

# Direct & Indirect Dark Matter Detection in the NUHM



The  
University  
Of  
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# OUTLINE

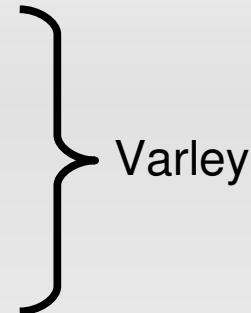
Bayesian theory

Brief introduction to NUHM

Dark matter direct detection

Dark matter indirect detection

Conclusion



Varley



Tsai

# Bayes' Theorem

Probability density

posterior

prior

likelihood

Credit:

Roberto Ruiz

prior      likelihood

$$P(M|D, H) = \frac{P(M|H)P(D|M, H)}{P(D|H)}$$

posterior

evidence

# Priors

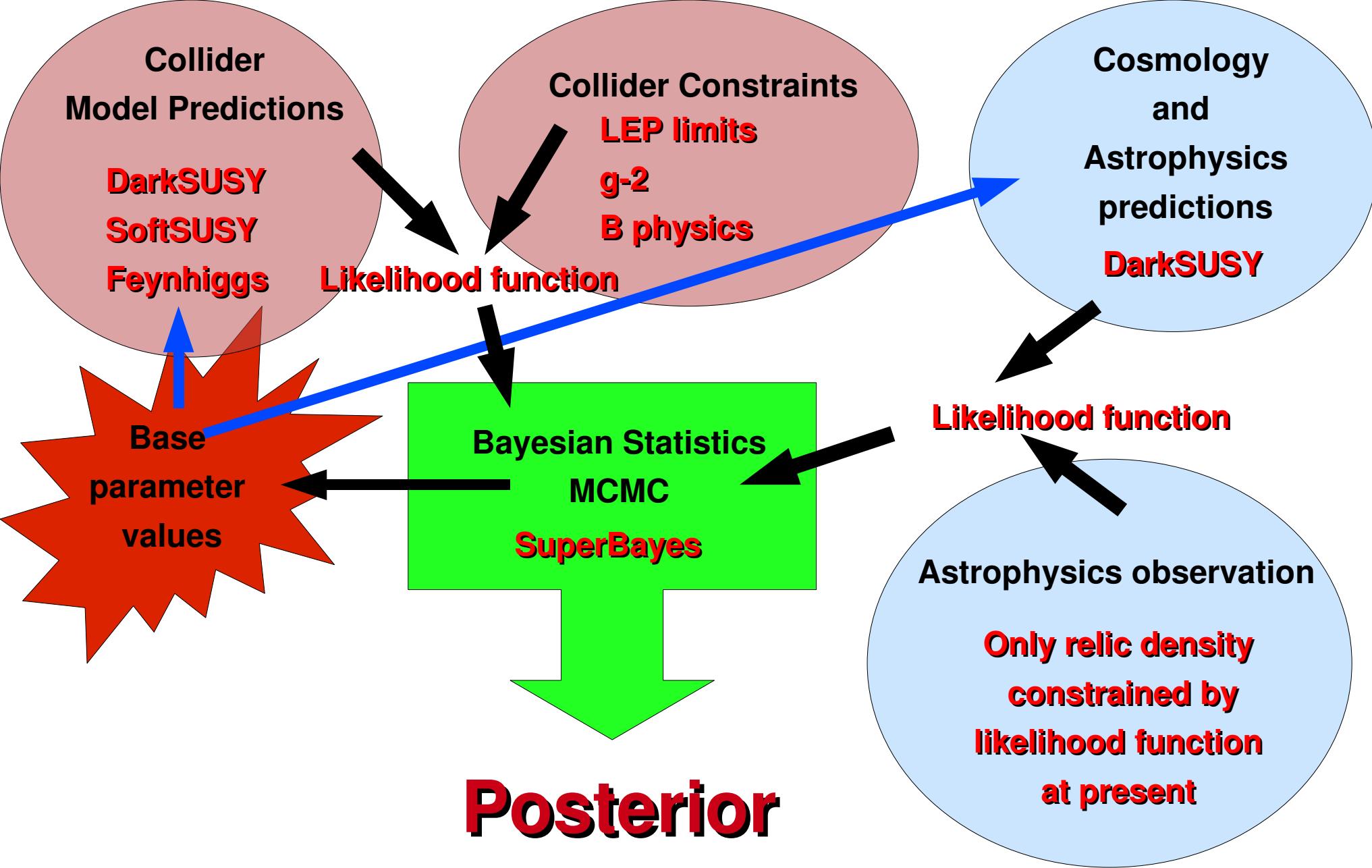
**Flat Prior:**

$$\pi(m_0) = \frac{1}{(m_{max} - m_{min})}$$

**Log Prior:**

$$\pi(m_0) = \frac{1}{m_0 \log[\frac{m_{max}}{m_{min}}]}$$

- Priors arbitrary in Bayesian formalism, so issues of prior dependence and choice of priors important.
- In case of perfect data, our conclusions should be stable regardless of prior.



# NUHM

Many papers (e.g. Berezinsky *et al.*, Ellis *et al.*, Baer *et al.*, Nath & Arnowitt, Cerdeño & Muñoz.....)

- Difference from CMSSM: Allow soft higgs masses to vary independently.
- Can give us a very different picture.
- Parameters:  $(m_0, m_{\frac{1}{2}}, A_0, \tan \beta, m_{H_u}, m_{H_d})$ .
- Log prior applies only to mass parameters.

# NUHM: experimental constraints

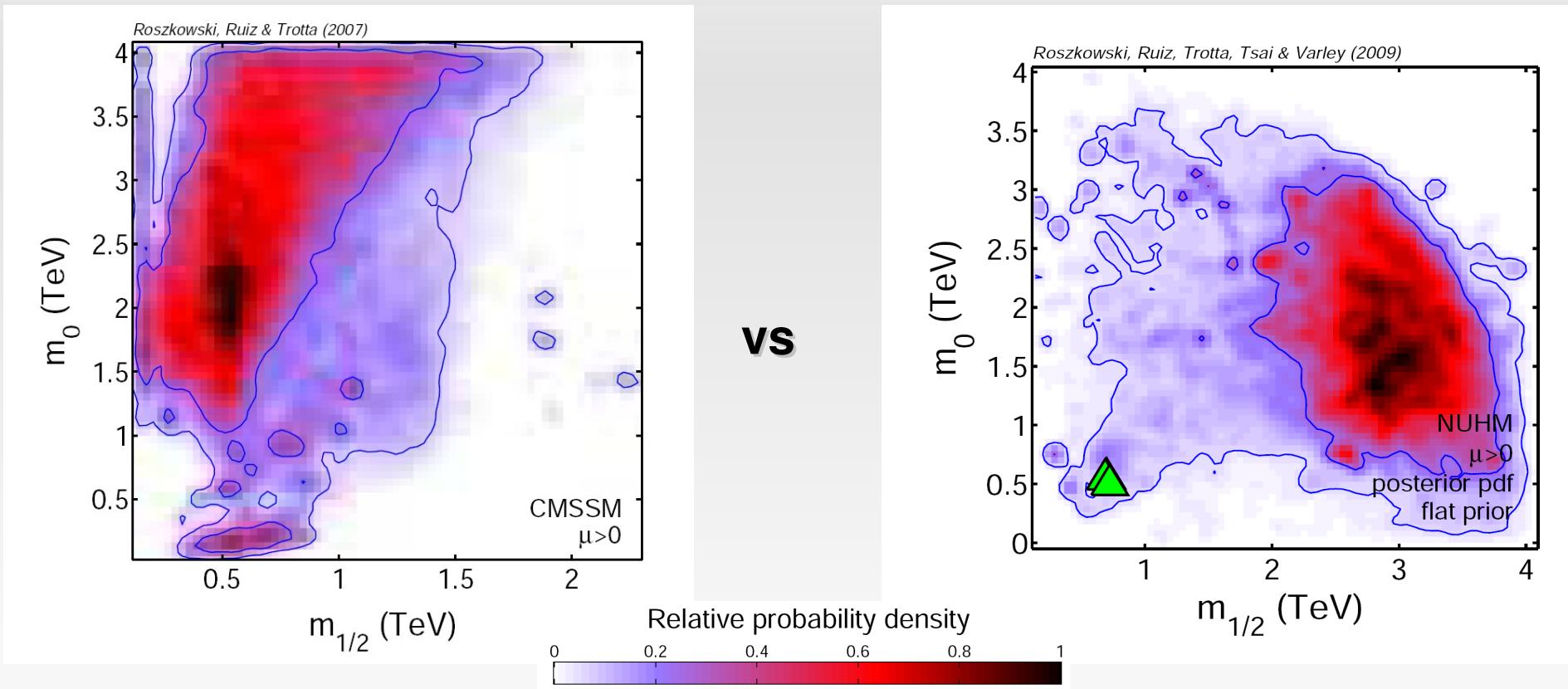
Observable	Mean value $\mu$	Uncertainties		ref.
		$\sigma$ (exper.)	$\tau$ (theor.)	
$M_W$	80.392 GeV	29 MeV	15 MeV	[21]
$\sin^2 \theta_{\text{eff}}$	0.23153	$16 \times 10^{-5}$	$15 \times 10^{-5}$	[21]
$\delta a_\mu^{\text{SUSY}} \times 10^{10}$	27.5	8.4	1	[9]
$BR(\overline{B} \rightarrow X_s \gamma) \times 10^4$	3.55	0.26	0.21	[5]
$\Delta M_{B_s}$	$17.33 \text{ ps}^{-1}$	$0.12 \text{ ps}^{-1}$	$4.8 \text{ ps}^{-1}$	[22]
$\Omega_\chi h^2$	0.109	0.006	$0.1 \Omega_\chi h^2$	[23]
	Limit (95% CL)		$\tau$ (theor.)	ref.
$BR(\overline{B}_s \rightarrow \mu^+ \mu^-)$	$< 5.8 \times 10^{-8}$		14%	[24]
$m_h$	$> 114.4 \text{ GeV}$ (91.0 GeV)		3 GeV	[25]
$\zeta_h^2$	$f(m_h)$		negligible	[25]
sparticle masses	See table 4 in ref. [14].			

