

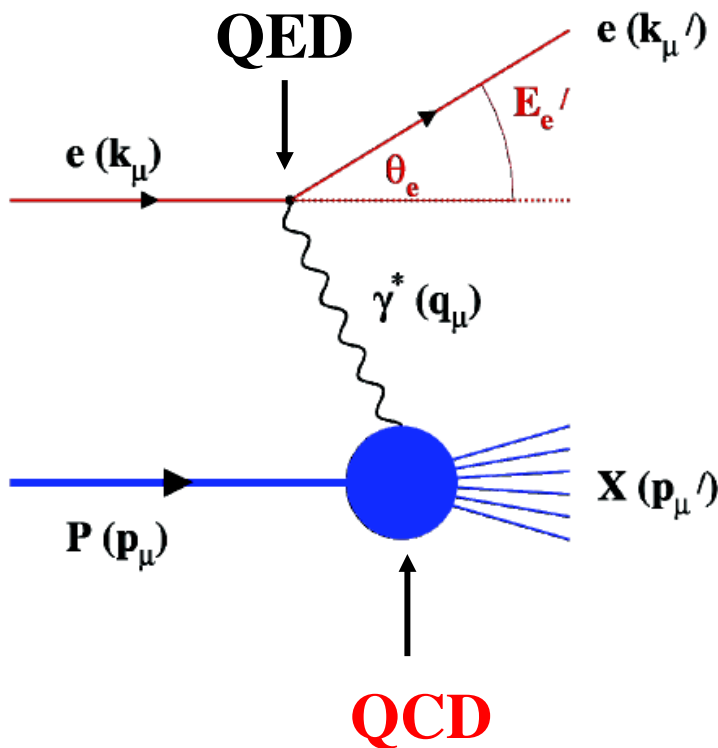
# **QCD evolution equations at small $x$ :** pomeron, odderons and more

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*Guido Altarelli memorial session*  
5<sup>th</sup> International Conference on New Frontiers in Physics  
6-14 July, 2016, Crete, Greece

# Deeply Inelastic Scattering (DIS)

## probing hadron structure



### Kinematic Invariants

$$Q^2 = -q^2 = -(k_\mu - k'_\mu)^2$$

$$Q^2 = 4E_e E'_e \sin^2\left(\frac{\theta'_e}{2}\right)$$

$$y = \frac{pq}{pk} = 1 - \frac{E'_e}{E_e} \cos^2\left(\frac{\theta'_e}{2}\right)$$

