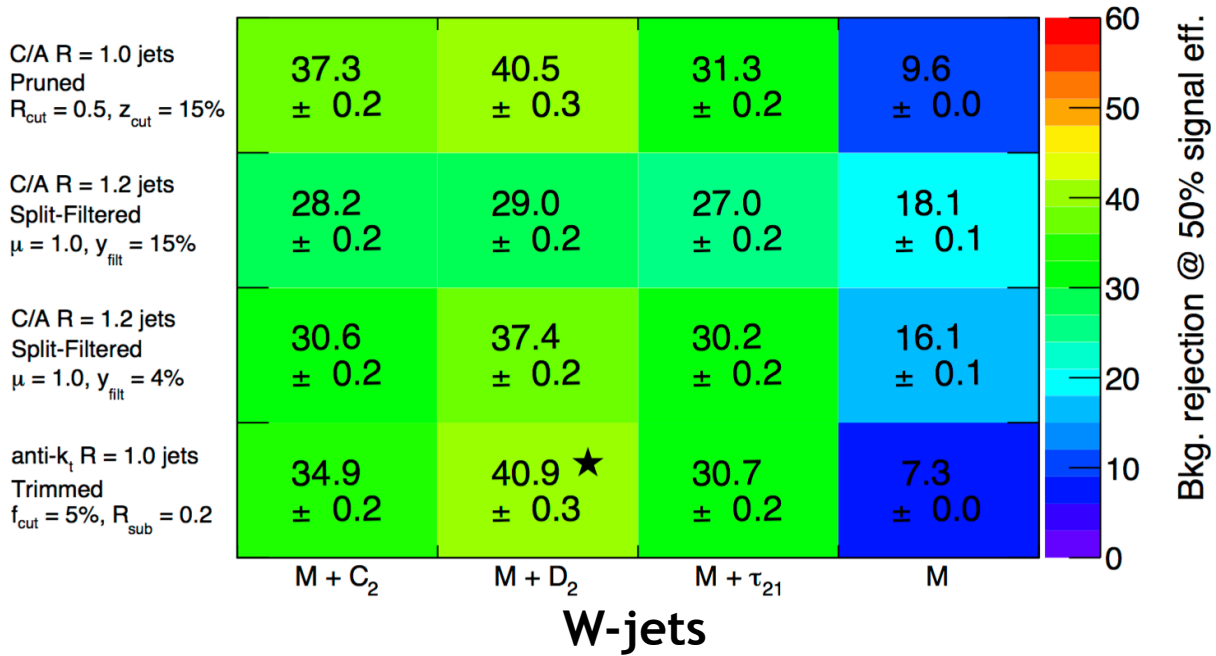


## High Pt

**ATLAS** Simulation Preliminary

$\sqrt{s} = 13$  TeV ★ = Optimal grooming + tagging combination

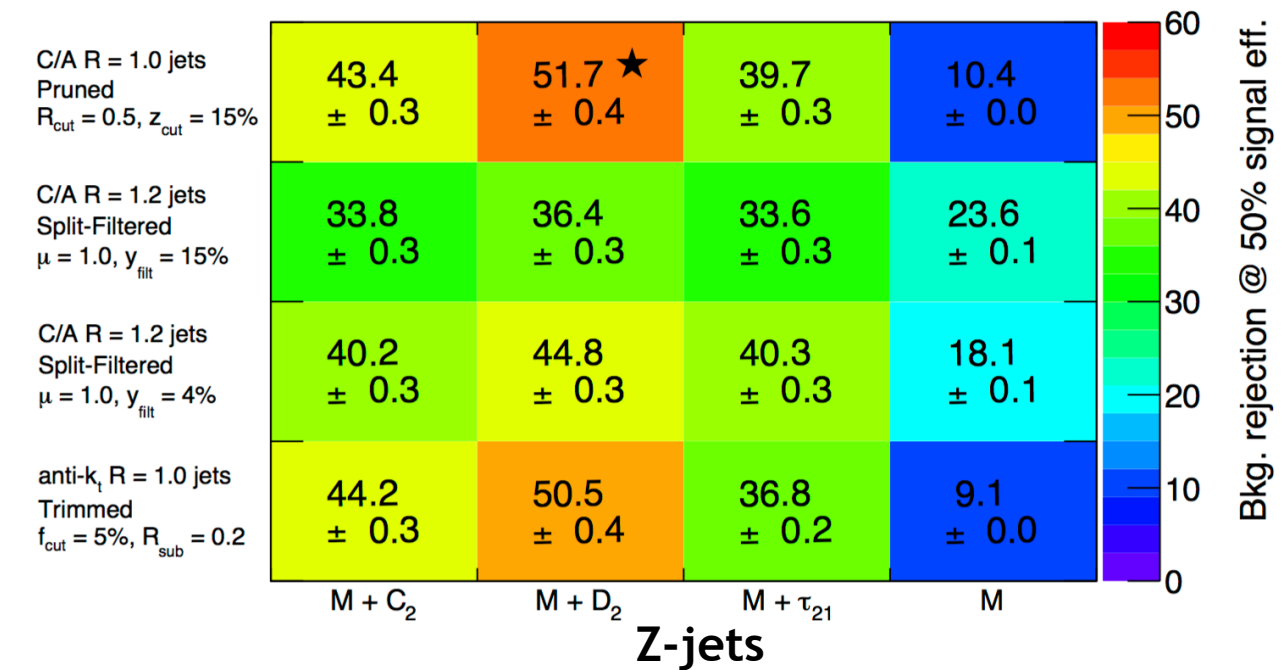
$|\eta^{\text{Truth}}| < 2.0, 1500 < p_T^{\text{Truth}} < 2000$  GeV,  $M^{\text{Reco}}$  Cut W-jets



**ATLAS** Simulation Preliminary

$\sqrt{s} = 13$  TeV ★ = Optimal grooming + tagging combination

$|\eta^{\text{Truth}}| < 2.0, 1500 < p_T^{\text{Truth}} < 2000$  GeV,  $M^{\text{Reco}}$  Cut Z-jets



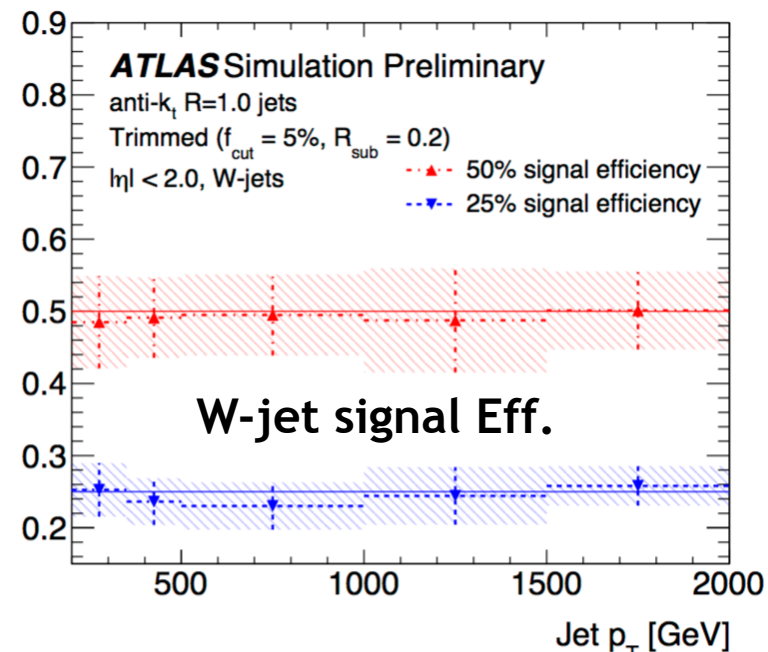
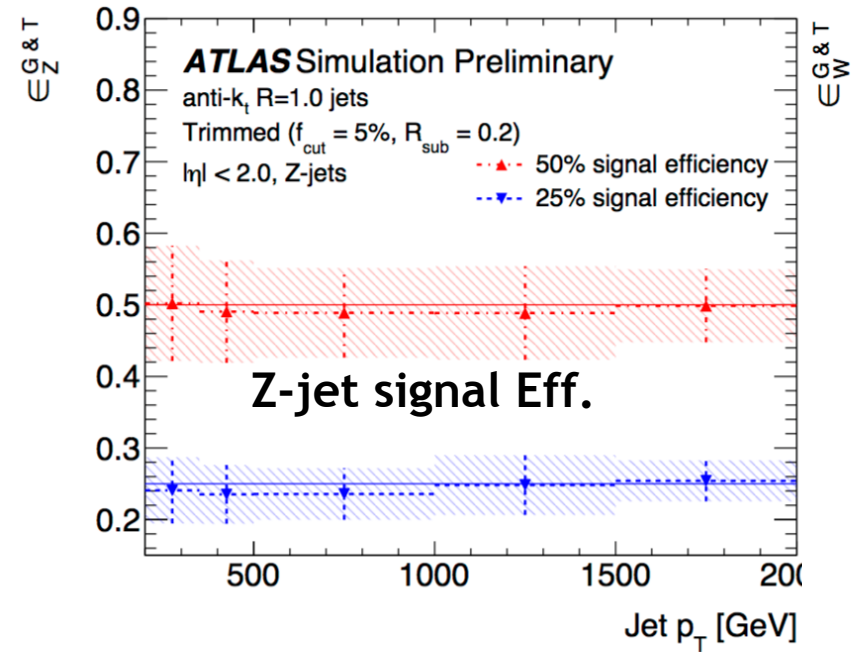
- $M+C_2, M+D_2, M+t_{21}$ : substructure + Mass, M: mass window (narrower) requirement only,  $1.5 < Pt < 2$  TeV, 50% signal eff.
- At high  $p_T$ , better background rejection for Z-jets than W-jets.

	Best G&T combinations in ranges of $p_T^{\text{Truth}}$ [GeV]				
	250 – 350	350 – 500	500 – 1000	1000 – 1500	1500 – 2000
<b>W-jets</b>					
Best ‘Medium’ tagger $1/\epsilon_{QCD}^{\text{G&T}}$	R2 $D_2$ “trim” $55.3 \pm 0.9$	P515 $C_2$ “prun” $71.5 \pm 0.8$	P515 $C_2$ “prun” $70.0 \pm 0.7$	P515 $D_2$ “prun” $55.2 \pm 0.5$	R2 $D_2$ “trim” $40.9 \pm 0.3$
Best ‘Tight’ tagger $1/\epsilon_{QCD}^{\text{G&T}}$	R2 $D_2$ $215 \pm 7$	P515 $C_2$ $271 \pm 5$	P515 $D_2$ $274 \pm 6$	R2 $D_2$ $254 \pm 4$	R2 $D_2$ $188 \pm 3$
<b>Z-jets</b>					
Best ‘Medium’ tagger $1/\epsilon_{QCD}^{\text{G&T}}$	R2 $D_2$ “trim” $47.8 \pm 0.8$	P515 $C_2$ “prun” $65.4 \pm 0.8$	P515 $C_2$ “prun” $67.1 \pm 0.7$	P515 $D_2$ “prun” $65.3 \pm 0.6$	P515 $C_2$ “prun” $51.7 \pm 0.4$
Best ‘Tight’ tagger $1/\epsilon_{QCD}^{\text{G&T}}$	R2 $D_2$ $179 \pm 5$	P515 $C_2$ $255 \pm 5$	P515 $D_2$ $259 \pm 6$	R2 $D_2$ $270 \pm 5$	R2 $D_2$ $232 \pm 4$

- P515 and R2 with C2 and D2 perform particularly well
- Signal Efficiency:
  - Medium: 50%
  - Tight: 25%
- G&T: grooming and tagging
- Grooming algorithms
  - R2
  - P515
- Tagging variables
  - D2
  - C2
- $1/\epsilon_{QCD}$ : background rejection for QCD multijets
- Only stat uncertainties shown



