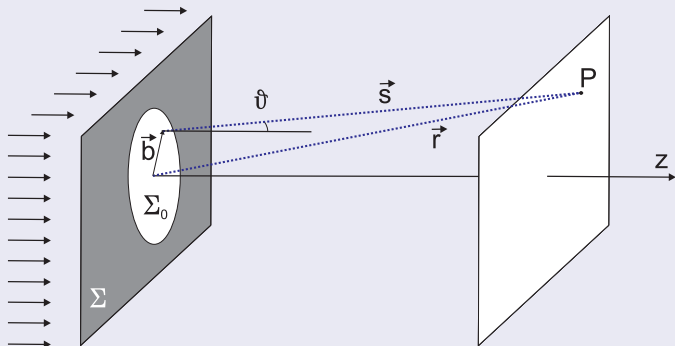


The Bialas-Bzdak model of elastic pp scattering

Results with proton-antiproton elastic scattering data

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Light scattered by a hole Σ_0


The Kirchhoff integral (U_0 , k are the amplitude and momentum of the incoming wave):

$$U(x, y, z) = -\frac{ik}{4\pi} U_0 \int_{\Sigma_0} d^2\vec{b} (1 + \cos\vartheta) \frac{e^{iks}}{s} \quad (1)$$

From the Kirchhof integral to Fraunhofer diffraction

Far from the obstacle (Fraunhofer limit):

$$U(x, y, z) = -\frac{ik}{2\pi} U_0 \frac{e^{ikr}}{r} \int d^2\vec{b} S(\vec{b}) e^{-i\vec{q}\cdot\vec{b}} \quad (\text{disk } \Sigma_0) \quad (2)$$

where $S(\vec{b}) = 1 - \Gamma(\vec{b})$ and the profile function:

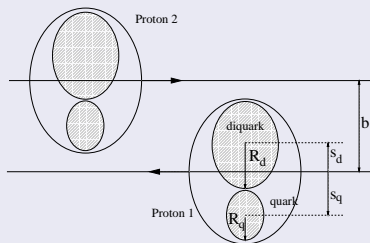
$$\Gamma(\vec{b}) = \begin{cases} 1 & \text{if } \vec{b} \in \Sigma_0, \\ 0 & \text{otherwise.} \end{cases} \quad (3)$$

The scattering amplitude:

$$f(q) = ik \int_0^\infty db b J_0(qb) \Gamma(b) \quad (4)$$

$$\frac{d\sigma}{d\Omega} = |f(\vec{q})|^2 \quad (5)$$

High-energy particle diffraction: light waves \rightarrow wave functions, screen \rightarrow silicon detector



The scattering amplitude in the eikonal approximation:

$$T(s, \Delta) = 2\pi \int_0^{\infty} J_0(\Delta \cdot b) t_{\text{el}}(s, b) b db, \quad (6)$$

where $t_{\text{el}}(s, b) = i [1 - e^{-\Omega(s, b)}]$ and $t \simeq -\Delta^2$. See Eq. (4)!

Cross-sections and others

Differential cross-section:

$$\frac{d\sigma}{dt} = \frac{1}{4\pi} |T(s, \Delta)|^2 \quad (7)$$

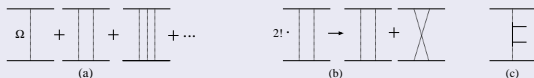
Total cross-section:

$$\sigma_{\text{tot}} = 2 \operatorname{Im} T(s, \Delta)|_{t=0} \quad (8)$$

The real to imaginary part ratio:

$$\rho = \frac{\operatorname{Re} T(s, 0)}{\operatorname{Im} T(s, 0)} \quad (9)$$

The (seemingly!) simple equations in QFT: multi meson exchange



The profile function in the BB model:

Partial wave equation \rightarrow impact parameter b ! Unitarity:

$$2 \operatorname{Im} t_{\text{el}}(s, b) = |t_{\text{el}}(s, b)|^2 + \tilde{\sigma}_{\text{inel}}(s, b) \quad (10)$$

The profile function becomes the function of the inelastic scattering probability:

$$t_{\text{el}}(s, b) = i \left[1 - \sqrt{1 - \tilde{\sigma}_{\text{inel}}(s, b)} \right] \quad (11)$$

With real part:

$$t_{\text{el}}(s, b) = i \left[1 - e^{-i \operatorname{Im} \Omega(s, b)} \sqrt{1 - \tilde{\sigma}_{\text{inel}}(s, b)} \right] \quad (12)$$

The inelastic scattering probability function in the BB model:

