

# ***SEARCH FOR INTRINSIC CHARM IN HARD HADRON PROCESSES AT ATLAS: STATUS AND PERSPECTIVES***



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***<sup>2</sup>JINR, Dubna***

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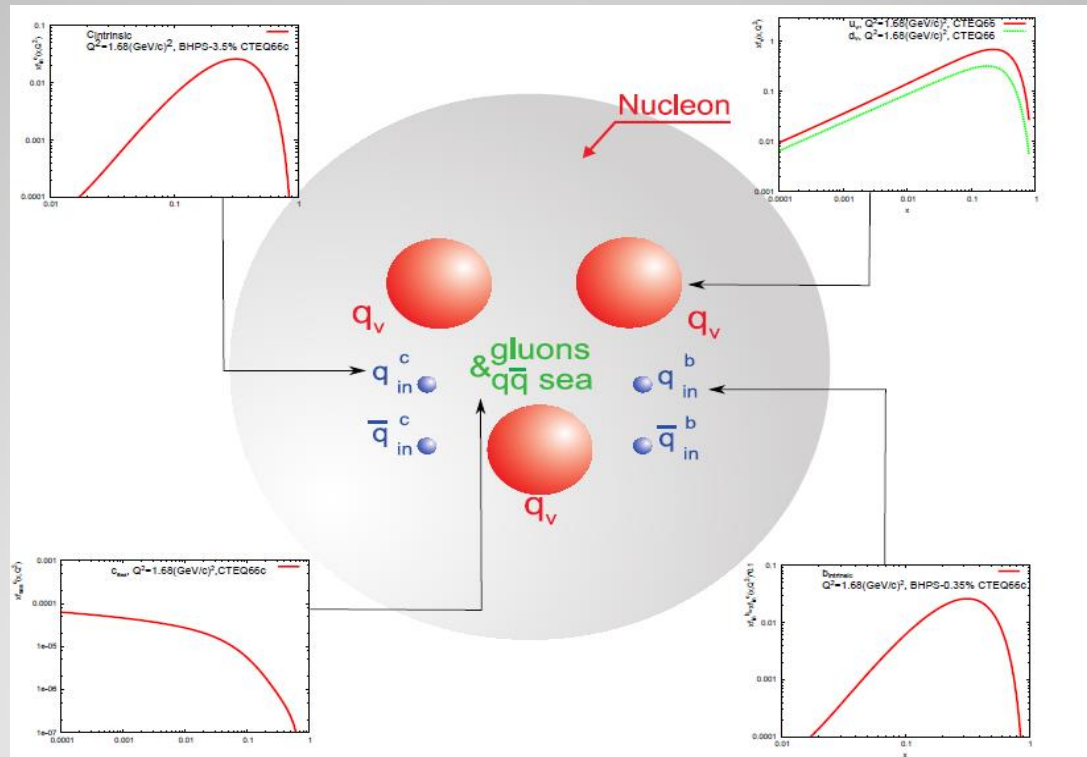


## *OUTLINE*

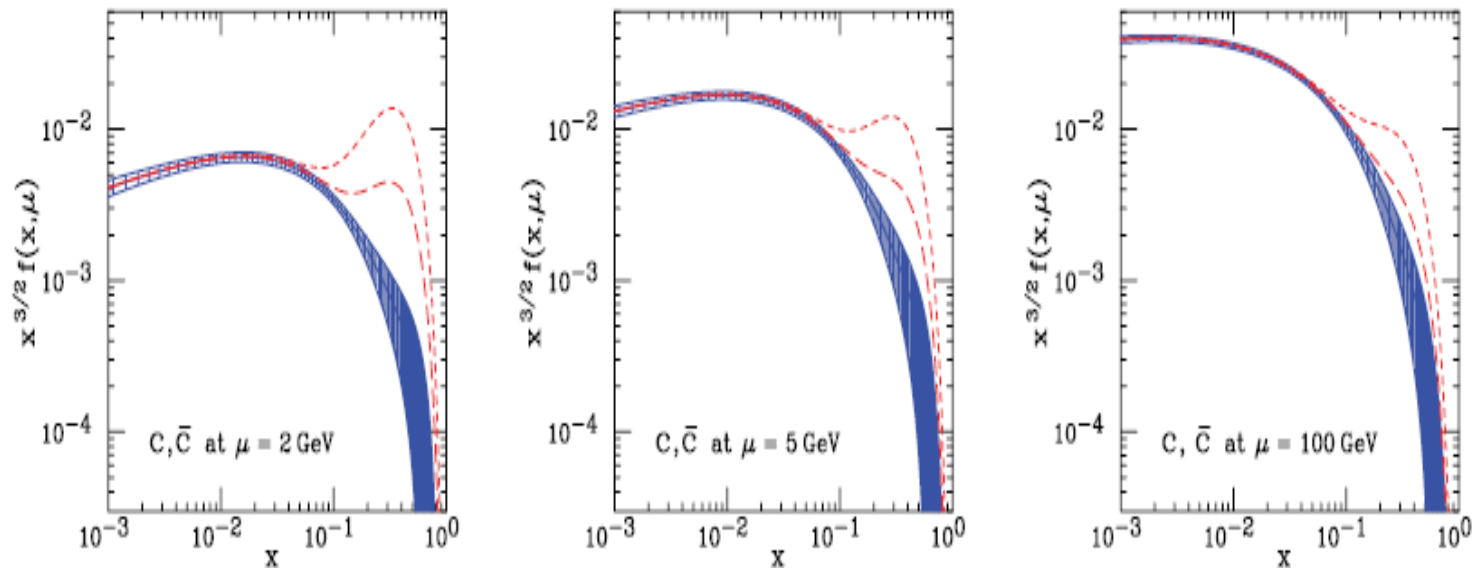
1. Intrinsic heavy flavour in proton
2. Main goal of our study
3. New theoretical predictions on the search of intrinsic charm (*IC*) signal in production of  $\gamma + c(b)$  and  $Z/W+c(b)$  in p-p at  $s^{1/2} = 8, 13$  TeV
4. Status of **ATLAS** experiments on possible observation of *IC* signal in these hard processes
5. Perspectives of searching of the *IC* signal at **ATLAS** experiments

**BHPS** model: S.J. Brodsky, P. Hoyer, C. Peterson and N. Sakai, Phys.Lett.B9(1980) 451; S.J. Brodsky, S.J. Peterson and N. Sakai, Phys.Rev. D23 (1981) 2745.

## Intrinsic $Q\bar{Q}$ in proton



## CHARM QUARK DISTRIBUTIONS IN PROTON



Charm quark distributions within the BHPS model. The three panels correspond to the renormalization scales  $\mu = 2, 5, 100$  GeV respectively. The long-dashed and the short-dashed curves correspond to  $\langle x_{c\bar{c}} \rangle = 0.57\%, 2.0\%$  respectively using the PDF CTEQ66c. The solid curve and shaded region show the central value and uncertainty from CTEQ6.5, which contains no **IC**.

**There is an enhancement at  $x > 0.1$  due to the IC contribution**

*Main goal: searching for the signal of the intrinsic charm (IC) contribution in proton from the analysis of the prompt photon or Z/W boson production in p-p collision accompanied by heavy c(b)-jet.*

## STATUS OF SEARCHING FOR INTRINSIC CHARM AT ATLAS EXPERIMENTS

**I.** We have predictions on  
*PP*-> $\gamma + c + X$  V.A.Bednyakov, M.A.Demichev, G.L., T.Stavreva,  
M.Stockton, *Phys.Lett. B728, 602 (1914)*  
and *PP*-> $Z/W + c(b) + X$  H.Beauchemin, V.A.Bednyakov, G.L.,  
Yu. Yu. Stepanenko, *Phys.Rev.D92, 034014 (2015)*

## Status of this research

**II.** Data analysis on  $\gamma + c$  production in p-p at  $s^{1/2} = 8$  TeV is performed by many people, for example,

*Juraj Smiesko from Bratislava University*

*Robert Keys from McGill University, Canada*

There are already very preliminary data which were presented at CERN informal meetings.

**III.** Data analysis on  $\gamma + b$  produced in p-p at 8 TeV.

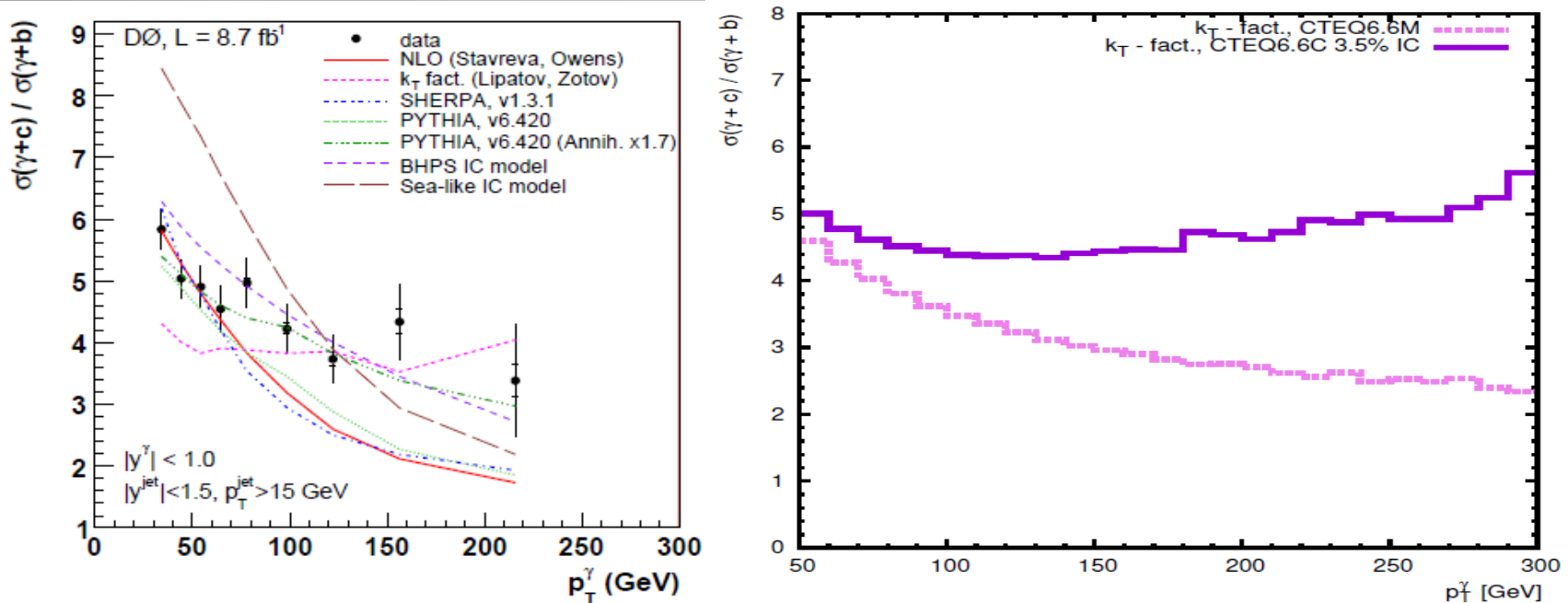
**IV.** Data analysis on **Z and W** bosons accompanied by heavy flavour jets **c** and **b** is started by