$$x = \frac{Q^2}{2pq} = \frac{Q^2}{sy}$$

$$s \equiv (p + k)^2$$

Measure of  
resolution  
power

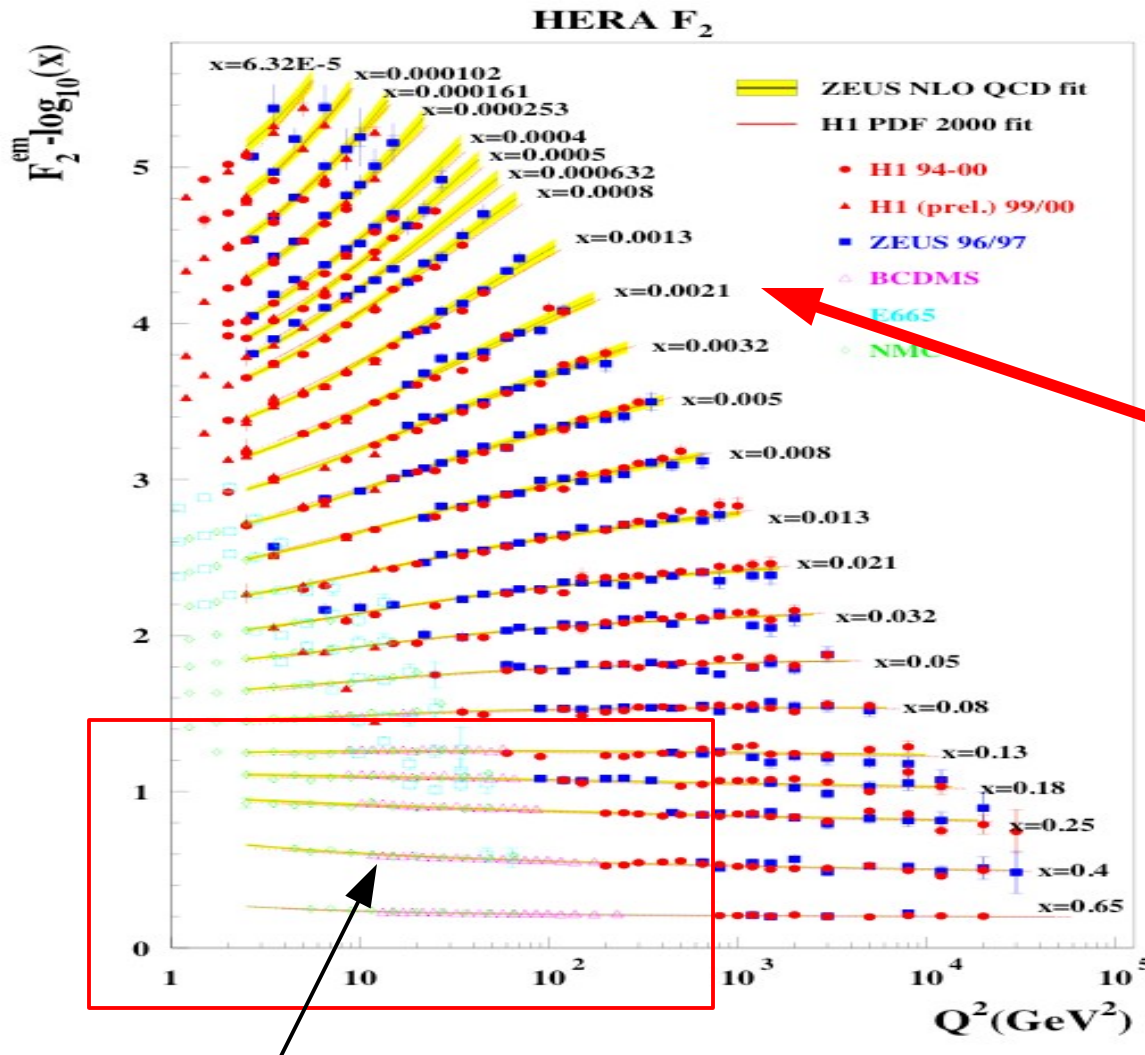
Measure of  
inelasticity

Measure of  
momentum  
fraction of  
struck quark

structure functions  $F_1, F_2$

$$F_2(x, Q^2) \equiv \sum_f e_f^2 x [q_f(x, Q^2) + \bar{q}_f(x, Q^2)]$$

# Deeply Inelastic Scattering $e(k) p(p) \rightarrow e(k') X$



$$F_2 \equiv \frac{Q^2}{4\pi^2 \alpha_{\text{em}}} \sigma^{\gamma^* p}$$

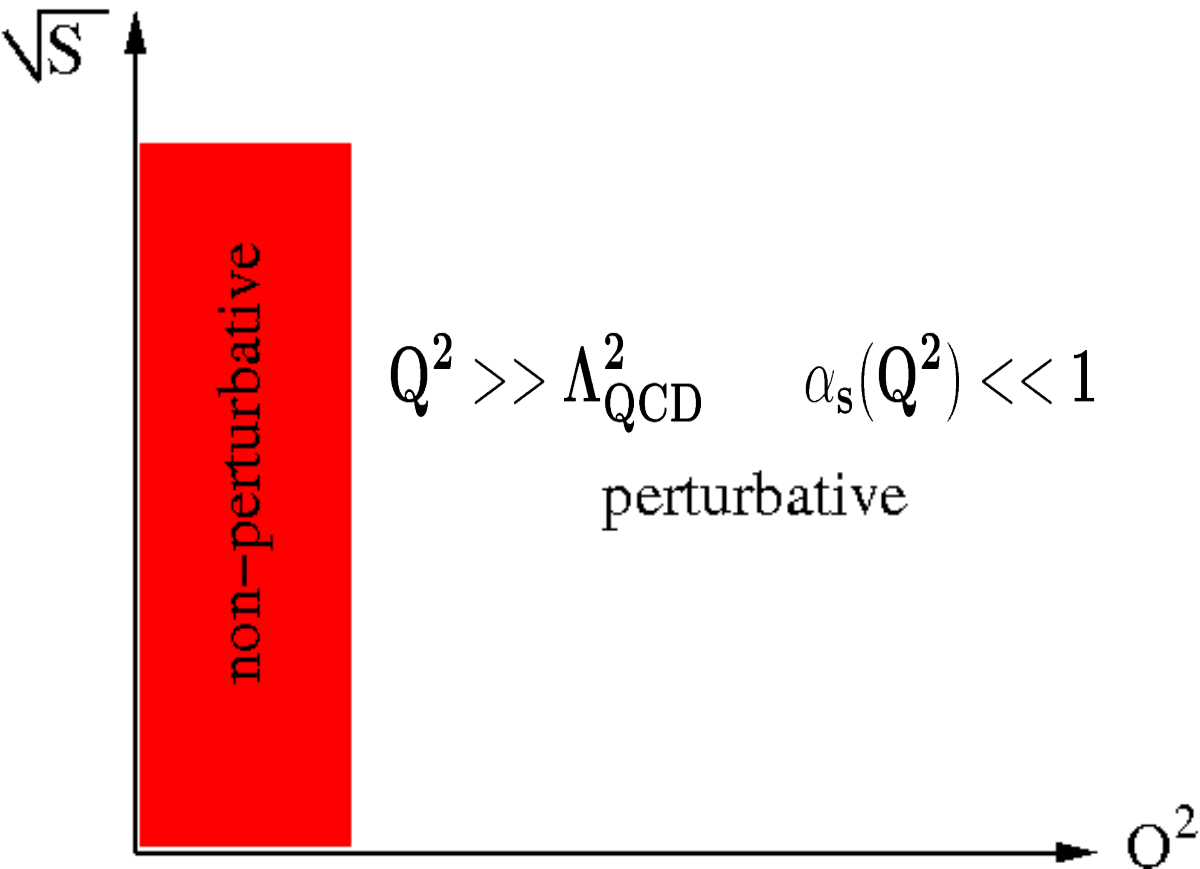
scaling violations

parton model of a hadron leads to scale invariance

**DGLAP** ( $Q^2$ )  
evolution of quark  
and gluon  
distribution  
functions in pQCD

early experiments (SLAC,...):  
scale invariance of hadron structure

