Search for BSM physics in fermionic final states Alexander Schmidt on behalf of the ATLAS and CMS collaborations



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1106-

LHC 27 km

Vauxhall Astra BSM Diecast Model Car by Vanguards



Universität Hamburg

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ATLAS

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ALICE

typical signatures



search for exotic resonances



non-SUSY BSM

- similar motivations as SUSY
- examples:
 - extra dimensions: warped extra dimension models where fermions propagate in the bulk
 - composite Higgs: Heavy Vector
 - Triplet model, with new W'[±], Z' states
 - contributions to S and T parameters must be small



collider reach

parton luminosity scaling:



collider reach

parton luminosity scaling:



• we do better at high masses

topology of a tt resonance

light resonances (M ~ 500 GeV):

"classical" event topology



- decay products p_T~1TeV
- large ¥ factor (>5 -10)
- jets overlap and merge
- special reconstruction techniques needed !







modeling of boosted variables at 13 TeV:

semi-leptonic top selection



 after initial difficulties at the start of Run II, things look very good now

EXPERIMENT IT RESONANCES: SEMI-leptonic



SATLAS TT RESONANCES: Semi-leptonic

[ATLAS-CONF-2016-014]



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- top tagging
- subjet b-tagging



background from data:

- invert substructure selection criterion on one jet (QCD region)
- measure mistagging probability of other jet
- parameterised in bins of b-tag and p_{T}
- apply mistag rate in single t-tagged sample





tt resonances: hadronic

[CMS-PAS-B2G-15-003]







vector-like quarks: pair production

X^{5/3} in same-sign di-lepton final state:

- first VLQ result at 13 TeV
- clean same-sign lepton selection
- backgrounds from opposite-sign events with charge mis-measurements





VALUAS vector-like quarks: pair production



Vector-like quarks: pair production

[ATLAS-CONF-2016-032]



YATLAS Vector-like quarks: pair production

T^{2/3} in lepton+jets final state:

- optimised for T→tH + X
 (with H→bb)
- signal characterised by high (b)-jet multiplicity
- 11 categories based on
 - jet multiplicity (>6 for signal)
 - multiplicity of mass-tagged large-R jets(m_>100GeV)
 - b-tag multiplicity
 - mass of two close b-jets
 - high mass m_{bb} >100 GeV
 - low mass mbb<100 GeV

also sensitive to 4-top production!



 $BR(T \rightarrow Ht)$

[ATLAS-CONF-2016-013]



m_{KK} [GeV]

vector-like quarks: pair production



vector-like quarks: single production

[CMS-PAS-B2G-15-008]



vector-like quarks: single production

[CMS-PAS-B2G-16-005]



experimental challenges:

- lepton reconstruction at high momenta
- efficiency and fake rate calibration (extrapolation from low to high momenta)
- background estimation (extrapolation)
- 2.6 fb⁻¹ (13 TeV) Events 10⁵ Events / GeV Data 10⁶ ATLAS Preliminary CMS 10⁴ Data Z/γ^* $\sqrt{s} = 13 \text{ TeV}, 3.2 \text{ fb}^{-1}$ Preliminary γ/Z→e⁺e⁻ **Dilepton Search Selection** 10⁵ Top Quarks 10³ $t\bar{t}$, tW, WW, WZ, ZZ, $\tau\tau$ Diboson 10² Jets 10⁴ Multi-Jet & W+Jets Narrow Z' ($M_{2} = 2 \text{ TeV}$) Z'_{γ} (3 TeV) 10 10^{3} $\Lambda_{11} = 20 \text{ TeV}$ 10² 10^{-1} 10 10⁻² 10⁻³ 10^{-4} 10- 10^{-5} 10⁻² 10^{-6} 100 200 300 1000 70 2000 Data / Bkg m(ee) [GeV] highest mass event: mee=2.9 TeV 0.6 2000 3000 200 300 400 100 1000 (0.3 events expected above 2 TeV) **Dielectron Invariant Mass [GeV]** (0.08 events expected above 2.5 TeV)



NEW



search for di-tau resonances:

- first di-tau result at 13 TeV
- four final states considered: $\tau_h \tau_{h}$, $\tau_\mu \tau_{h}$, $\tau_e \tau_{h}$, $\tau_\mu \tau_e$



high-mass dijet resonances



high-mass dijet resonances





MATLAS high-mass di-b-jet resonances



. .



