

Exploring supersymmetry at the LHC using the MoEDAL detector

Vasiliki A. Mitsou (*IFIC – CSIC / Valencia Univ.*)

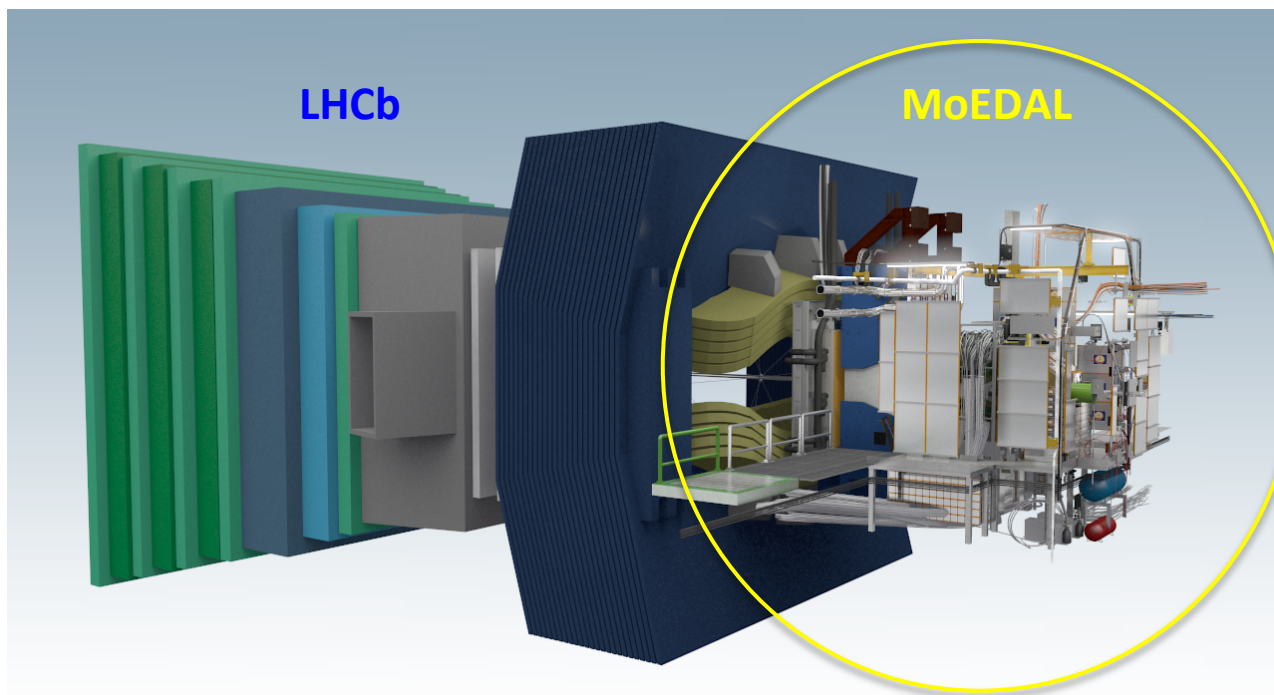
for the MoEDAL Collaboration

5th International Conference on New Frontiers in Physics
ICNFP 2016

6 – 14 July 2016, Kolymbari, Crete, Greece



The MoEDAL detector



DETECTOR SYSTEMS

- ① Low-threshold NTD (LT-NTD) array
 - $Z/\beta > \sim 5$
- ② Very High Charge Catcher NTD (HCC-NTD) array
 - $Z/\beta > \sim 50$
- ③ TimePix radiation background monitor
- ④ Monopole Trapping detector (MMT)

MoEDAL is unlike any other LHC experiment:

- mostly **passive detectors**; no trigger; no readout
- the largest deployment of passive **Nuclear Track Detectors (NTDs)** at an accelerator
- the 1st time **trapping detectors** are deployed as a detector

Beyond magnetic monopoles

- Key feature: **high ionisation**

MoEDAL detectors have a threshold of $z/\beta \sim 5$

$$-\frac{dE}{dx} = K \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta}{2} \right]$$

Electric charge
Bethe-Bloch formula

charge
velocity: $\beta = v/c$ = z/β

- Achieved, e.g., by magnetic monopoles due to ionisation 68.5^2 times higher than minimum ionising particle
- Actually any (meta-)stable massive charged particle, (M)SMCP, (hence slow moving) with electric charge should give a track in Nuclear Track Detectors
- For singly charged particles to be detected in NTDs, velocity should be $\beta \lesssim 0.2$
- Moreover in some cases they may lose all of their momentum, mainly from ionisation energy loss, and come to rest within the magnetic trappers, MMTs
 - if metastable, they may be monitored with a dedicated detector system in an underground laboratory over a long period for their decay

Stable charged particles in BSM scenarios

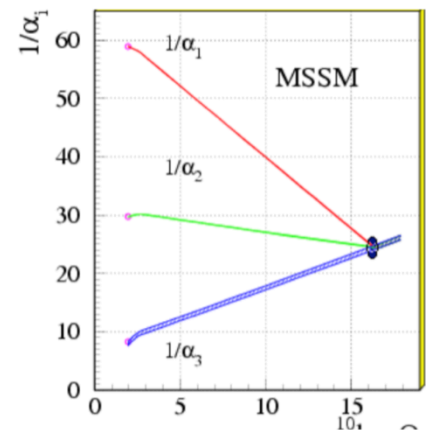
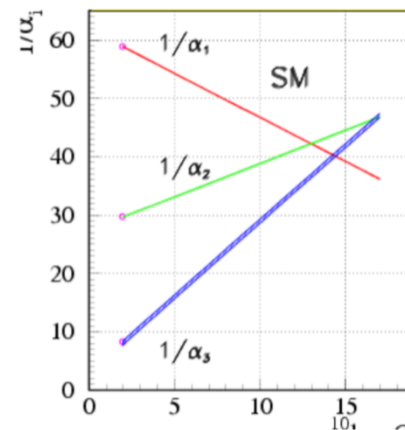
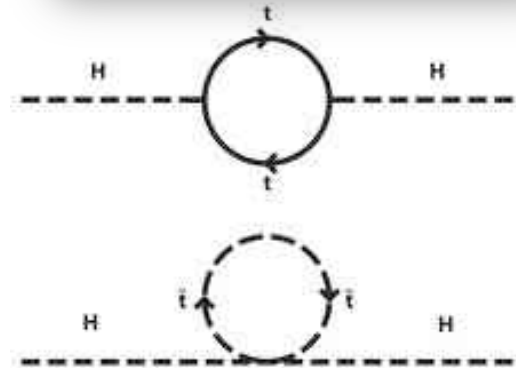
- Scenarios with Extra Dimensions
 - long-lived microscopic black holes
 - microscopic black hole remnants
 - long-lived Kaluza-Klein particles from UED
- D matter
 - electrically-charged D-particles
- Long-lived heavy quarks
- Fourth-generation fermions
- Multiparticle excitations
 - Q-balls
 - strangelets
 - quirks
- ...
- **Supersymmetry**