The average (over configurations) of inelastic probabilities:

$$\tilde{\sigma}_{\text{inel}}(b) = \left\langle \sigma(h; \vec{b}) \right\rangle_H = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} dh p(h) \cdot \sigma(h; \vec{b}) \quad (13)$$

Inelastic scatt. prob. of a composite particle (Glauber):

$$\sigma(h; \vec{b}) = 1 - \prod_a \prod_b \left[1 - \sigma_{ab}(\vec{b} + \vec{s}'_a - \vec{s}_b) \right] \quad (14)$$

and the elementary inelastic scattering probabilities:

$$\sigma_{ab}(\vec{s}) = A_{ab} e^{-s^2/S_{ab}^2}, \quad S_{ab}^2 = R_a^2 + R_b^2 \quad (15)$$

Distributions

The probability density of collisions:

$$p(h) = D(\vec{s}_q, \vec{s}_d) \cdot D(\vec{s}'_q, \vec{s}'_d). \quad (16)$$

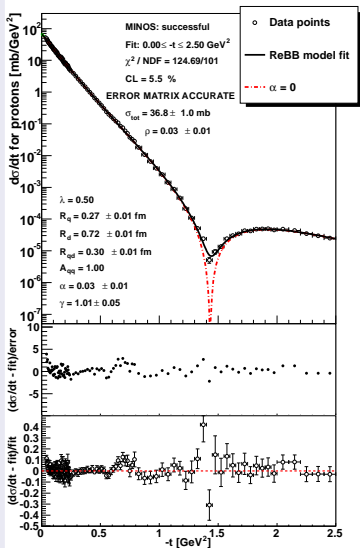
Quark-diquark probability distribution:

$$D(\vec{s}_q, \vec{s}_d) = \frac{1 + \lambda^2}{R_{qd}^2 \pi} e^{-(s_q^2 + s_d^2)/R_{qd}^2} \delta^2(\vec{s}_d + \lambda \vec{s}_q), \quad (17)$$

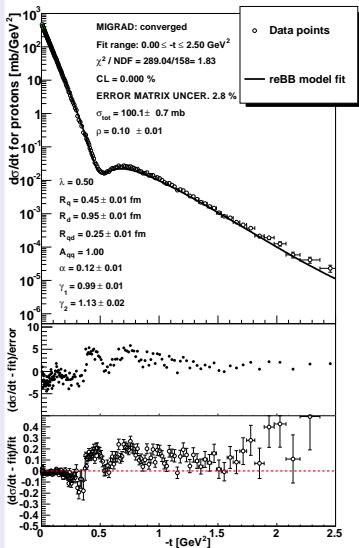
$$\lambda = \frac{m_q}{m_d} \quad (18)$$

ISR and TOTEM data fits (ReBB paper with Tamás and Máté)