- data scouting (trigger-object-level) search in mass region 450 - 900 GeV
- avoids high trigger pre-scales for full analysis
- dedicated trigger-level jet calibration





results:

Events

10⁶

10⁵

Significance

400

p⊤>185 (85) GeV

500

600

- empirical function to fit the background spectrum
- analyse all possible mass intervals for excess
- most discrepant interval 574-685 GeV (0.8 σ)
- excludes gaussian excess with cross sections 3 pb (450 GeV) to 0.7 pb (900 GeV)





[ATLAS-CONF-2016-030]





SATLAS IOW-MASS di-b-jet resonances

[ATLAS-CONF-2016-031]

[ATLAS-CONF-2016-029]

NEW

alternative approach:

 exploit initial state photon for triggering dijets (isolated photon p_T>120 GeV)





ATLAS Exotics Searches* - 95% CL Exclusion

Status: March 2016

ATLAS Preliminary

 $\int \mathcal{L} dt = (3.2 - 20.3) \text{ fb}^{-1}$ $\sqrt{s} = 8, 13 \text{ TeV}$

	 	-

	Model	<i>ℓ</i> ,γ	Jets†	$\mathbf{E}_{\mathrm{T}}^{\mathrm{miss}}$	∫£ dt[fb	¹] Limit	Reference
Extra dimensions	ADD $G_{KK} + g/q$ ADD non-resonant $\ell\ell$ ADD QBH $\rightarrow \ell q$ ADD QBH ADD QBH ADD BH high $\sum p_T$ ADD BH multijet RS1 $G_{KK} \rightarrow \ell\ell$ RS1 $G_{KK} \rightarrow \gamma\gamma$ Bulk RS $G_{KK} \rightarrow WW \rightarrow qq\ell\nu$ Bulk RS $G_{KK} \rightarrow HH \rightarrow bbbb$ Bulk RS $g_{KK} \rightarrow tt$ 2UED / RPP	$ \begin{array}{c} - \\ 2 e, \mu \\ 1 e, \mu \\ - \\ \geq 1 e, \mu \\ 2 e, \mu \\ 2 \gamma \\ 1 e, \mu \\ - \\ 1 e, \mu \\ 1 e, \mu \\ \end{array} $	$\geq 1j$ $-$ $1j$ $2j$ $\geq 2j$ $\geq 3j$ $-$ $1J$ $4b$ $\geq 1b, \geq 1J$ $\geq 2b, \geq 4$	Yes Yes j Yes	3.2 20.3 20.3 3.6 3.2 3.6 20.3 20.3 3.2 3.2 20.3 3.2 20.3 3.2	M _D 6.86 TeV $n = 2$ M _S 4.7 TeV $n = 3$ HLZ M _{th} 5.2 TeV $n = 6$ M _{th} 8.3 TeV $n = 6$ M _{th} 8.2 TeV $n = 6$ M _{th} 9.55 TeV $n = 6$, $M_D = 3$ TeV, rot BH M _{th} 9.55 TeV $n = 6$, $M_D = 3$ TeV, rot BH M _{th} 9.55 TeV $n = 6$, $M_D = 3$ TeV, rot BH M _{th} 9.55 TeV $n = 6$, $M_D = 3$ TeV, rot BH M _{th} 9.55 TeV $n = 0.1$ K/M p _l = 0.1 $k/\overline{M}_{Pl} = 0.1$ K/M mass 1.06 TeV $k/\overline{M}_{Pl} = 1.0$ G _{KK} mass 2.2 TeV BR = 0.925 KK mass 1.46 TeV Tier (1,1), BR(A^{(1,1)} \to tt) = 1	Preliminary 1407.2410 1311.2006 1512.01530 ATLAS-CONF-2016-006 1512.02586 1405.4123 1504.05511 ATLAS-CONF-2015-075 ATLAS-CONF-2016-017 1505.07018 ATLAS-CONF-2016-013