# QCD: the standard paradigm



pQCD tools:  
twist expansion,  
collinear factorization

$$\sqrt{S} \rightarrow \infty$$

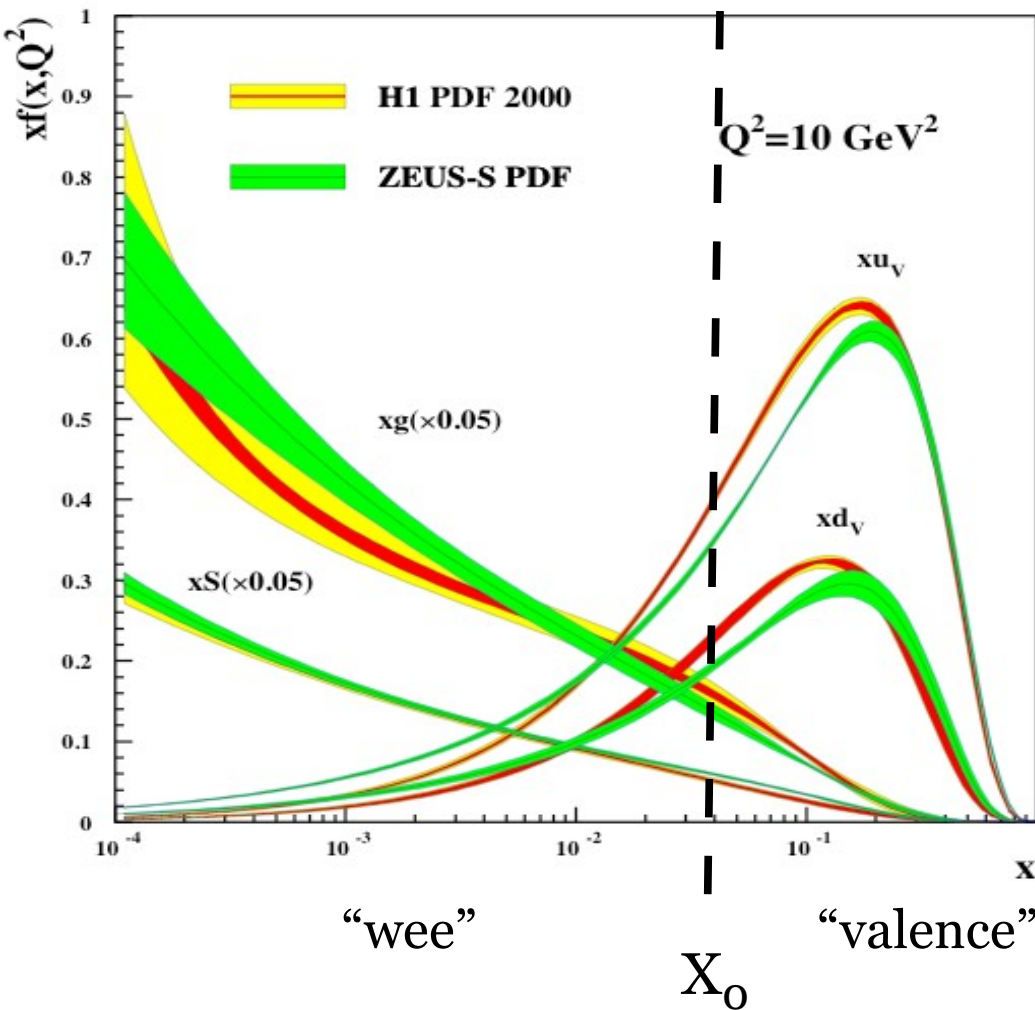
$$Q^2 \rightarrow \infty$$

*but bulk of QCD phenomena happens at low  $p_t$*

$$x \sim \frac{p_t}{\sqrt{S}} e^{-y} \rightarrow 0$$

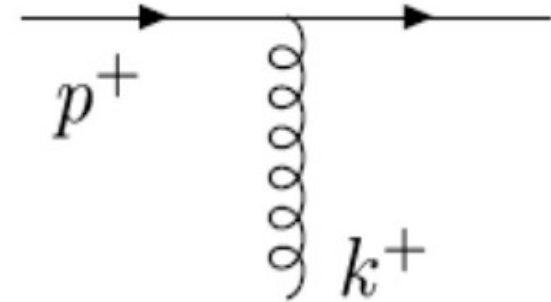
*QCD cross sections  
at small  $x$*

# A hadron at small $x$



recall for any  
4-momentum

$$p^2 = 2p^+ p^- - p_t^2$$



small  $x$ :  $k^+ \ll p^+$

$$\tau_{\text{val}} \sim \frac{1}{p^-} = \frac{2p^+}{k_t^2} \gg \tau_{\text{wee}} \sim \frac{1}{k^-} = \frac{2k^+}{k_t^2}$$

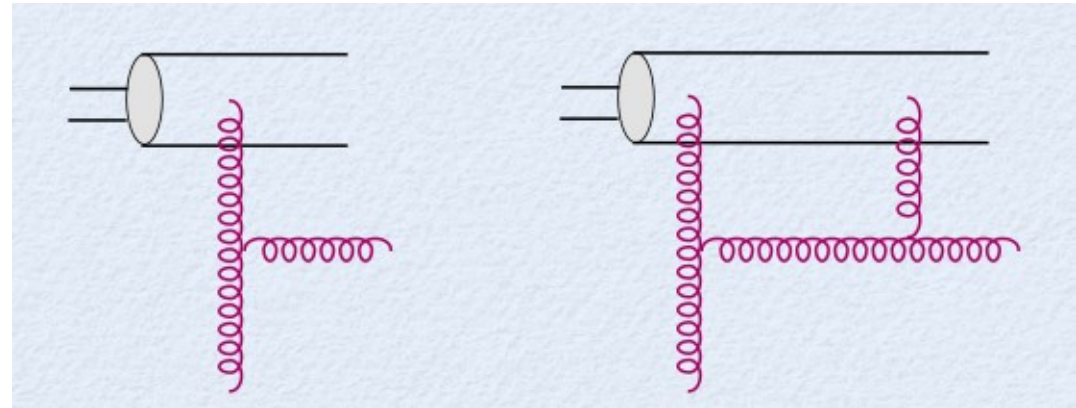
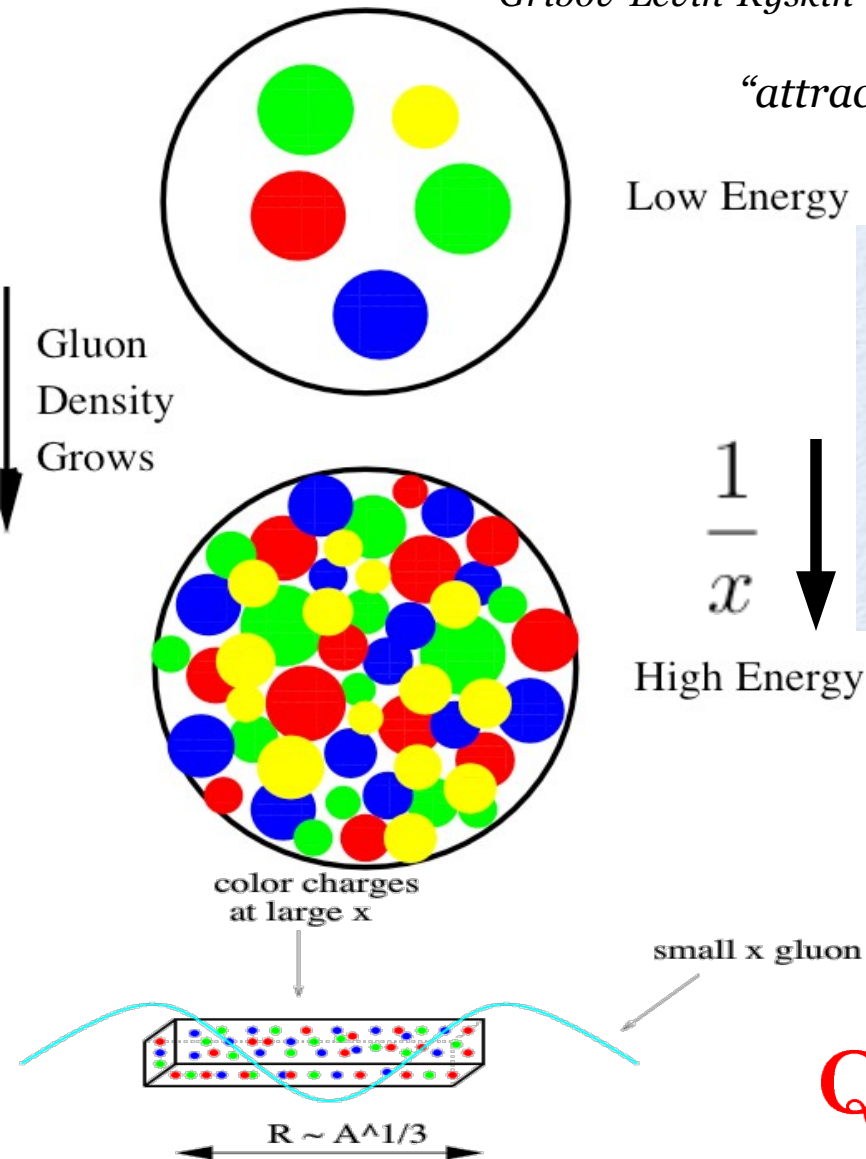
*high  $x$  partons as static color charges  $\rho$*

