BROOKH/VEN NATIONAL LABORATORY

SEARCH FOR THE RARE DECAYS $K+\rightarrow \mu+ X IN E949 EXPERIMENT$

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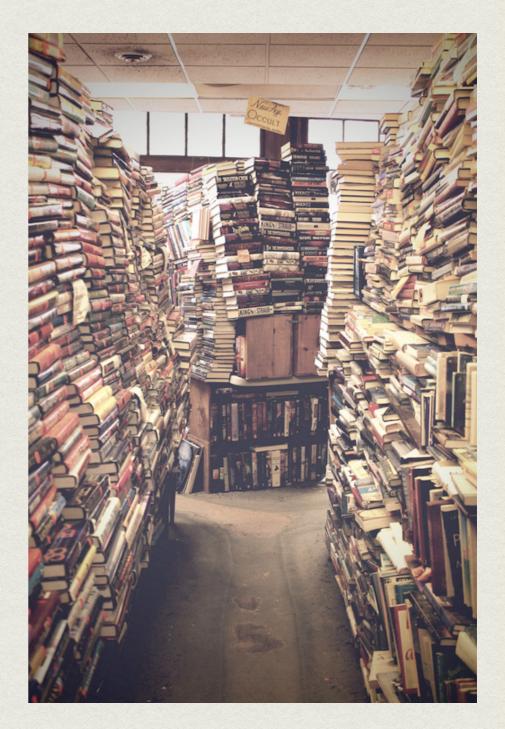
KAON2016 14–17 September, University of Birmingham, UK

OUTLINE

- Search for heavy neutrinos: $K^+ \rightarrow \mu^+ \nu_H$, Phys. Rev. D 91, 052001 (2015)
- Search for rare decay K⁺→µ⁺vvv,
 Phys. Rev. D 94, 032012 (2016)

INTRODUCTION

- 1. Motivation
- 2. Experiment BNL E949
 - The E949 detector
 - Data samples

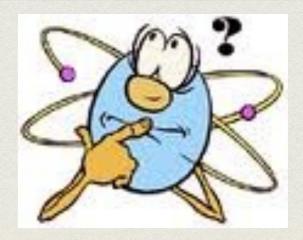


BEYOND THE STANDARD MODEL



 ν MSM: SM + 3 right-handed neutrinos m₁ ~ 10 keV m_{2,3}~ 100 MeV - 100 GeV There is new physics beyond the Standard Model, but we don't know exactly what is it

T. Asaka and M. Shaposhnikov Phys. Lett. B620, 17 (2005).



HOW TO FIND HEAVY NEUTRINOS?

• <u>Meson decays</u>

Search for extra peaks in lepton distributions (momentum, energy, missing mass, ...)

 $\Gamma(M^+ \to l^+ \nu_H) = \rho \times \Gamma(M^+ \to l^+ \nu_l) \times |U_{lH}|^2$ R.E. Shrock, Phys. Rev. D24, 1232 (1981)

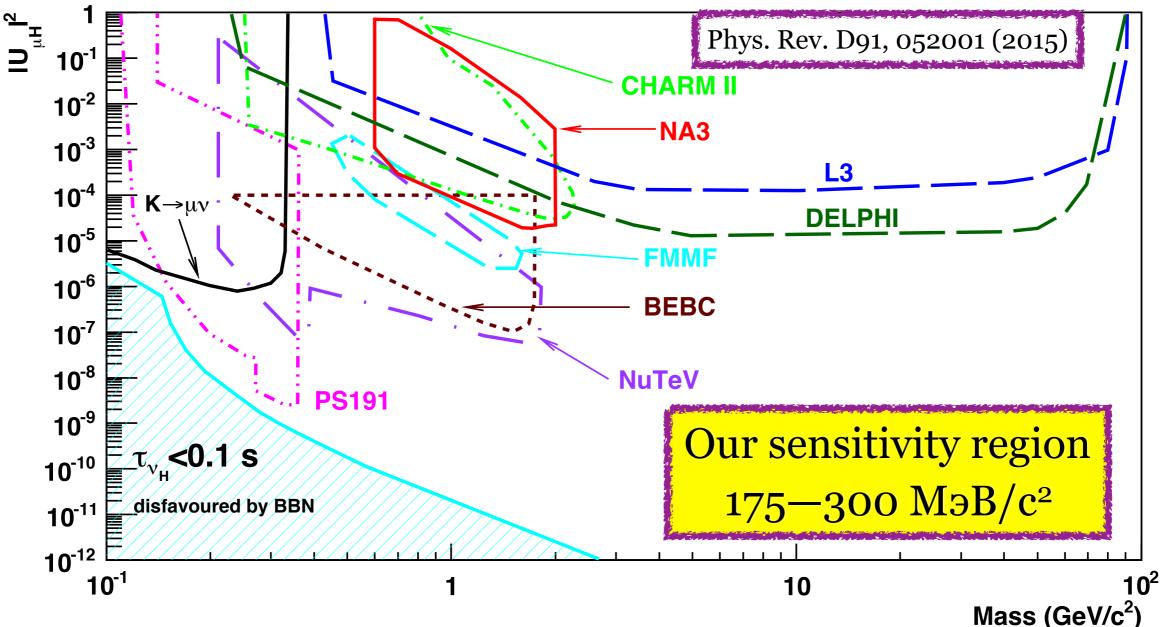
Heavy neutrino decays

"Nothing" > leptons and hadrons $\nu_H \to e^+ e^- \nu_{\alpha}, \nu_H \to \mu^{\pm} e^{\mp} \nu_{\alpha}, \nu_H \to \mu^+ \mu^- \nu_{\alpha},$ $\nu_H \to \pi^0 \nu, \pi e, \pi \mu, K e, K \mu, ...$

CURRENT LIMITS(E949 RESULT IS OMITTED)

limits Belle, BaBaR, LHCb for heavy neutrino mass up to 5 GeV/c^2 are not shown here (Nuclear Physics B Proceedings Supplement oo (2014) 1–4, see also arXiv:1502.00477),

limits from CMS for mass $50-500 \text{ GeV/c}^2$ are also not shown (JHEP 04 (2016) 169)



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EXPERIMENT BNL-E949



Search for ultra-rare decay $K^+ \to \pi^+ \nu \bar{\nu}$ Stopped kaons were used $BR_{SM}(K^+ \to \pi^+ \nu \bar{\nu}) =$ $(9.11 \pm 0.72) \times 10^{-11}$

Buras, A.J., Buttazzo, D. & Knegjens, R., JHEP 1511 (2015) 166

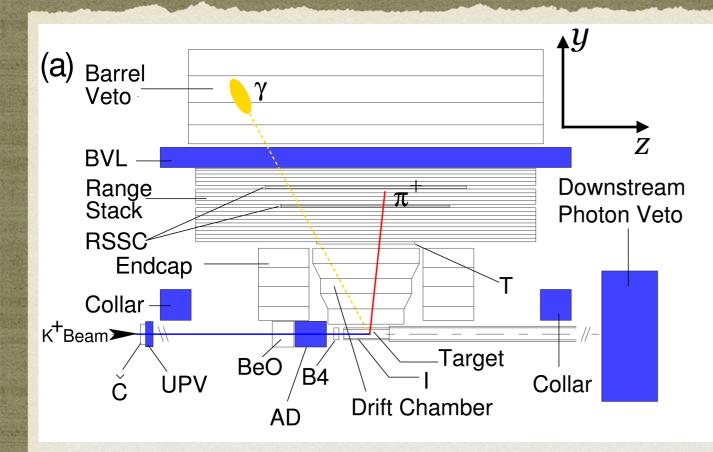
E949+E787

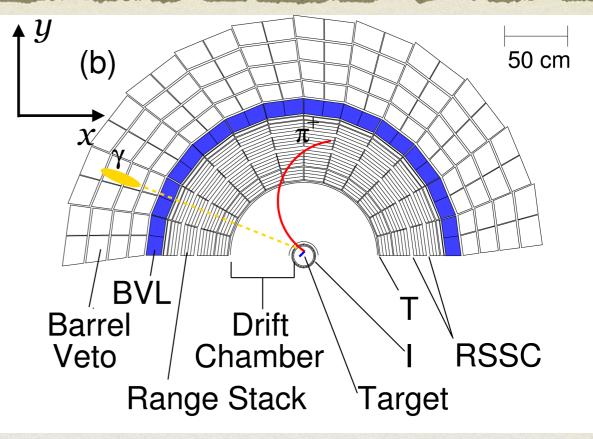
4+3(E787)=7 events

 $BR(K^+ \to \pi^+ \nu \bar{\nu}) =$ $(1.73^{+1.15}_{-1.05}) \times 10^{-10}$

Phys. Rev. D79, 092004 (2009)

THE E949 DETECTOR





- Incoming 710 MeV/c kaons were identified by Čerenkov counter and slowed down by degraders.
- Kaons came to rest in the centre of the target, which was made of 413 5 mm square scintillating fibres.
- The whole spectrometer was in a 1 Tl magnetic field. Daughter particle momentum was measured in the Drift Chamber, energy and range in the Range Stack.
- Hermetic system of photon veto detectors.

