



# Unit 16

## Degradation and training

### Episode I

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# Outline



1. Introduction
2. Quench of a superconductor
3. Classification of quenches
4. Conductor limited quenches
  - Measurements of critical surfaces
  - Short sample current
  - Degradation
5. Energy deposited quenches
  - Distributed disturbances
    - Minimum quench energy density
  - Point disturbances
    - Minimum quench energy
    - Minimum propagation velocity
6. Practical example: SSC and HERA wires
7. Cryogenic stabilization
8. Practical example: LHC wire
9. Conclusions



# References



- [1] A. Devred, "*Quench origins*", AIP Conference Proceedings 249, edited by M. Month and M. Dienes, 1992, p. 1309-1372.
- [2] M. Wilson, "*Superconducting magnets*", Oxford UK: Clarendon Press, 1983.
- [3] A. Godeke, "*Performance boundaries in Nb<sub>3</sub>Sn superconductors*", PhD thesis, 2005.
- [4] A. Devred, "*Practical low-temperature superconductors for electromagnets*", CERN-2004-006, 2006.
- [5] Y. Iwasa, "*Mechanical disturbances in superconducting magnets – A review*", IEEE Trans. Magn., Vol. 28, No. 1, January 1992, p. 113-120.
- [6] L. Rossi, "*Superconducting Magnets*", CERN Academic Training, 15-18 May 2000.
- [7] K.-H. Mess, P. Schmuser, S. Wolff, "*Superconducting accelerator magnets*", Singapore: World Scientific, 1996.
- [8] D. Leroy, "*Review of the R&D and supply of the LHC superconducting cables*", IEEE Trans. Appl. Supercond. , Vol. 16, No. 2, June 2006, p. 1152- 1159.
- [9] A.K. Ghosh, et al., "*Minimum quench energies of rutherford cables and single wires*", IEEE Trans. Appl. Supercond. , Vol. 7, No. 2, June 1997, p. 954-957.
- [10] B.J. Maddok, et al., "*Superconductive composites: heat transfer and steady state stabilization*", Cryogenics 9, 1969, p. 261-273.



# 1. Introduction

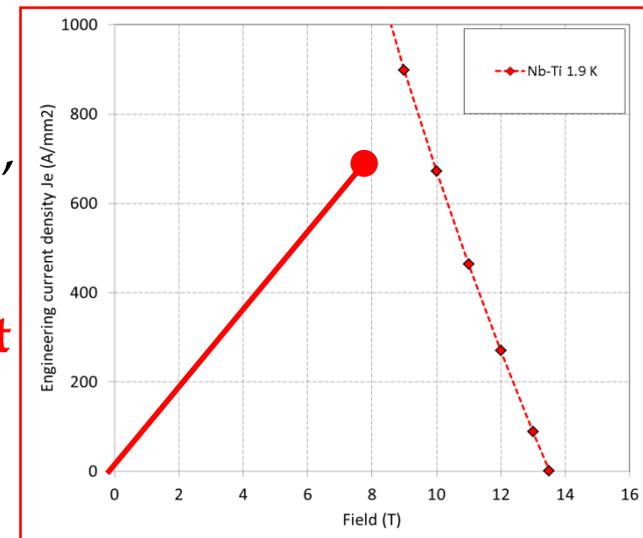
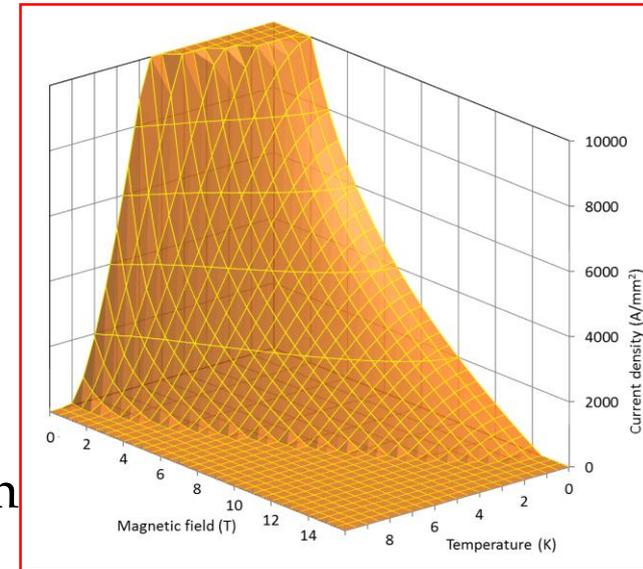


- In this unit we will address the following questions
  - What is a quench?
  - How can we classify a quench according to the different causes?
  - What is the maximum current that a wire can carry and how can we measure it?
  - Which are the volumes and the energies involved in the quench phenomenon?
  - What is the role of the stabilizer and of the liquid helium?



## 2. Quench of a superconductor

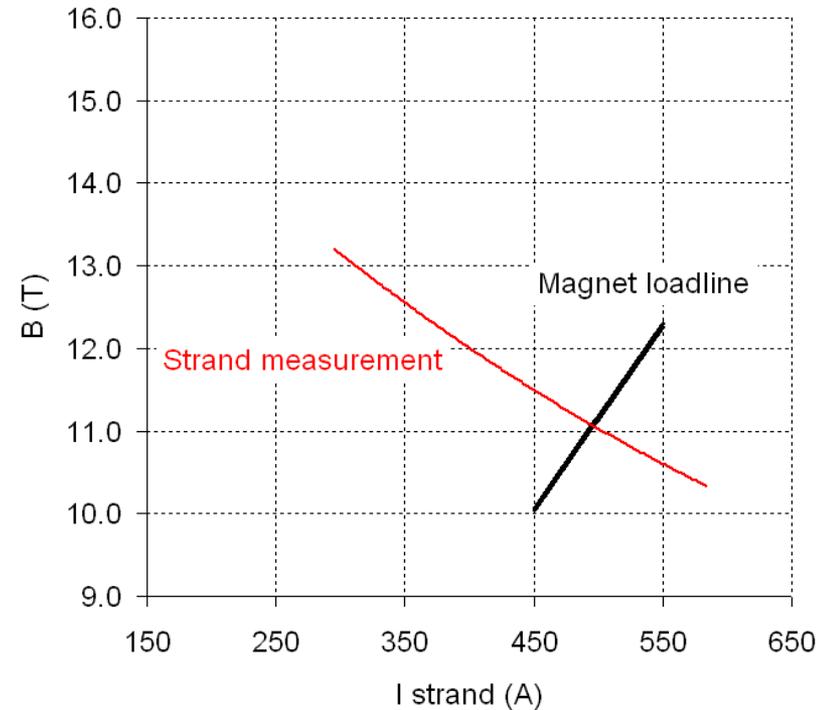
- The superconducting state is defined by the **critical surface**
  - $B$  (T),  $J$  (A/mm<sup>2</sup>),  $T$  (K)
- A superconducting magnet operates in conditions corresponding to a point located beneath the critical surface, and defined by  $T=T_{op}$ ,  $J=J_{op}$ , and  $B=B_{op}$ .
- Let's assume (A. Devred, [1]) that, starting from the operational conditions, we **increase the current** in the magnet, and, as a result, the magnetic field.
- At a certain point, the critical surface is crossed, and a small volume  $V$  of superconductor **becomes normal**.
- Therefore, the volume  $V$  starts **dissipating heat** because of Joule effect, and its temperature increases.





## 2. Quench of a superconductor

- Due to **thermal diffusion**, also the surrounding volume  $dV$  undergoes an increase in temperature.
- If the heat dissipated by  $V$  is **enough**,  $dV$  reaches the critical temperature, becomes normal conducting, and dissipates heat.
- In certain conditions, the normal zone propagates through the coil: this phenomenon is called **quench**.
- Depending on its causes, a quench can be classified and defined in different ways.





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### 3. Classification of quenches