See previous talk by
Nick Mavromatos

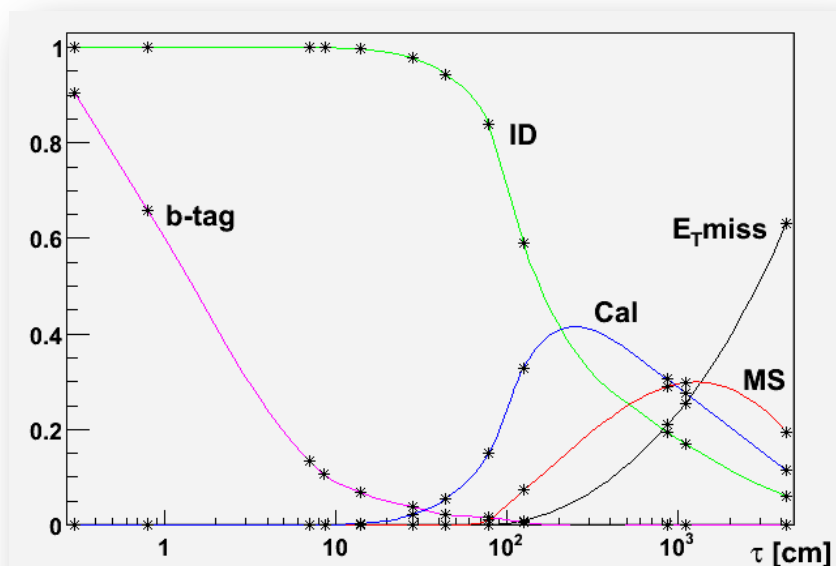
Supersymmetry (SUSY)

- Global symmetry between fermions & bosons
- Motivation
 - Higgs mass stabilisation against loop corrections (**fine-tuning** problem)
 - **unification** of gauge couplings at single scale
 - **dark matter** candidate
- Particle stability mechanisms
 - a) lightest state (LSP) carrying a conserved quantum number:
R-parity: $R = (-1)^{3(B-L)+2s}$
 - b) suppressed (effective) coupling
 - c) lack of phase space for decay, e.g. mass degeneracies



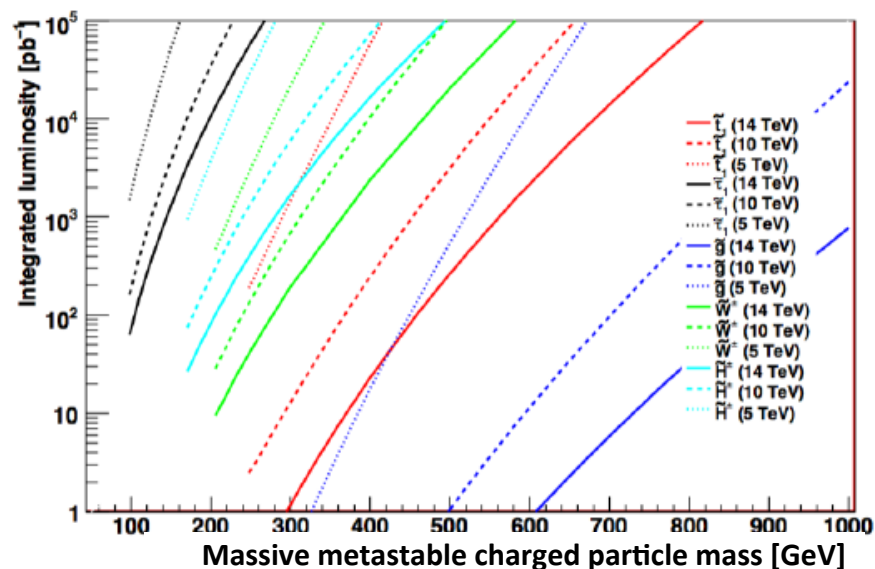
LHC sensitivity to sparticle direct production

- Metastable particles \equiv they live long enough to pass through detector
- Detection in ATLAS and CMS
 - large ionisation energy loss dE/dx , e.g. time-over-threshold in ATLAS Transition Radiation Tracker
 - nuclear interactions (R-hadron) in calorimeters
 - delay (time of flight) reconstructed in muon chambers



Integrated luminosities needed for discovery at LHC at 14 TeV (solid), 10 TeV (dashed) and 5 TeV (dotted)

- signal efficiency of 20% (5%) for electrically charged (strongly interacting) SMCPs
- 1 bkg event for 100 pb^{-1}



Long-lived particles in SUSY scenarios

- GMSB:** NLSP decays to gravitino LSP only via (small) gravitational coupling

- $N_{\text{mes}} = 1$: non-pointing photons

$$\tilde{\chi}_1^0 \rightarrow \tilde{G} + \gamma$$

- $N_{\text{mes}} > 1$: **penetrating sleptons**

$$\tilde{\ell} \xrightarrow{\text{long}} \tilde{G} + \ell$$

- Split SUSY:** squarks heavy, suppressing gluino decays \rightarrow colored heavy particles