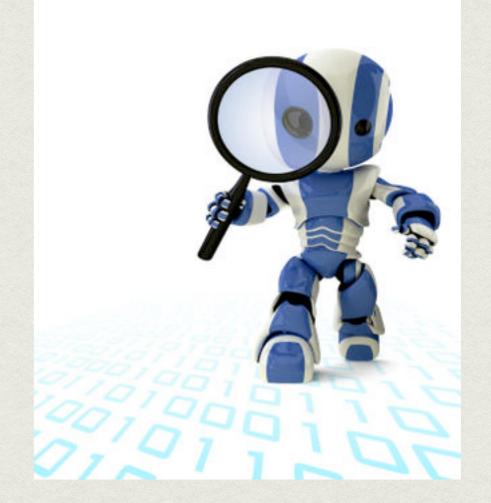
DATA SAMPLES

- The main trigger. Huge statistics: 1.70×10¹² stopped kaons
 - → ✓ Pros signal decay $K^+ \rightarrow \mu^+ \nu_H$ is similar to $K^+ \rightarrow \pi^+ \nu \nu$: one single charge track and no other detector activity
 - □ Cons it was designed to select pions from the $K^+ \rightarrow \pi^+ \nu \nu$, rejected muons
 - Monitor triggers
 - ✓ Pros simple selection, no muons rejection
 - \Box Cons much less stopped kaons (~10⁷)

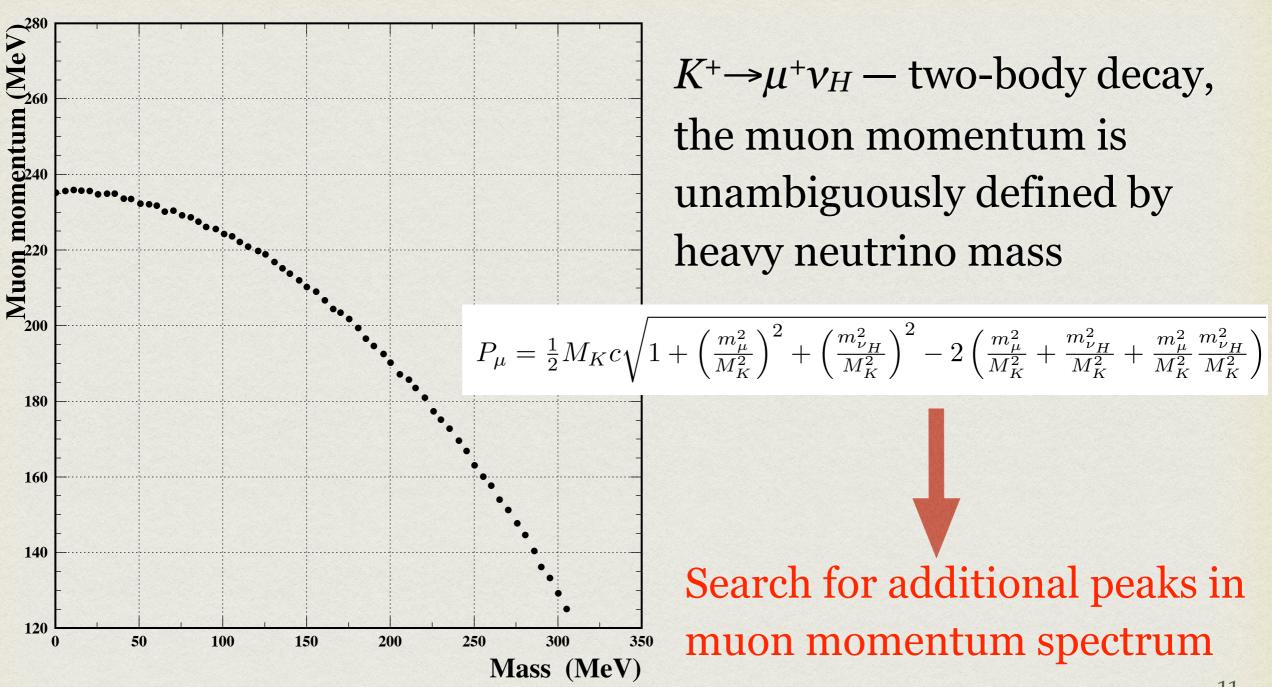
Used for data quality assessment, calibrations of the detector subsystems and acceptance measurement for the $K^+ \rightarrow \mu^+ v_H$ decay.

DATA ANALYSIS

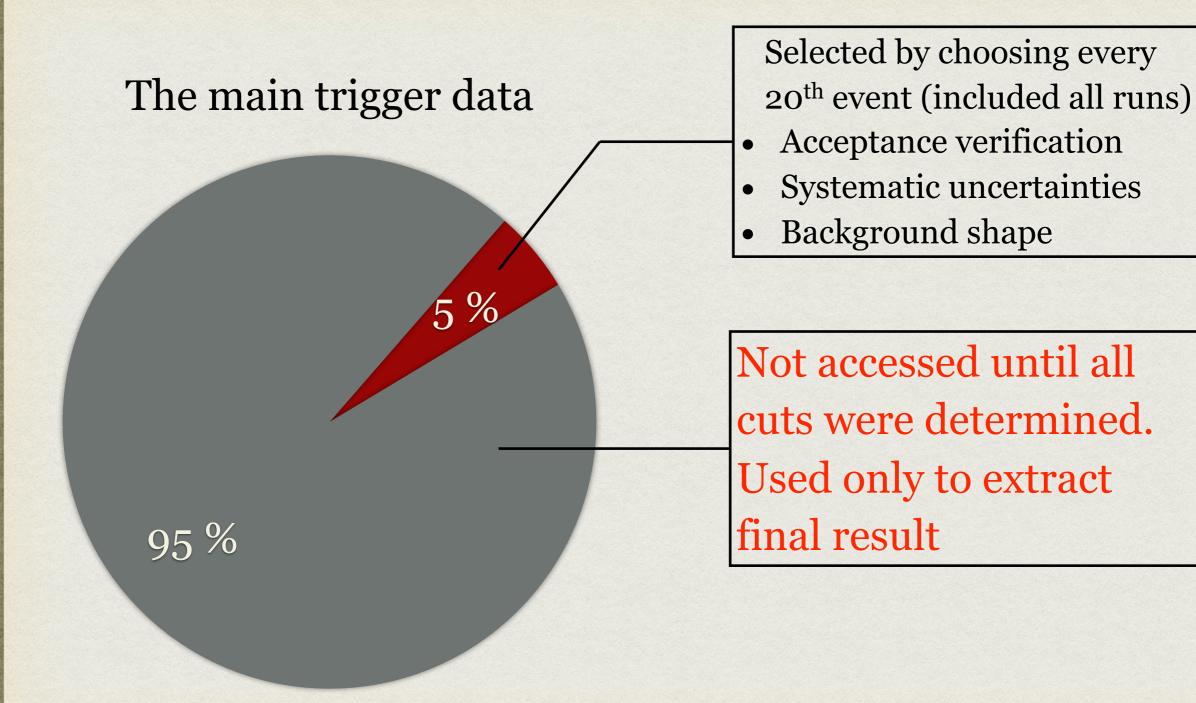
- 1. Single muon selection criteria
- 2. Acceptance measurement
- 3. Study of the 1/20 data
 - Acceptance verification
 - Systematic uncertainties
 - Background study
- 4. Momentum resolution for the different muon momentum
- 5. Peak search method



THE STRATEGY



THE STRATEGY



SELECTION CRITERIA

The $K^+ \rightarrow \mu^+ v_H$ signature — single muon track

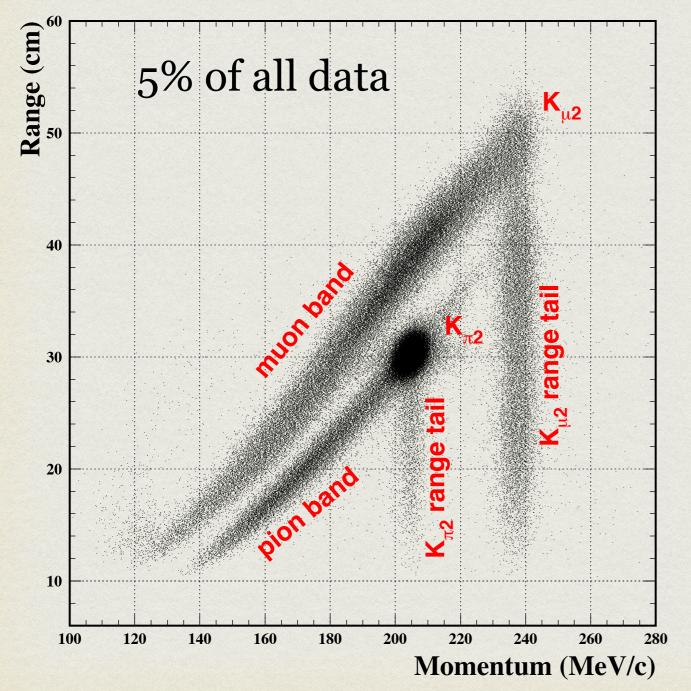
Trigger

- Kaon stopped in target
- Charged track in the fiducial volume(1)
- Delay coincidence (rejected kaon decays-in-flight)
- Rejected long track
- Photon veto
- Pion identification in the stopping counter: requested decay chain $\pi \rightarrow \mu$

Offline

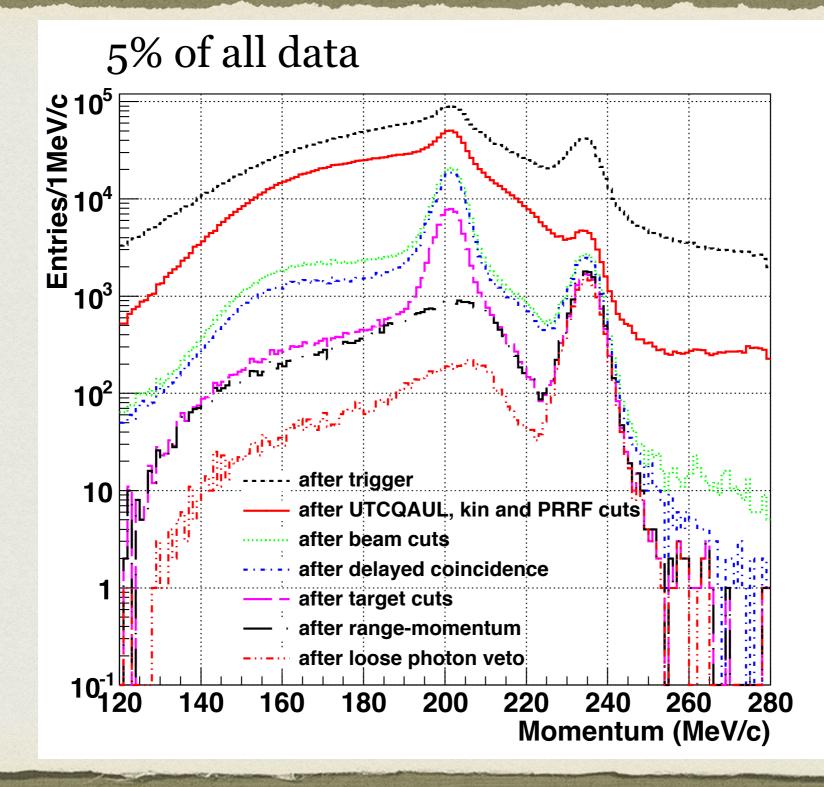
- Track was reconstructed in the drift chamber
- Charged track in the fiducial volume(2)
- Rejected kaon decays-in-flight
- Rejected extra beam particles
- Good track in the target and Range Stack (numerous requirements)
- Muon identification
- Photon veto (loose or tight)

