Unit 12
Protection of Superconducting Accelerator Magnets – Episode I

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European Organization for Nuclear Research (CERN)

Soren Prestemon   (With many thanks to **Helene Felice)**

Lawrence Berkeley National Laboratory (LBNL)

- What is quench protection?

- What are the critical parameters?

- How does the quench propagate?

- Protection Methods

- Electrical integrity

- Detection

Why do we need to protect the magnets ?

- A superconducting accelerator magnet has a **large magnetic stored energy** $\frac{1}{2}LI^2$

- A quench produces a resistive zone
- Current is flowing through the magnet

  Joule Heating Voltages (R and L)

- The challenge of the protection is to provide a **safe conversion** of the magnetic energy to heat in order to minimize
  - Peak temperature ("hot spot") and temperature gradients in the magnet
  - Peak voltages

- The final goal being to **avoid any magnet degradation**
  - High temperature => damage to the insulation or stabilizer
  - Large temperature gradient => damage to the conductor due to differential thermal expansion of materials

Quench

↓

Normal zone growth

↓

Detection

↓ ↓

Power supply switched off

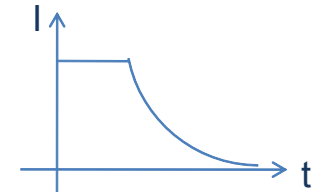Trigger protection options

Protection heaters | Extraction | Quenchback

Current decay in the magnet

Magnetic energy $\frac{1}{2}LI^2$

Converted to heat by Joule heating

$\int_0^\tau R(t)I(t)^2\, dt$

The faster this chain happens the safer is the magnet

Current decay/Hot spot correlation: the MIITs

- **Basic energy conservation**
  - Energy deposited by Joule heating inducing an increase in temperature based on the specific heat of the materials
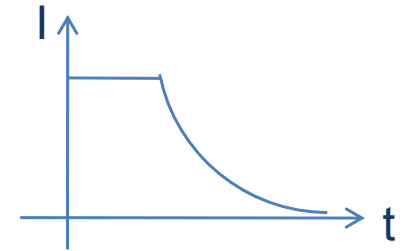
  $$\gamma C_p(T)dT = \rho J(t)^2 dt$$

  - With the assumptions that:
    - Joule heating only produced by stabilizer (Cu)
    - Heat capacity of conductor and insulation
    - Adiabatic condition: no longitudinal heat transfer

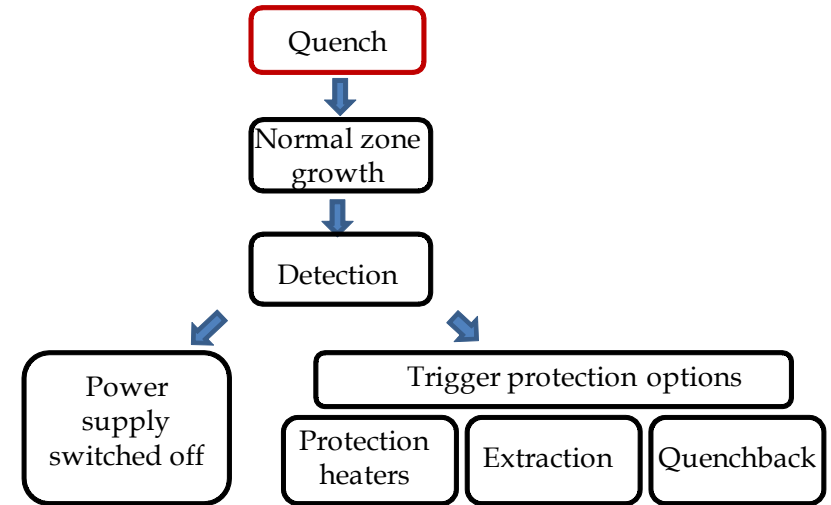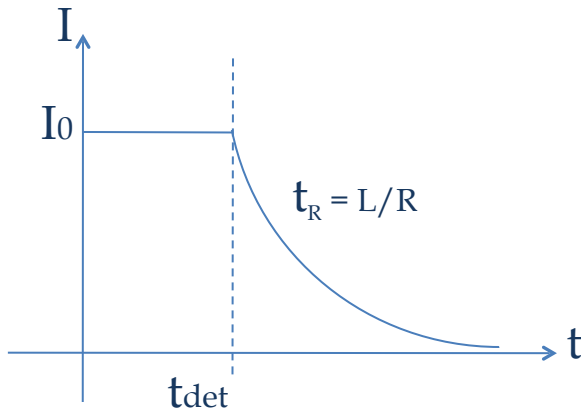 $$A_{tot} \sum_k \gamma_k v_k C_{p,k}(T)dT = \frac{\rho_{Cu}(T,B)}{A_{Cu}} I(t)^2 dt$$

  $$10^6 MIITs = \int_0^\infty I(t)^2 dt = A_{tot} A_{Cu} \int_{T_0}^{T_{max}} \frac{\sum_k \gamma_k v_k C_{p,k}}{\rho_{Cu}(T,B)} dT$$

By evaluating the MIITs from the current decay, we obtain an estimate of the hot spot temperature T_{max}

What are the critical parameters?

- **Typical scenario**
  - The current in the magnet is I_0, a quench occurs
  - Normal zone propagation and voltage build-up
  - Detection system based on some voltage threshold
  - Magnet current starts to decay with a time constant tau=L/R where
    - L is the total inductance of the system
    - R is the total resistance in the system: normal zone + external

Quench propagation velocity

- Once the normal zone has started to grow, it will continue to expand under the combined actions of Joule heating and heat conduction. The rate of this propagation is called quench propagation velocity.

- Using some approximation, it can be derived from the heat equation

$$\gamma C_p \frac{\partial T}{\partial t} = \nabla \cdot (k(T)\vec{\nabla}T) + Q$$

  - In case of an adiabatic coil (heat transfer only by conduction – impregnated coils)
  - assuming a constant propagation velocity
  - assuming a 1D propagation
  - assuming constant material properties wrt temperature

  The quench propagation velocity v_{long} can be estimated by:

$$v_{long} \approx v_{adia} \approx \frac{J_e}{\gamma C_{p,av}} \left(\frac{\rho_{Cu} k_{av}}{\Delta T} \right)^{1/2}$$

  - where $C_{p,av}$ and k_{av} include only copper
  - where $\Delta T = T_s - T_0$ with T_s average temperature between critical temperature and current sharing temperature

- The heat is also propagating from turn-to-turn: it is the transverse propagation
- But is it a much slower propagation, about 10 to 100 times slower
- Since the time scales are very different, including the transverse propagation does not affect the longitudinal propagation velocity estimate
- The same formula as for longitudinal propagation can be used to estimate the transverse propagation velocity.
  - Insulation plays a key role so the $C_{p,av}$ needs to include the insulation C_p
  - The thermal conductivity also needs to be changed

