# ReD — Recoil Directionality Studies in Two-Phase Liquid Argon TPC Detectors

Deller

Matteo Cadeddu on behalf of the ReD collaboration

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## The RED Collaboration

#### Italy (INFN):

Scoil Directionality University of Napoli Federico II, University of Roma, University of Cagliari, University of Genova, University of Pisa, LNGS, LNS, TIFPA

















#### **International Collaborators:**

APC-IN2p3, Princeton University, Temple, UCLA











#### Facing the problems of past and future WIMP detection

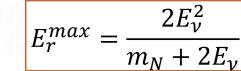
Upcoming direct dark matter detection experiments will have sensitivity to detect neutrinos from several astrophysical sources (Sun, atmosphere, and diffuse Supernovae).

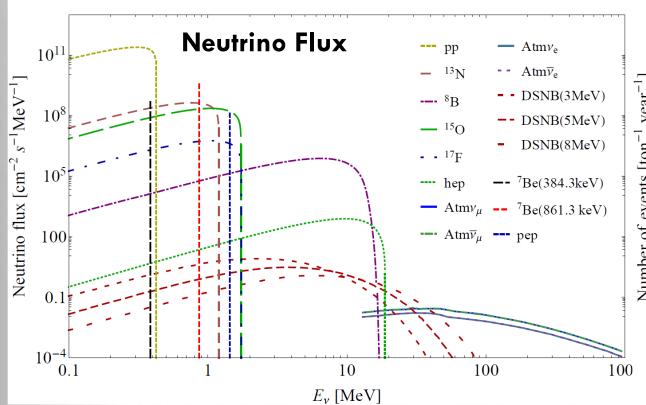
#### **Coherent neutrino scattering on Nucleus (CNS)**

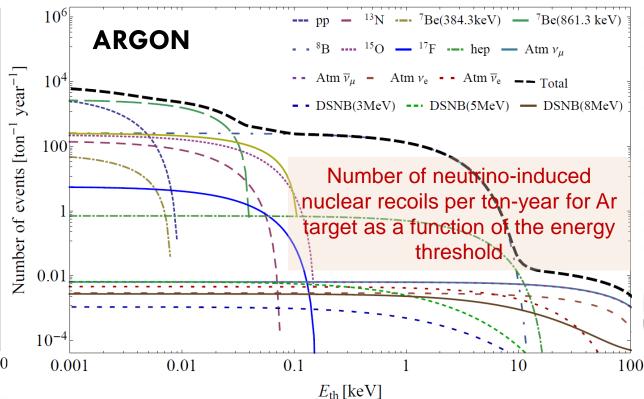
$$\nu_{\chi} + (A, Z) \rightarrow \nu_{\chi} + (A, Z)$$

$$\frac{d\sigma^{CNS}(E_{\nu}, E_{r})}{dE_{r}} = \frac{G_{f}^{2}}{4\pi} Q_{w}^{2} m_{N} \left( 1 - \frac{m_{N} E_{r}}{2E_{\nu}^{2}} \right) F^{2}(E_{r})$$

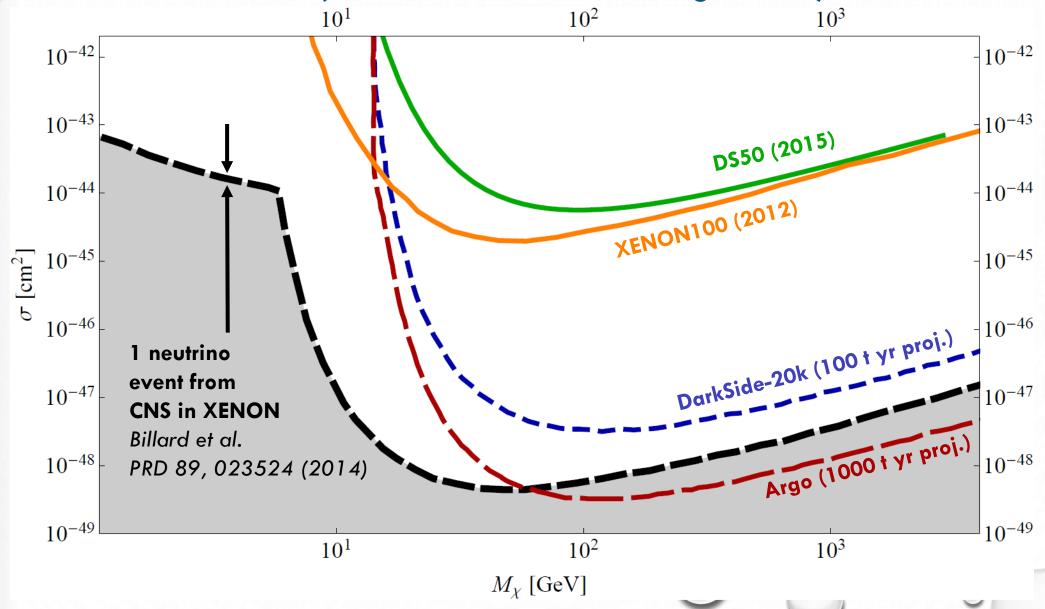
$$Q_W = N - (1 - 4\sin^2\theta_W)Z$$







## The neutrino floor (Neutrinos as a background)



## Searching for the Wimp wind

Non rotating Wimp Halo + Barionic matter rotation

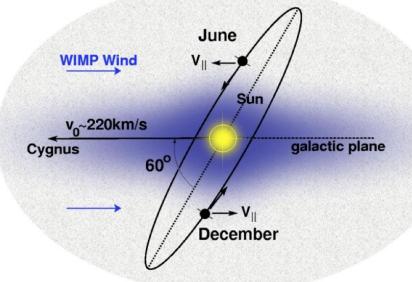


Solar System orbit at  $v_0 \sim \! 220 \ \mathrm{km/s}$  around the galactic center

#### Standard technique:

Annual rate modulation

Earth orbits at  $v_E\cong \pm 30~\mathrm{km/s}$  (few % effect)

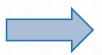


# BUT: Background may be also annual modulated!

#### Innovative technique:

Sidereal direction modulation

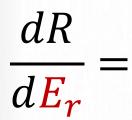
Measuring the angle between WIMP and Earth gives a directionality signature unique to WIMPs.



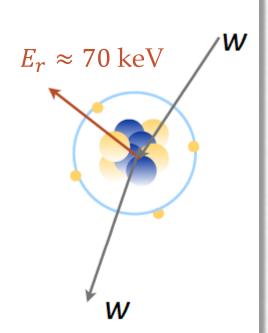
Directionality may be the most robust signature of the WIMP nature of DM

#### The standard WIMP recoil spectrum

Standard non-directional recoil spectrum, i.e. differential event rate per unit detector mass:



#### Recoiling nucleus



$$\frac{\sigma_{w-n}}{2M_w \,\mu_n^2}$$

#### **Physics**

 $\sigma_{w-n} \rightarrow \text{WIMP-}$ nucleon cross section

 $M_w \rightarrow \text{WIMP mass}$ 

 $\mu_n \rightarrow \text{WIMP-nucleon}$ reduced mass



#### Detector

 $A \rightarrow$  atomic mass of target material

 $F(E_r) \rightarrow \text{The finite}$ size of the nucleus is implemented with **Helm form Factor** 

$$\frac{\sigma_{w-n}}{2M_w \,\mu_n^2} \times \left(A^2 \,F^2 \,(E_r)\right) \times \left(\rho \,\int_{v>v_n} \frac{f(v)}{v} d^3v\right)$$

#### **Astrophysics (DM halo** properties)

 $\rho \rightarrow WIMP$  mass density

 $f(v) \rightarrow \text{WIMP}$  velocity distribution

 $v_n \rightarrow \text{minimum WIMP speed}$ required to transfer an energy  $E_r$  to the nucleus of mass  $M_n$  in the detector.

