

# Detection of mirror-meson decays at CERN

George Triantaphyllou

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# Motivation

- What hides behind the **BEH** mechanism?
- What causes the  $M_{Planck}/M_{BEH} \sim 10^{17}$  **hierarchy**?
- The **LHC** has started producing interesting results

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- 1 Motivation
- 2 Inductive approach
  - Hierarchy stabilization and extrapolation to  $M_{Planck}$
  - Mirror fermions: a promising alternative
- 3 Katopton phenomenology
  - General considerations
  - Mirror meson decays at the LHC
- 4 Conclusions

## Some known solutions

- Large extra dimensions  
Stabilizer: **size** of extra dimensions
- Known particles have spin-zero partners (SUSY)  
Stabilizer: **space-time symmetry** - Interaction: **weak**
- Known particles have mirror partners (Katoptrons)  
Stabilizer: **gauge symmetry** - Interaction: **strong**

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## Previous work

- T.D. Lee and C.N. Yang (1956): **Mirror** fermions
- J.C. Pati and A. Salam (1973): Coupling unification
- S. Weinberg (1976), L. Susskind (1979):  
Dynamical BEH mechanism  
Universe: a "**superconductor**",
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# The Katoptron Lagrangian

$$\mathcal{L} = \mathcal{L}_{YM} + \mathcal{L}_{int}$$

$$\mathcal{L}_{YM} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu} - \frac{1}{4} G_{\mu\nu}^e G^{e\mu\nu} - \frac{1}{4} G_{\mu\nu}^K G^{K e\mu\nu}$$

$$\mathcal{L}_{int} = i \sum_{j,k} \left[ (\bar{\Psi}_u^{j|k}, \bar{\Psi}_d^{j|k}) \gamma_\mu \mathcal{D}_k^\mu \begin{pmatrix} \Psi_u^{j|k} \\ \Psi_d^{j|k} \end{pmatrix} + (\bar{\hat{\Psi}}_u^{j|k}, \bar{\hat{\Psi}}_d^{j|k}) \gamma_\mu \hat{\mathcal{D}}_k^\mu \begin{pmatrix} \hat{\Psi}_u^{j|k} \\ \hat{\Psi}_d^{j|k} \end{pmatrix} \right]$$

$$\Psi_u^{j|k} = (N_L^j, U_L^j, N_R^K \delta^{3j}, U_R^K \delta^{3j})$$

$$\Psi_d^{j|k} = (E_L^j, D_L^j, E_R^K \delta^{3j}, D_R^K \delta^{3j})$$

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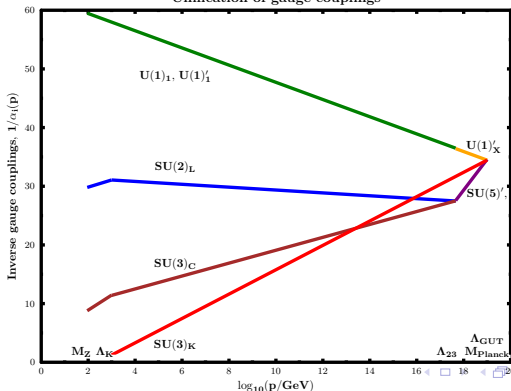
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# Gauged mirror-family symmetry: strong at 1 TeV $\sim M_{\text{Planck}} \exp(-1/\alpha_{\text{GUT}})$ : G.T., *EJTP* 10 (2013) 135.

$$E_8 \times E'_8(M_{\text{Planck}}) \rightarrow SU(3)_C \times SU(2)_L \times U(1)_1 \times SU(3)_K(1 \text{ TeV}) \rightarrow \text{Standard Model}$$

Unification of gauge couplings



# The (parity-odd) spectrum

- **Color singlets:**

$\pi^{K 2 0}$  and  $\pi^{K 2 \pm}$ , "eaten" by  $Z^0$ ,  $W^\pm$

$\pi^{K 1 0}$ ,  $\pi^{K 1 \pm}$ ,  $\pi^{K 1 0'}$ ,  $\eta^K$  (spin - 0)  
 $\rho^{K 1,2 0}$ ,  $\rho^{K 1,2 \pm}$ ,  $\rho^{K 0'}$ ,  $\omega^K$  (spin - 1)

- **Color triplets (leptoquarks):**

$\pi_3^{K 1,2,2',5}$ ,  $\bar{\pi}_3^{K 1,2,2',5}$  (spin - 0)  
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- Katopton-family group self breaks:  
**Doubling** of the mirror-meson spectrum  $\rightarrow$   
heavy (group "B") and light (group "A") mirror mesons:

$$r = M_B/M_A \sim \exp\left(3(C_2(SU(3)_K) - C_2(SU(2)_K))\right) \sim 5.75$$

G.T., *Mod. Phys. Lett. A* **16** (2001) 53.