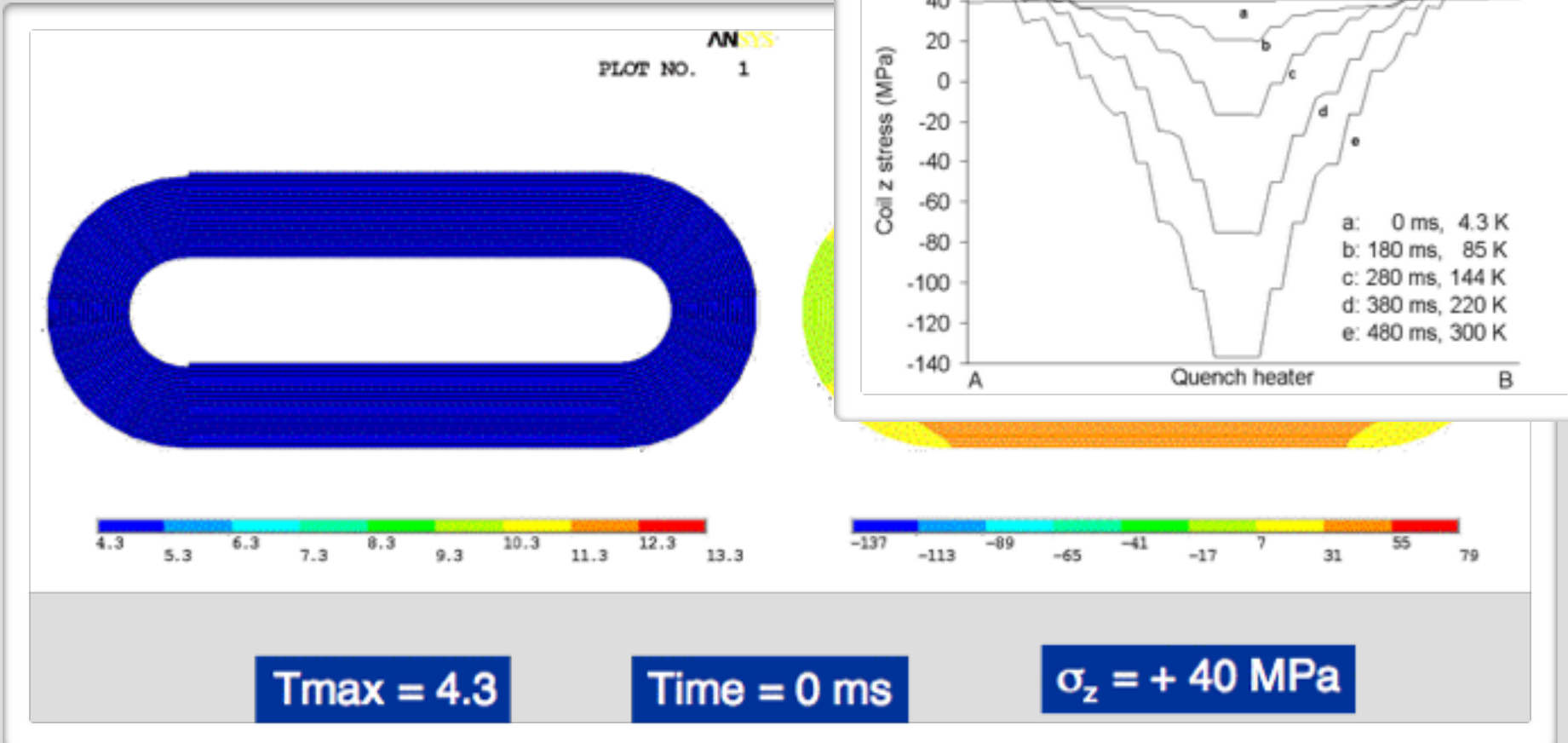$$v_{transv} \approx v_{long} \left(\frac{k_{transv}}{k_{long}}\right)^{1/2} \frac{C_{p,av}}{C_{p,all}}$$

5 ms
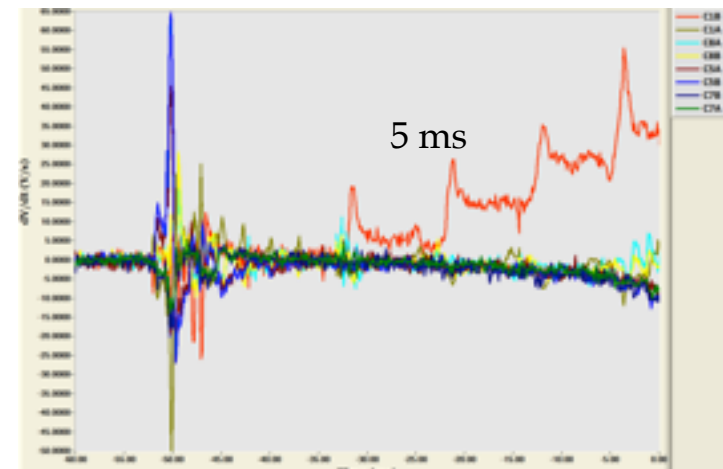
Courtesy of Juan Lizarazo, LBNL
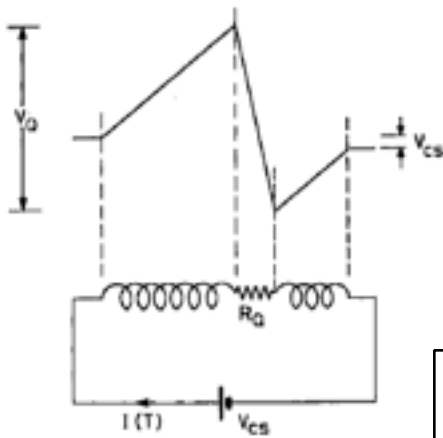
Quench propagation

Caspi et al, TAS 2003
Ferracin et al, TAS 2004

- In absence of external resistance in the circuit,
  - the voltage comes first from the power supply (Vcs).
  - After the detection of the quench the power supply is turned off and the current starts to decay.
  - Some internal voltages due to the magnet inductance develop in the magnet

$$V_Q(t) = R_Q(t)I(t) - M\frac{dI(t)}{dt}$$

$$L\frac{dI(t)}{dt} - R_Q(t)I(t) = 0$$

Mutual inductance between the normal zone and the rest of the coil

Self-inductance of the system

$$V_Q(t) = R_Q(t)I(t)(1 - \frac{M}{L})$$

M.K. Wilson, Superconducting Magnets

- As the normal zone propagates, the I(t) decays and M increases. The internal voltage will rise to a peak and decreases.
- A widely spread normal zone will mitigate the peak voltage

$$10^6 MIITs = \int_0^\infty I(t)^2 dt = A_{tot} A_{Cu} \int_{T_0}^{T_{max}} \frac{\sum_k \gamma_k v_k C_{p,k}}{\rho_{Cu}(T,B)} dT$$

Assuming all energy
deposited in the coil

Fast normal zone growth

τ = L/R gets smaller

Reduced Local Joule heating in dV

Stored energy better
distributed in the coil volume
Low hot spot temperature

Slow normal zone growth

Long time constant τ = L/R
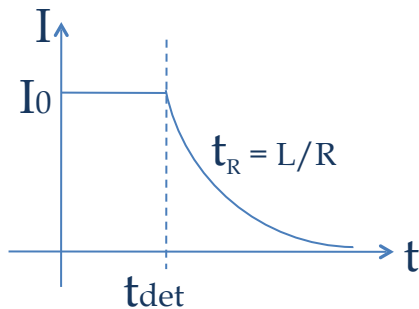
Increased Local Joule heating in dV

Increased hot spot temperature

How much resistance do we need?