 $E_r \rightarrow \text{Recoiling nucleus energy}$ 

## A closer look to directionality

Double differential directional recoil spectrum:

$$\frac{dR}{dE_r \ d\Omega(\theta, \phi)} = \frac{\sigma_{w-n}}{4\pi \ M_w \ \mu_n^2} A^2 F^2 (E_r) \rho \int \delta(\boldsymbol{v} \cdot \boldsymbol{w} - v_n) f(\boldsymbol{v}) d^3 \boldsymbol{v}$$

If we assume the Standard Halo Model (SHM), i.e., an isotropic Maxwell-Boltzmann WIMP velocity distribution of width  $\sigma_v$  in a inertial reference frame at rest with respect to the Galactic center.

Radon transform  $\equiv \hat{f}(v_n, \mathbf{w})$   $\hat{f}(v_n, \mathbf{w}) = \frac{1}{\sqrt{2\pi\sigma_n^2}} \exp\left[-\frac{1}{2} \frac{(v_n - \mathbf{w} \cdot \mathbf{V})^2}{\sigma_n^2}\right]$ 

WIMP

recoiling

nucleus

target

Kinematics of a WIMP-nucleus

elastic scattering

incoming

WIMP

Here V is the average velocity of the WIMPs with respect to the detector:  $V = -V_{SC} - V_{FS}$ 

- ullet  $V_{SG}$ : velocity of the Sun relative to the Galactic center.
- V<sub>ES</sub>: velocity of the center of mass of the Earth relative to the Sun

To evaluate the Radon transform we had to calculate explicitly the scalar products  $\mathbf{W} \cdot \mathbf{V}_{ES}$  and  $\mathbf{W} \cdot \mathbf{V}_{SG}$  in a defined **reference frame**.

#### Assumptions and choice of parameters

#### **Astrophysics:**

$$ho=0.3~{
m GeV}~c^{-2}cm^{-3}$$
 (Dark Matter density) ;  $v_0=220~km~s^{-1}$ ;

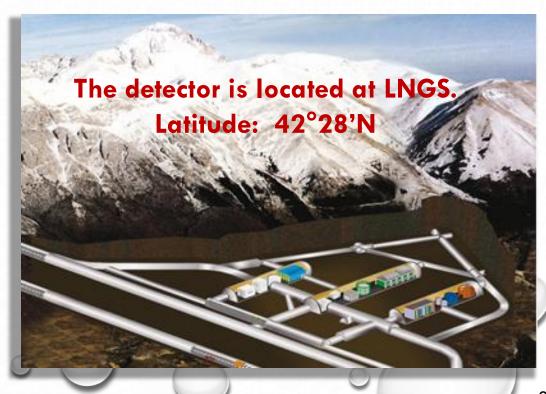


$$\sigma_{W-n}=10^{-46}~cm^2$$
 (Wimp-nucleon cross section) ;  $M_W=200~GeV~c^{-2}$  (Wimp mass) ;

#### **Detector:**

$$E_{min}^{th}=50\ keV$$
 (Minimum threshold Energy) ;  $E_{max}^{th}=200\ keV$  (Maximum Energy) ; Exposure=  $100\ tonne$  year ;





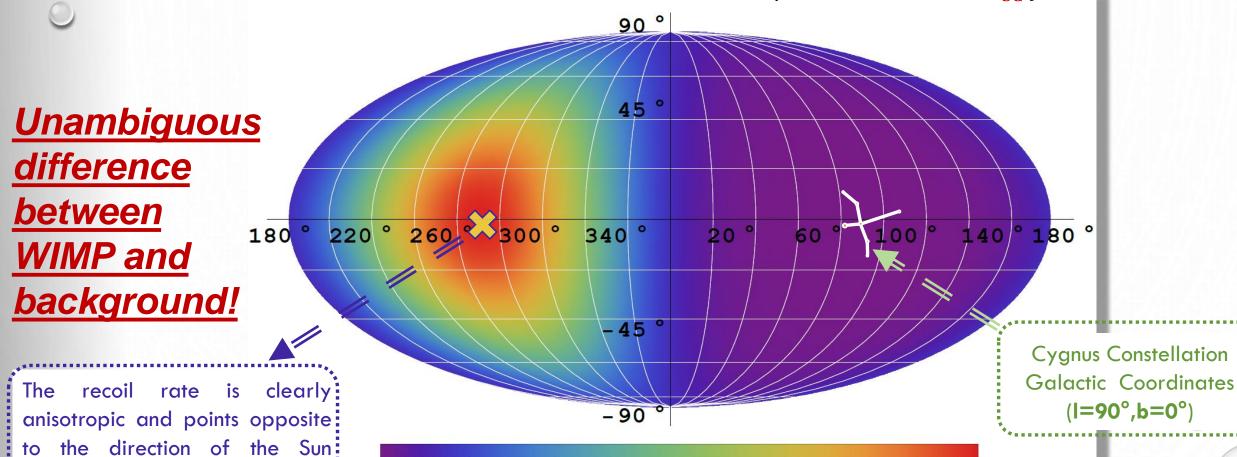
## Wimp directionality in Galactic frame

motion  $(I = 270^{\circ}; b = 0^{\circ}).$ 

 $\frac{dR}{d\cos\vartheta\,d\varphi} = \int_{50\;keV}^{200\;keV} \frac{dR}{dE_r d\cos\vartheta\,d\varphi} dE_r$ 

0.06

In the Galactic coordinate system x points from the Sun towards the Galactic center, y in the direction of the Solar motion and z towards the Galactic north pole; therefore,  $V = V_{SG} y$ .



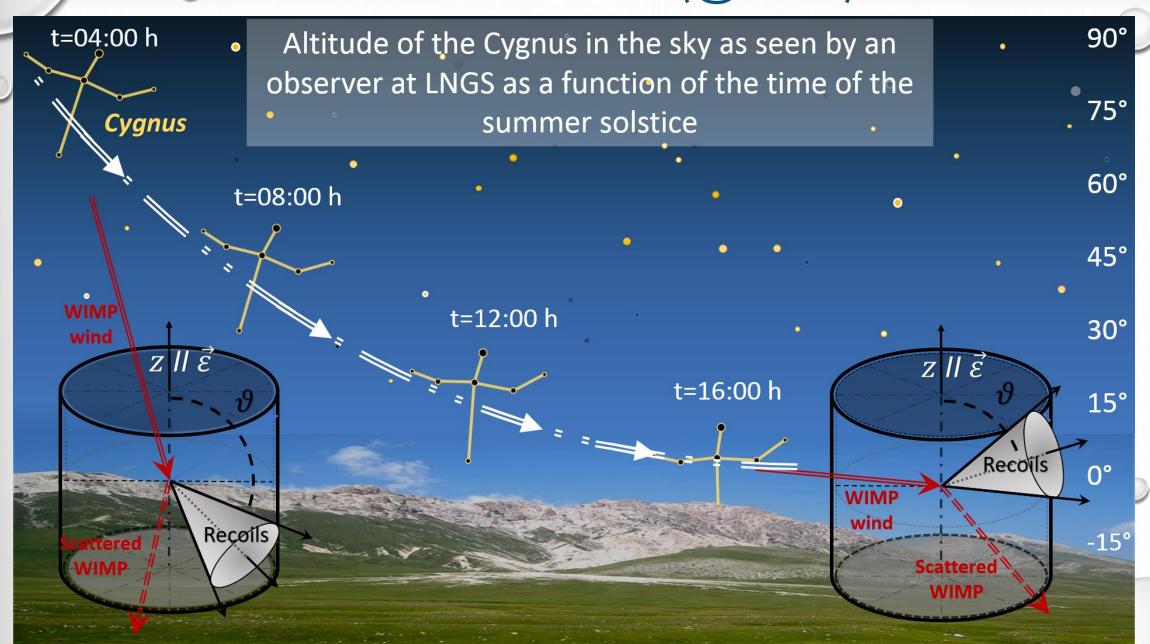
Angular distribution of Argon recoils in Mollweide equal area projection maps of the celestial sphere in galactic coordinates (I, b)