- **First** signature of  $r$ -hierarchy  $\rightarrow$  LHC excess of diphoton events at **750 GeV** corresponding to a new scalar  $\sigma_B^K$   
**6** times heavier than  $\sigma_A^K$  (known **Higgs** boson).

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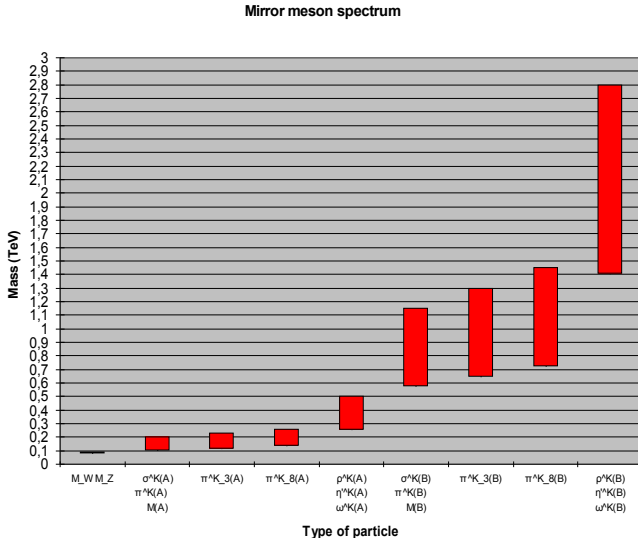
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# Group A & B mirror-meson mass spectra





# Dynamical Masses and the S parameter

$$\langle \bar{\Psi}_{u,d}^{3|3,4} \hat{\Psi}_{u,d}^{3|3,4} + \text{h.c.} \rangle \approx -\frac{N_i}{4\pi^2} \int dp^2 M_i(p^2)$$

$$m_f = \begin{pmatrix} m_{SM} & m \\ m & M \end{pmatrix}, \text{ with } M \equiv \begin{pmatrix} M_A & 0 \\ 0 & M_B \end{pmatrix} \text{ and } m \equiv \begin{pmatrix} m_{AA} & m_{AB} \\ m_{AB} & m_{BB} \end{pmatrix}$$

$$S = S_A + S_B = 4\pi \Sigma_A \left( \frac{F_{\rho_A^K}^2}{M_{\rho_A^K}^2} - \frac{F_{a_A^K}^2}{M_{a_A^K}^2} \right) + 4\pi \Sigma_B \left( \frac{F_{\rho_B^K}^2}{M_{\rho_B^K}^2} - \frac{F_{a_B^K}^2}{M_{a_B^K}^2} \right)$$

$$S_B \approx 4\pi(v/M_{\rho_B^K})^2 \lesssim 0.122 \text{ for } M_{\rho_B^K} > 2.5 \text{ TeV}$$

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# Deviations from SM due to radiative corrections

Importance of heavy fermions:

- **top** (and bottom)-quark **left-right** asymmetries
- $V_{tb}$  not assuming  $3 \times 3$  CKM matrix unitarity
- lepton universality

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# Most promising processes

- $gg \longrightarrow \pi_{8(A)B}^{K 0'} \longrightarrow (\bar{b}b) \bar{t}t$
- $gg \longrightarrow \sigma_{(A)B}^{K 0} \longrightarrow (\bar{b}b) \bar{t}t, \gamma\gamma$
- **Acollinear**  $\bar{f}_i f_i$  jets:

$$gg \longrightarrow \text{direct, } g, \rho_{8B}^{K 0'} \longrightarrow \pi_{8B}^{K +} \pi_{8B}^{K -} \longrightarrow \bar{b}t + \bar{t}b$$

$$gg \longrightarrow \text{direct, } g, \rho_{8B}^{K 0'} \longrightarrow \pi_{3B}^{K 5} \bar{\pi}_{3B}^{K 5} \longrightarrow \bar{\tau}t + \bar{t}\tau$$

$$gg \longrightarrow \text{direct, } g, \rho_{8B}^{K 0'} \longrightarrow \pi_{3B}^{K 2} \bar{\pi}_{3B}^{K 2} \longrightarrow \bar{\nu}_\tau t + \bar{t} \nu_\tau \text{ or } \bar{\tau}b + \bar{b}\tau$$