- The normal zone has to grow fast to minimize the peak temperature and the peak voltage

- In rare cases the normal zone propagation is fast enough to provide enough resistance => the magnet is then self-protected

$$10^6 MIITs = \int_0^\infty I(t)^2 dt = \int_0^{t_{det}} I_0^2 \, dt + \int_0^\infty \left(I_0 e^{-\frac{R(t)}{L}t} \right)^2 dt$$

- By assuming that the resistance is constant over time, we can obtain a estimate of what is the required resistance

$$R = \frac{1}{2} \left[\frac{LI^2}{10^6 MIITS - I_0^2 t_{det}} \right]$$

- In reality, the normal zone resistance is function of time, temperature and magnetic field through the **magneto-resistance**

F.R. Fickett, Transverse magnetoresistance of oxygen free copper, IEEE Trans. On Magn., vol 24, No 2, March 1988

- The stabilizer plays a major role in the normal zone resistance:

  - If the stabilizer resistance is large, the current decays faster

  - Opposite requirements for protection and stability
    - Protection would tend to request low RRR to hasten the current decay
    - Dynamic stability asks for a high RRR
    - In high operating field magnets, the magneto-resistance will lower the RRR
   - High impact for high RRR
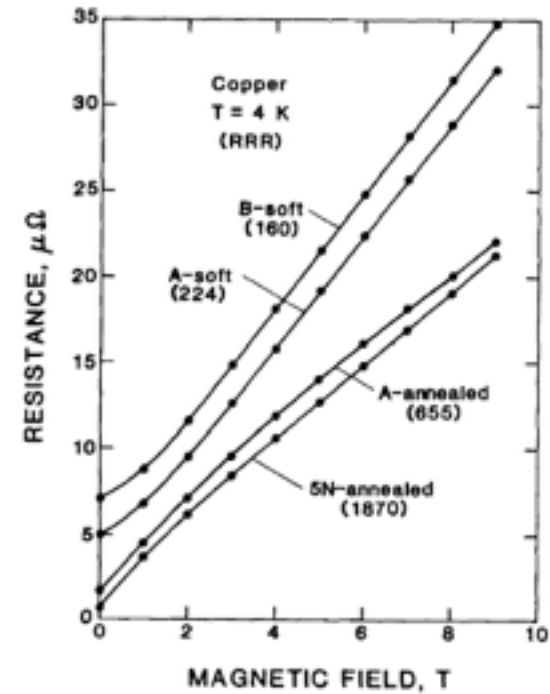
- In most accelerator magnets, the "natural" resistance growth is insufficient to provide a good protection => Need to enhance the resistance

- *Method 1:*
  - Add an external resistor

- *Method 2:*
  - Apply additional heat to the coils to force them to become more resistive
    - Protection heaters mounted on the coils
    - Optimized to minimize thermal diffusion time

- *Method 3:*
  - Use of couple secondary circuits
    - Can be external or internal to the coil
    - Energy dissipation in the secondary circuit through the voltage generated by the mutual inductance between the coil and the secondary
    - Can lead to quenchback in cases where secondary becomes a heat source

- *Method 4*:
  - Coil subdivision

- The system total resistance can be simply increased by the addition of an external resistor R_{ext} also called extraction or dump resistor.

- R_{ext} is connected to the magnet after detection => active switching

- Advantage: heat dissipated in R_{ext} does not affect the cryogenics

Cold

- Ideally we want $R_{ext} > R_n$ to force a fast current decay
  - but voltage develops across R_{ext} and across the magnet $V_{dump} = R_{ext} \cdot I_0$
  - Typical V_{dump} is below 1 kV
  - If V_{dump} is too large:
    - Difficulty to switch
    - The magnet electrical integrity can be compromised

M.K. Wilson, Superconducting Magnets

- R_{ext} is usually small with respect to R_n because most of the magnets operate at high current.
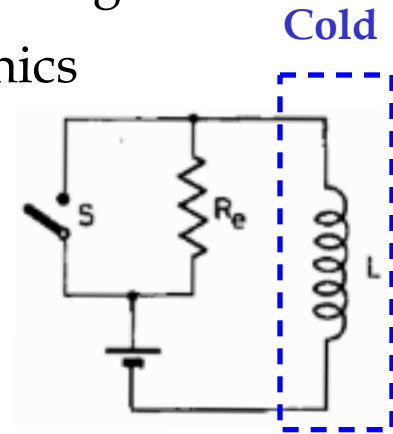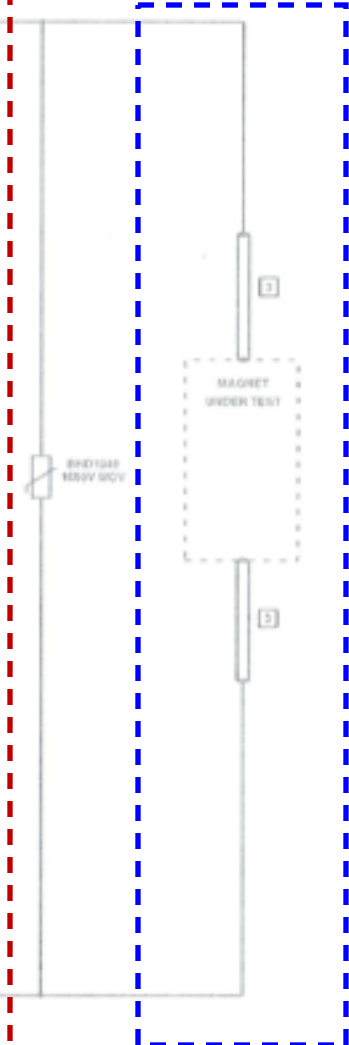
Example: LBNL extraction rack schematics

To power supply
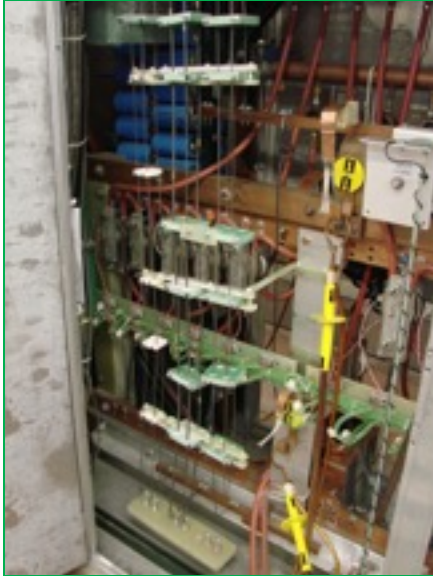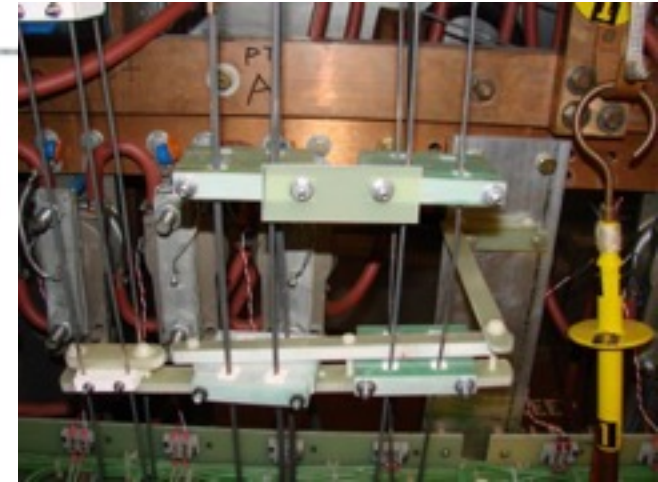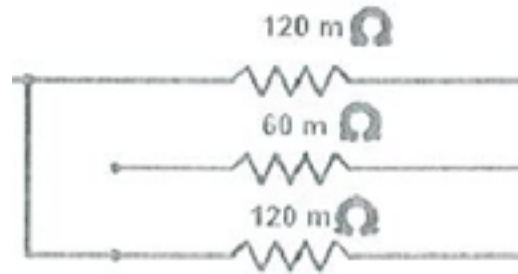
EXTRACTION RACK
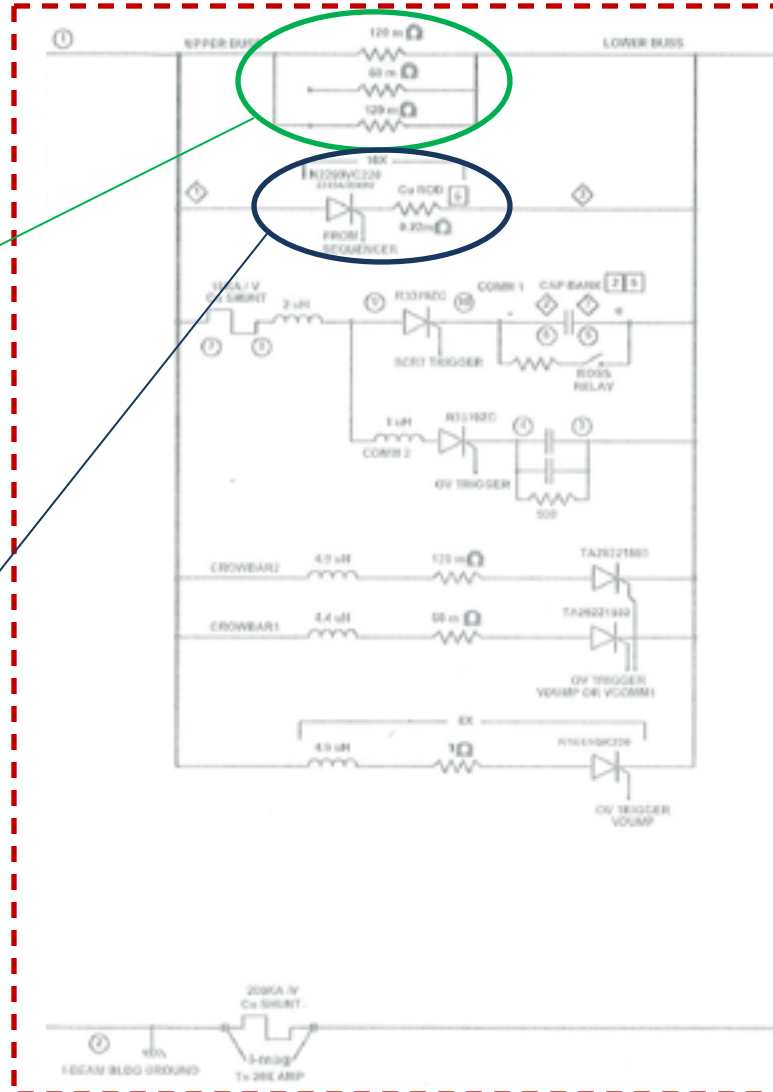
CRYOSTAT