$p+p \rightarrow p+p$, diquark as a single entity at $\sqrt{s}=23.5$ GeV

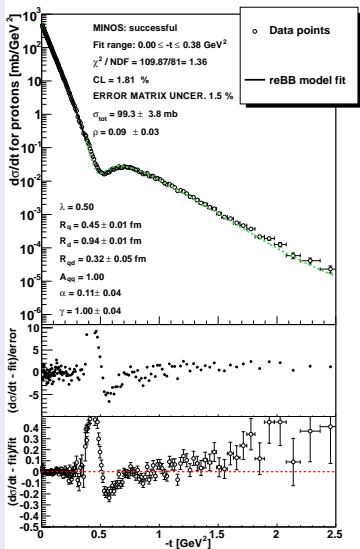


$p+p \rightarrow p+p$, diquark as a single entity at $\sqrt{s}=7000.0$ GeV

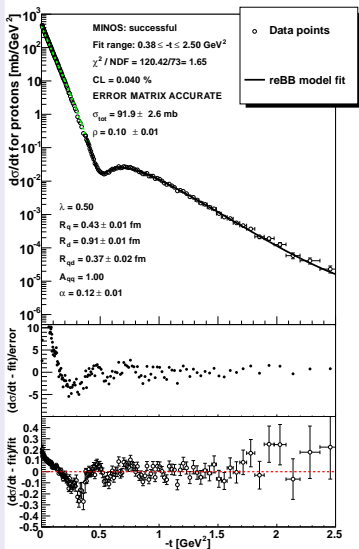


TOTEM data fits with two ranges

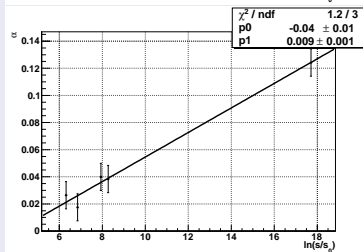
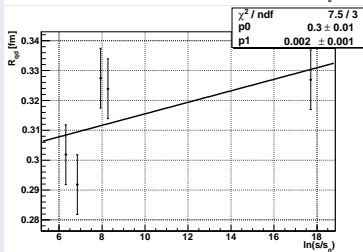
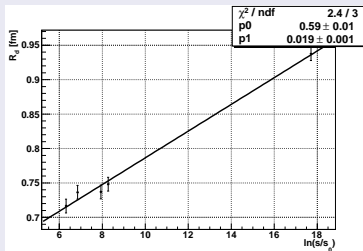
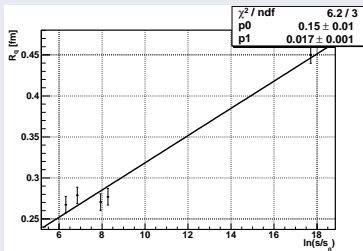
$p+p \rightarrow p+p$, diquark as a single entity at $\sqrt{s}=7000.0$ GeV



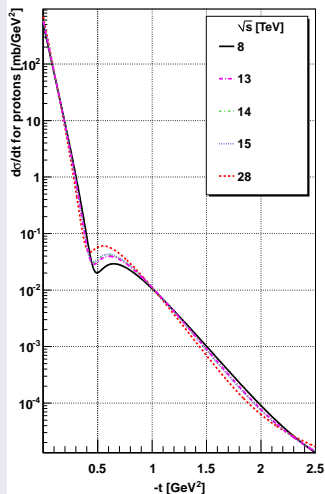
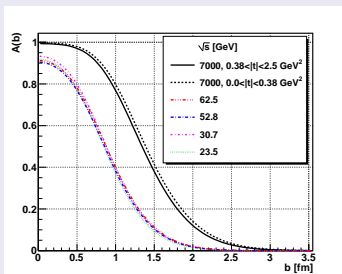
$p+p \rightarrow p+p$, diquark as a single entity at $\sqrt{s}=7000.0$ GeV



The excitation functions of the parameters



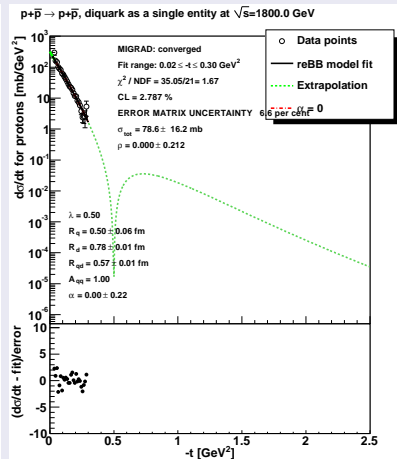
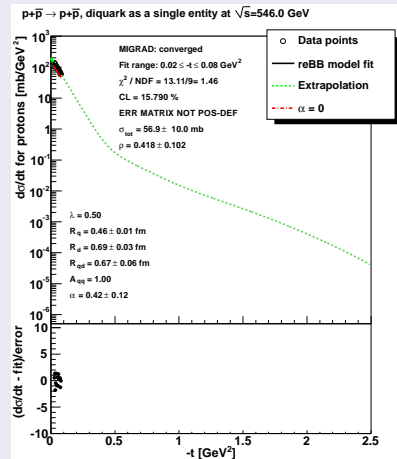
Shadow pr. + extrapolations to future LHC energies and beyond



The increasing total cross-section with increasing \sqrt{s}

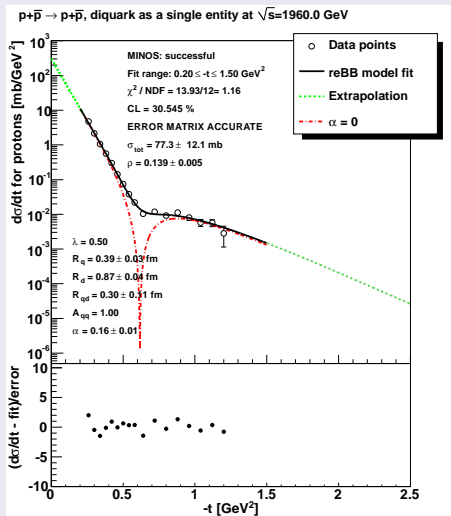
- Pomeron trajectory !? BFKL dynamics: emission and absorption of parton cascades: leads to diffusion like expansion of the proton, see the Lund model (next talk)
- QED with massive photon in 2 space dimensions (b space): expanding fermion (proton)

