

Effective actions for the high energy scattering in QCD and gravity

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1 Gluon reggeization

QCD Born amplitude at high energies $s \gg t$

$$M|_{Born} = 2g T_{A'A}^c \delta_{\lambda_{A'}\lambda_A} \frac{s}{t} g T_{B'B}^c \delta_{\lambda_{B'}\lambda_B}, [T^a, T^b] = if_{abc}T^c$$

Regge behavior in Leading Logarithmic Approximation

$$M(s, t) = M|_{Born} s^{\omega(t)}, \alpha_s \ln s \sim 1, \alpha_s = \frac{g^2}{4\pi} \ll 1$$

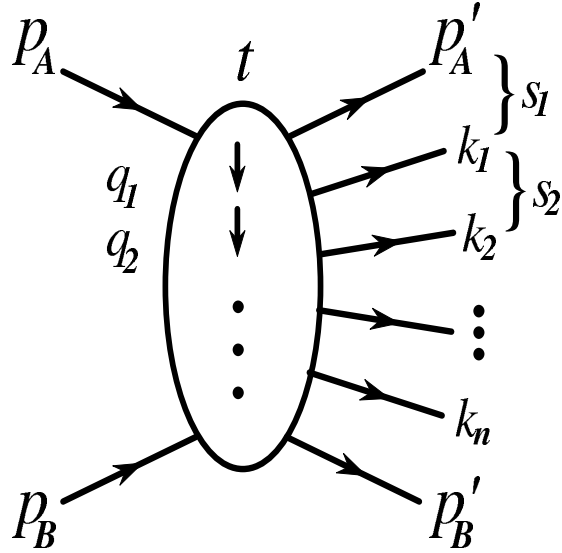
Gluon scattering vertex ($n^+ = 2p_A/\sqrt{s}$, $n^- = 2p_B/\sqrt{s}$)

$$\gamma_{\mu'\mu}^B = -\delta_{\mu\mu'} + p_{\mu'} \frac{n_{\mu}^+}{p^+} + p'_{\mu} \frac{n_{\mu'}^+}{p^+} + q^2 \frac{n_{\mu}^+ n_{\mu'}^+}{2(p^+)^2} \rightarrow \delta_{\lambda'\lambda}$$

Gluon Regge trajectory in LLA

$$\omega(-|q|^2) = -\frac{\alpha_s N_c}{4\pi^2} \int d^2k \frac{|q|^2}{|k|^2 |q-k|^2} \approx -\frac{\alpha_s N_c}{2\pi} \ln \frac{|q|^2}{\lambda^2}$$

2 Gluon production at high energies



$$M_{2 \rightarrow 2+n}^{FKL} \sim \frac{s_1^{\omega_1}}{|q_1|^2} g T_{c_2 c_1}^{d_1} C(q_2, q_1) \frac{s_2^{\omega_2}}{|q_2|^2} \dots g T_{c_{n+1} c_n}^{d_n} C(q_{n+1}, q_n) \frac{s_{n+1}^{\omega_{n+1}}}{|q_{n+1}|^2},$$

$$C_\mu = -q_{1\mu}^\perp - q_{2\mu}^\perp + n_\mu^+ \left(k_1^- + \frac{q_1^2}{k_1^+} \right) - n_\mu^- \left(k_1^+ + \frac{q_2^2}{k_1^-} \right) \rightarrow C(q_2, q_1) = \frac{q_2^\perp q_1^{\perp*}}{k_1^{\perp*}}$$

3 BFKL Pomeron

Balitsky-Fadin-Kuraev-Lipatov equation (1975)

$$E \Psi(\vec{\rho}_1, \vec{\rho}_2) = H_{12} \Psi(\vec{\rho}_1, \vec{\rho}_2), \quad \sigma_t \sim s^\Delta, \quad \Delta = \frac{4\alpha N_c}{\pi} \ln 2$$

BFKL Hamiltonian in the operator form

$$H_{12} = \frac{1}{p_1 p_2^*} \ln |\rho_{12}|^2 p_1 p_2^* + \frac{1}{p_1^* p_2} \ln |\rho_{12}|^2 p_1^* p_2 + \ln |p_1 p_2|^2 - 4\psi(1),$$

Holomorphic separability of H_{12} (L. (1986))

$$H_{12} = h_{12} + h_{12}^*, \quad [h_{12}, h_{12}^*] = 0, \quad \rho_{12} = \rho_1 - \rho_2, \quad \rho_r = x_r + iy_r$$