R-hadrons

$$R = \tilde{g}q\bar{q}, \tilde{g}qqq, \tilde{g}g$$

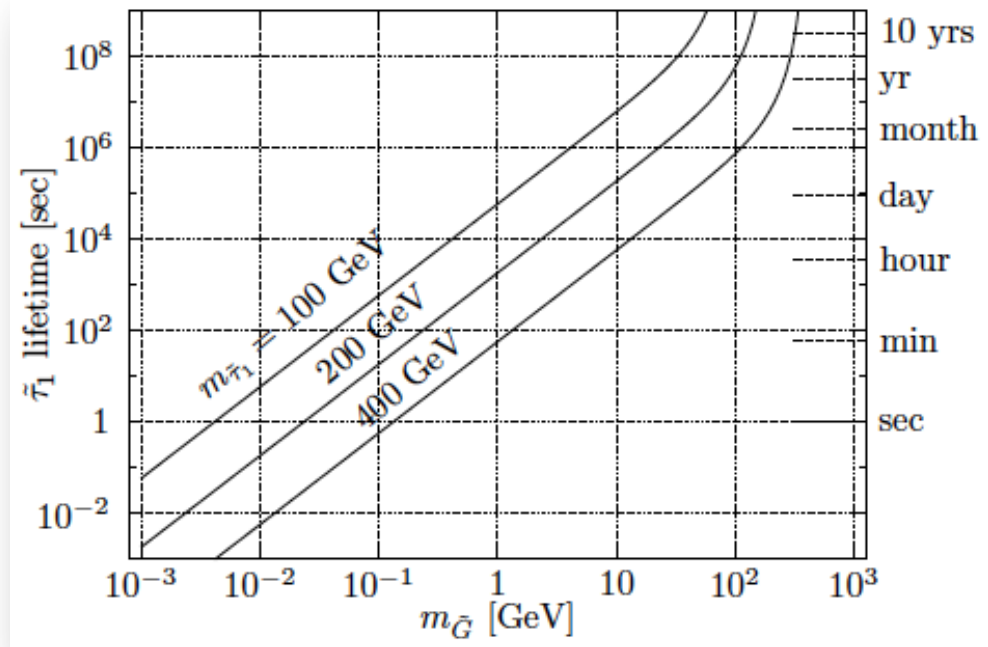
- AMSB:** $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^0$ are mass degenerate
 - long-lived chargino (\rightarrow kink track)

$$\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 + \pi^\pm$$

| SMP | LSP | Scenario | Conditions |
|------------------|---------------------------------------|-----------------|--|
| $\tilde{\tau}_1$ | $\tilde{\chi}_1^0$ | MSSM | $\tilde{\tau}_1$ mass (determined by $m_{\tilde{\tau}_{L,R}}^2, \mu, \tan \beta$, and A_τ) close to $\tilde{\chi}_1^0$ mass. |
| | \tilde{G} | GMSB | Large N , small M , and/or large $\tan \beta$. |
| | \tilde{g} | GMSB | No detailed phenomenology studies, see [20]. |
| | | SUGRA | Supergravity with a gravitino LSP, see [21]. |
| $\tilde{\tau}_1$ | | MSSM | Small $m_{\tilde{\tau}_{L,R}}$ and/or large $\tan \beta$ and/or very large A_τ . |
| | | AMSB | Small m_0 , large $\tan \beta$. |
| | | \tilde{g} MSB | Generic in minimal models. |
| \tilde{e}_{11} | \tilde{G} | GMSB | $\tilde{\tau}_1$ NLSP (see above). \tilde{e}_1 and $\tilde{\mu}_1$ co-NLSP and also SMP for small $\tan \beta$ and μ . |
| | | \tilde{g} MSB | \tilde{e}_1 and $\tilde{\mu}_1$ co-LSP and also SMP when stau mixing small. |
| | $\tilde{\chi}_1^+$ $\tilde{\chi}_1^0$ | MSSM | $m_{\tilde{\chi}_1^+} - m_{\tilde{\chi}_1^0} \lesssim m_{\pi^+}$. Very large $M_{1,2} \gtrsim 2 \text{ TeV} \gg \mu $ (Higgsino region) or non-universal gaugino masses $M_1 \gtrsim 4M_2$, with the latter condition relaxed to $M_1 \gtrsim M_2$ for $M_2 \ll \mu $. Natural in O-II models, where simultaneously also the \tilde{g} can be long-lived near $\delta_{\text{GS}} = -3$. |
| | | AMSB | $M_1 > M_2$ natural. m_0 not too small. See MSSM above. |
| | \tilde{g} $\tilde{\chi}_1^0$ | MSSM | Very large $m_{\tilde{q}}^2 \gg M_3$, e.g. split SUSY. |
| | \tilde{G} | GMSB | SUSY GUT extensions [22–24]. |
| | \tilde{g} | MSSM | Very small $M_3 \ll M_{1,2}$, O-II models near $\delta_{\text{GS}} = -3$. |
| | | GMSB | SUSY GUT extensions [22–26]. |
| \tilde{t}_1 | $\tilde{\chi}_1^0$ | MSSM | Non-universal squark and gaugino masses. Small $m_{\tilde{q}}^2$ and M_3 , small $\tan \beta$, large A_t . |
| \tilde{b}_1 | | | Small $m_{\tilde{q}}^2$ and M_3 , large $\tan \beta$ and/or large $A_b \gg A_t$. |

Long-lived sleptons – GMSB

- Gauge-mediated Supersymmetry-Breaking (GMSB)
- Stau NLSP decays via gravitational interaction to gravitino LSP
 - naturally long lifetime
 - LSP dark matter candidate
- Long-lived staus
 - may be slow-moving when produced at LHC
 - → high ionisation



Hamaguchi, Nojiri, De Roeck,
JHEP 0703 (2007) 046 [hep-ph/0612060]

$$\Gamma(\tilde{l} \rightarrow l\tilde{G}) = \frac{1}{48\pi M_*^2} \frac{m_{\tilde{l}}^5}{m_{\tilde{G}}^2} \left[1 - \frac{m_{\tilde{G}}^2}{m_{\tilde{l}}^2} \right]^4$$

average distance
travelled

$$L = \frac{1}{\kappa_\gamma} \left(\frac{100\text{GeV}}{m} \right)^5 \left(\frac{\sqrt{F/k}}{100\text{TeV}} \right)^4 \sqrt{\frac{E^2}{m^2} - 1} \times 10^{-2} \text{cm} \sqrt{F} \gtrsim 10^6 \text{ GeV}$$