# Gluon saturation/CGC

Gribov-Levin-Ryskin

$$S \rightarrow \infty, \quad Q^2 \text{ fixed}, \quad x_{Bj} \equiv \frac{Q^2}{S} \rightarrow 0$$

*“attractive” bremsstrahlung vs. “repulsive” recombination*



$$\frac{\alpha_s x G(x, b_t, Q^2)}{S_\perp Q^2} \sim 1$$

$$Q_s^2(x, b_t, A) \sim A^{1/3} \left(\frac{1}{x}\right)^{0.3}$$

# ***MV effective Action + RGE***

$$S[\mathbf{A}, \rho] = -\frac{1}{4} \int d^4x F_{\mu\nu}^2 + \frac{i}{N_c} \int d^2x_t dx^- \delta(x^-) \text{Tr}[\rho(x_t) U(\mathbf{A}^-)]$$

Large  $x$ : color source  $\rho$

small  $x$ : gluon field  $\mathbf{A}^\mu$

$$U(\mathbf{A}^-) = \hat{P} \text{Exp} \left[ ig \int dx^+ \mathbf{A}_a^- T_a \right]$$

$$\mathbf{Z}[\mathbf{j}] = \int [\mathbf{D}\rho] \mathbf{W}_{\Lambda^+}[\rho] \left[ \frac{\int^{\Lambda^+} [\mathbf{D}\mathbf{A}] \delta(\mathbf{A}^+) e^{iS[\mathbf{A}, \rho] - \int \mathbf{j} \cdot \mathbf{A}}}{\int^{\Lambda^+} [\mathbf{D}\mathbf{A}] \delta(\mathbf{A}^+) e^{iS[\mathbf{A}, \rho]}} \right]$$

weight functional:

$\mathbf{W}_{\Lambda^+}[\rho]$  probability distribution of color sources  $\rho$   
at longitudinal scale  $\Lambda^+$

# Wilsonian RG: JIMWLK evolution eq

invariance under change of  $\Lambda^+$   $\longrightarrow$  RGE for  $W_{\Lambda^+}[\rho]$

Field  $A_\mu^a$

source  $\rho^a$

weight functional  $W_x[\rho]$

0

x

1

$$x \equiv \frac{\Lambda^+}{P^+}$$



# Wilsonian RG: JIMWLK evolution eq

resum  $\alpha_s \log \frac{1}{x}$

Field  $A_\mu^a$

source  $\rho^a$

weight functional  $W_x[\rho]$

0

x

1

$$A^\mu = A_{\text{class}}^\mu + \delta A^\mu$$

integrate out field fluctuations quadratically

$$\rho \rightarrow \rho' = \rho + \delta \rho \quad \mathbf{W}[\rho] \rightarrow \mathbf{W}[\rho']$$

# Wilsonian RG: JIMWLK evolution eq

resum  $\alpha_s \log \frac{1}{x}$

Field  $A_\mu^a$

source  $\rho^a$

weight functional  $W_x[\rho]$

0

$x$

1

$$\frac{\partial}{\partial \ln 1/x} \langle O \rangle = \left\langle \frac{1}{2} \int d^2 x_t d^2 y_t \frac{\delta}{\alpha_a(x_t)} \eta^{ab}(x_t, y_t) \frac{\delta}{\delta \alpha_b(y_t)} O \right\rangle$$

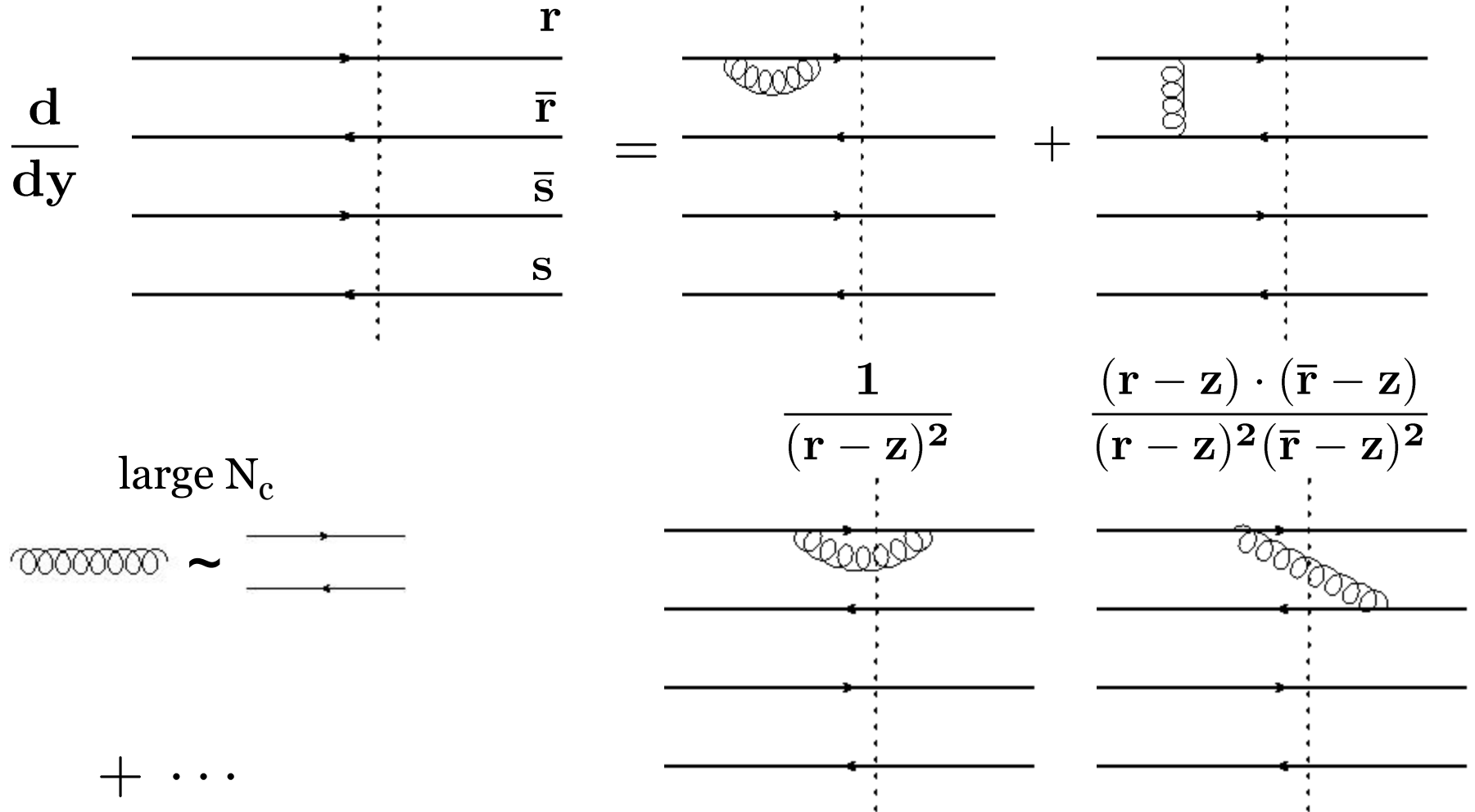
$$\eta^{ab}(x_t, y_t) = K \left[ 1 + U^\dagger(x_t) U(y_t) - U^\dagger(x_t) U(z_t) - U^\dagger(z_t) U(y_t) \right]^{ab}$$