- "homemade"
- 1/8" diameter steel wire (~ 5 mm)
- Range of 30 to 120 mΩ
- G10 spacer to prevent motion and contact due to magnetic forces
- Note that the value of the dump resistor increases during extraction due to elevation of the steel wire temperature

- Made of 304 stainless steel
- Range of 10 to 120 mΩ

- Design choices are detailed in:

Fermi National Accelerator Laboratory

TM-1611

Design Note of a 10,000 Amp 2 MJoules Dump Resistor for the Magnet Test Facility

A. T. Visser
Fermi National Accelerator Laboratory
P.O. Box 500
Batavia, Illinois 60510

March 1990

Courtesy of Guram Chlachidze, FNAL

- Another active protection scheme consists in adding heaters to the coil
- Advantages:
  - Design flexibility
  - Optimized placement on the coil to maximize efficiency
  - Fast response time

- Disadvantages:
  - Powering of the heaters put an additional electrical strain on the system
  - Balance needs to be found between optimized thermal diffusion time and dielectric strength
  - Attention must be paid to the electrical integrity of the "coil + protection heater" system

- Protection heaters (PH) are usually made of thin stainless steel strips
  - as thin as 25 microns

 - $$R_{PH} = \frac{\rho_{ss} L_{PH}}{A_{PH}}$$

- The heaters should provide
  - The coverage of several turns
  - The coverage of a full cable twist pitch
- The design should minimize the thermal diffusion time (factor limiting the heater efficiency)
- The heaters are powered by the discharge of a capacitor
  - The $R_{PH}C$ constant determines the time constant of the discharge τ_{PH}
  - The voltage across the capacitor determines the voltage across the heaters V_{PH}
  - R_{PH}, V_{PH} and τ_{PH} define the energy deposited in the heater

- The efficiency of the heaters depend on the heat diffusion time from the protection heater to the conductor
  - For a given operating current, the diffusion time is primary dictated by the **amount of insulation between heaters and coil**
    - The more insulation between the PH and the coil, the slower will be the diffusion
    - Trade-off between efficiency and electrical integrity
  - **Cooling condition** might also affect the heater efficiency
  - **The temperature margin** in the magnet
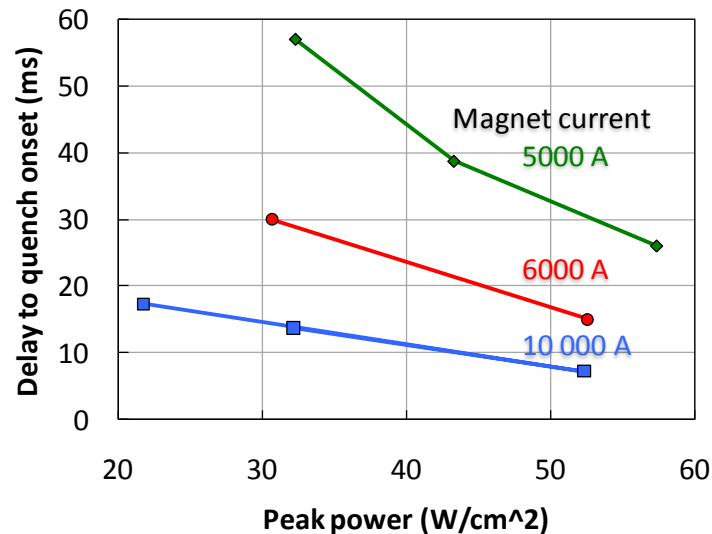
Courtesy of Tiina Salmi, LBNL

- A secondary circuit is couple to the magnet via the mutual inductance M

- Magnetic flux is shared between the magnet winding and the secondary. A voltage is generated in the secondary, allowing for energy dissipation.

$$L_p \frac{dI_p(t)}{dt} + R_Q(t)I_p(t) + M\frac{dI_s(t)}{dt} = 0$$

$$L_s \frac{dI_s(t)}{dt} + R_s(t)I_s(t) + M\frac{dI_p(t)}{dt} = 0$$
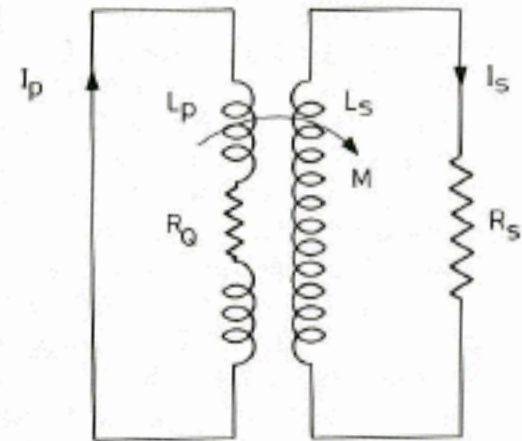
- Protection method used in magnets producing long-term steady-state fields such as detectors
  - Note: the coupling produces heat in the secondary during ramping.

$$L_p(1-k^2)\frac{dI_p(t)}{dt} + R_Q(t)I_p(t) - I_s(t)\frac{MR_s(t)}{L_s} = 0$$

Coupling coefficient

$$k = \frac{M}{\sqrt{L_p L_s}}$$

🔴 If the current in the secondary is small, the system behaves as if the inductance of the magnet was smaller:

$$L_{eff} = L_p(1 - \frac{M^2}{L_p L_s})$$

🔴 Which leads to a smaller time constant => faster current decay

$$\tau_{eff} = \frac{L_{eff}}{L_p}\tau_p$$

- Due to eddy currents, the secondary circuit becomes a heat source

- If the secondary is in thermal contact with the coil winding, the heat generated by the secondary can be used to quench the main coil faster. This is called **quench-back**

- We also refer to **quench-back** when any other mechanisms relying on dI/dt contribute to quenching the coil
  - Such as AC losses

- Some examples in the second protection session

- What are the basic requirements for quench-back to occur?