Its Möbius invariance (L. (1986))

$$\rho_k \rightarrow \frac{a\rho_k + b}{c\rho_k + d}, \quad \Psi = \left(\frac{\rho_{12}}{\rho_{10}\rho_{20}} \right)^m \left(\frac{\rho_{12}^*}{\rho_{10}^*\rho_{20}^*} \right)^{\tilde{m}}, \quad m = \gamma + \frac{n}{2}, \quad \gamma = \frac{1}{2} + i\nu$$

4 BKP equation in LLA

Bartels-Kwiecinski-Praszalowicz equation (1980)

$$E \Psi(\vec{\rho}_1, \dots, \vec{\rho}_n) = H \Psi(\vec{\rho}_1, \dots, \vec{\rho}_n), \quad H = \sum_{k < l} \frac{\vec{T}_k \vec{T}_l}{-N_c} H_{kl}$$

Holomorphic separability at large N_c (L. (1988))

$$H = \frac{1}{2} (h + h^*), \quad [h, h^*] = 0, \quad h = \sum_{k=1}^n h_{k,k+1}$$

Monodromy matrix and integrability (L. (1993))

$$\prod_{k=1}^n \begin{pmatrix} u + \rho_k p_k & p_k \\ -\rho_k^2 p_k & u - \rho_k p_k \end{pmatrix} = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix}, \quad [h, A(u) + D(u)] = 0$$

Integrability of equations for adjoint composite states (L. (2009))

5 Gluon and reggeized gluon fields

Locality of interactions in the rapidity space

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Gauge transformations for quark and gluon fields

$$\psi(x) \rightarrow U\psi(x), \quad v_\mu(x) \rightarrow \frac{1}{g}UD_\mu U^{-1}, \quad D_\mu = \partial_\mu + gv_\mu(x)$$

Gauge invariance of reggeized gluons

$$\delta A^\pm(x) = 0, \quad U(\infty) = 1, \quad A^\pm = A^0 \pm A^3$$

Kinematical constraints for reggeon fields

$$\partial_\pm A^\pm(x) = 0$$

6 Effective theory for high energy QCD

Lagrangian for reggeized gluon interactions (L. (1995))

$$L_{eff} = L_{QCD} + Tr(V_+ \partial_\mu^2 A^+ + V_- \partial_\mu^2 A^-) + 2Tr \partial_\sigma^\perp A_+ \partial_\sigma^\perp A_-$$

Eikonal representation for effective currents

$$V_\pm = \frac{1}{g} \partial_\pm O(x^\pm), \quad O(x^\pm) = -\frac{1}{D_\pm} \overleftarrow{\partial}_\pm$$

Principal value prescription for propagators

$$\frac{1}{D_\pm} = P \frac{e^{-\frac{g}{4} \int_{-\infty}^{x^\pm} dx^\pm v_\pm}}{e^{-\frac{g}{4} \int_{x^\pm}^{\infty} dx^\pm v_\pm}} \frac{1}{\partial_\pm} \bar{P} \frac{e^{-\frac{g}{4} \int_{x^\pm}^{\infty} dx^\pm v_\pm}}{e^{-\frac{g}{4} \int_{-\infty}^{x^\pm} dx^\pm v_\pm}}$$

Effective currents in terms of P -exponents

$$O(x^\pm) = P \frac{e^{-\frac{g}{4} \int_{-\infty}^{x^\pm} d\tilde{x}^\pm v_\pm}}{e^{-\frac{g}{4} \int_{x^\pm}^{\infty} d\tilde{x}^\pm v_\pm}} \frac{P e^{\frac{g}{4} \int_{-\infty}^{\infty} d\tilde{x}^\pm v_\pm} + \bar{P} e^{-\frac{g}{4} \int_{-\infty}^{\infty} d\tilde{x}^\pm v_\pm}}{2}$$

7 Classical equations for effective QCD

Euler-Lagrange equation for high energy QCD

$$[D_\mu, G^{\mu\nu}] = j^\nu, \quad j^\pm = O(x^\pm)(\partial_\sigma^2 A^\pm)O^+(x^\pm), \quad j_\perp^\nu = 0$$