THE MAIN TRIGGER DATA



- Selected by choosing every 20th event (included all runs)
- Muon band: muons from $K_{\mu\nu\gamma}$ and $K_{\mu3}$ decays
- Pion band: pions from *K*_{π2γ} decay and from *K*_{π2} decay in which the pions scattered in the target or Range Stack and beam pions that scattered in the target
- Both *K_{µ2}* and *K_{π2}* range tails elastic or inelastic scattering in the Range Stack

SELECTION CRITERIA

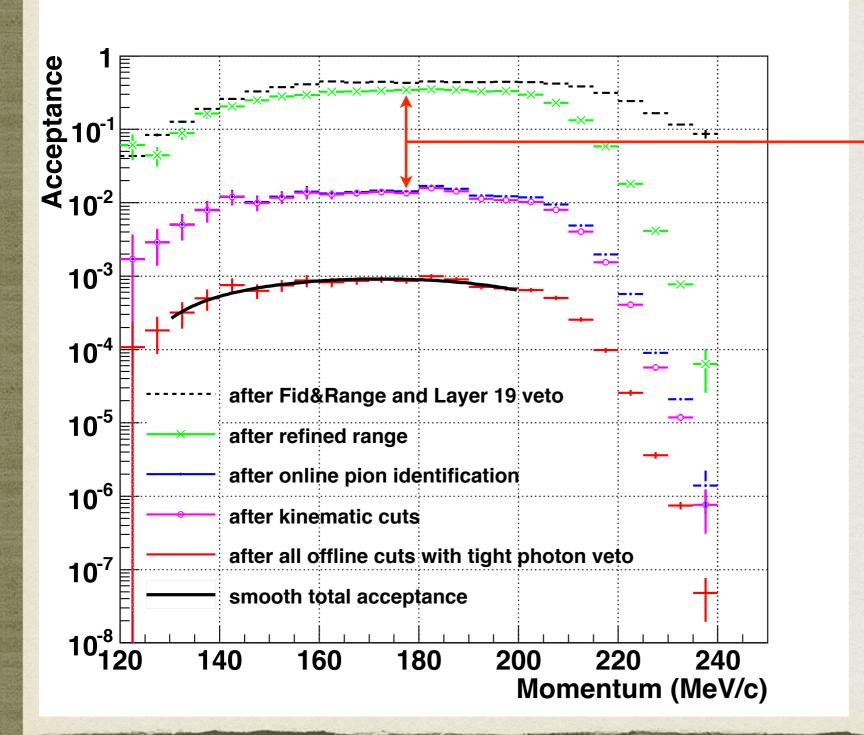


ACCEPTANCE MEASUREMENT

Total: 73 selection criteria

- Muons from $K^+ \rightarrow \mu^+ \nu_{\mu}$ and $K^+ \rightarrow \mu^+ \nu_{\mu} \gamma$ decays. Used monitor triggers (63 selection criteria)
- Monte-Carlo simulation of the $K^+ \rightarrow \mu^+ v_H$ decay (fiducial volume) (6 selection criteria)
- Used values from the main E949 analysis $(K^+ \rightarrow \pi^+ vv^- \text{decay})$ (4 selection criteria)

ACCEPTANCE MEASUREMENT



The largest acceptance lost
was due to pion
identification online (~20).
It was implemented in the
trigger, cannot be turned
off

Single events sensitivity for the $K^+ \rightarrow \mu^+ \nu_H$ decay:

$$S.E.S = \frac{1}{KBlive \times Acc} =$$

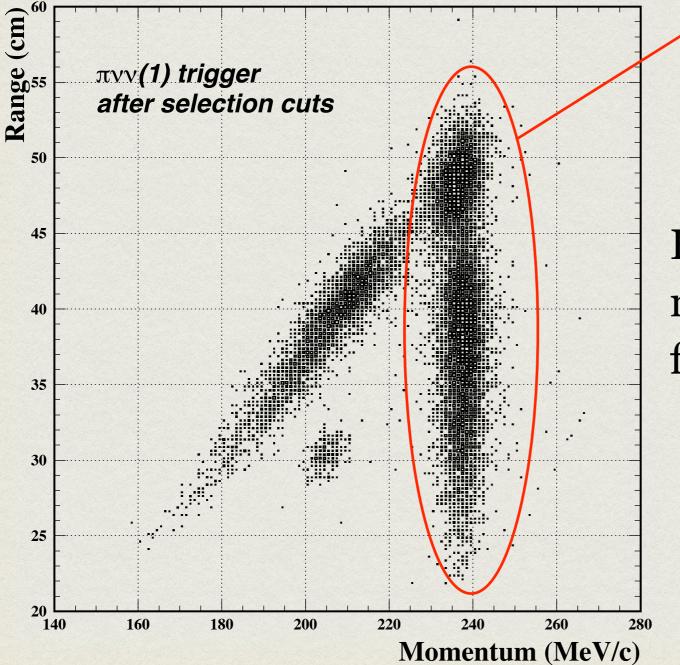
 7.35×10^{-10}

ACCEPTANCE VERIFICATION

* The K⁺ \rightarrow µ⁺v_µ branching ratio measurement, 5% of all data was used

* The K⁺ \rightarrow µ⁺v_µγ branching ratio measurement (140<p_µ<200 MeV/c), 5% of all data was used

$K^+ \rightarrow \mu^+ \nu_{\mu} DECAY$



Events for the K_{μ2}
 branching ratio
 measurement

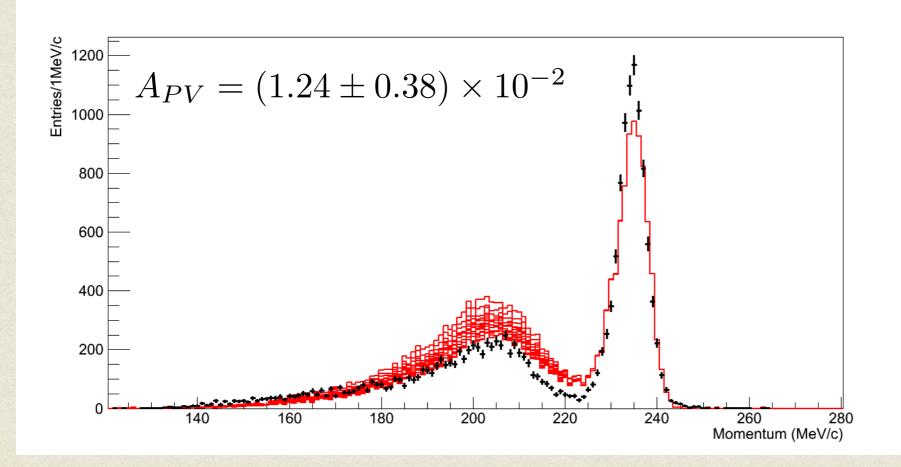
For the given muon momentum some acceptance factors should be corrected

 $BR(K_{\mu 2}) = 0.54 \pm 0.15$ $BR^{PDG}(K_{\mu 2}) = 0.6355 \pm 0.0011$

$K^+ \rightarrow \mu^+ \nu_{\mu} \gamma DECAY$

Branching ratio measurement in the muon momentum region 140< p_{μ} <200 MeV/c. We measured all acceptance factors, except the photon veto (one photon in the final state)

Used MC simulation and single photon inefficiency table.



Varied photon veto detectors thresholds from zero to the nominal values. The photon veto acceptance (A_{PV}) — the mean acceptance of all red histograms.

$K^+ \rightarrow \mu^+ \nu_{\mu} \gamma DECAY$

The branching ratio for the muon momentum $140 < p_{\mu} < 200 \text{ MeV/c}$ was measured to be:

$$BR(K_{\mu\nu\gamma}) = (1.3 \pm 0.4) \times 10^{-3}$$

The measured uncertainty is the systematic uncertainty for the acceptance of the $K^+ \rightarrow \mu^+ \nu_H$ decay.

 $BR^{PDG}(K_{\mu\nu\gamma}) = (6.2 \pm 0.8) \times 10^{-3}, p_{\mu} < 231.5 \text{ MeV/c}$ Used MC simulation to determine branching ratio in the selected muon momentum region:

$$BR^{PDG}(K_{\mu\nu\gamma}) = (1.4 \pm 0.2) \times 10^{-3}$$

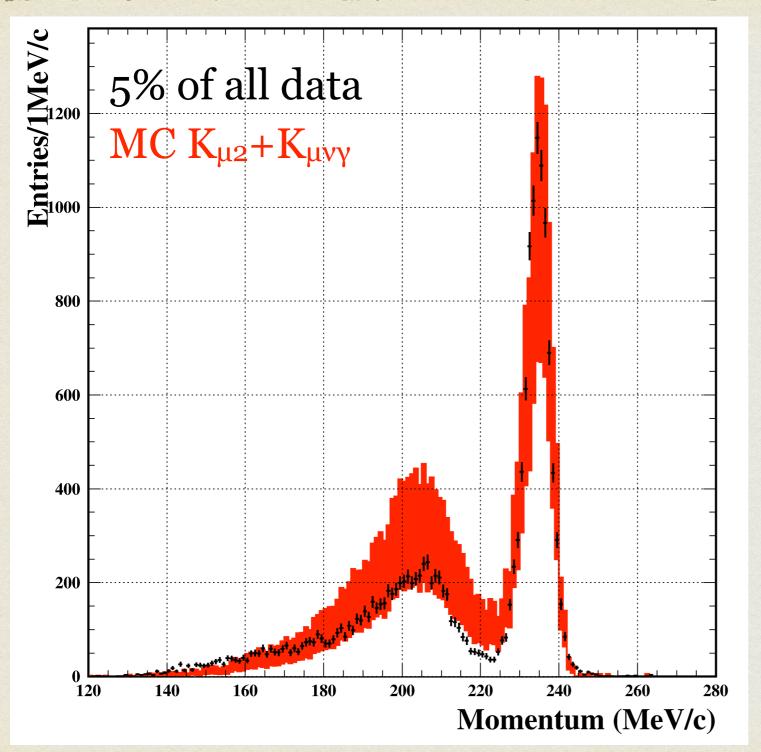
BACKGROUND PROCESSES

Process	Trigger+cuts rej	BR	Total rejection
$K_{\mu\nu\gamma}$	$\sim 10^4$	6.2×10^{-3}	$\sim 10^7$
$K_{\mu3}$	$\sim 10^7$	3.35×10^{-2}	$\sim 10^9$
Only $\pi \nu \nu (1+2)$ trigger			
$K_{\pi 2\gamma}$	$\sim 5 \times 10^4$	2.75×10^{-4}	$\sim 2 \times 10^9$

Used MC simulation . The $K_{\pi_{2\gamma}}$ decay was ignored due to three photons in the final state and large pion rejection (removed pion band).