# NUHM

**CMSSM**

**flat prior**

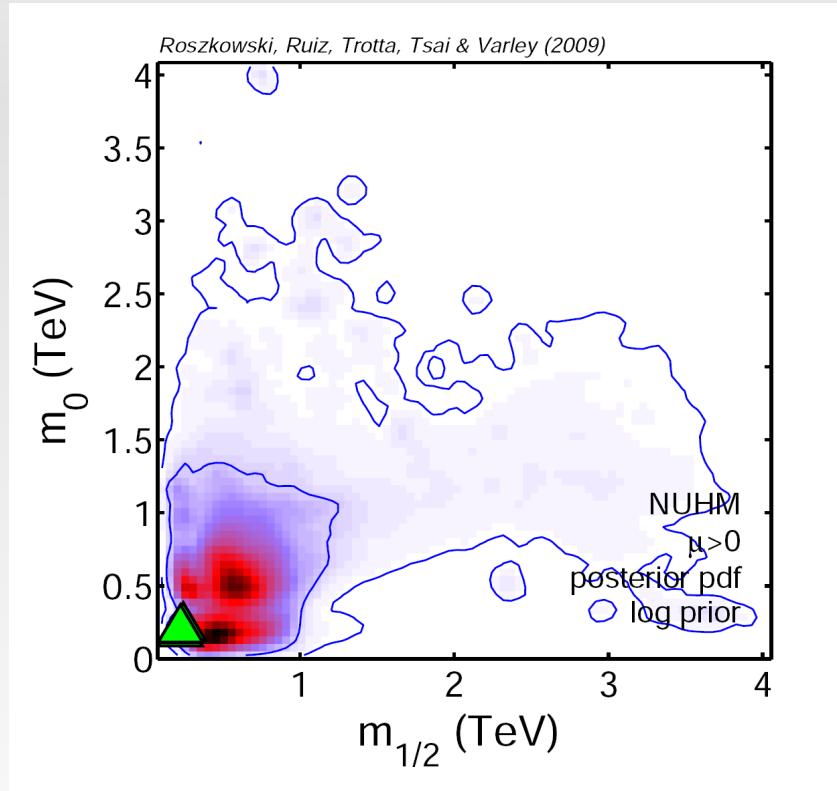
**NUHM**



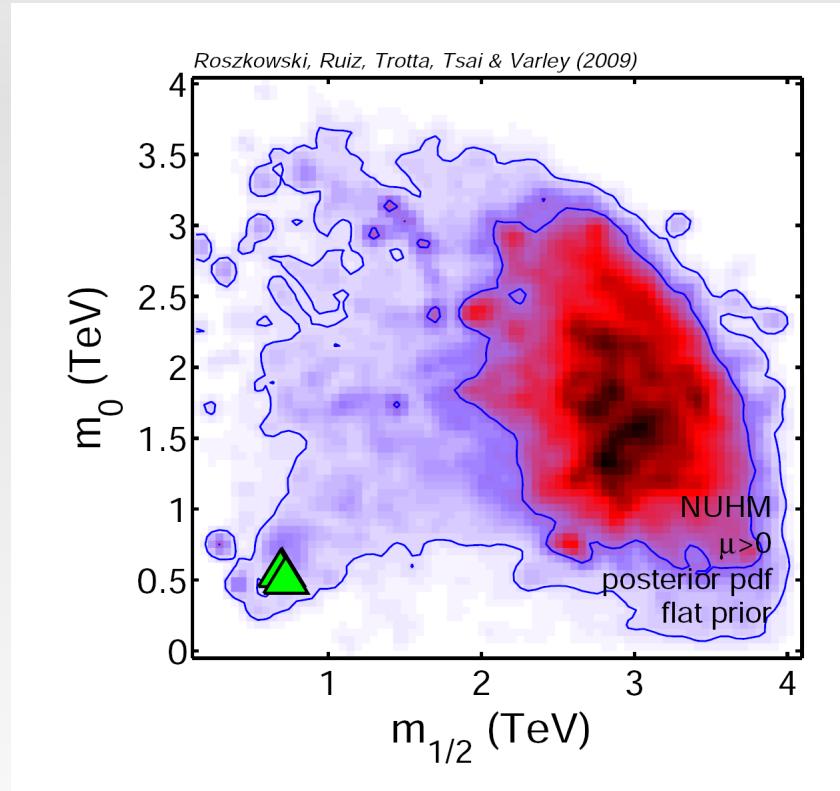
- NUHM gives us a very different picture here.

# NUHM

## Log prior



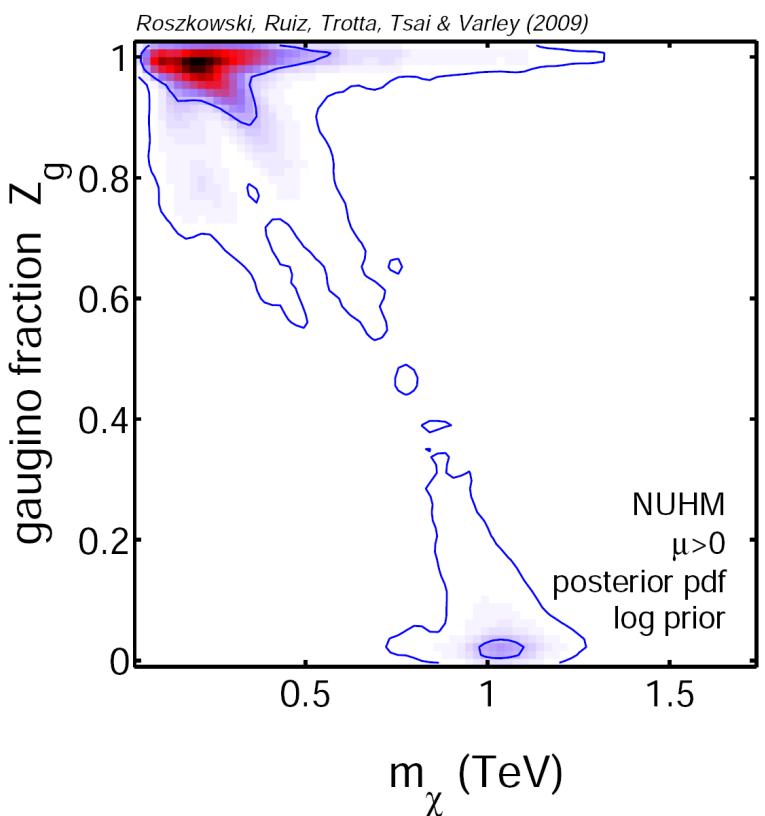
## Flat prior



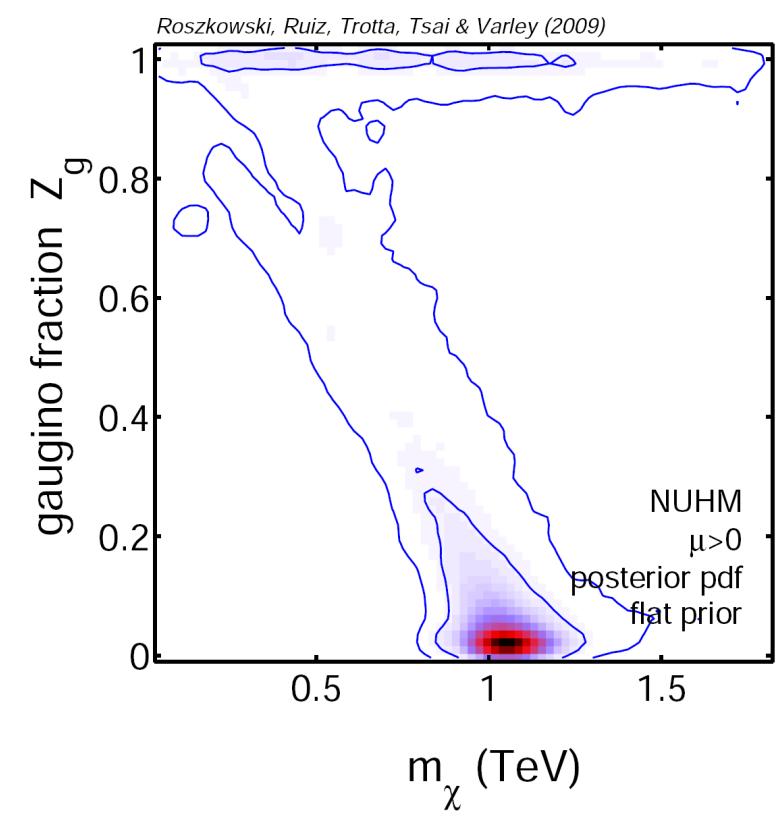
- Prior dependence indicates the model is underconstrained.

# Neutralino masses in NUHM

Log prior



Flat prior



- Can have a heavy higgsino LSP at around 1 TeV but note prior dependence here.

# Direct Detection

## Spin-independent component

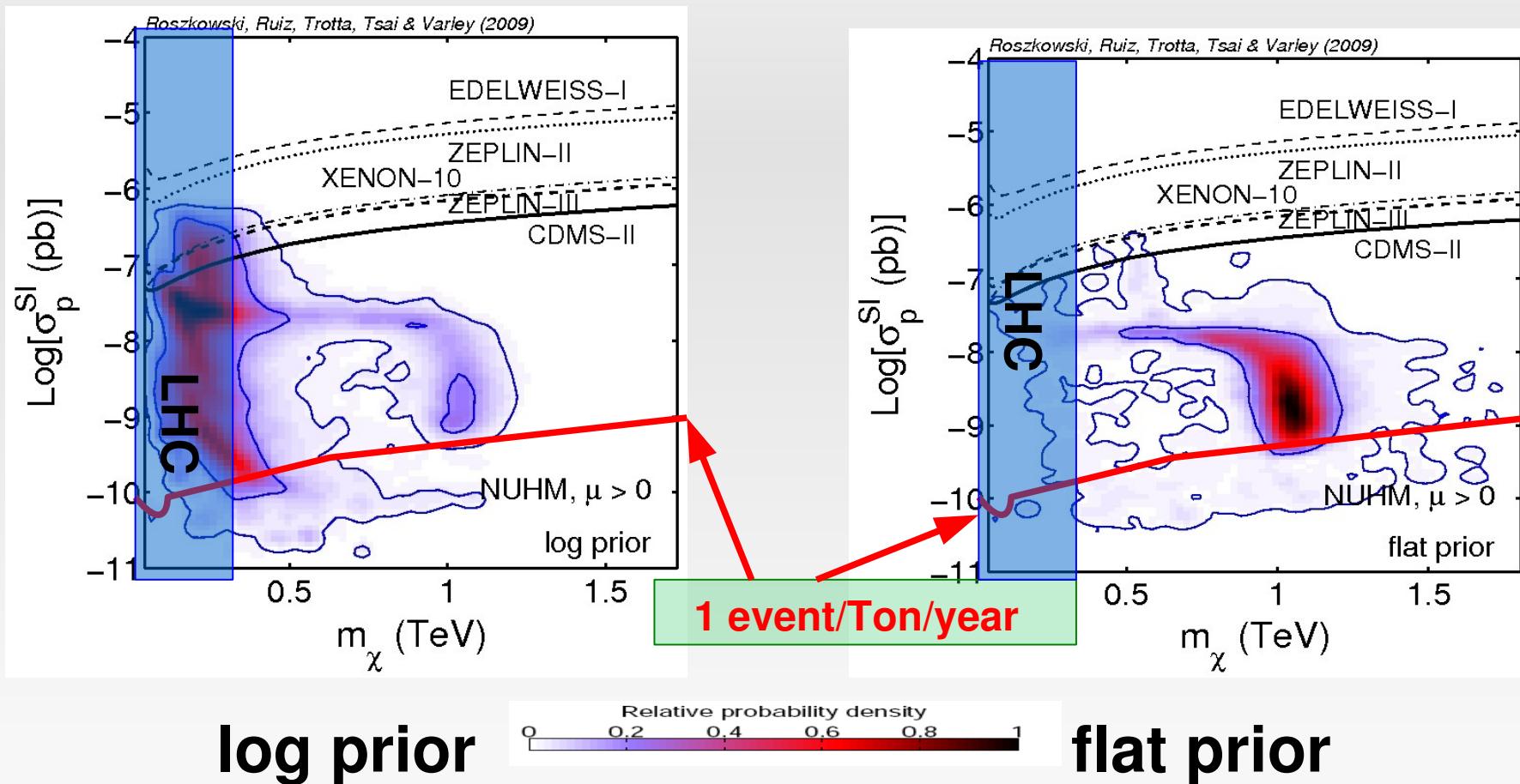
$$\frac{d\sigma^{SI}}{d|\mathbf{q}|} = \frac{1}{\pi v^2} [Z f_p + (A - Z) f_n]^2 F^2(|\mathbf{q}|) \approx \frac{1}{\pi v^2} A^2 f_p^2 F^2(|\mathbf{q}|)$$

