## Observables, forgó megoldások X = Z = R, Y esetében

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# Our starting equations

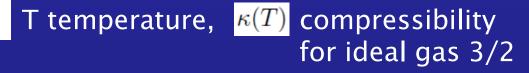
Consider the non-relativistic hydrodynamical problem, as given by the continuity, Euler and energy equations:

> $\partial_t n + \nabla \cdot (n\mathbf{v}) = 0,$  $\partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -(\nabla p)/(mn),$  $\partial_t \epsilon + \nabla \cdot (\epsilon \mathbf{v}) = -p \nabla \cdot \mathbf{v},$

v flow velocity field, n particle number, p pressure, m particle mass,  $\epsilon$  energy density

## The closing Equation of State

p = nT,  $\epsilon = \kappa(T)nT$ ,



# The Ansatz for the solution

self-similar, ellipsoidally symmetric density and flow profiles

$$T(\mathbf{r},t) = T(t),$$
  

$$n(\mathbf{r},t) = n_0 \frac{V_0}{V} \exp\left\{-\frac{r_x^2}{2X^2} - \frac{r_y^2}{2Y^2} - \frac{r_z^2}{2Z^2}\right\}$$
  

$$\mathbf{v}(\mathbf{r},t) = \left(\frac{\dot{X}}{X}r_x + \omega r_z, \frac{\dot{Y}}{Y}r_y, \frac{\dot{Z}}{Z}r_z - \omega r_x\right)$$

where

$$(X, Y, Z) = (X(t), Y(t), Z(t)) \qquad V = XYZ,$$

### is known from a long time

S.V. Akkelin, T. Csörgő, B. Lukács, Yu. M. Sinyukov and M. Weiner, Phys. Lett. **B505** (2001) 64.

# The final ordinary differentialequation for the time propagation

$$X = Z \equiv R, \quad \dot{X} = \dot{Z} \equiv \dot{R},$$
$$R\ddot{R} - R^{2}\omega^{2} = \ddot{Y}Y = \frac{T}{m_{0}},$$
$$\dot{T}\frac{d}{dT}(T\kappa(T)) + T\frac{\dot{V}}{V} = 0,,$$

 $\omega(t) = \omega_0 \frac{R_0^2}{R^2(t)}.$ 

The time dependence of the temperature is

$$\frac{V_0}{V} = \exp\left[\kappa(T) - \kappa(T_0)\right] \exp\left(\int_{T_0}^T \frac{dT'}{T'} \kappa(T')\right)$$

Inital conditions

$$V_0 = V(t_0)$$
$$T_0 = T(t_0)$$

$$T = T_0 \left(\frac{V_0}{V}\right)^{1/\kappa}$$

# The observables and the new solutions

*the emission function at a constant freeze-out temperature* 

$$S(t, \mathbf{r}', \mathbf{k}') \propto e^{-\frac{(x\mathbf{k}'-m\mathbf{v}')^2}{2mT_f} - \frac{r'_x^2}{2X_f^2} - \frac{r'_y^2}{2Y_f^2} - \frac{r'_z^2}{2Z_f^2}} \delta(t - t_f)$$

#### single particle spectrum

$$S \propto \exp\left\{-\frac{T_x r_x^2}{2T_f X_f^2} + \frac{p_x \dot{X}_f - p_z \omega_f X_f}{X_f T_f} r_x - \frac{p_x^2}{2m T_f}\right\} \times \\ \times \exp\left\{-\frac{T_y r_y^2}{2T_f Y_f^2} + \frac{p_y \dot{Y}_f}{Y_f T_f} r_y - \frac{p_y^2}{2m T_f}\right\} \times \\ \times \exp\left\{-\frac{T_z r_z^2}{2T_f Z_f^2} + \frac{p_z \dot{Z}_f + p_x \omega_f Z_f}{Z_f T_f} r_z - \frac{p_z^2}{2m T_f}\right\}$$