Classical equation for a quasi-elastic kinematics

$$[D_\mu, G^{\mu\nu}] = O(x^+) \partial_{\perp\sigma}^2 A^+(x^-, x_\perp) O^+(x^+) \delta_+^\nu$$

Transformation to the light-cone gauge $v'_+ = 0$

$$v'_\mu = V^{-1}(v_+)(v_\mu + \frac{\partial_\mu}{g})V(v_+), \quad V(v_+) = P \frac{e^{-\frac{g}{4} \int_{-\infty}^{x^\pm} d\tilde{x}^\pm v_\pm}}{e^{-\frac{g}{4} \int_{x^\pm}^{\infty} d\tilde{x}^\pm v_\pm}}$$

Classical solution as a superposition of shock waves

$$\tilde{v}_\nu(x) = \delta_\nu^- \int \frac{d^2 z}{4\pi} \ln(|x - z|^2) \partial_\sigma^{\perp 2} A^+(x^-, z^\perp) = \delta_\nu^- A^+(x^-, x_\perp)$$

8 BFKL equation in $N = 4$ SUSY

BFKL kernel eigenvalue in two loops (F., L. (1998))

$$\omega = 4 \hat{a} \chi(n, \gamma) + 4 \hat{a}^2 \Delta(n, \gamma), \quad \hat{a} = g^2 N_c / (16\pi^2), \quad \gamma = i\nu + 1/2$$

Hermitian separability in $N = 4$ SUSY (K.,L. (2000))

$$\Delta(n, \gamma) = \phi(M) + \phi(M^*) - \frac{\rho(M) + \rho(M^*)}{2\hat{a}/\omega}, \quad M = \gamma + \frac{|n|}{2},$$

$$\rho(M) = \beta'(M) + \frac{1}{2}\zeta(2), \quad \beta'(z) = \frac{1}{4} \left[\Psi'\left(\frac{z+1}{2}\right) - \Psi'\left(\frac{z}{2}\right) \right]$$

Maximal transcendentality (K.,L. (2002))

$$\phi(M) = 3\zeta(3) + \Psi''(M) - 2\Phi(M) + 2\beta'(M) \left(\Psi(1) - \Psi(M) \right),$$

$$\Phi(M) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k+M} \left(\Psi'(k+1) - \frac{\Psi(k+1) - \Psi(1)}{k+M} \right)$$

9 Pomeron and graviton in N=4 SUSY

Diffusion approximation for the BFKL kernel

$$j = 2 - \Delta - \Delta \nu^2, \quad \gamma = \frac{j}{2} + i\nu$$

AdS/CFT relation with the graviton Regge trajectory

$$j = 2 + \frac{\alpha'}{2} t, \quad t = E^2/R^2, \quad \alpha' = \frac{R^2}{2} \Delta, \quad \lambda = g^2 N_c$$

Large coupling expansion (KLOV, BPST, KL, GKL)

$$\gamma = 1 - \sqrt{1 + (j - 2)/\Delta}, \quad \Delta = 2\lambda^{-1/2} + \lambda^{-1} - 1/4 \lambda^{-3/2} - 2(1 + 3\zeta_3)\lambda^{-2}$$

Exact expression for the slope of γ (KLOV, V., Basso)

$$\gamma'(2) = -\frac{\lambda}{24} + \frac{1}{2} \frac{\lambda^2}{24^2} - \frac{2}{5} \frac{\lambda^3}{24^2} + \frac{7}{20} \frac{\lambda^4}{24^4} - \frac{11}{35} \frac{\lambda^5}{24^5} + \dots = -\frac{\sqrt{\lambda}}{4} \frac{I_3(\sqrt{\lambda})}{I_2(\sqrt{\lambda})}$$

10 Fields in high energy gravity

Locality in the rapidity space

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Metric tensor and its coordinate transformation

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad \delta g_{\mu\nu} = D_\mu \chi_\nu + D_\nu \chi_\mu$$

Coordinate invariance of reggeized graviton fields

$$\delta A^{\pm\pm}(x) = 0$$

Kinematical constraints for reggeon fields

$$\partial_\pm A^{\pm\pm}(x) = 0$$

11 Reggeized graviton interactions

Action for the high energy gravity (L. 2011)

$$S = -\frac{1}{2\kappa^2} \int d^4x (\sqrt{-g} R + L_{ind})$$

Induced lagrangian

$$L_{ind} = j_{++} \partial_\sigma^2 A^{++} + j_{--} \partial_\sigma^2 A^{--} + \partial_\sigma A^{++} \partial_\sigma A^{--}$$

