

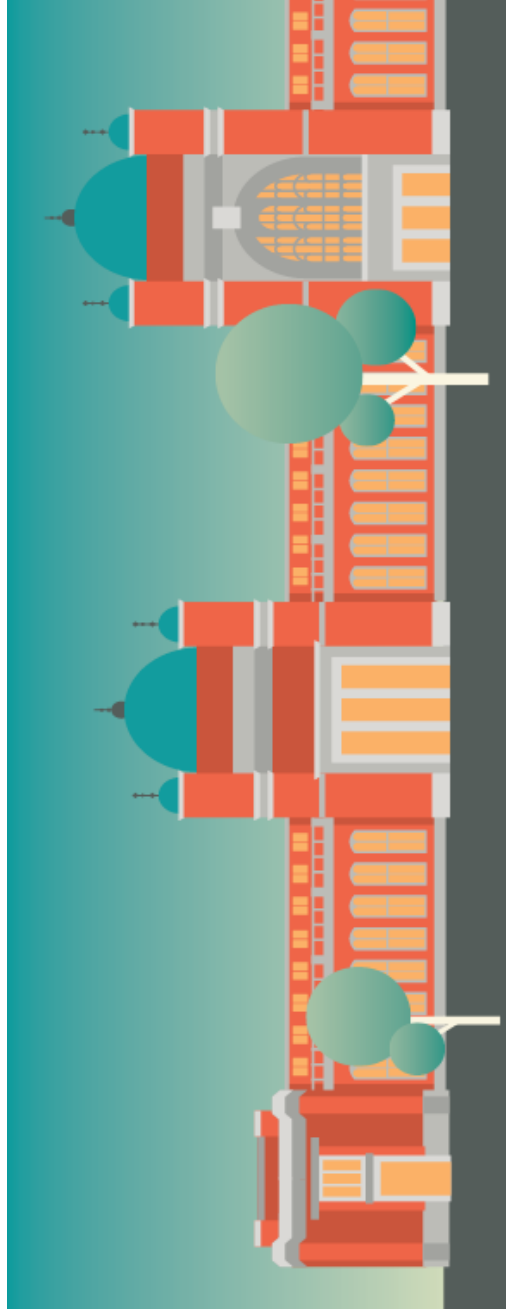
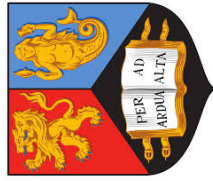
---

# Probing CPT in transitions with entangled neutral kaons



Antonio Di Domenico

Dipartimento di Fisica, Sapienza Università di Roma  
and INFN sezione di Roma, Italy



**KAON 2016 Conference**  
**University of Birmingham, UK, September 14-17, 2016**

---

# Testing CPT: introduction

The three discrete symmetries of QM, C (charge conjugation:  $q \rightarrow -q$ ), P (parity:  $x \rightarrow -x$ ), and T (time reversal:  $t \rightarrow -t$ ) are known to be violated in nature both singly and in pairs. Only CPT appears to be an exact symmetry of nature.

**CPT theorem holds for any QFT formulated on flat space-time which assumes: (1) Lorentz invariance (2) Locality (3) Unitarity (i.e. conservation of probability).**

Extension of CPT theorem to a theory of quantum gravity far from obvious. (e.g. CPT violation appears in several QG models)

huge effort in the last decades to study and shed light on QG phenomenology  $\Rightarrow$  Phenomenological CPTV parameters to be constrained by experiments

Consequences of CPT symmetry: equality of masses, lifetimes,  $|q|$  and  $|\mu|$  of a particle and its anti-particle.

Neutral meson systems offer unique possibilities to test CPT invariance; e.g. taking as figure of merit the fractional difference between the masses of a particle and its anti-particle:

neutral K system	neutral B system	proton- anti-proton
$\left  m_{K^0} - m_{\bar{K}^0} \right  / m_K < 10^{-18}$	$\left  m_{B^0} - m_{\bar{B}^0} \right  / m_B < 10^{-14}$	$\left  m_p - m_{\bar{p}} \right  / m_p < 10^{-8}$

Other interesting CPT tests: e.g. the study of anti-hydrogen atoms, etc..

# Testing CPT: introduction

The three discrete symmetries of QM, C (charge conjugation:  $q \rightarrow -q$ ), P (parity:  $x \rightarrow -x$ ), and T (time reversal:  $t \rightarrow -t$ ) are known to be violated in nature both singly and in pairs. Only CPT appears to be an exact symmetry of nature.

**CPT theorem holds for any QFT formulated on flat space-time which assumes:**  
**(1) Lorentz invariance (2) Locality (3) Unitarity (i.e. conservation of probability).**

Extension of CPT theorem to a theory of quantum gravity far from obvious.  
(e.g. CPT violation appears in several QG models)

huge effort in the last decades to study and shed light on QG phenomenology  
 $\Rightarrow$  Phenomenological CPTV parameters to be constrained by experiments

Consequences of CPT symmetry: equality of masses, lifetimes,  $|q|$  and  $|\mu|$  of a particle and its anti-particle.

Neutral meson systems offer unique possibilities to test CPT invariance; e.g. taking as figure of merit the fractional difference between the masses of a particle and its anti-particle:

neutral K system

$$\left| \frac{m_{K^0} - m_{\bar{K}^0}}{m_K} \right| < 10^{-18}$$

neutral B system

$$\left| \frac{m_{B^0} - m_{\bar{B}^0}}{m_B} \right| < 10^{-14}$$

proton- anti-proton

$$\left| \frac{m_p - m_{\bar{p}}}{m_p} \right| < 10^{-8}$$

Other interesting CPT tests: e.g. the study of anti-hydrogen atoms, etc..

## CPT violation: standard picture

$$|K_{S,L}\rangle \propto \left[ (1 + \varepsilon_{S,L}) |K^0\rangle \pm (1 - \varepsilon_{S,L}) |\bar{K}^0\rangle \right]$$

CP violation:

$$\varepsilon_{S,L} = \varepsilon \pm \delta$$

T violation:

$$\varepsilon = \frac{H_{12} - H_{21}}{2(\lambda_S - \lambda_L)} = \frac{-i\Im M_{12} - \Im\Gamma_{12}/2}{\Delta m + i\Delta\Gamma/2}$$

**CPT violation:**

$$\delta = \frac{H_{11} - H_{22}}{2(\lambda_S - \lambda_L)} = \frac{1}{2} \frac{(m_{\bar{K}^0} - m_{K^0}) - (i/2)(\Gamma_{\bar{K}^0} - \Gamma_{K^0})}{\Delta m + i\Delta\Gamma/2}$$

- $\delta \neq 0$  implies CPT violation
- $\varepsilon \neq 0$  implies T violation
- $\varepsilon \neq 0$  or  $\delta \neq 0$  implies CP violation

$$\Delta m = m_L - m_S, \quad \Delta\Gamma = \Gamma_S - \Gamma_L$$

$$\Delta m = 3.5 \times 10^{-15} \text{ GeV}$$

$$\Delta\Gamma \approx \Gamma_S \approx 2\Delta m = 7 \times 10^{-15} \text{ GeV}$$

(with a phase convention  $\Im\Gamma_{12} = 0$ )

# CPT violation: standard picture

$$|K_{S,L}\rangle \propto \left[ (1 + \varepsilon_{S,L}) |K^0\rangle \pm (1 - \varepsilon_{S,L}) |\bar{K}^0\rangle \right]$$

CP violation:

$$\varepsilon_{S,L} = \varepsilon \pm \delta$$

T violation:

$$\varepsilon = \frac{H_{12} - H_{21}}{2(\lambda_S - \lambda_L)} = \frac{-i\Im M_{12} - \Im\Gamma_{12}/2}{\Delta m + i\Delta\Gamma/2}$$

CPT violation:

$$\delta = \frac{H_{11} - H_{22}}{2(\lambda_S - \lambda_L)} = \frac{1}{2} \frac{(m_{\bar{K}^0} - m_{K^0}) - (i/2)(\Gamma_{\bar{K}^0} - \Gamma_{K^0})}{\Delta m + i\Delta\Gamma/2}$$

huge amplification factor!!