  - The secondary must heat above the current sharing temperature of the coil (+margin for thermal diffusion)

  - The secondary must be in good thermal contact with the magnet

  - The thermal rise time for the secondary must be significantly shorter than the magnets characteristic time constant
    - this is the time constant associated with combined normal zone growth, external resistor, and the inductive contribution from the secondary

- The coil is subdivided in several sections
- when one section quenches, the current of the other section is shunted through a resistance
- If the resistance are in thermal contact with the coil sections, they can be used as passive heaters as they heat up
  - Similar to quench-back

- The technique can be expanded to a large number
of subdivisions

- Diodes can be integrated to the circuit
  - Diodes are nonlinear circuit elements that carry
  current only when subjected to voltages greater than
  a threshold voltage
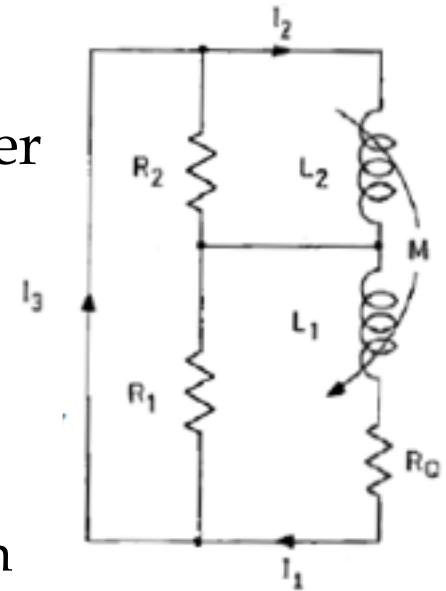
- Quench and protection are source of voltages in the magnet. Several possible failure points:
  - Coil to ground (or to structure): ground insulation
  - Internal failure: layer to layer, turn-to-turn
  - Coil to Protection heaters
- Need to provide a reliable insulation scheme
 - Example of a TQ LHC Accelerator Research Program (LARP)

bronze or titanium outer pole
5 mil dry S2 glass sheet
5 mil Kapton trace
20 mils
10 mils
10 mils
5 mil G-10 shim
bronze or titanium inner pole
5 mil Kapton trace
10 mil ceramic and binder sheet
5 mil dry S2 glass sheet
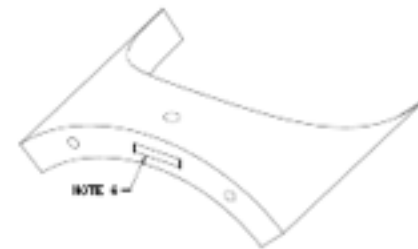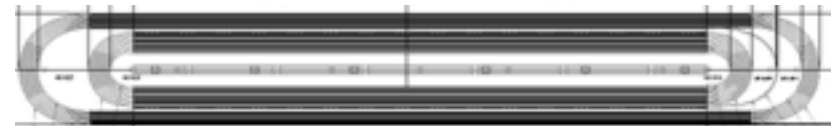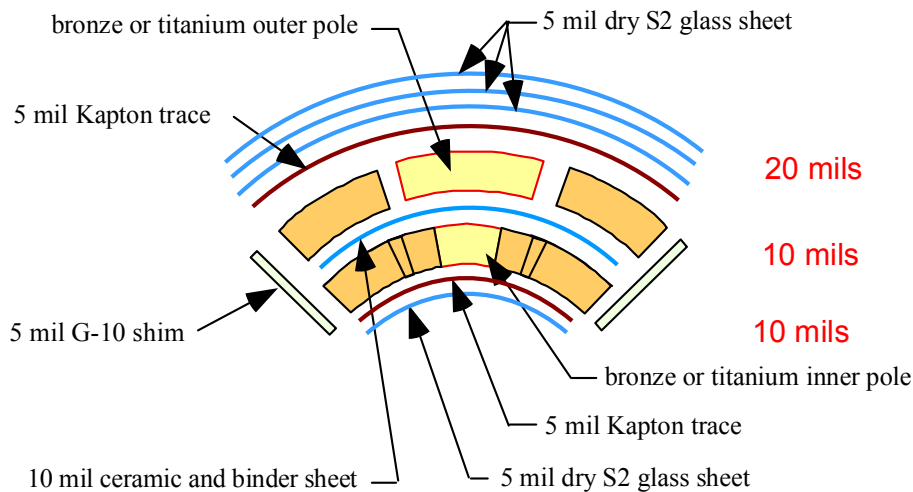
Courtesy of Rodger Bossert, FNAL

- Quality control is important to insure proper electrical integrity
  - Hipot
    - Look for leakage current between coil and magnet components under high DC voltage
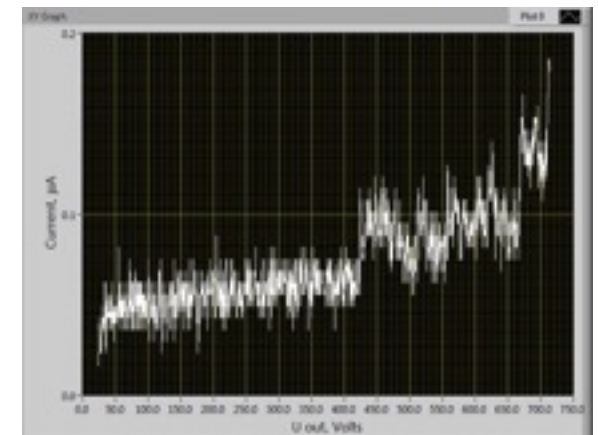   - Verify that the coil sustains the required voltage limits without breakout
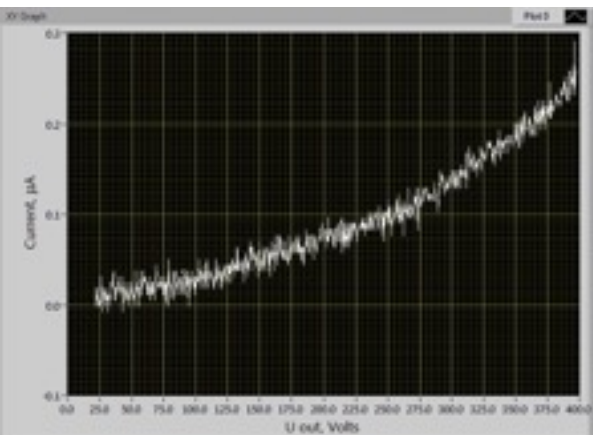
Coil to end-shoe
Non-linear current rise indicates **decrease of resistance** with the applied voltage. This is a "warning sign" sign of potential breakthrough failure.
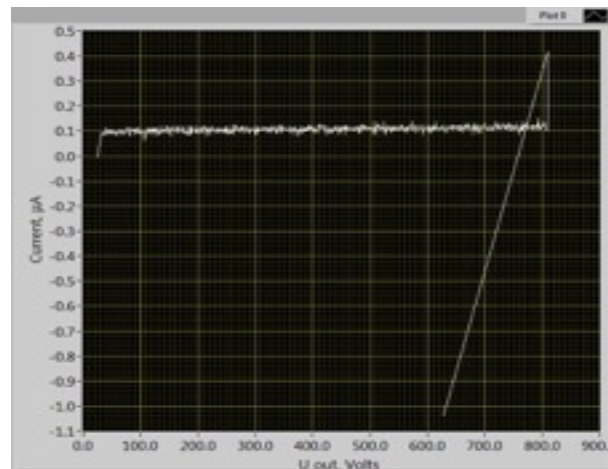
Coil to structure
Failure occurred with no "pre-cursor"