*Juri Stepanenko, Evelin Meoni (CERN), Hugo*

*Beauchemin (Tifs University, USA)*



# New results on $\gamma + c$ production in p-p



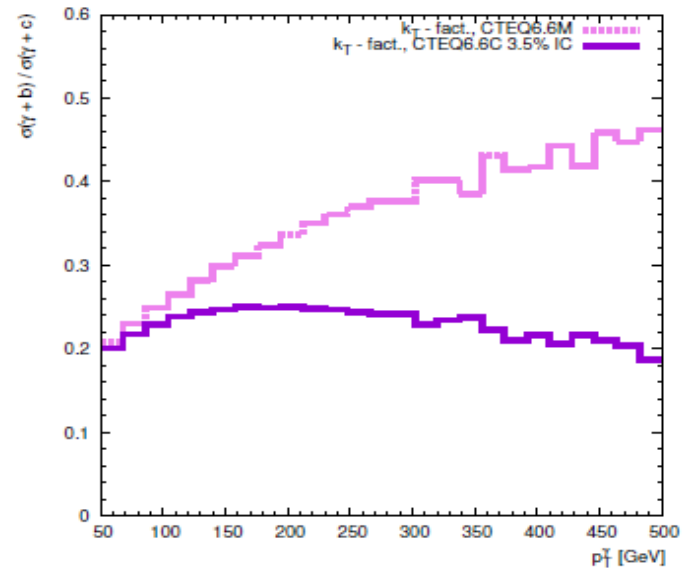
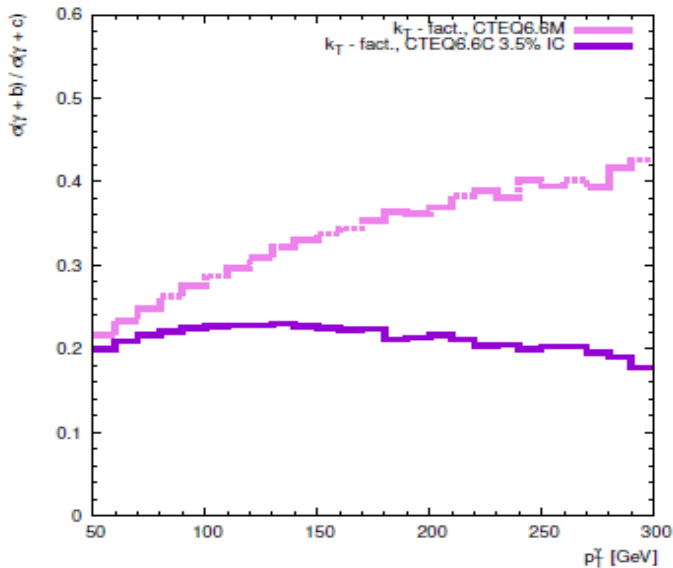
*Ratio of differential cross sections of  $\gamma + c$  and  $\gamma + b$  productions as function of photon transverse momentum*

*Left :  $D0$  data from TEVATRON and theoretical calculations*

*D0 Collaboration, Phys. Lett. B 719, 354 (2013); arXiv:1210.5033 [hep-ex].*

*Right: our latest predictions with the IC (top line) and without it (bottom line)*

# $pp \rightarrow \gamma + c(b) + X$

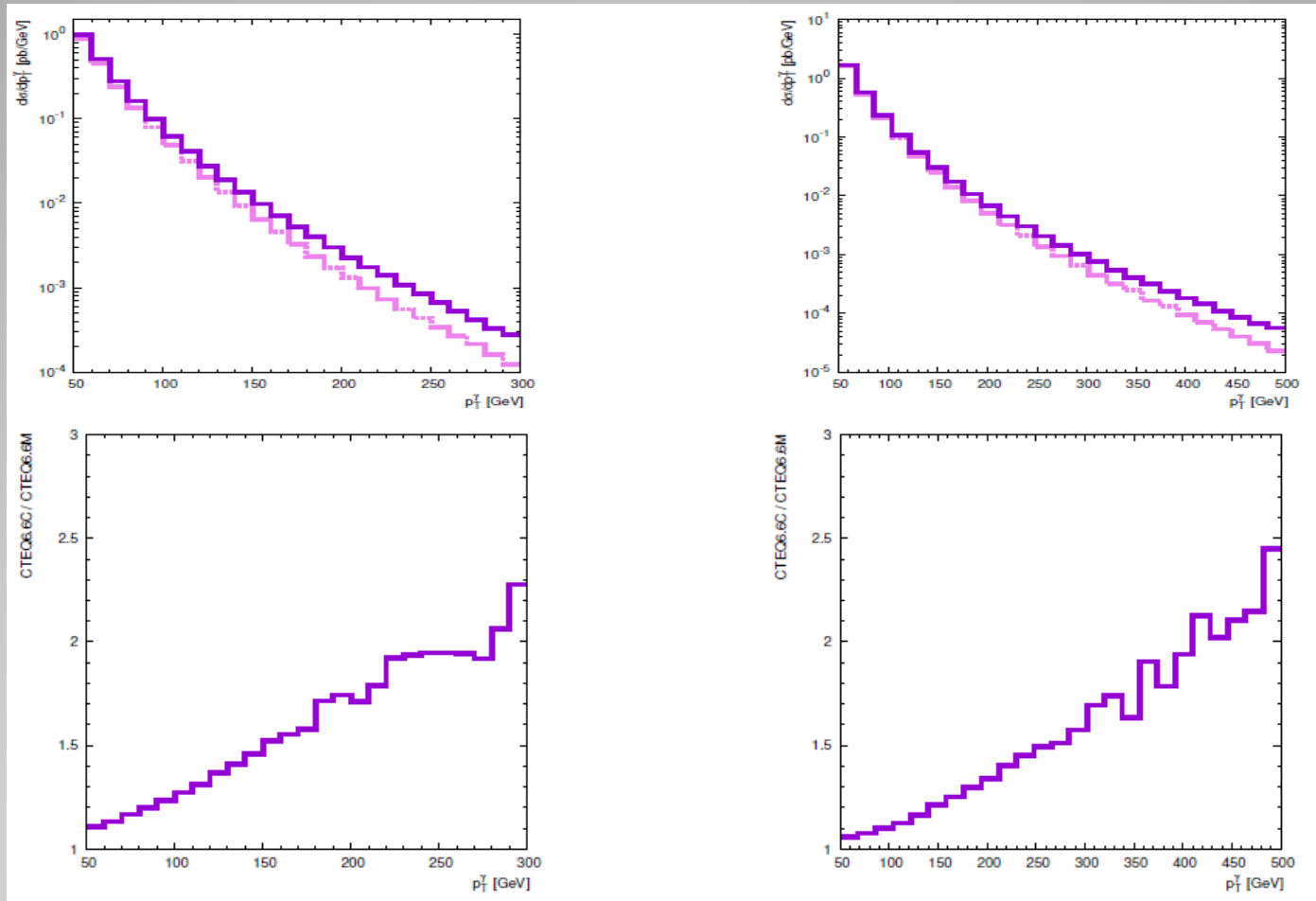


The ratio  $\sigma(\gamma + b)/\sigma(\gamma + c)$  as a function of the photon transverse momentum at  $1.5 < |y^\gamma| < 2.4$ ,  $|y^{\text{jet}}| < 2.4$ ,  $p_T^{\text{jet}} > 20$  GeV,  $\sqrt{s} = 8$  TeV (left panel) and  $\sqrt{s} = 13$  TeV (right panel)

**One can see a grow of this ratio by increasing of  $p_T^\gamma$ , when the intrinsic charm contribution (IC) in PDF is ignored. This ratio is approximately flat or slowly decreasing when the IC about 3.5 % is included. That is an additional IC signal, which could be measured at ATLAS.**

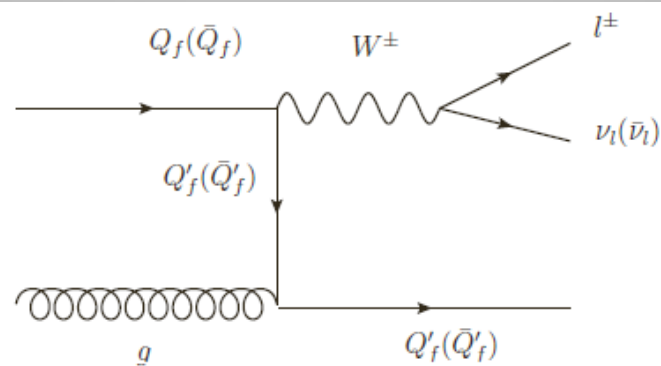


# $pp \rightarrow \gamma + D^* + X$

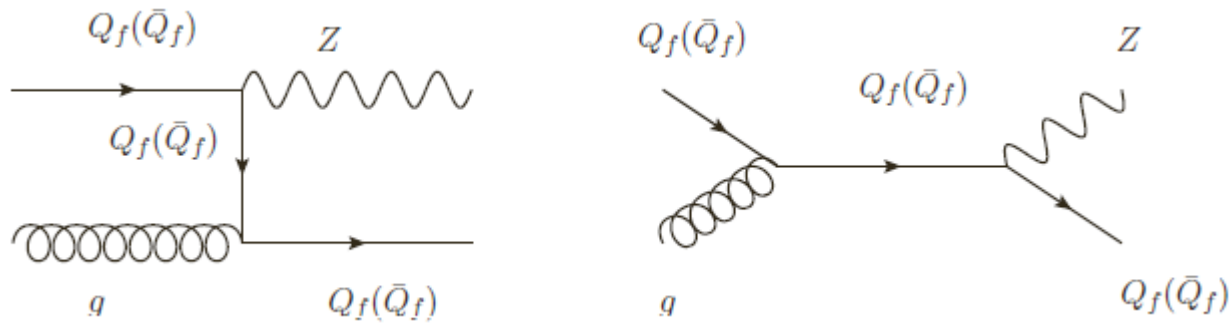


**Top:**  $p_T$ –spectra of a photon at  $\sqrt{s} = 8$  TeV (left) and 13 TeV (right)  
**Bottom:** ratio of  $p_T$ –spectra with IC (3.5%) and without it at 8 TeV (left) and 13 TeV (right).

# $pp \rightarrow Z/W + \text{heavy flavour jets}$

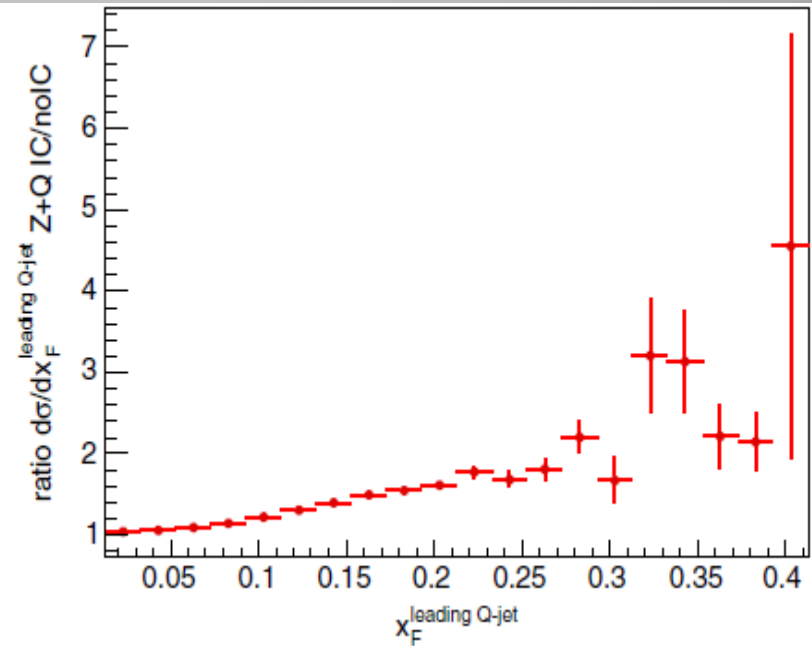
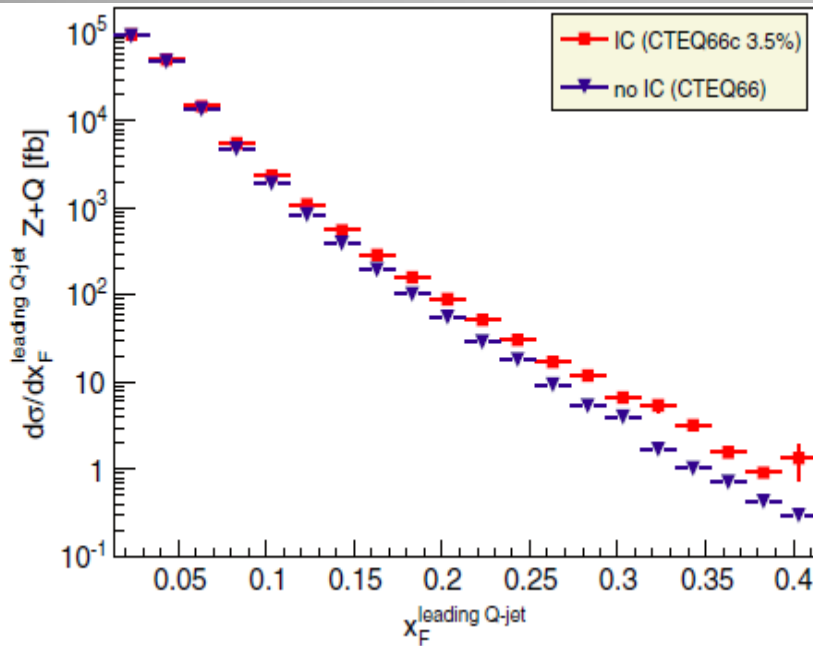


The LO Feynman diagrams for the process  $Q_f(\bar{Q}_f)g \rightarrow W^\pm Q'_f(\bar{Q}'_f)$ , where  $Q_f = c, b$  and  $Q'_f = b, c$  respectively.