- $\delta \neq 0$  implies CPT violation
- $\varepsilon \neq 0$  implies T violation
- $\varepsilon \neq 0$  or  $\delta \neq 0$  implies CP violation

$$\Delta m = m_L - m_S, \quad \Delta\Gamma = \Gamma_S - \Gamma_L$$

$$\Delta m = 3.5 \times 10^{-15} \text{ GeV}$$

$$(\text{with a phase convention } \Im\Gamma_{12} = 0) \quad \Delta\Gamma \approx \Gamma_S \approx 2\Delta m = 7 \times 10^{-15} \text{ GeV}$$

## neutral kaons vs other oscillating meson systems

	$\langle m \rangle$ (GeV)	$\Delta m$ (GeV)	$\langle \Gamma \rangle$ (GeV)	$\Delta\Gamma/2$ (GeV)
$K^0$	0.5	$3 \times 10^{-15}$	$3 \times 10^{-15}$	$3 \times 10^{-15}$
$D^0$	1.9	$6 \times 10^{-15}$	$2 \times 10^{-12}$	$1 \times 10^{-14}$
$B^0_d$	5.3	$3 \times 10^{-13}$	$4 \times 10^{-13}$	$O(10^{-15})$ (SM prediction)
$B^0_s$	5.4	$1 \times 10^{-11}$	$4 \times 10^{-13}$	$3 \times 10^{-14}$

# “Standard” CPT test

Comparing “survival” probabilities of  $K^0$  and  $\bar{K}^0$  measuring

semileptonic decays vs time:  $\Re\delta = (3.0 \pm 3.3 \pm 0.6) \times 10^{-4}$

CPLEAR

PLB444 (1998) 52

using the unitarity constraint  
(Bell-Steinberger relation)

$$\text{Im } \delta = (-0.7 \pm 1.4) \times 10^{-5}$$

PDG fit (2014)

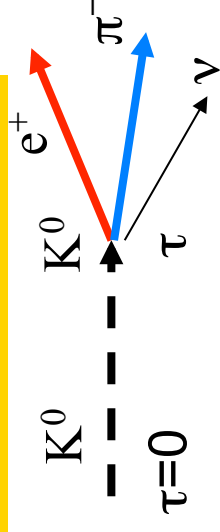
$$\delta = \frac{1}{2} \frac{(m_{\bar{K}^0} - m_{K^0}) - (i/2)(\Gamma_{\bar{K}^0} - \Gamma_{K^0})}{\Delta m + i\Delta\Gamma/2}$$

Combining  $\text{Re}\delta$  and  $\text{Im}\delta$  results

Assuming  $(\Gamma_{\bar{K}^0} - \Gamma_{K^0}) = 0$ , i.e. no CPT viol. in decay:

$$|m_{\bar{K}^0} - m_{K^0}| < 4.0 \times 10^{-19} \text{ GeV}$$

at 95% c.l.

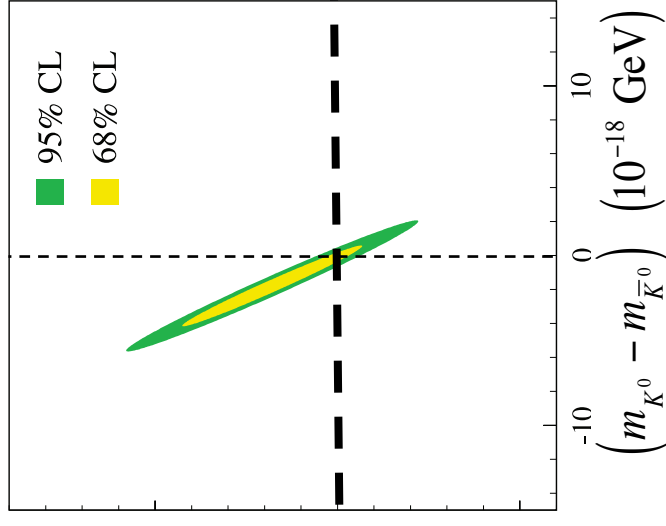


$$2\Re\delta = \Im[\langle K_L | K_S \rangle] = \Im \left[ \frac{\sum_f \langle f | T | K_S \rangle \langle f | T | K_L \rangle^*}{i(\lambda_S - \lambda_L^*)} \right]$$

PDG fit (2014)

$$(\Gamma_{K^0} - \Gamma_{\bar{K}^0})$$

(10<sup>-18</sup> GeV)



---

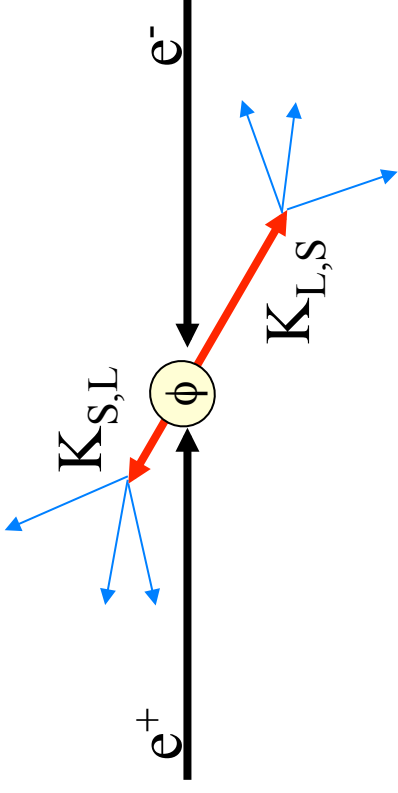
# Entangled neutral kaon pairs



# Neutral kaons at a $\phi$ -factory

Production of the vector meson  $\phi$  in  $e^+e^-$  annihilations:

- $e^+e^- \rightarrow \phi$     $\sigma_\phi \sim 3 \mu\text{b}$
- $W = m_\phi = 1019.4 \text{ MeV}$
- $\text{BR}(\phi \rightarrow K^0\bar{K}^0) \sim 34\%$
- $\sim 10^6$  neutral kaon pairs per  $\text{pb}^{-1}$  produced in an antisymmetric quantum state with  $J^{PC} = 1^{--}$  :



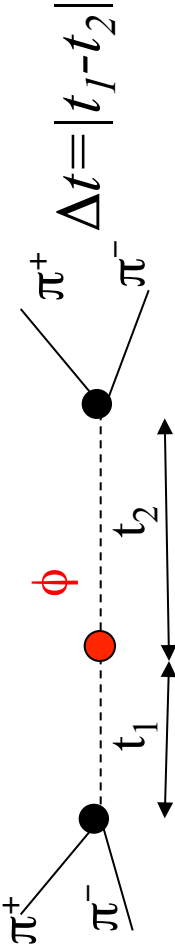
$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[ |K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{N}{\sqrt{2}} \left[ |K_S(\vec{p})\rangle |K_L(-\vec{p})\rangle - |K_L(\vec{p})\rangle |K_S(-\vec{p})\rangle \right]
 \end{aligned}$$

$$N = \sqrt{(1 + |\varepsilon_S|^2)(1 + |\varepsilon_L|^2)} / (1 - \varepsilon_S \varepsilon_L) \cong 1$$

$$\mathbf{p_K = 110 \text{ MeV}/c} \quad \mathbf{\lambda_S = 6 \text{ mm}} \quad \mathbf{\lambda_L = 3.5 \text{ m}}$$

# Neutral kaon interferometry: $\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$

$$|i\rangle = \frac{1}{\sqrt{2}} [ |K^0\rangle |\bar{K}^0\rangle - |\bar{K}^0\rangle |K^0\rangle ]$$



Most precise test of quantum coherence in an entangled system (amplification mechanism due to CP suppression)

$$\zeta_{00} = (1.4 \pm 9.5_{\text{STAT}} \pm 3.8_{\text{SYST}}) \times 10^{-7}$$

$\zeta$  decoherence parameter (QM predicts  $\zeta=0$ )

**KLOE results:** [PLB 642\(2006\) 315](#)

[Found. Phys. 40 \(2010\) 852](#)

Quantum gravity effects might induce:

1) decoherence and CPT violation

(at most  $\gamma = O(m_K^2/M_{\text{Planck}}) \sim 2 \times 10^{-20}$  GeV)