Long-lived sleptons – CMSSM

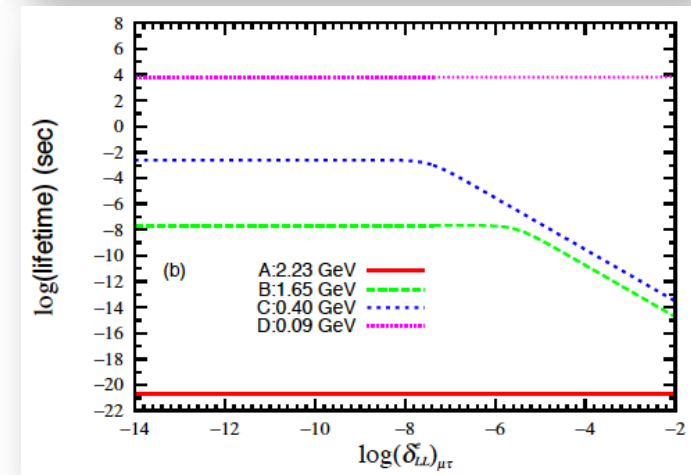
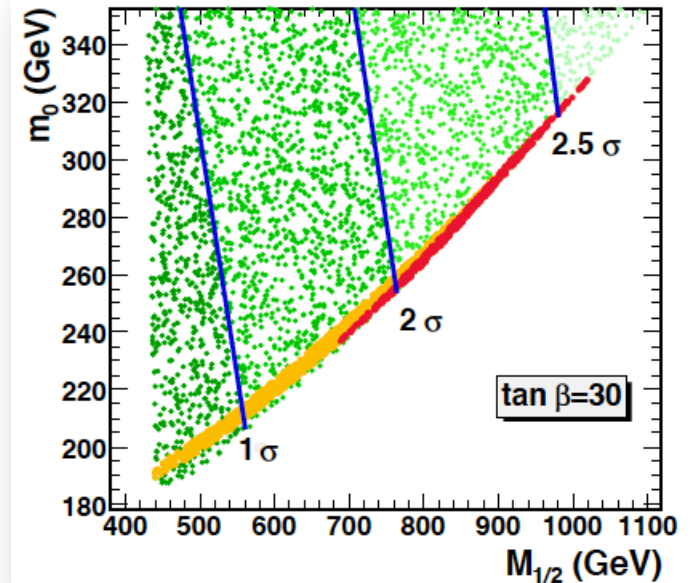
- Stau becomes long lived in MSSM when $m(\tilde{\tau}) - m(\tilde{\chi}_1^0) < m(\tau)$
- Coannihilation region in CMSSM
- Consistent with cosmological constraints
- Lepton Flavour Violating (LFV) elements in slepton mass matrix may decrease stau lifetime

$$(\delta_{RR/LL}^e)_{\alpha\beta} = \frac{\Delta M_{RR/LL}^{e2}}{M_{R/L\alpha}^e M_{R/L\beta}^e},$$

- Stau remains metastable in large regions of parameter space


$$\Gamma_{2\text{-body}} = \frac{g_2^2}{2\pi m_{\tilde{\tau}_1}} (\delta m)^2 (|g_{1\alpha 1}^L|^2 + |g_{1\alpha 1}^R|^2),$$

Kaneko, Sato, Shimomura, Vives, Yamanaka,
PRD87 (2013) 039904 [arXiv:0811.0703]



R-hadrons

- Gluinos in Split Supersymmetry

- long-lived because squarks very heavy
- possible gluino hadrons: $R = \tilde{g}q\bar{q}, \tilde{g}qqq, \tilde{g}g$
- gluino hadrons may flip charge as they pass through matter
 - e.g., $\tilde{g}u\bar{u} + uud \rightarrow \tilde{g}uud + u\bar{u}$
 -  may be missed by ATLAS and CMS

- R-parity violating SUSY

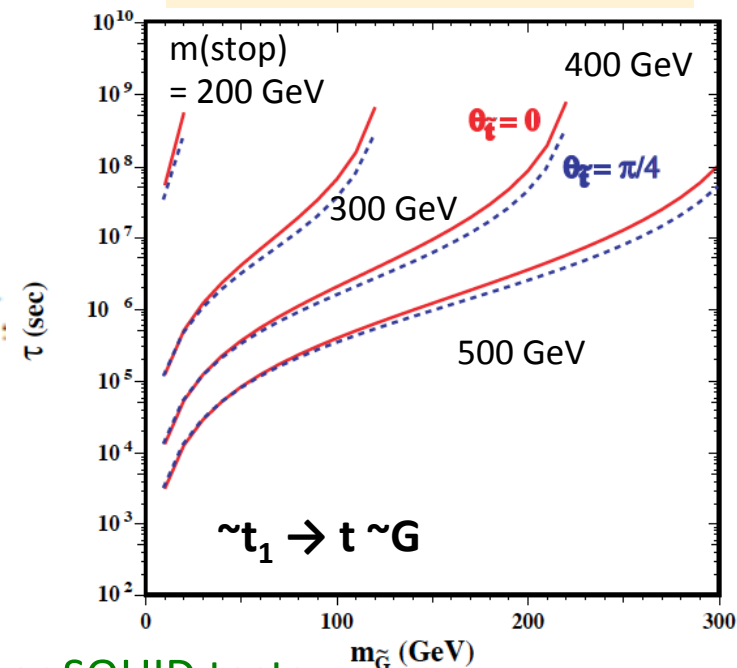
$$W_{RV} = \lambda''_{ijk} \bar{U}_i \bar{D}_j \bar{D}_k + \lambda'_{ijk} L_i Q_j \bar{D}_k + \lambda_{ijk} L_i L_j \bar{E}_k + \mu_i L_i H$$

- if λ' or $\lambda'' \neq 0$, stop NLSP case \rightarrow stop R-hadron
 - \rightarrow metastable charged particle in material
 - \rightarrow detection in MoEDAL, if sufficiently slow

- Moreover R-hadrons may be “trapped” in MMTs and decay at later times \rightarrow monitoring of MMTs after SQUID tests

$$\tau \simeq 8 \left(\frac{m_S}{10^9 \text{ GeV}} \right)^4 \left(\frac{1 \text{ TeV}}{m_{\tilde{g}}} \right)^5 \text{ s}$$

Diaz-Cruz et al, JHEP 0705 (2007) 003



Why MoEDAL when searching MSPs?