# The observables and the new solutions

Using the new notation:

$$T_{x} = T_{f} + m \left( \dot{X}_{f}^{2} + \omega_{f}^{2} Z_{f}^{2} \right),$$
  

$$T_{y} = T_{f} + m \dot{Y}_{f}^{2},$$
  

$$T_{z} = T_{f} + m \left( \dot{Z}_{f}^{2} + \omega_{f}^{2} X_{f}^{2} \right).$$
  
new terms from the rotation

# The observables II

$$\begin{split} C(\mathbf{K}',\mathbf{q}') &= 1 + \lambda \exp\left(-q_x'^2 R_x'^2 - q_y'^2 R_y'^2 - q_z'^2 R_z'^2\right), \\ \mathbf{K}' &= \mathbf{K}'_{12} = 0.5(\mathbf{k}'_1 + \mathbf{k}'_2), \\ \mathbf{q}' &= \mathbf{q}'_{12} = \mathbf{k}'_1 - \mathbf{k}'_2 = (q_x', q_y', q_z'), \\ R_x'^{-2} &= X_f^{-2} \left(1 + \frac{m}{T_f} (\dot{X}_f^2 + Y_f^2 \omega^2)\right), \\ R_y'^{-2} &= Y_f^{-2} \left(1 + \frac{m}{T_f} \dot{Y}_f^2\right), \\ R_z'^{-2} &= Z_f^{-2} \left(1 + \frac{m}{T_f} (\dot{Z}_f^2 + X_f^2 \omega^2)\right). \end{split}$$

the two-particle BECF new terms from the rotation

$$C_{2}(\mathbf{K}, \mathbf{q}) = 1 + \lambda \exp\left(-\sum_{i,j=s,o,l} q_{i}q_{j}R_{ij}^{2}\right),$$

$$R_{s}^{2} = R_{y}^{\prime 2}\cos^{2}\phi + R_{x}^{2}\sin^{2}\phi,$$

$$R_{o}^{2} = R_{x}^{2}\cos^{2}\phi + R_{y}^{\prime 2}\sin^{2}\phi + \beta_{t}^{2}\Delta t^{2},$$

$$R_{l}^{2} = R_{z}^{\prime 2}\cos^{2}\theta + R_{x}^{\prime 2}\sin^{2}\theta + \beta_{l}^{2}\Delta t^{2},$$

$$R_{ol}^{2} = 0 + \beta_{t}\beta_{l}\Delta t^{2},$$

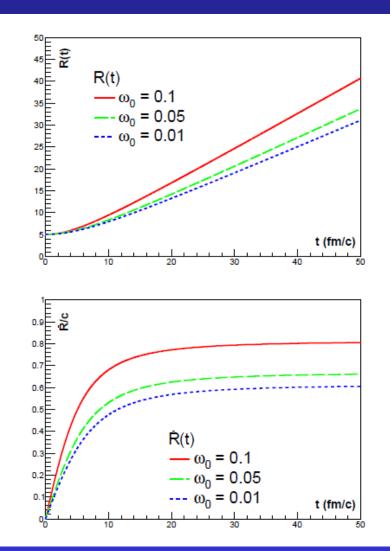
$$R_{os}^{2} = (R_{x}^{2} - R_{y}^{\prime 2})\cos\phi\sin\phi,$$

$$R_{sl}^{2} = 0.$$

the transverse momentum of the measured pair

# Results I

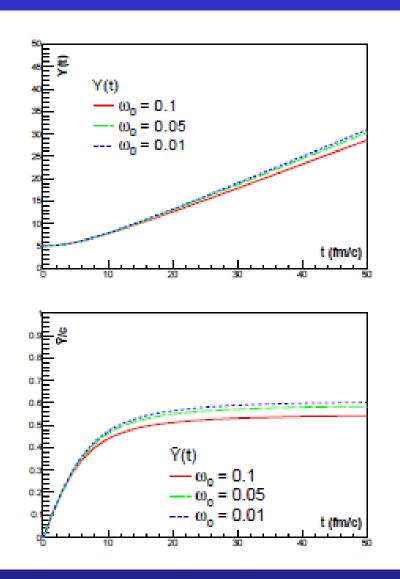
#### *Time evolutions of the axes* X(t) = Z(t) = R(t)