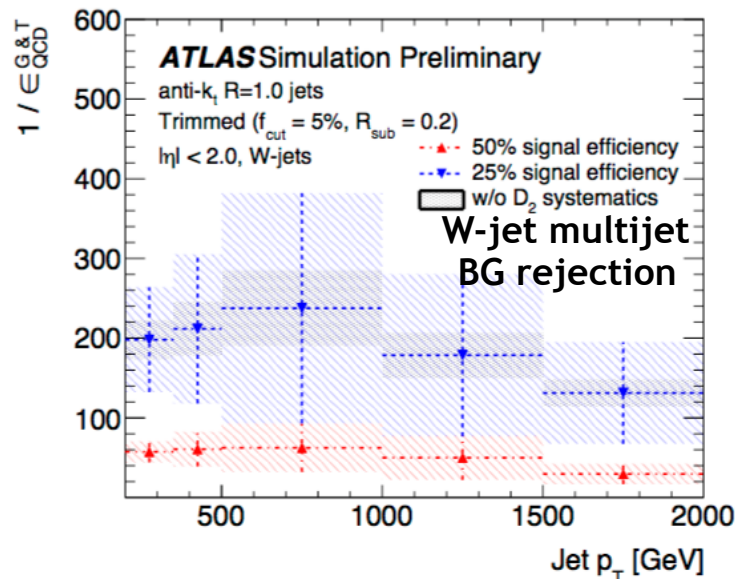
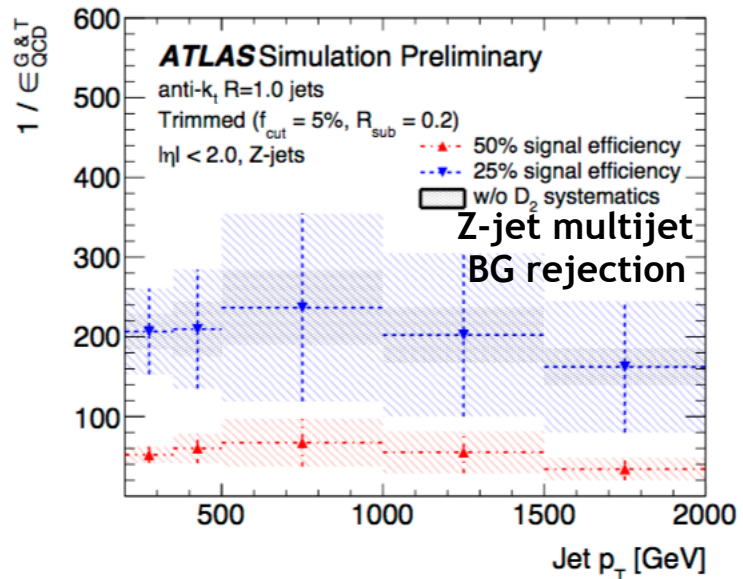
# Systematic uncertainties on Signal efficiency and Background rejection



Signal efficiency and background rejection in MC only for two tagger working points:

- medium (50% signal efficiency)
- Tight (25% signal efficiency)

$$\epsilon_{MC} = \frac{N_{\text{Boosted-W}}^{\text{tagged}}}{N_{\text{Boosted-W}}^{\text{pre-tagged}}}$$

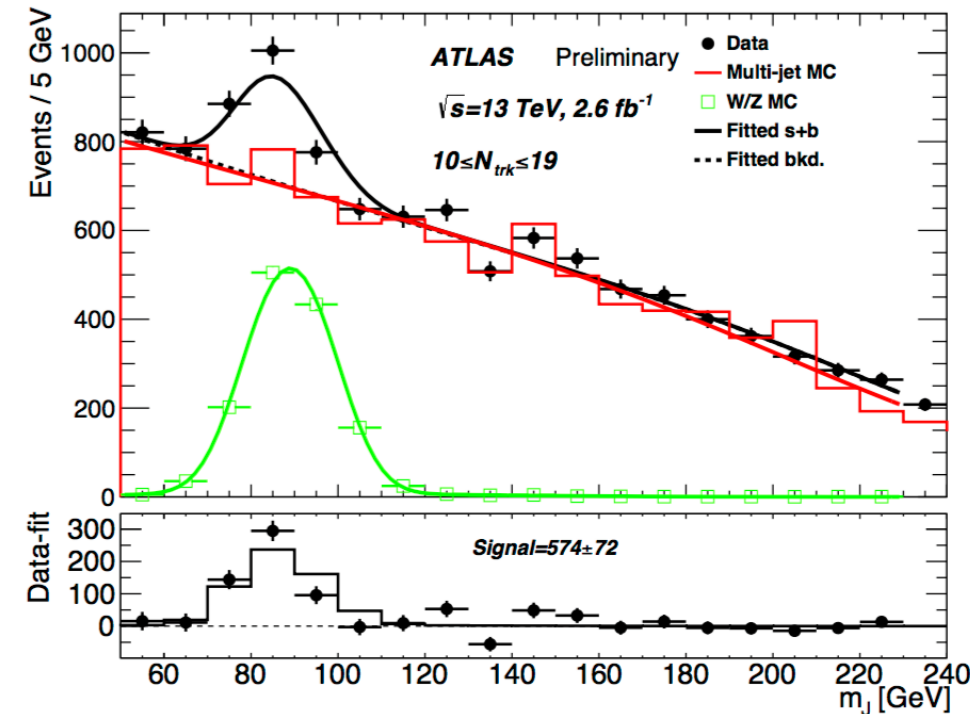


$D_2$  tagger related uncertainties dominate at high Pt.

- Gluon-induced jets contain more charged hadrons than quark-induced jets. The number of tracks associated with the boson jet versus the background QCD jet can be used to provide further discrimination than the standard tagging requirements → BG has a significant gluon contribution
- This is applied by cutting on the number of associated tracks with a jet : tracks are matched to the jet via ghost association.
- Requirement that  $N_{trk} < 30$  after grooming and substructure variable related cuts.

Ghost association: Add multiple constituents for jet finding with negligible momentum and with same direction as track, for each track.

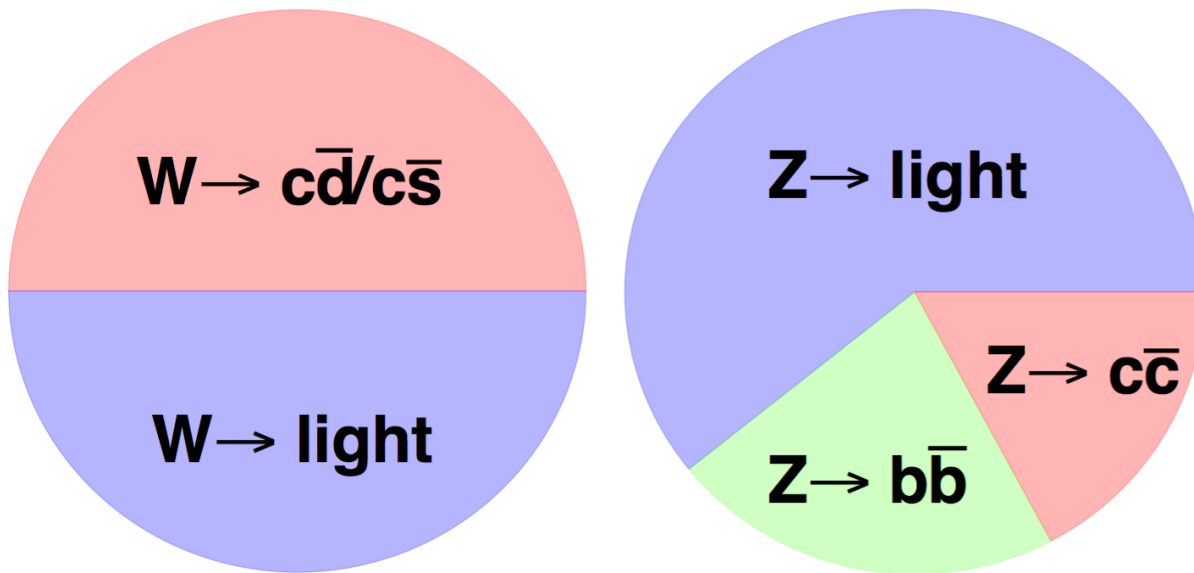
optimized medium (50% signal eff.) D2 selection for the R2 configuration is used.



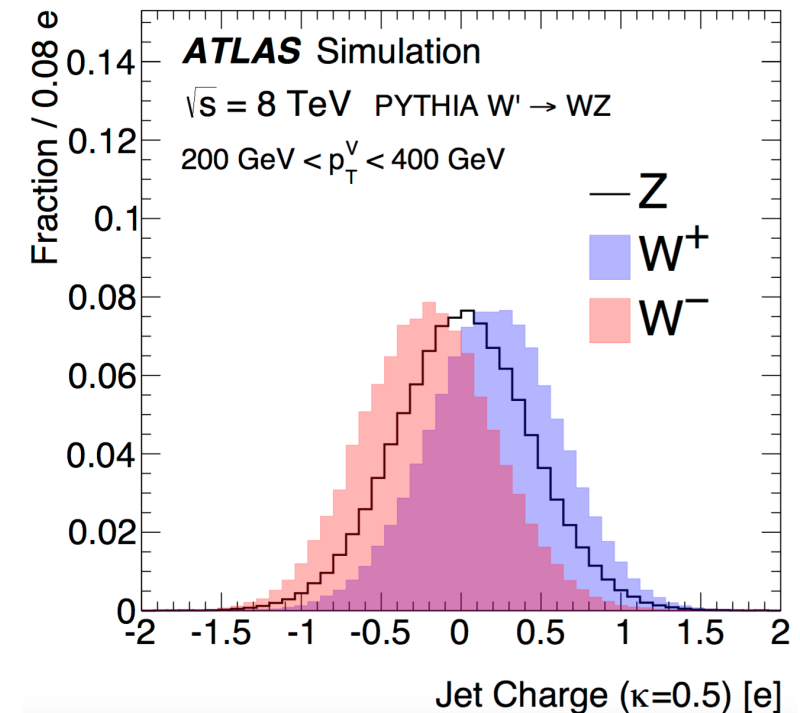
Validation of ntrack selection in data in  $10 < N_{trk} < 19$  bin

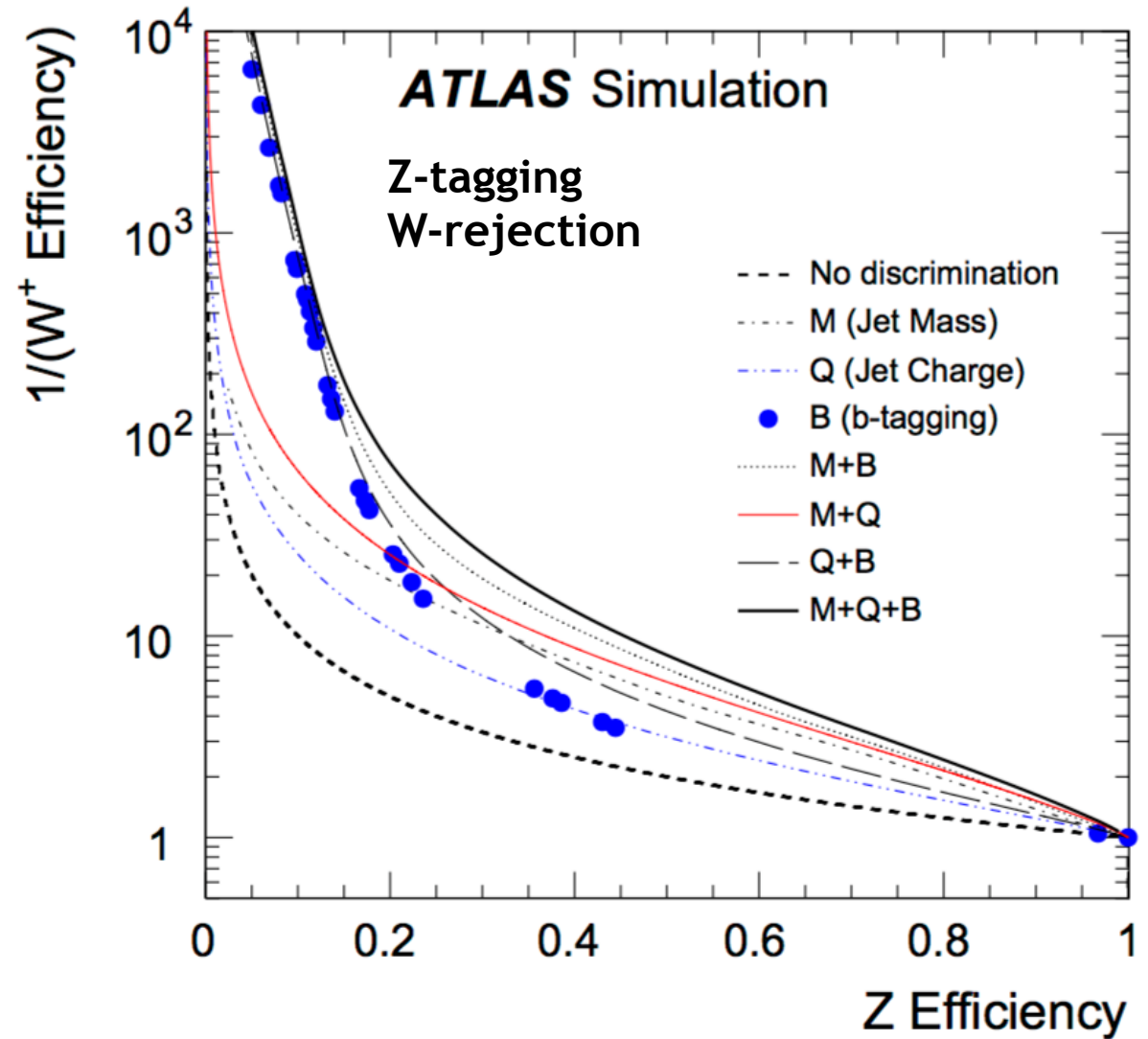
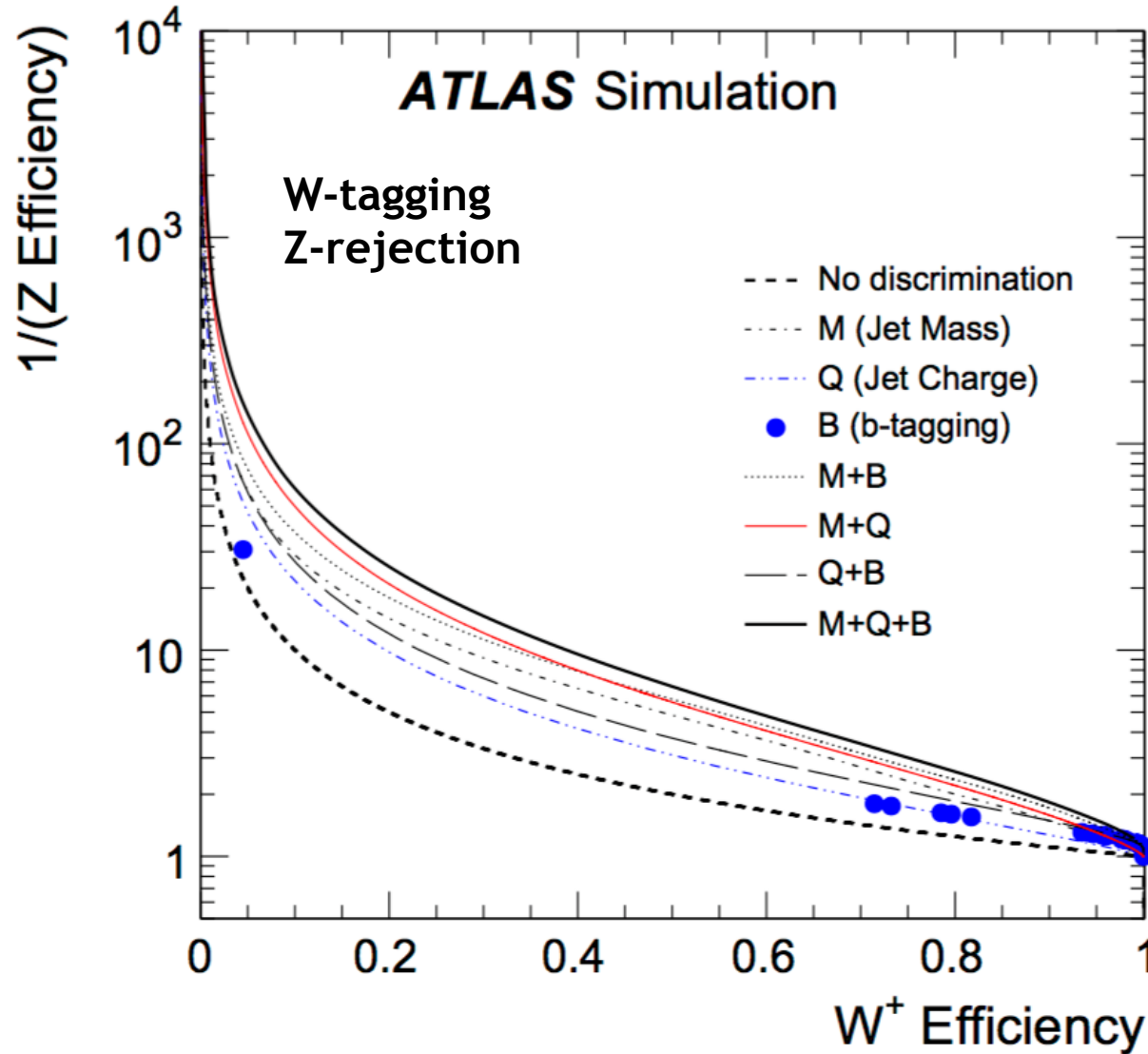
expected sensitivity of the analysis is improved by about 30% by using the additional  $N_{trk}$  selection.