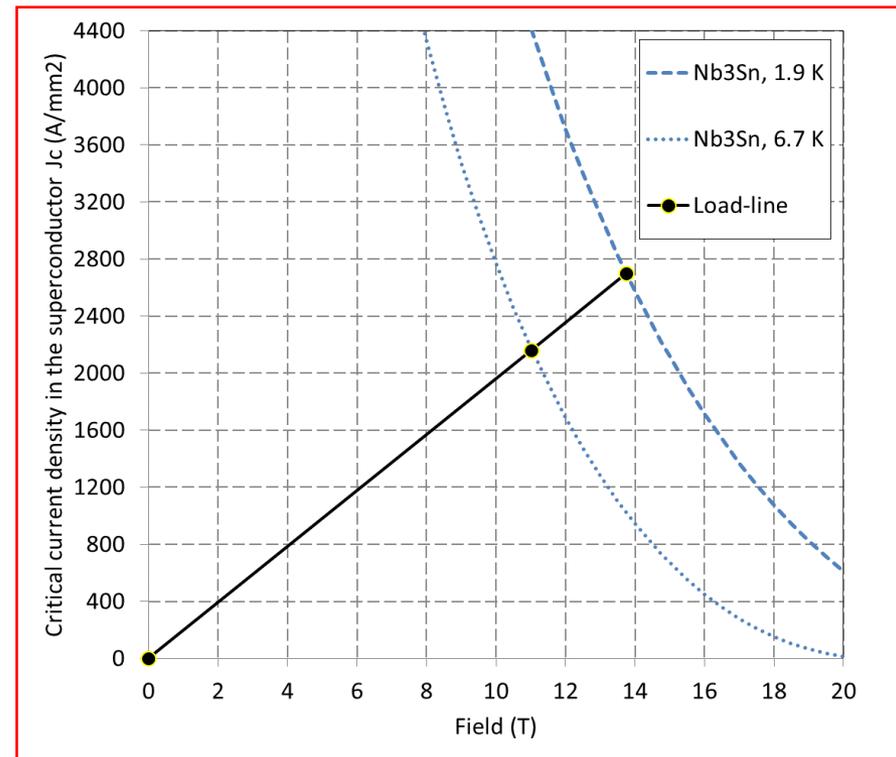
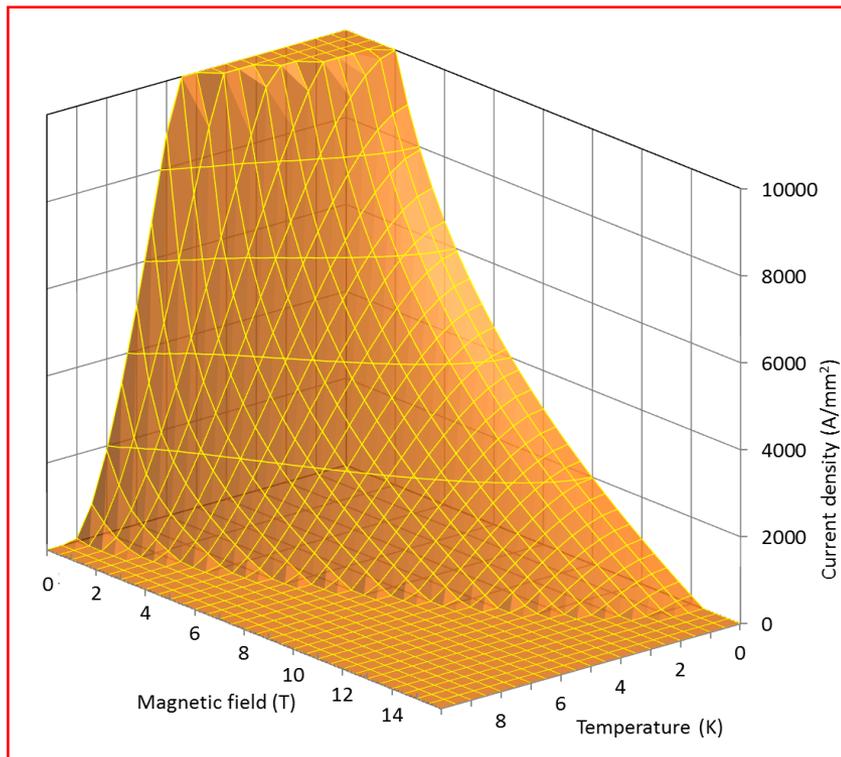
- A first classification (A. Devred, [1])
  - The maximum field seen by the conductor in the coil is usually called **peak field** ( $B_{peak}$ ).
  - At a given temperature  $T_0$ , the maximum current that the conductor can reach will be  $I_{max} = I_c(B_{peak}, T_0)$ , where  $I_c$  is the critical current at  $B_{peak}$  and  $T_0$ .
  - When the magnet quenches, we can have either  $I_{quench} = I_{max}$  or  $I_{quench} < I_{max}$ .
  - In the first case we have a **conductor-limited quench** (or **exhausted margin quench**).
    - If the maximum current is consistent with measurements performed on wire short samples, we have a **short-sample quench** (success!!)
    - If the current is lower, we have **degradation**.
  - In the second case, a quench occurs because of a release of energy that increases the temperature of the conductor beyond the critical temperature. These are called **energy-deposited quenches** or, also, **premature quenches**.



# 3. Classification of quenches



- In other words, in the first case the critical surface is crossed because of an **increase of  $I$  (and  $B$ )**; in the second case the critical surface is crossed because of an **increase of temperature**.





### 3. Classification of quenches



- A second, more detailed classification (M. Wilson, [2])
  - We can define a **spectrum of disturbances**, which classifies the energy disturbances along two dimensions: time and space.

		Space	
		Point	Distributed
Time	Transient	J	J/m <sup>3</sup>
	Continuous	W	W/m <sup>3</sup>

- **Continuous disturbances** are due to a steady power dissipation in the coil
  - Point: ramp splice with high resistance joint
  - Distributed: a.c. losses in the conductor, thermal leak of the cryogenic system.
- They are usually well understood disturbances.



### 3. Classification of quenches



		Space	
		Point	Distributed
Time	Transient	J	$J/m^3$
	Continuous	W	$W/m^3$

- **Transient quenches** are due to a sudden release of energy, either over a small volume (J) or over a large volume ( $J/m^3$ )
  - **Flux jumps**: dissipative redistribution of magnetic field within the superconductor. It can be eliminated with small filaments.
  - **Mechanical disturbances**: wire frictional motion, epoxy cracking. They are less predictable and difficult to avoid, since they are related to mechanical design, material properties, fabrication processes, etc.



# Outline

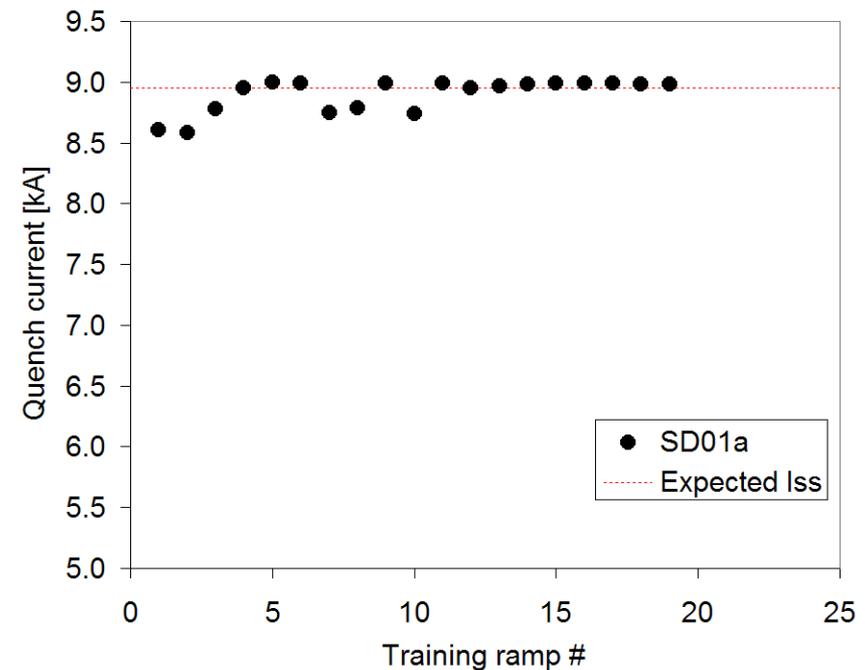


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## 4. Conductor limited quenches

- Conductor limited quenches are usually very stable.
- A series of conductor limited quench in a “quench current vs. quench number” graph appears as a stable plateau. For these reasons they are also called *plateau quenches*.
- After having reached the maximum current of the magnet, we have to compare it with the critical current measured on a short sample of the conductor.
- If the two coincide, we say that the magnet reached *short sample current  $I_{ss}$* .





## 4. Conductor limited quenches

### Measurements of the conductor critical current



- The critical current of a conductor is measured by winding a sample of the wire around a **sample holder**.
- To avoid premature quenching induced by Lorentz forces during ramping, the wire must be well supported
  - Stycast glue may be used to constrain the wire around the holder
- In case of  $\text{Nb}_3\text{Sn}$  wires, a sample holder made of titanium is used.
- Once the wire is cooled-down and placed in a given magnetic field, the current is increased until the **transition occurs**.



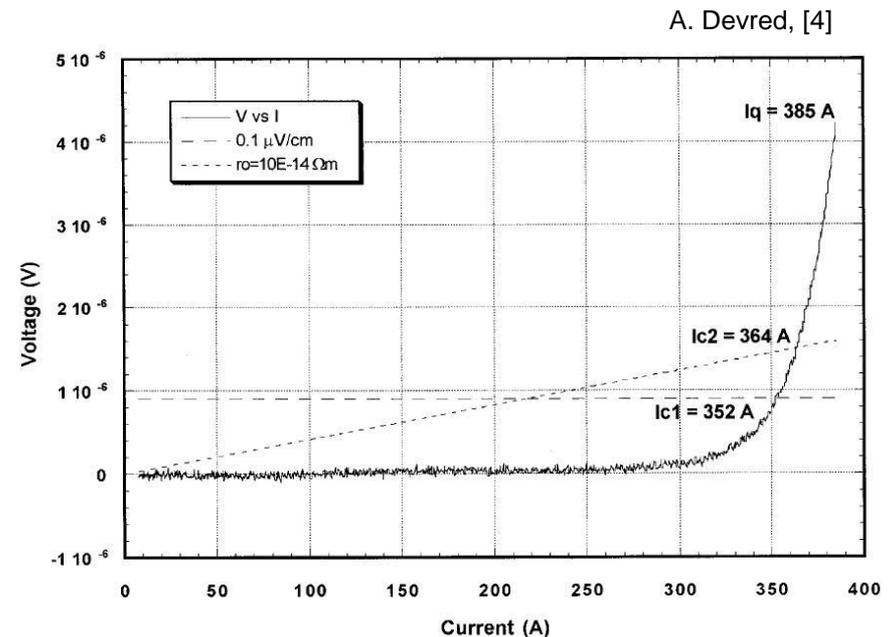
A. Devred, [4]



## 4. Conductor limited quenches

### Measurements of the conductor critical current

- The transition from superconducting to normal state is observed through an **increase of voltage**.
- The increase is initially slow and reversible. By further increasing the current, the **voltage rise becomes more steep**, until it takes off and becomes irreversible: we refer in this case to short sample quench current  $I_{qss}$ .
- In different conditions (wire as a part of a cable in a magnet, in different cooling conditions) the quench may occur at a different current.