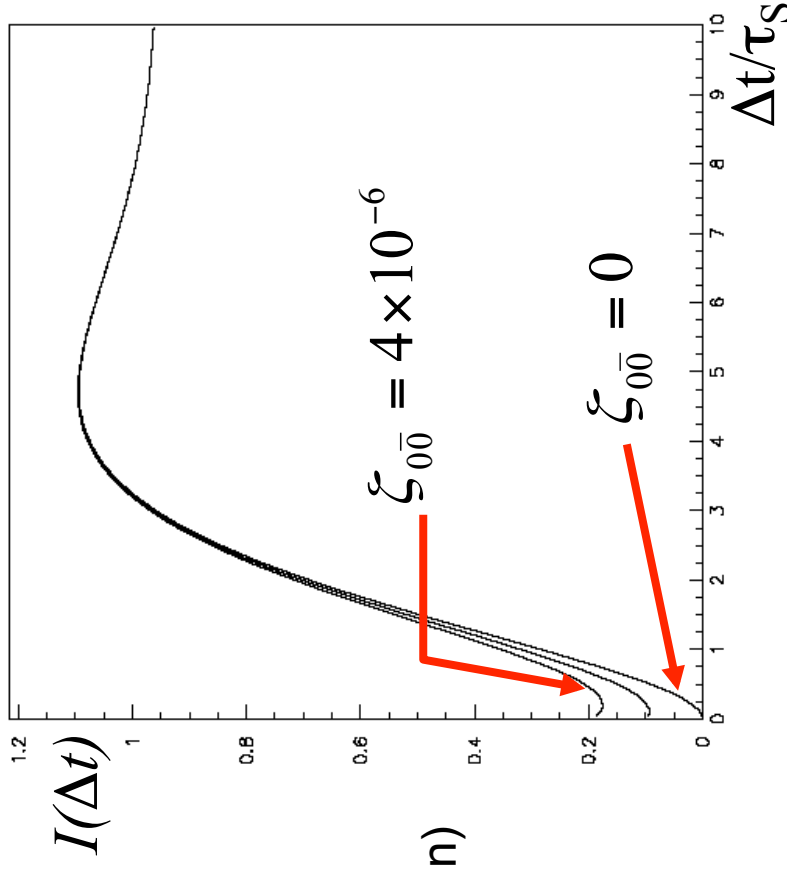
Ellis et al. PRD53 (1996) 3846

2) decoherence and CPT violation induce

modification of the initial correlation of the kaon

pair (at most  $\omega = O(m_K^2/M_{\text{Planck}}/\Delta\Gamma) \sim 1 \times 10^{-3}$ )

Bernabeu, et al. PRL 92 (2004) 131601



$$\gamma = (0.7 \pm 1.2_{\text{STAT}} \pm 0.3_{\text{SYST}}) \times 10^{-21} \text{ GeV}$$

$$|i\rangle \propto (K^0 \bar{K}^0 - K^0 \bar{K}^0) + \omega (K^0 \bar{K}^0 + K^0 \bar{K}^0)$$

$$\Re \omega = (-1.6^{+3.0}_{-2.1 \text{ STAT}} \pm 0.4_{\text{SYST}}) \times 10^{-4}$$

$$\Im \omega = (-1.7^{+3.3}_{-3.0 \text{ STAT}} \pm 1.2_{\text{SYST}}) \times 10^{-4}$$

# Probing CPT in transitions with entangled neutral kaons

---

## Motivations:

- test the CPT symmetry directly in transition processes between kaon states, rather than comparing masses, lifetimes, or other intrinsic properties of particle and anti-particle states.
- CPT violating effects may not appear at first order in diagonal mass terms (survival probabilities) while they can manifest at first order in transitions (non-diagonal terms).
- In standard WWA the test is related to  $\text{Re}\delta$ , a genuine CPT violating effect independent of  $\Delta\Gamma$ , i.e. not requiring the decay as an essential ingredient.
- Clean formulation required. Possible spurious effects induced by CP violation in the decay and/or a violation of the  $\Delta S = \Delta Q$  rule have to be well under control.

**Probing CPT: J. Bernabeu, A.D.D., P. Villanueva, JHEP 10 (2015) 139**  
**Time-reversal violation: J. Bernabeu, A.D.D., P. Villanueva, NPB 868 (2013) 102**

# Definition of states

Let us also consider the states  $|K_+\rangle$ ,  $|K_-\rangle$  defined as follows:  $|K_+\rangle$  is the state filtered by the decay into  $\pi\pi$  ( $\pi^+\pi^+$  or  $\pi^0\pi^0$ ), a pure CP = +1 state; analogously  $|K_-\rangle$  is the state filtered by the decay into  $3\pi^0$ , a pure CP = -1 state. Their orthogonal states correspond to the states which cannot decay into  $\pi\pi$  or  $3\pi^0$ , defined, respectively, as

$$\begin{aligned} |\tilde{K}_-\rangle &\equiv \tilde{N}_- [|K_L\rangle - \eta_{\pi\pi} |K_S\rangle] \\ |\tilde{K}_+\rangle &\equiv \tilde{N}_+ [|K_S\rangle - \eta_{3\pi^0} |K_L\rangle] \end{aligned}$$

$$\begin{aligned} \eta_{\pi\pi} &= \frac{\langle \pi\pi | T | K_L \rangle}{\langle \pi\pi | T | K_S \rangle} \\ \eta_{3\pi^0} &= \frac{\langle 3\pi^0 | T | K_S \rangle}{\langle 3\pi^0 | T | K_L \rangle} \end{aligned}$$

Orthogonal bases:  $\{K_+, \tilde{K}_-\}$   $\{\tilde{K}_+, K_-\}$

Even though the decay products are orthogonal, the filtered  $|K_+\rangle$  and  $|K_-\rangle$  states can still be nonorthogonal.

Condition of orthogonality:

$$\eta_{\pi\pi} + \eta_{3\pi^0}^* = \epsilon_L + \epsilon_S^* \longrightarrow \begin{array}{l} |K_+\rangle \equiv |\tilde{K}_+\rangle \\ |K_-\rangle \equiv |\tilde{K}_-\rangle \end{array}$$

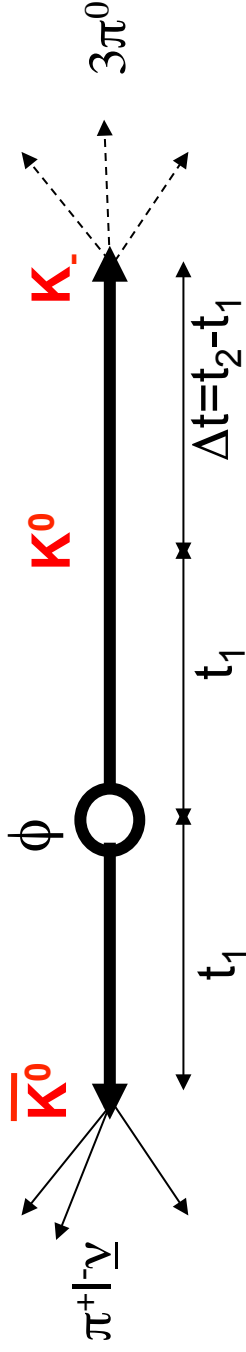
Neglect direct CP violation. Similarly any  $\Delta S = \Delta Q$  rule violation for  $|K^0\rangle$  and  $|\bar{K}^0\rangle$

# Entanglement in neutral kaon pairs

- EPR correlations at a  $\phi$ -factory (or B-factory) can be exploited to study other transitions involving also orthogonal “CP states”  $K_+$  and  $K_-$

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[ |K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[ |K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_- (\vec{p})\rangle |K_+(-\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement -> preparation of state

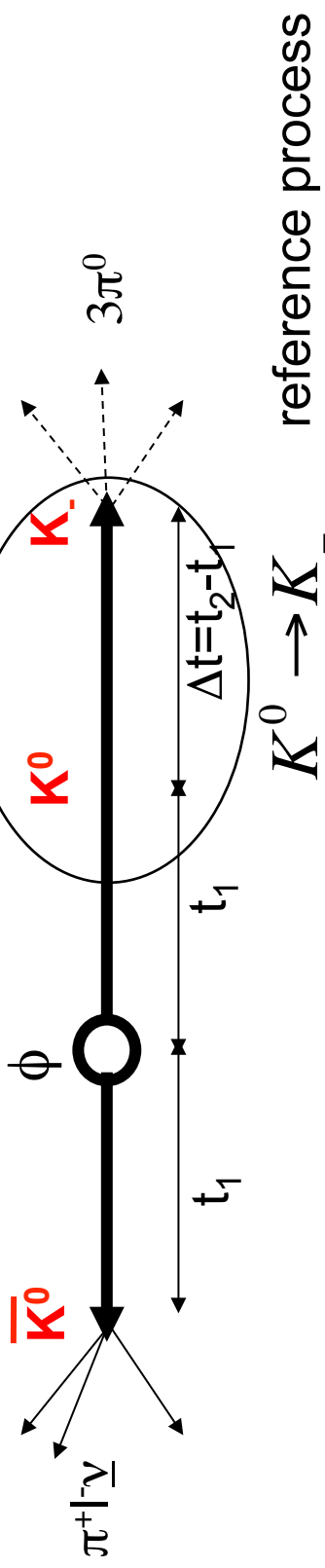


# Entanglement in neutral kaon pairs

- EPR correlations at a  $\phi$ -factory (or B-factory) can be exploited to study other transitions involving also orthogonal “CP states”  $K_+$  and  $K_-$