- Differences in production XS and subtle differences in differential decay.
- Differentiating features: mass, charge and branching ratios.
- Variables considered for likelihood tagger for W-jet and Z-jet separation
  - Jet mass - sensitive to boson mass
  - Jet charge - sensitive to boson charge
  - B-tagging discriminant - sensitive to heavy flavor decay branching fractions of the bosons



hadronic branching fractions of the  $W^+$  boson, and Z boson

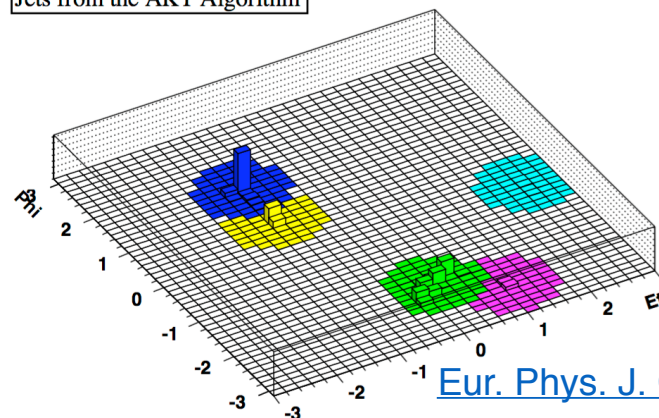




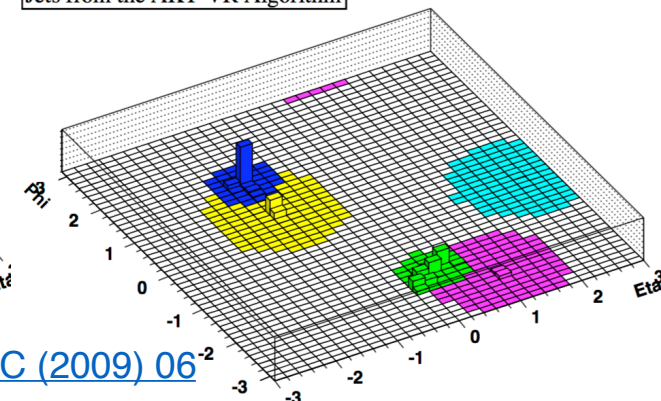
At low Z-tagging efficiencies, large W rejection due to  $Z \rightarrow bb$  tagging

- Size of a variable- $R$  jet depends on its transverse momentum, making these jets ideally suited for boosted decay topologies - separation between decay products of a heavy particle decreases with momentum.  $R_0 \rightarrow R_{eff}(p_{T,i}) = \frac{\rho}{p_{T,i}}$ .
- Significant performance improvements are found with Variable- $R$  jets for a number of discriminating variables.

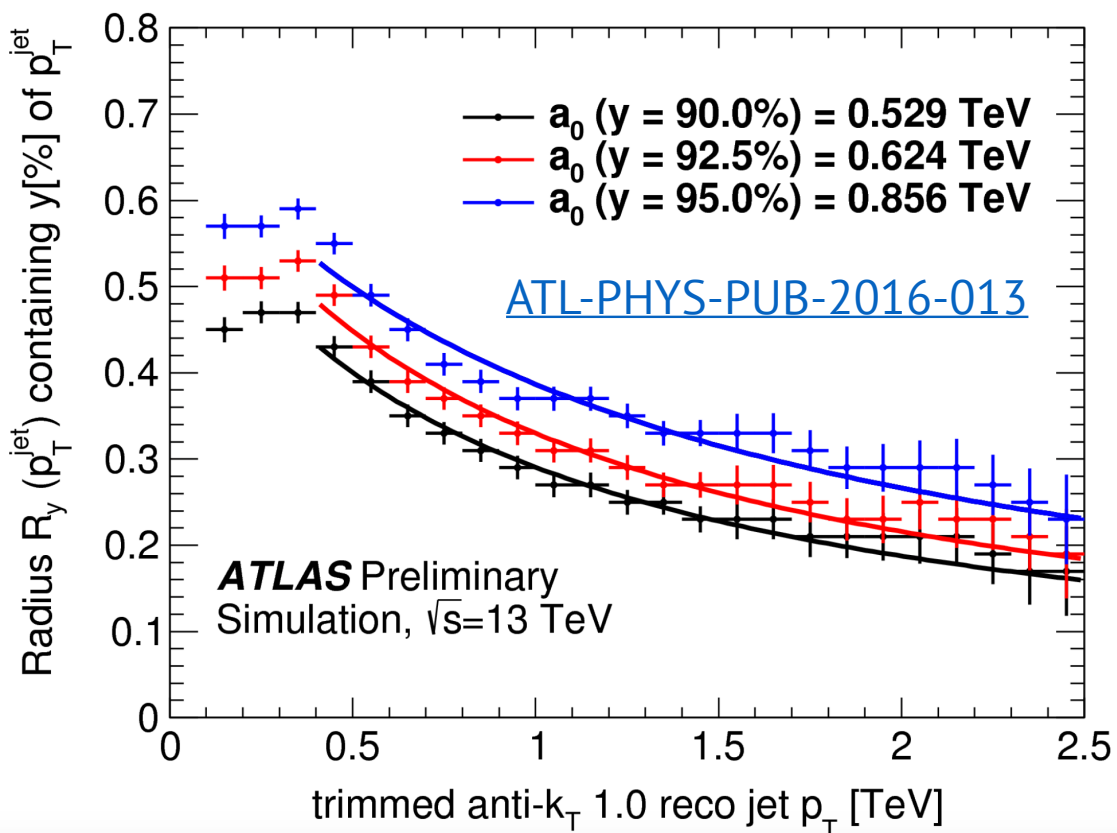
Jets from the AKT Algorithm



Jets from the AKT-VR Algorithm



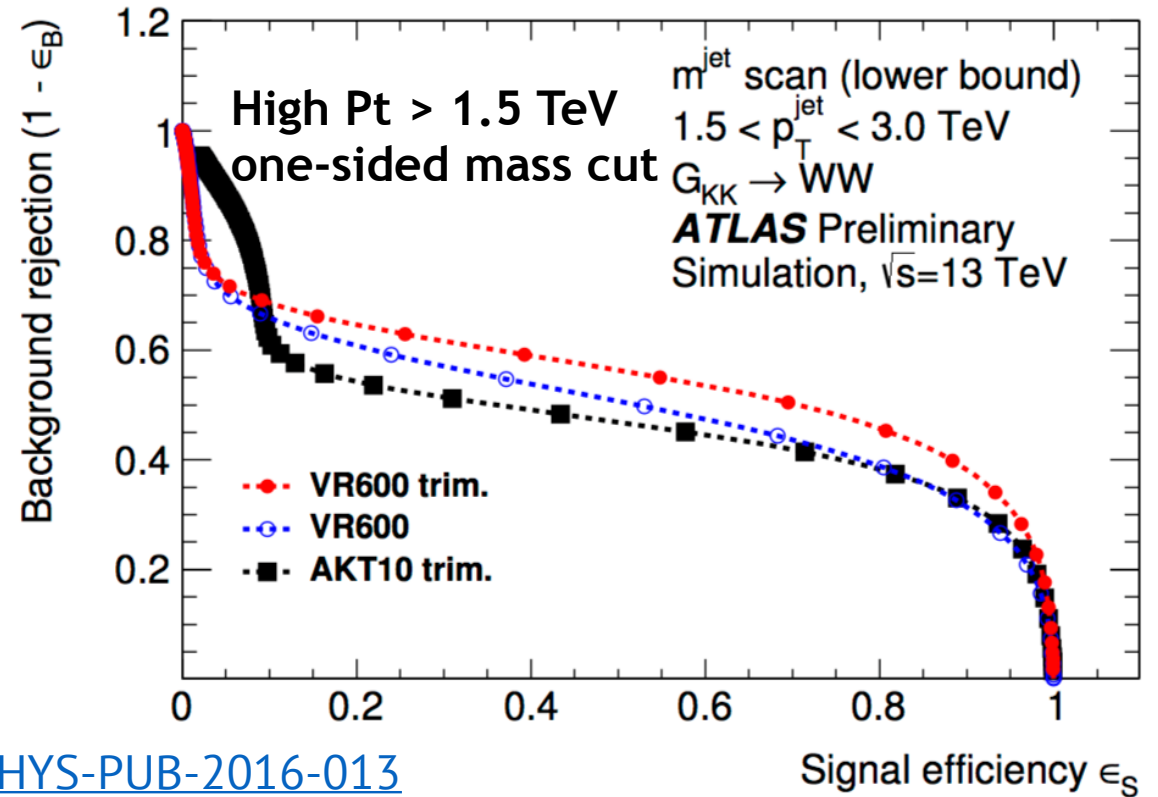
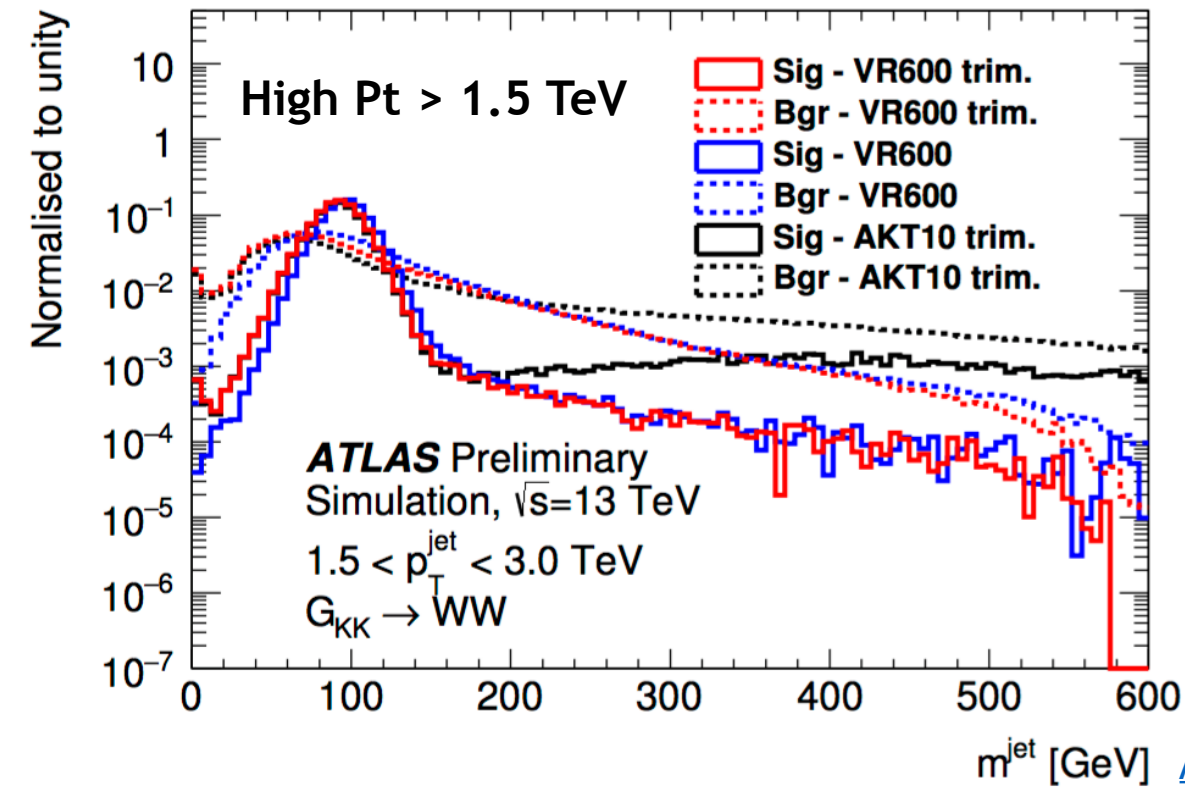
[Eur. Phys. J. C \(2009\) 06](#)