Perturbative expansion of effective currents

$$j_{\pm\pm} = \partial_\pm j^\mp = h_{\pm\pm} - \left(h_{\rho\pm} - \frac{1}{2} \frac{\partial_\rho}{\partial_\pm} h_{\pm\pm} \right)^2 + \dots$$

Hamilton-Jacobi equation

$$j^\mp = 2x^\mp - \omega^\mp, \quad g^{\mu\nu} \partial_\mu \omega^\pm \partial_\nu \omega^\pm = 0$$

12 Global light-cone time systems

Coordinate transformation to the global light-cone time system

$$g'^{\pm\pm} = 0, \quad g^{\mu\nu} \partial_\mu x'^{\pm} \partial_\nu x'^{\pm} = 0, \quad x'^{\pm} = \omega^{\pm}$$

Coordinate transformation to the system with the global time

$$g'^{00} = 1, \quad g^{\mu\nu} \partial_\mu x^0 \partial_\nu x^0 = 0$$

Schwarzschild metric with the time \tilde{t}

$$dx^2 = \left(1 - \frac{2M}{r}\right) d\tilde{t}^2 - \frac{dr^2}{1 - \frac{2M}{r}} - r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

Schwarzschild metric with the global time (Painleve (1921))

$$dx^2 = \left(1 - \frac{2M}{r}\right) dt^2 + 2\sqrt{\frac{2M}{r}} dt dr - dr^2 - r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

13 Classical equation for effective action

Einstein-Hilbert equation for effective gravity

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu} = \theta_{\mu\nu}, \quad \theta_{\mu\nu} = c_{\mu\nu}\partial_\sigma^2 A^{--} + d_{\mu\nu}\partial_\sigma^2 A^{++}$$

Coordinate transformation to the metrics $g'^{\pm\nu} = \eta^{\pm\nu}$

$$g'^{\rho\sigma} = g^{\mu\nu}\partial_\mu x'^\rho\partial x'^\sigma, \quad T'^{\rho\sigma} = \theta^{\mu\nu}\partial_\mu x''^\rho\partial_\nu x''^\sigma$$

Stress tensor in an arbitrary coordinate system

$$\theta_{\mu\nu} = \partial_\mu x'^+\partial_\nu x'^+\partial_\chi^2 A^{--}(x') + \partial_\mu x'^-\partial_\nu x'^-\partial_\chi^2 A^{++}(x')$$

Solution of equations for a quasi-elastic kinematics

$$\tilde{g}_{\mu\nu}(x) = \eta_{\mu\nu} + \partial_\mu x'^+\partial_\nu x'^+ A^{--}(x'),$$

$$\tilde{g}_{\mu\nu}^{AS} = \eta_{\mu\nu} + a\delta_\mu^+\delta_\nu^+ \ln|x_\perp|\delta(x^+)$$

14 Multi-Regge processes in gravity

Multi-graviton production amplitude (L. (1982))

$$M_{2 \rightarrow 2+n} \sim s^2 \kappa \delta_{\lambda_A \lambda_{A'}} \frac{s_1^{\omega_1}}{|q_1|^2} \kappa C(q_2, q_1) \dots \kappa C(q_{n+1}, q_n) \frac{s_{n+1}^{\omega_{n+1}}}{|q_{n+1}|^2} \kappa \delta_{\lambda_B \lambda_{B'}}$$

Graviton-graviton-reggeized graviton vertex

$$\gamma_{\mu'\nu',\mu\nu}^{++} = \gamma_{\mu'\mu}^+ \gamma_{\nu'\nu}^+ \rightarrow \delta_{\lambda\lambda'}$$

Reggeized graviton-reggeized graviton-graviton vertex

$$\gamma_{\mu\nu} = \gamma_\mu \gamma_\nu - q_1^2 q_2^2 \left(\frac{n^+}{k^+} - \frac{n^-}{k^-} \right)_\mu \left(\frac{n^+}{k^+} - \frac{n^-}{k^-} \right)_\nu \rightarrow C(q_2, q_1)$$