## Spin-dependent component

$$\frac{d\sigma^{SD}}{d|\mathbf{q}|} = \frac{8G_F^2}{\pi v^2} \Lambda^2 J(J+1) F^2(|\mathbf{q}|)$$

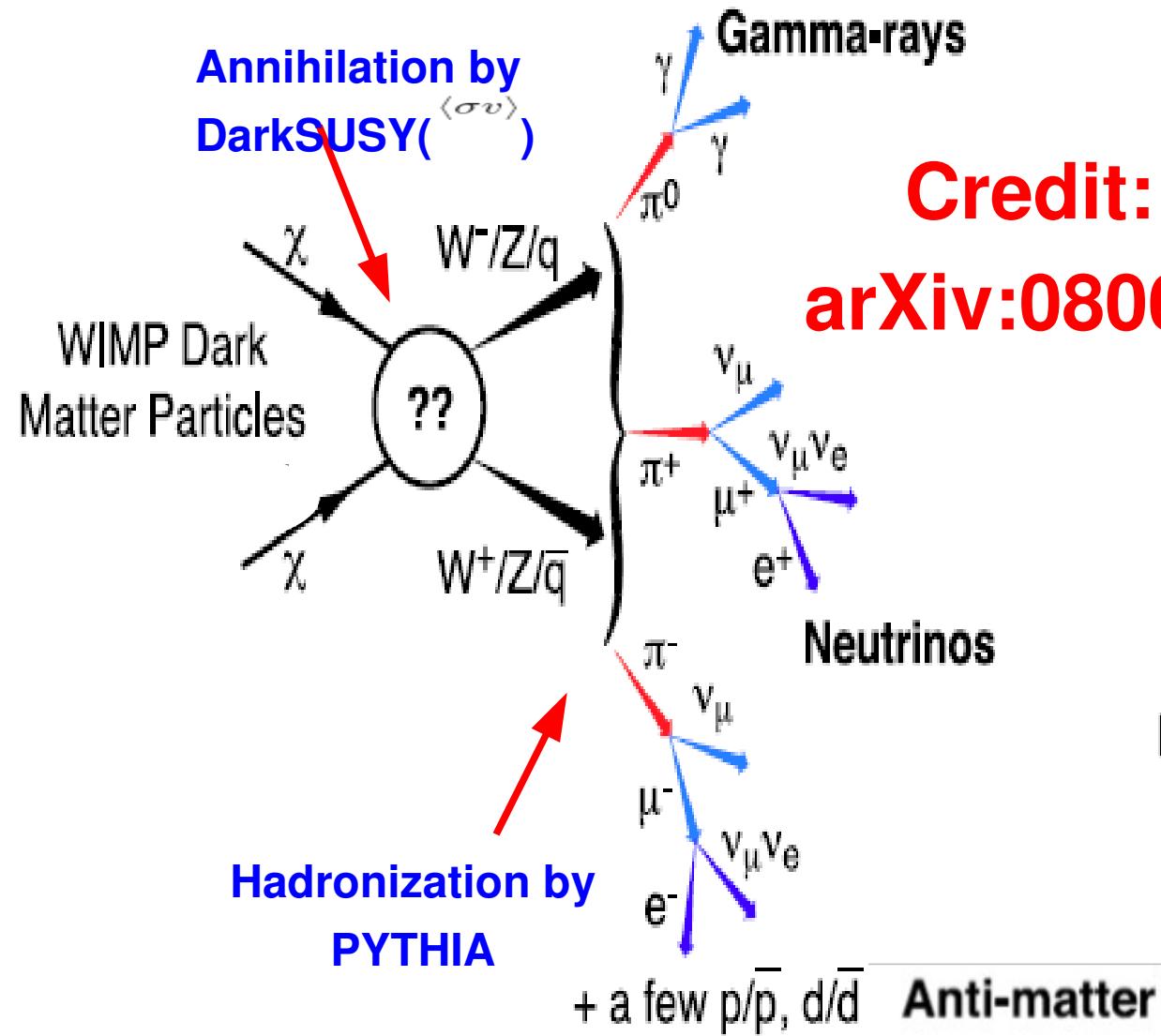
G. Jungman, M. Kamionkowski and K. Griest, “Supersymmetric dark matter,” Phys. Rept. **267**, 195 (1996) [arXiv:hep-ph/9506380].

# The spin-independent cross section

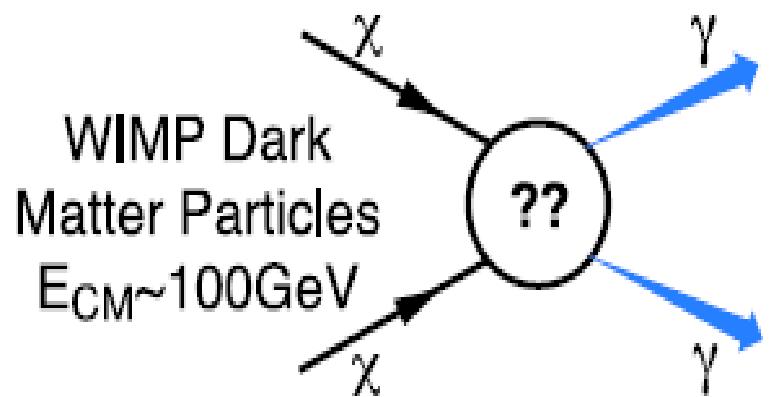


Limit on cross section not very dependent on prior.

# Indirect-detection



Credit: E. A. Baltz et al.  
arXiv:0806.2911v2 [astro-ph]



# Anti-matter or Gamma rays from GC

Dark matter Halo Model:

Klypin et al. (2001)

NFW

Isothermal cored

Galactic Center

Solar system

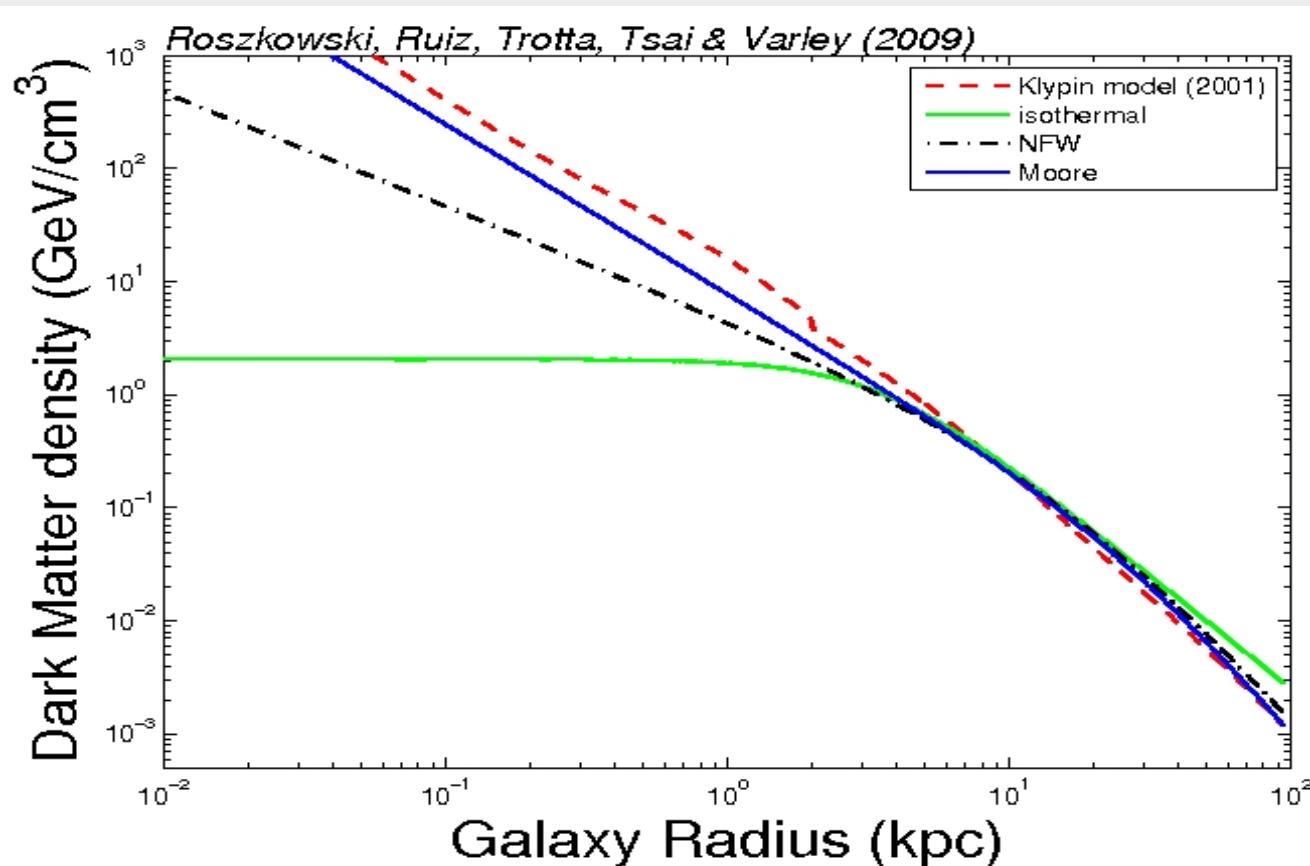
Propagation model:

Same as DarkSUSY

Edward A. Baltz and Joakim Edsjo  
astro-ph/9808243v1

$$\frac{\partial}{\partial t} \frac{dn}{d\varepsilon} = \vec{\nabla} \cdot \left[ K(\varepsilon, \vec{x}) \vec{\nabla} \frac{an}{d\varepsilon} \right] + \frac{\sigma}{\partial \varepsilon} \left[ b(\varepsilon, \vec{x}) \frac{an}{d\varepsilon} \right] + Q(\varepsilon, \vec{x}),$$