Events/(100 tonne day sr)

0.04

0.02

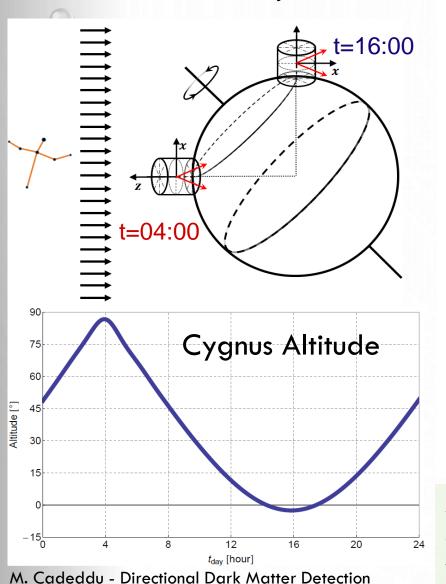
## As seen in a Earth-bound detector (@LNGS)

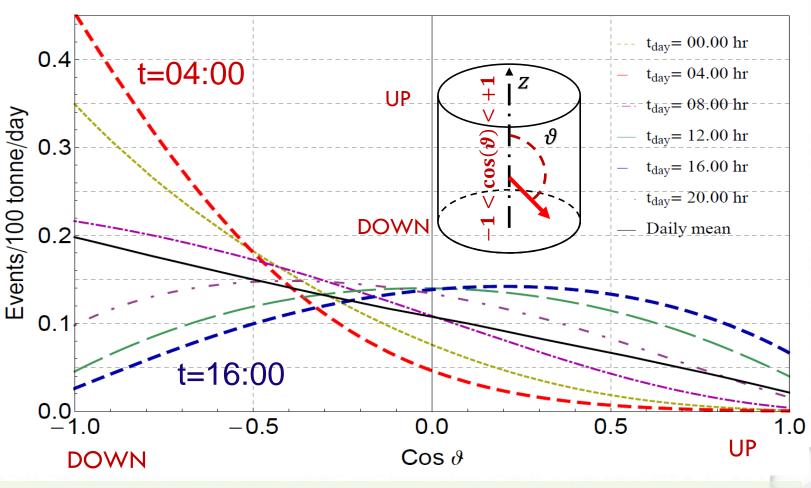


## As seen in the laboratory (LNGS)

\*The azimuthal  $\phi$  angle is integrate  $[0, 2\pi]$ 

For an Earth-bound laboratory the velocity **V** can be decomposed as  $V = -V_{SG} - V_{ES}$ , where  $V_{ES}$  is  $V = -V_{SG} + V_{ES} + V_{$ 

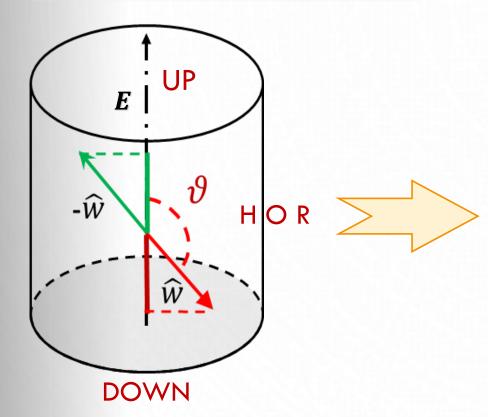




Strong angular dependence of the event rate with respect to the z-axis of the detector as a function of the time of the day. The dependence remains also when it is mediated over the full day (black line: "Daily average").

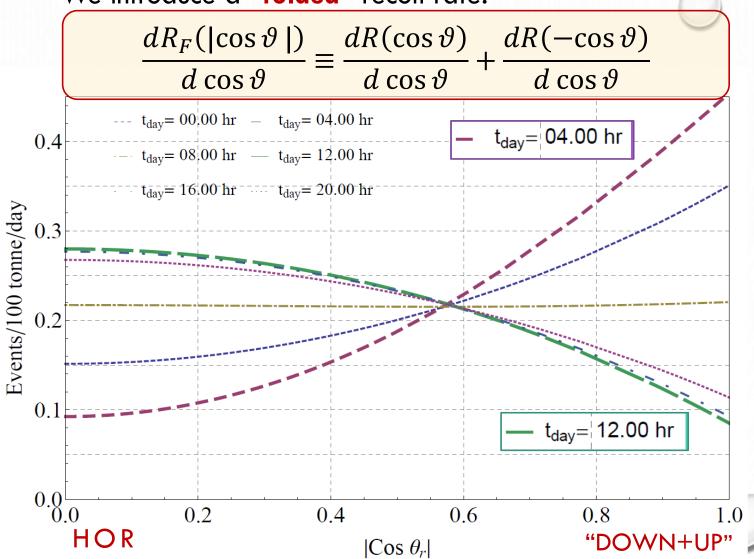
#### Columnar recombination detector

Recoils at 180° give the same signal in detectors based on CR



Columnar Recombination may display a sensitivity to the angle between nuclear recoil direction and drift field *E* in a LAr TPC.

We introduce a "folded" recoil rate:



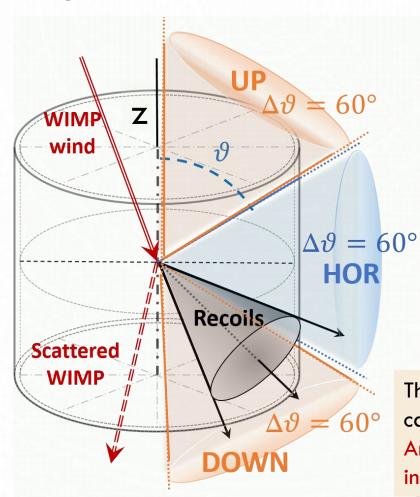
Strong angular dependence of the event rate with respect to the z-axis of the detector as a function of the time of the day

#### The division of the events

#### In case of **minimal** angular resolution

A way to categorize the events is to introduce the UP events corresponding to ( $\vartheta < 120^\circ$ ). The horizontal events (HOR), corresponding to  $|\cos\vartheta\>| < 0.5~(60^\circ < \vartheta < 120^\circ)$  and DOWN events ( $\vartheta > 120^\circ$ ).

A CR based detector should be unable to discriminate between UP and DOWN. We introduce the vertical events (VER), corresponding to  $|\cos\vartheta| > 0.5$  ( $\vartheta < 60^\circ$  and  $\vartheta > 120^\circ$ ).



#### Ratio HOR/VER

From an experimental point of view, ratios of event rates are useful quantities in order to keep the systematic uncertainties budget under control.

An interesting observables is the ratio between HOR and VER events

$$R = \frac{HOR}{UP + DOWN} = \frac{HOR}{VERTICAL}$$

This ratio should be exactly equal to one in the case of isotropic signal.

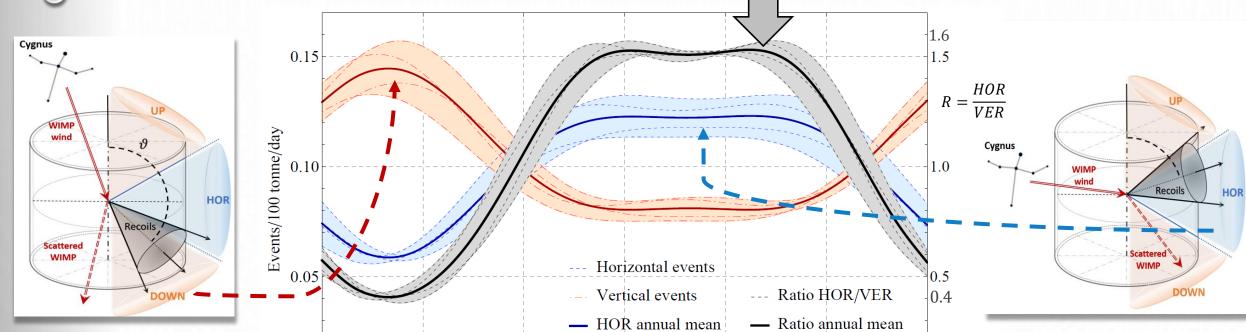
Any significant deviation would be a strong indication in favor of a genuine WIMP signal.