Feynman diagram for the process  $Q_f(\bar{Q}_f)g \rightarrow Z Q_f(\bar{Q}_f)$

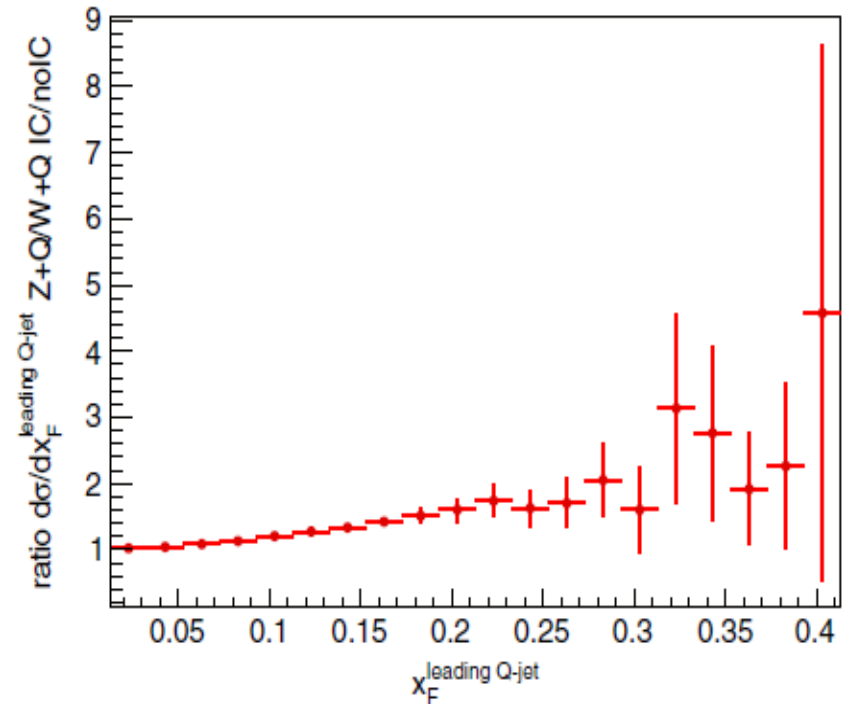
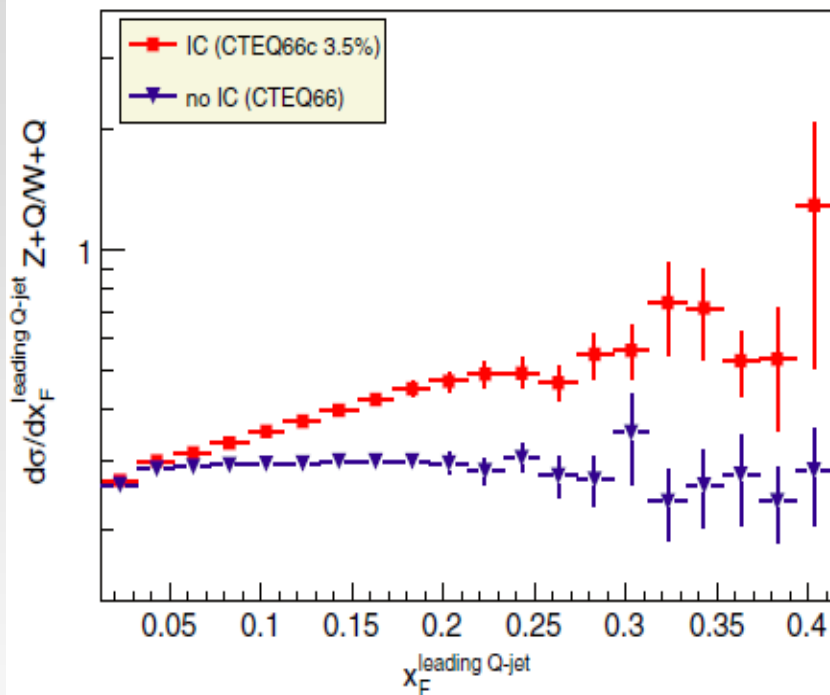
# New results on $Z+Q$ production in $pp$ at $s^{1/2} = 8$ TeV ( $Q=c,b$ )



**Left:**  $x_F$  –spectra of leading  $c$  or  $b$  jet with the  $IC$  (red line) and without it (blue line).

**Right:** ratio of  $Z+Q$  spectra with  $IC$  and without it as a function of  $x_F$ .

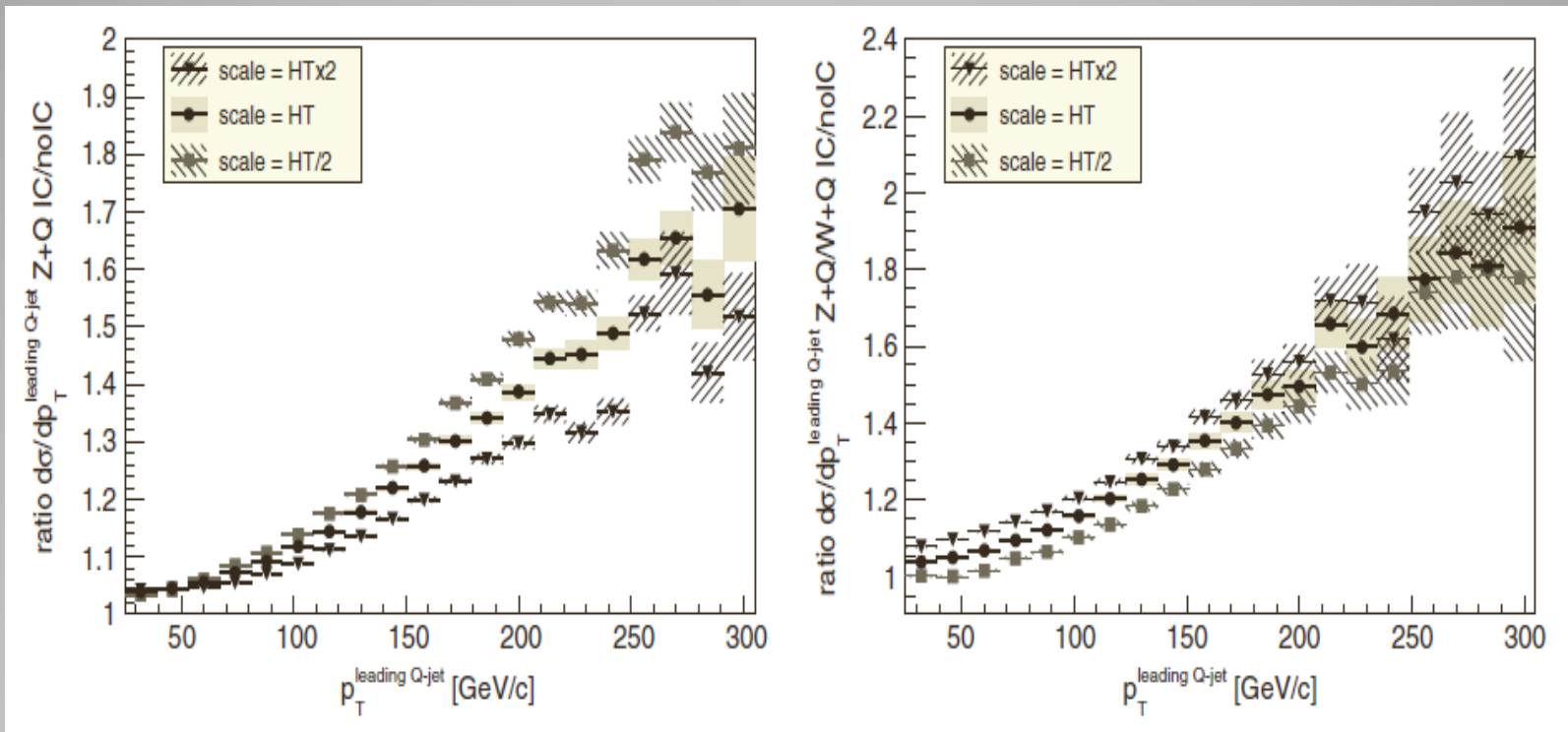
# New results on $Z+Q/W+Q$ at $s^{1/2} = 8$ TeV



**Left:** ratio  $Z+Q/W+Q$  as a function of  $x_F$  of leading  $c$  or  $b$  jet with the  $IC$  (red line) and without it (blue line).

**Right:** double ratio of  $Z+Q/W+Q$  spectra with  $IC$  and without it as a function of  $x_F$ .

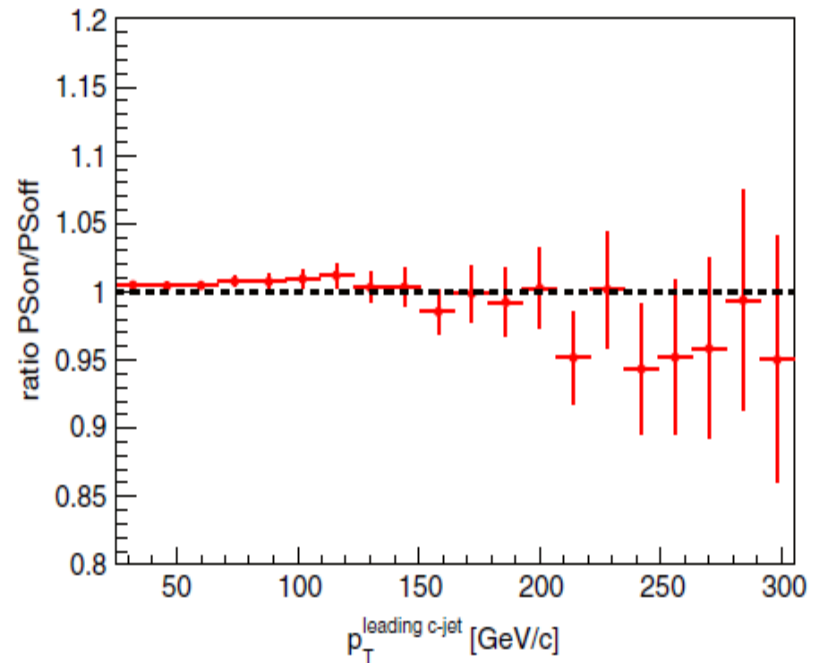
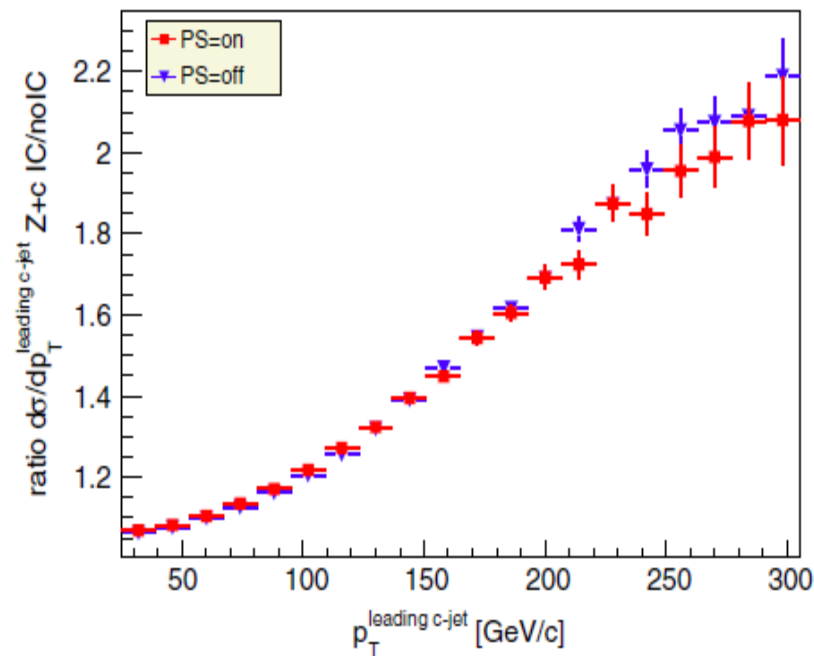
## Scale uncertainty for Z+Q and Z+Q/W+Q