$$K = \frac{1}{\pi} \int \frac{d^2 z_t}{(2\pi)^2} \frac{(x_t - z_t) \cdot (y_t - z_t)}{(x_t - z_t)^2 (y_t - z_t)^2}$$

# Evolution of quadrupole from JIMWLK

$$Q(\mathbf{r}, \bar{\mathbf{r}}, \bar{\mathbf{s}}, \mathbf{s}) \equiv \frac{1}{N_c} \langle \text{Tr } V(\mathbf{r}) V^\dagger(\bar{\mathbf{r}}) V(\bar{\mathbf{s}}) V^\dagger(\mathbf{s}) \rangle$$

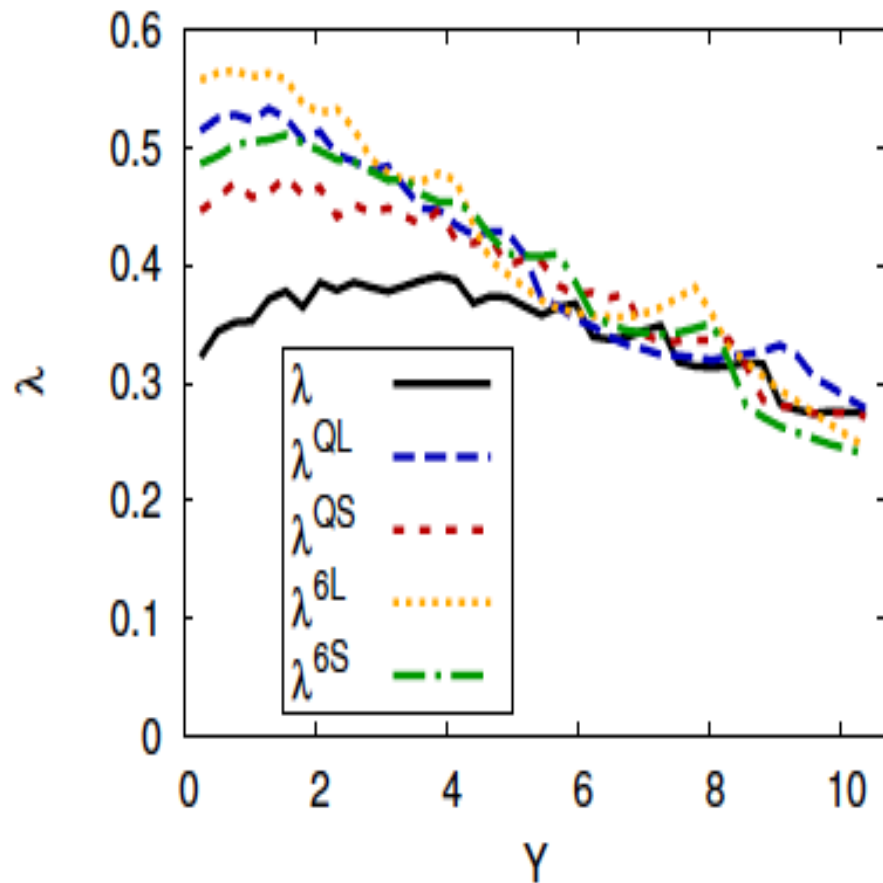
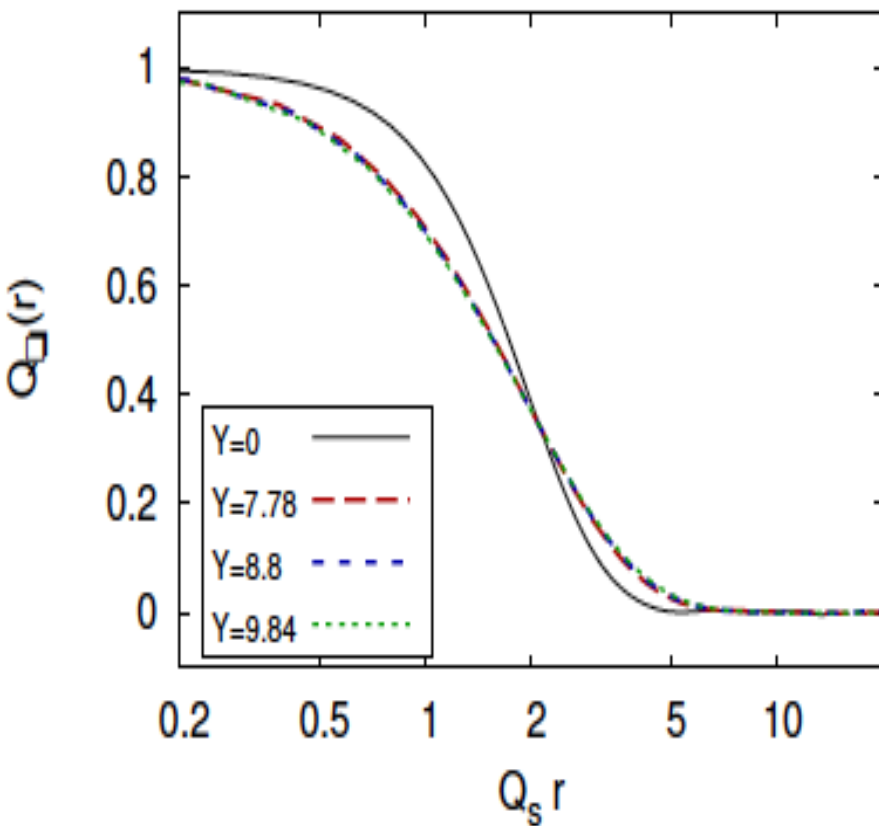
radiation kernels  
as in dipole