## Event ratio Horizontal/Vertical

$$R = \frac{HOR}{UP + DOWN} = \frac{HOR}{VERTICAL}$$

 $0.00^{\perp}_{0}$ 

The ratio of horizontal to vertical events is shown by the black line. It exhibits a huge variation of *a factor 4* during the day.



Cygnus close to the zenith: vertical events are greater than horizontal ones at the beginning of the day, until 8:00 a.m.

The single horizontal (vertical) component shows a huge variation within a sidereal day, more precisely a 38% (30%) effect with respect to the average.

12

 $t_{\rm day}$  [hour]

— VER annual mean

16

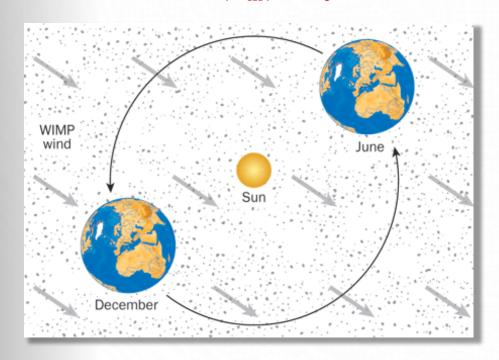
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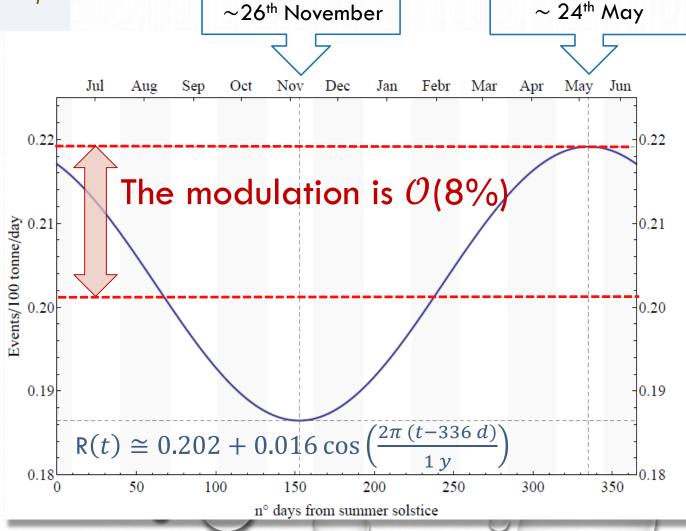
Cygnus close to the horizon: horizontal events are greater than vertical ones in the middle of the day, after 8:00 a.m.

## Annual rate modulation without directionality

$$R(t) = \int_0^{2\pi} \int_{-1}^1 \int_{E_{th}=50 \ keV}^{E_{th}=200 \ keV} \frac{dR}{dE_r d\cos\theta \ d\varphi} dE_r d\cos\theta \ d\varphi$$

The amplitude of the modulation is small and, up to the first approximation, the event rate can be written as a Taylor series  $R(t)\cong R_0+R_m\cos\left(\frac{2\pi\,(t-t_0)}{T}\right)$ ; with the condition that  $|R_m|\ll R_0$ 

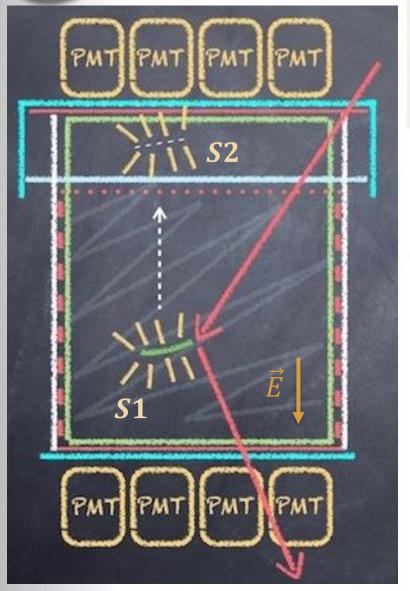




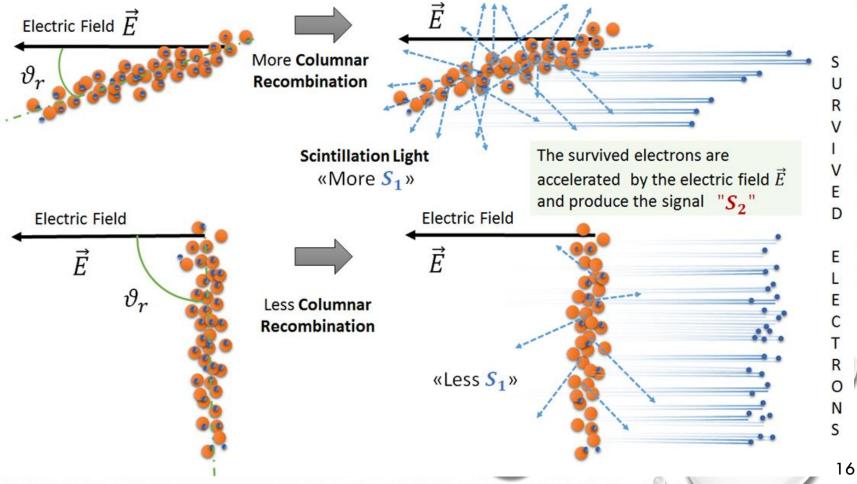
Minimum at

Maximum at

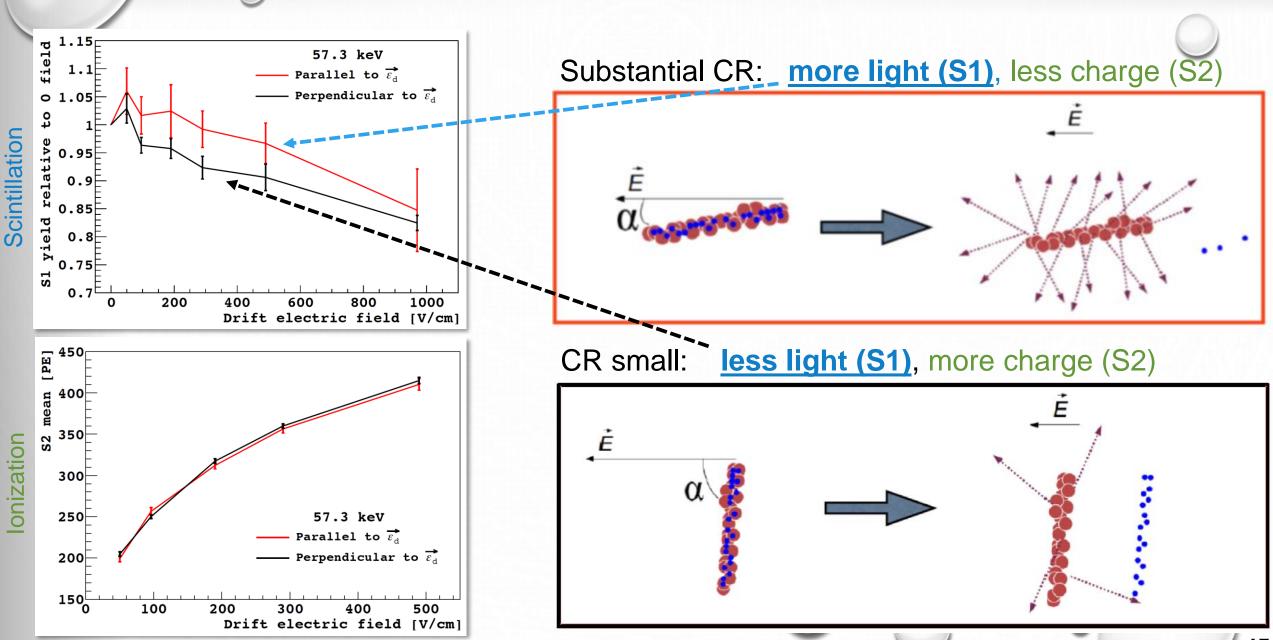
## Two-phase LAr detector: S1 and S2 signal