- ATLAS and CMS triggers have to
 - rely on other “objects”, e.g. E_T^{miss} , that accompany MSCP, thus limiting the reach of the search
 - final states with associated object present
 - trigger threshold set high for high luminosity
 - develop specialised triggers
 - dedicated studies needed
 - usually efficiency significantly less than 100%
- Timing: signal from (slow-moving) MSCP should arrive within the correct bunch crossing
- MoEDAL mainly constrained by its geometrical acceptance
- When looking for trapped particles
 - monitoring of detector volumes in an underground laboratory has less background than using empty bunches in LHC cavern

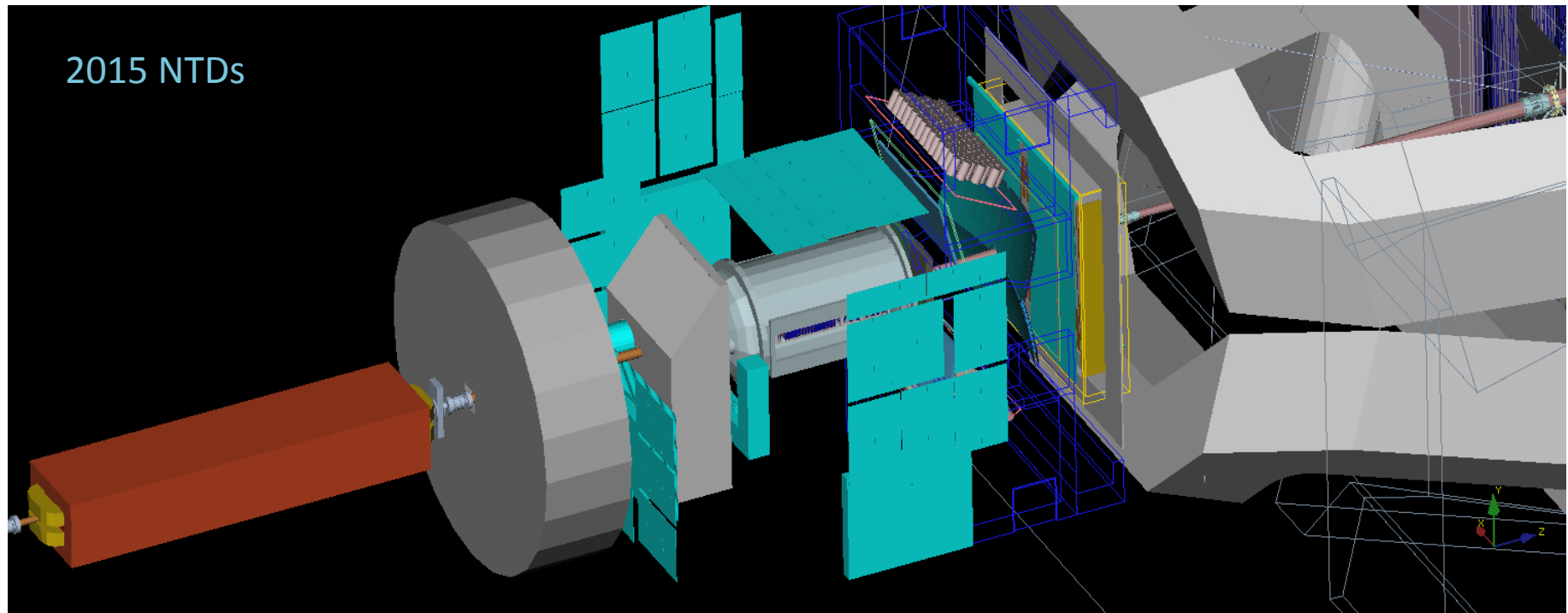
Slepton searches comparison*

| | ATLAS / CMS | MoEDAL | comments |
|-------------------|--|---|--|
| Velocity | $\beta > 0.2$ Constrained by LHC bunch pattern | $\beta < 0.2$ Constrained by NTD Z/ β threshold | Complementarity 😊 |
| Efficiency | $\epsilon \times A$ order of 20% <i>See limitations in previous slide</i> | $\sim 100\%$ (if $\beta < 0.2$) | 😐 |
| Acceptance | | $\sim 50\%$ for 2015 Scalable to higher coverage | |
| Background | May be considerable or difficult to estimate | Practically zero | For same signal yield, MoEDAL should obtain better cross-section upper limits 😊 |
| Luminosity | high | factor of 10 less | 😐 |