# Evolution of **quadrupole** from JIMWLK

$$\begin{aligned}
 & \frac{d}{dy} \langle Q(r, \bar{r}, \bar{s}, s) \rangle \\
 = & \frac{N_c \alpha_s}{(2\pi)^2} \int d^2 z \left\{ \left\langle \left[ \frac{(r - \bar{r})^2}{(r - z)^2 (\bar{r} - z)^2} + \frac{(r - s)^2}{(r - z)^2 (s - z)^2} - \frac{(\bar{r} - s)^2}{(\bar{r} - z)^2 (s - z)^2} \right] Q(z, \bar{r}, \bar{s}, s) S(r, z) \right. \right. \\
 + & \left[ \frac{(r - \bar{r})^2}{(r - z)^2 (\bar{r} - z)^2} + \frac{(\bar{r} - \bar{s})^2}{(\bar{r} - z)^2 (\bar{s} - z)^2} - \frac{(r - \bar{s})^2}{(r - z)^2 (\bar{s} - z)^2} \right] Q(r, z, \bar{s}, s) S(z, \bar{r}) \\
 + & \left[ \frac{(\bar{r} - \bar{s})^2}{(\bar{r} - z)^2 (\bar{s} - z)^2} + \frac{(s - \bar{s})^2}{(s - z)^2 (\bar{s} - z)^2} - \frac{(\bar{r} - s)^2}{(s - z)^2 (\bar{r} - z)^2} \right] Q(r, \bar{r}, z, s) S(\bar{s}, z) \\
 + & \left[ \frac{(r - s)^2}{(r - z)^2 (s - z)^2} + \frac{(s - \bar{s})^2}{(s - z)^2 (\bar{s} - z)^2} - \frac{(r - \bar{s})^2}{(r - z)^2 (\bar{s} - z)^2} \right] Q(r, \bar{r}, \bar{s}, z) S(z, s) \\
 - & \left[ \frac{(r - \bar{r})^2}{(r - z)^2 (\bar{r} - z)^2} + \frac{(s - \bar{s})^2}{(s - z)^2 (\bar{s} - z)^2} + \frac{(r - s)^2}{(r - z)^2 (s - z)^2} + \frac{(\bar{r} - \bar{s})^2}{(\bar{r} - z)^2 (\bar{s} - z)^2} \right] Q(r, \bar{r}, \bar{s}, s) \\
 - & \left[ \frac{(r - s)^2}{(r - z)^2 (s - z)^2} + \frac{(\bar{r} - \bar{s})^2}{(\bar{r} - z)^2 (\bar{s} - z)^2} - \frac{(\bar{r} - s)^2}{(\bar{r} - z)^2 (s - z)^2} - \frac{(r - \bar{s})^2}{(r - z)^2 (\bar{s} - z)^2} \right] S(r, s) S(\bar{r}, \bar{s}) \\
 - & \left. \left[ \frac{(r - \bar{r})^2}{(r - z)^2 (\bar{r} - z)^2} + \frac{(s - \bar{s})^2}{(s - z)^2 (\bar{s} - z)^2} - \frac{(r - \bar{s})^2}{(r - z)^2 (\bar{s} - z)^2} - \frac{(\bar{r} - s)^2}{(\bar{r} - z)^2 (s - z)^2} \right] S(r, \bar{r}) S(\bar{s}, s) \right\rangle \Bigg\} \\
 \frac{d}{dy} \mathbf{Q} = & \int \mathbf{P}_1 [\mathbf{Q} \mathbf{S}] - \mathbf{P}_2 [\mathbf{Q}] + \mathbf{P}_3 [\mathbf{S} \mathbf{S}] \quad \text{with} \quad P_1 - P_2 + P_3 = 0
 \end{aligned}$$

# Quadrupole evolution: JIMWLK



*geometric scaling also present in quadrupoles*

*energy dependence of saturation scale*

## ***JIMWLK evolution in the linear regime***

$$V(\mathbf{x}_t) \equiv \text{[Diagram: A thick black vertical line with a horizontal line at the top] } \equiv \text{[Diagram: Two vertical lines, each with a series of loops, and a horizontal line at the top with an arrow pointing right] } \sim 1 + \mathcal{O}(g A) + \mathcal{O}(g^2 A^2)$$