**Left:** Z+Q with IC and without IC at different scales

**Right:** Z+Q/W+Q with IC and without IC at different scales

# Inclusion of parton shower by Z+Q production in pp at $s^{1/2} = 8$ TeV



- Left:**  $p_T$  – spectra with parton shower (red points) and without it (blue points) using the PDF of type CTEQ66c (3.5% of *IC*)
- Right:** ratio of red points to blue points



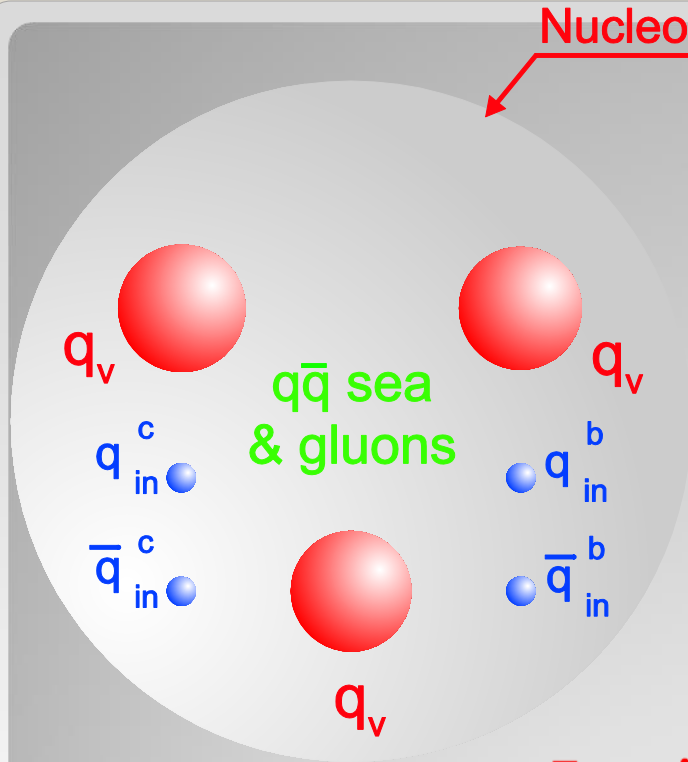
# Perspectives of search for intrinsic charm at ATLAS

1. Informal photon group at ATLAS is going to finish the data analysis on  $\gamma + c(b)$  production in p-p at 8 TeV and prepare a paper on it in the end of 2015, and to do the similar analysis at 13 TeV in 2016.
2. We are doing the theoretical predictions on  $p_T$  – spectra of photons within different MC generators and QCD models including and ignoring the *IC* contributions with different values of the *IC* probability. Then, we will compare our calculations to data.
3. From this comparison we intend to extract the *IC* probability from the ATLAS data.
4. There is a plan to finish data analysis of  $Z/W+c(b)$  production in p-p at 8 TeV and prepare a paper in March 2016, and turn to the similar analysis at 13 TeV.

**THANK YOU VERY MUCH FOR  
YOUR ATTENTION !**

**BACK UP**

Nucleon



## BHPS model

S.J. Brodsky, P. Hoyer, C. Peterson and N. Sakai, Phys.Lett. B93 (1980) 451;  
S.J. Brodsky, S.J. Peterson and N. Sakai, Phys.Rev. D23 (1981) 2745.

Intrinsic  $Q\bar{Q}$  in proton

$Q\bar{Q}$  is  $u\bar{u}, d\bar{d}, s\bar{s}, c\bar{c}, b\bar{b}, t\bar{t}$

J.Pumplin, H.L. Lai and W.K.Tung, Phys.Rev.D75 (2007) 054029

# INTRINSIC HEAVY QUARK STATES

Two types of parton contributions

The **extrinsic** quarks and gluons are generated on a short time scale in association with a large transverse-momentum reaction.

The **intrinsic** quarks and gluons exist over a time scale independent of any probe momentum, they are associated with the bound state hadron dynamics.

$$P(x_1, \dots, x_5) = N_5 \delta\left(1 - \sum_{i=1}^5 x_i\right) \left[ M_p^2 - \sum_{i=1}^5 \frac{m_i^2}{x_i} \right]^2$$

# INTRINSIC HEAVY QUARK DISTRIBUTION IN PROTON

Integrating  $P(x_1, \dots, x_5)$  over  $dx_1 \dots dx_4$  and neglecting of all quark masses except the charm quark mass we get

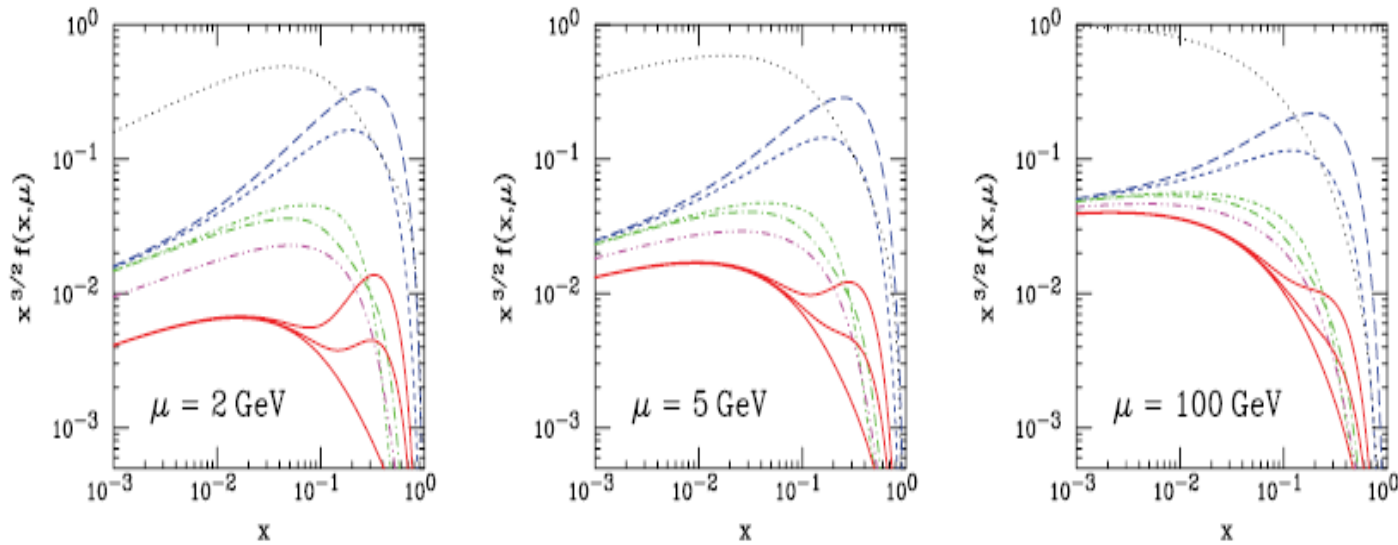
$$P(x_5) = \frac{1}{2} \bar{N}_5 x_5^2 \left[ \frac{1}{3} (1 - x_5)(1 + 10x_5 + x_5^2) + 2x_5(1 + x_5) \ln\left(\frac{1}{x_5}\right) \right]$$

Where  $\bar{N}_5 = N_5 / m_{4,5}^4$  normalization constant. Here  $m_4 = m_5 = m_c = m_{\bar{c}}$  is the bar mass of the charmed quark.  $N_5$  determines some probability  $w_{10}$  to find the Fock state  $|uud\bar{Q}Q\rangle$  in the proton.

**One can see qualitatively that  $P(x_5)$  vanishes at  $x_5 \rightarrow 0$  and  $x_5 \rightarrow 1$  and has an enhancement at  $0 < x_5 < 1$**



# COMPARISON OF LIGHT AND HEAVY QUARK DISTRIBUTIONS IN PROTON



The dotted line is the gluon distribution, the blue long-dashed curve is the valence  $\bar{u}$ -distribution, the blue short-dashed line is the valence  $\bar{d}$ -distribution, the green long-dashed-dotted line is the **intrinsic**  $\bar{u}$ , the short dashed-dotted line is the **intrinsic**  $\bar{d}$  distribution, the dashed-dot-dotted is the **intrinsic**  $\bar{s} = \bar{s}$  and the solid curves are  $\bar{c} = \bar{c}$  with **no IC** (lowest) and with **IC**,  $\langle x_{\bar{c}} \rangle = 0.57\%, 2.0\%$  respectively. It is shown that **IC** contribution is larger than  $\bar{u}, \bar{d}, \bar{s}$  at  $x > 0.2$

## PRODUCTION OF HEAVY FLAVOURS IN HARD P-P COLLISIONS

$$E \frac{d\sigma}{d^3p} = \sum_{i,j} \int d^2k_{iT} \int d^2k_{jT} \int_{x_i^{\min}}^1 dx_i \int_{x_j^{\min}}^1 dx_j f_i(x_i, k_{iT}) f_j(x_j, k_{jT}) \frac{d\sigma_{ij}(\hat{s}, \hat{t})}{d\hat{t}} \frac{D_{i,j}^h(z_h)}{\pi z_h}$$

$$x_i^{\min} = \frac{x_T \cot(\frac{\theta}{2})}{2 - x_T \tan(\frac{\theta}{2})}$$