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} [ |K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle ] \\
 &= \frac{1}{\sqrt{2}} [ |K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(\vec{p})\rangle |K_+(-\vec{p})\rangle ]
 \end{aligned}$$

- decay as filtering measurement
- entanglement  $\rightarrow$  preparation of state

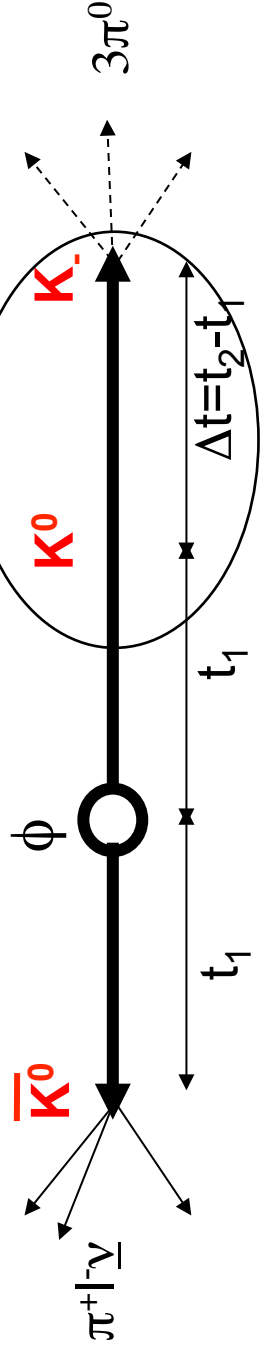


# Entanglement in neutral kaon pairs

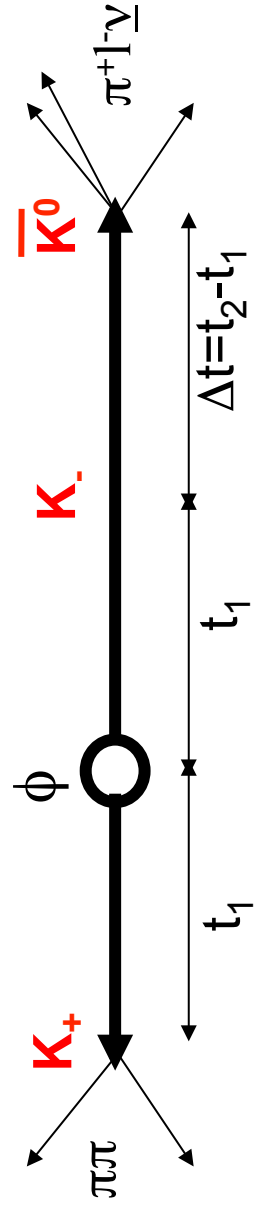
• EPR correlations at a  $\phi$ -factory (or B-factory) can be exploited to study other transitions involving also orthogonal “CP states”  $K_+$  and  $K_-$

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} [ |K^0(\vec{p})\rangle | \bar{K}^0(-\vec{p})\rangle - | \bar{K}^0(\vec{p})\rangle | K^0(-\vec{p})\rangle ] \\
 &= \frac{1}{\sqrt{2}} [ |K_+(\vec{p})\rangle | K_-(-\vec{p})\rangle - |K_-(\vec{p})\rangle | K_+(-\vec{p})\rangle ]
 \end{aligned}$$

- decay as filtering measurement
- entanglement -> preparation of state



$K^0 \rightarrow K_-$  reference process

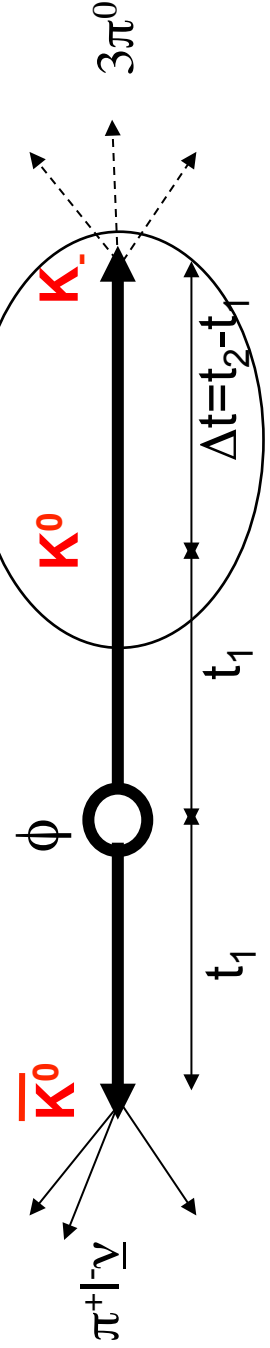


# Entanglement in neutral kaon pairs

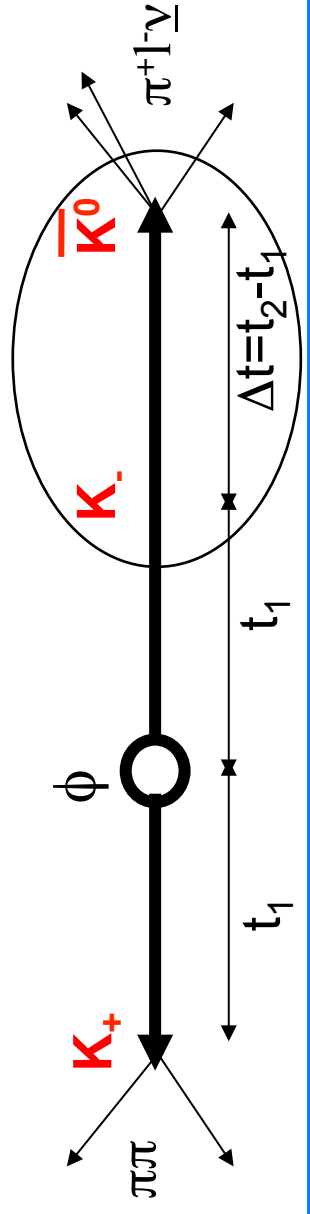
- EPR correlations at a  $\phi$ -factory (or B-factory) can be exploited to study other transitions involving also orthogonal “CP states”  $K_+$  and  $K_-$

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} [ |K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle ] \\
 &= \frac{1}{\sqrt{2}} [ |K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(\vec{p})\rangle |K_+(-\vec{p})\rangle ]
 \end{aligned}$$

- decay as filtering measurement
- entanglement -> preparation of state



$K_- \rightarrow \bar{K}^0$  CPT-conjugated process

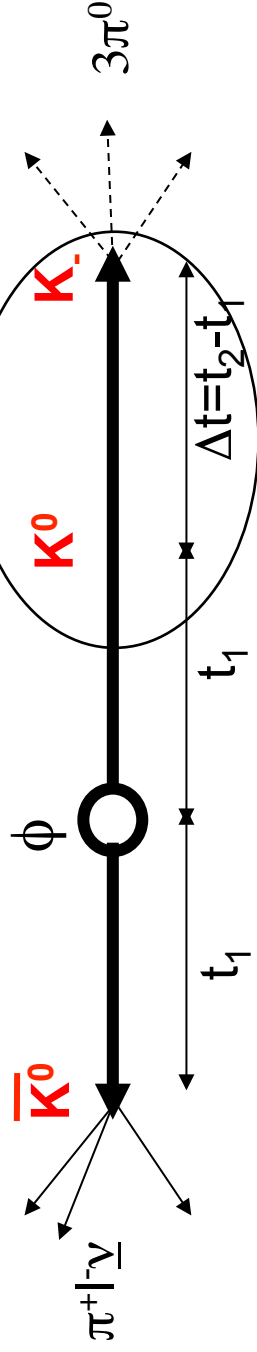




# Entanglement in neutral kaon pairs

- EPR correlations at a  $\phi$ -factory (or B-factory) can be exploited to study other transitions involving also orthogonal “CP states”  $K_+$  and  $K_-$