Inital conditions:  $R_0 = 5 \text{ fm},$  $Y_0 = 5 \text{ fm}$ 

## Results II

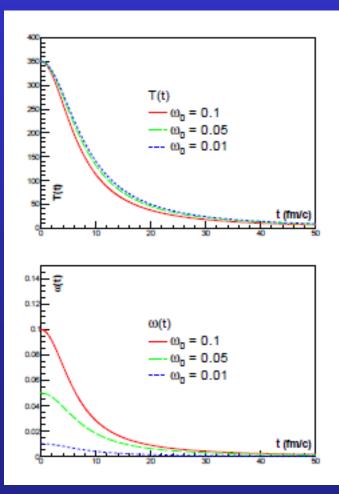
## *Time evolutions of the axes Y(t)*



Inital conditions:  $R_0 = 5 \text{ fm},$  $Y_0 = 5 \text{ fm}$ 

# Results III

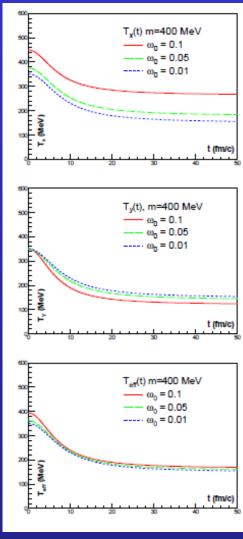
*The time evolution of the temperature and the angular velocity* 



Inital conditions:  $R_0 = 5 \text{ fm},$   $Y_0 = 5 \text{ fm}$  Kappa = 3/2  $T_f = 140 \text{ MeV}$ at 8 - 10 fm/c

# Results IV

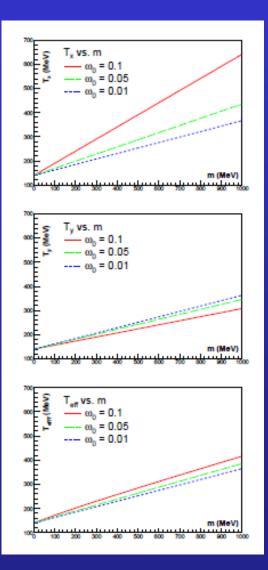
# Freeze-out time dependence of the slope parameters $T_x = T_z$



Inital conditions:  $R_0 = 5 \text{ fm},$   $Y_0 = 5 \text{ fm}$  Kappa = 3/2 $T_f = 140 \text{ MeV}$ 

# Results V

### The slope parameters $T_x = T_z$ and the effective temperature

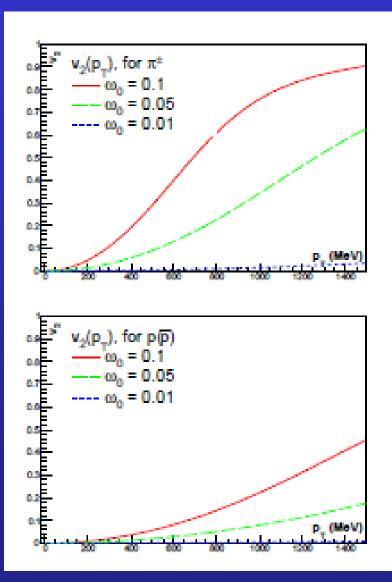


$$\frac{1}{T_{\rm eff}} = \frac{1}{2} \left( \frac{1}{T_x} + \frac{1}{T_y} \right)$$

Inital conditions:  $R_0 = 5 \text{ fm},$   $Y_0 = 5 \text{ fm}$  Kappa = 3/2 $T_f = 140 \text{ MeV}$ 

# Results VI

## $V_2$ for pions (upper) and protons (lower)

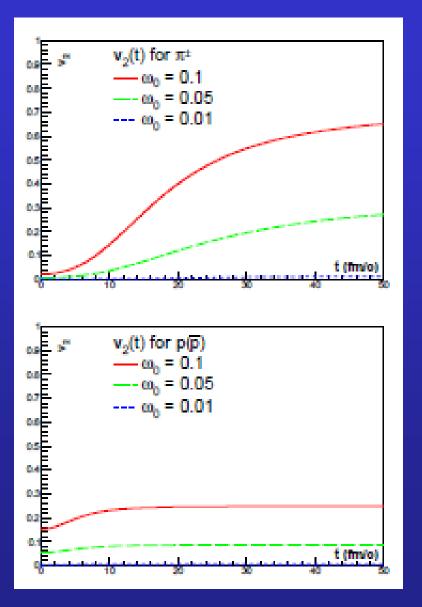


Inital conditions:  $R_0 = 5 \text{ fm},$   $Y_0 = 5 \text{ fm}$  Kappa = 3/2 $T_f = 140 \text{ MeV } 0$ 