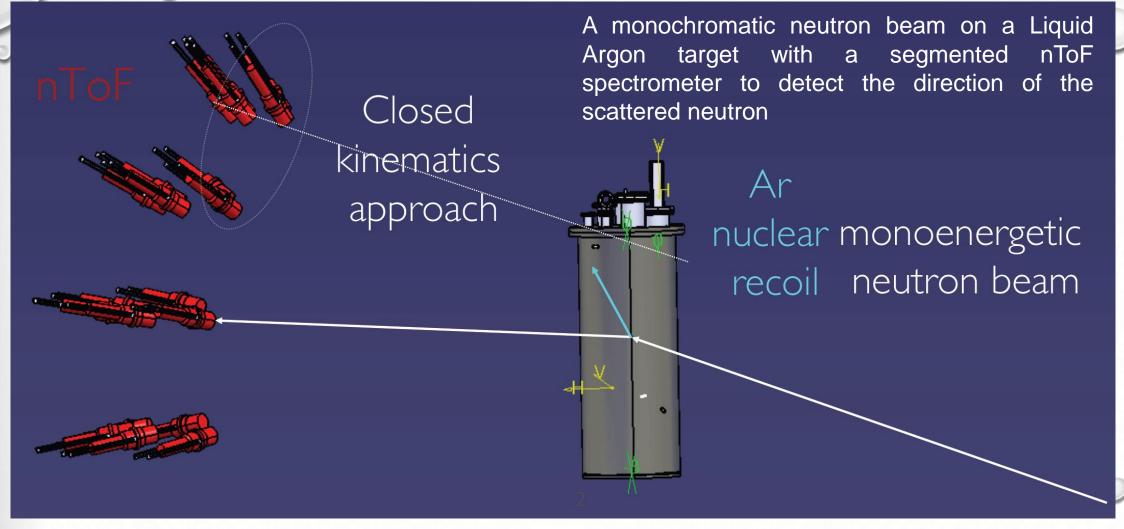
The basic idea of Columnar Recombination: When a nuclear recoil is parallel to the electric field, there will be more electronion recombination since the electrons pass more ions as they drift through the chamber.



## Moving from theory to experiment (SCENE experiment)



#### Essential Setup of the ReD experiment



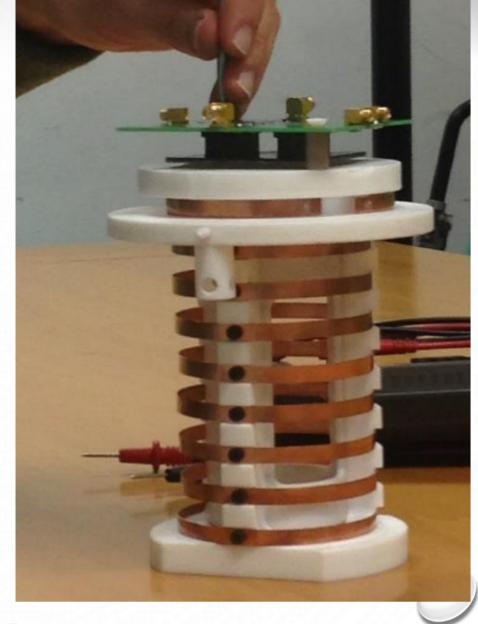
- Monochromatic proton beam on an appropriate target will produce a beam of ~1-4MeV neutrons
- The TPC is placed at a chosen angle to intercept neutrons of the desired energy
- Recoils (10-200 keV) in the LAr produce \$1 and \$2 signals detected by \$iPMs.

#### RED@Unina: the GAP-TPC

- GAP-TPC can have improved light yield using Silicon
   Photomultipliers
- Higher PDE compared to PMTs
- Individual readout of top SiPMs provides improved spatial resolution of the X-Y position of the S2 signal
- Low noise at cryogenic temperatures
- Low radioactivity (for DM application)



The current prototype uses two commercial SensLJ-series TSV arrays of 64 SiPMs



G. Fiorillo, B. Rossi, P. Trinchese, S. Walker

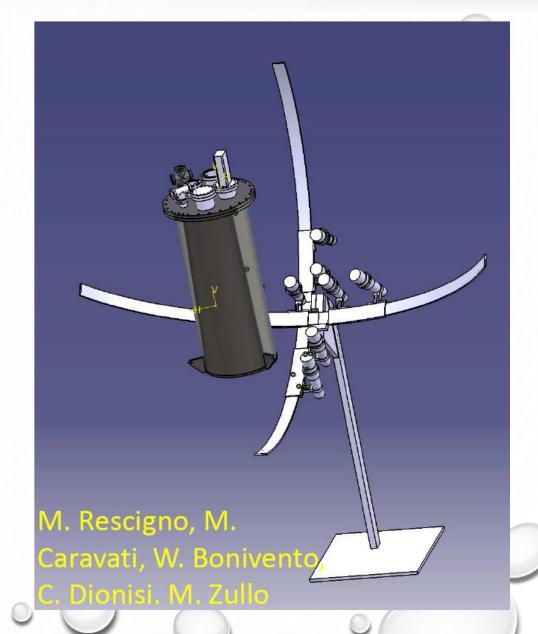
## Timeline of the RED experiment

# Phase 1: aiming at confirming the results obtained by the SCENE experiment

- Data taking campaign on the neutron beam line with a preliminary experimental configuration provisional neutron beam setup;
- GAP-TPC prototype installed in the final cryogenic system; preliminary nToF spectrometer (few LSci detectors).

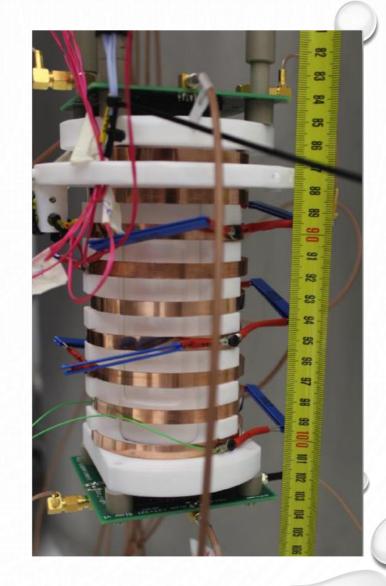
#### Phase 2: directionality studies

- Detectors in final configuration and neutron/gamma data taking optimised beam target + collimator,
- final GAP-TPC with ancillary systems
- full size nToF spectrometer.



## Preliminary Data Taken has already Started





#### Conclusions

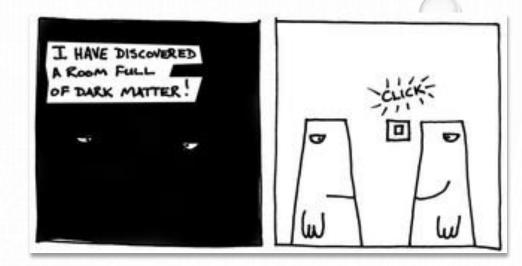
Directional detection is a very promising Dark Matter search strategy thanks to the expected strong direction dependence of the WIMP signal.

- ReD will test the feasibility of using columnar recombination to detect the direction of possible WIMP signals using tonscale liquid Argon TPCs
- Experimental results can be used to verify and/or improve the theoretical models for columnar recombination

Bringing this new tool into the DarkSide program will make the future very exciting!

In the near future, from the phenomenological side:

 A study to explore the impact of direction-sensitivity on the neutrino floor in dark matter direct detection.