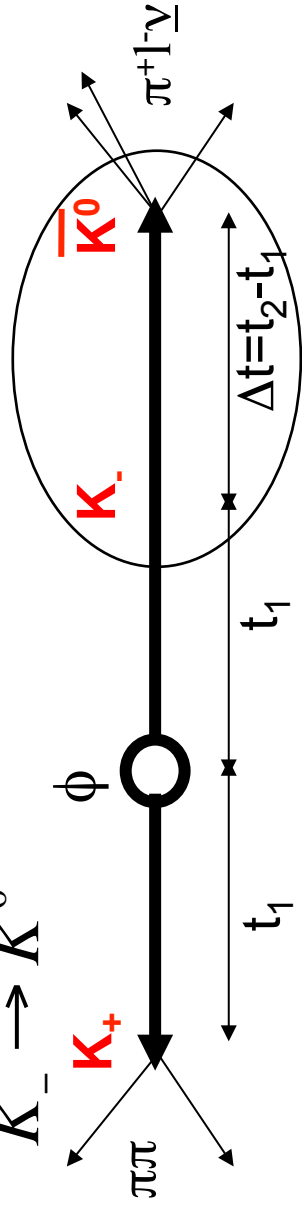
$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} [ |K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle ] \\
 &= \frac{1}{\sqrt{2}} [ |K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(\vec{p})\rangle |K_+(-\vec{p})\rangle ]
 \end{aligned}$$



Note: CP and T conjugated process



CPT-conjugated process



# Direct test of CPT symmetry in neutral kaon transitions

## CPT symmetry test

Reference	<i>CPT</i> -conjugate		
Transition	Decay products	Transition	Decay products
$K^0 \rightarrow K_+$	$(\ell^-, \pi\pi)$	$K_+ \rightarrow \bar{K}^0$	$(3\pi^0, \ell^-)$
$K^0 \rightarrow K_-$	$(\ell^-, 3\pi^0)$	$K_- \rightarrow \bar{K}^0$	$(\pi\pi, \ell^-)$
$\bar{K}^0 \rightarrow K_+$	$(\ell^+, \pi\pi)$	$K_+ \rightarrow K^0$	$(3\pi^0, \ell^+)$
$\bar{K}^0 \rightarrow K_-$	$(\ell^+, 3\pi^0)$	$K_- \rightarrow K^0$	$(\pi\pi, \ell^+)$

One can define the following ratios of probabilities:

$$\begin{aligned} R_{1,CPT}(\Delta t) &= P [K_+(0) \rightarrow \bar{K}^0(\Delta t)] / P [K^0(0) \rightarrow K_+(\Delta t)] \\ R_{2,CPT}(\Delta t) &= P [K^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow \bar{K}^0(\Delta t)] \\ R_{3,CPT}(\Delta t) &= P [K_+(0) \rightarrow K^0(\Delta t)] / P [\bar{K}^0(0) \rightarrow K_+(\Delta t)] \\ R_{4,CPT}(\Delta t) &= P [\bar{K}^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow K^0(\Delta t)] \end{aligned}$$

Any deviation from  $R_{i,CPT}=1$  constitutes a violation of CPT-symmetry

# Direct test of CPT symmetry in neutral kaon transitions

Two observable ratios of double decay intensities

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^-, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^-; \Delta t)}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^+, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^+; \Delta t)}$$

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = R_{2,\text{CPT}}(\Delta t) \times D_{\text{CPT}}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = R_{4,\text{CPT}}(\Delta t) \times D_{\text{CPT}}$$

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = R_{1,\text{CPT}}(|\Delta t|) \times D_{\text{CPT}}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = R_{3,\text{CPT}}(|\Delta t|) \times D_{\text{CPT}}$$

for  $\Delta t > 0$

for  $\Delta t < 0$

$$D_{\text{CPT}} = \frac{\text{BR}(K_L \rightarrow 3\pi^0) \Gamma_L}{\text{BR}(K_S \rightarrow \pi\pi) \Gamma_S}$$

with  $D_{\text{CPT}}$  constant

# Direct test of CPT symmetry in neutral kaon transitions

Explicitly in standard Wigner Weisskopf approach for  $\Delta t > 0$ :

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = \frac{P[K^0(0) \rightarrow K_-(\Delta t)]}{P[K_-(0) \rightarrow \bar{K}^0(\Delta t)]} \times D_{\text{CPT}} \simeq |1 - 2\delta|^2 \left| 1 + 2\delta e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = \frac{P[\bar{K}^0(0) \rightarrow K_-(\Delta t)]}{P[K_-(0) \rightarrow K^0(\Delta t)]} \times D_{\text{CPT}} \simeq |1 + 2\delta|^2 \left| 1 - 2\delta e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}}$$

For comparison the ratio of survival probabilities:

$$\frac{I(\ell^-, \ell^+; \Delta t)}{I(\ell^+, \ell^-; \Delta t)} = \frac{P[K^0(0) \rightarrow K^0(\Delta t)]}{P[\bar{K}^0(0) \rightarrow \bar{K}^0(\Delta t)]} \simeq |1 - 4\delta|^2 \left| 1 + \frac{8\delta}{1 + e^{+i(\lambda_S - \lambda_L)\Delta t}} \right|^2$$

Vanishes for  $\Delta\Gamma \rightarrow 0$

As an illustration of the different sensitivity: it vanishes up to second order in CPTV and decoherence parameters  $\alpha, \beta, \gamma$

# Impact of the approximations

Direct CP (CPT) violation

$$\eta_{\pi\pi} = \epsilon_L + \epsilon'_{\pi\pi}$$

$$\eta_{3\pi^0} = \epsilon_S + \epsilon'_{3\pi^0}$$

Orthogonal bases

$$\{K_+, \tilde{K}_-\}$$

$$\{\tilde{K}_0, K_0\}$$

CPT cons. and CPT viol.  
 $\Delta S = \Delta Q$  violation

$$x_+, x_-$$

Explicitly for  $\Delta t > 0$ :

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = \frac{P[\tilde{K}_0(0) \rightarrow K_-(\Delta t)]}{P[\tilde{K}_-(0) \rightarrow K_0(\Delta t)]} \times D_{\text{CPT}}$$

$$= |1 - 2\delta + 2x_+^* - 2x_-^*|^2 \left| 1 + (2\delta + \epsilon'_{3\pi^0} - \epsilon'_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = \frac{P[\tilde{K}_0(0) \rightarrow K_-(\Delta t)]}{P[\tilde{K}_-(0) \rightarrow K_0(\Delta t)]} \times D_{\text{CPT}}$$

$$= |1 + 2\delta + 2x_+ + 2x_-|^2 \left| 1 - (2\delta + \epsilon'_{3\pi^0} - \epsilon'_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}}$$

## Impact of the approximations

$$\begin{aligned}\frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t)} &\simeq (1 - 8\Re\delta - 8\Re x_-) \left| 1 + 2(\eta_{3\pi^0} - \eta_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \\ &= (1 - 8\Re\delta - 8\Re x_-) \left| 1 + 2(2\delta + \epsilon'_{3\pi^0} - \epsilon'_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2\end{aligned}$$