- Can build Variable- $R$  versions of the AntiKt (kt) and C/A algorithms.
- At high  $P_t$ , most of the jet  $P_t$  is contained within a  $dR$  smaller than the  $R=1$  jet used in tagging.
- Using VR jets with smaller effective area reduces contaminations from PU/UE and ISR.



# Tagging using variable-R jets: performance for WW final state

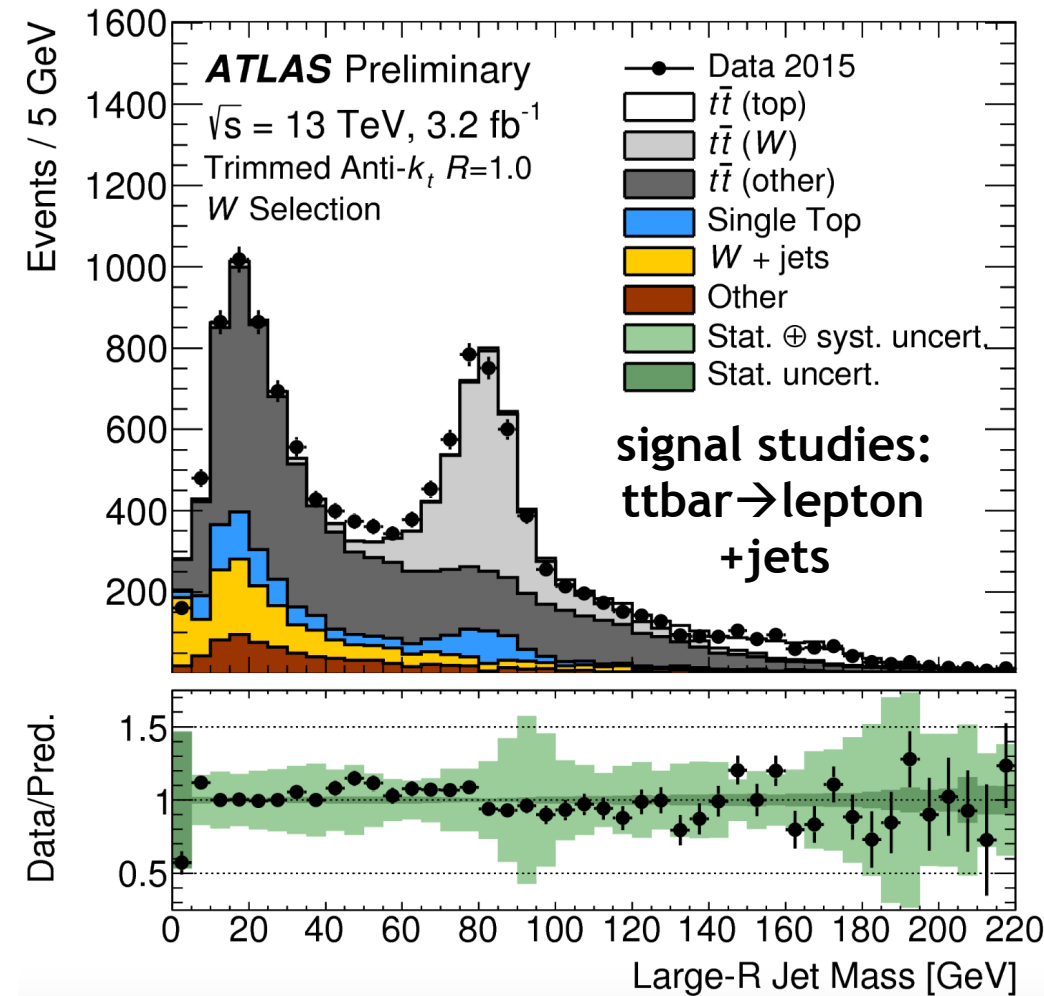


- Trimmed VR jets behave similarly to AntiKt  $R = 1.0$  jets below 1 TeV, while outperforming AntiKt  $R = 1.0$  jets for  $Pt > 1$  TeV.
- With the two-sided mass window cut, variable-R jets has slightly worse performance than traditional AntiKt  $R=1.0$  jets as a result of lower ISR contamination (high mass tail) for VR jets
  - Only small reduction (2-4%) in  $S/\sqrt{B}$  for  $Pt > 1.0$  TeV
- Notable performance improvements for substructure variables like first Kt splitting scale.

[ATL-PHYS-PUB-2016-013](https://arxiv.org/abs/1603.07513)



## Calorimeter jet mass spectrum for leading jet Pt in 13 TeV data and MC for signal efficiency (left) and background rejection studies (right)

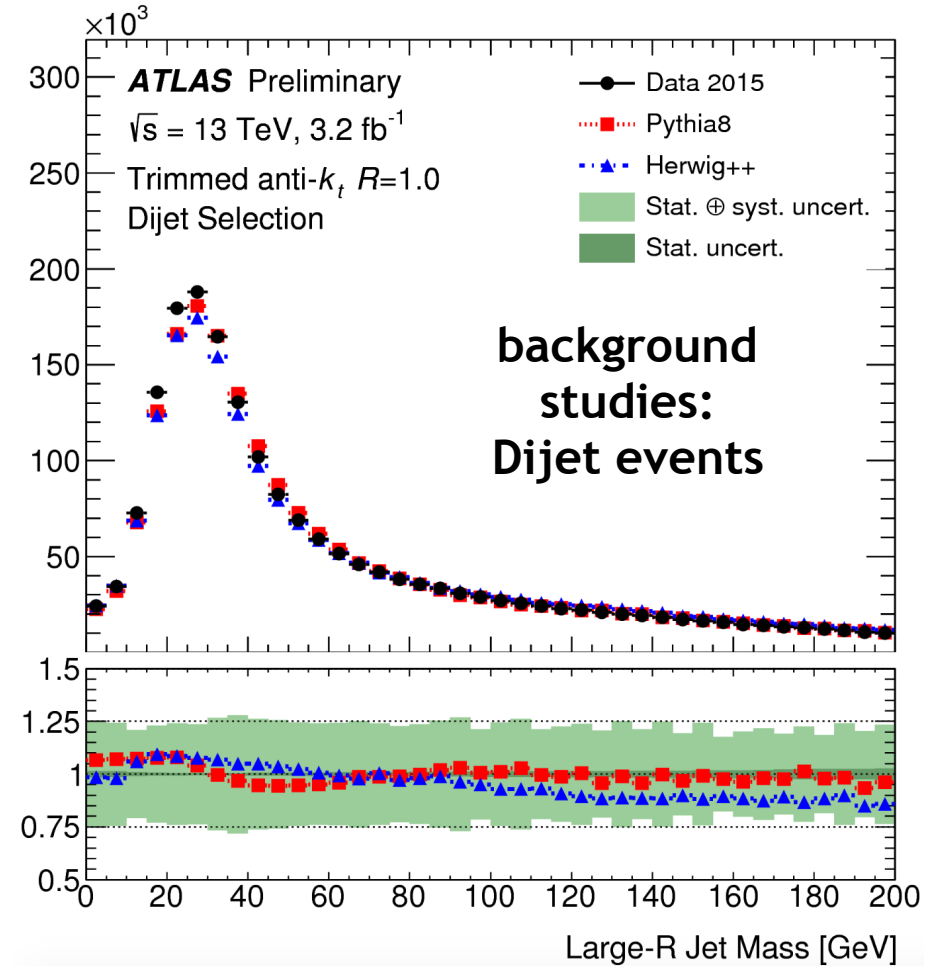


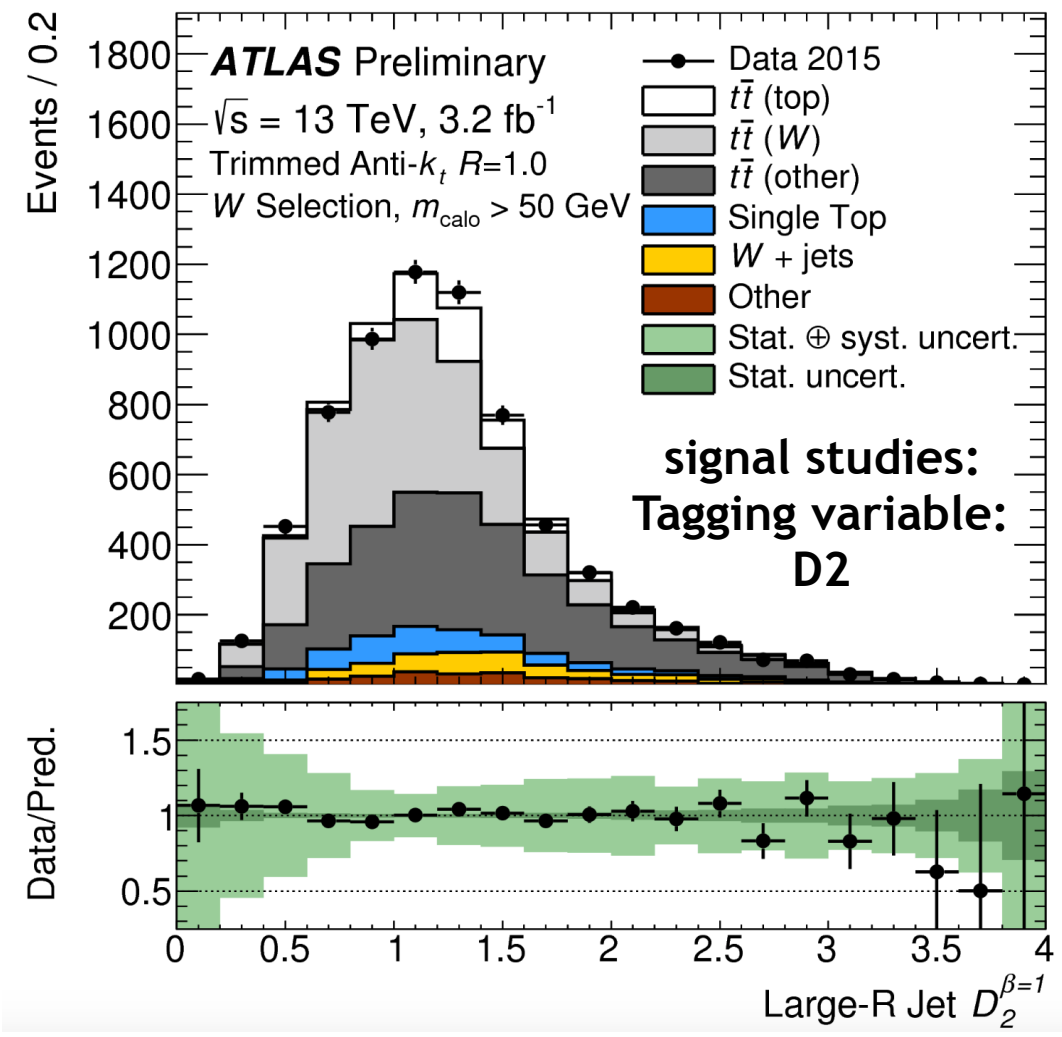
### signal studies (left):

- $t\bar{t}$  sample is split into a “W -matched” part by requiring the decay products of W boson to be within the large-R jet.
- “top-matched” part requires the qqb particles from the top-decay to be within the large-R jet.