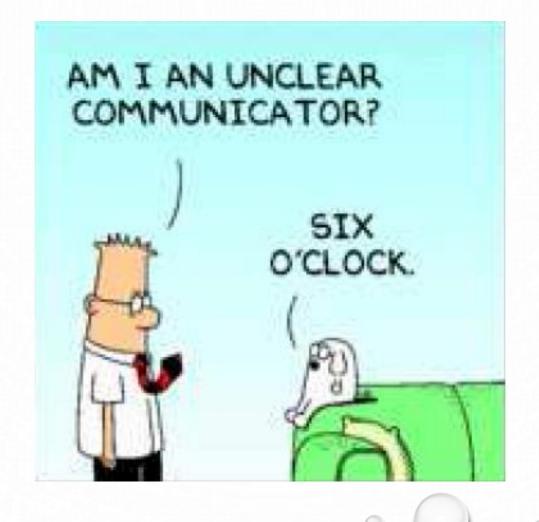
results demonstrate The that directional detectors offer the experimental most promising technique for the future WIMPs discovery and Columnar recombination provides the best possibility for combining directional sensitivity with the ton-scale size of experiment needed to this search.

# Thanks for your attention





# BACKUP



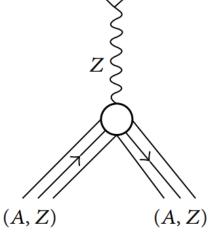
#### Facing the problems of past and future WIMP detection



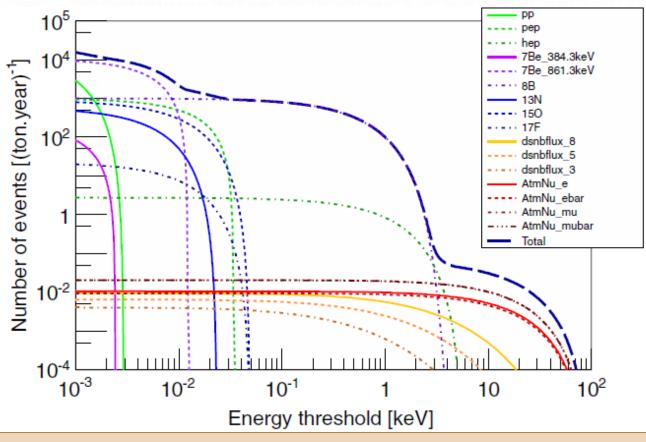
Upcoming direct dark matter detection experiments will have sensitivity to detect neutrinos from several astrophysical sources (Sun, atmosphere, and diffuse Supernovae).

## Coherent neutrino scattering (CNS)

$$\nu_{\chi} + (A, Z) \rightarrow \nu_{\chi} + (A, Z)$$

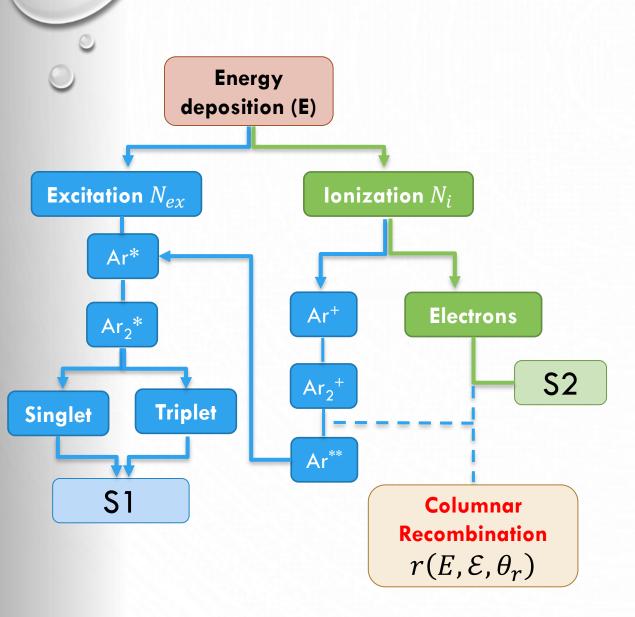


$$\frac{d\sigma^{CNS}(E_{\nu}, E_{r})}{dE_{r}} = \frac{G_{f}^{2}}{4\pi} Q_{w}^{2} m_{N} \left(1 - \frac{m_{N} E_{r}}{2E_{\nu}^{2}}\right) F^{2}(E_{r})$$



Number of neutrino-induced nuclear recoils per ton-year for Xe target as a function of the energy threshold

### Simulating \$1 and \$2 signal with the Columnar Recombination



A simple first-order three-parameters approximation that maintains the main features of more detailed Columnar Recombination theoretical models is

$$\frac{r_0}{r(E, \mathcal{E}, \theta_r)} = 1 + \frac{\mathcal{E}}{\mathcal{E}_0} (\xi + \sin \theta_r)$$

where the intrinsic recombination probability in absence of field,  $r_0$ , the electric field scale,  $\mathcal{E}_0$ , and the relative strength of the isotropic recombination,  $\xi$ , depend on the energy deposit E.

This parameterization ensures that the recombination probability satisfies  $0 \le r \le r_0 \le 1$  and that r is linear in the component of the field orthogonal to the track,  $\mathcal{E}\sin\theta_r$ , to first order.

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08/07/2016

## Simulating \$1 and \$2 signal

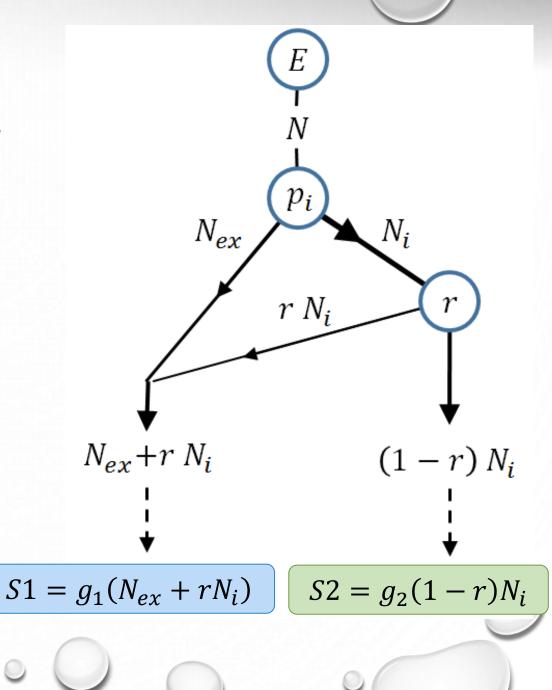
We study the possibility of discriminating events corresponding to different **recombination probability**, **r**, in the plane (S2, S1).

It is already very interesting to distinguish recoils that are parallel to the electric field  $\mathcal{E}$  (vertical events,  $\sin\theta_r=0$ ) from recoils that are perpendicular to  $\mathcal{E}$  (horizontal events,  $\sin\theta_r=1$ ). We consider two recombination probabilities: for parallel,  $r_{\parallel}$ , and perpendicular  $r_{\perp}$ , recoils.

$$\frac{r_0}{r(E, \mathcal{E}, \theta_r)} = 1 + \frac{\mathcal{E}}{\mathcal{E}_0} (\xi + \sin \theta_r)$$

with some typical field value  $\mathcal{E} = \mathcal{E}_0$ , we use

$$r_{\parallel} = \frac{r_0}{1+\xi} = 0.4$$
 where  $r_0 = 0.8$  and  $\xi = 1$ .  $r_{\perp} = \frac{r_0}{2+\xi} \approx 0.27$ 