The $K_{\mu\nu\gamma}$ decay is dominant background source.

BACKGROUND SHAPE

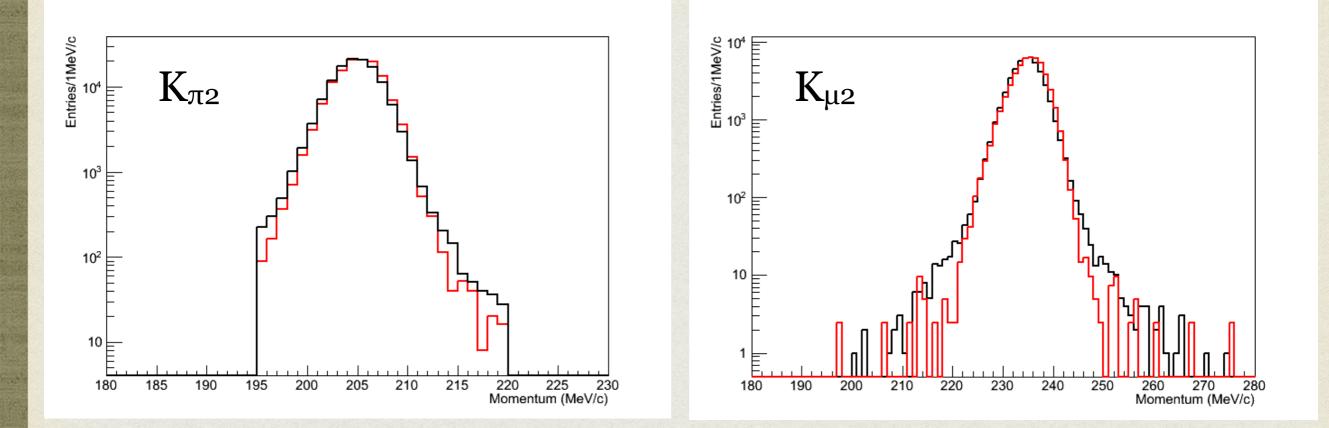


$$k_{K_{\mu 2}} = \frac{(KBlive)_{1/20}}{N_{sim,K_{\mu 2}}} \times \frac{A_{K_{\mu 2}}^{Exp}}{A_{K_{\mu 2}}^{MC}} \times BR(K_{\mu 2})$$

 $k_{K_{\mu\nu\gamma}} = \frac{(KBlive)_{1/20}}{N_{sim,K_{\mu\nu\gamma}}^{140$

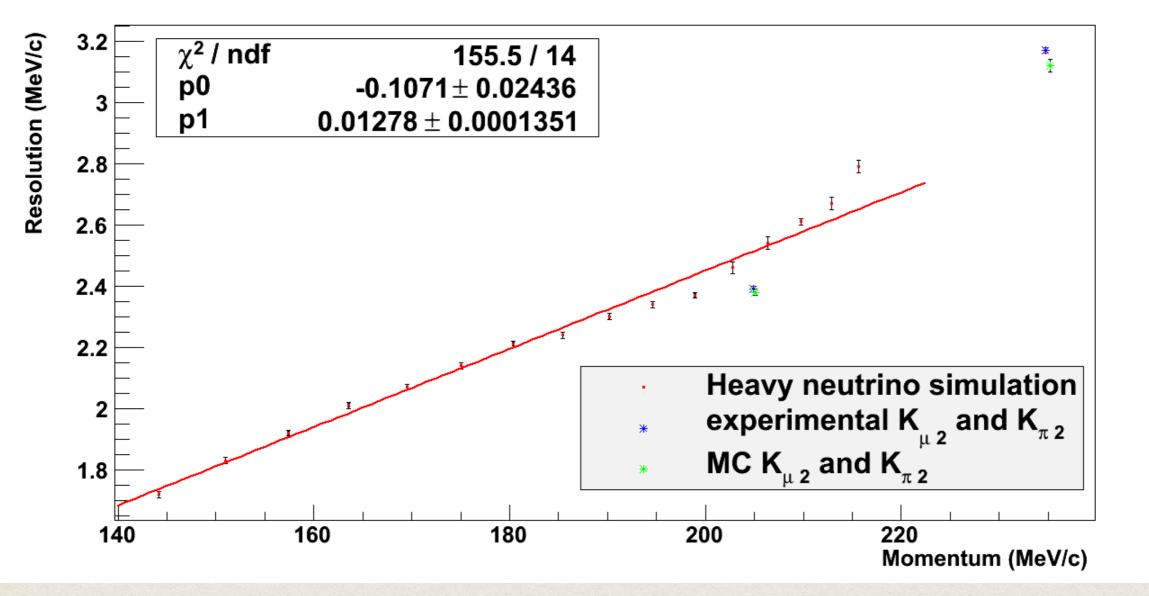
Obviously, MC is not the best fit of data, but we are not going to use the simulated background shape in the further analysis. MC shape does not have obvious bumps or valleys, so the only conclusion from simulation is to assume that experimental background shape also should be smooth, but we do not know exact background shape.

MOMENTUM RESOLUTION



MC simulation was consistent with data, so we used the $K^+ \rightarrow \mu^+ \nu_H$ simulation to measure muon momentum resolution

MOMENTUM RESOLUTION



 $\sigma(p)(MeV/c) = (-0.1071 + 0.01278 \times p(MeV/c)) \pm 0.14 \pm 0.05$

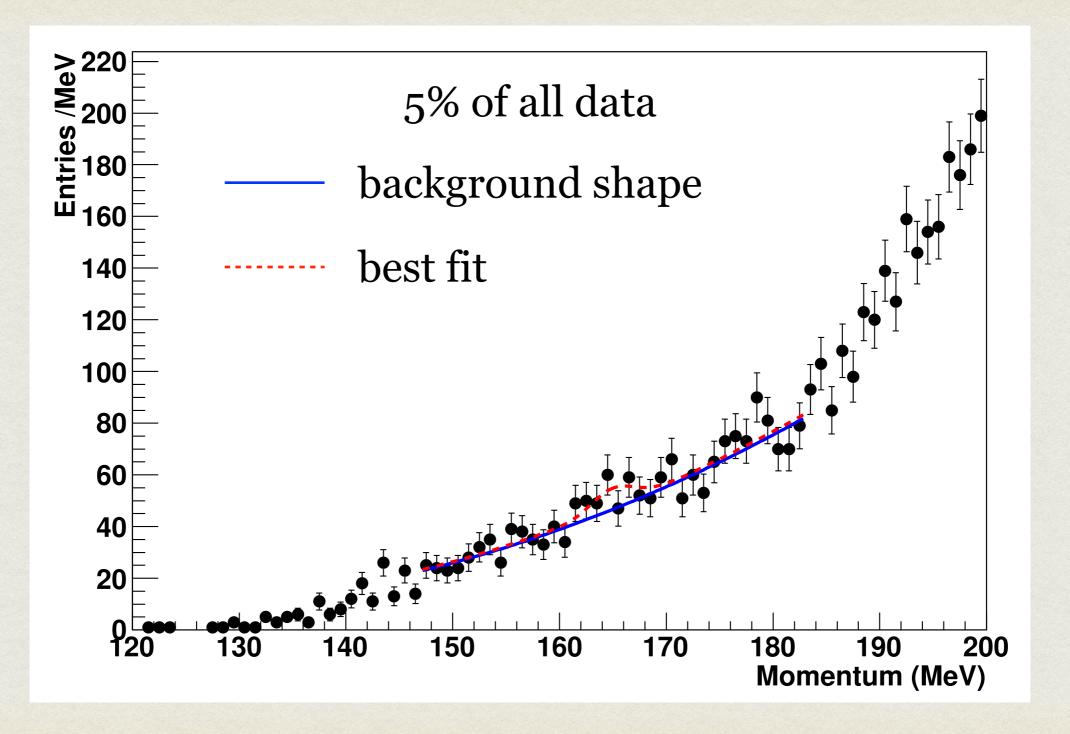
- We don't know exact background shape due to very low acceptance for background. But it is not really necessary because the background fitting is data-driven
- Define shape locally: choose $\pm 6\sigma$ region around the point of interest (σ is momentum resolution for the selected point) and fit it by 2 order polynomial function
- Use Gaussian with well known sigma as a signal
- Use likelihood approach which is approved by Higgs search to get upper limit (Eur.Phys.J.C71:1554,2011)