# 4. Conductor limited quenches

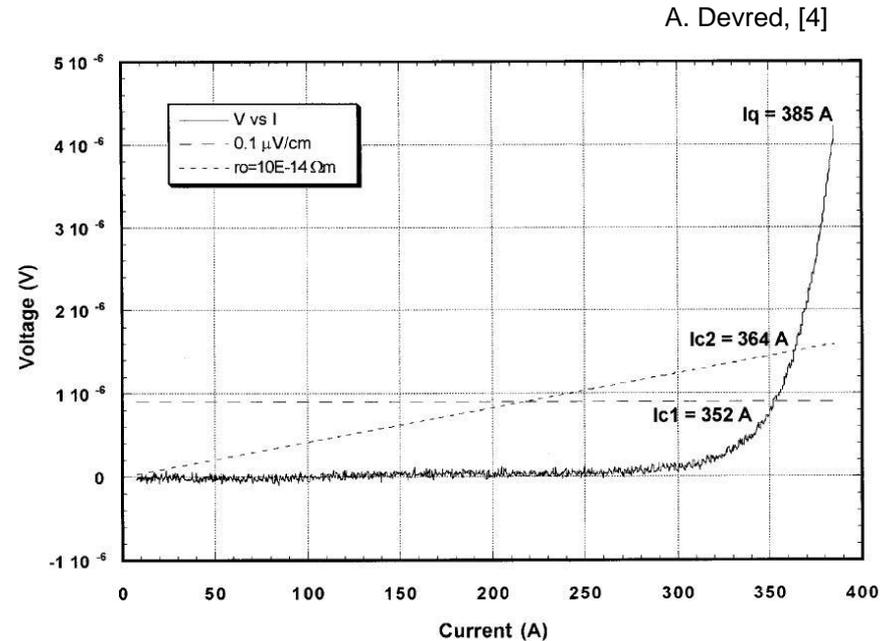
## Measurements of the conductor critical current

- The conventions to determine the quench current are the following [4].
  - Let's assume that we have a wire of length  $L$  and cross-sectional area  $S$ .
  - Being  $\lambda$  the copper to supercond. or copper to non-copper ratio, we define as **apparent resistivity of the superconductor and apparent electrical field**

$$\rho_{sc} = \frac{1}{1 + \lambda} \frac{V S}{L I} \quad E = \frac{V}{L}$$

where  $V$  is the measured voltage across the wire.

- The critical current  $I_c$  is defined as the current where  $\rho_{sc} = 10^{-14} \Omega\text{m}$  or  $E = 0.1 \mu\text{V/cm}$ .



“It has been verified, in particular for accelerator magnets, that the critical currents defined above can be used to make fairly accurate estimations of the maximum quench current of a superconducting magnet” [4].



## 4. Conductor limited quenches Degradation



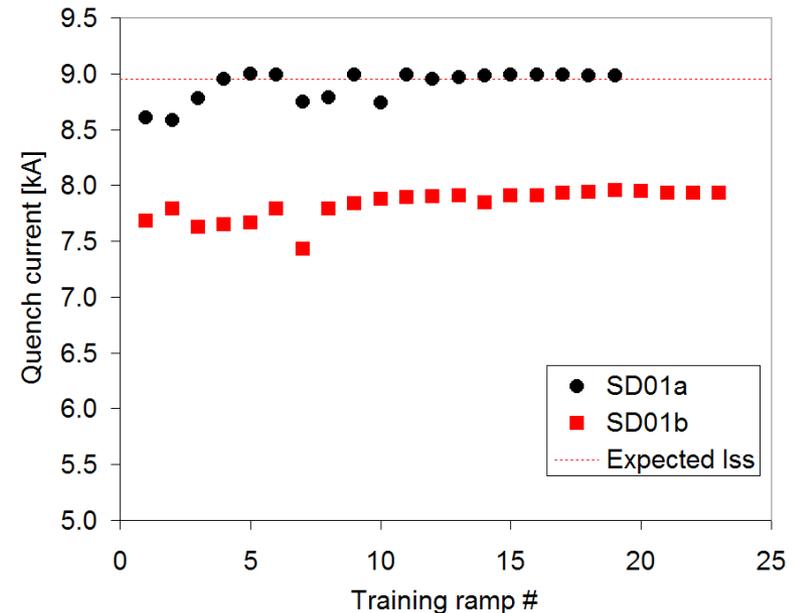
- The critical current is measured in few different conditions of temperature and field. By fitting the data with known **parameterizations** (see Unit 3), the entire critical surface can be reconstructed.
- If the magnet reaches the maximum current computed through the intersection of the measured critical surface and the load line, i.e.  $I_{max} = I_{ss}$ , one can declare victory (at least from the quench performance point of view).
- If the magnet maximum current  $I_{max}$  is lower than  $I_{ss}$ , the quench performance is expressed in term of **fraction of short sample** ( $I_{max}/I_{ss}$ ).
- A conductor-limited quench or a plateau at a level lower than the expected short sample is an indication of **conductor degradation**



# 4. Conductor limited quenches Degradation



- **Degradation** in a superconductor can be due to
  - Conductor damage or error in cable manufacturing
  - Stress
- ....but, sometimes, magnet performance can be interpreted as degraded in case of
  - **Errors** in the evaluation of maximum current
    - Computation of peak field
    - Measurements of temperature
    - Difference between coil and witness sample reaction temperature and time.
      - During the reaction of Nb<sub>3</sub>Sn coils, the temperature in the oven is not uniform.
- In general, there is about **5% uncertainty** in the short sample current estimate.





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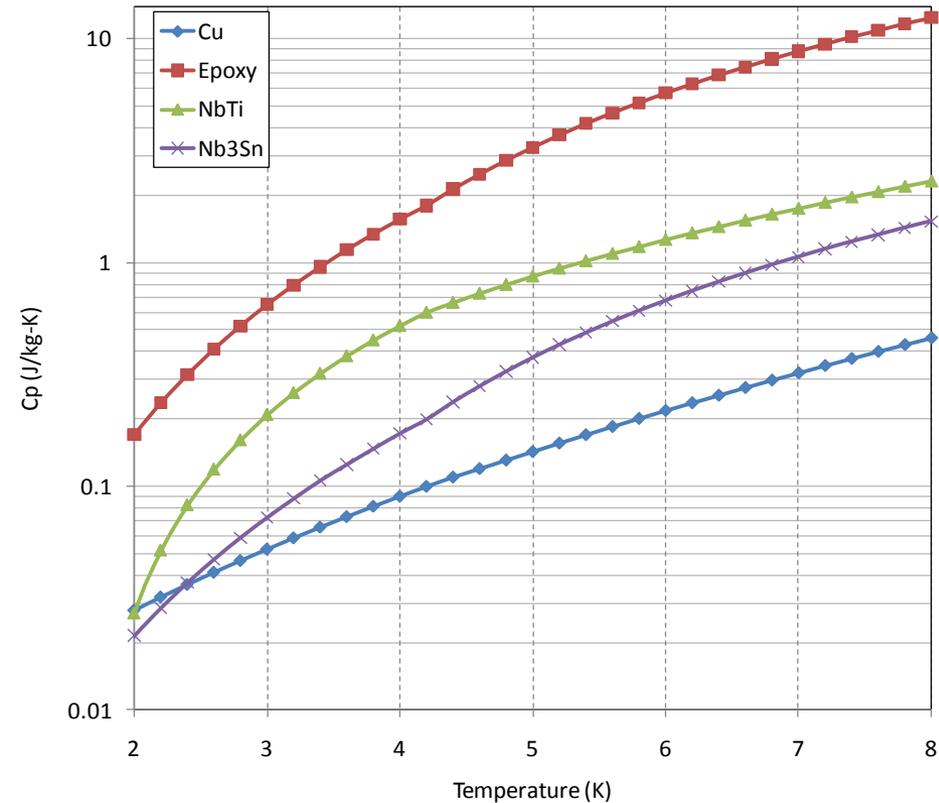


# 5. Energy deposited quenches

## Distributed disturbances



- Let's consider a **distributed disturbance**, that is a release of energy uniformly distributed. This situation corresponds to an **adiabatic condition**.
  - The temperature increase is uniform and no heat is conducted along the coil.
- In this conditions, the temperature rise depends only on the **specific heat  $C_p$**  (J/kg-K)
- Energy at low temperature are small since  $C_p$  are about  $10^{-3}$  of the room temperature values
  - For copper,  $C_p$  from  $28e-3$  J/kg-K at 1.8 K to 380 J/kg-K at 300 K





## 5. Energy deposited quenches Distributed disturbances

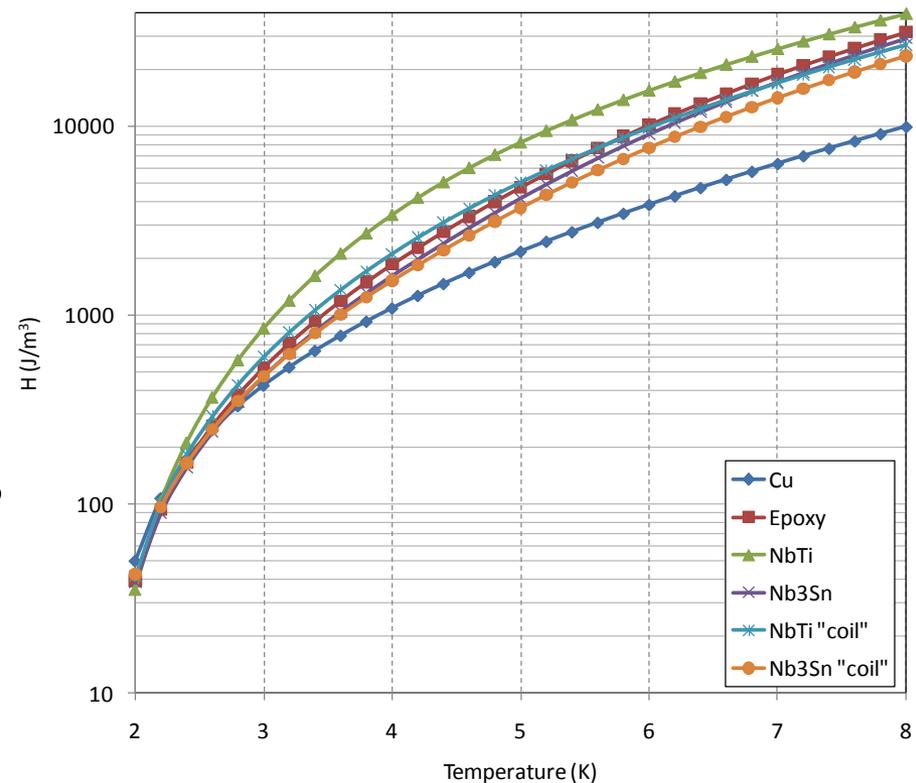


- Because of the change in  $C_p$  with temperature, for the computation of the energy density required to quench a coil, it is convenient to use the **volumetric specific enthalpy  $H$**  ( $\text{J}/\text{m}^3$ )

$$H(\theta) = \gamma \int_{1.8}^{\theta} C(\theta) d\theta$$

where  $\gamma$  ( $\text{kg}/\text{m}^3$ ) is the density of the material.