# Klypin et al. Halo Model (2001)

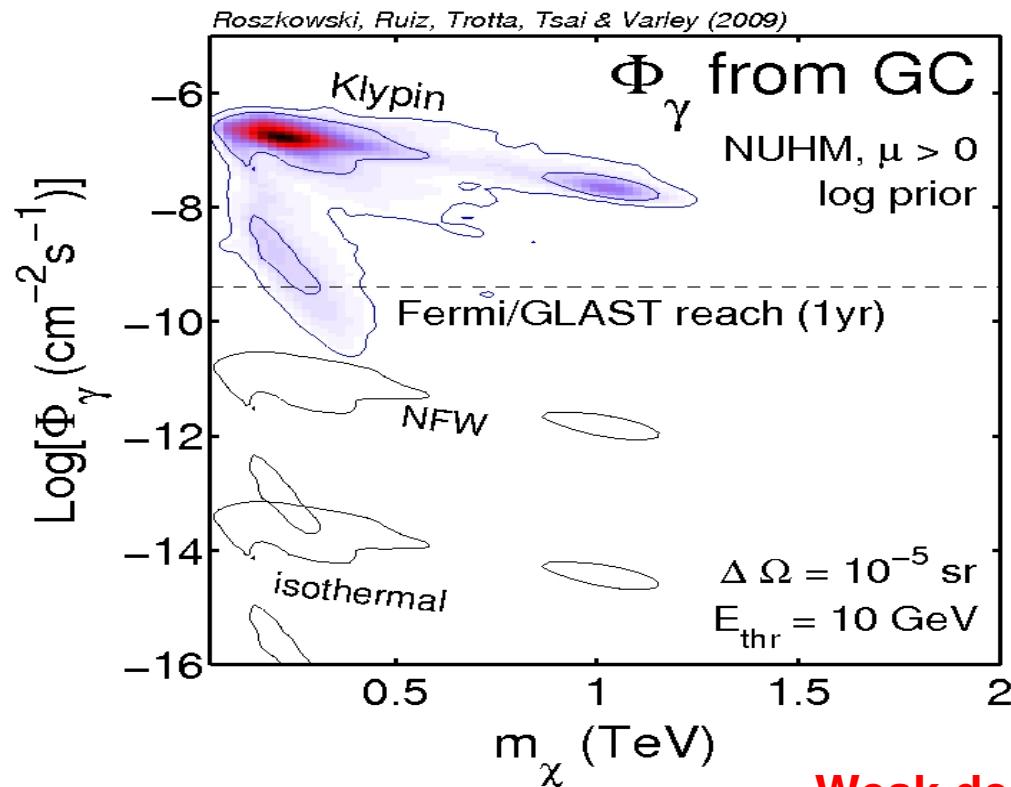


$M_{vir} (\text{M}_\odot)$	$M_{disk} (\text{M}_\odot)$	$M_{bulge} (\text{M}_\odot)$	$C$	$r_0 (\text{kpc})$	$r_s (\text{kpc})$
$8.57 \times 10^{11}$	$4.2 \times 10^{10}$	$8 \times 10^9$	11	8.0	22.27
$r_{vir} (\text{kpc})$	$r_{disk} (\text{kpc})$	$r_{bulge} (\text{kpc})$	$\gamma$	$\rho_0 (\text{GeV}/\text{cm}^3)$	$J (10^{-5} \text{sr})$
244.93	3.5	0.5	1.8	0.3	$1.95 \times 10^8$

# The gamma-ray from GC

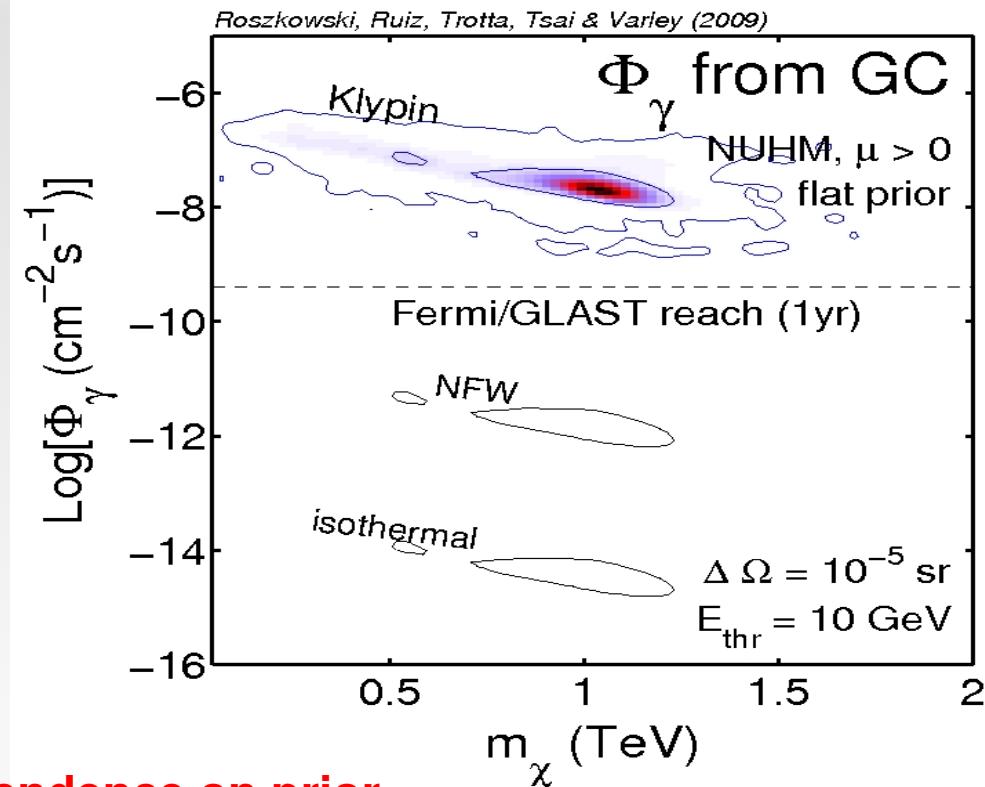
$$\frac{d\Phi_\gamma}{dE_\gamma}(E_\gamma, \psi) = \sum_i \frac{\sigma_i v}{8\pi m_\chi^2} \frac{dN_\gamma^i}{dE_\gamma} \int_{\text{l.o.s.}} dl \rho_\chi^2(r(l, \psi)).$$

$$\Phi_\gamma(\Delta\Omega) = \int_{E_{th}}^{m_\chi} \frac{d\Phi_\gamma}{dE_\gamma} dE_\gamma.$$



**Weak dependence on prior,  
strong dependence on halo model**

**log prior**

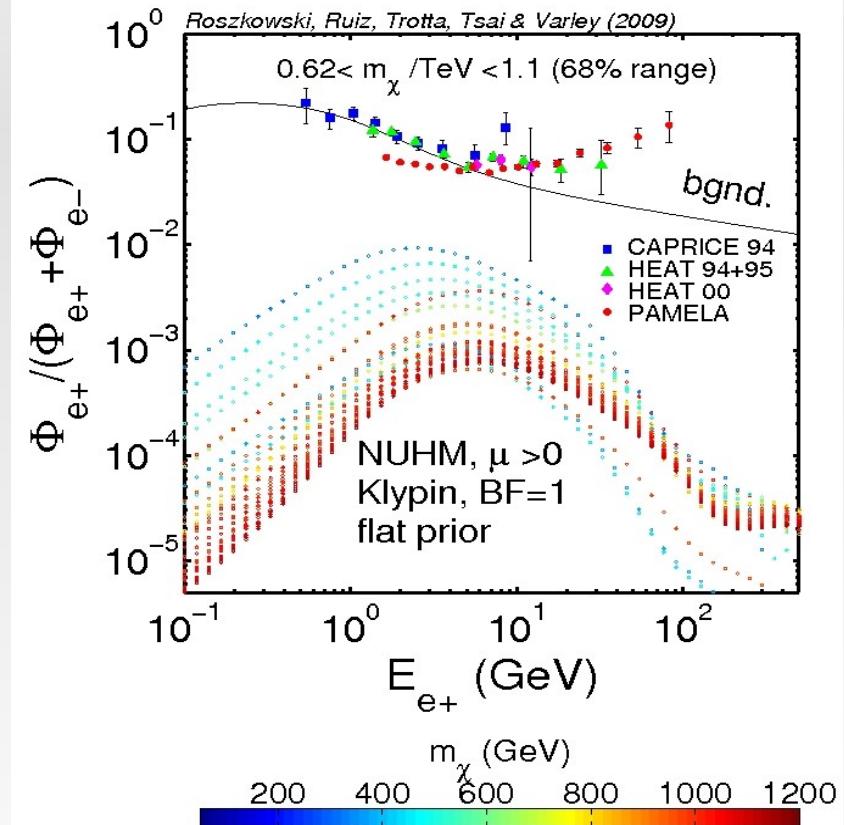
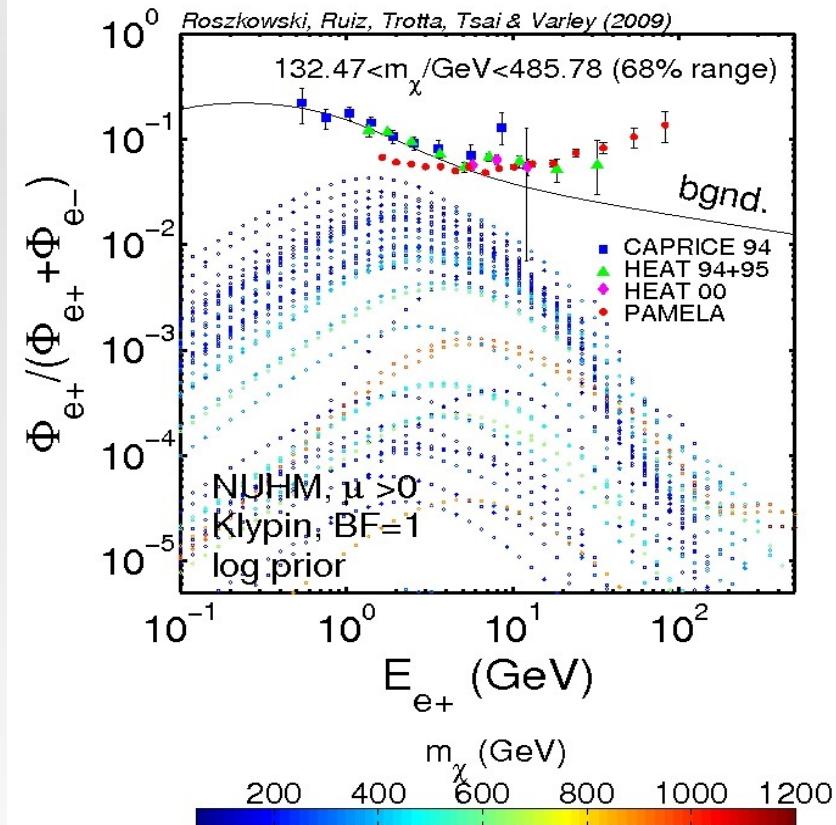


**Weak dependence on prior,  
strong dependence on halo model**

**flat prior**

# Positron fraction

Example: Klypin et al. (2001) halo model (NFW and iso are similiar.)