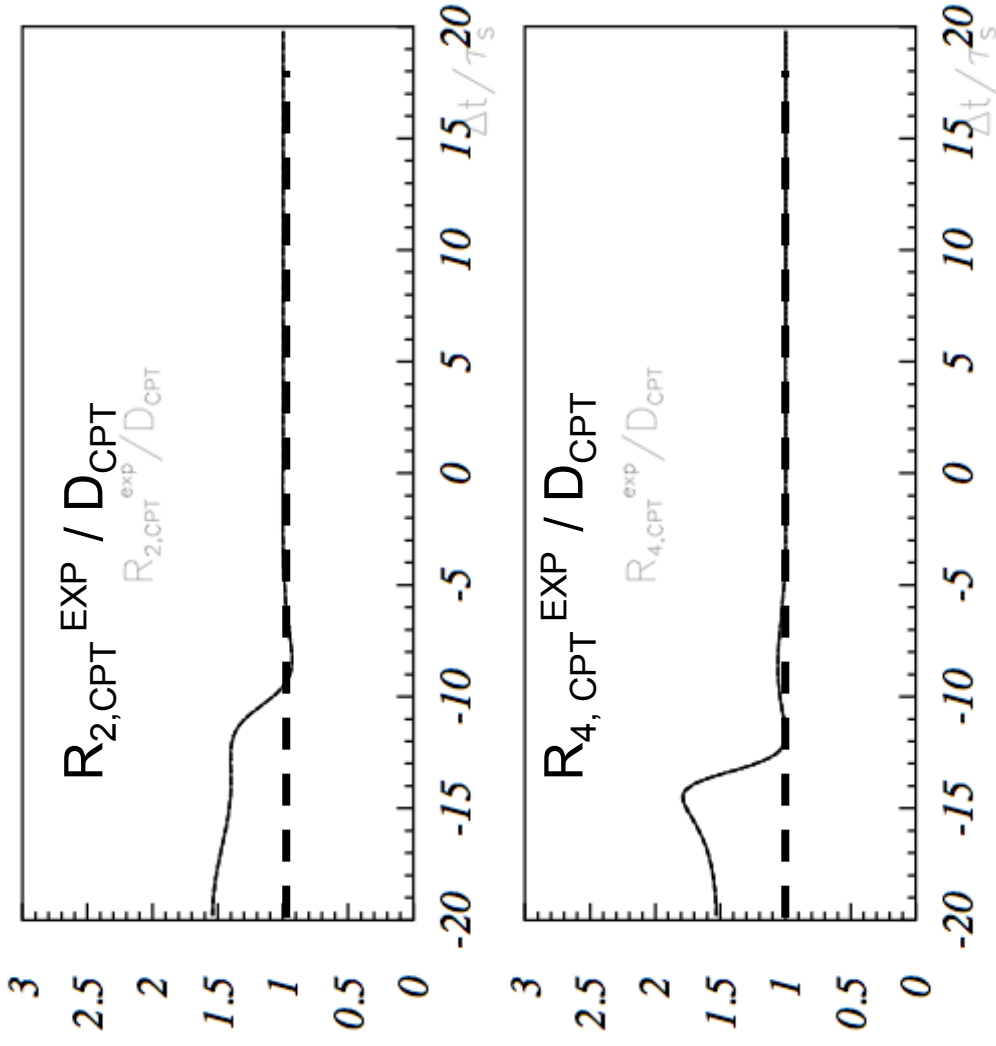
The double ratio constitutes one of the most robust observables for the proposed CPT test. In the limit  $\Delta t \gg \tau_S$  it exhibits a pure and genuine CPT violating effect, even without assuming negligible contaminations from direct CP violation and/or  $\Delta S = \Delta Q$  rule violation.

$$\frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)} = 1 - 8\Re\delta - 8\Re x_-$$

# Direct test of CPT symmetry with neutral kaons

for visualization purposes, plots with

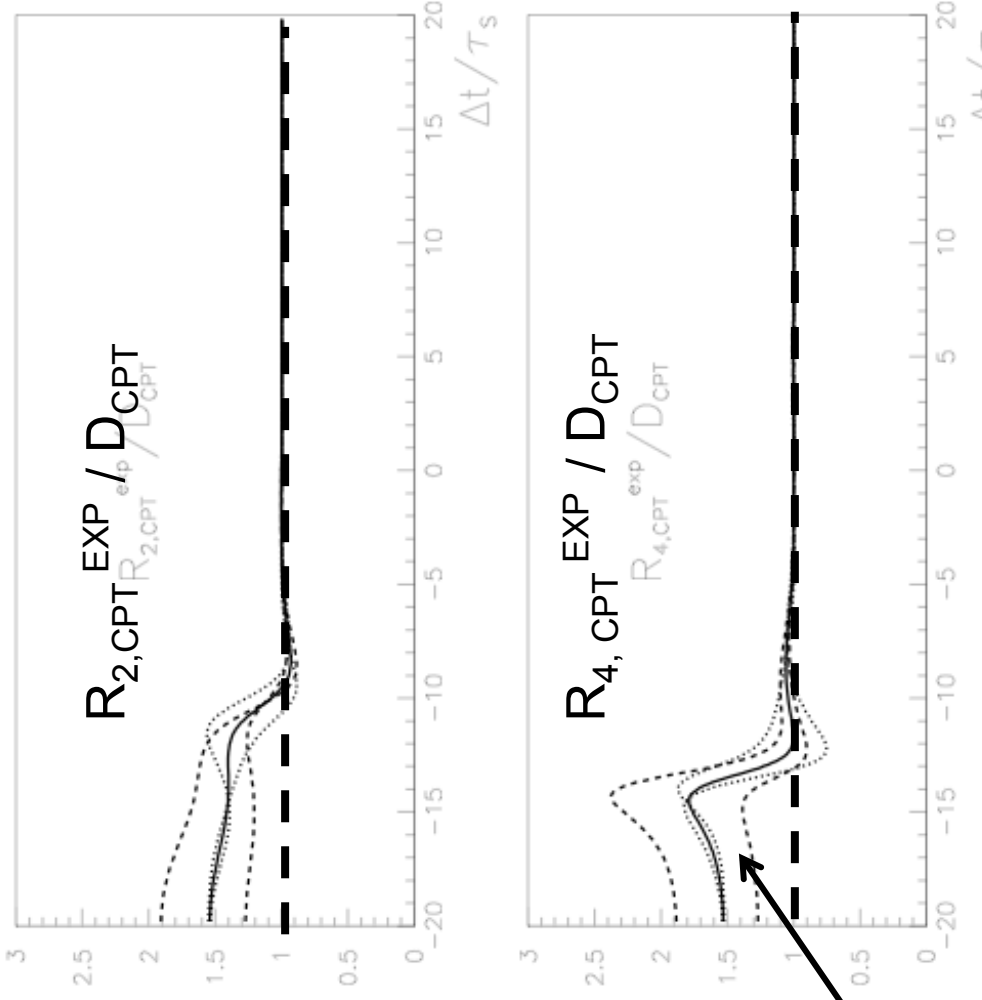
$$\text{Re}(\delta)=3.3 \cdot 10^{-4} \quad \text{Im}(\delta)=1.6 \cdot 10^{-5}$$



# Direct test of CPT symmetry with neutral kaons

for visualization purposes, plots with

$$\text{Re}(\delta)=3.3 \cdot 10^{-4} \quad \text{Im}(\delta)=1.6 \cdot 10^{-5}$$



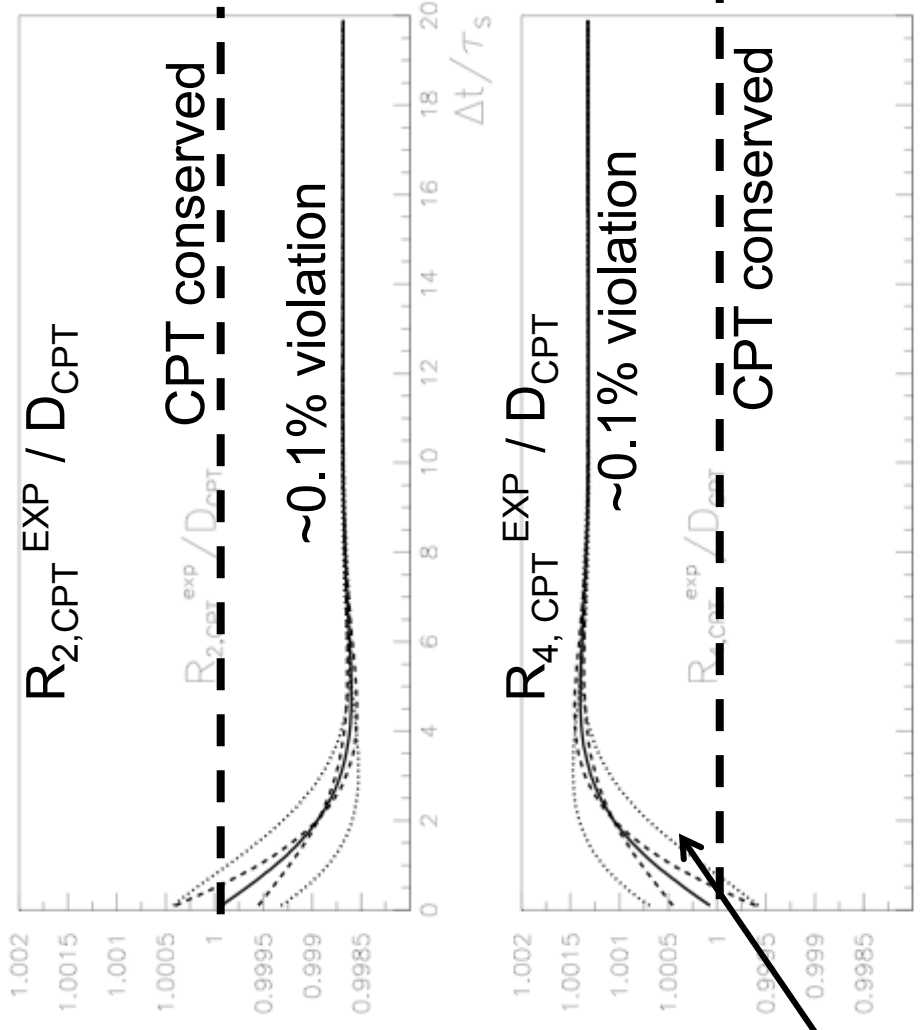
Modifications due to direct CP violation effects (unrealistically amplified  $\sim \times 100$ )