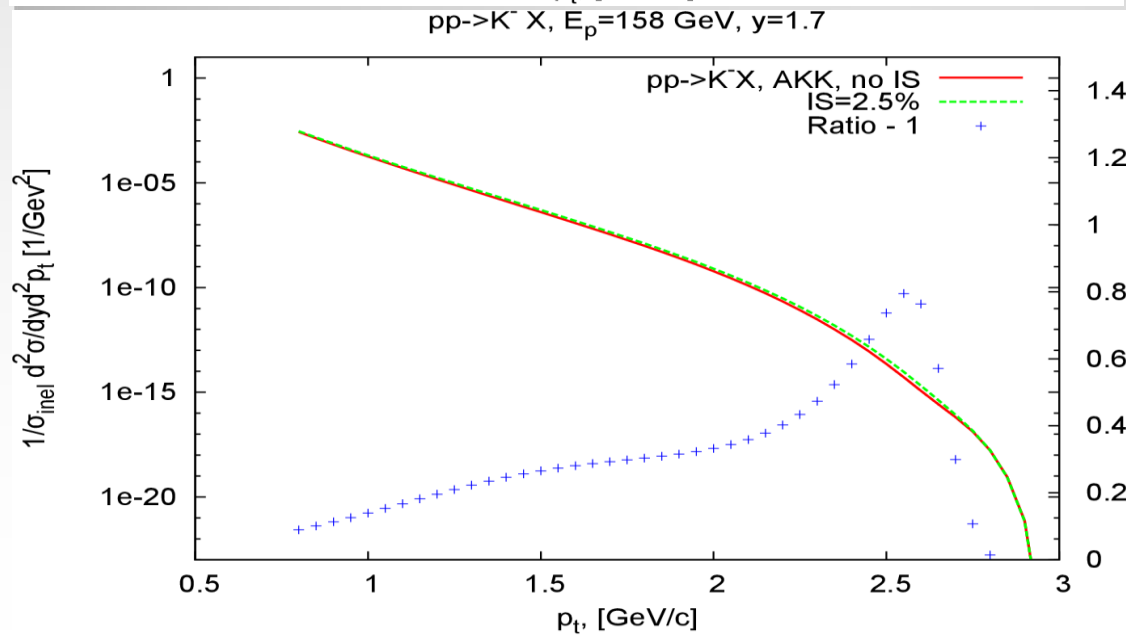
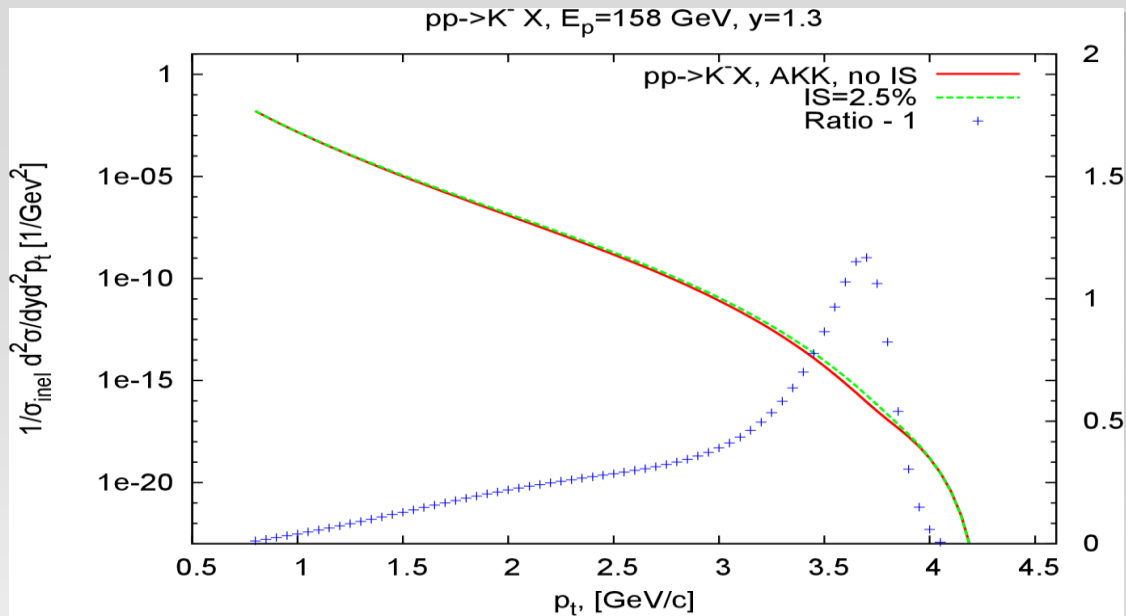
$$x_F \equiv \frac{2p_z}{\sqrt{s}} = \frac{2p_T}{\sqrt{s}} \frac{1}{\tan \theta} = \frac{2p_T}{\sqrt{s}} \sinh(\eta)$$

$$x_i^{\min} = \frac{x_R + x_F}{2 - (x_R - x_F)}$$

$$x_R = 2p/\sqrt{s}$$

One can see that  $x_i \geq x_F$ . If  $x_F > 0.1$  then,  $x_i > 0.1$  and the **conventional sea** heavy quark (extrinsic) contributions are suppressed in comparison to the **intrinsic** ones.

$x_F$  is related to  $p_T$  and  $\eta$ . So, at certain values of these variables, in fact, there is **no conventional sea** heavy quark (**extrinsic**) contribution. And we can study the **IQ contributions** in hard processes at the **certain** kinematical region.



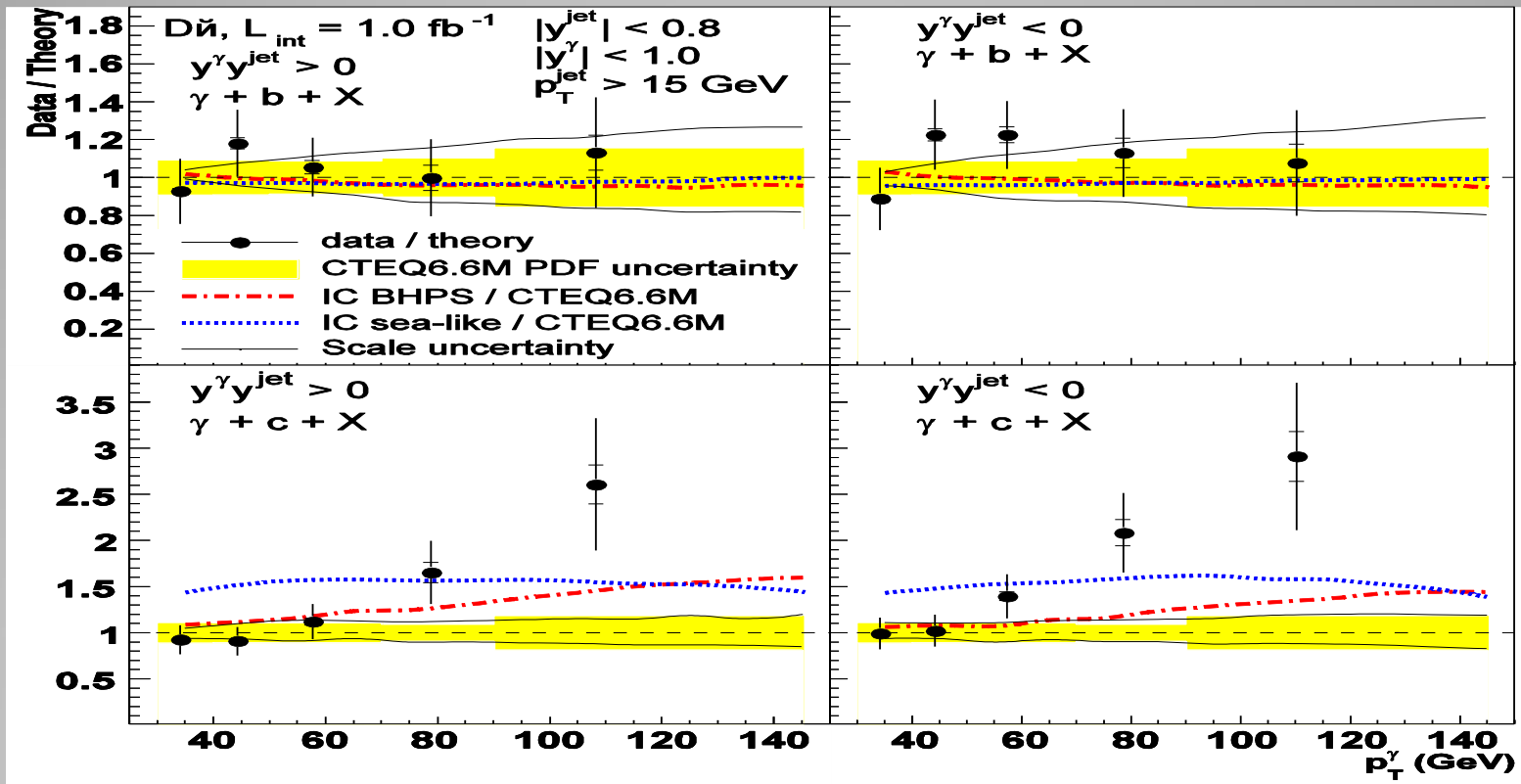
*G.L., A.A.Grinyuk,  
I.V.Bednyakov, Proc.  
Baldin Conference,  
Dubna, Sept. 2012; C12-  
09-10.4;  
arXiv:1212.6381 [hep-ph].*

The red line is the  $p_T$  spectrum of  $K^-$  mesons produced in p-p at  $E_p = 158$  GeV,  $y=1.3$  (top) and  $y=1.7$  (bottom) without IC; the green curve is the same as the red one but with the IC contribution, its probability is about 2.5%.

The dotted line corresponds to the ratio of the spectra with IC and without IC minus 1.

The IC signal is about 200% at high transverse momenta

$p\bar{p} \rightarrow \gamma + c(b) + X$  D0 experiment at Tevatron  $s^{1/2} = 1.96\text{TeV}$



The data-to-theory ratio of cross sections as a function of  $p_T^\gamma$  for  $p\bar{p} \rightarrow \gamma + c(b) + X$ . There is the **three time excess** of the data above the theory for  $\gamma + c$  at  $p_T > 150 \text{ GeV}/c$ . **It stimulates us to study**  $pp \rightarrow \gamma + c(b) + X$

# PHOTON (DI-LEPTON) AND c(b)-JETS PRODUCTION IN P-P

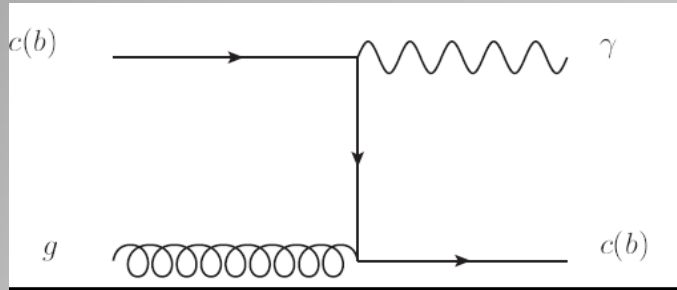


Fig.a. Feynman diagram for the process  $c(b) + g \rightarrow \gamma + c(b)$

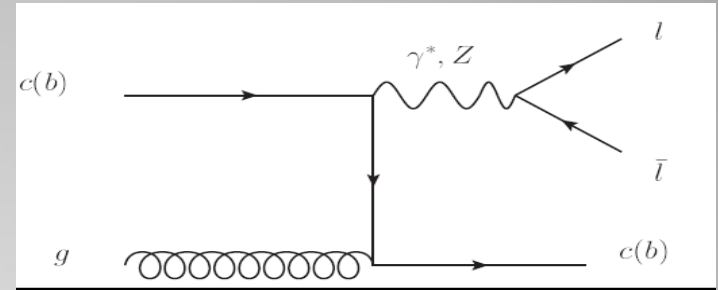


Fig.b. Feynman graph for the process  $c(b) + g \rightarrow \gamma/Z^0 + c(b)$

$$x_F = \frac{2p_T}{s^{1/2}} sh(\eta); p_{T\gamma} = -p_{Tc} \quad x_{c(b)} = \frac{m_{l^+l^-}^2}{x_g s} + x_{c(b)}^f$$