Can't produce PAMELA excess!!

log prior

flat prior

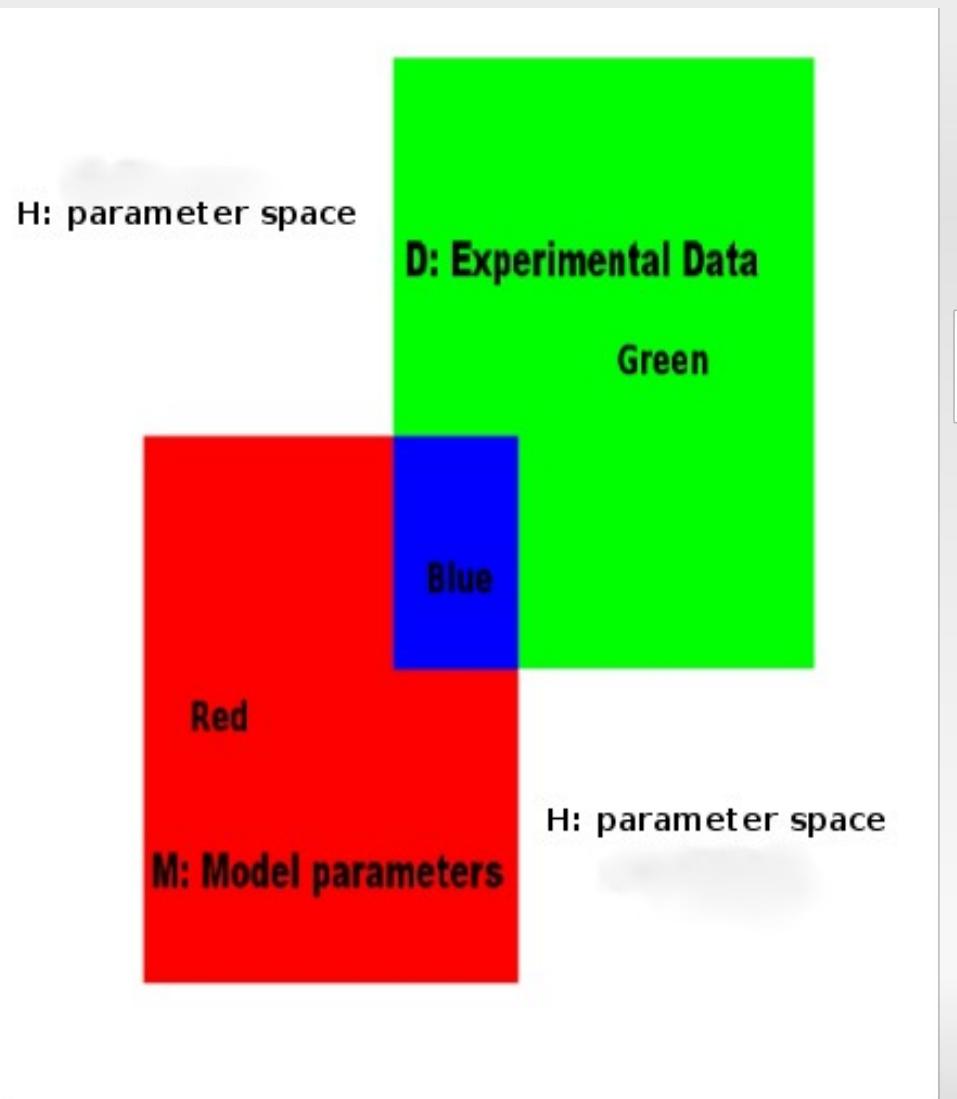
# Conclusion

- NUHM an interesting alternative to CMSSM. Analysed with flat and log prior in Bayesian formalism.
- As experimental uncertainties are large, the spin-independent cross section cannot offer a very effective constraint on  $m_0$  vs  $m_{1/2}$  plane.
- Prospects for FERMI strongly depend on how cuspy the halo model is in galactic center.
- No reasonable halo model can reproduce PAMELA excess here.

**Thank you very much for your attention.**

# **BACKUP MATERIAL**

# The Bayesian theory



$$P(M|H)P(D|M, H) = P(D|H)P(M|D, H)$$

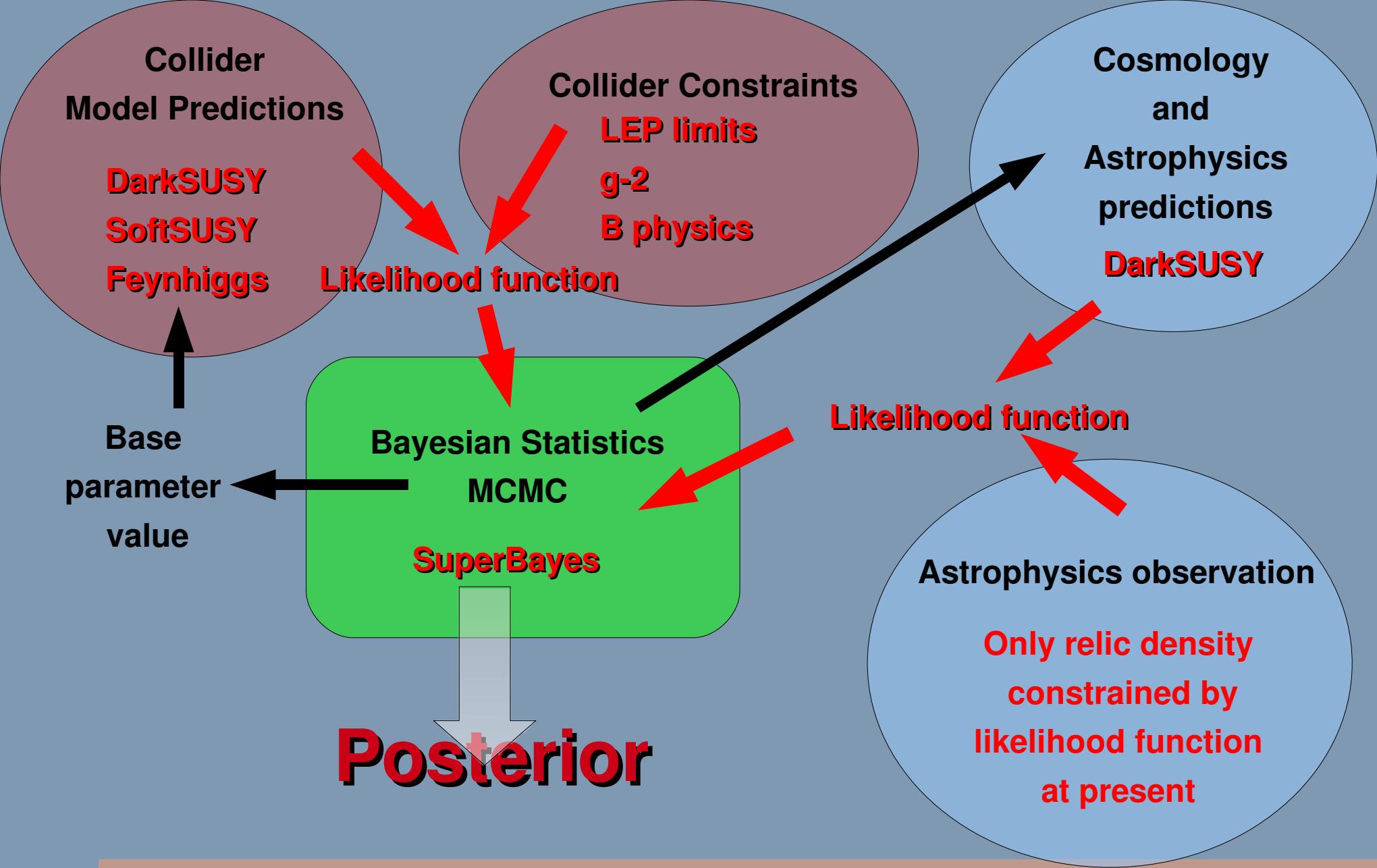
$$P(M|D, H) = \frac{P(M|H)P(D|M, H)}{P(D|H)}$$

# Markov Chain Monte Carlo

- Want to obtain posterior, how do we do this?
- Take random walk through parameter space.
- Probability of accepting next point depends on relative posteriors of both; Metropolis-Hastings algorithm. For two points  $x$  and  $y$ , accept with probability:

$$P_{x \rightarrow y} = \min(p(y)/p(x), 1)$$

- Iterate: as  $t$  goes to infinity recover posterior we are interested in studying.



# The Neutralino

$$\frac{1}{2} \begin{pmatrix} \tilde{B} & \tilde{W}^0 & \tilde{H}_d^0 & \tilde{H}_u^0 \end{pmatrix} \begin{pmatrix} M_1 & 0 & \frac{-g'v_d}{\sqrt{2}} & \frac{g'v_u}{\sqrt{2}} \\ 0 & M_2 & \frac{gv_d}{\sqrt{2}} & \frac{-gv_u}{\sqrt{2}} \\ \frac{-g'v_d}{\sqrt{2}} & \frac{gv_d}{\sqrt{2}} & 0 & -\mu \\ \frac{g'v_u}{\sqrt{2}} & \frac{-gv_u}{\sqrt{2}} & -\mu & 0 \end{pmatrix} \begin{pmatrix} \tilde{B} \\ \tilde{W}^0 \\ \tilde{H}_d^0 \\ \tilde{H}_u^0 \end{pmatrix}$$