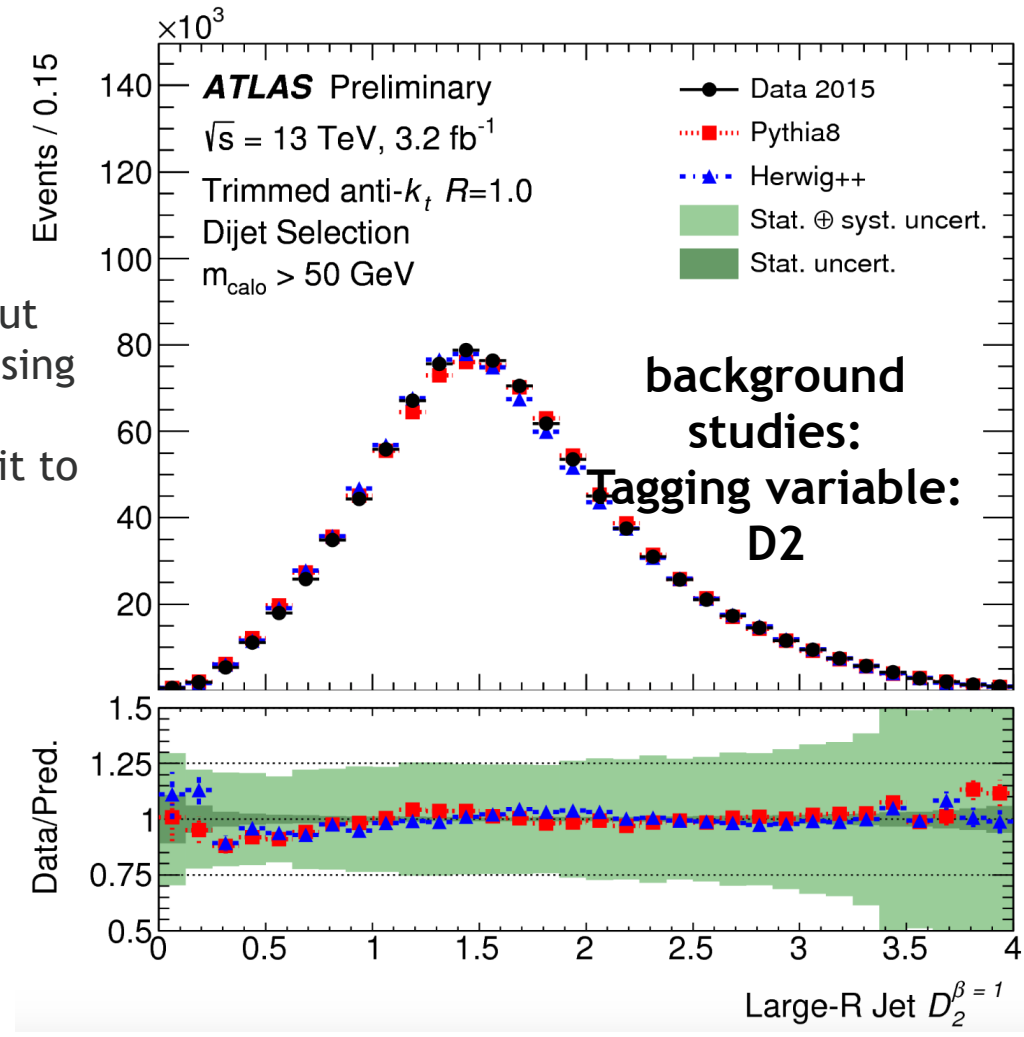
### background studies:

- large-R jet trigger requirement with  $ET > 360 \text{ GeV}$ .
- At least 2 large-R jets with  $Pt > 200 \text{ GeV}$ ,
- leading jet  $Pt > 500 \text{ GeV}$



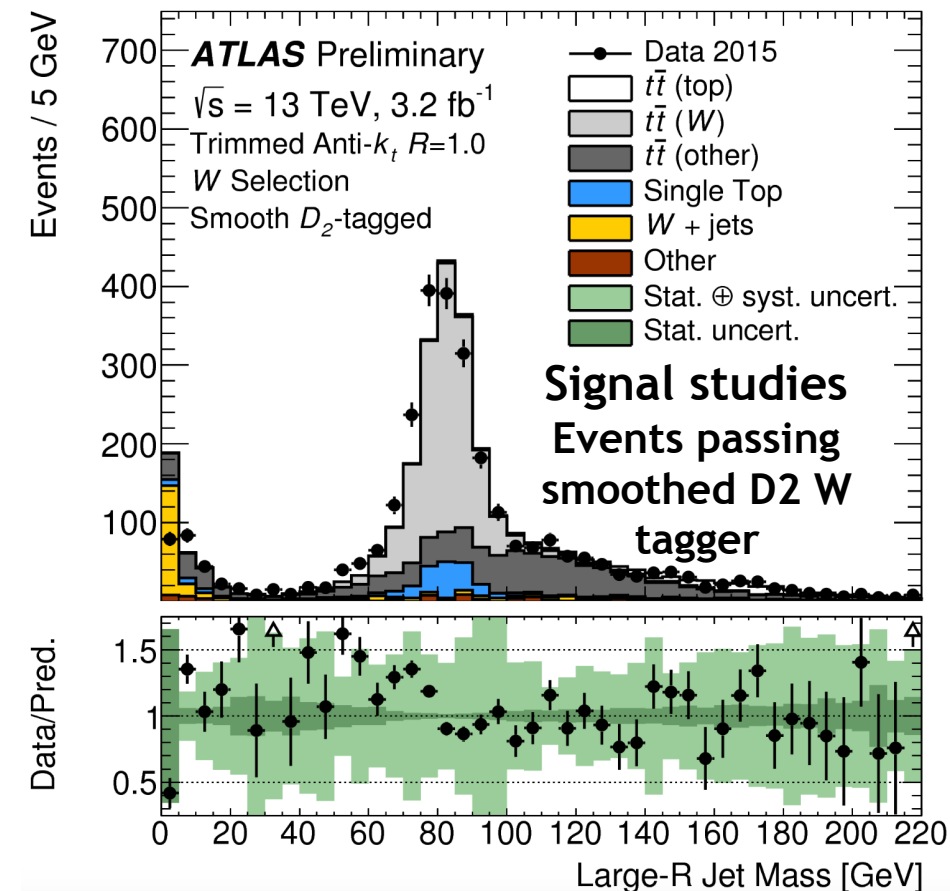


- D2 smoothed tagging algorithm applies a cut based on the jet pt using a fourth-order polynomial formula fit to define a smoothed selection.

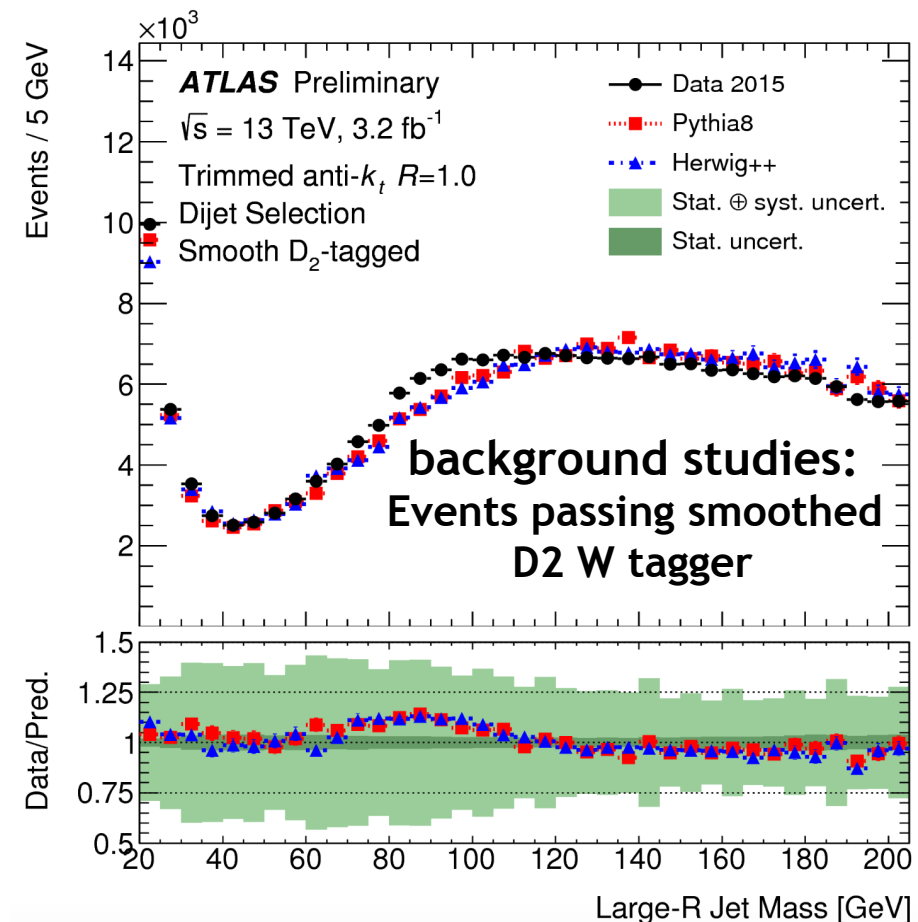


D2 observable in 13 TeV data and MC for signal efficiency (left) and background rejection studies (right)

D2-tagged events in 13 TeV data and MC for signal efficiency (left) and background rejection studies (right). 50% working point is shown.



- Systematic uncertainties are composed of:
- scale and resolution uncertainties on large-R jet energy, mass and substructure observables
- Ttbar modeling uncertainties on parton shower, hadronization model
- Initial-state and final-state radiation.
- Subdominant scale and resolution uncertainties on small-R jet energy and lepton kinematics

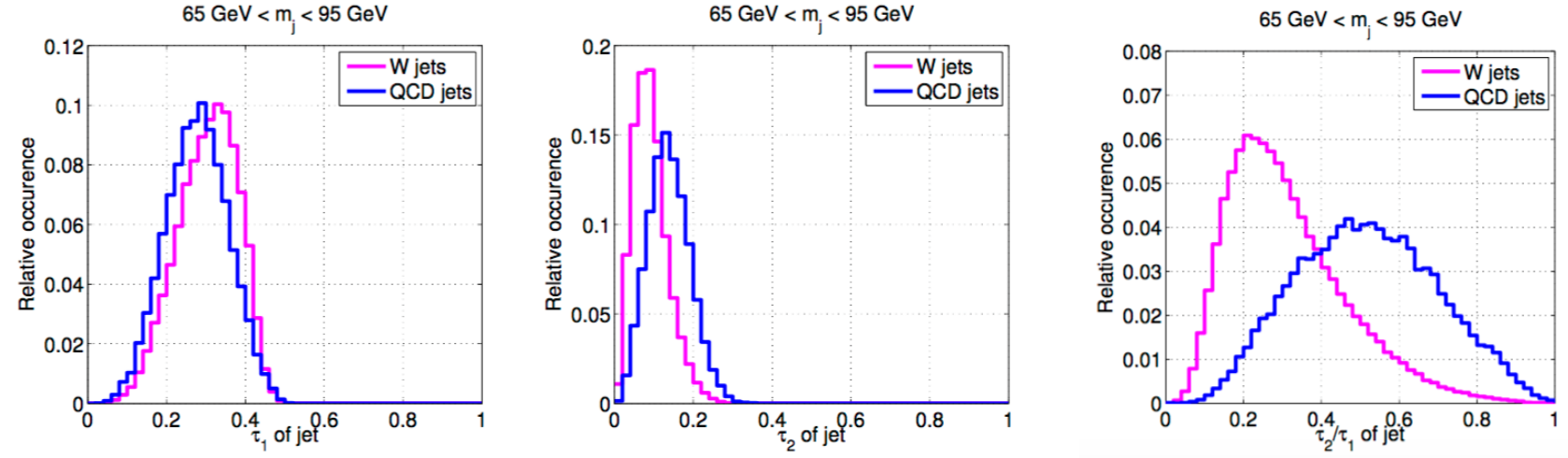
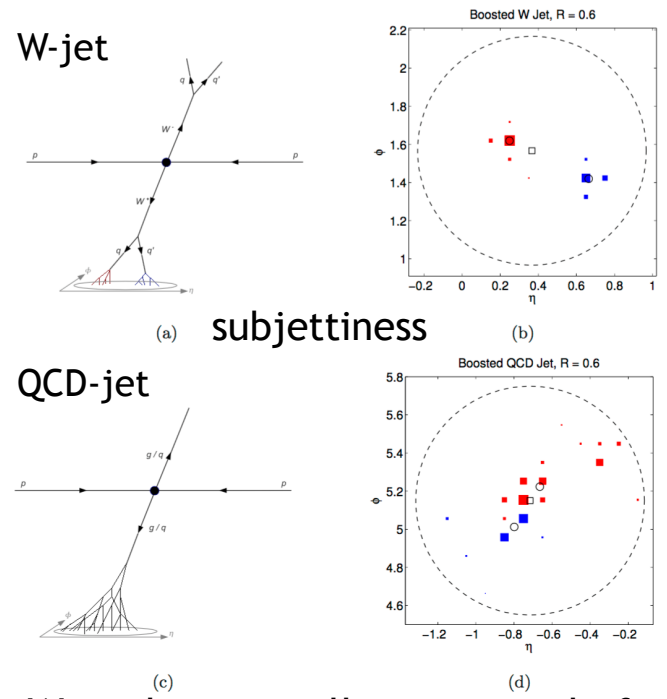