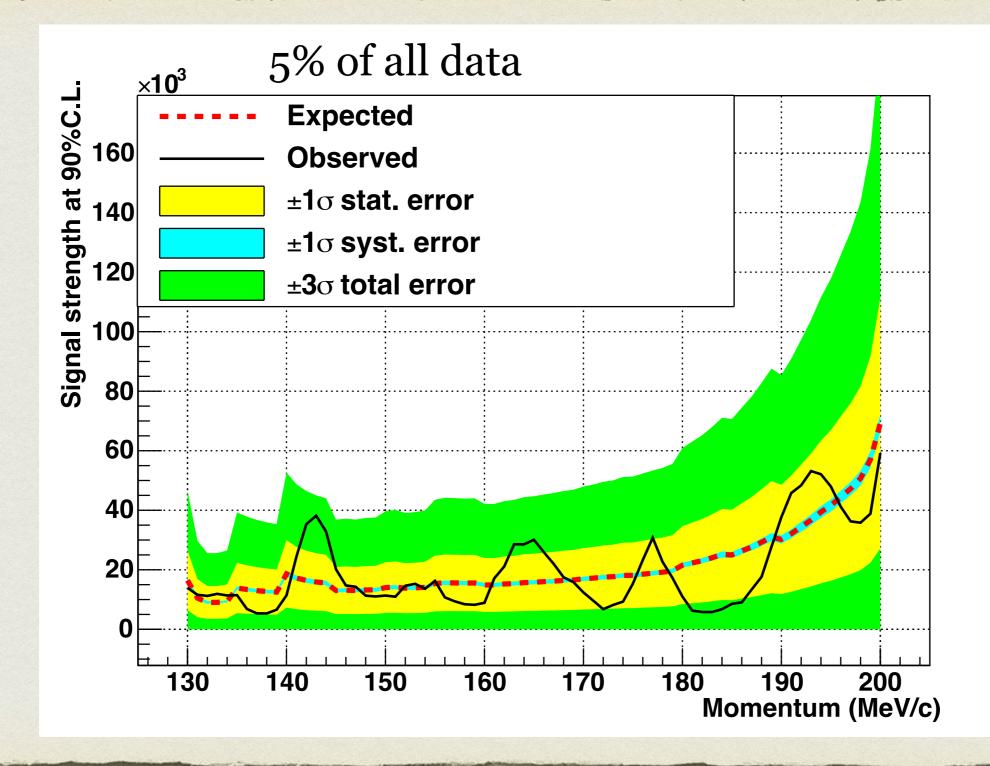
• Construct the following likelihood function

$$L(\mu, \theta) = \{\prod_{i=1}^{Nbins} \frac{(\mu \cdot \beta s_i + \theta b_i)^n}{n!} e^{-(\mu \cdot \beta s_i + \theta b_i)}\} \times Gauss(\beta; \beta_{peak}, \sigma_{\beta_{peak}})$$

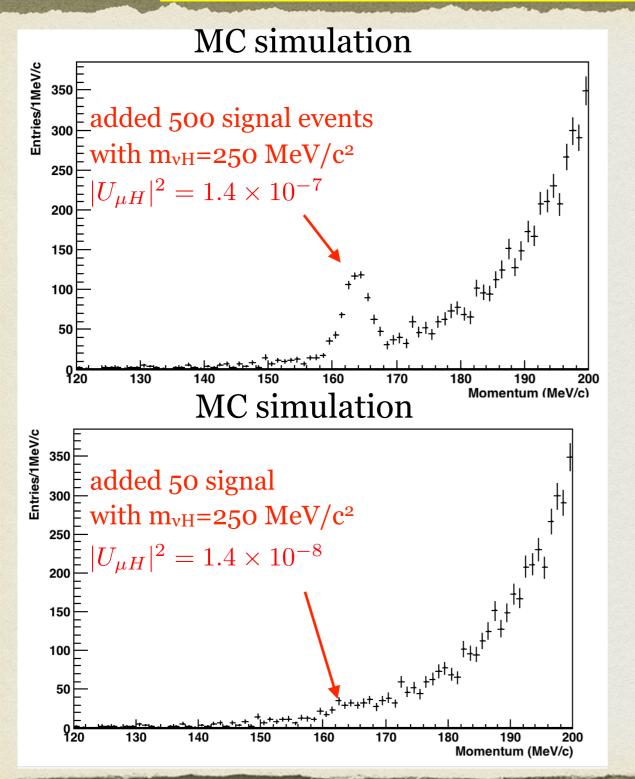
s and b – signal and background distributions. Since background distribution is taken from data fit, $\theta = 1$, μ is signal strength parameter; s – gauss. n_i – number of observed events in each bin; β takes into account acceptance of the point of interest (β_{peak}) and its total error (σ_{β})

- "Asimov" data set based on background shape was generated to calculate expected upper limit (background only hypothesis) at 90%C.L.
- Use the same upper limit calculation method but with experimental data to get observed upper limit at 90%C.L.

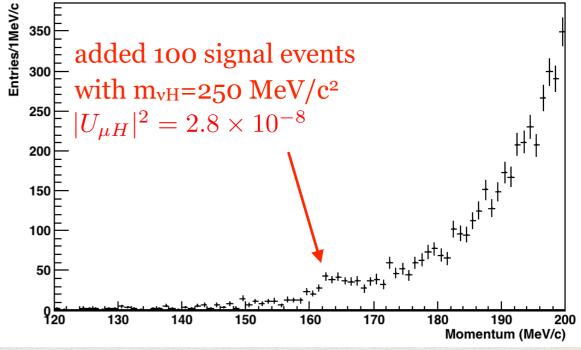




SIGNAL ILLUSTRATION THERE ARE NO REAL DATA POINTS!!!

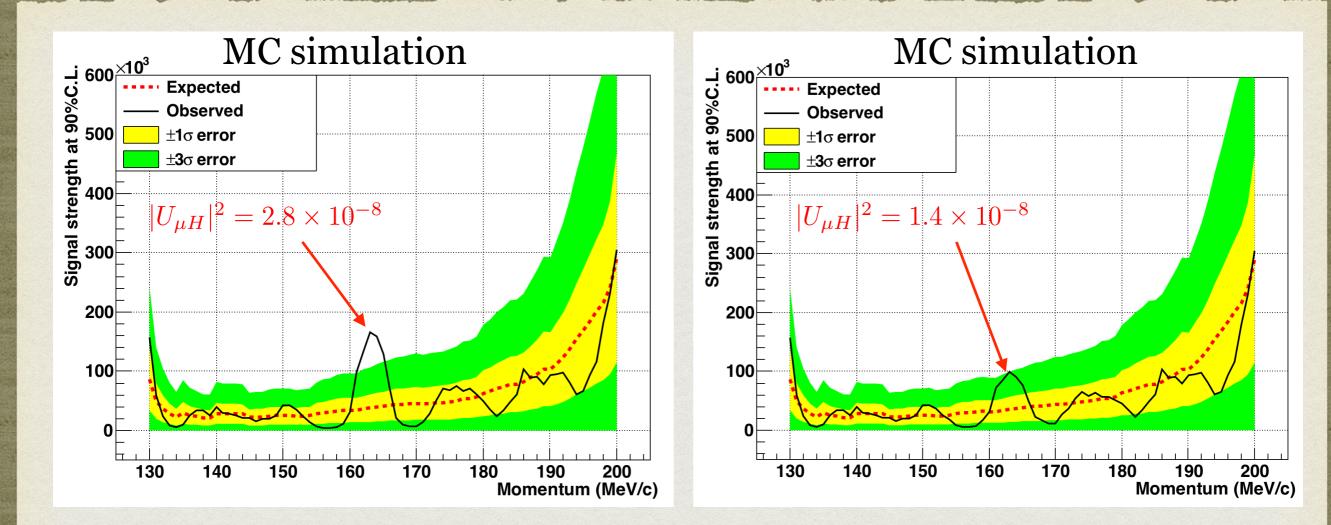


MC simulation



The 5% of all data was fitted by exponent in the momentum region 120—200 MeV/c. The background was generated in ROOT by using the fit parameters. The signal events were added to the generated background

SIGNAL ILLUSTRATION THERE ARE NO REAL DATA POINTS!!!



Yes! We found the peak at 5.6σ

added 100 signal events with $m_{\nu H}$ =250 MeV/c²

We are not able to find the peak

added 50 signal events with $m_{\nu H}$ =250 MeV/c²

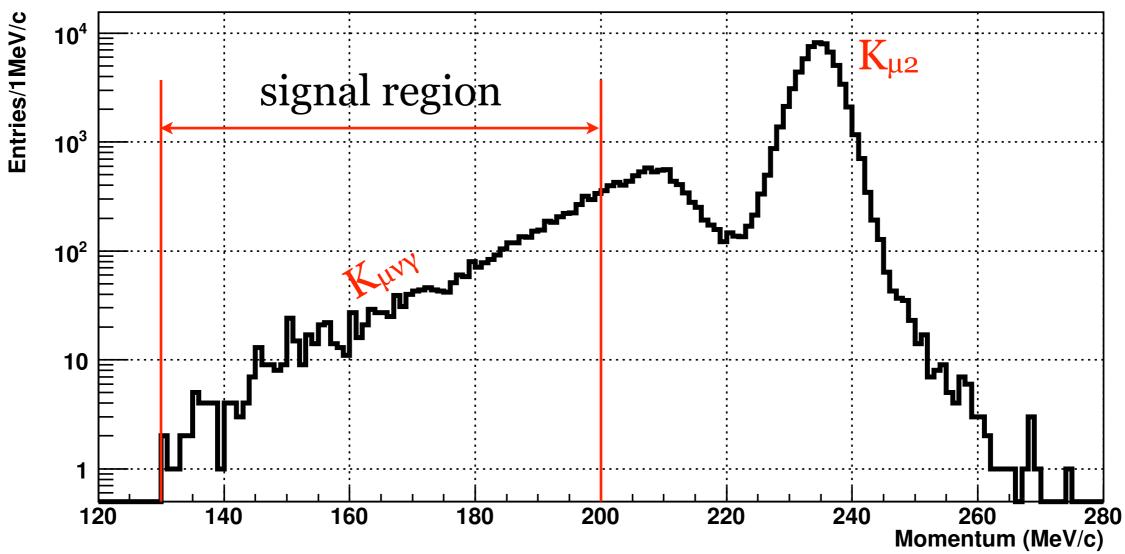
RESULTS

- 1. Full data analysis
- 2. Interpretations
- 3. Conclusions

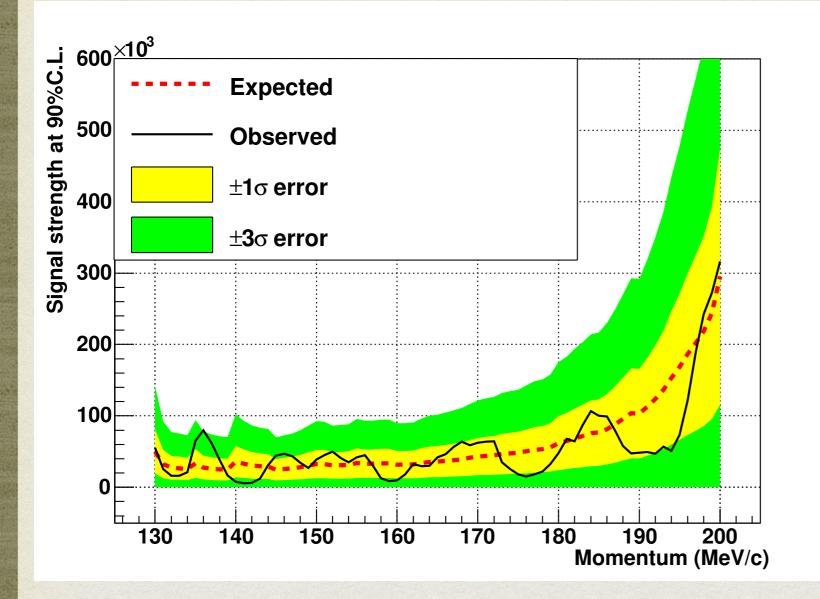


FULL DATA SAMPLE

after all selection criteria (100% of data)