Coil to protection heater
Steep current rise is another "precursor" of failure

Courtesy of Maxim Marchevsky, LBNL

- The purpose is to test for shorts and insulation breakdown, turn-to turn and coil–to-coil parts

- A steep voltage impulse is sent to the winding
  - Generated by the discharge of a capacitor
  - This non-linear voltage creates a turn-to-turn voltage difference
  - Damped oscillation of the voltage with frequency:

$$f = \frac{1}{2\pi}\sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}}$$

Here L is the inductance of the coil and C is the net capacitance (including inter-turn and coil–to-coil parts capacitance)

Impulse test failure/breakout

300 V, PASS

1000 V, FAIL

Courtesy of Maxim Marchevsky, LBNL
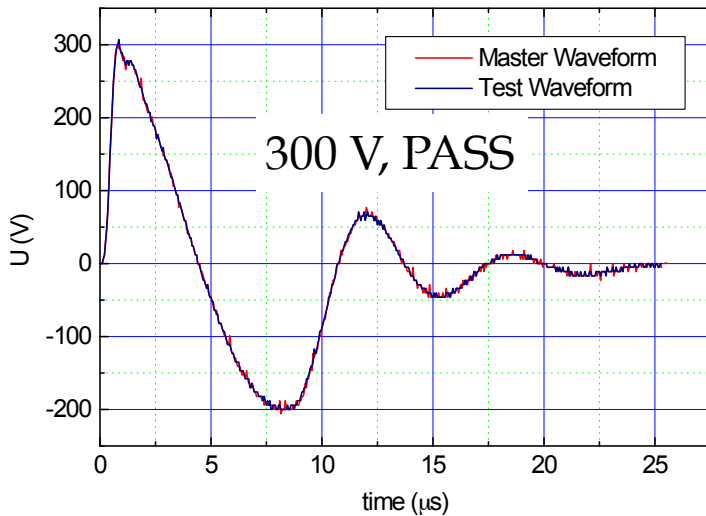
- A short through the **turn-to-turn** insulation will create a parallel path in the winding and reduce the effective inductance, thus **<u>increasing</u>** the oscillation frequency

- A short through the insulation between **winding and parts** (end parts, central pole, etc…) will increase the effective capacitance, thus **<u>decreasing</u>** the oscillation frequency.

- A discharge (arcing), either inside the winding (between the layers) or outside the winding (coil to parts), will cause a characteristic high frequency (> 1 MHz) noise to appear near the peaks of the oscillatory waveform.
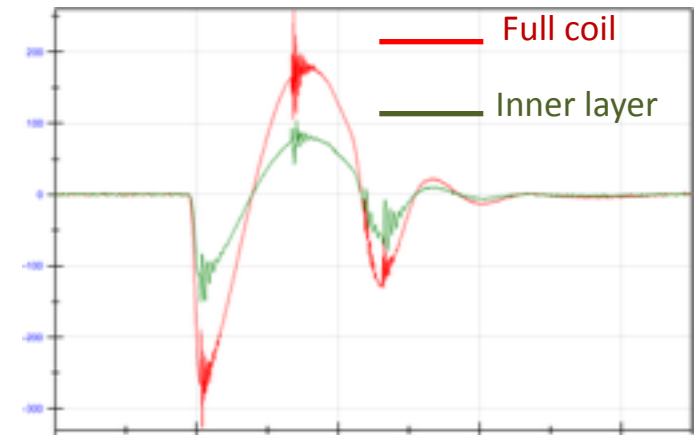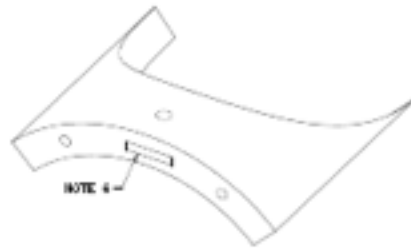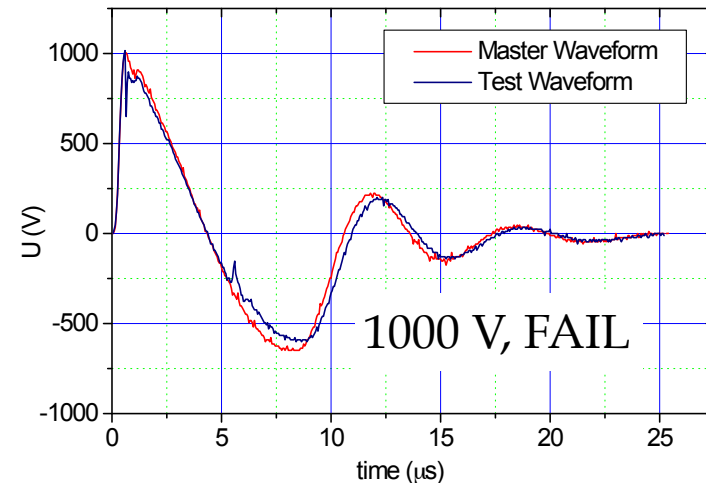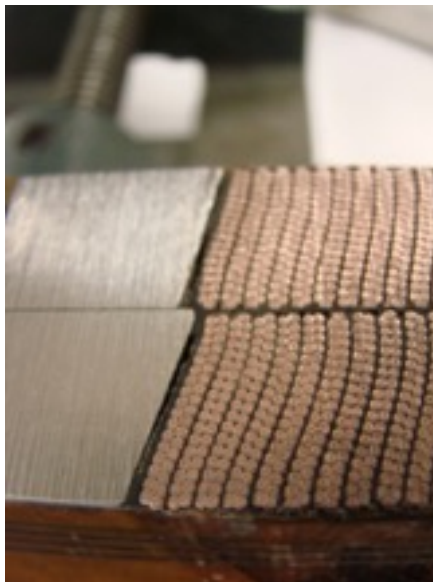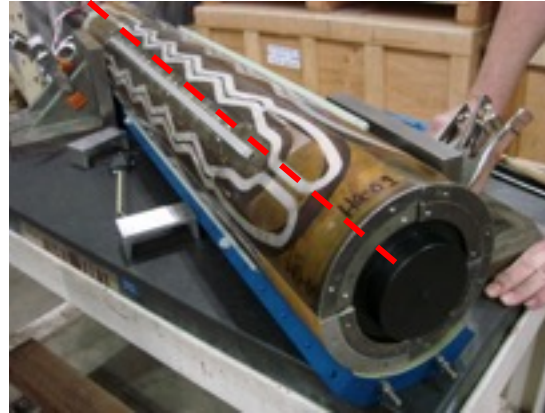

Full coil

Inner layer

Layer-to-layer failure

- Need to detect the quench as soon as possible to **avoid MIITs build-up**
  - Choice of voltage thresholds to be compared with magnet signals
  - Optimization of these thresholds

- Principle: looking for voltage imbalances in the magnet
  - Difference between 2 voltages which are supposed to be the same
    - Cancelation of the inductive component
    - Allowing detection of resistive voltage growth