# Direct test of CPT symmetry with neutral kaons

for visualization purposes, plots with

$$\text{Re}(\delta)=3.3 \cdot 10^{-4} \quad \text{Im}(\delta)=1.6 \cdot 10^{-5}$$

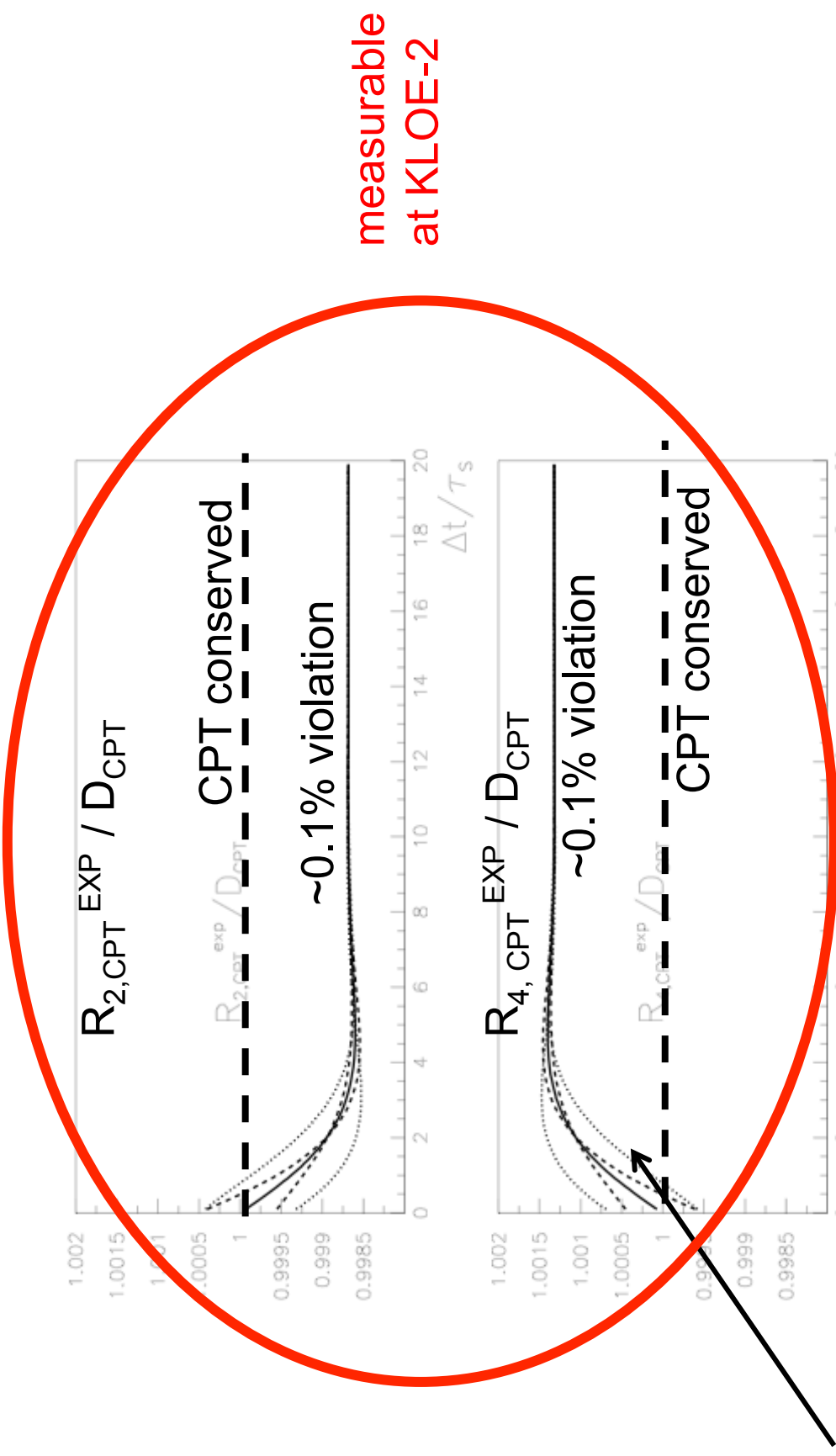


Modifications due to direct CP violation effects (unrealistically amplified  $\sim \times 10$ )

# Direct test of CPT symmetry with neutral kaons

for visualization purposes, plots with

$$\text{Re}(\delta)=3.3 \cdot 10^{-4} \quad \text{Im}(\delta)=1.6 \cdot 10^{-5}$$



measurable  
at KLOE-2

Modifications due to direct CP violation effects (unrealistically amplified  $\sim \times 10$ )

# Conclusions

- It is possible to directly test the CPT symmetry in transition processes.
- The proposed CPT test for neutral kaons is model independent and fully robust. (It can then be translated in terms of  $\delta$ ,  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\Delta a_\mu$  etc..).
- In standard WWA the test is related to  $\text{Re}\delta$ , a genuine CPT violating effect independent of  $\Delta\Gamma$  and not requiring the decay as an essential ingredient.
- Possible spurious effects induced by CP violation in the decay and/or a violation of the  $\Delta S = \Delta Q$  rule have been shown to be well under control.
- Perfect anti-correlation of the initial state is assumed (to be evaluated the impact of possible loss of coherence).
- There exists a connection with charge semileptonic asymmetries of  $K_S$  and  $K_L$

$$\frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)} = \frac{1 + A_L}{1 - A_L} \times \frac{1 - A_S}{1 + A_S} \simeq 1 + 2(A_L - A_S)$$

- KLOE data analysis ongoing (see Daria Kaminska's talk)
- KLOE-2 can reach a statistical sensitivity of  $O(10^{-3})$  on this new observable.
- In B meson system similar test (see Miyashita's talk); less precise, more assumptions needed

---

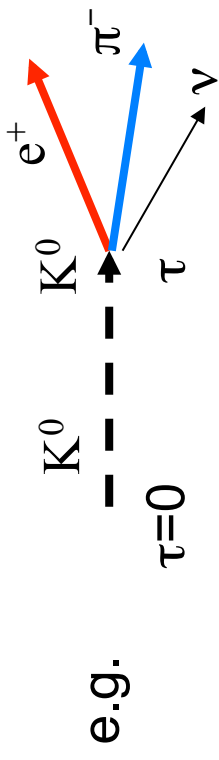
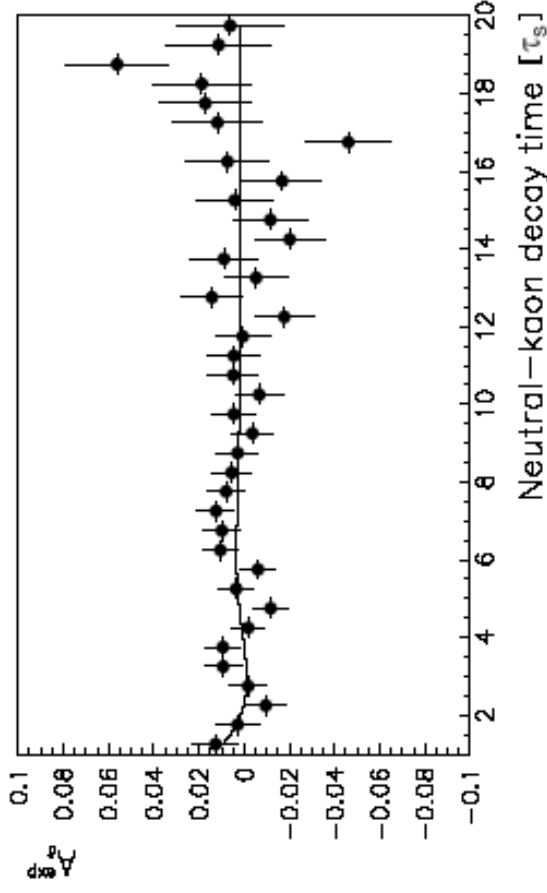
---