$$\frac{dn}{2\pi p_t dp_t dp_z} \propto \exp\left(-\frac{p_t^2}{2mT_{\text{eff}}} - \frac{p_z^2}{2mT_z}\right) I_0(w)$$
$$\frac{dn}{dp_z p_t dp_t d\phi} = \frac{dn}{2\pi dp_z p_t dp_t} \left[1 + 2\sum_{n=1}^{\infty} v_n \cos(n\phi)\right]$$

$$v_{2n+1} = 0, \quad v_{2n} = \frac{I_n(w)}{I_0(w)}.$$

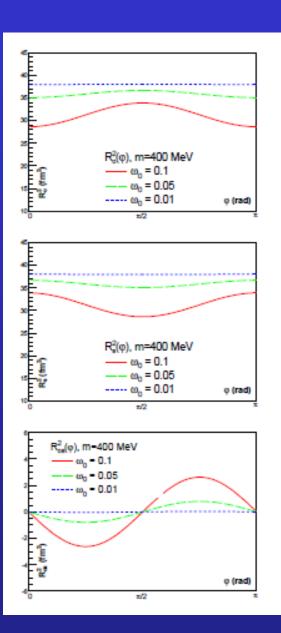
# Results VII



The freeze-out time dependence of  $v_2$  for pions (upper) and for protons (lower) at fixed  $p_T$ (300 MeV/c) and (1000 MeV/c)

Inital conditions:  $R_0 = 5 fm$ ,  $Y_0 = 5 fm$ Kappa = 3/2

# Results VIII



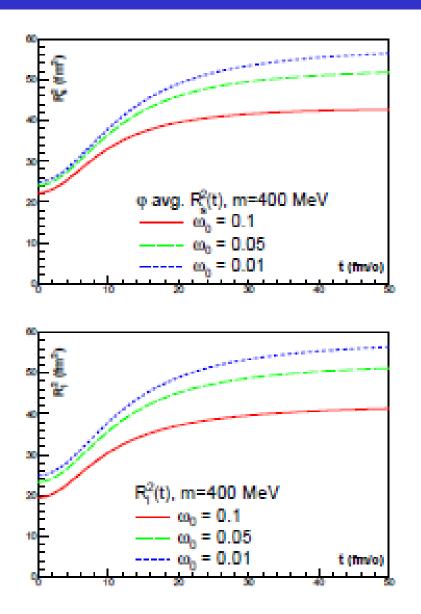
The HBT radius parameters  $R_0^2$  (upper panel)  $R_s^2$  (middel panel)  $R_{os}^2$  (lower panel)

$$C (\mathbf{K}, \mathbf{q}) = 1 + \lambda \exp \left\{ -q_x^2 R_x^2 - q_y^2 R_y^2 - q_z^2 R_z^2 \right\}$$
$$\mathbf{K} = \mathbf{K}_{12} = \frac{1}{2} \left( \mathbf{p}_1 + \mathbf{p}_2 \right),$$
$$\mathbf{q} = \mathbf{q}_{12} = \left( \mathbf{k}_1 - \mathbf{k}_2 \right),$$

$$\begin{aligned} C_{2} \left( \mathbf{K}, \mathbf{q} \right) &= 1 + \lambda \exp \left( -\sum_{i,j=s,o,1} q_{i}q_{j}R_{ij}^{2} \right) \\ R_{s}^{2} &= R_{y}^{2}\cos^{2}\phi + R_{x}^{2}\sin^{2}\phi \,, \\ R_{o}^{2} &= R_{x}^{2}\cos^{2}\phi + R_{y}^{2}\sin^{2}\phi + \beta_{t}^{2}\Delta t^{2} \,, \\ R_{1}^{2} &= R_{x}^{2} + \beta_{1}^{2}\Delta t^{2} \,, \\ R_{ol}^{2} &= \beta_{t}\beta_{l}\Delta t^{2} \,, \\ R_{os}^{2} &= \left( R_{x}^{2} - R_{y}^{2} \right)\cos\phi\sin\phi \,, \\ R_{sl}^{2} &= 0 \,. \end{aligned}$$

for m = 400 MeV

# Results IX



Freeze-out time dependence of the azimuthal average of  $R_2^s$  (upper) and  $R_2^l$  (lower)

 $m = 400 \; MeV$ 

Thank You for Your Attention!

Questions, Remarks, Comments?...