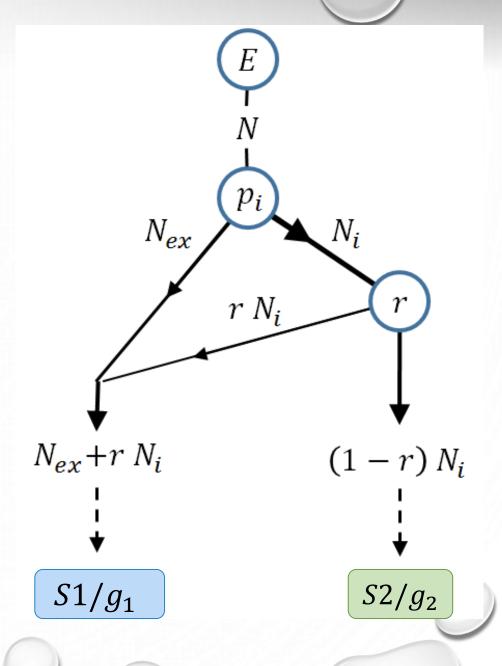
## Simulating \$1 and \$2 signal

#### **Steps of the simulations**

- 1. A number N of primaries is extracted from a Poissonian distribution with mean N.
- 2. These primaries are split into  $N_{ex}$  primary scintillation photons and Ni electrons according to a binomial distribution with ionization probability  $p_i = 0.5$ .
- 3. The  $N_i$  electrons split on average into  $rN_i$  excitons from recombination and  $N_{ie} = (1-r) \, N_i$  remaining electrons according to a binomial distribution with recombination probabitities  $r_{\parallel} = 0.4$  for vertical recoils and  $r_{\perp} = 0.27$  for horizontal recoils. The exictons from recombination contribute to the total number of scintillation photons

$$N_{ph} = N_{ex} + r N_i$$

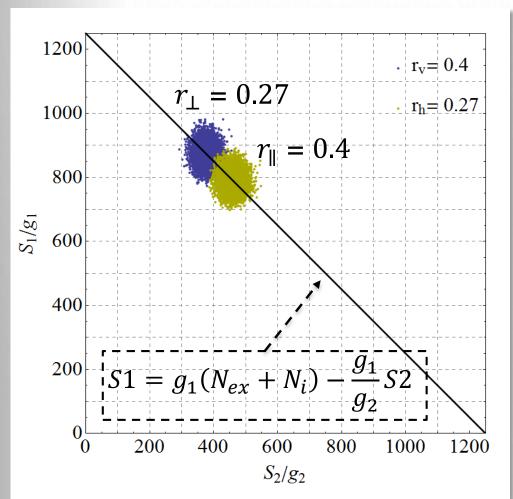
The resulting signals are  $S1/g_1$  and  $S2/g_2$ 

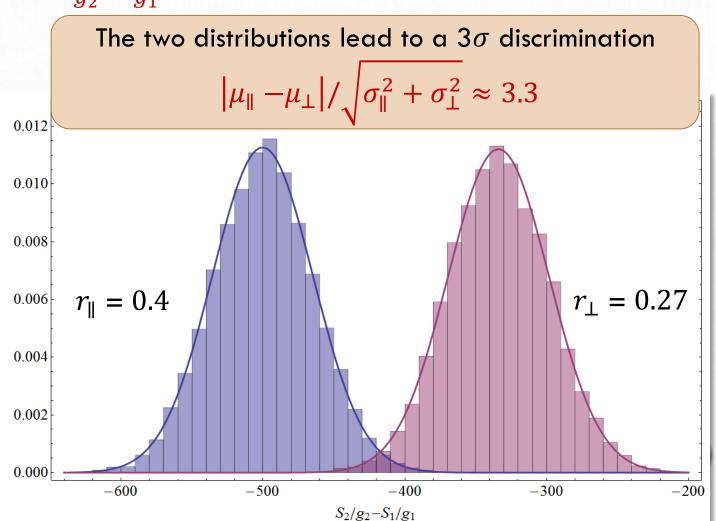


## Simulating \$1 and \$2 signal

Vertical and horizontal events are very well discriminated. On the right figure the distributions of the same events are shown as a function of  $\Delta_{2-1} = \frac{S2}{g_2} - \frac{S1}{g_1}$ , the best variable to discriminate different

values of r.





#### Seasonal effects

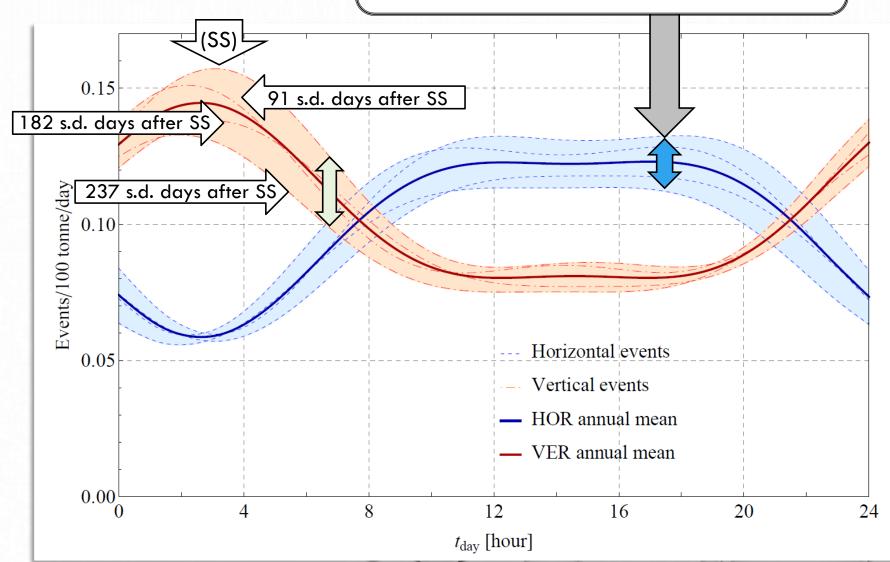
The overall variation registered in the HOR components among the different sidereal days of the year is shown by the blue band.

#### Sidereal day VS solar day

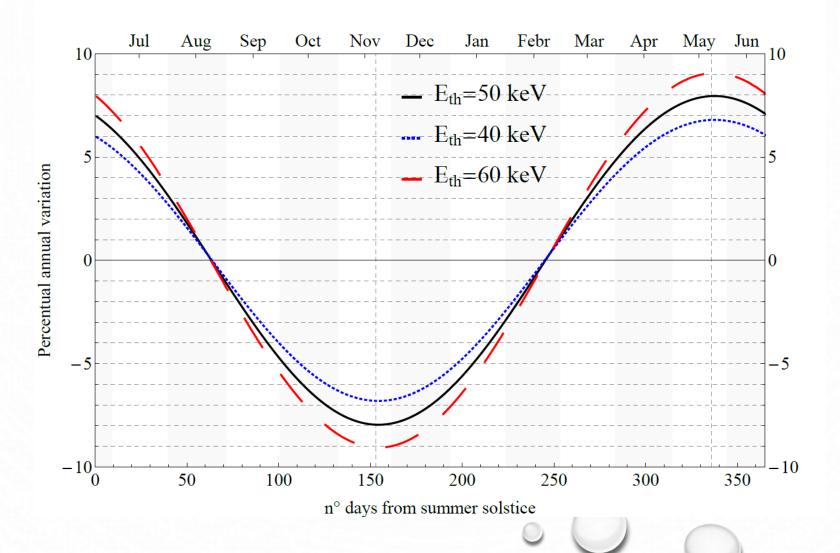
Use of the **sidereal day** (s.d.) as a unit is convenient because after a sidereal day the Cygnus constellation returns exactly in the same position in the sky

#### **Seasonal variation**

Here, the different colored dotted lines refer to the rate evaluated at summer solstice (\$\$) and \$91, 182 and 237 sidereal days after it, while the solid black line is the annual mean.