## Back up slides

# CPT test at CPLEAR

Test of **CPT** in the time evolution of neutral kaons using the semileptonic asymmetry

Comparing “survival”  $K^0 \rightarrow K^0$  probabilities:  $\bar{K}^0 \rightarrow \bar{K}^0$



$$A_\delta(\tau) = \frac{\bar{R}_+(\tau) - \alpha R_-(\tau)}{\bar{R}_+(\tau) + \alpha R_-(\tau)} + \frac{\bar{R}_-(\tau) - \alpha R_+(\tau)}{\bar{R}_-(\tau) + \alpha R_+(\tau)}$$

$$R_{+(-)}(\tau) = R \left( K^0_{t=0} \rightarrow (e^{+(-)} \pi^{-(+)} \nu)_{t=\tau} \right)$$

$$\bar{R}_{- (+)}(\tau) = R \left( \bar{K}^0_{t=0} \rightarrow (e^{- (+)} \pi^{+(-)} \nu)_{t=\tau} \right)$$

$$\alpha = 1 + 4\Re \varepsilon_L$$

$$\Re \delta = (0.30 \pm 0.33 \pm 0.06) \times 10^{-3}$$

$$A_\delta(\tau \gg \tau_S) = 8\Re \delta$$

CPLEAR PLB444 (1998) 52

# The Bell-Steinberger relationship



J. Bell

(1965)



J. Steinberger

Unitarity constraint:

$$|K\rangle = a_S |K_S\rangle + a_L |K_L\rangle$$

$$\left( -\frac{d}{dt} \|K(t)\|^2 \right)_{t=0} = \sum_f |a_S \langle f|T|K_S\rangle + a_L \langle f|T|K_L\rangle|^2$$

Sum over all possible decay products  
(sum over few decay products for kaons;  
many for B and D mesons => not easy to evaluate)

yields two trivial relations:

$$\Gamma_{S,L} = \sum_f |\langle f|T|K_{S,L}\rangle|^2$$

and a not trivial one, i.e. the B-S relationship:

All observables quantities

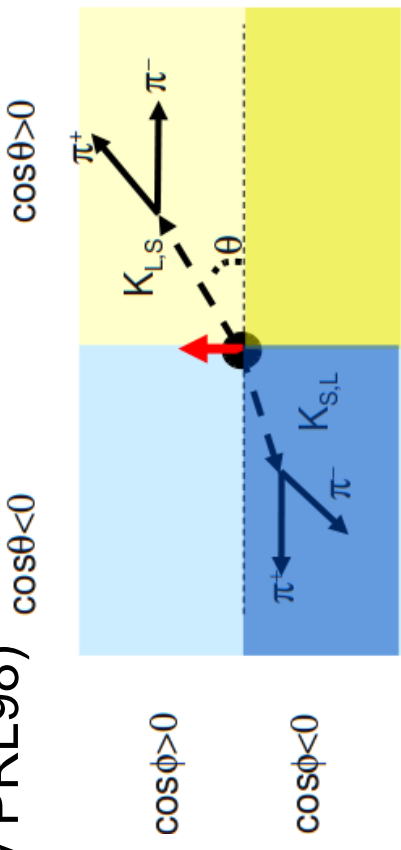
$$\langle K_L | K_S \rangle = 2(\Re \varepsilon + i\Im \delta) = \frac{\sum_f \langle f|T|K_S\rangle \langle f|T|K_L\rangle^*}{i(\lambda_S - \lambda_L^*)}$$

# Search for CPT and Lorentz invariance violation at KLOE

$$\delta = i \sin \phi_{SW} e^{i\phi_{SW}} \gamma_K \left( \Delta a_0 - \vec{\beta}_K \cdot \Delta \vec{a} \right) / \Delta m$$

Standard Model extension (SME)  
(Kostelecky PRL98)

Data divided in  
4 sidereal time bins x 2 angular bins  
Simultaneous fit of the  $\Delta t$  distributions  
to extract  $\Delta a_{\nu}$  parameters



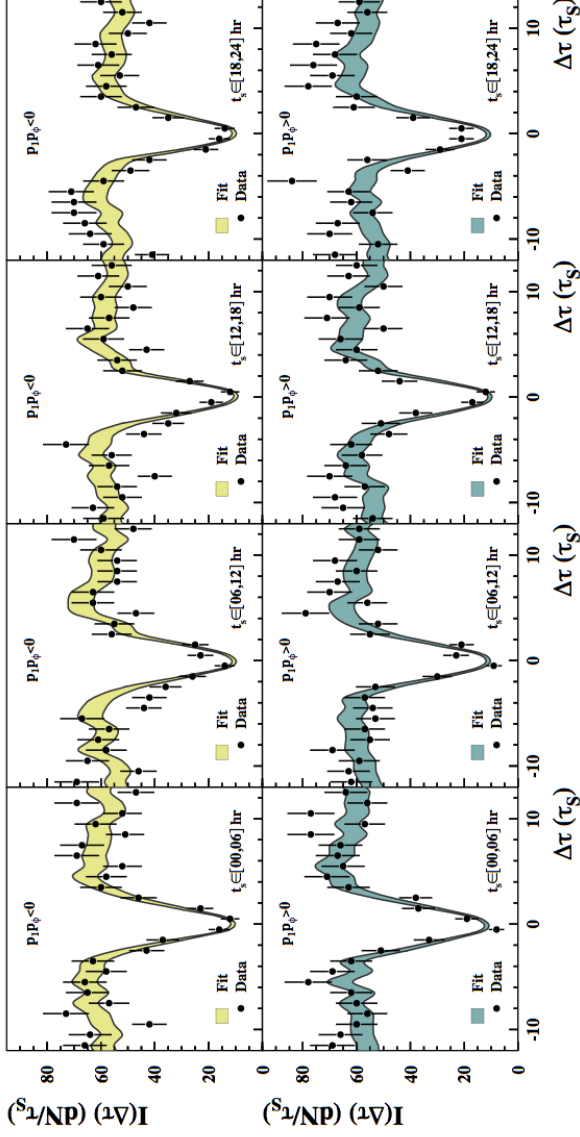
KLOE PLB 730 (2014) 89 with  $L=1.7 \text{ fb}^{-1}$

$$\Delta a_0 = \left( -6.0 \pm 7.7_{\text{STAT}} \pm 3.1_{\text{SYST}} \right) \times 10^{-18} \text{ GeV}$$

$$\Delta a_X = \left( 0.9 \pm 1.5_{\text{STAT}} \pm 0.6_{\text{SYST}} \right) \times 10^{-18} \text{ GeV}$$

$$\Delta a_Y = \left( -2.0 \pm 1.5_{\text{STAT}} \pm 0.5_{\text{SYST}} \right) \times 10^{-18} \text{ GeV}$$

$$\Delta a_Z = \left( -3.1 \pm 1.7_{\text{STAT}} \pm 0.6_{\text{SYST}} \right) \times 10^{-18} \text{ GeV}$$



presently the most precise measurements in the quark sector of the SME

B meson system:  $\Delta a_{X,Y}^B \sim O(10^{-15} \text{ GeV})$   $\Delta a^B_{\text{perp}} \sim O(10^{-13} \text{ GeV})$  [LHCb PRL116 (2016) 241601]

D meson system:  $\Delta a_{X,Y}^D$ ,  $(\Delta a^D_0 - 0.6 \Delta a^D_Z) \sim O(10^{-13} \text{ GeV})$  [Focus PLB 556 (2003) 7]