To observe the IC

for Fig.a

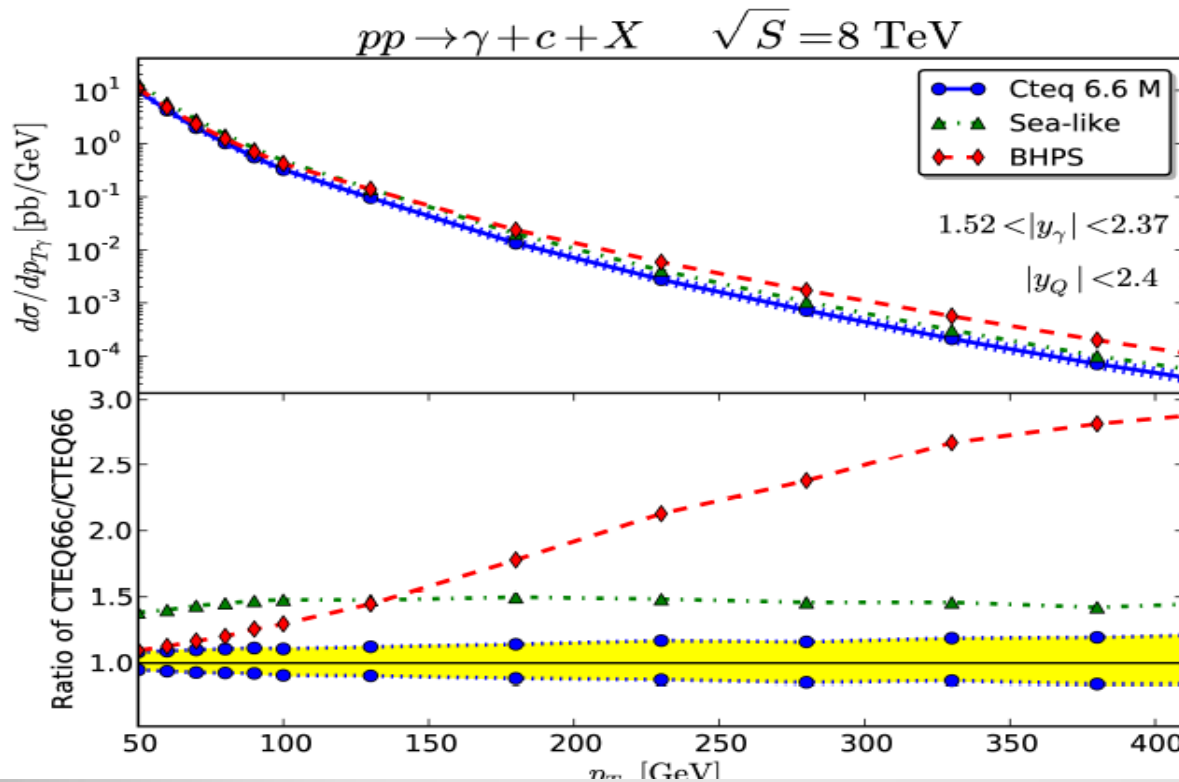
$$x_c \geq x_F > 0.1$$

for Fig.b

$$x_{c(b)} = \frac{m_{l^+l^-}^2}{x_g s} + x_{c(b)}^f > 0.1$$

# IC signal in $pp \rightarrow \gamma + c(\text{jet}) + X$

V.A.Bednyakov,  
M.A.Demichev,  
G.L., T.Stavreva,  
M.Stockton,  
hep-ph/1305.3548  
Phys.Lett. B728  
(2014) 602.

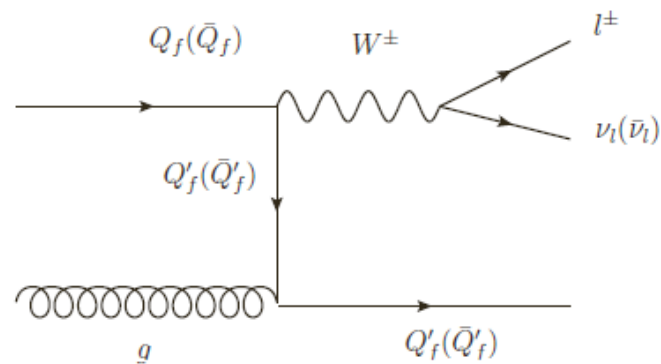


$p_T$  - distribution of photons produced in  $pp \rightarrow \gamma + c(\text{jet}) + X$

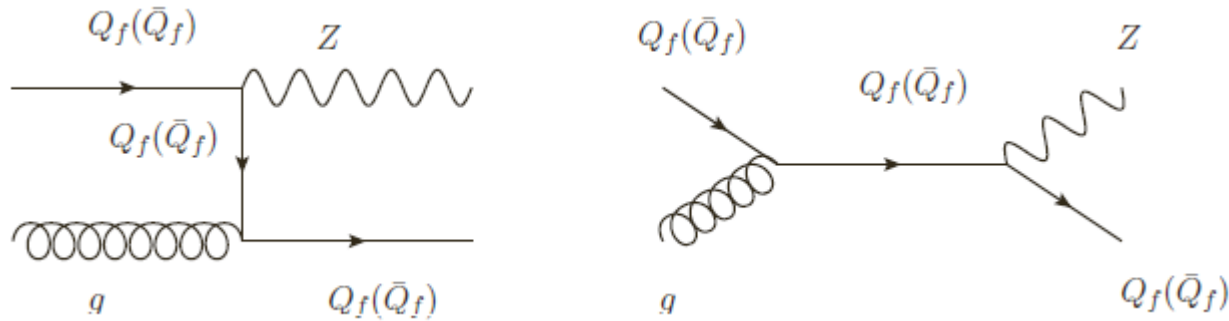
The blue line is calculation without the **IC**. The red curve includes the **IC**, its **probability is about 3.5 %** (top). The ratio of spectra with and without the **IC** **The IC signal is about 200%-250% at  $p_T \sim 150 - 200 \text{ GeV}/c$**  where the cross section is about 20-80 fb (400-3200 events) and can be measured



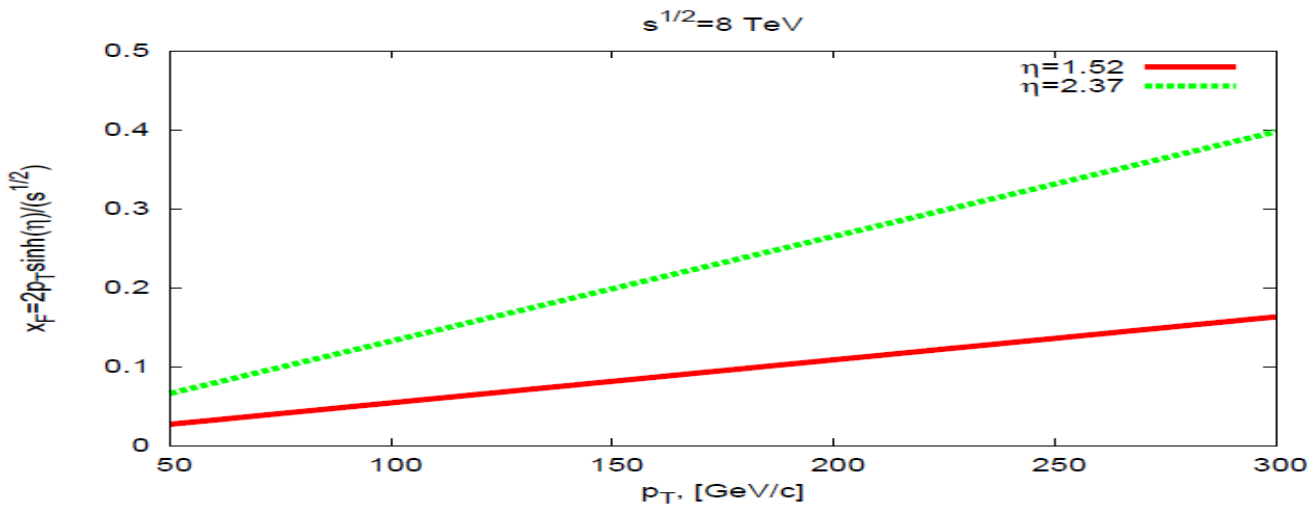
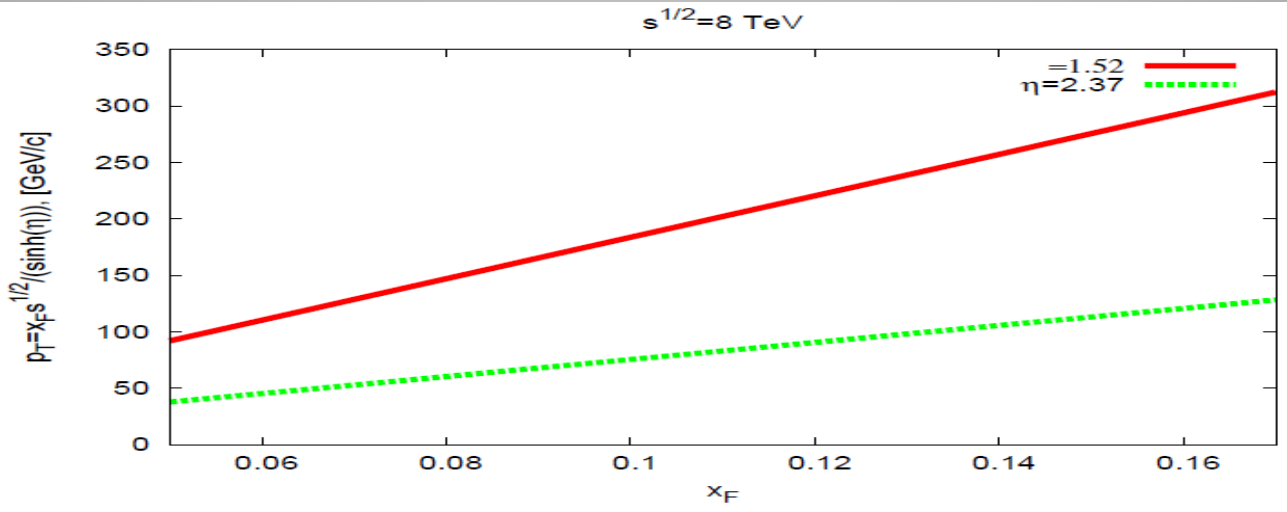
## pp $\rightarrow$ W/Z+heavy flavour jets



The LO Feynman diagrams for the process  $Q_f(\bar{Q}_f)g \rightarrow W^\pm Q'_f(\bar{Q}'_f)$ , where  $Q_f = c, b$  and  $Q'_f = b, c$  respectively.



Feynman diagram for the process  $Q_f(\bar{Q}_f)g \rightarrow ZQ_f(\bar{Q}_f)$



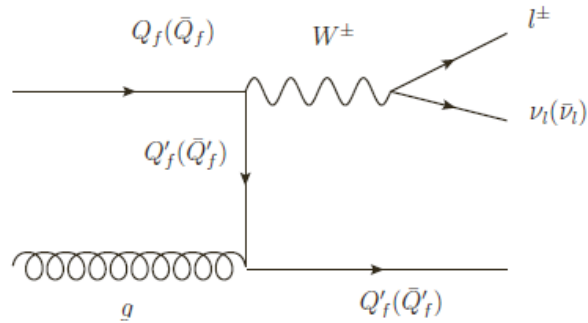


Figure 2: The LO Feynman diagrams for the process  $Q_f(\bar{Q}_f)g \rightarrow W^\pm Q'_f(\bar{Q}'_f)$ , where  $Q_f = c, b$  and  $Q'_f = b, c$  respectively.