- Assuming a volume fraction of 1/3 of superconductor, 1/3 of copper, and 1/3 of epoxy, we can estimate a Nb-Ti and Nb<sub>3</sub>Sn coil volumetric specific enthalpy.
- Then, it is possible to compute the energy density to quench, but we need the **temperature margins**.

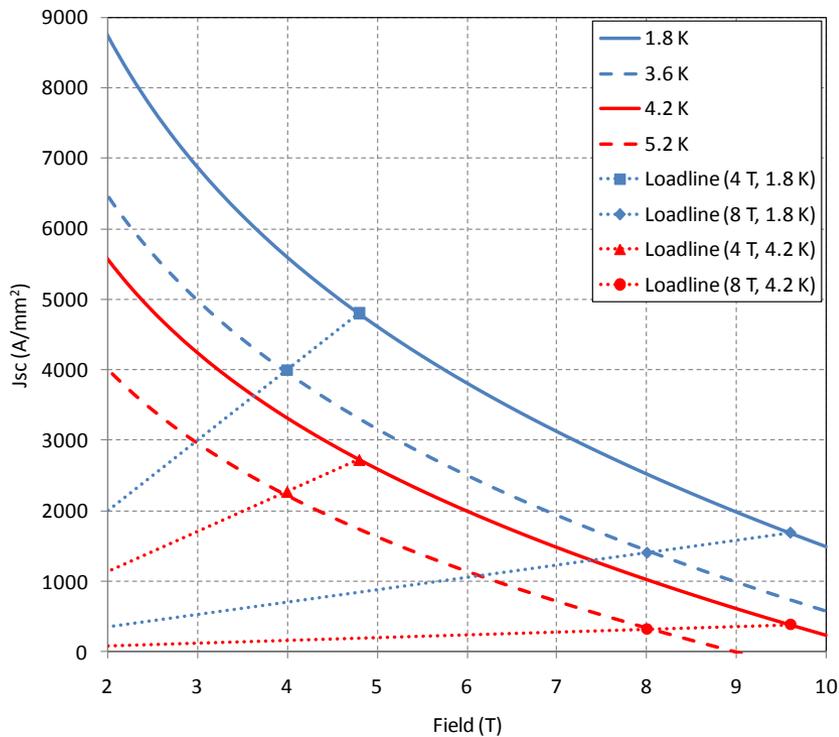




# Temperature margins at 80% of $I_{ss}$

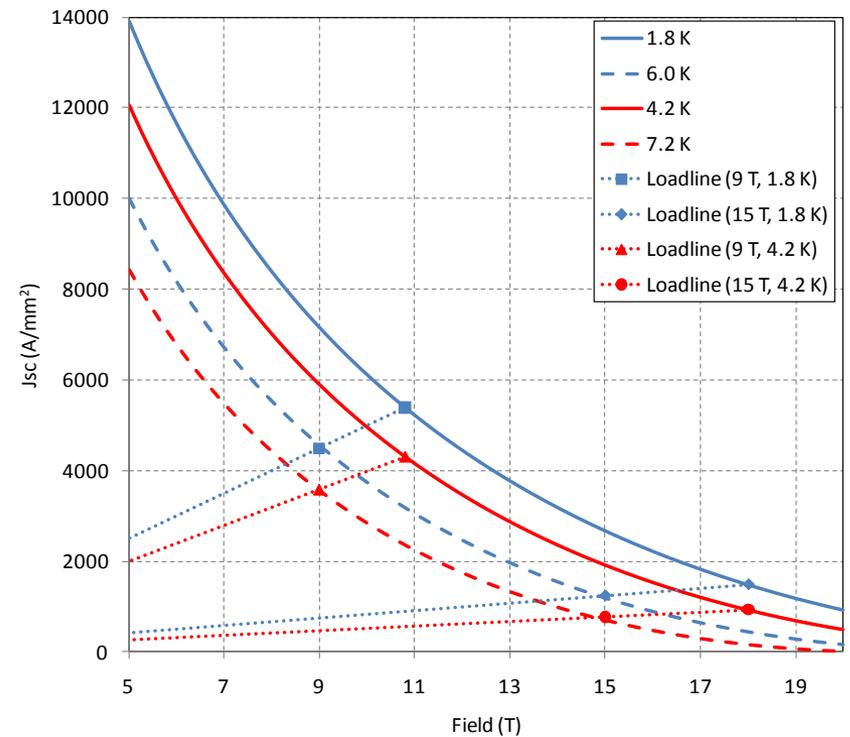
## NbTi

- At 1.8 K
  - 1.8 K of margin
- At 4.2 K
  - 1.0 K of margin



## Nb<sub>3</sub>Sn

- At 1.8 K
  - 4.2 K of margin
- At 4.2 K
  - 3.0 K of margin

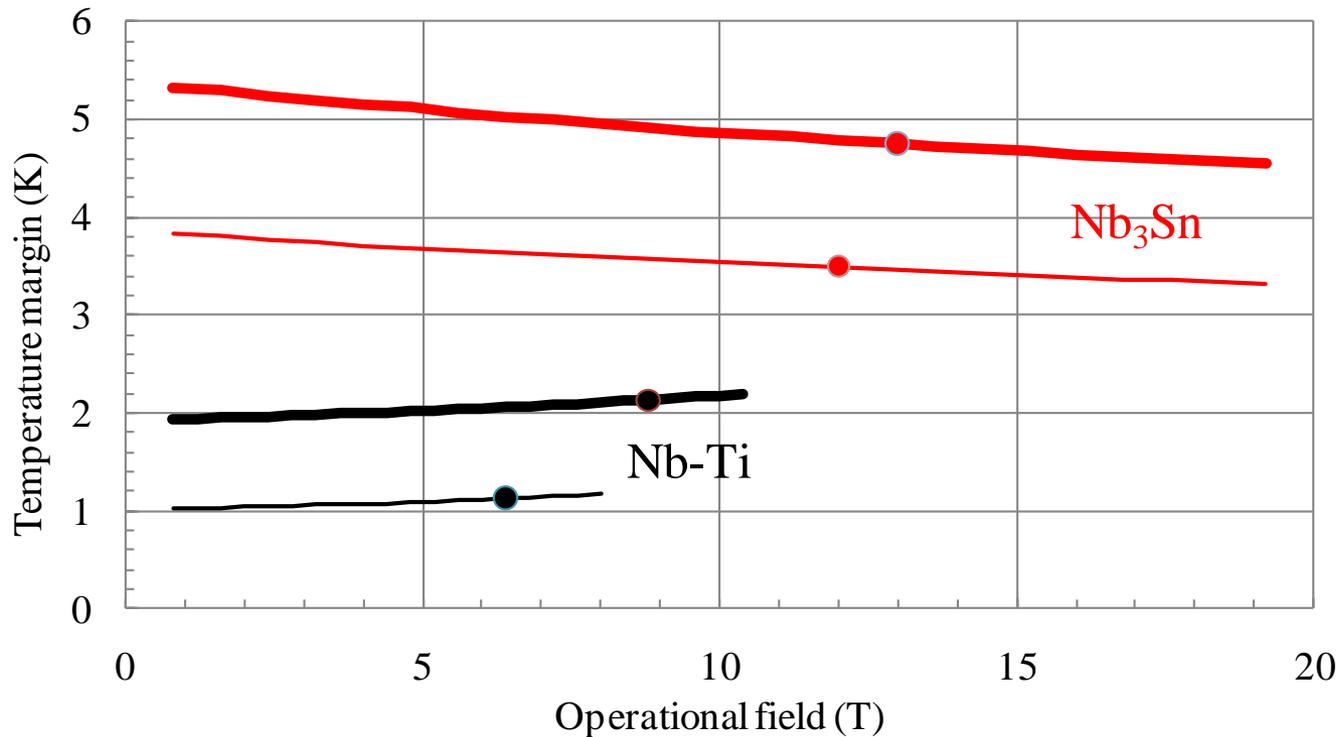




# Temperature margins at 80% of $I_{ss}$



- Margin roughly independent on operational field



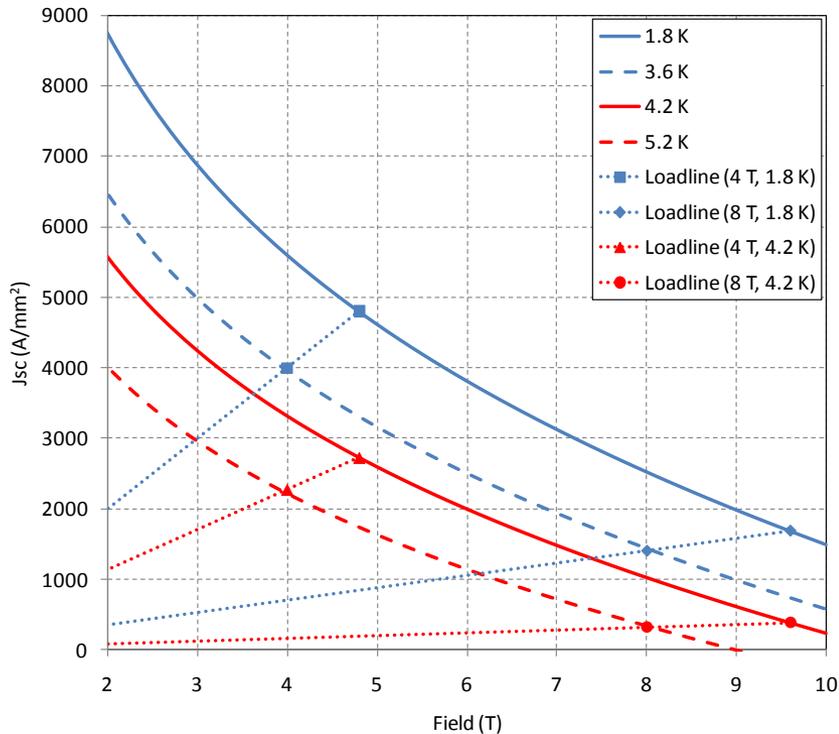


# Quench energy densities (Impregnated Nb-Ti and Nb<sub>3</sub>Sn coils)



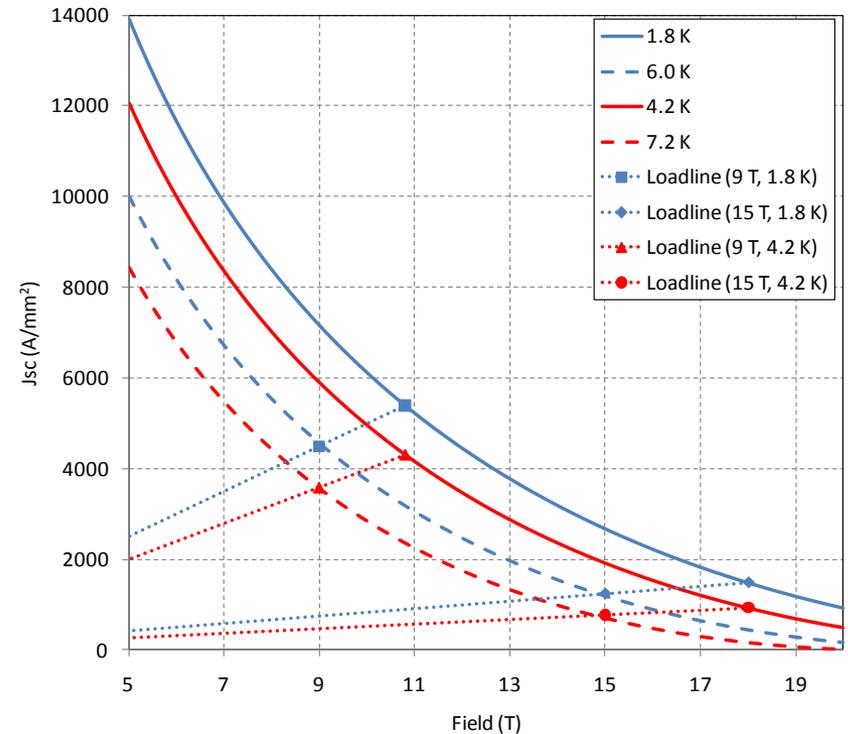
## ● NbTi

- At 1.8 K
  - 1.8 K of margin:  $1.4 \cdot 10^3 \text{ J/m}^3$
- At 4.2 K
  - 1.0 K of margin:  $3.3 \cdot 10^3 \text{ J/m}^3$



## ● Nb<sub>3</sub>Sn

- At 1.8 K
  - 4.2 K of margin:  $7.7 \cdot 10^3 \text{ J/m}^3$
- At 4.2 K
  - 3.0 K of margin:  $13.9 \cdot 10^3 \text{ J/m}^3$





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# 5. Energy deposited quenches

## Point disturbances



- Let's imagine that, in a coil operating in the superconducting state, a certain amount of **energy  $E$  is released**. This energy will bring a volume  $V$  of the superconductor to a temperature  $T \geq T_c$ .
- One can imagine that if  $E$  or  $V$  are not large enough, the temperature of the superconductor will decrease to below the critical temperature, because of cooling and or thermal conductivity.
- If, on the other hand,  $E$  or  $V$  are large enough, the normal zone will increase and a quench propagate.
- We therefore define
  - as **minimum quench energy  $MQE$** , the minimum energy necessary to initiate a quench
  - as **minimum propagation zone  $MPZ$** , the minimum volume of superconductor that must be brought beyond the critical temperature in order to initiate a quench.