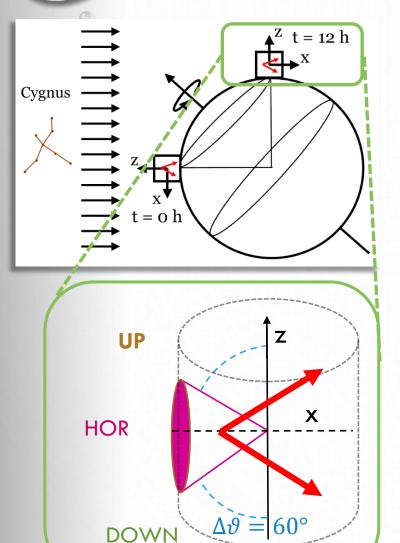


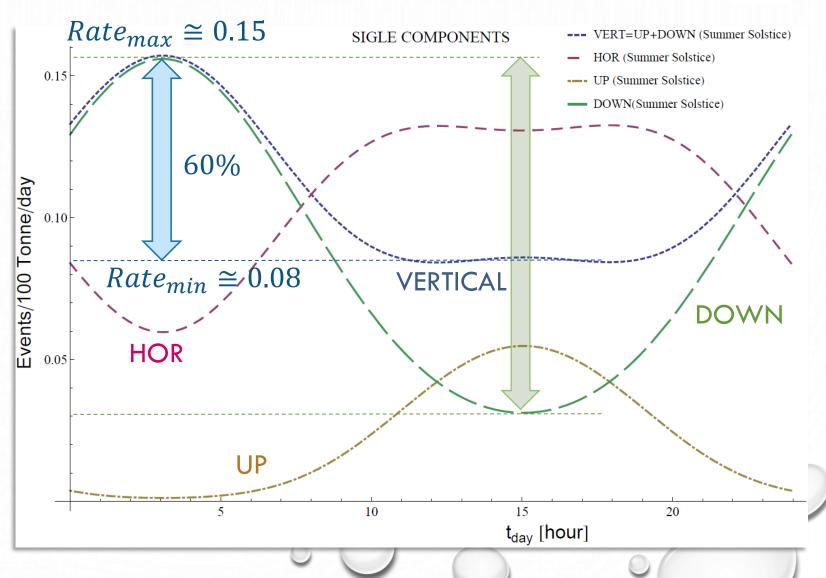
#### \*Energy threshold effects on annual modulation



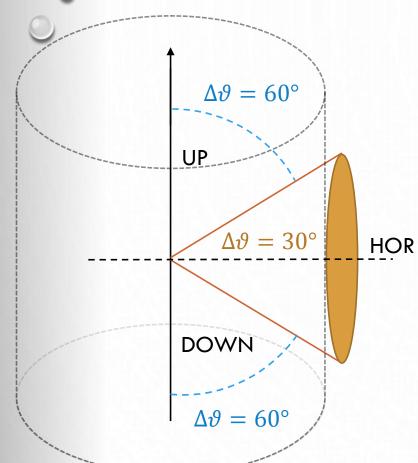
## Diurnal modulation (with directionality)

\* The  $\phi$  angle is integrate  $[0, 2\pi]$ 

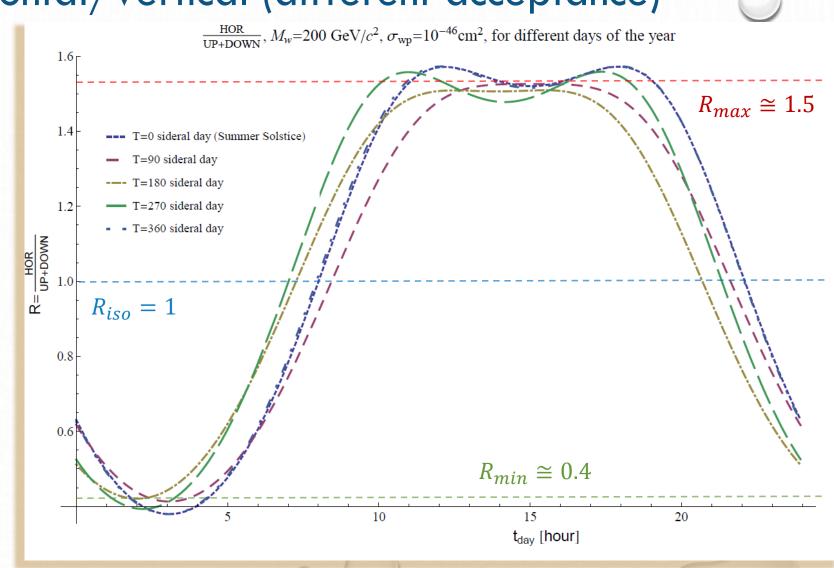




## Event ratio Horizontal/Vertical (different acceptance)



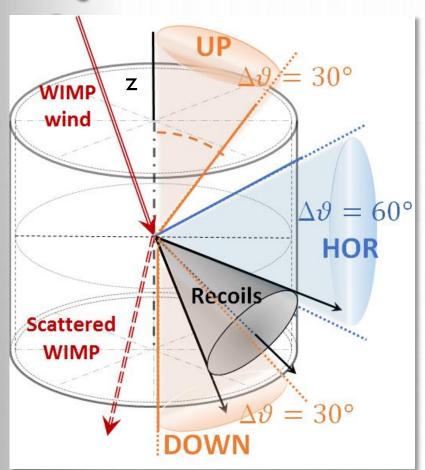
$$R = \frac{HOR}{UP + DOWN} = \frac{HOR}{VERTICAL}$$

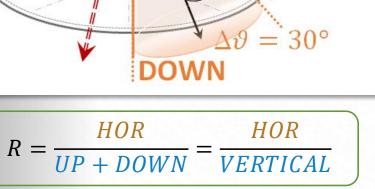


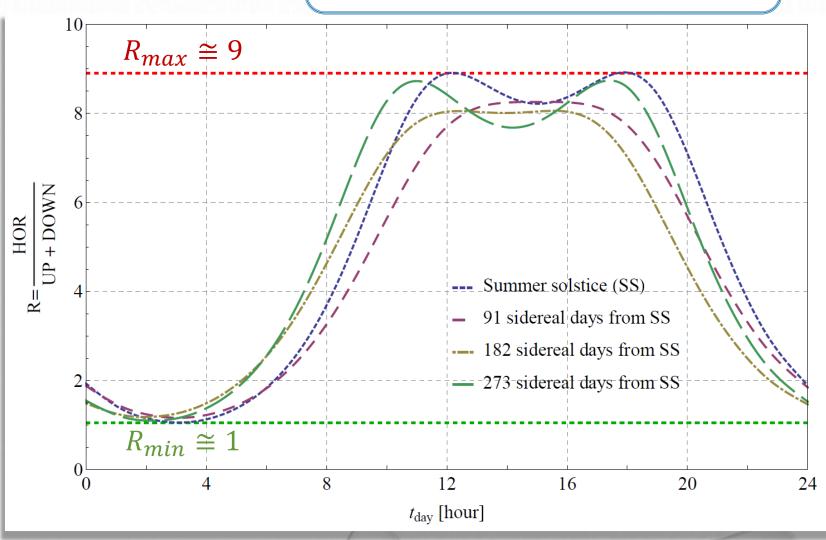
08/07/2016

## Angular resolution effect ratio HOR/VER

Ratio of horizontal WIMPs induced Ar recoils to vertical ones varies by a factor  $f \sim 10$  over the day (acceptance  $\pm 30^{\circ}$ )

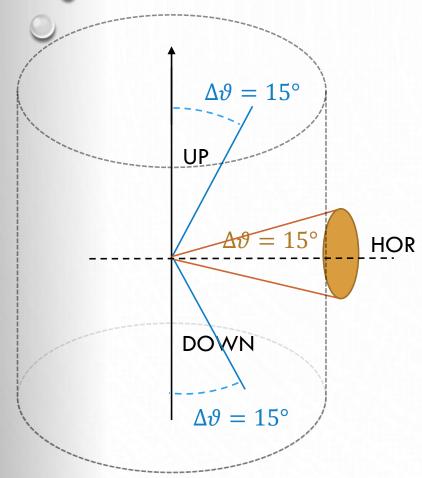






## Event ratio Horizontal/Vertical (±15° acceptance)





$$R = \frac{HOR}{UP + DOWN} = \frac{HOR}{VERTICAL}$$