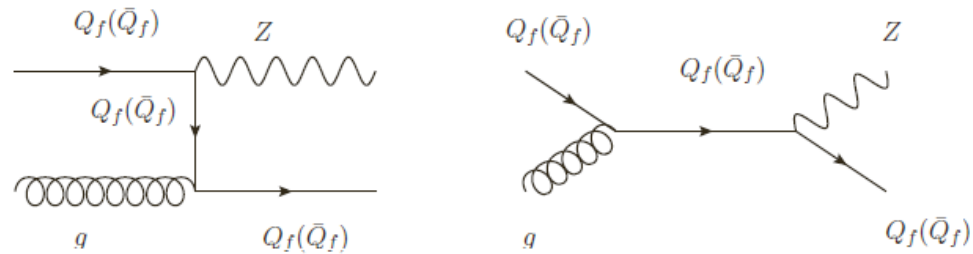
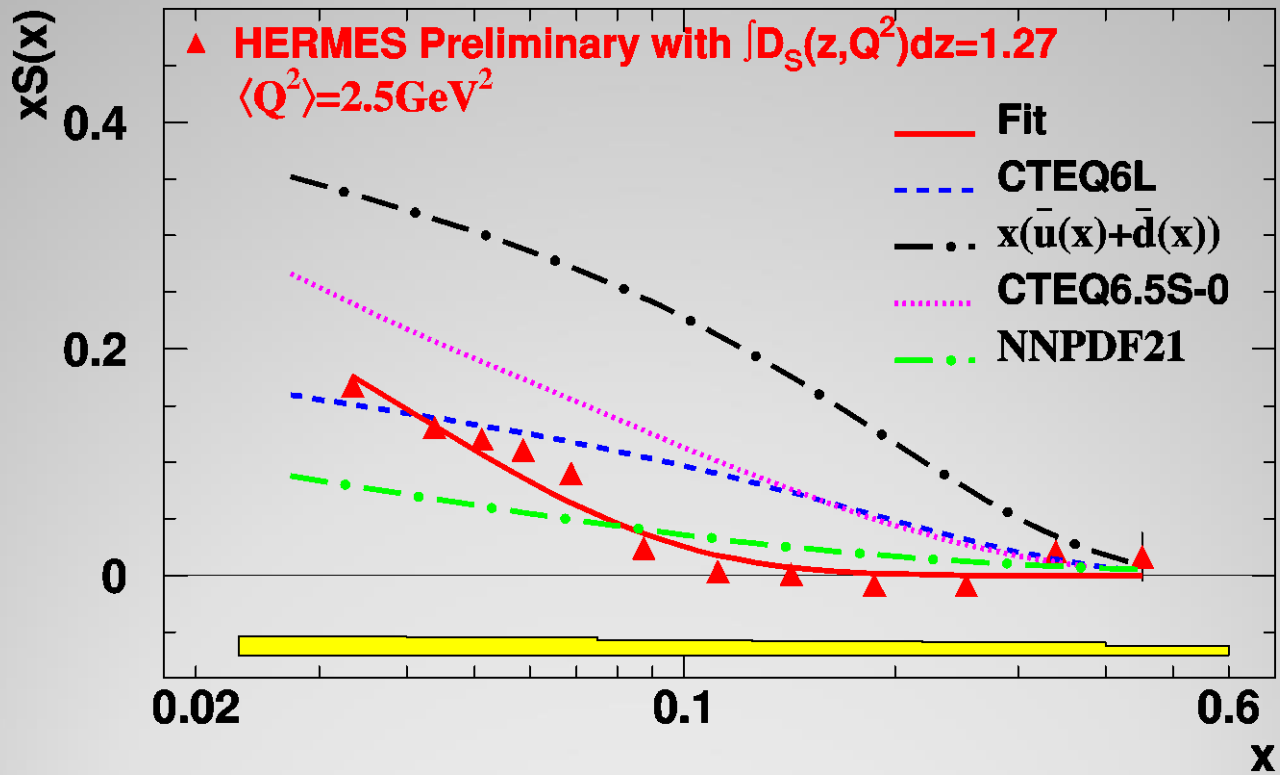


Figure 3: Feynman diagram for the process  $Q_f(\bar{Q}_f)g \rightarrow Z Q_f(\bar{Q}_f)$

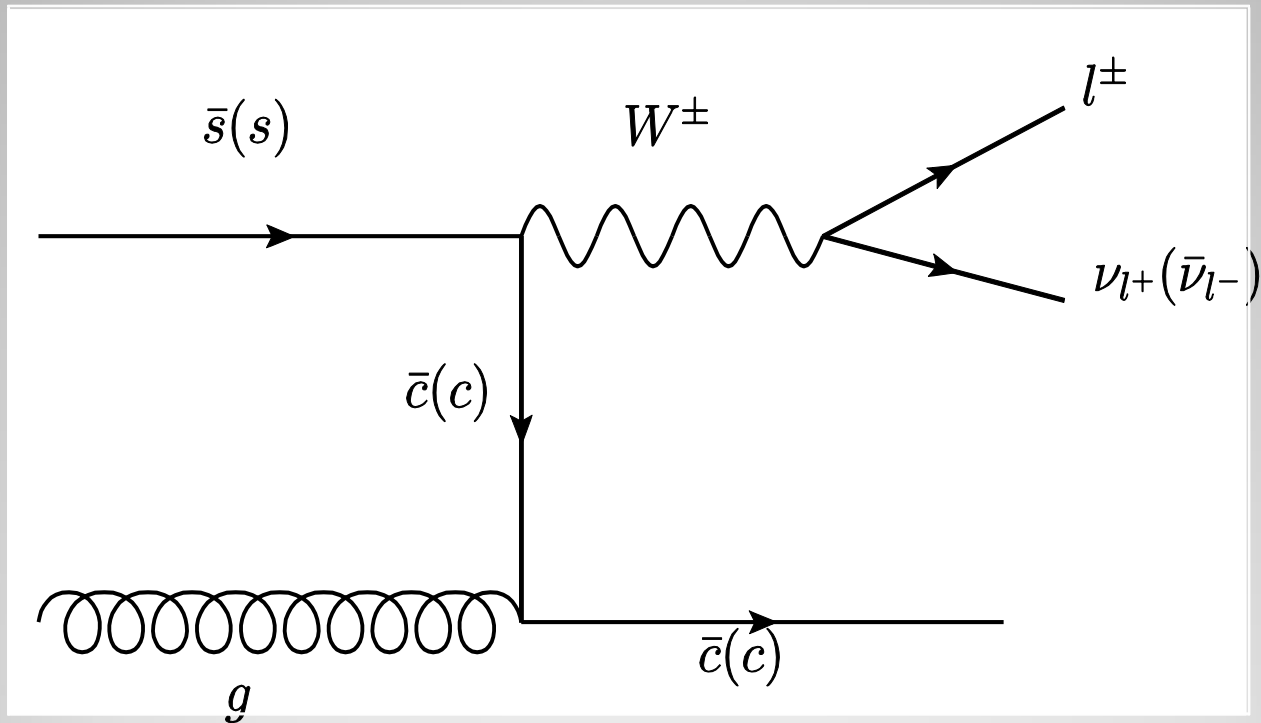
## SUMMARY

1. It is shown that at  $x_q > 0.1$  the contribution of the conventional (**extrinsic**) sea heavy quark distributions is negligibly small in comparison to the **intrinsic** one. It does not contribute to the heavy flavour production in p-p collisions at high energies.
2. The signal of the intrinsic charm (**IC**) and strangeness (**IS**) in proton can be studied in the inclusive open charm and open strangeness production in p-p at the LHC. The **IC** and **IS** signal can be about 200 % -300% at **high  $y$  and  $p_t$**
3. These **intrinsic heavy quark** contributions to the PDF can be studied also in the hard **SM** processes of production of  $\gamma$  and W/Z associated with the heavy flavour **c-** and **b-jets**.
4. The **IC** and **IS contributions** can be about also 250%-300 % at certain values of rapidities and transverse momenta of photons or vector bosons. They can be measured at LHC



# WHAT we are doing now ?

$$pp \rightarrow W + c - jet + X$$





# SEARCH FOR INTRINSIC STRANGENESS IN P-P

$$pp \rightarrow K^{+, -, 0} X$$

At  $x_F = \frac{2p_t}{\sqrt{s}} \sinh(\eta)$  above 0.1

there can be an enhancement due to the **IS**.

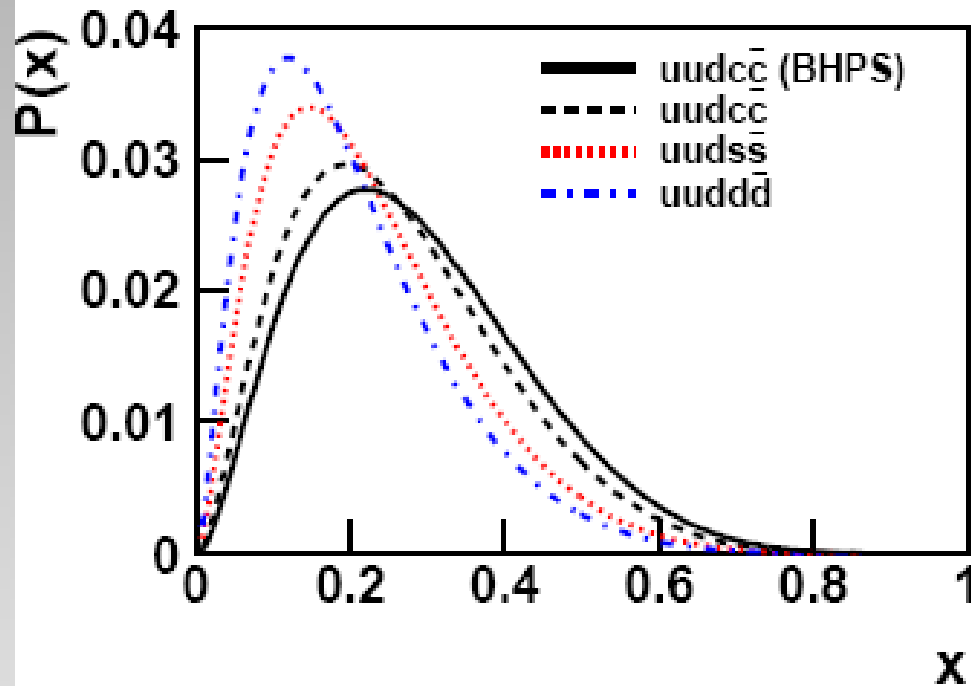
It means that the possible IS signal depend on  $\frac{p_T}{\sqrt{s}}$   
and does not depend on  $\sqrt{s}$

$$K^+ (u \bar{s}); K^- (\bar{u} s)$$

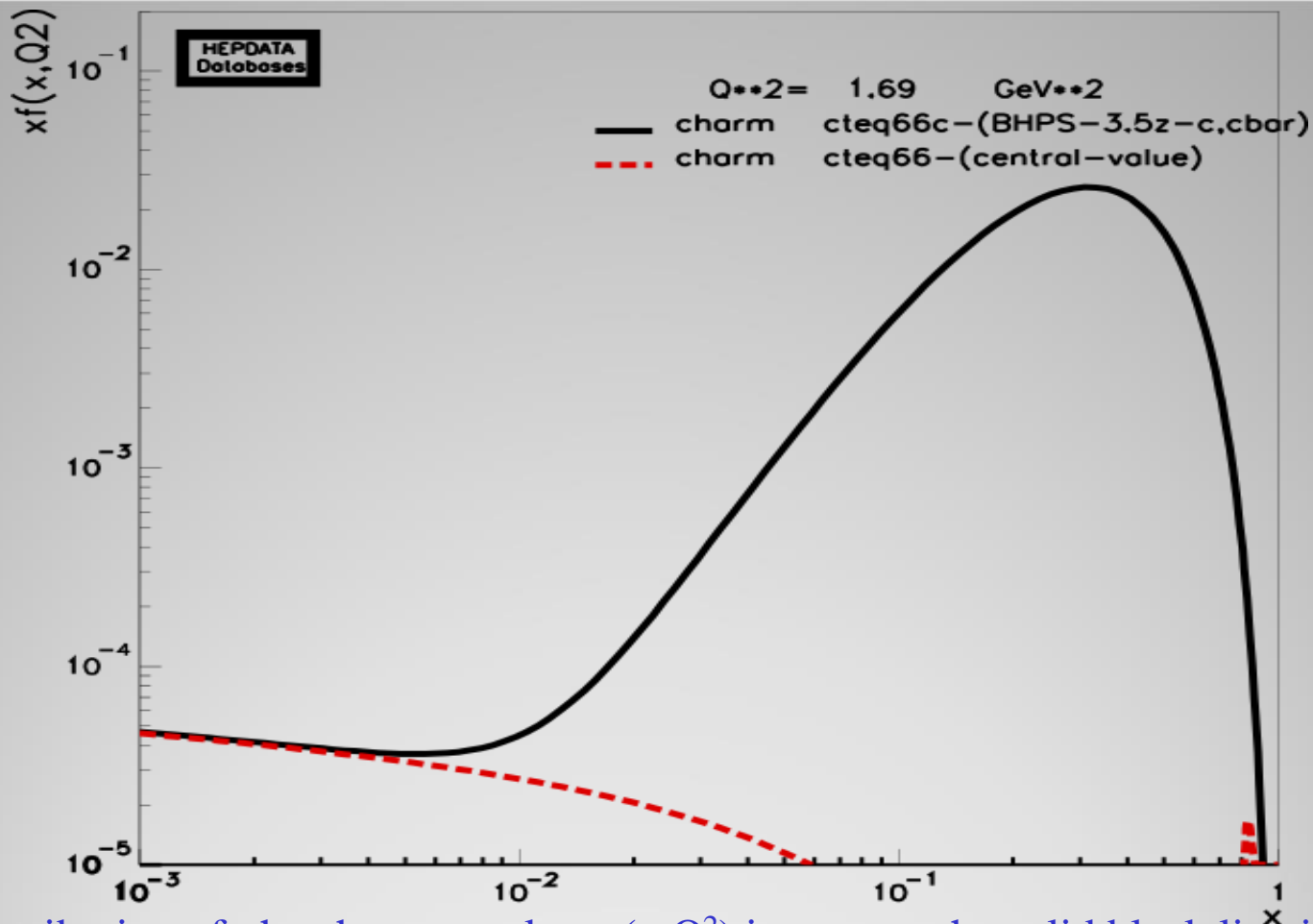
Therefore, it makes the certain sense to  
measure  $K^-$  mesons in p-p collisions at