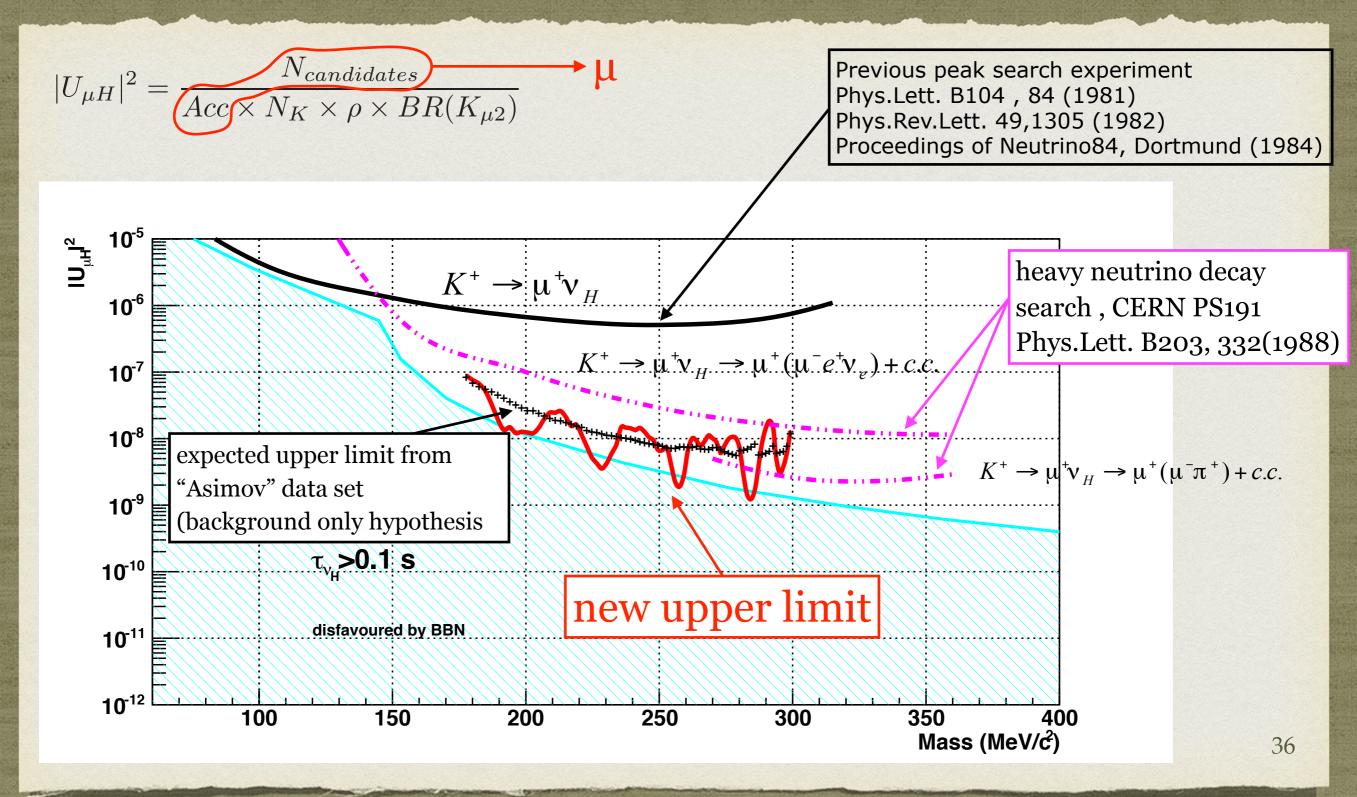
PEAK SEARCH



No peaks were found No evidence for the heavy neutrino signal Set upper limit on the

Set upper limit on the mixing matrix element between muon and heavy neutrino

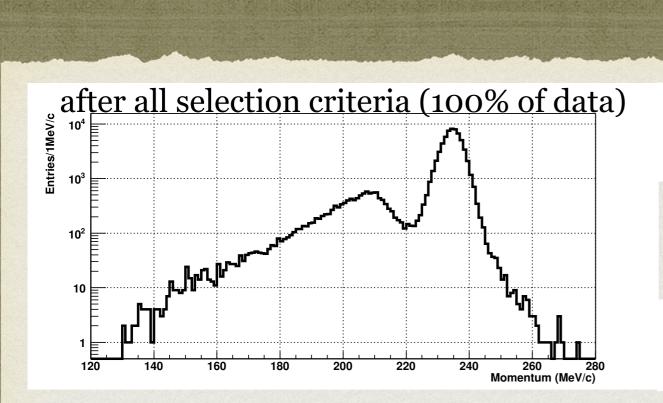
MIXING MATRIX ELEMENT



$K^+ \rightarrow \mu^+ X DECAY$

- X is invisible set of neutral particles
- We are able to set a limit on any decay with a single muon
- For example, the K⁺ \rightarrow µ⁺vvv decay, which hasn't been probed since 1973
 - BR ~10⁻¹⁶ in SM (recent calculation, arXiv:1605.08077)
 - Can provide information on two effects: neutrino-neutrino interaction (Phys. Lett. B32, 121–124 (1970), arXiv:hep-ph/ 9908272) and six fermion interaction (Phys. Rev. 133, B130–B131 (1964))
 - PDG value 6.0×10⁻⁶ based on neutrino-neutrino interaction model (1973!)

$K^+ \rightarrow \mu^+ X DECAY$



$$BR_P(K^+ \to \mu^+ + X) < \frac{1}{N_K} \sum_{i=130}^{P} \frac{N_i}{Acc_i} + 1.28\sigma,$$

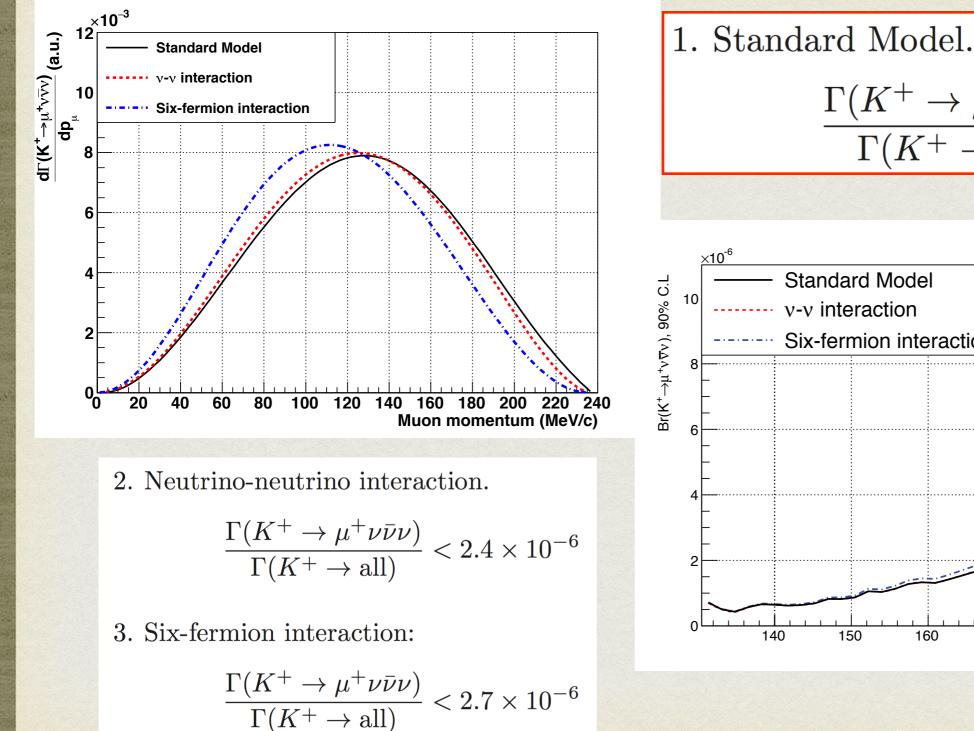
Acceptance for single muon selection is already known

 $\begin{array}{l} BR(K^+ \to \mu^+ + X, 130 < p_\mu < 175 \ {\rm MeV/c}) < \\ < 7.5 \times 10^{-7} \end{array}$