* Numbers are indicative

Nuclear Track Detectors coverage

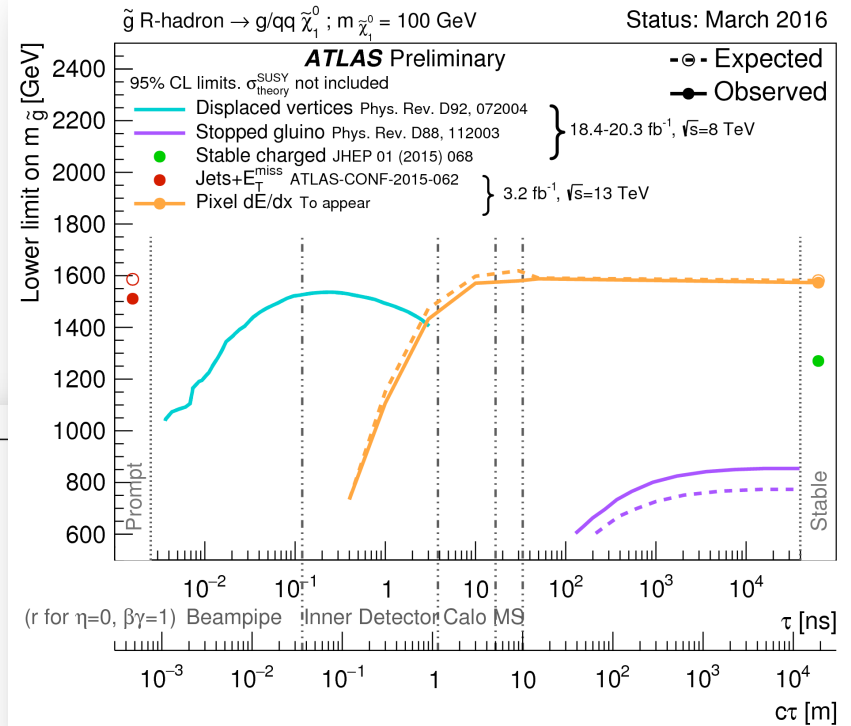
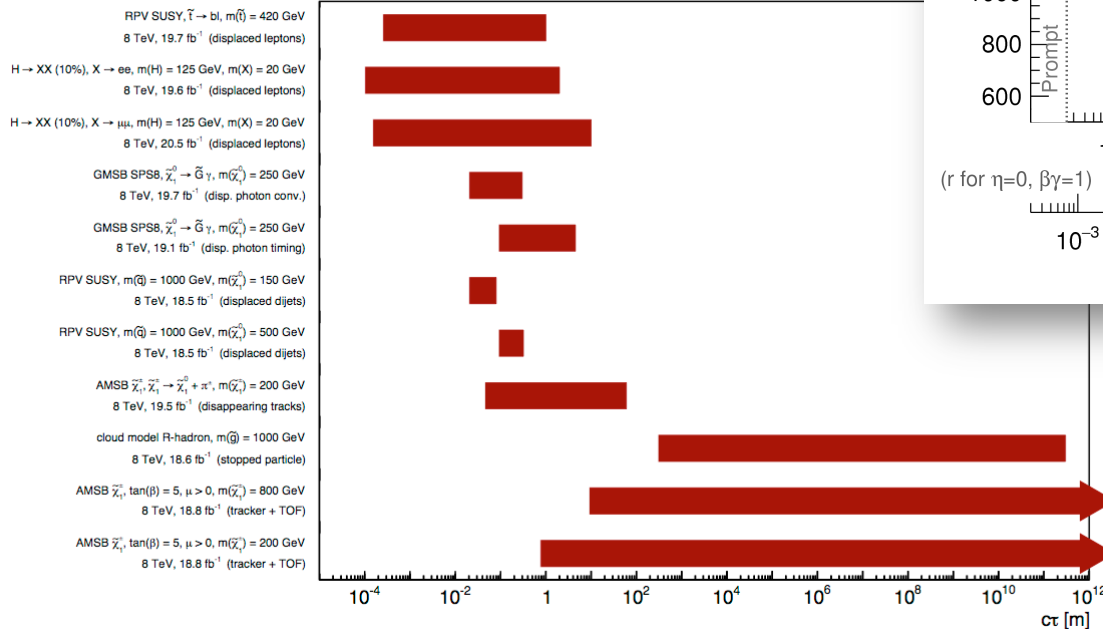
- High acceptance in central region $\eta \sim 0$
 - back-to-back pair production means probability $> \sim 70\%$ for at least one SMCP to hit NTD
- For particles over Z/β threshold, detection efficiency practically 100%



ATLAS & CMS limits on LL SUSY particles

1. Stable charged (dE/dx)
2. Stopped gluinos

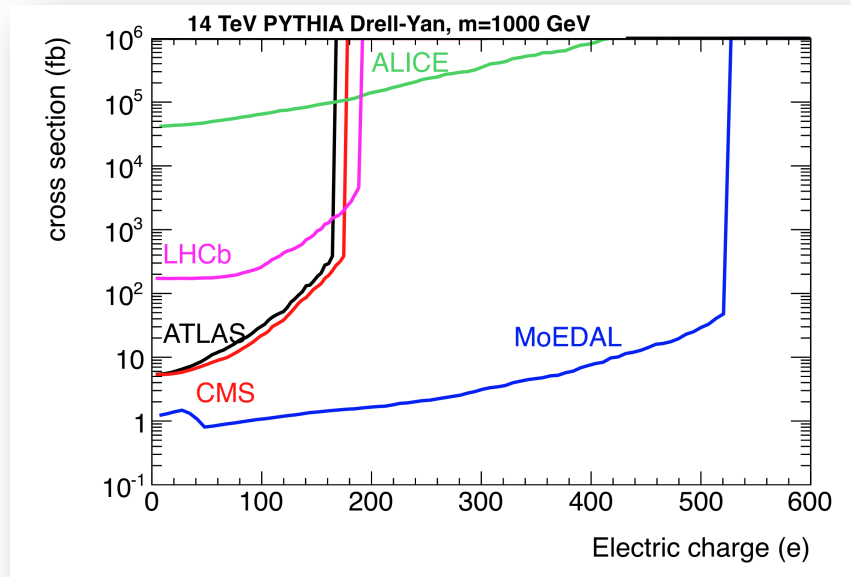
CMS long-lived particle searches, lifetime exclusions at 95% CL



MoEDAL sensitivity in SMCPs

Cross-section limits for electric charge assuming that:

- ~ one MoEDAL event is required for discovery and ~100 events in the other LHC detectors
- integrated luminosities correspond to about two years of 14 TeV run



De Roeck, Katre, Mermoud, Milstead, Sloan, EPJC72 (2012) 1985 [arXiv: 1112.2999]

- MoEDAL offers robustness against timing and well-estimated signal efficiency
- These results can be propagated to SUSY models featuring SMCPs in order to assess the discovery reach