*BJKP: evolution of  $n$ -Reggeized gluons in a singlet state*

**O(A<sup>2</sup>)**    2-gluon exchange (BFKL pomeron)

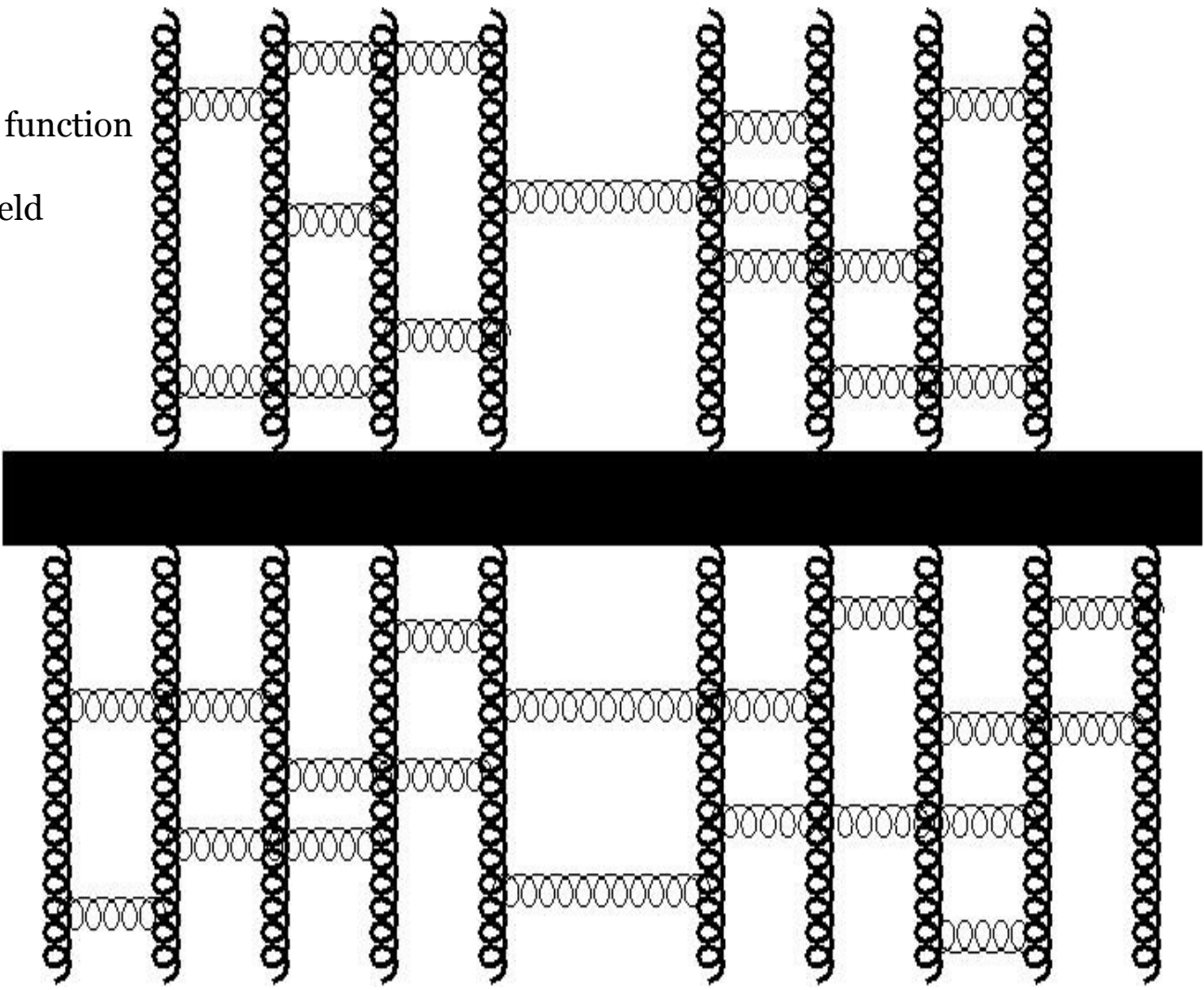
**O(A<sup>3</sup>)**    3-gluon exchange (odderon)  
                   i) a dipole  
                   ii) 3-quarks

**O(A<sup>4</sup>)** 4-gluon exchange

**JIMWLK (linear) and BJKP eqs. agree for  $n=2,3,4,\dots,n$**

# JIMWLK in non-linear regime: $n \rightarrow n+1$ “pomeron” vertices

Follow the same strategy:  
Evolution equation for  $2n$ -pt function  
Keep the non-linear terms  
Expand in powers of gluon field  
Extract the vertex



**1  $\rightarrow$  2 vertex**  
**(triple pomeron)**

**2  $\rightarrow$  3 vertex**

**$n \rightarrow n + 1$  vertex**

# ***CGC signatures?***

## ***two main effects:***

*multiple scatterings*  
*evolution with  $x$  (rapidity)*

### ***dense-dense (AA, pA, pp) collisions***

*initial conditions*

### ***dilute-dense (pA, forward pp ) collisions***

*multiplicities*

*$p_t$  spectra*

*angular correlations*

### ***spin asymmetries***

### ***DIS***

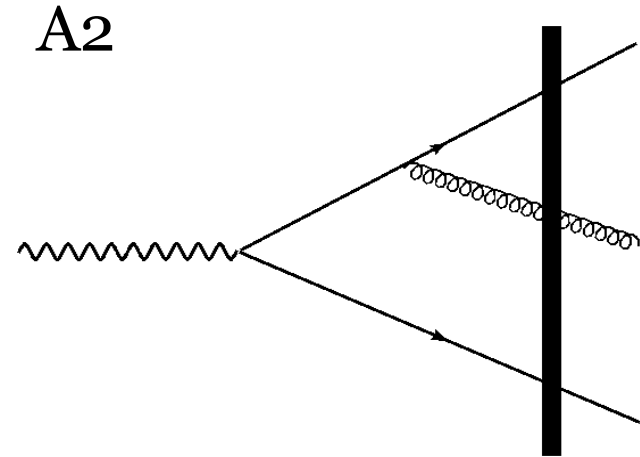
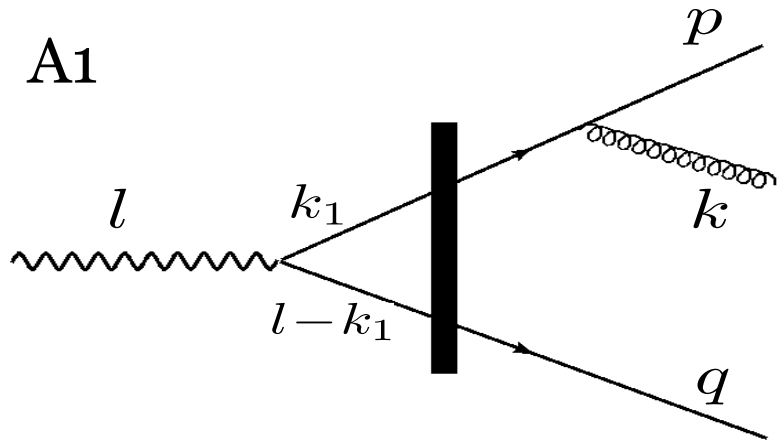
*structure functions (diffraction)*