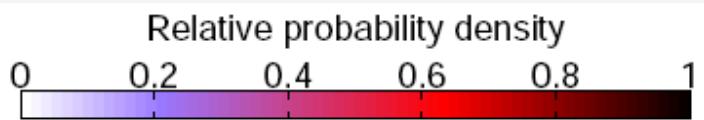
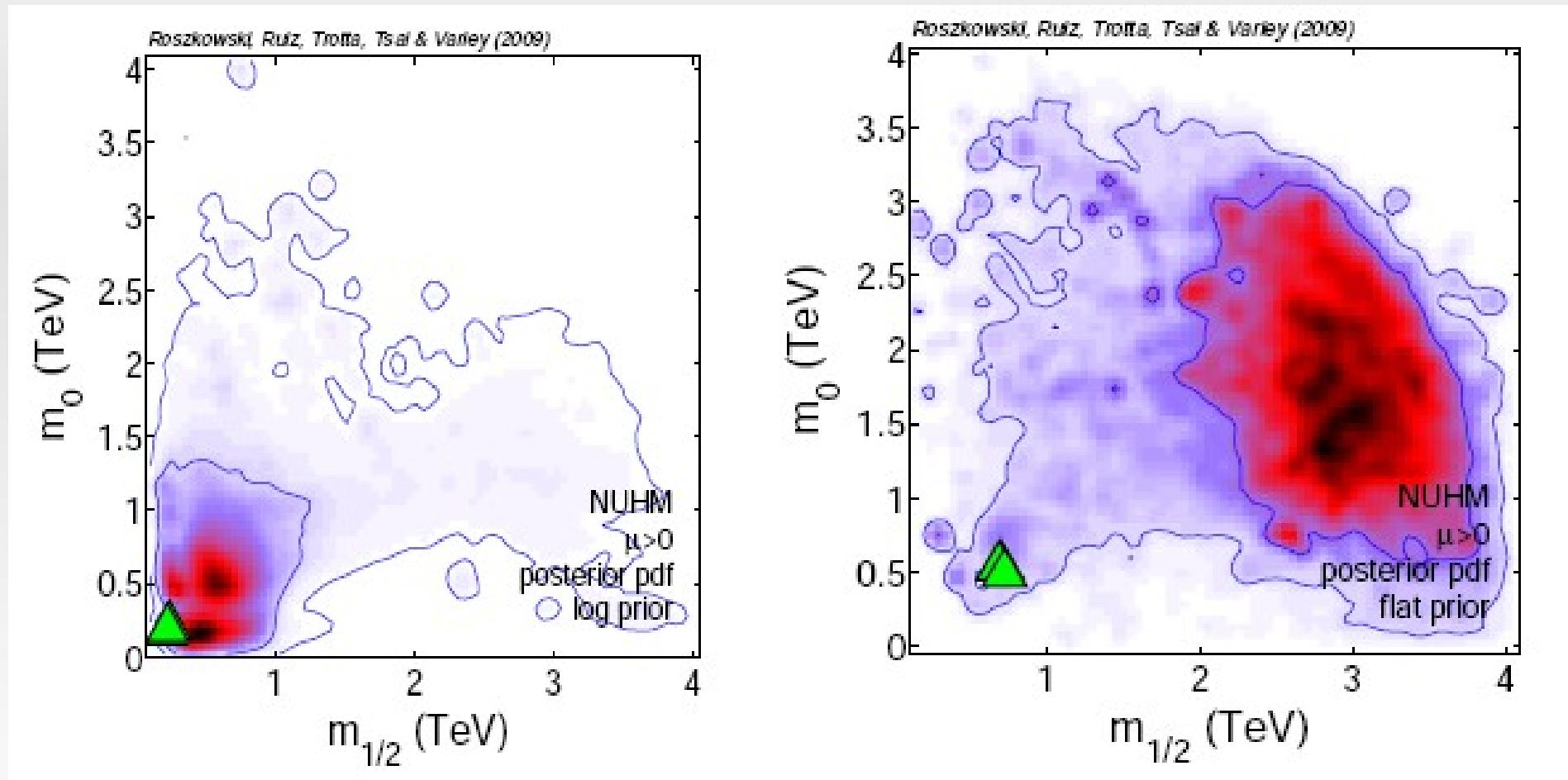
LSP

$$\begin{pmatrix} m_{\chi_1^0} & 0 & 0 & 0 \\ 0 & m_{\chi_2^0} & 0 & 0 \\ 0 & 0 & m_{\chi_3^0} & 0 \\ 0 & 0 & 0 & m_{\chi_4^0} \end{pmatrix}$$

higgsino LSP  $\mu < M_1$

$$R \equiv (-1)^{3(B-L)+2S}$$

# Some plots - effect of prior

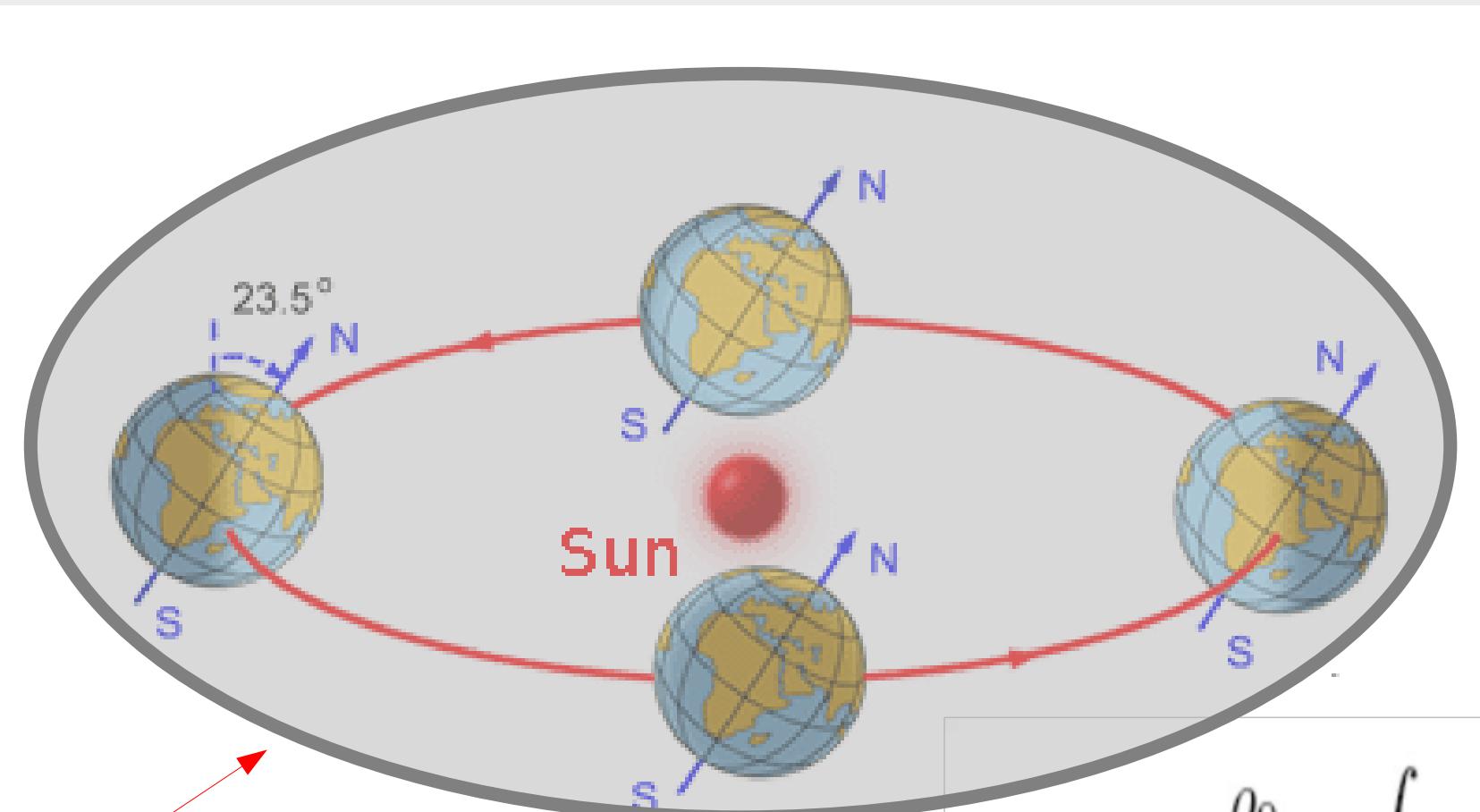


# More NUHM

Berezinsky *et al.*, Ellis *et al.*, Baer *et al.*, Nath & Arnowitt, Cerdeño & Muñoz.....

- Difference from CMSSM: Allow soft higgs masses to vary independently.
- Parameters:  $(m_0, m_{\frac{1}{2}}, A_0, \tan\beta, m_{H_u}, m_{H_d}).$
- Can also be parametrised as having  $\mu$  and the pseudoscalar mass as free parameters.
- S parameter- feature of RGEs, non-zero here.
- New features: A funnel moves around, varied NLSP and coannihilation.
- Can be a strong higgsino component, cross over region can be important.

# Direct Detection



Dark matter local halo

$$dR = \frac{\rho_0}{m_\chi m_N} \int v f(v) d\sigma dv$$

# The neutralino-nucleon effective lagrangian

$$\mathcal{L} = \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q}_i \gamma_\mu (\alpha_{1i} + \alpha_{2i} \gamma^5) q_i$$

$$+ \alpha_{3i} \bar{\chi} \chi \bar{q}_i q_i + \alpha_{4i} \bar{\chi} \gamma^5 \chi \bar{q}_i \gamma^5 q_i$$

$$+ \alpha_{5i} \bar{\chi} \chi \bar{q}_i \gamma^5 q_i + \alpha_{6i} \bar{\chi} \gamma^5 \chi \bar{q}_i q_i$$

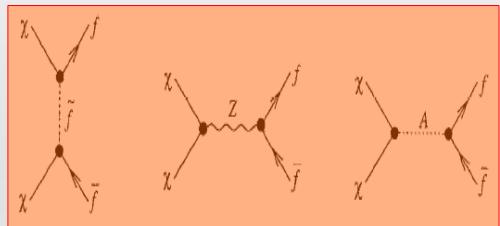
SI

SD

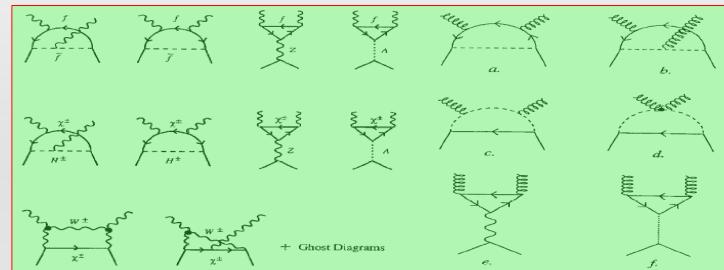
Velocity dependent elastic cross section



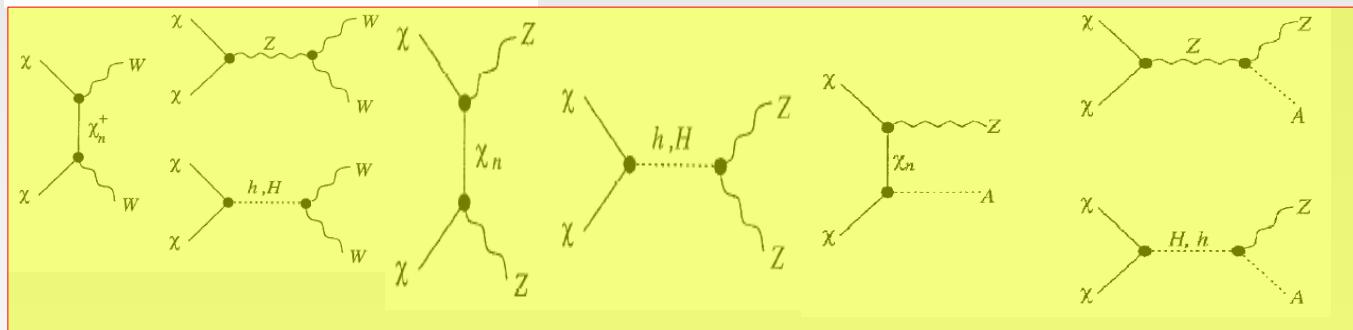
# Contributions to Neutralino Annihilation