- The two parameters are connected:  $MQE$  is the energy necessary to create  $MPZ$ , i.e [5]

$$E_{MQE} = V_{MPZ} \int_{\theta_0}^{\theta_c} \gamma C(\theta) d\theta$$

where  $C$  is the specific heat [ $\text{J kg}^{-3} \text{K}^{-1}$ ] and  $\gamma$  is the density [ $\text{kg m}^{-3}$ ].



# 5. Energy deposited quenches Point disturbances

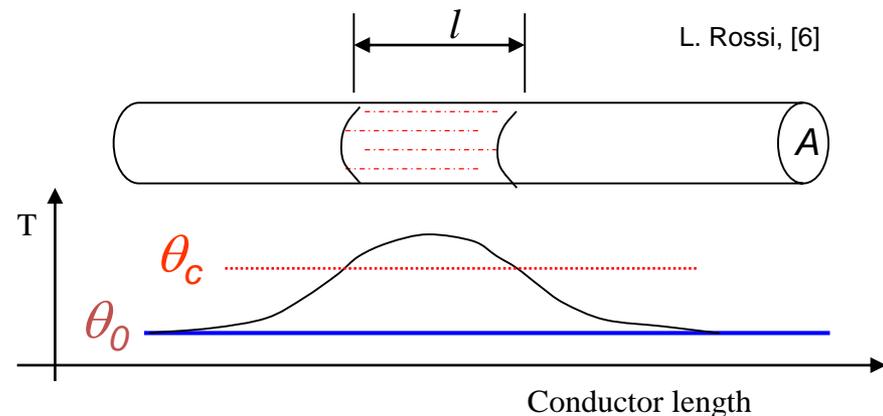


- We start considering a **wire** made purely of superconductor.
- Let's assume that a certain amount of energy  $E$  increased the temperature of the superconductor beyond  $\theta_c$  over a length  $l$ . The segment  $l$  of superconductor is **dissipating power** given by  $J_c^2 \rho A l$  [W].
- Part (or all) of the heat is conducted out of the segment because of the thermal gradient, which can be approximated as  $(\theta_c - \theta_0)/l$ . Therefore, being  $k$  the conductivity [ $\text{W m}^{-1} \text{K}^{-1}$ ], when the power dissipated equals the power conducted away

$$\frac{2kA(\theta_c - \theta_0)}{l} = J_c^2 \rho A l$$

which results in

$$l = \sqrt{\frac{2k(\theta_c - \theta_0)}{J_c^2 \rho}}$$





# 5. Energy deposited quenches

## Point disturbances



- The **length  $l$**  defines the MPZ (and MQE).
  - A normal zone longer than  $l$  will keep growing (quench). A normal zone shorter than  $l$  will collapse.
- An example [2]
  - A typical Nb-Ti 6 T magnet has the following properties
    - $J_c = 2 \times 10^9 \text{ A m}^{-2}$
    - $\rho = 6.5 \times 10^{-7} \Omega \text{ m}$
    - $k = 0.1 \text{ W m}^{-1} \text{ K}^{-1}$
    - $\theta_c = 6.5 \text{ K}$
    - $\theta_0 = 4.2 \text{ K}$
  - In this case,  $l = 0.5 \mu\text{m}$  and, assuming a 0.3 mm diameter, the required energy to bring to  $\theta_c$  is  $10^{-9} \text{ J}$ .
- A wire made purely of superconductor, without any stabilizer (like copper) around, **would quench with nJ of energy**.
  - In order to increase  $l$ , since we do not want to reduce  $J_c$ , we have to increase  $k/\rho$ : we need a composite conductor!