Summary

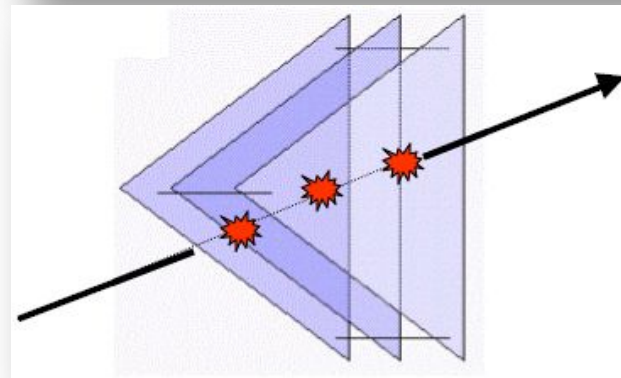
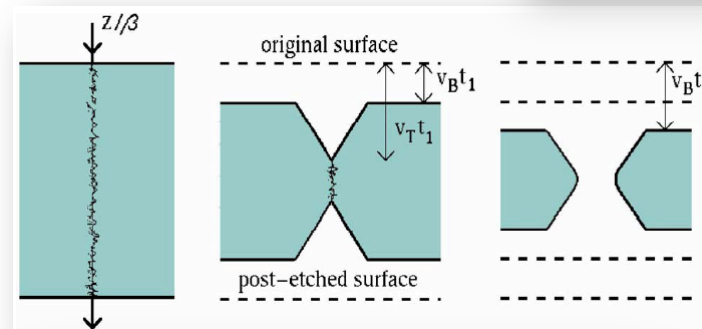
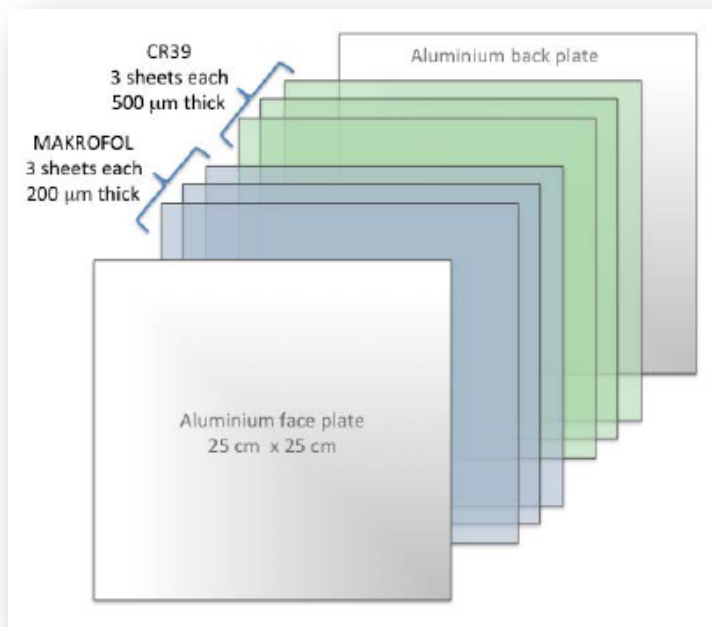
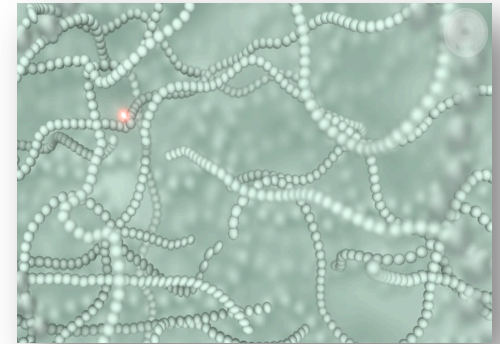
- Apart from magnetic monopoles, MoEDAL is also searching for **(meta)stable *electrically-charged* massive particles**
- Such particles arise in numerous supersymmetric scenarios (an not only)
- MoEDAL should extend significantly the discovery reach of the LHC w.r.t. such states thanks to its trigger-free concept of passive detectors
- Studies to assess discovery potential in specific SUSY models currently underway



Spares

① & ② HI particle detection in NTDs

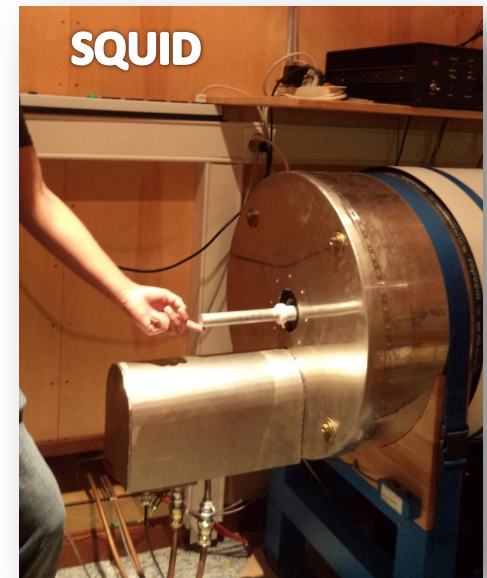
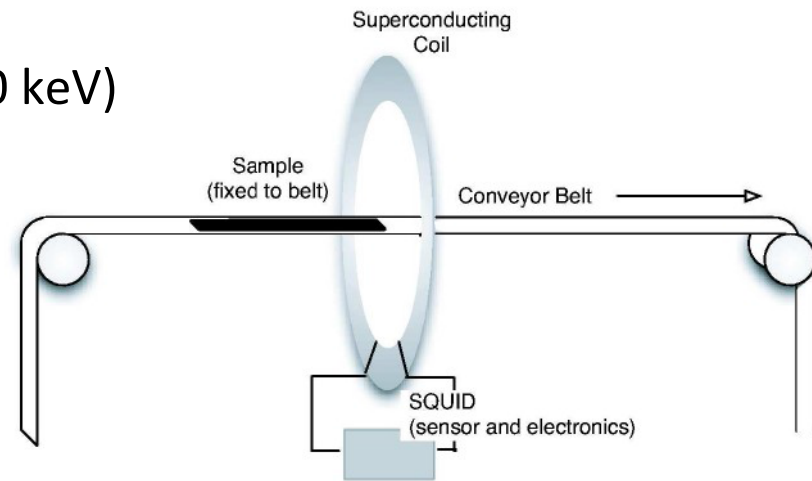
- The passage of a highly ionising particle through the plastic track-etch detector (e.g. CR39[®]) is marked by an invisible damage zone (“latent track”) along the trajectory
- The damage zone is revealed as a cone-shaped etch-pit when the plastic detector is etched in a controlled manner using a hot sodium hydroxide solution



Looking for
aligned etch pits
in multiple sheets

4 MMT: Magnetic Monopole Trapper

- Binding energies of monopoles in nuclei with finite magnetic dipole moments $\mathcal{O}(100 \text{ keV})$
- MMTs analysed with superconducting quantum interference device (SQUID)
- Material: Aluminium
 - large nuclear dipole moment
 - relatively cheap
- Disadvantage: rather low geometrical acceptance
- Advantages:
 - speed: SQUID measurements & analysis take ~ 2 weeks
 - complementarity: totally different concept from NTDs
→ different systematic uncertainties
 - magnetic charge measurement with $< 5\%$ precision
 - **Bonus:** monitoring for decay products of trapped electrically-charged particles at underground laboratory