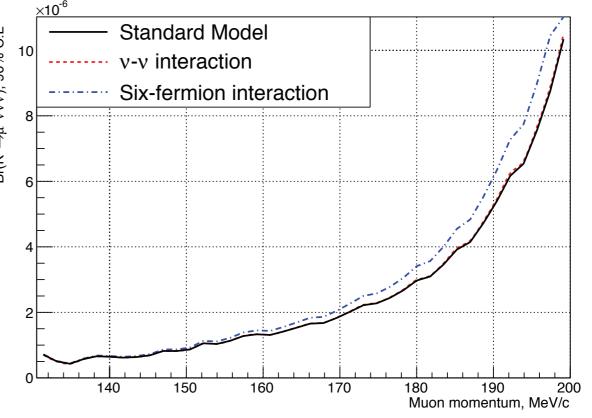
To get total decay rate we need to know the signal shape

$$\frac{\Gamma(K^{+} \to \mu^{+} \nu \bar{\nu} \nu)}{\Gamma(K^{+} \to \text{all})} = 7.5 \times 10^{-7} \times \frac{\int_{0}^{p_{\mu}^{max}} (d\Gamma/dp_{\mu}) dp_{\mu}}{\int_{130}^{175} (d\Gamma/dp_{\mu}) dp_{\mu}} \text{ here can be any set of neutral particles which are invisible in the detector}$$
More general, if you do not like 130–175 MeV/c region:
$$\frac{\Gamma(K^{+} \to \mu^{+} \nu \bar{\nu} \nu)}{\Gamma(K^{+} \to \text{all})} = BR_{P}(K^{+} \to \mu^{+} + X) \times \frac{\int_{0}^{p_{\mu}^{max}} (d\Gamma/dp_{\mu}) dp_{\mu}}{\int_{130}^{P} (d\Gamma/dp_{\mu}) dp_{\mu}}$$
38

$K^+ \rightarrow \mu^+ \nu \nu \nu DECAY$



$$\frac{\Gamma(K^+ \to \mu^+ \nu \bar{\nu} \nu)}{\Gamma(K^+ \to \text{all})} < 2.4 \times 10^{-6}$$



CONCLUSIONS

- ★ The heavy neutrino existence in the mass region 175-300 MeV/c² was tested using E949 experimental data set
- \star No evidence was found
- ★ Previous best constraints from CERN PS191 were improved by order of magnitude
- ★ New mixing matrix element $|U_{\mu H}|^2$ upper limit is varying between 10⁻⁹ and 10⁻⁸
- ★ In contrast to CERN PS191 or BBN lower limit our result is modelindependent.
- ★ No evidence for $K^+ \rightarrow \mu^+ \nu \nu \nu$ decay was found
- ★ Improved the current limit by factor ~3. For the first time limit is presented in the framework of the standard model
- ★ Presented procedure to calculate upper limit on any possible kaon decay to muon and invisible set of neutral particles for any assumed muon momentum spectrum

THANK YOU!



EXTRA SLIDES

PHOTON INEFFICIENCY

	260							
Energy (MeV)	240			(19±19) ×10 ⁻⁶		(75±37) ×10 ⁻⁶		
	220		(11±6)	(76+17)	(44±11)	(70+16)	(21±21)	
	200		×10 ⁻⁴ (17±12)	×10 ⁻⁶ (60±15)	×10 ⁻⁶ (64±13)	×10 ⁻⁶ (98±19)	×10 ⁻⁵ (54±18)	
	180		×10 ⁻⁵	×10 ⁻⁶	×10 ^{-₀}	×10 ⁻	×10 ⁻⁵	
			(26±9) ×10 ⁻⁵	(12±2) ×10 ⁻⁵	(85±16) ×10 ⁻⁶	(12 <u>±</u> 2) ×10 ⁻⁵	(55±12) ×10 ⁻⁵	(23 <u>±</u> 8) ×10 ⁻⁴
	160	(17±12) ×10 ⁻⁵	(25±7) ×10 ⁻⁵	(14±3) ×10⁵	(83±17) ×10 ⁻⁶	(15±3) ×10⁵	(81±11) ×10 ⁻⁵	(12±2) ×10 ⁻⁴
	140	(54±13) ×10 ⁻⁵	(31±6) ×10 ⁻⁵	(18±3) ×10 ⁻⁵	(11±2) ×10 ⁻⁵	(20±3) ×10 ⁻⁵	(89±10) ×10 ⁻⁵	(17±2) ×10 ⁻⁴
Ener	120	(25±2 ×10 ⁻⁴	(25±5) ×10 ⁻⁵	(24±4) ×10 ⁻⁵	(20±3) ×10 ⁻⁵	(26±4) ×10 ⁻⁵	(10±1) ×10 ⁻⁴	(16±1) ×10 ⁻⁴
	100	(72±2) ×10 ⁻⁴	(62±8) ×10 ⁻⁵	(47±5) ×10 ⁻⁵	(46±5) ×10 ⁻⁵	(38±5) ×10 ⁻⁵	(18±1) ×10 ⁻⁴	(25±1) ×10 ⁻⁴
	80	(219+1)	(16±1)	(11±1)	(96± <u>8</u>)	(92± <u>8</u>)	(27±1)	(60±1)
	60	×10 ⁻⁴ (451±5	×10 ⁻⁴ ′ (48±2)	×10 ⁻⁴ ′ (20±1)	×10 ⁻⁵ (17±1)	×10 ⁻⁵ (24±1)	×10 ⁻⁴ (56±2)	×10 ⁻⁴ (129±2)
	40	×10 ⁻⁴	×10 ⁻⁴	×10 ⁻⁴	×10⁻⁴	×10 ⁻⁴ ′ (104±2)	×10 ⁻⁴	×10 ⁻⁴
	20	(603±9) ×10 ⁻⁴	(252±2) ×10 ⁻⁴	(144±2) ×10 ⁻⁴	(106±2) ×10 ⁻⁴	(104±2) ×10 ⁻⁴	(142±2) ×10 ⁻⁴	(239±1) ×10 ⁻⁴
		(70±57) ×10⁻⁴	(11±7) ×10 ⁻³	(63±52) ×10 ^{−3}				(53±46) ×10 ⁻⁴
	0	1	-0.5		0		0.5	1
					$\cos \theta_{\gamma}$			

Obtained ned by fitting $K^+ \rightarrow \pi^+\pi^0 \rightarrow \pi^+\gamma\gamma$ where the π^+ and one γ are measure. This determines the second γ 's energy and direction.

MUON SAMPLES

MC simulation

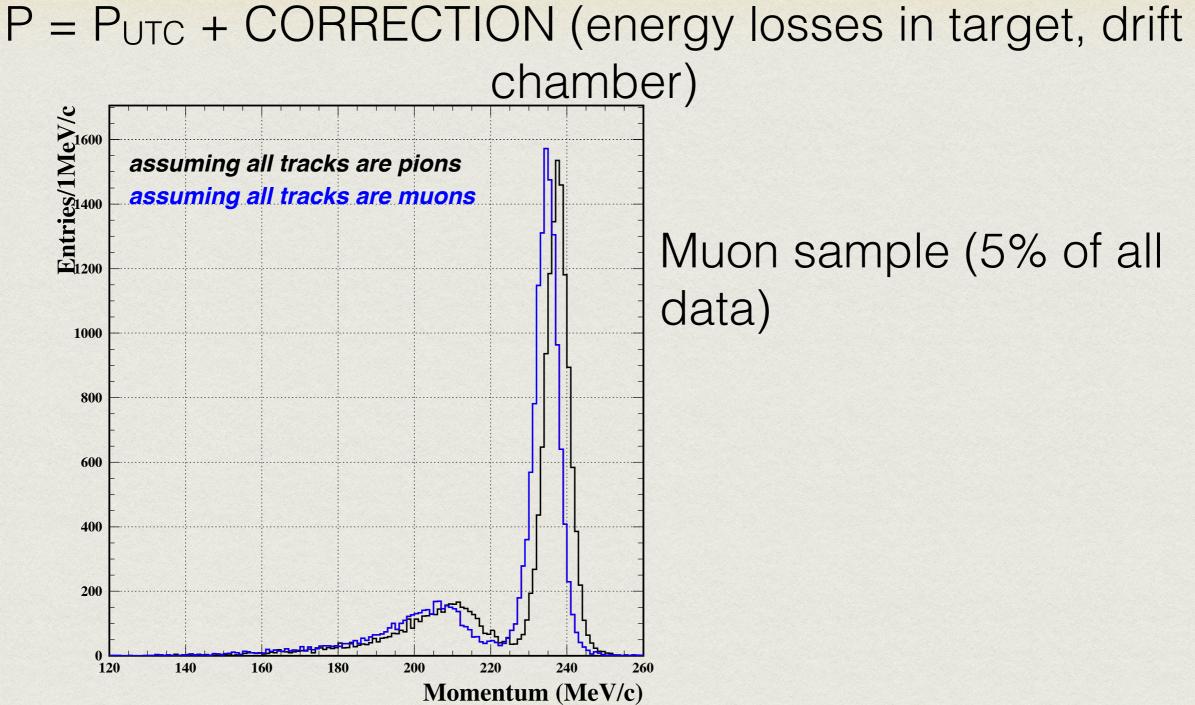
Roles		$K^+ \to \mu^+ \nu_\mu$	$K^+ \to \mu^+ \nu_\mu \gamma$	$K^+ \to \mu^+ \nu_H$
Acceptance	trigger			\checkmark
measurement for	L1n		\checkmark	
the $K^+ \to \mu^+ \nu_H$ decay	Refined Range		\checkmark	
	kin			\checkmark
	beam⌖	\checkmark		
	RNGMOM		\checkmark	
	PV	\checkmark		
Background study		\checkmark		
Resolution	\checkmark		\checkmark	

We used muons from the $K_{\mu 2}$ decay to measure the acceptance, in general. Muons from this decay have certain momentum. To study acceptance vs momentum we used $K_{\mu \nu \gamma}$ decay.