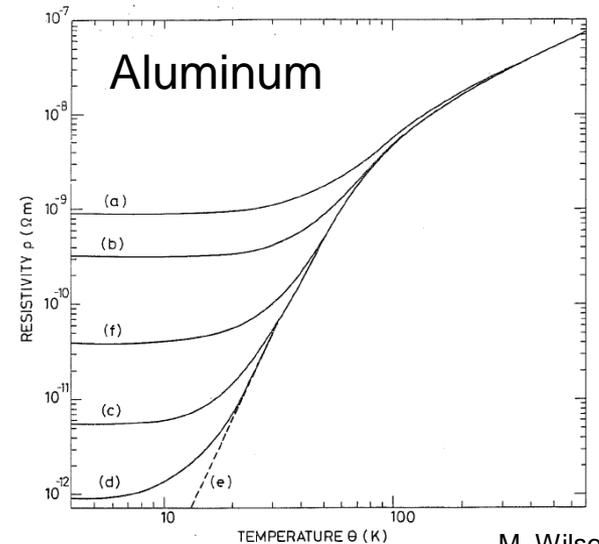
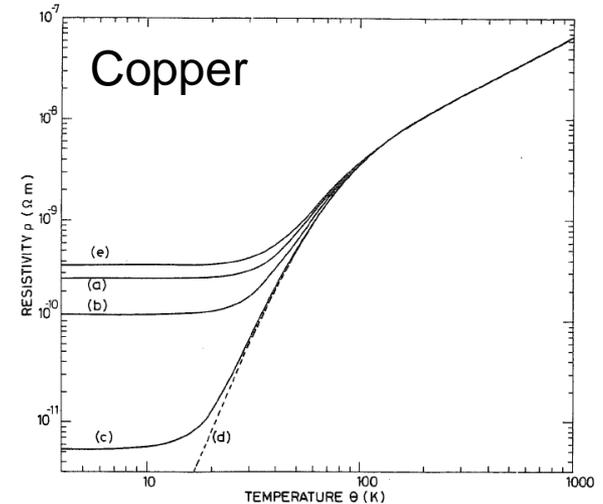
$$l = \sqrt{\frac{2k(\theta_c - \theta_0)}{J_c^2 \rho}}$$



# 5. Energy deposited quenches

## Point disturbances

- We now consider the situation where the superconductor is surrounded by material with **low resistivity and high conductivity**.
- Copper can have at 4.2 K
  - resistivity  $\rho = 3 \times 10^{-10} \Omega \text{ m}$  (instead of  $6.5 \times 10^{-7} \Omega \text{ m}$  for NbTi)
  - $k = 350 \text{ W m}^{-1} \text{ K}^{-1}$  (instead of  $0.1 \text{ W m}^{-1} \text{ K}^{-1}$  for NbTi).
- We can therefore **increase  $k/\rho$**  by almost a factor  $10^7$ .
- A significant improvement was achieved in the early years of superconducting magnet development after the introduction of composite conductor
  - Both for flux jump and stability viewpoint



M. Wilson, [2]

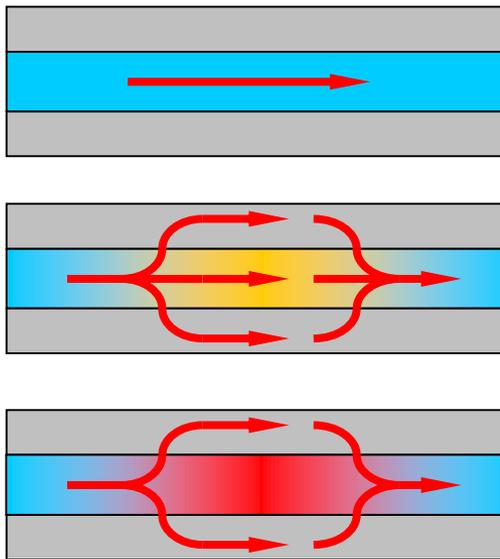


# 5. Energy deposited quenches

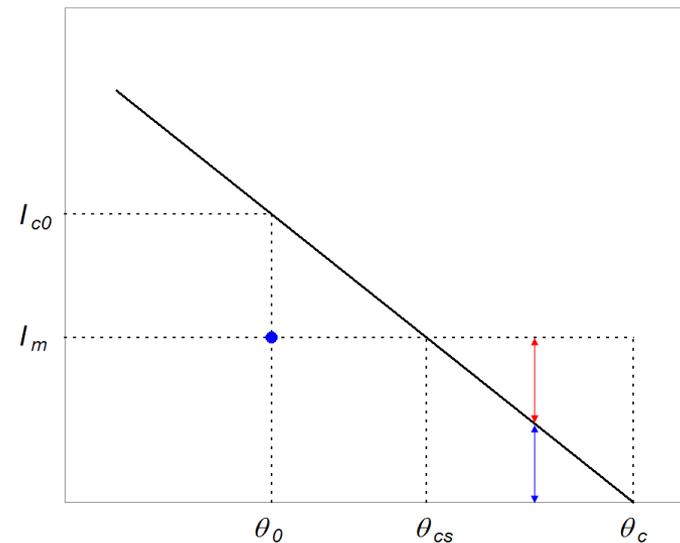
## Point disturbances



- In a composite superconductor, when a transition from normal to superconducting state occurs, the heat dissipated per unit volume of the entire wire (stabilizer and supercond.)  $G$  [ $\text{W m}^{-3}$ ], can be subdivided in three parts
  - **All** the current flows in the **superconductor**
  - The **current is shared** by the superconductor and the stabilizer
  - **All** the current flows in the **stabilizer**.



by L. Bottura





# 5. Energy deposited quenches

## Point disturbances



- The heat dissipated per unit volume of the entire wire (stabilizer and supercond.)  $G$  [ $\text{W m}^{-3}$ ] has a linear increase from  $\theta_{cs}$  and  $\theta_c$ .

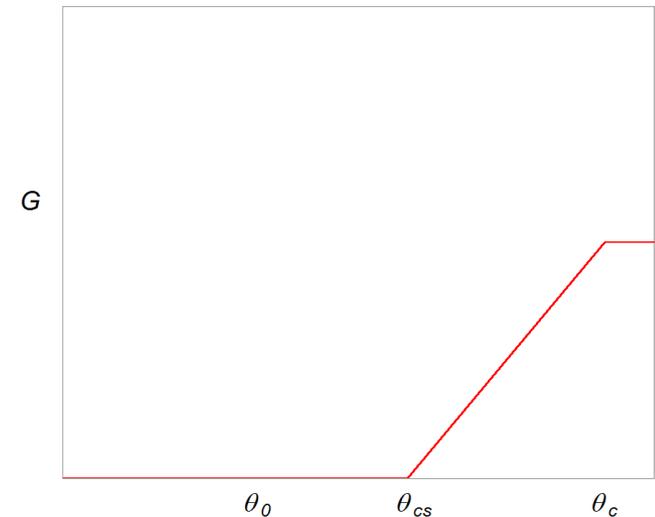
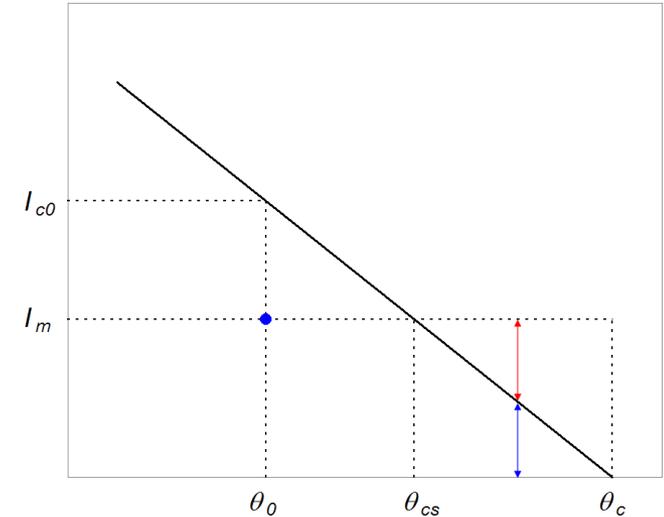
$$\theta < \theta_{cs} \longrightarrow G = 0$$

$$\theta > \theta_c \longrightarrow G = G_c = \rho_{stab} \frac{\lambda^2 J_m^2}{1 - \lambda}$$

$$\theta_{cs} < \theta < \theta_c \longrightarrow G = G_c \frac{(\theta - \theta_{cs})}{(\theta_c - \theta_{cs})} = \rho_{stab} \frac{\lambda^2 J_m^2}{1 - \lambda} \frac{(\theta - \theta_{cs})}{(\theta_c - \theta_{cs})}$$

$$\lambda = \frac{A_{sc}}{A_{tot}}$$

$$J_m = \frac{I_m}{A_{sc}}$$





## 5. Energy deposited quenches Point disturbances



- Once we provided an expression for  $G(\theta)$ , we modify the one-dimensional equation considered for the pure superconductor case, with a **three-dimensional equation** which includes the transverse conductivity.
- Assuming the coil as an isotropic continuum with two-direction conduction, the steady state equation of heat conduction becomes

$$\frac{1}{r} \frac{\partial}{\partial r} \left( r k_r \frac{\partial \theta}{\partial r} \right) + \frac{\partial}{\partial z} \left( k_z \frac{\partial \theta}{\partial z} \right) + \lambda_w G(\theta) = 0$$

where

- $k_r$  is the conductivity along the wire
- $k_z$  is the conductivity transverse to the wire
- $\lambda_w$  is the fraction in volume of the composite conductor (both superconductor and stabilizer) over the coil.



## 5. Energy deposited quenches Point disturbances

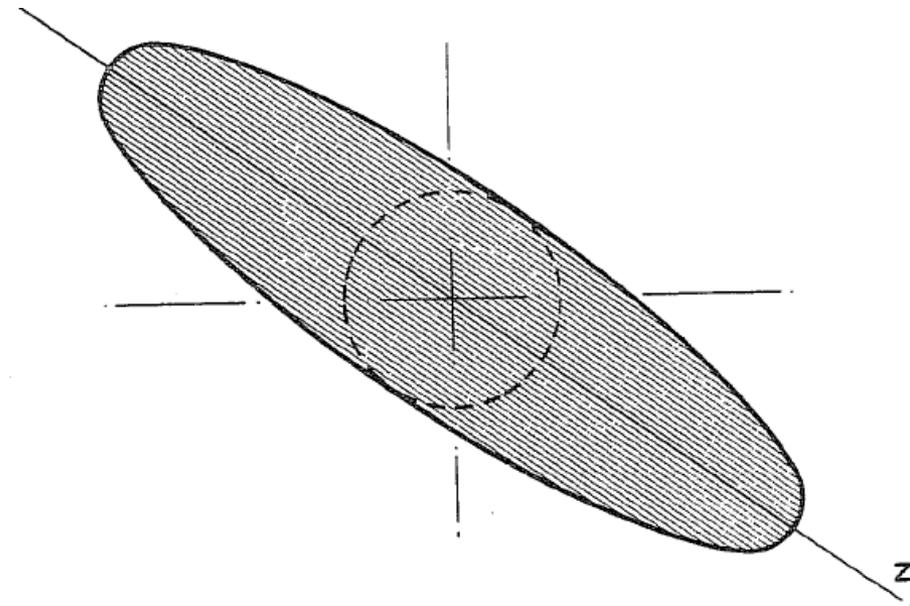
- The solution provides as MPZ an **ellipsoid elongated in the z direction** (along the cable), with a semi-axis  $R_g$  in  $z$

$$R_g = \pi \sqrt{\frac{k_z (\theta_c - \theta_{cs})}{\lambda_w G_c}}$$

and a circular cross-section in the transverse plane with radius  $r_g$

$$r_g = R_g \sqrt{\frac{k_r}{k_z}}$$

- The MQE necessary to MPZ is significantly increased, from the nJ level to the **10-100  $\mu$ J level**.