- Excess in plot on right around W mass is covered by systematics

- Best performing grooming techniques and tagger variables from Run 1 studies have been studied and validated in Run 2.
  - grooming with AntiKt  $R=1.0$  trimmed  $f_{cut}=5\%$ ,  $R_{subjet}=0.2$  was best performing
  - $\tau_{21}$ , C2 and D2 were best performing variables
- Early W- and Z-tagger results in Run 2 studied for medium and tight working points
  - Better background rejection for Z+jets overall
  - At low Pt, W-tagger has better background rejection with tagger+mass window cut selection.
- Track association to large R jet has shown to further help distinguish boson-jet signal and QCD jet background in the VVJJ analysis
- Variable-R jets have been studied for the first time as an alternate method for defining the boosted object compared to traditional methods, MC based studies:
  - Pt dependent jet radius expected to give better hold over pileup, comparison with  $R=1$  groomed jet shows better performance at high Pt
  - VR jets worth investigating for W tagging: notable performance improvement in first kT splitting scale
- Smoothed D2 tagger for W-jet tagging has been studied in 13 TeV data and MC



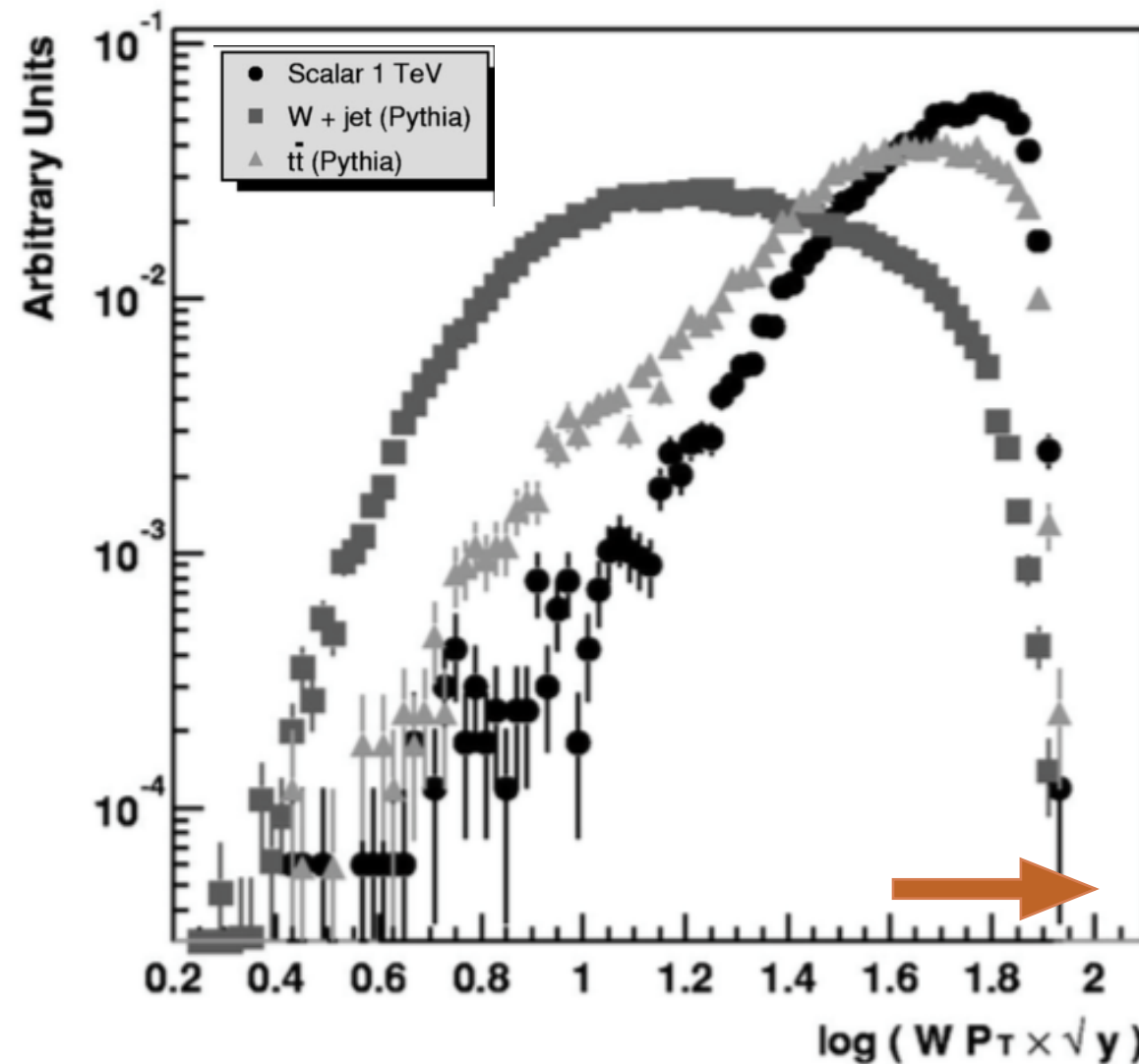
N-subjettiness is an inclusive jet shape that offers a direct measure of how well jet energy is aligned into subjets, and is therefore an excellent starting point for boosted object identification.



- W-jet his typically composed of two lobes of energy while QCD-jet acquires invariant mass through multiple splittings
- open squares (right plots) indicate overall jet energy direction, open squares represent the subject energy directions
- discriminating variable  $\tau_2/\tau_1$  measures the relative alignment of the jet energy along the open circles compared to the open square

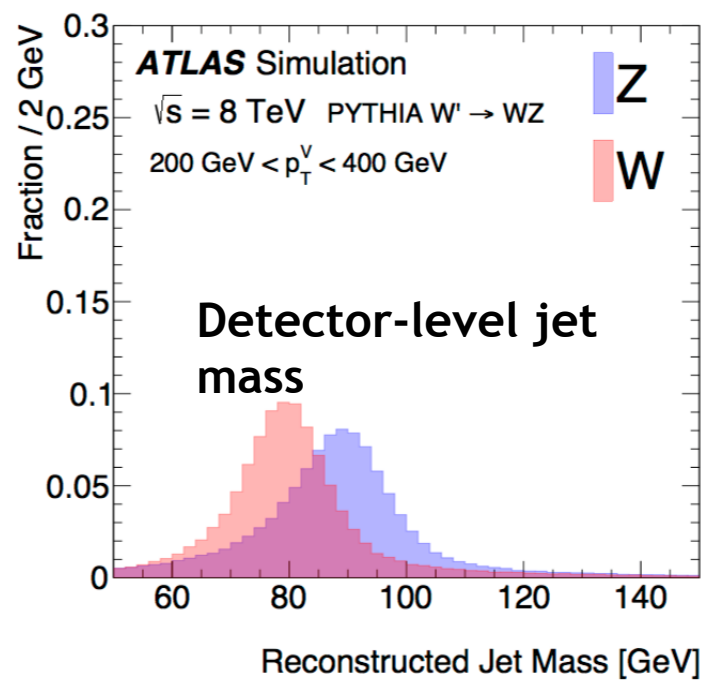
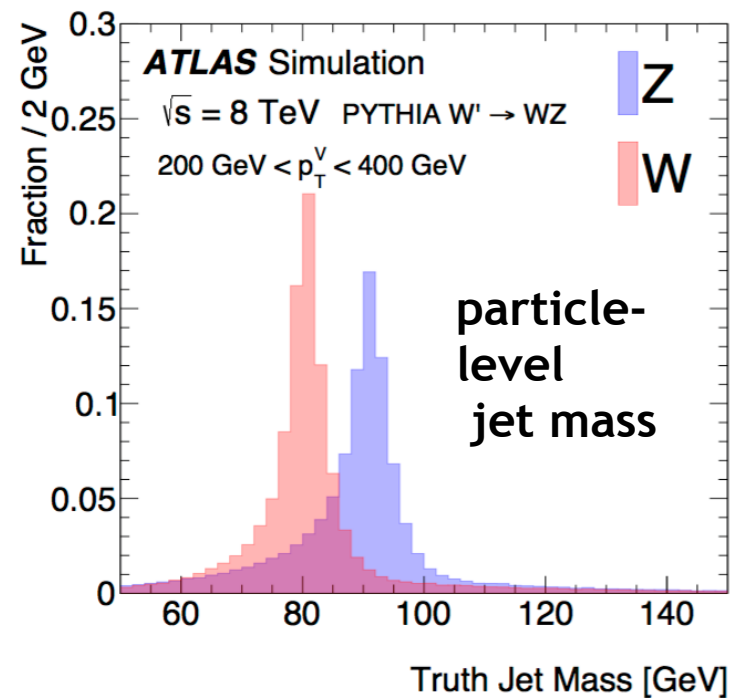
- W-jet is expected to have lower values of tau2 but a QCD-jet can also have low value of tau2 (middle)
- A larger value of tau1 (left) is likely for a W-jet but a QCD-jet with a diffuse spray of large-angle radiation can also have large tau1
- But QCD-jets with large values of tau1 typically have large values of tau2 as well. so we can use tau2/tau1 to discriminate between W-jet and QCD-jet.

- Jet is split into subjets using reclustering with Kt algorithm
- extra information gained from the subjet decomposition are the  $y$  cut at which the subjets are defined and the four-vectors of the subjets.
- For a genuine  $W$  decay the expectation is that the scale at which the jet is resolved into subjets i.e.  $y_{Pt^2}$  will be  $O(M^2_W)$ .
- scale of the splitting is indeed high in the signal and softer in the  $W$  jets background, where the hadronic  $W$  is in general a QCD jet rather than a genuine second  $W$ .
- $1.6 < \log(Pt y^{1/2}) < 2.0$  is a powerful cut for reducing QCD  $W$ +jets background.
- effect on the  $t\bar{t}$  background, which more often contains two real  $W$  bosons, is less marked.



distribution of  $\log(p_T y^{1/2})$

- W boson can radiate a Z boson making the separation challenging.
- Differentiating features: mass, charge and branching ratios.
- Differences in production XS and subtle differences in differential decay.
- Variables considered
  - Jet mass - sensitive to boson mass
  - Jet charge - sensitive to boson charge
  - B-tagging discriminant - sensitive to heavy flavor decay branching fractions of the bosons