**NA61, CBM & NICA**

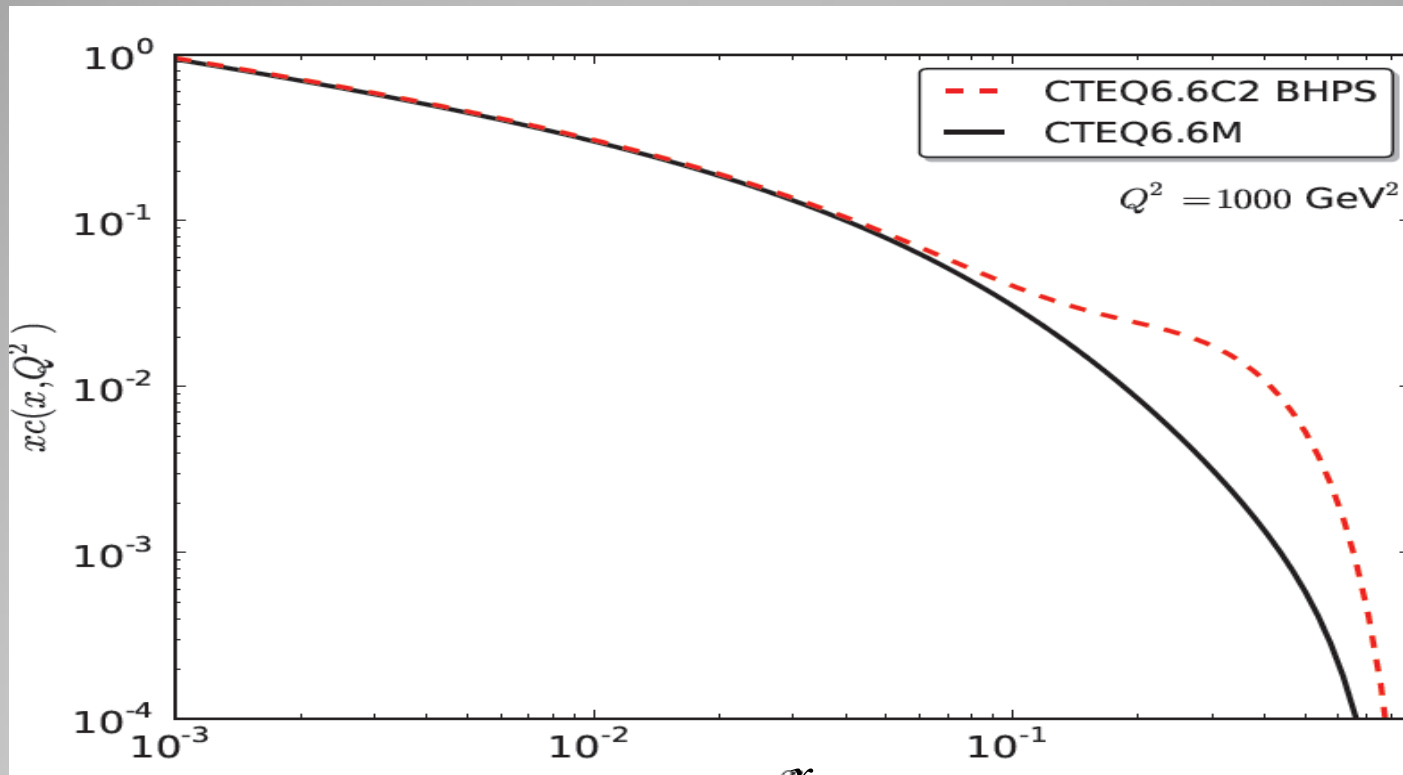
to observe a possible **intrinsic** strangeness in the proton



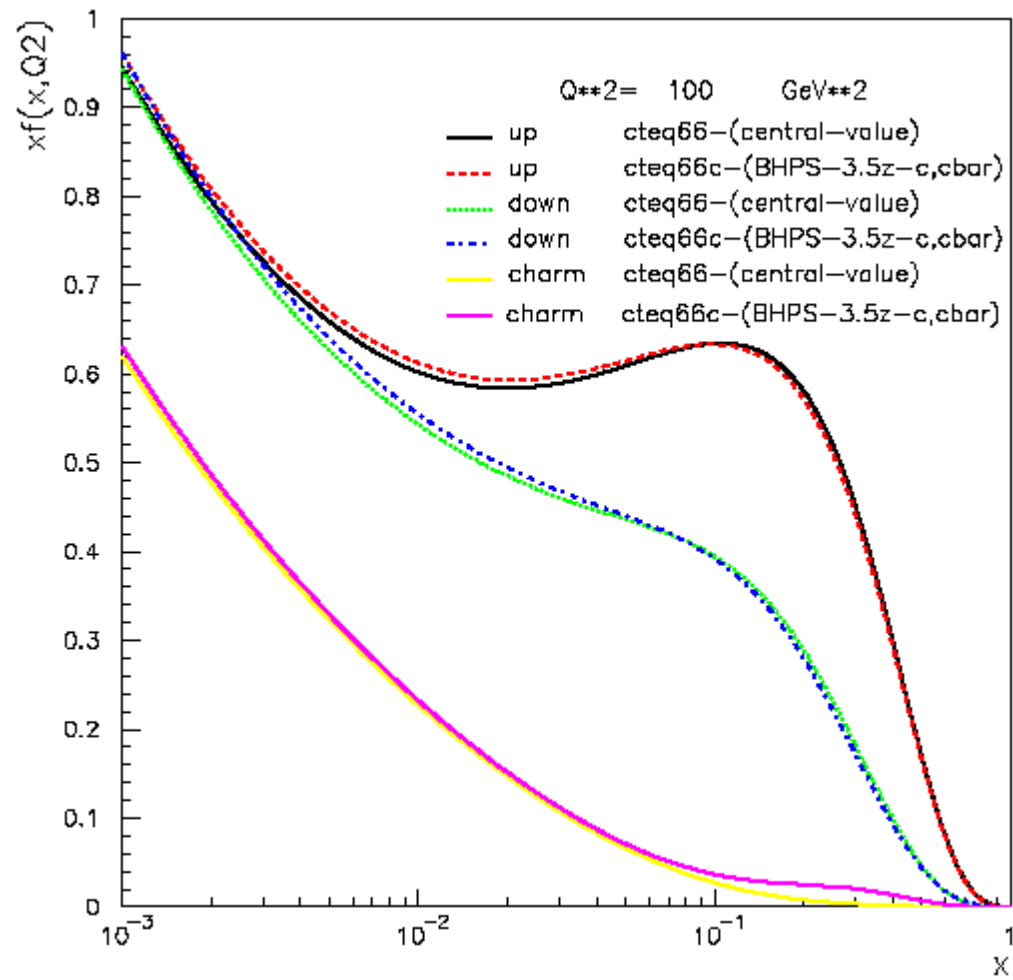
The  $x$ -distribution of the intrinsic  $Q$  calculated within the BHPS model. **There is an enhancement at  $x > 0.1$**   
 Jen-Chieh Peng & We-Chen Chang, hep-ph/1207.2193.

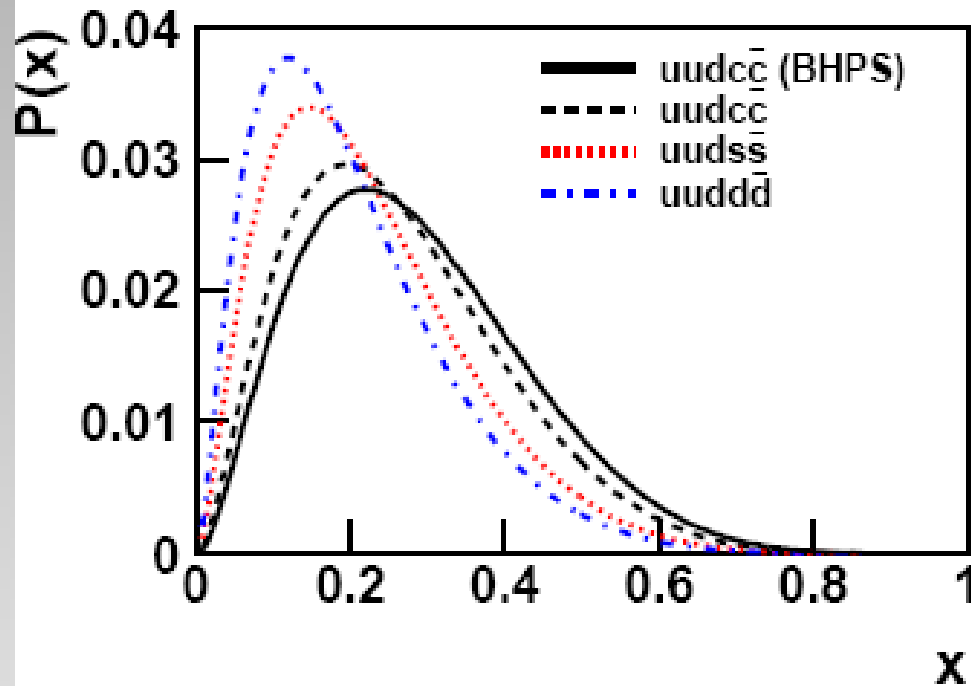


The  $x$ -distribution of the charm quarks  $x c(x, Q^2)$  in proton; the solid black line is the IC contribution with its probability about 3.5 %, the dash green curve is the sea charm quark contribution  $x c_{\text{sea}}(x, Q^2)$  at  $Q^2 = 1.69 \text{ GeV}^2$ . There is enhancement at  $x > 0.1$ .



The  $x$ -distribution of the charm quarks  $x_C(x, Q^2)$  in the proton at  $Q^2=1000$   $\text{GeV}^2$ ; the solid black line is the radiatively generated charm density  $x_{C_{rg}}(x, Q^2)$  distribution only, whereas the dashed curve is the sum of  $x_{C_{rg}}(x, Q^2)$  and the intrinsic charm density  $x_{C_{in}}(x, Q^2)$  with its probability about 3.5%. **There is the sizable enhancement at  $x > 0.1$ .**





The  $x$ -distribution of the intrinsic  $Q$  calculated within the BHPS model. **There is an enhancement at  $x > 0.1$**   
 Jen-Chieh Peng & We-Chen Chang, hep-ph/1207.2193.