Fulfilment of the Steinmann relation

$$\Delta_{s_1} \Delta_{s_2} M_{2 \rightarrow 3} = 0$$

15 Graviton trajectory at supergravity

One loop graviton Regge trajectory (L. (1982))

$$j = 2 + \omega, \quad \omega(q^2) = \frac{\alpha}{\pi} \int \frac{q^2 d^2k}{k^2(q-k)^2} f(k, q), \quad \alpha = \frac{\kappa^2}{8\pi^2},$$

$$f(k, q) = (k, q-k)^2 \left(\frac{1}{k^2} + \frac{1}{(q-k)^2} \right) - q^2 + \frac{N}{2} (k, q-k)$$

Gravitino action

$$S_{3/2} = -\frac{1}{2} \epsilon^{\mu\nu\rho\sigma} \int d^4x \sum_{r=1}^N \bar{\psi}_\mu^r \gamma_5 \gamma_\nu \partial_\rho \psi_\sigma^r$$

Divergencies of the graviton Regge trajectory

$$\omega(q^2) = -\alpha |q|^2 \left(\ln \frac{|q|^2}{\lambda^2} + \frac{N-4}{2} \ln \frac{|\Lambda|^2}{|q|^2} \right)$$

16 Double-logarithms in gravity

Mellin representation for the scattering amplitude

$$A(s, t) = A_{Born} s^{-\alpha|q|^2 \ln \frac{|q|^2}{\lambda^2}} \Phi(\xi), \quad \Phi(\xi) = \int_{a-i\infty}^{a+i\infty} \frac{d\omega}{2\pi i \omega} \left(\frac{s}{|q|^2} \right)^\omega f_\omega$$

Infrared evolution equation for super-gravity (BLS (2012))

$$f_\omega = 1 + \alpha|q|^2 \left(\frac{d}{d\omega} \frac{f_\omega}{\omega} - \frac{N-6}{2} \frac{f_\omega^2}{\omega^2} \right), \quad \xi = \alpha|q|^2 \ln^2 \frac{s}{|q|^2}$$

Its solution in terms of the parabolic cylinder function

$$\frac{f_\omega}{\omega} = \frac{2}{6-N} \frac{1}{\sqrt{b}} \frac{d}{dx} \ln \left(e^{\frac{x^2}{4}} D_{\frac{6-N}{2}}(x) \right), \quad x = \frac{\omega}{\sqrt{b}}$$

Perturbative expansion

$$\Phi(\xi) = 1 - \frac{N-4}{2} \frac{\xi}{2} + \frac{(N-4)(N-3)}{2} \frac{\xi^2}{4!} - \frac{N-4}{8} (5N^2 - 26N + 36) \frac{\xi^3}{6!} + \dots$$

17 Double-logarithmic eikonal picture

Scattering amplitude in the eikonal approximation

$$A_{DL}(s, t) = -2is s^{-\alpha|q|^2 \ln \frac{|q|^2}{\mu^2}} \int d^2\rho e^{i\vec{q}\vec{\rho}} \left(e^{i\delta_{DL}(\vec{\rho}, \ln s)} - 1 \right),$$

Double logarithmic approximation for the phase (BLS)

$$\delta_{DL}(\vec{\rho}, \ln s) = \frac{s}{2} \frac{\kappa^2}{(2\pi)^2} \int \frac{d^2q}{|q|^2} e^{-i\vec{q}\vec{\rho}} \Phi(\xi)$$

Eikonal phase in $N = 8$ SUSY at small impact parameters

$$\delta_{DL}^{N=8}(\vec{\rho}, \ln s) = \frac{s}{2} \frac{\kappa^2}{(2\pi)^2} \pi \ln \frac{1}{\lambda^2 \alpha \ln^2(\rho^2 s)}, \quad \rho^2 \ll \alpha \ln^2(\rho^2 s)$$

Agreement with exact calculations (BLS)

$$A_4^{N=8} = \frac{\kappa^2 s^2}{|q|^2} (-i\pi s) \alpha^2 |q|^2 \frac{\ln^3 \frac{s}{|q|^2}}{3}$$

18 Discussion

1. Gluon and graviton reggeization
2. Möbius invariance in coordinate and momentum spaces
3. BFKL dynamics and integrability
4. Effective action for high energy QCD
5. Euler-Lagrange equation for effective QCD
6. Pomeron and reggeized graviton in $N = 4$ SUSY
7. Effective action for reggeized gravitons
8. Euler-Lagrange equation for high energy gravity
9. Graviton Regge trajectory in supergravities
10. Scattering amplitudes in the DL approximation.