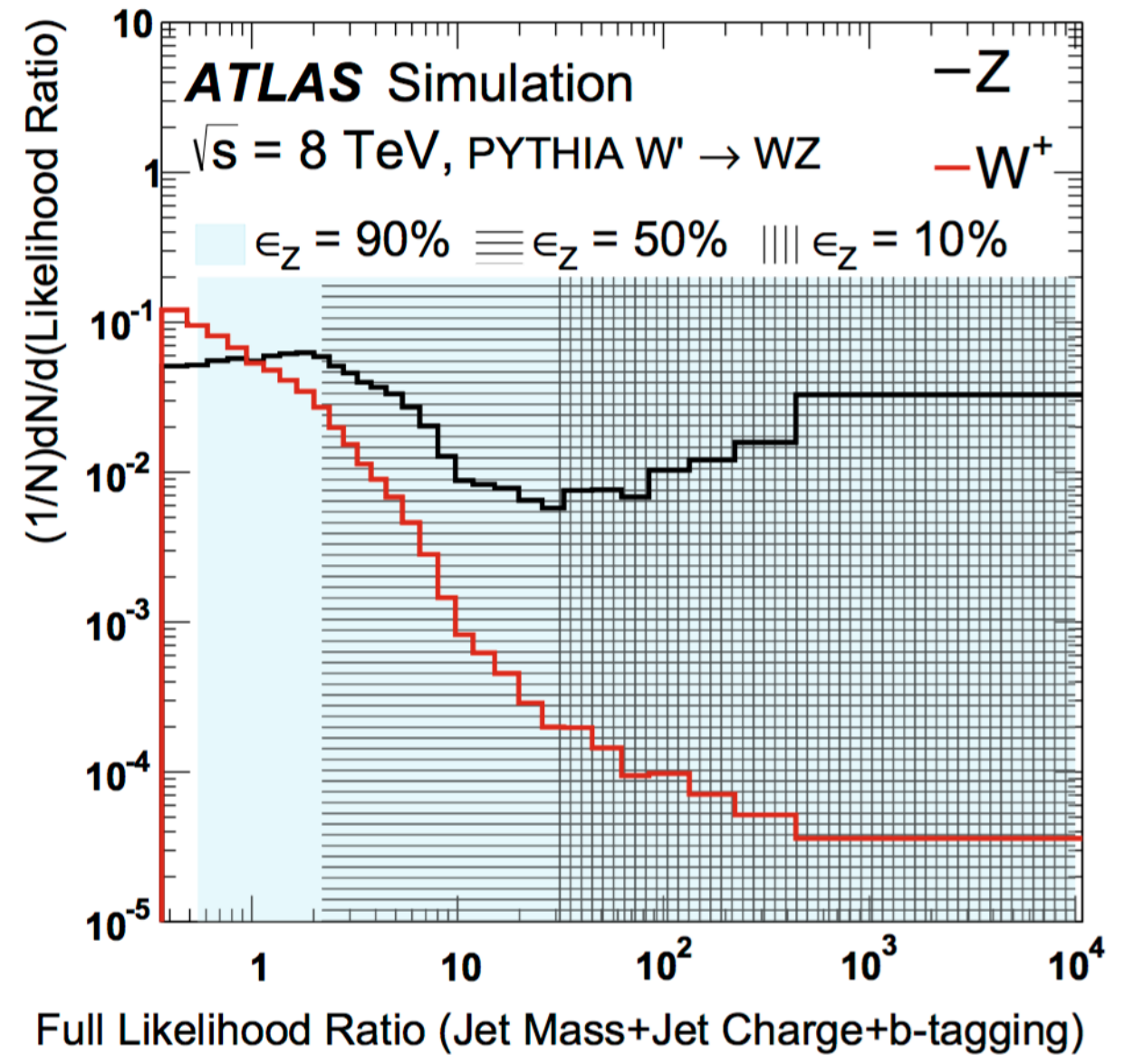
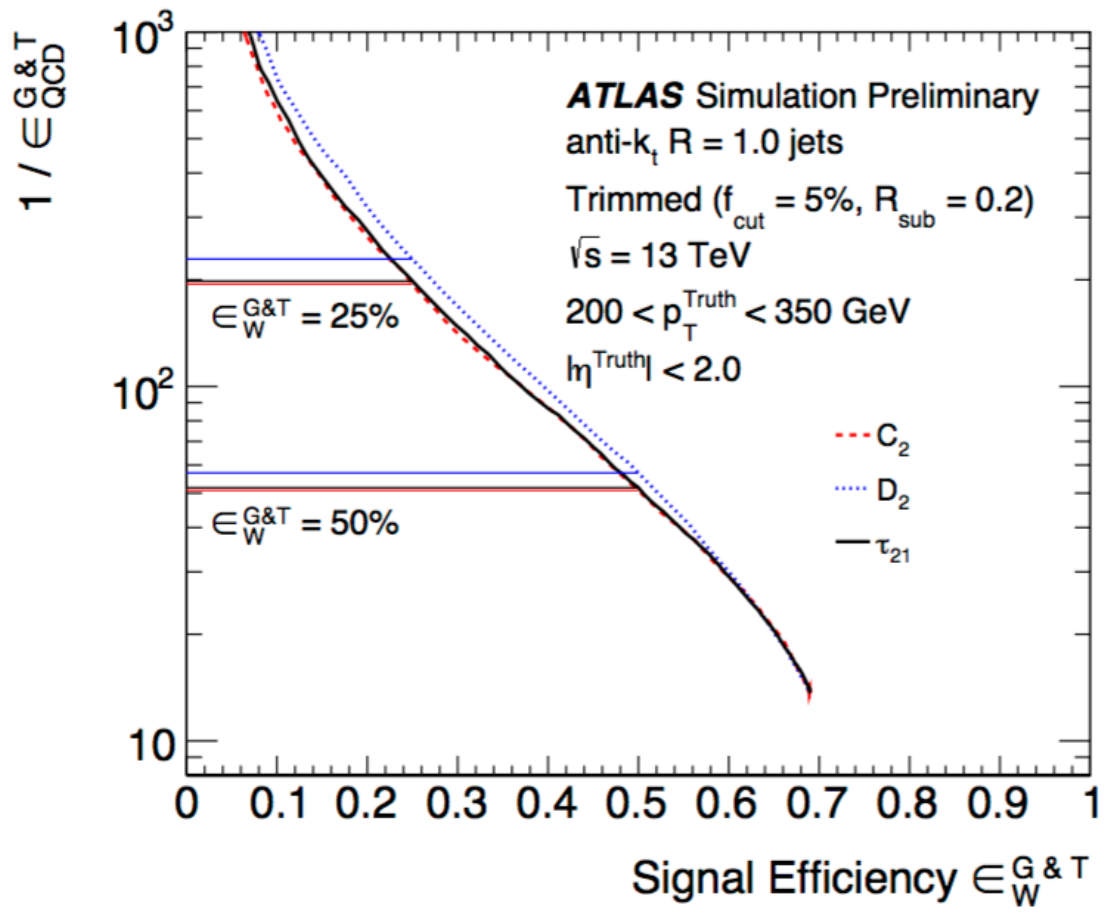


Despite resolution worsening at detector level, the jet mass is still a useful variable for distinguishing between W- and Z-jet

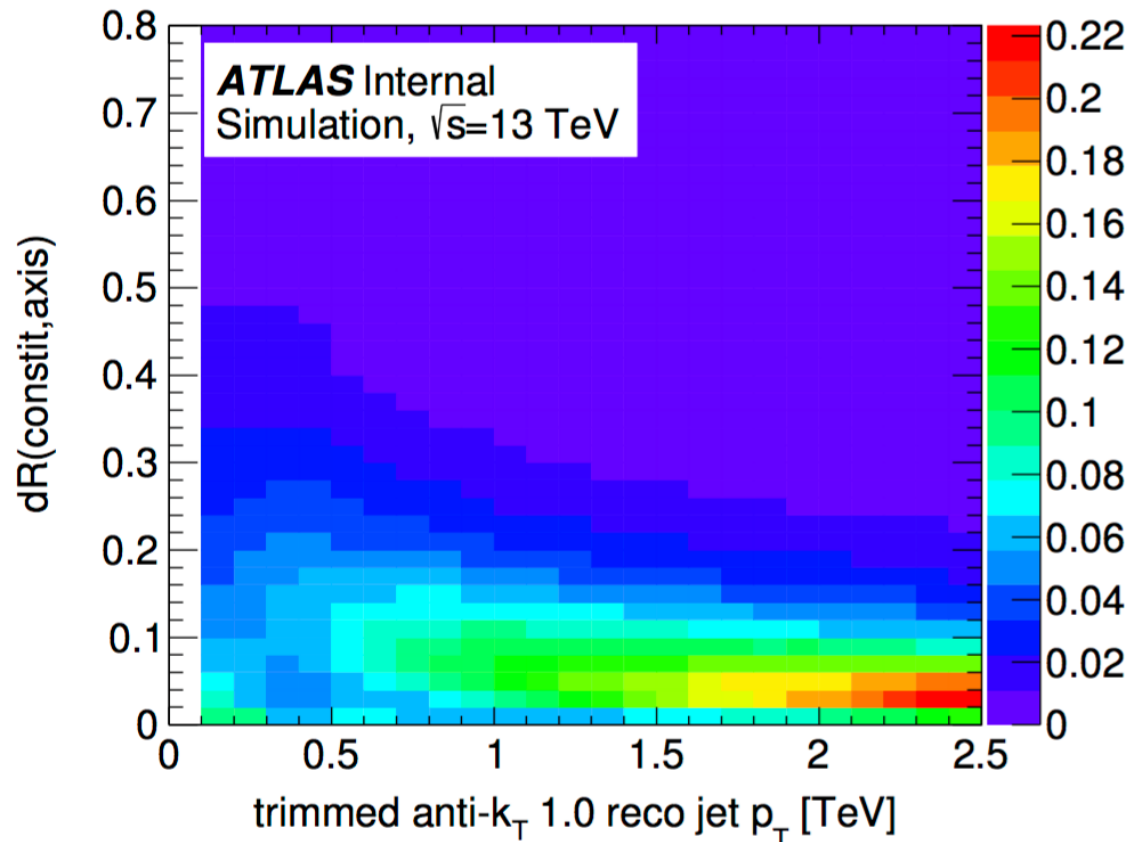




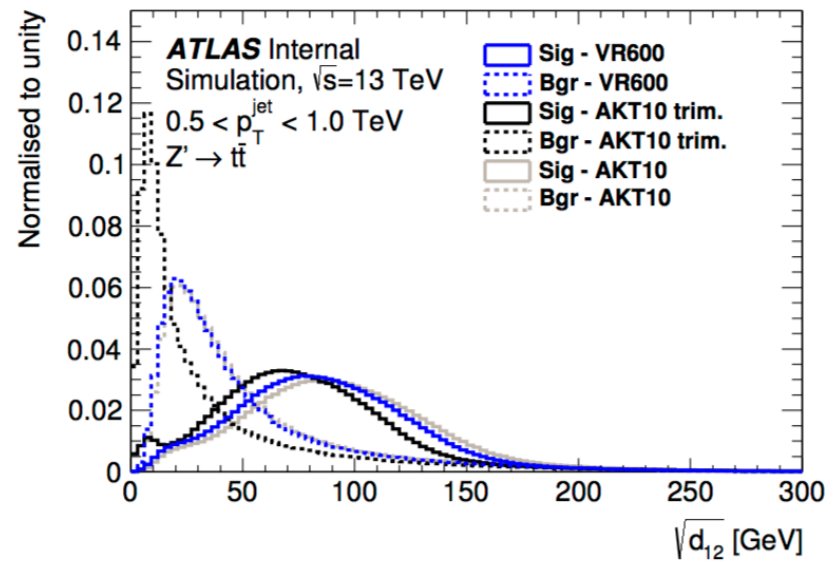
## Systematic uncertainties on efficiency



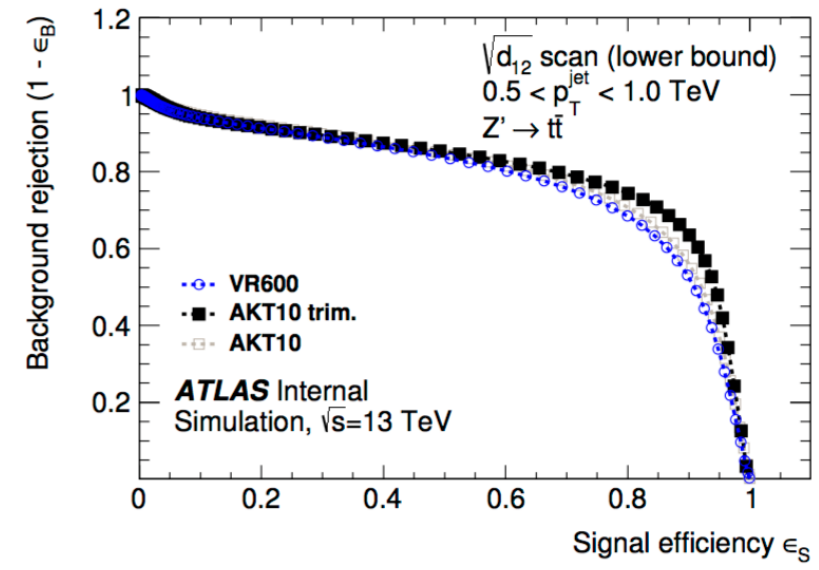
- Size of a variable- $R$  jet depends on its transverse momentum, making these jets ideally suited for boosted decay topologies - separation between decay products of a heavy particle decreases with momentum
- Significant performance improvements are found with Variable- $R$  jets for a number of discriminating variables, most notably the jet mass



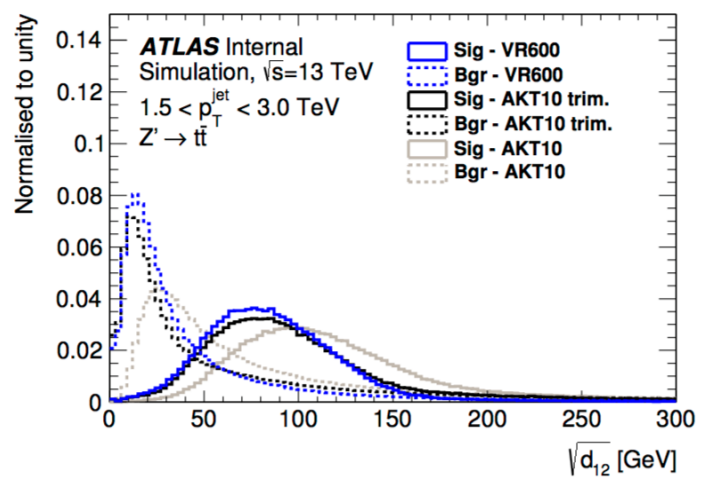
- Can build Variable- $R$  versions of the anti- $k_T$ , the C/A and the  $k_T$  algorithm.
- At high  $p_T$ , most of the jet constituents are at very small  $dR$  values with respect to the jet axis