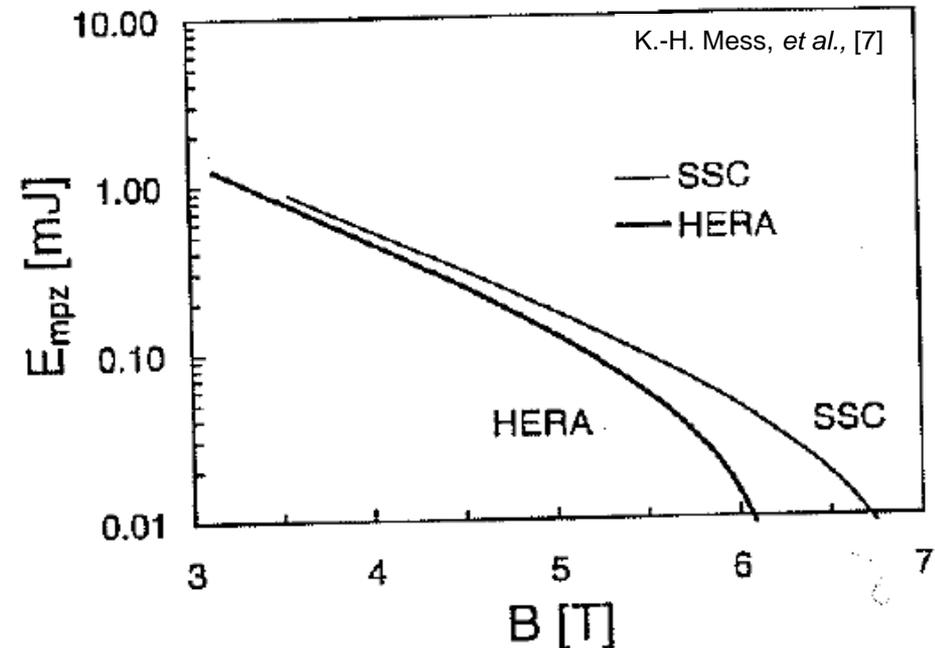
M. Wilson, [2]



## 6. Practical example I SSC and HERA wires [7]



- SSC
  - The design field ( $> 6.5$  T) is close to the short sample field
  - MQE =  $10 \mu\text{J}$
  - MPZ is less than strand diameter
    - “Adiabatic” condition (no heat exchanged with liquid helium) is a good estimate
- HERA
  - The design field (5 T) is less close to the short sample field
  - MQE =  $150 \mu\text{J}$
  - MPZ is several mm
    - Helium cooling will increase the energy threshold.





# Outline

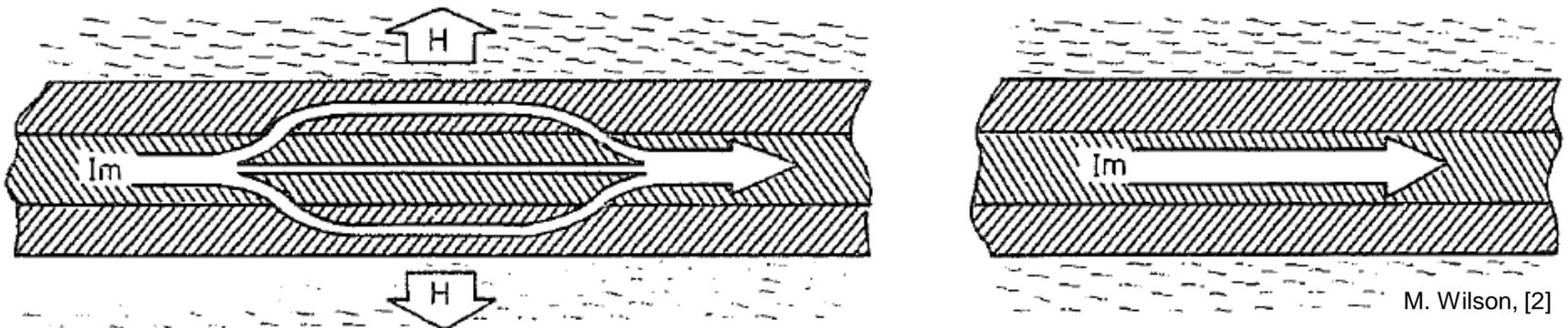


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- 7. Cryogenic stabilization**
8. Practical example: LHC wire
9. Conclusions



## 7. Cryogenic stabilization

- When a ***distributed disturbance*** (large volume) produces a transition in the superconductor in the normal zone, or when the **MPZ is significantly larger** than the strand diameter, we abandon the adiabatic conditions described in Sec. 5.
- In this case, we have to take into account the **cryogenic liquid** (liquid Helium) in contact with the wire.





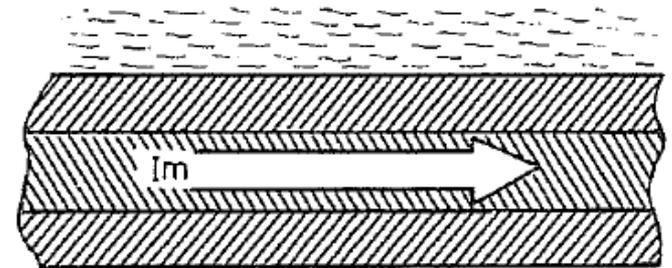
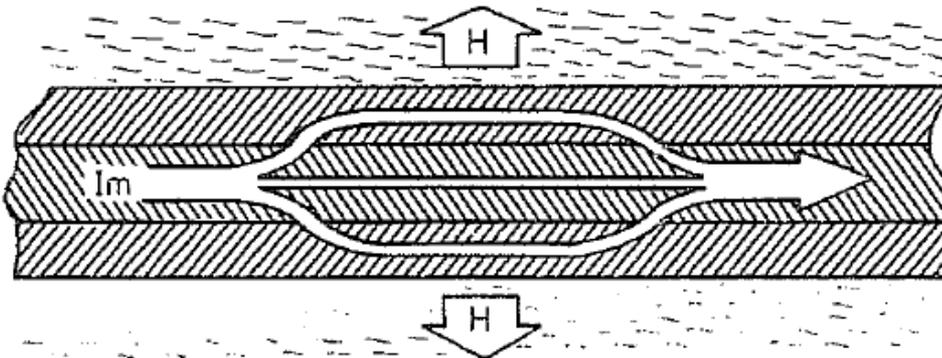
## 7. Cryogenic stabilization



- The **heat generation per unit of cooled area** is given by

$$G(\theta) = \frac{\lambda^2 J_c \rho}{(1-\lambda)} \frac{(\theta - \theta_0)}{(\theta_c - \theta_0)} \frac{A}{P} = G_c \frac{(\theta - \theta_0)}{(\theta_c - \theta_0)} \frac{A}{P}$$

where we assumed that  $\theta_0 = \theta_{cs}$  (i.e.  $J_m = J_C$ ), and **A** is the cross-sectional area of the conductor and **P** is the cooled perimeter.



M. Wilson, [2]



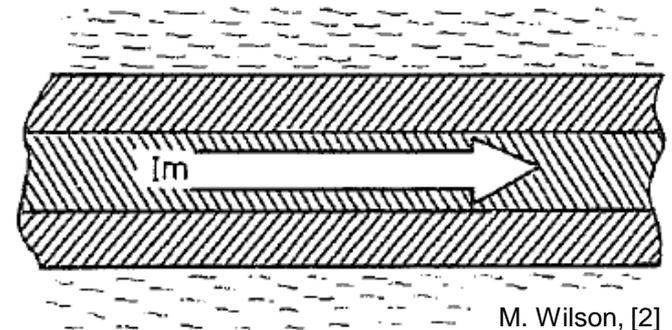
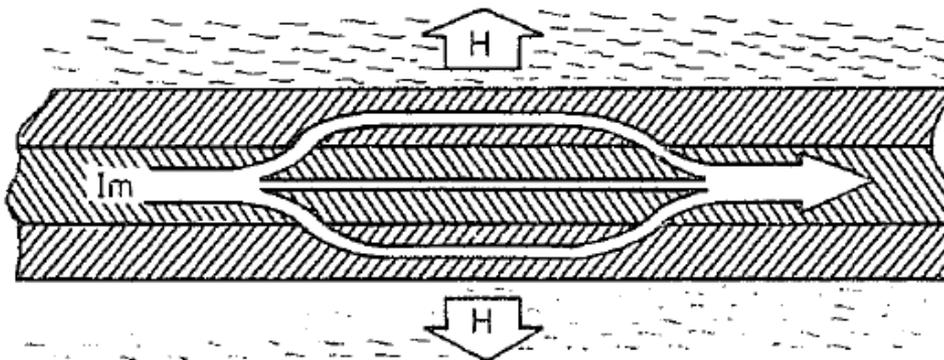
# 7. Cryogenic stabilization



- The **cooling** per unit of cooled area is given by

$$h(\theta - \theta_0)$$

where  $h$  [ $\text{W m}^{-1} \text{K}^{-1}$ ] is the heat transfer coefficient (we assume it constant) and  $\theta_0$  the temperature of the bath.