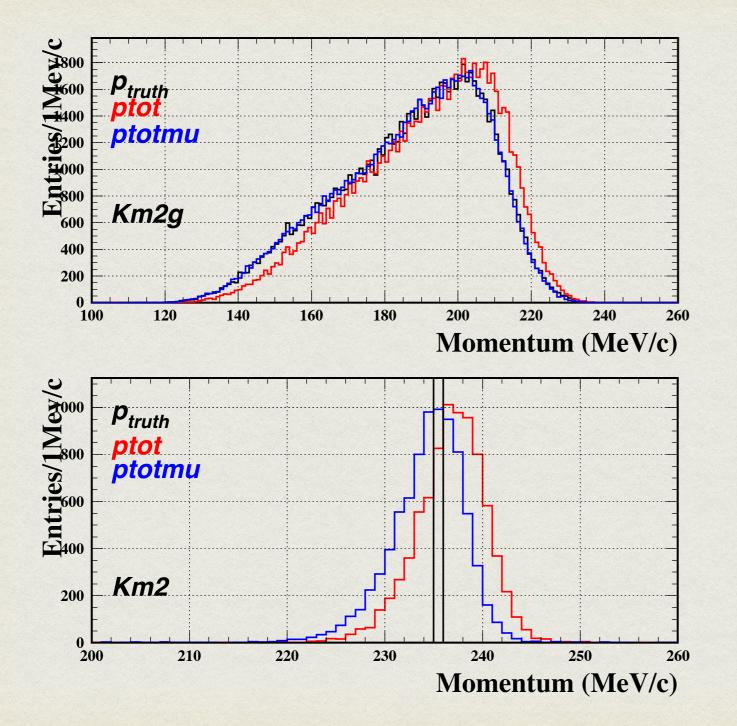
ACCEPTANCE MEASUREMENT

Acceptance factors	$K^+ \to \mu^+ \nu_H$	$K^+ \to \mu^+ \nu_\mu \gamma$	$K^+ \to \mu^+ \nu_\mu$
f_s	Κπ		
A_{PRRF}	π		
$A_{UTCQUAL}$	π_{scat}		
$\epsilon_{T \bullet 2}$	Kb		
$A_{Fid\&Range}$			
$A_{Kinematic}$	N		
$A_{\overline{19_{ct}}}$		Κμ2	
$A_{RefinedRange}$			$K\pi 2(1)$
$A_{\pi \to \mu}$	$K\pi$		
ARNGMOM		not applied	
$A_{Beam\&Target}$	Κμ2		
A_{DELC}			
A_{PV} Loose	not applied	MC	
Tight	$K\mu 2$	not applied	not applied

MOMENTUM MEASUREMENT

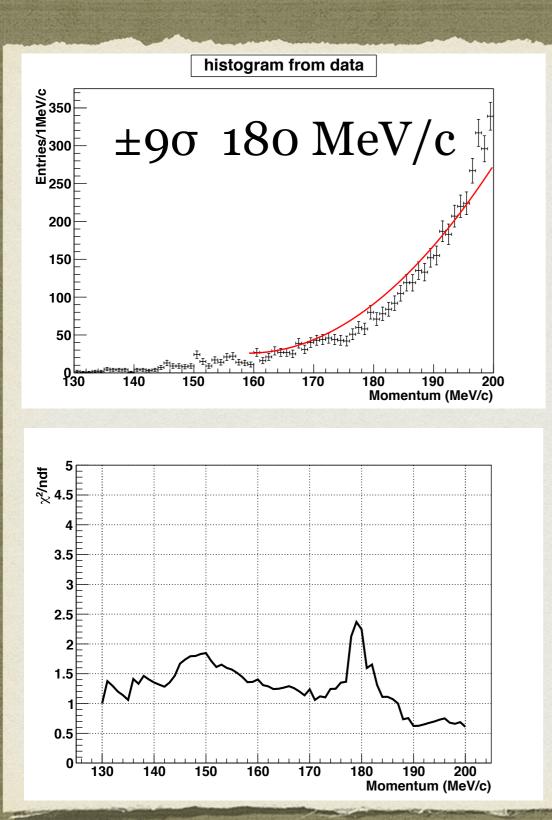


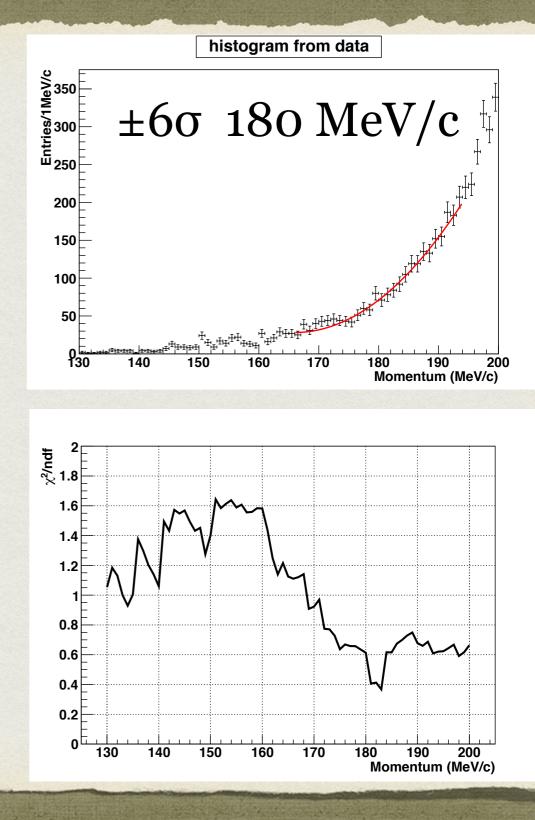
MOMENTUM MEASUREMENT



MC simulation

$\pm 9\sigma \rightarrow \pm 6\sigma$

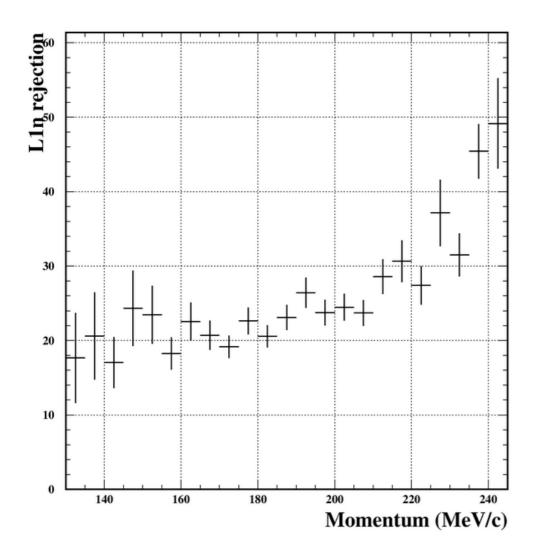




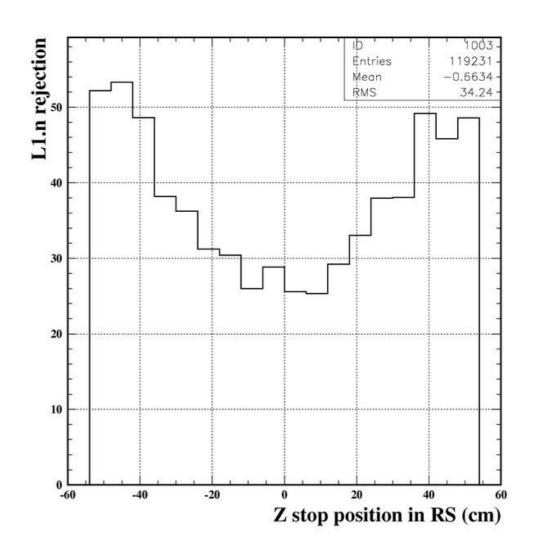
ONLINE PION IDENTIFICATION (L1N)

- Online $\pi \rightarrow \mu$ decay requirement:
 - The pulse height (PH) and the area (PA) of the pulse(s) recorded by TDs in the stopping counter are compared. The ratio PH/PA would be smaller for double pulses than for a single pulse
- Requirements for stopping counter finding:
 - Events are rejected if coincident hits are detected in a RS near the stopping counter
 - Events are rejected if one of two adjacent hextant hits is not due to a charged track

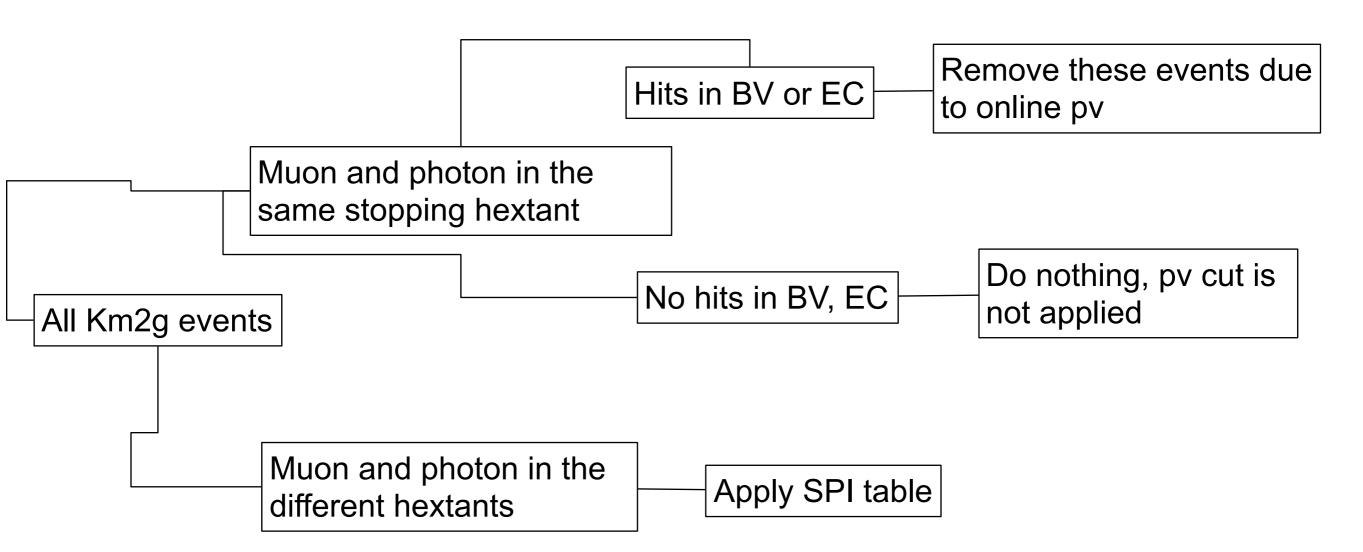
ONLINE PION IDENTIFICATION (L1N)



<u>We use Z-dependence</u> of L1n cut to emulate it in MC Use monitor trigger (KP21) and some offline cuts to select pure muon sample: KP21, UTCQUAL, COS3D, ZFRF, ZUTOUT, PRRF, B4DEDX, PV90, RNGMOM. <u>No TD cuts</u>



PV acceptance measurement



Use MC simulation of the Km2g decay within pnn1 or pnn2 or pnn12 trigger w/o online photon veto

Thresholds in BV and EC are 5 MeV and 20 MeV respectively. For 140<ptotmu<200 MeV/c gamma energy E>20 MeV