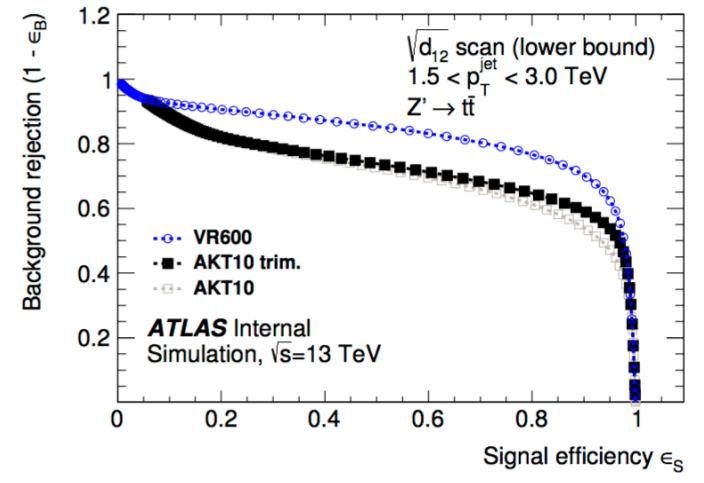
(a)  $0.5 \text{ TeV} < p_T^{\text{jet}} < 1.0 \text{ TeV}$



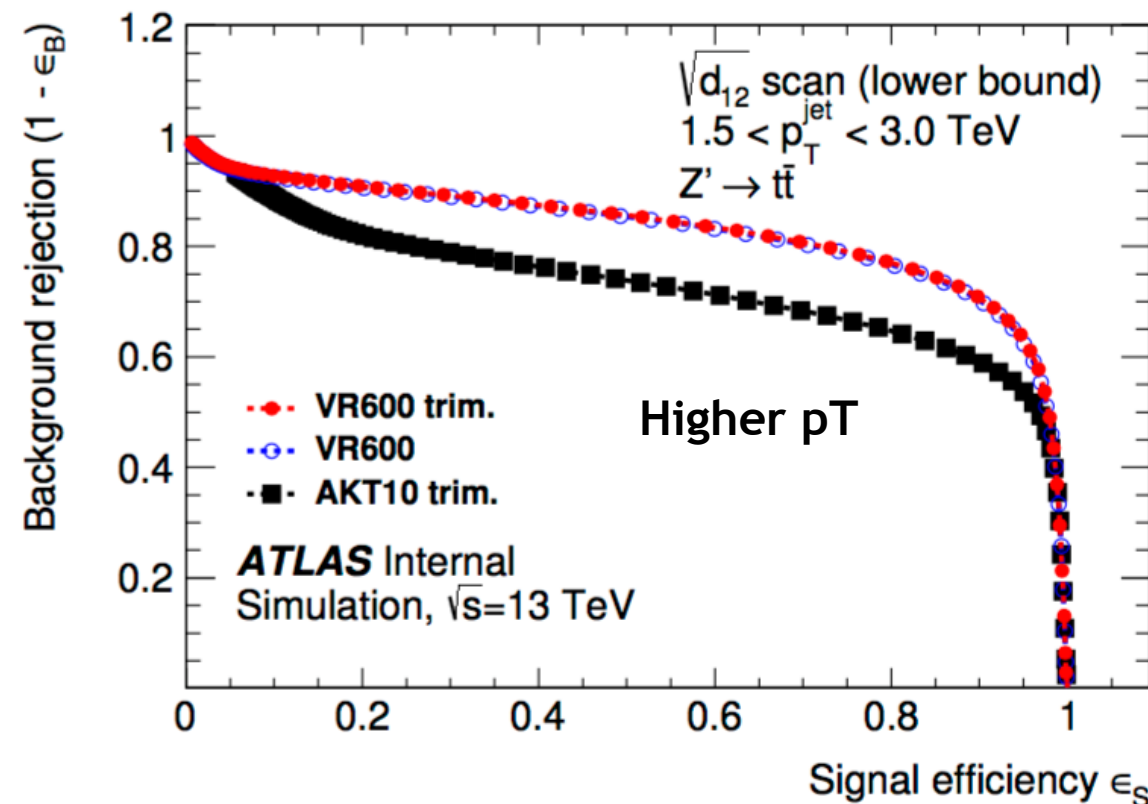
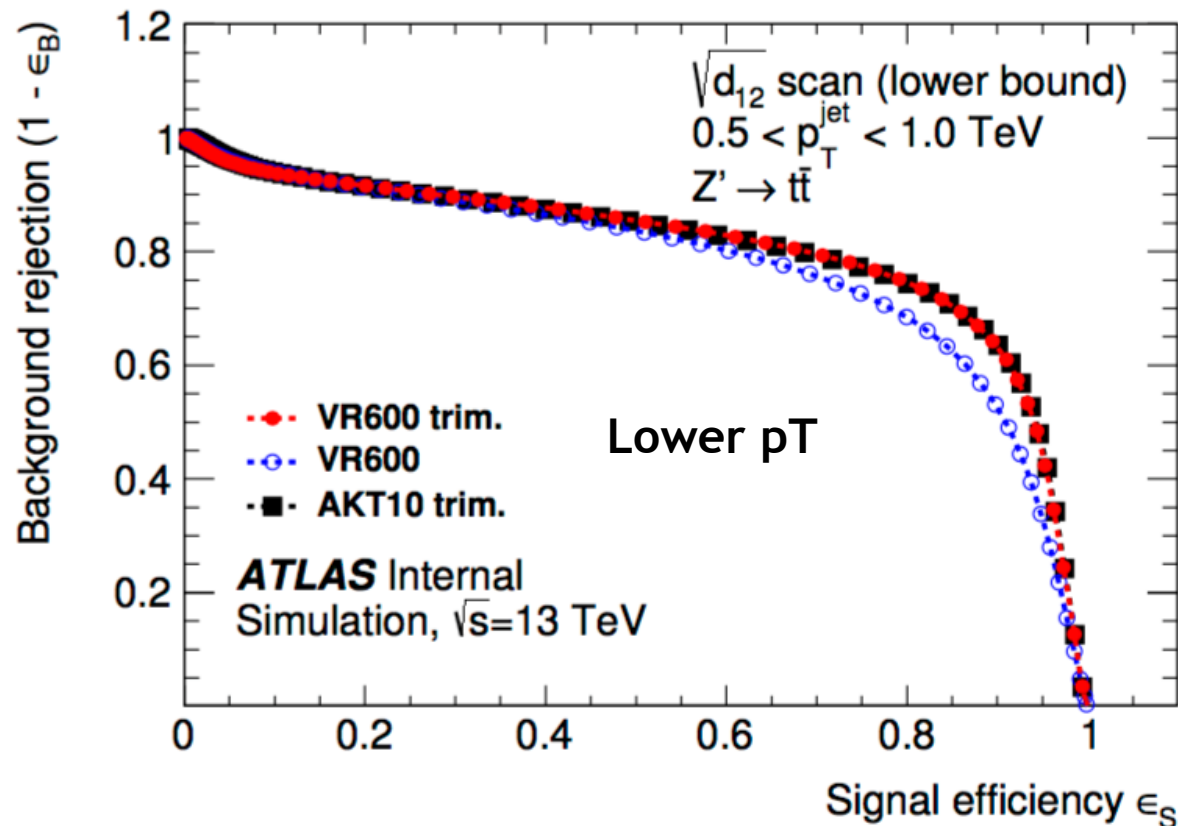
(b)  $0.5 \text{ TeV} < p_T^{\text{jet}} < 1.0 \text{ TeV}$



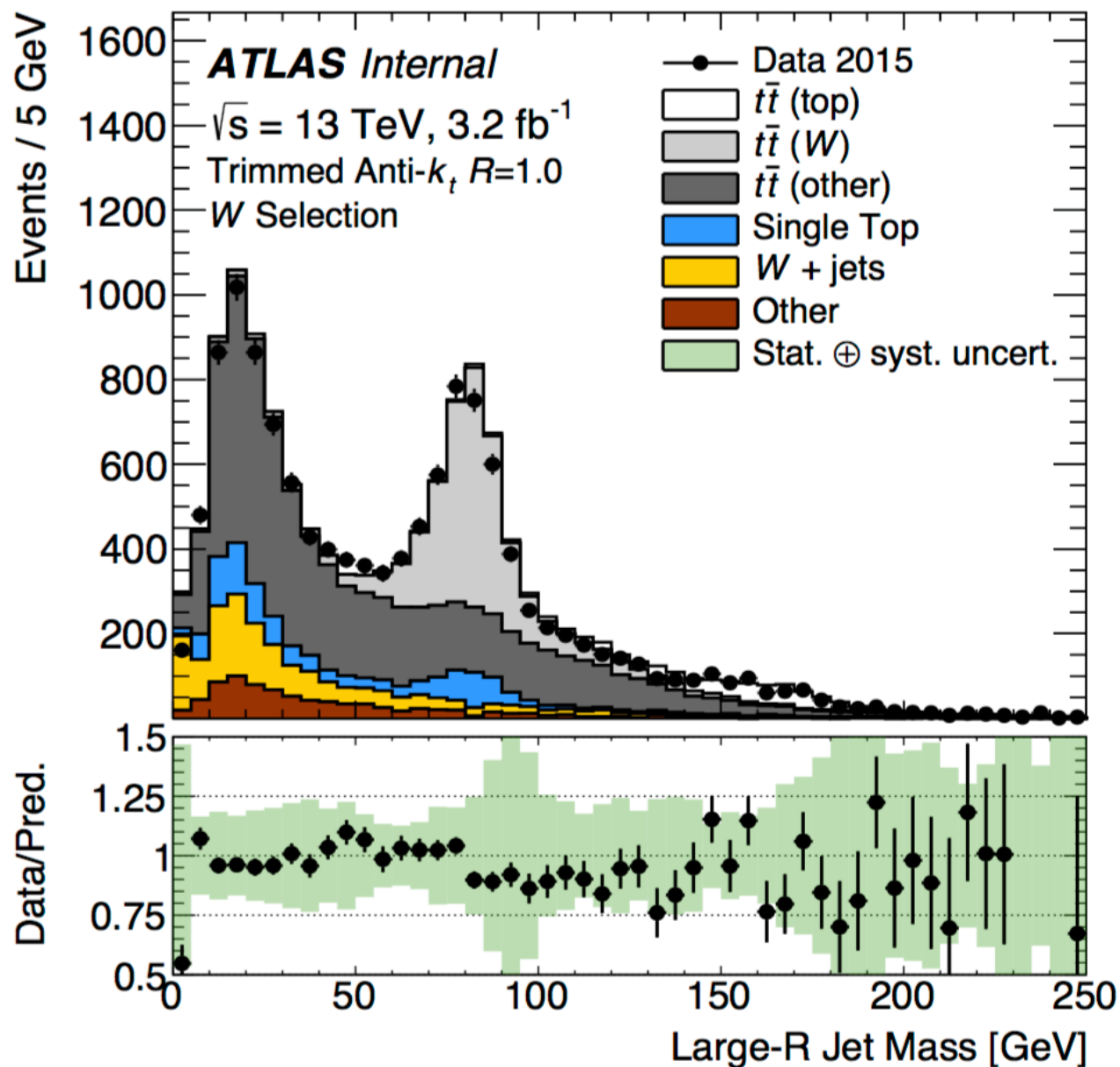
(e)  $1.5 \text{ TeV} < p_T^{\text{jet}} < 3.0 \text{ TeV}$



(f)  $1.5 \text{ TeV} < p_T^{\text{jet}} < 3.0 \text{ TeV}$



- Trimmed VR jets behave similarly to AntiKt  $R = 1.0$  jets in the low-Pt regime, while outperforming AntiKt  $R = 1.0$  jets in the medium- and high-Pt regions.
- With the two-sided mass window cut, variable-R jets has slightly worse performance than the standard W-tagger.



- Lepton+jets selection in  $t\bar{t}$  events
- Leading- $p_T$  jet mass in 13 TeV data and MC simulation.
- Anti- $k_T$   $R=1$  jets calibrated at LCW+JES scale
- Trimmed jet:  $f_{\text{cut}}=5\%$ ,  $R_{\text{subjet}}=0.2$
- $T\bar{t}$  sample is split into a “W -matched” part by requiring the decay products of W boson to be within the large-R jet.
- “top-matched” part requires the  $q\bar{q}b$  particles from the top-decay to be within the large-R jet.
- In addition to enhance purity of jets matched to a W boson, the events is required to have a  $R=0.4$  b-tagged jet but outside the large-R jet cone.
- Systematic uncertainties are composed of
  - scale and resolution uncertainties on large-R jet energy, mass and substructure observables
  - $T\bar{t}$  modeling uncertainties on parton shower, hadronization model
  - Initial-state and final-state radiation.
  - Subdominant scale and resolution uncertainties on small-R jet energy and lepton kinematics