**NLO** di-hadron correlations

***3-hadron/jet correlations***



# 3-jet production in DIS: azimuthal correlations



$$A_{1,hg}^L = \sqrt{2Q^2} e^{ix_t(k_t+p_t)+iq_t y_t} K_0 \left[ \sqrt{Q^2 x_{12}^2 z_2 (z_1 + z_3)} \right] \cdot a_{1,hg}^L$$

with  $x_{12}^2 \equiv (x_t - y_t)^2$

$$a_{1,++}^L = \frac{z_1 z_2 \sqrt{z_1 z_2} (z_1 + z_3)}{z_3 e^{-i\theta_p} |p_t| - z_1 e^{-i\theta_k} |k_t|}$$

$$a_{1,-+}^L = \frac{\sqrt{z_1 z_2} z_2 (z_1 + z_3)^2}{z_3 e^{-i\theta_p} |p_t| - z_1 e^{-i\theta_k} |k_t|}$$

$$a_{1,--}^L = (a_{1,++}^L)^*$$

$$a_{1,+-}^L = (a_{1,+ - +}^L)^*$$

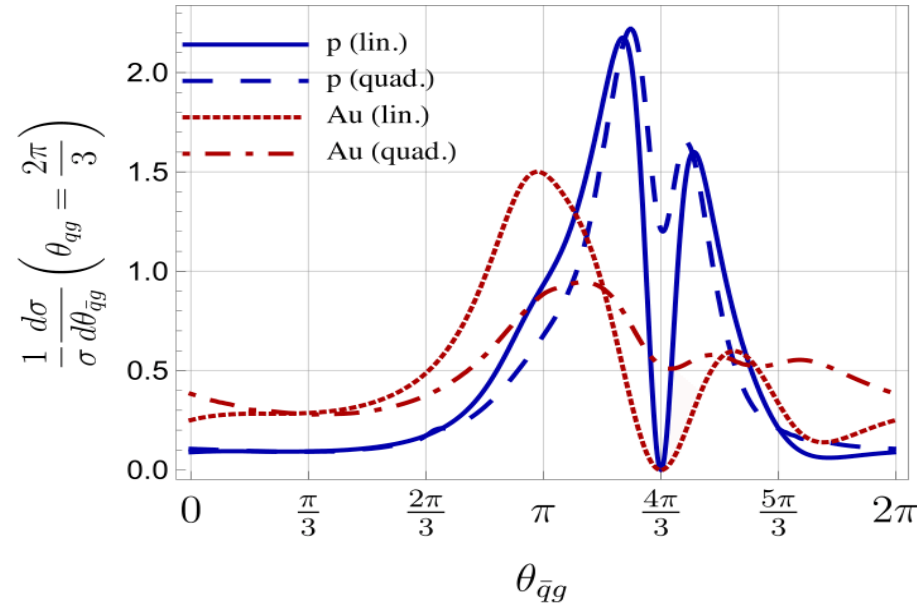
spinor helicity methods

longitudinal photons

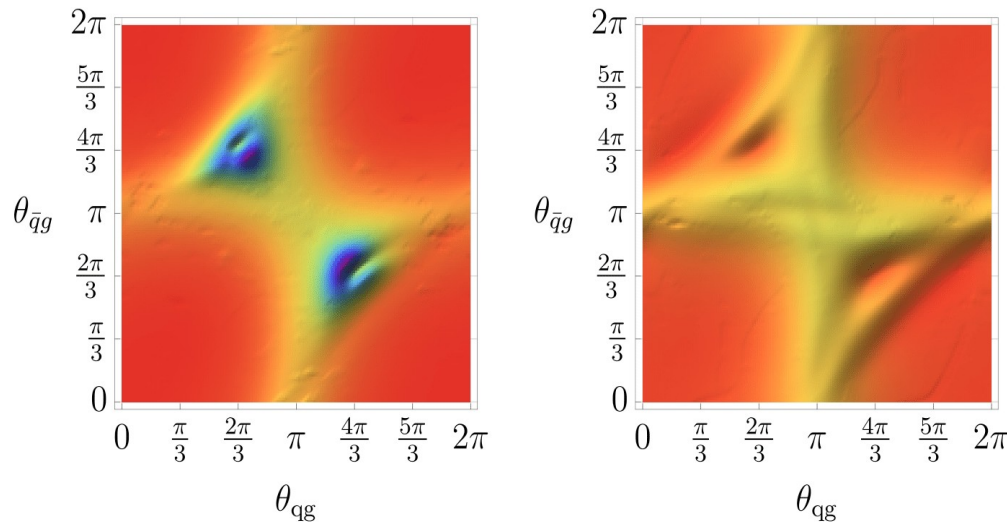
gluon, quark helicity

$$g, h = \pm$$

# 3-parton azimuthal angular correlations in DIS



*disappearance of the away  
side peaks*



*Ayala, Hentschinski, Jalilian-Marian,  
Tejeda-Yeomans, arXiv:1604:08526*

# Possible extensions to other processes

real photons:  $Q^2 \rightarrow 0$

**ultra-peripheral nucleus-nucleus collisions**

crossing symmetry:

$$\gamma^{(\star)} T \longrightarrow q \bar{q} g X \longleftrightarrow \left\{ \begin{array}{l} q T \longrightarrow q g \gamma^{(\star)} X \\ \bar{q} T \longrightarrow \bar{q} g \gamma^{(\star)} X \\ g T \longrightarrow q \bar{q} \gamma^{(\star)} X \end{array} \right\}$$

**proton-nucleus collisions** (collinear factorization in proton?)

$$pA \longrightarrow h_1 h_2 \gamma^{(\star)} X$$

**MPI** (double/triple parton interactions)

$$\gamma^{(\star)} T \longrightarrow q \bar{q} g X \longleftrightarrow \left\{ \begin{array}{l} q \bar{q} T \longrightarrow g \gamma^{(\star)} X \\ g \bar{q} T \longrightarrow \bar{q} \gamma^{(\star)} X \\ g q T \longrightarrow q \gamma^{(\star)} X \end{array} \right\}$$

# ***SUMMARY***

***CGC is a systematic approach to high energy collisions***

***JIMWLK eq reduces to BJKP eq in the linear regime:***

***Pomerons, odderons,....***

***$n \rightarrow n+1$  “pomeron” vertices from non-linear terms in JIMWLK***

***CGC has been used to fit a wealth of data:***

***ep, eA, pp, pA, AA***

***Azimuthal angular correlations offer a unique probe of CGC:***

***3-hadron/jet correlations should be even more discriminatory***

***Must generalize CGC to include high  $p_t$  (DGLAP) physics***

***signs of trouble at  $p_t > Q_s$  ?***