$(e^+ e^-, \mu^+ \mu^-, \tau^+ \tau^-, q\bar{q})$



$(gg, \gamma\gamma, Z^0\gamma)$

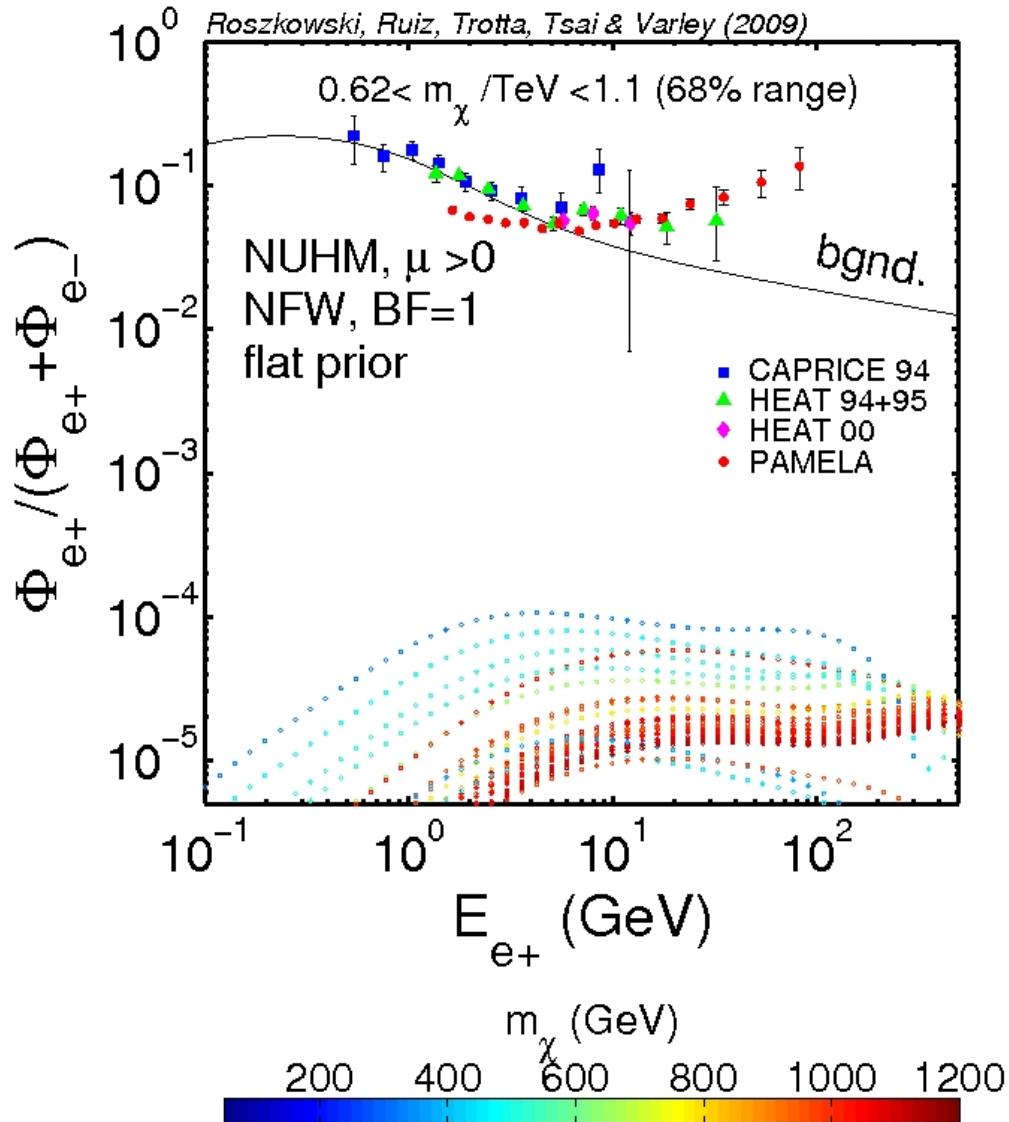
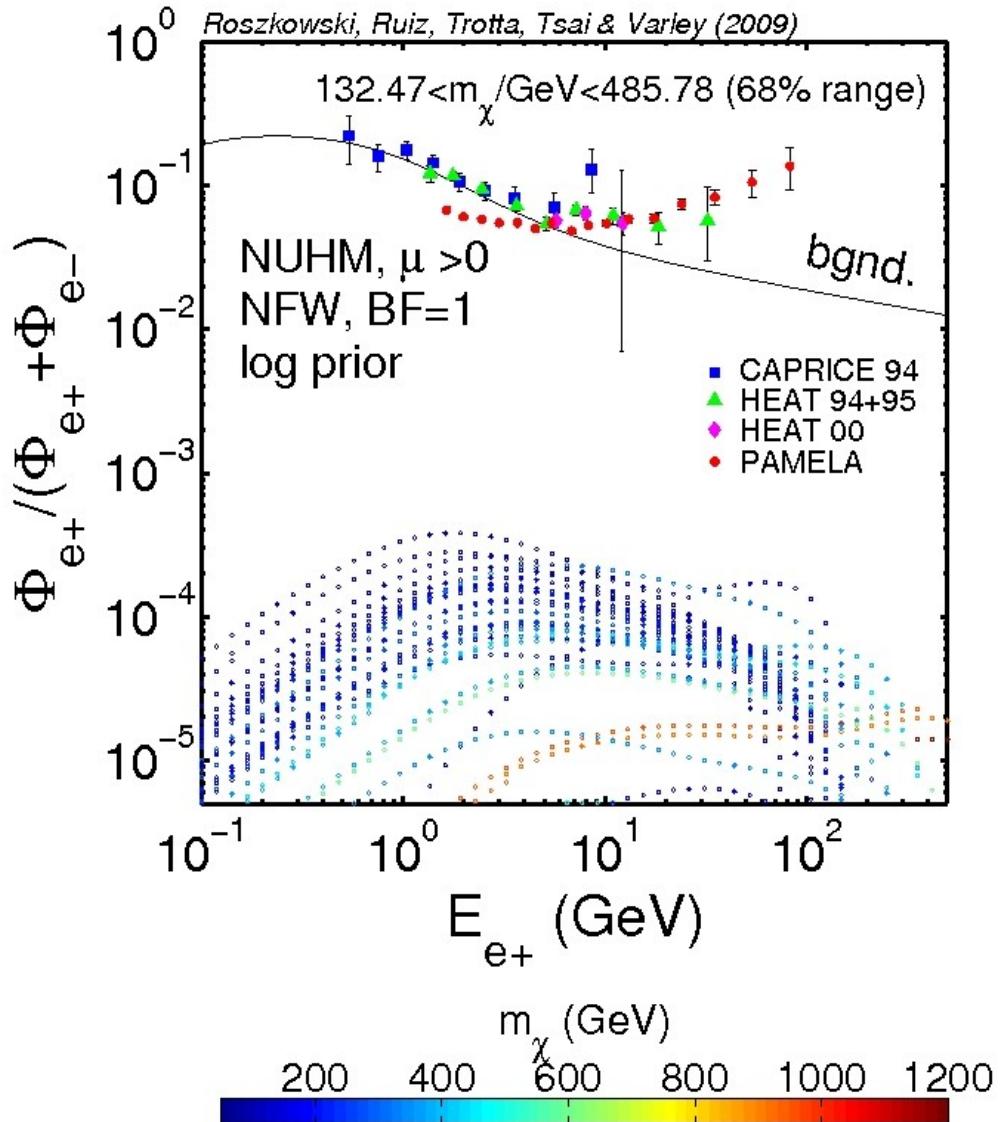


All the  $\sigma v$  are considered under the limit:

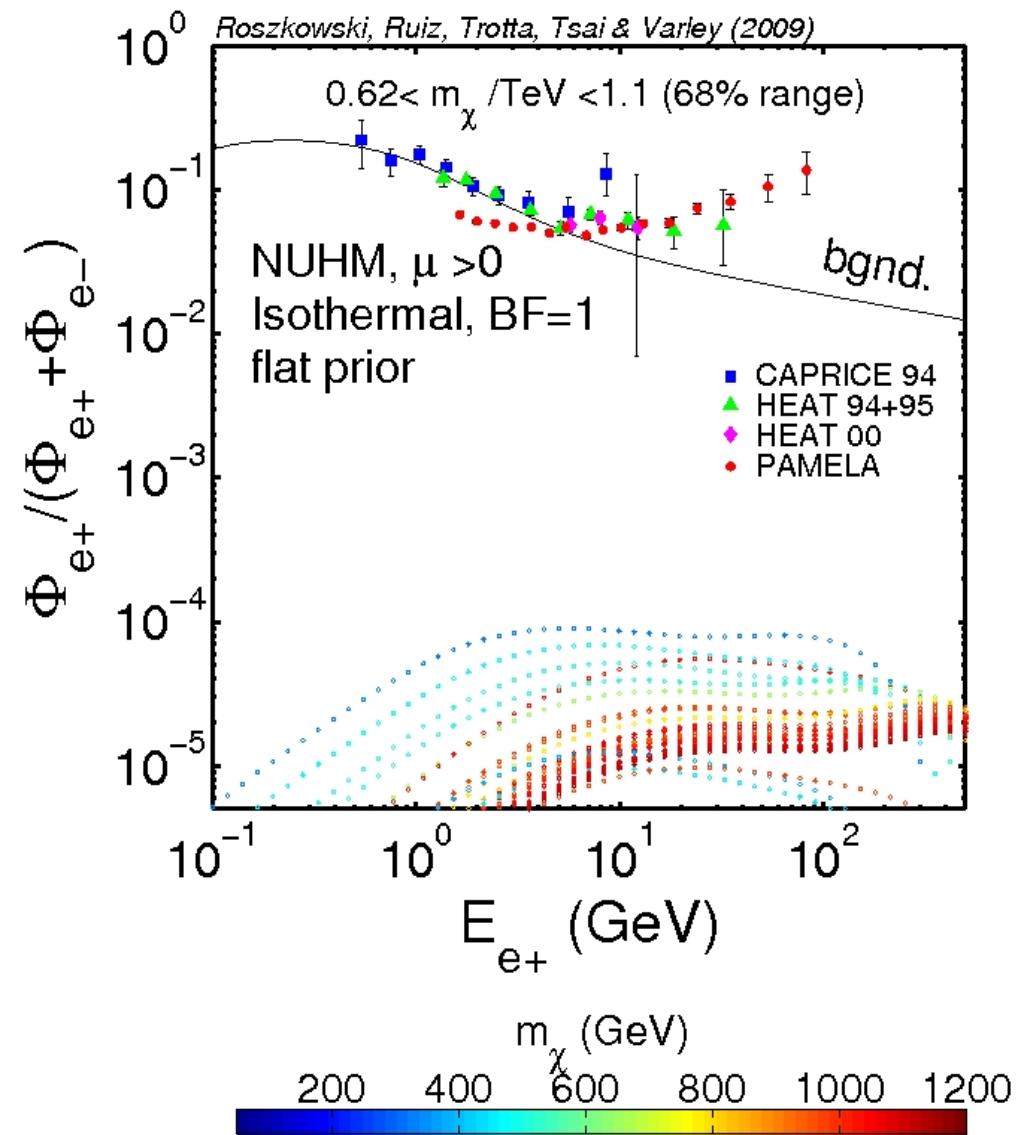
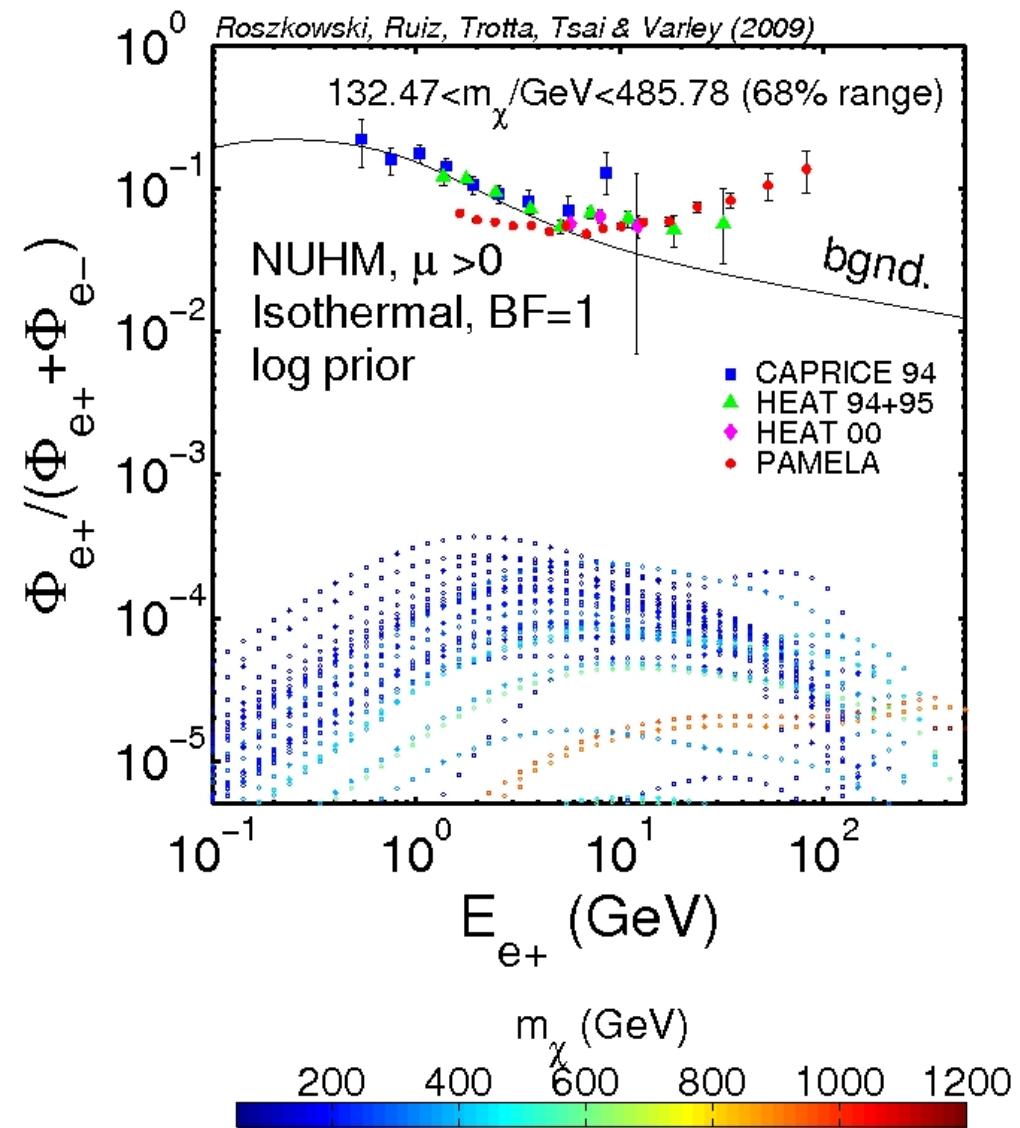
$(Z^0 A^0, Z^0 Z^0, W^+ W^-)$

$v \sim 0$

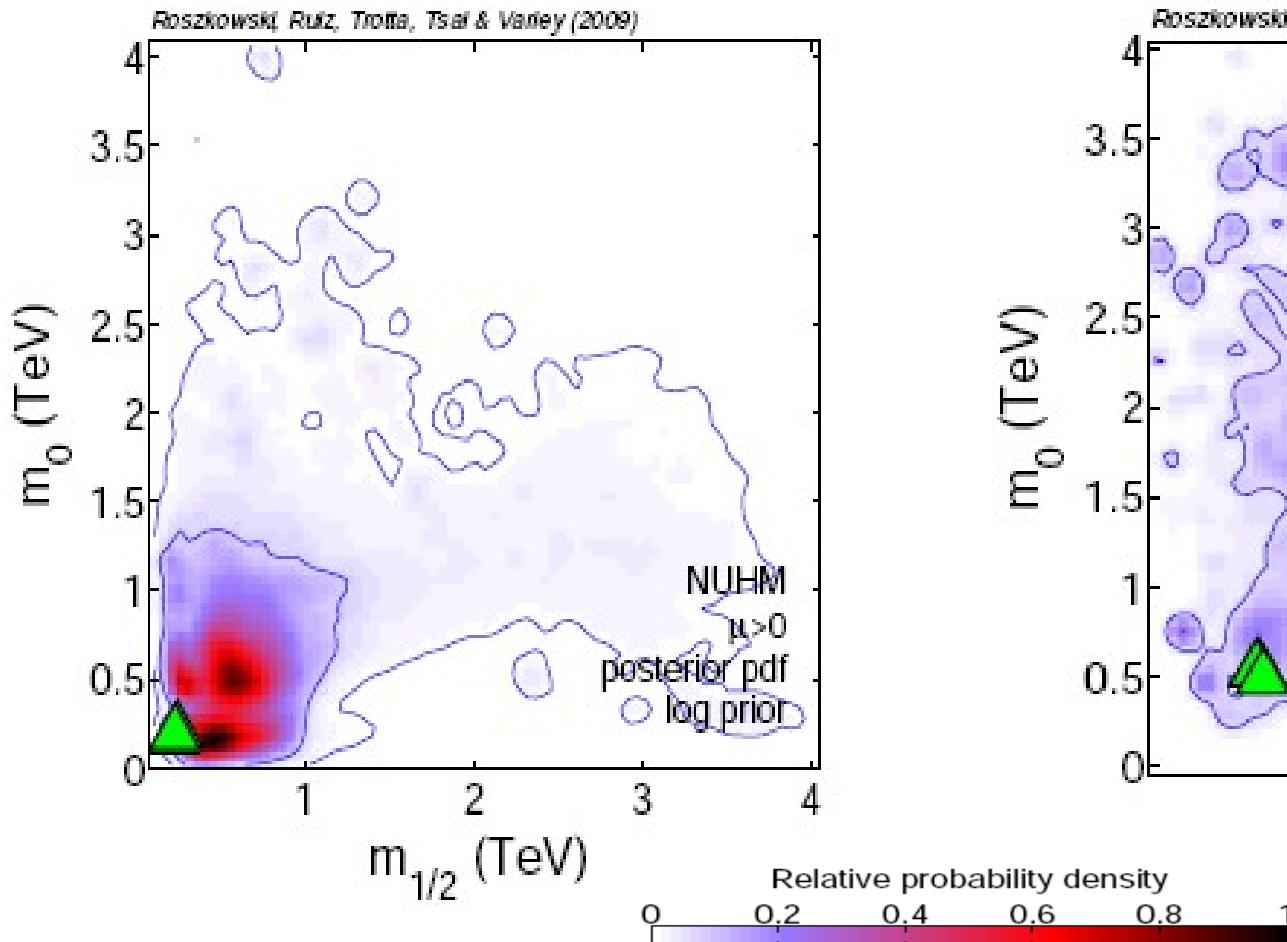
# Positron fraction (NFW)



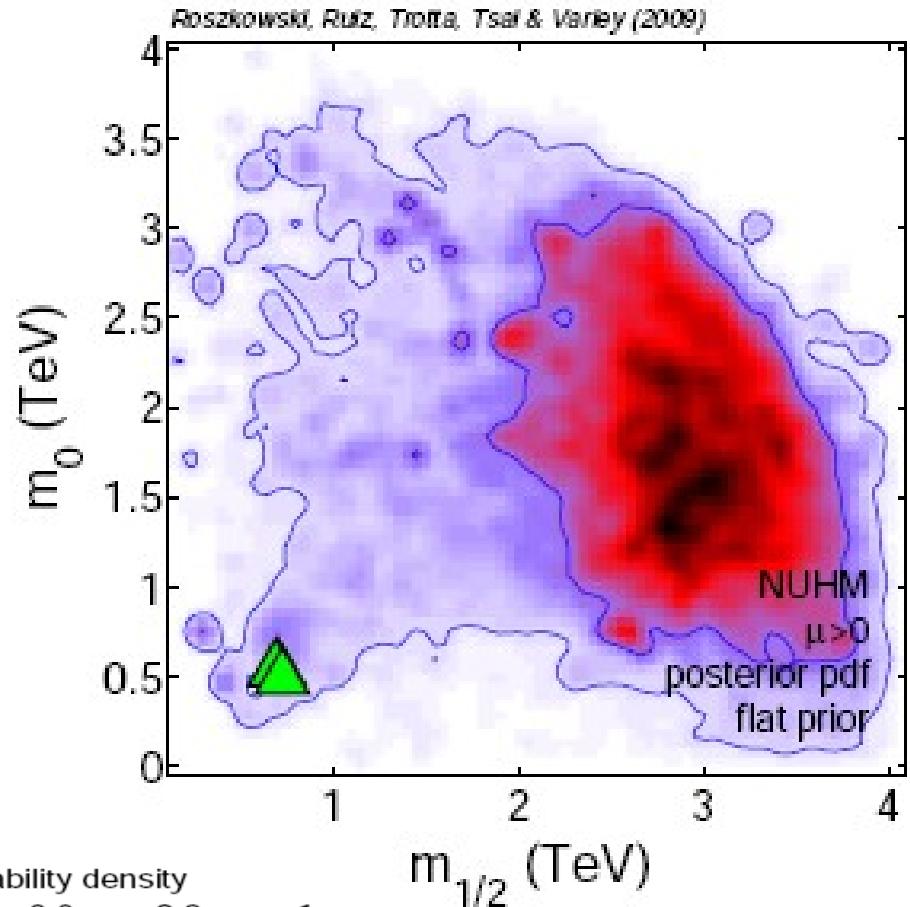
# Positron fraction (Isothermal cored)



# Does direct detection constrain $m_0$ - $m_{1/2}$ plane?



**log prior**



**flat prior**