- The detection triggers the extraction and the firing of the protection heaters
  - In case of failure, all the protection scheme is impaired
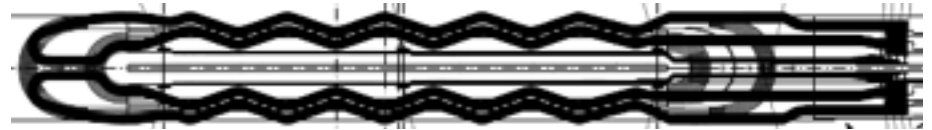 - Redundancy required to improve reliability
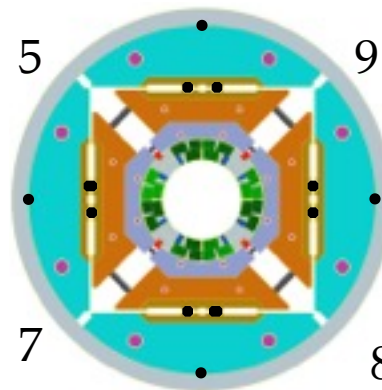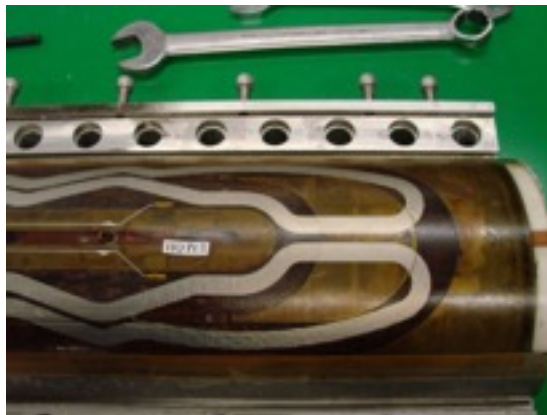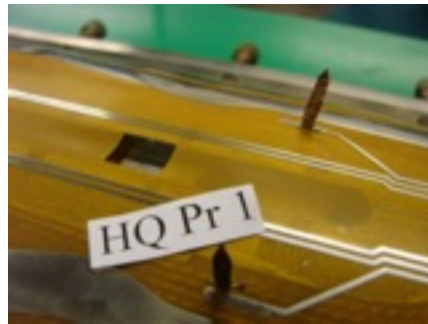

$t_R = L/R$

I_0

t_{det}
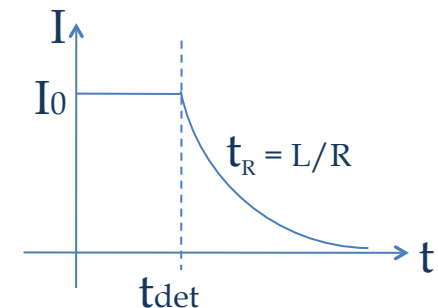
- Detection relies on instrumentation
  - Voltage taps in the magnet

  - Voltage taps installation on a Nb$_3$Sn magnet
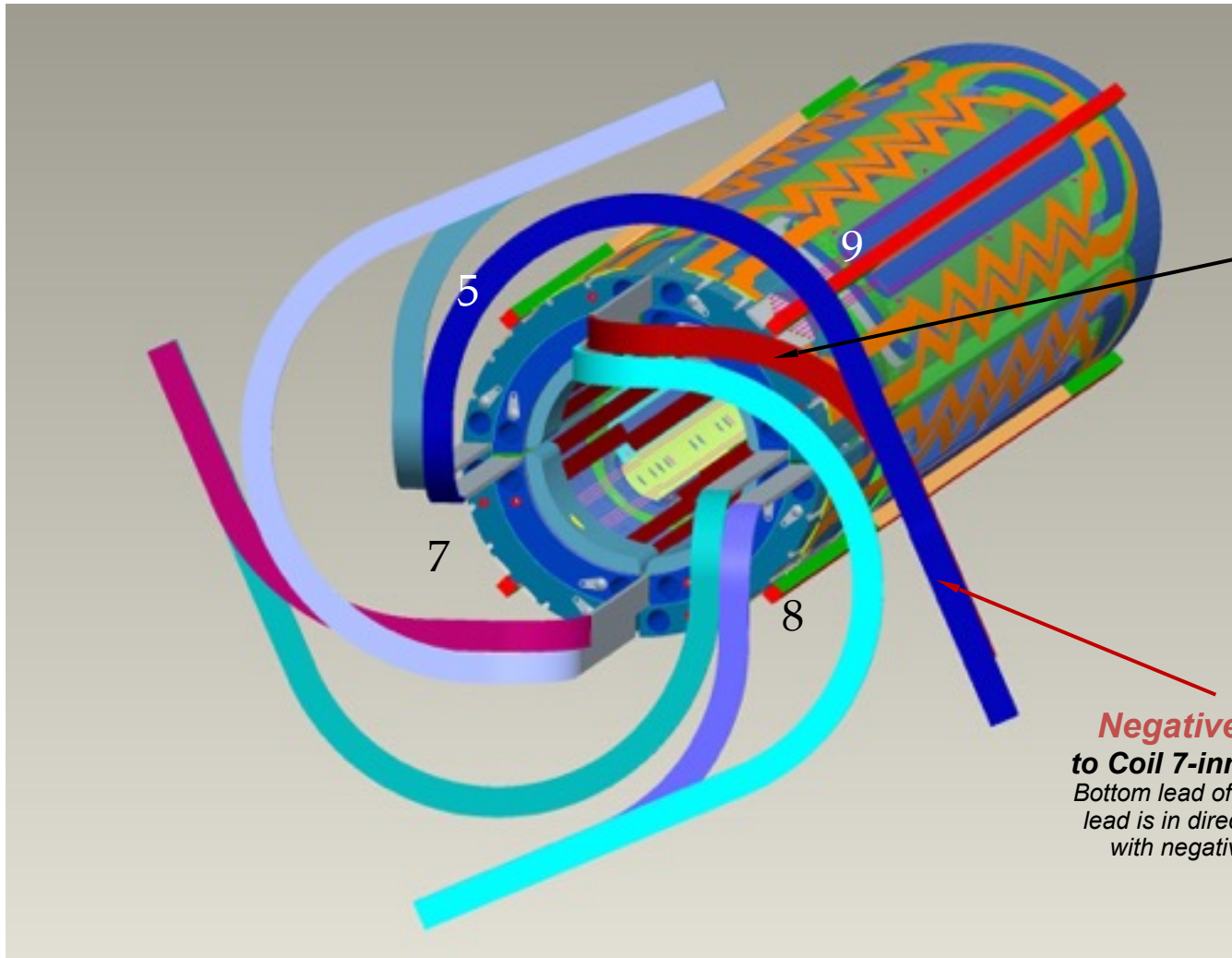
Ray Hafalia - LBNL

Positive lead
to Coil 9-outer layer
*Bottom lead of dual NbTi lead is in direct contact with positive flag.*

Negative lead
to Coil 7-inner layer
*Bottom lead of dual NbTi lead is in direct contact with negative flag.*

- Schematics representing all the voltage taps of the system
- Some vtaps are used to look at imbalances
- The rest of the vtaps aims at locating the quench origin

- Example of LBNL:
  - Imbalance 1:
    $$\frac{V_{mag}}{2} - V_{center}$$
  - Imbalance 2: difference between half magnet voltages

$$(V_{coil\,9} + V_{coil\,8}) - (V_{coil\,5} + V_{coil\,7})$$

  - Over-Voltage: absolute value of V_{mag}
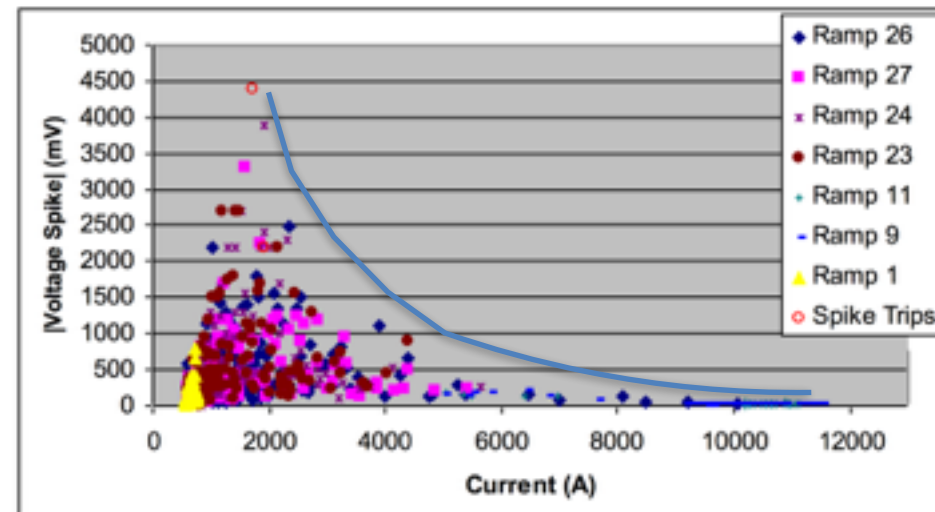