The physics programme of the MoEDAL experiment at the LHC

B. Acharya,^{1,2} J. Alexandre,¹ J. Bernab  ,³ M. Campbell,⁴ S. Cecchini,⁵
J. Chwastowski,⁸ M. De Montigny,⁹ D. Derendarz,⁸ A. De Roeck,⁴ J. R. Ellis,^{1,4}
M. Fairbairn,¹ D. Felea,¹⁰ M. Frank,¹¹ D. Frekers,¹² C. Garcia,³ G. Giacomelli,^{5,6,*}
J. Jak  bek,¹³ A. Katre,¹⁴ D.-W. Kim,¹⁵ M. G. L. King,³ K. Kinoshita,¹⁶
D. Lacarrere,⁴ S. C. Lee,¹⁵ C. Leroy,¹⁷ A. Margiotto,⁵ N. Mauri,^{5,6}
N. E. Mavromatos,^{1,4} P. Mermod,¹⁴ V. A. Mitsou,³ R. Orava,¹⁸ L. Pasqualini,^{5,6}
L. Patrizii,⁵ G. E. P  v  las,¹⁰ J. L. Pinfold,^{9,†} M. Platkevic,¹³ V. Popa,¹⁰
M. Pozzato,⁵ S. Pospisil,¹³ A. Rajantie,¹⁹ Z. Sahnoun,^{5,7} M. Sakellariadou,¹
S. Sarkar,¹ G. Semenoff,²⁰ G. Sirri,⁵ K. Sliwa,²¹ R. Soluk,⁹
M. Spurio,^{5,6} Y. N. Srivastava,²² R. Staszewski,⁸ J. Swain,²² M. Tenti,^{5,6}
V. Togo,⁵ M. Trzebinski,⁸ J. A. Tuszy  ski,⁹ V. Vento,³ O. Vives,³
Z. Vykydal,¹³ A. Widom²³ and J. H. Yoon²³

(for the MoEDAL Collaboration)

¹Theoretical Particle Physics and Cosmology Group,
Physics Department, King's College London, UK

²International Centre for Theoretical Physics, Trieste, Italy

³IFIC, Universitat de Val  ncia - CSIC, Valencia, Spain

⁴Physics Department, CERN, Geneva, Switzerland

⁵INFN, Section of Bologna, 40127, Bologna, Italy

⁶Department of Physics & Astronomy, University of Bologna, Italy

⁷Centre on Astronomy, Astrophysics and Geophysics, Algiers, Algeria

⁸Institute of Nuclear Physics Polish Academy of Sciences, Cracow, Poland

⁹Physics Department, University of Alberta, Edmonton Alberta, Canada

¹⁰Institute of Space Science, M  gurele, Romania

¹¹Department of Physics, Concordia University, Montreal, Quebec, Canada

¹²Physics Department, University of Muenster, Muenster, Germany

¹³IEAP, Czech Technical University in Prague, Czech Republic

¹⁴Section de Physique, Universit   de Gen  ve, Switzerland

¹⁵Physics Department, Gangneung-Wonju National University, Gangneung, Korea

¹⁶Physics Department, University of Cincinnati, Cincinnati OH, USA

¹⁷Physics Department, University de Montr  al, Montr  al, Qu  bec, Canada

¹⁸Physics Department, University of Helsinki, Helsinki, Finland

¹⁹Physics Department, Imperial College London, UK

²⁰Department of Physics, University of British Columbia, Vancouver BC, Canada

²¹Department of Physics and Astronomy, Tufts University, Medford MA, USA

²²Department of Physics, Northeastern University, Boston, USA

²³Physics Department, Konkuk University, Seoul, Korea

[†]jpinfold@ualberta.ca

Many more interesting theoretical scenarios relevant and accessible to MoEDAL not presented here:

- doubly-charged Higgs
- black-hole remnants
- quirks
- Q-balls
- CHAMPS
-

Complete and detailed review
on MoEDAL impact on
searches for exotic models

MoEDAL physics program:

IJMP A29 (2014) 1430050

[arXiv:1405.7662](https://arxiv.org/abs/1405.7662)

MoEDAL web page:

<http://moedal.web.cern.ch/>

Analysis procedure

✦ Track diameter:

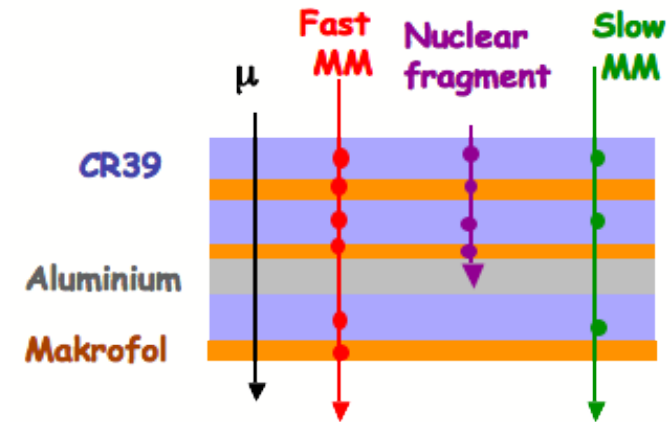
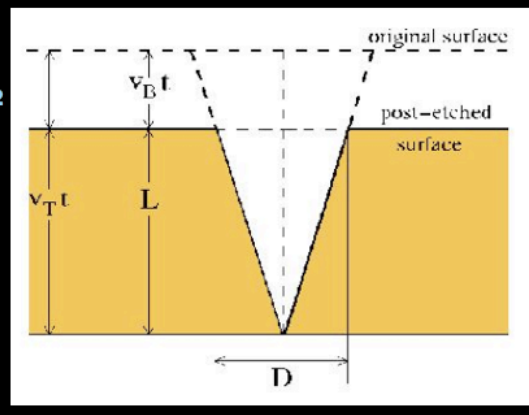
$$\star D = 2v_B[(v_T - v_B)/(v_T + v_B)]^{-1/2}$$

✦ Track depth:

$$\star L = (v_T - v_B) t$$

✦ Reduced etch rate:

$$\star p = v_T/v_B$$



- Electrically-charged particle: $dE/dx \sim \beta^{-2} \rightarrow$ slows down appreciably within NTD \rightarrow opening angle of etch-pit cone becomes **smaller**
- Magnetic monopole: $dE/dx \sim \ln \beta$
 - slow MM: slows down within an NTD stack \rightarrow its ionisation falls \rightarrow opening angle of the etch pits would become **larger**
 - relativistic MM: dE/dx essentially constant \rightarrow trail of equal diameter etch-pit pairs
- The reduced etch rate is simply related to the *restricted energy loss*
 $REL = (dE/dx)_{10\text{nm from track}}$