All the usual BB fit parameters are free ($\sqrt{s} = 546 \text{ GeV}, 1.8 \text{ TeV}$)



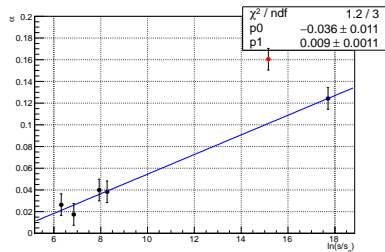
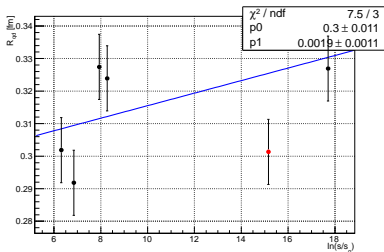
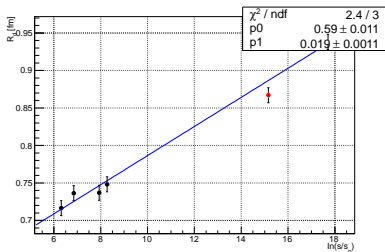
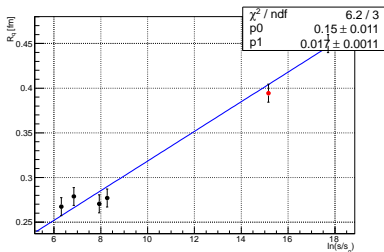
There are not enough points to pin down the shape.

All the usual BB fit parameters are free (1.96 TeV)

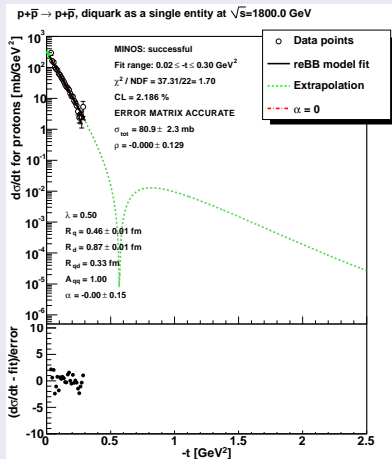
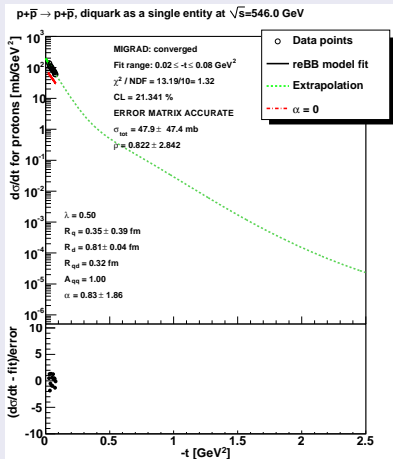


Ok.

The good fit at $\sqrt{s} = 1.96$ TeV compared with the extrapolations based only on pp fits of our ReBB paper

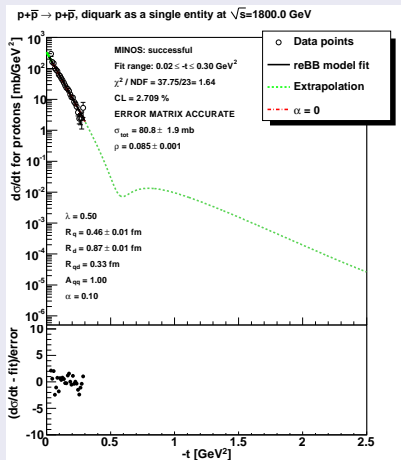
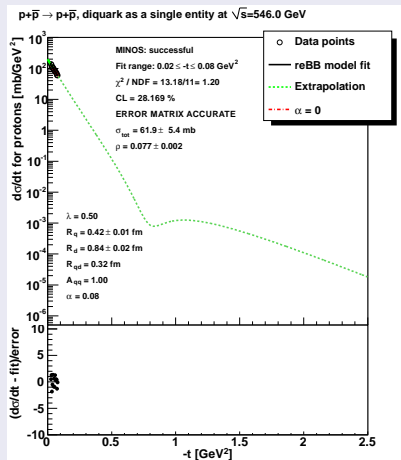


R_{qd} is fixed to the extrapolated value



Better fits.

Both R_{qd} and α is fixed to the extrapolated value



Ok.

Conclusions

- BB model is able to interpret $p\bar{p}$, $p\bar{p}$ (and πp) data
- Trends in energy are consistent (more or less)
- Provides solid baseline for future physics studies