M. Wilson, [2]



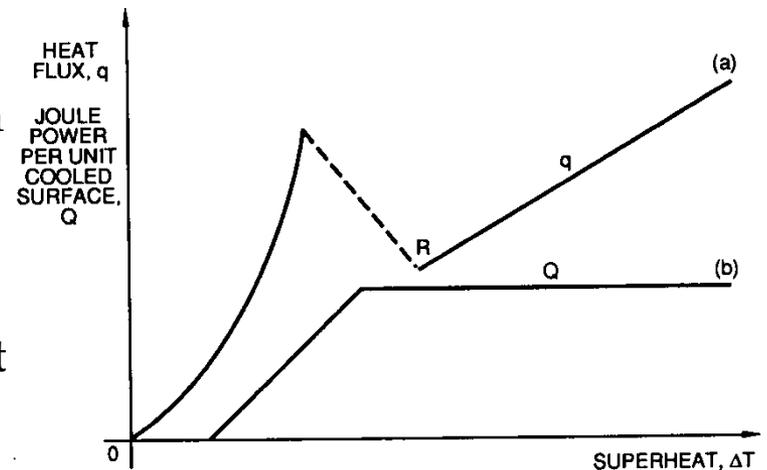
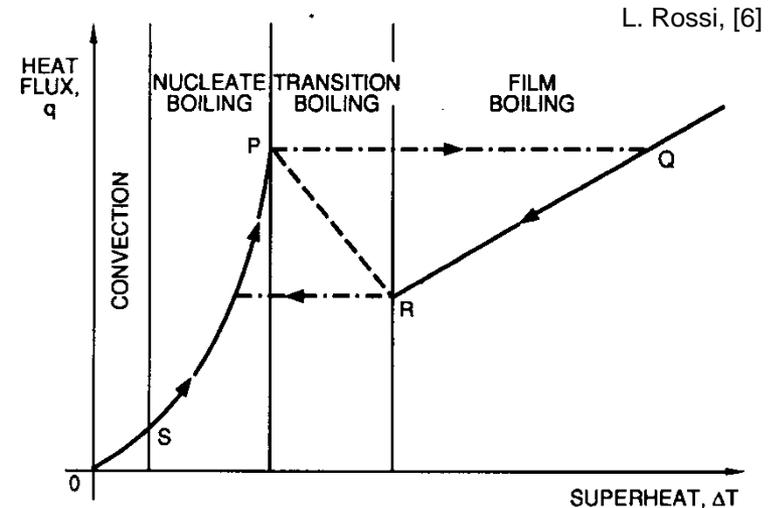
# 7. Cryogenic stabilization Stekly criteria



- The  **criterion for cryogenic stability** (Stekly and Zar) is therefore

$$\alpha = \frac{\lambda^2 J_c^2 \rho A}{(1-\lambda) Ph(\theta_c - \theta_0)} = \frac{G_c A}{Ph(\theta_c - \theta_0)} < 1$$

- In reality  $h$  is characterized by three zones
  - Nucleate boiling
  - Transition boiling
  - Film boiling
- In general, we have cryogenic stability when the curve of the dissipated power  $G(\theta)$  **stays below** the curve of the power removed by helium.
- Stekly criteria require an ohmic heat per unit area of cooled surface  $< 1.5\text{-}2.0 \text{ kW/m}^2$



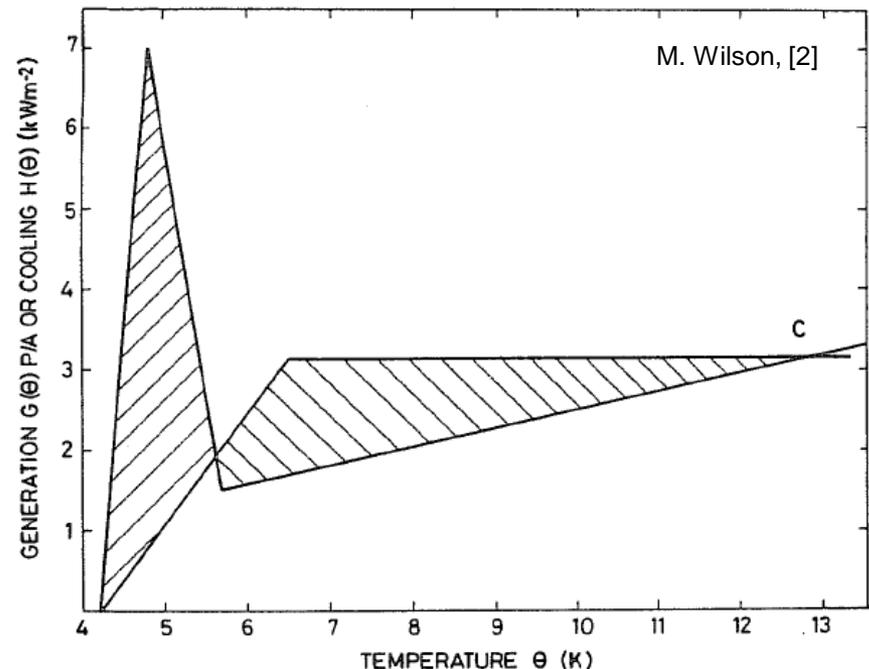


# 7. Cryogenic stabilization

## Equal-area theorem (Maddock criteria)

- The Stekly criteria neglects heat conduction.
- But, if the hot zone has finite dimensions and is surrounded by a cold zone, the **heat conduction** at the boundaries can improve the stability
  - Conductor can be stable also if  $G > H$
- Maddock, *et al.*, [10] proposed to **equal area theorem**
  - The surplus of heat generation must be balanced by a surplus of cooling.
- Maddock criteria requires an ohmic heat per unit area of cooled surface  $< 3 \text{ kW/m}^2$

$$\int_{\theta_0}^{\theta_1} \left\{ H(\theta) - \frac{A}{P} G(\theta) \right\} k(\theta) d\theta = 0$$





# Outline



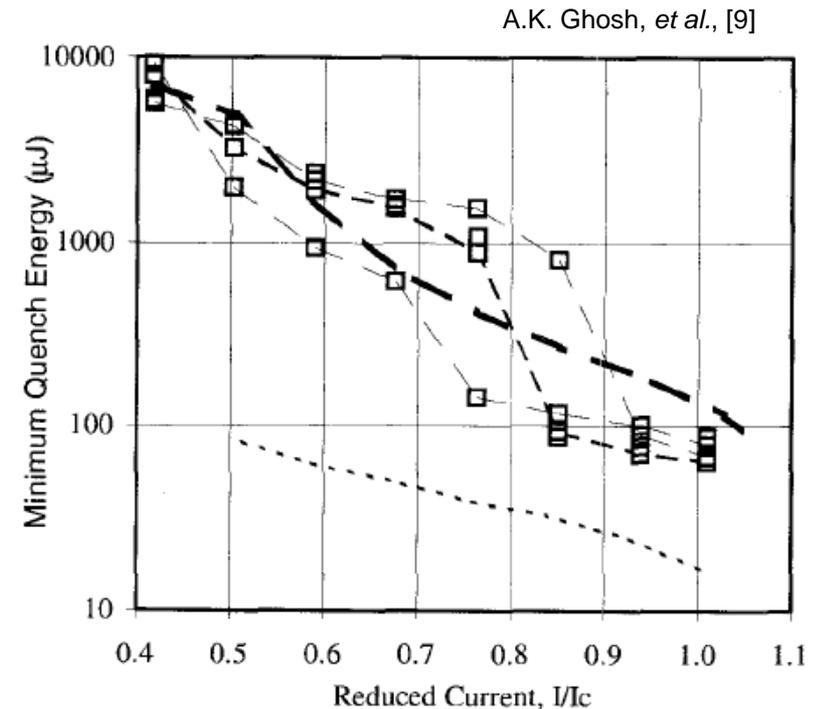
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# 8. Practical example II LHC wire [8]



- **Wire**
  - Measurements showed that the most important parameter is the amount of helium in contact with the strand.
  - The MQE at  $0.8 I_c$  is
    - $10 \mu\text{J}$  in adiabatic conditions
    - $1000 \mu\text{J}$  in open bath
- **Cable**
  - A “kink” is observed
    - Zone where the strands behave collectively
    - Zone where the strands behave individually
  - Above the kink
    - MQE is around  $100 \mu\text{J}$
  - Below the kink
    - MQE is  $> 1000 \mu\text{J}$





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## 9. Conclusions



- We defined and classified **quenches** in superconducting magnets
  - Conductor limited quenches
  - Energy deposited quenches
- We introduced the concept of **MPV and MQE**
  - Depending on the assumptions, the energy required to quench a magnet is
    - $10^{-9}$  J in the case of a “pure superconductor”
    - 10-100  $\mu$ J in the case of a composite wire
    - More than 1000  $\mu$ J in open bath
- In the case of distributed disturbance, the minimum quench energy density can be of about  $10^3$  J m<sup>-3</sup> (Nb-Ti magnets).
- Considering the effect of liquid He, a conductor can be stable up to an ohmic heat per unit area of cooled surface of about 3 kW/m<sup>2</sup>.



# Appendix I





## 4. Conductor limited quenches

### Measurements of the conductor critical current

- Both the increase in  $\rho_{sc}$  and  $V$  can be fitted by a scaling law like

$$\frac{\rho_{sc}}{\rho_c} = \left( \frac{I}{I_C} \right)^{N-1} \quad \frac{V}{V_c} = \left( \frac{I}{I_C} \right)^N$$

- $N$  is called the resistivity transition index, or  $N$ -value, and it is an indication of the sharpness of the transition from superconducting to normal state (see Unit 6).
- Usually, the better is the conductor, the higher ( $\geq 30$ ) is the  $N$ -value.