V_{center}
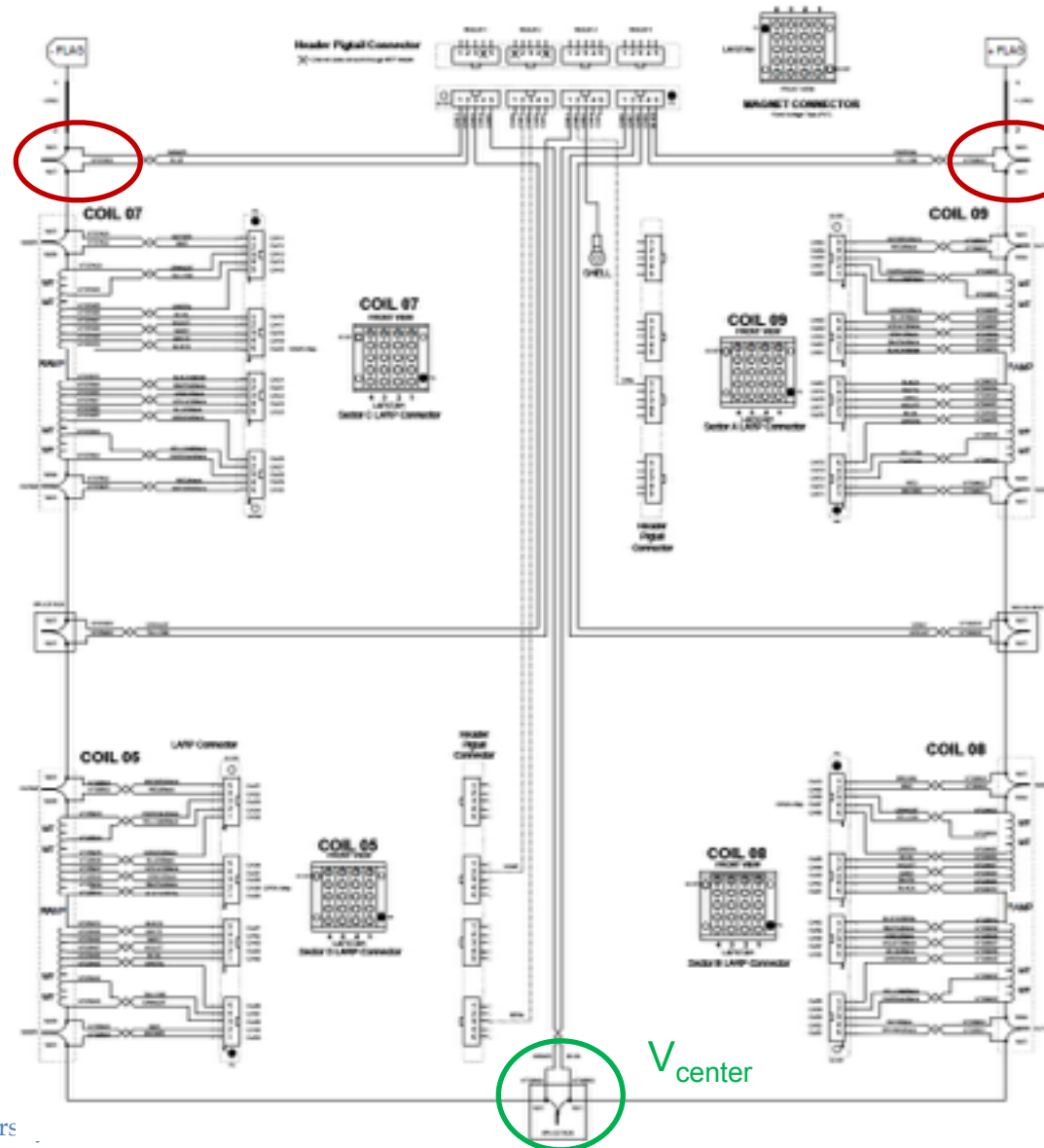
Courtesy of Xiaorong Wang, LBNL
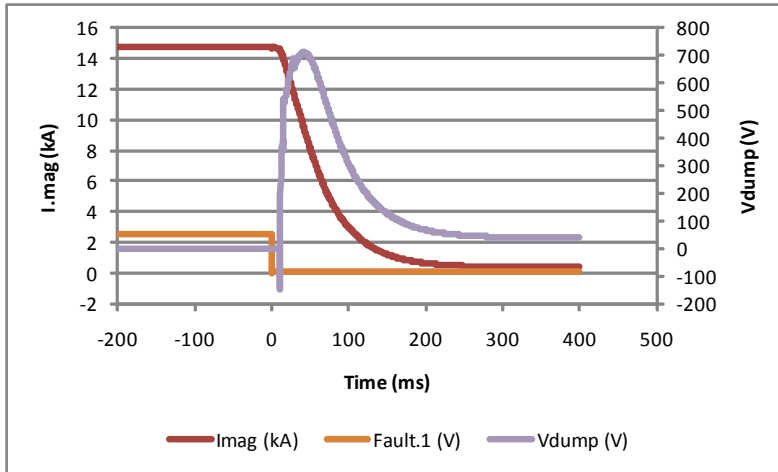
- The signals are then compared to threshold
- If threshold reached,
  - PS turned off ⎫
  - extraction   ⎬ Current decay
  - PH   ⎭

- In Nb$_3$Sn magnets, it is challenging to pick a safe threshold
  - Because of flux jumps
  - Requires to adapt threshold as the current is ramped up

FNAL TD-07-015 TQS02a Voltage spike analysis – C. Donnelly et al.

Summary

- The goal of quench protection is to convert - **safely -** magnetic energy into heat.
  - Minimizing hot spot temperature and peak voltages

- Protection strategies rely on various methods
  - Active protection: relies on external resistance and/or protection heaters
  - Passive protection: dissipation in a secondary circuit via magnetic coupling. It can include quench-back.

- Electrical integrity must be guaranteed
  - Importance of the QA methods: hipots and impulse tests

- Quench detection also plays a key role in the protection scheme

- Some examples of magnet systems protection

- An example of protection heater design

- Presentation of some codes

References

- M.K. Wilson, Superconducting Magnets, Oxford, Clarendon Press, 1983.
- A. Devred, General Formulas for the adiabatic propagation velocity of the normal zone, IEEE Trans on Magnetics, Vol. 25, No. 2, March 1989
- F.R. Fickett, Transverse magnetoresistance of oxygen free copper, IEEE Trans. On Magn., vol 24, No 2, March 1988
- G. Ambrosio et al, TQS01c test results summary, TD-07-007

- To Helene Felice, Xiaorong Wang, Maxim Marchevsky and John Joseph at LBNL

- To Al McInturff, Texas A&M

- To Giorgio Ambrosio and Guram Chlachidze at FNAL

- To Jerome Feuvrier and Marta Bajko at CERN