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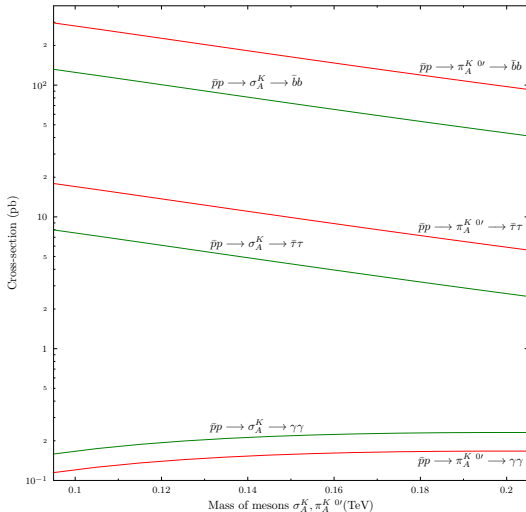
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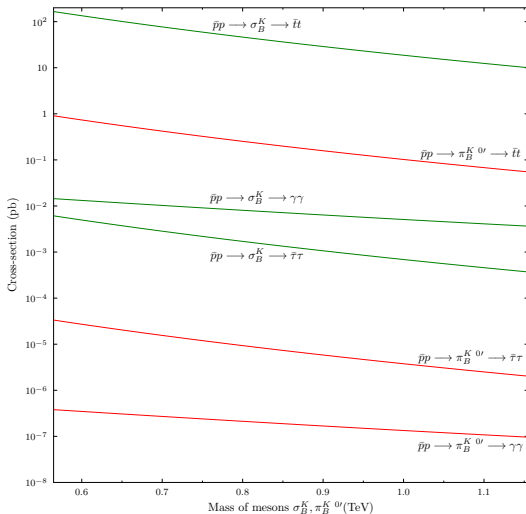
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# Important A-mirror-meson processes

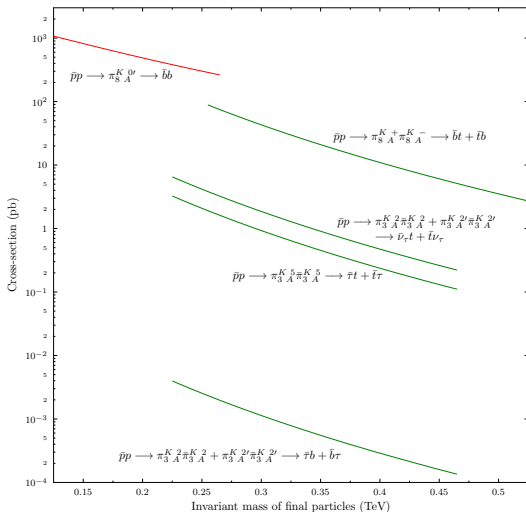




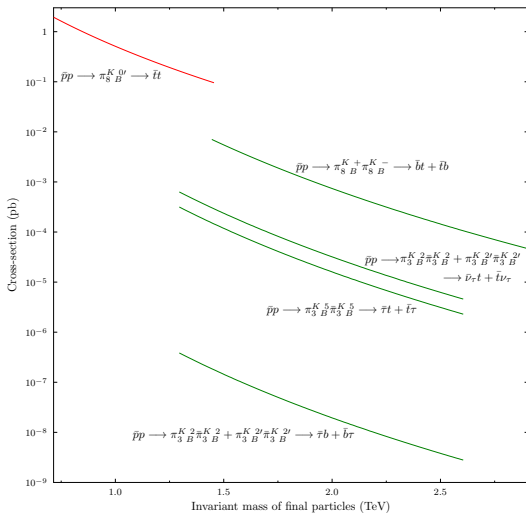
# Important B-mirror-meson processes



# Color-octet & color-triplet A-mirror mesons



# Color-octet & color-triplet B-mirror mesons



# Conclusions

- Viable effective BEH mechanism involving **strongly-interacting mirror fermions** (katoptrons)
- Rich **mirror-meson** LHC phenomenology → Holistic approach in order to **differentiate** competing models predicting similar signals
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