Gauge bosons	$\begin{array}{l} \operatorname{SSM} Z' \to \ell\ell \\ \operatorname{SSM} Z' \to \tau\tau \\ \operatorname{Leptophobic} Z' \to bb \\ \operatorname{SSM} W' \to \ell\nu \\ \operatorname{HVT} W' \to WZ \to qqv\nu \mbox{ model A} \\ \operatorname{HVT} W' \to WZ \to qqqq \mbox{ model A} \\ \operatorname{HVT} W' \to WH \to \ell\nu bb \mbox{ model B} \\ \operatorname{HVT} Z' \to ZH \to \nu\nu bb \mbox{ model B} \\ \operatorname{LRSM} W'_R \to tb \\ \operatorname{LRSM} W'_R \to tb \end{array}$	2 e,μ 2 τ - 1 e,μ 0 e,μ - 1 e,μ 0 e,μ 1 e,μ 0 e,μ	$\begin{array}{c} - \\ 2 \ b \\ - \\ 1 \ J \\ 2 \ J \\ 1 - 2 \ b, 1 - 0 \\ 1 - 2 \ b, 1 - 0 \\ 2 \ b, 0 - 1 \ j \\ \geq 1 \ b, 1 \ J \end{array}$	– Yes Yes j Yes j Yes J –	3.2 19.5 3.2 3.2 3.2 3.2 3.2 3.2 3.2 20.3 20.3	Z' mass 3.4 TeV Z' mass 2.02 TeV Z' mass 2.02 TeV Z' mass 1.5 TeV W' mass 4.07 TeV W' mass 1.6 TeV $g_V = 1$ W' mass 1.38-1.6 TeV $g_V = 3$ W' mass 1.76 TeV $g_V = 3$ W' mass 1.92 TeV $g_V = 3$	ATLAS-CONF-2015-070 1502.07177 Preliminary ATLAS-CONF-2015-063 ATLAS-CONF-2015-068 ATLAS-CONF-2015-073 ATLAS-CONF-2015-074 ATLAS-CONF-2015-074 1410.4103 1408.0886
CI	CI qqqq CI qqℓℓ CI uutt	_ 2 e, μ 2 e, μ (SS)	2 j _ ≥ 1 b, 1-4	– – j Yes	3.6 3.2 20.3	Λ 17.5 TeV $\eta_{LL} = -1$ Λ 23.1 TeV $\eta_{LL} = -1$ Λ 4.3 TeV $ C_{LL} = 1$	1512.01530 ATLAS-CONF-2015-070 1504.04605
DM	Axial-vector mediator (Dirac DM) Axial-vector mediator (Dirac DM) $ZZ_{\chi\chi}$ EFT (Dirac DM)	0 e, μ 0 e, μ, 1 γ 0 e, μ	≥1j 1j 1J,≤1j	Yes Yes Yes	3.2 3.2 3.2	m _A 1.0 TeV g_q =0.25, g_χ =1.0, $m(\chi)$ < 140 GeV m _A 650 GeV g_q =0.25, g_χ =1.0, $m(\chi)$ < 10 GeV M _* 550 GeV $m(\chi)$ < 150 GeV	Preliminary Preliminary ATLAS-CONF-2015-080
ΓG	Scalar LQ 1 st gen Scalar LQ 2 nd gen Scalar LQ 3 rd gen	2 e 2 μ 1 e, μ	≥ 2 j ≥ 2 j ≥1 b, ≥3	– – j Yes	3.2 3.2 20.3	LQ mass 1.07 TeV $\beta = 1$ LQ mass 1.03 TeV $\beta = 1$ LQ mass 640 GeV $\beta = 0$	Preliminary Preliminary 1508.04735
Heavy quarks	$\begin{array}{l} VLQ\ TT \to Ht + X\\ VLQ\ YY \to Wb + X\\ VLQ\ BB \to Hb + X\\ VLQ\ BB \to Zb + X\\ VLQ\ QQ \to WqWq\\ T_{5/3} \to Wt \end{array}$	$\begin{array}{c} 1 \ e, \mu \\ 1 \ e, \mu \\ 1 \ e, \mu \\ 2/{\geq} 3 \ e, \mu \\ 1 \ e, \mu \\ 1 \ e, \mu \\ 1 \ e, \mu \end{array}$	$\begin{array}{l} \geq 2 \ \text{b}, \geq 3 \\ \geq 1 \ \text{b}, \geq 3 \\ \geq 2 \ \text{b}, \geq 3 \\ \geq 2 \ \text{b}, \geq 3 \\ \geq 2/{\geq}1 \ \text{b} \\ \geq 4 \ \text{j} \\ \geq 1 \ \text{b}, \geq 5 \end{array}$	j Yes j Yes j Yes - Yes j Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3	T mass855 GeVY mass770 GeVB mass735 GeVB mass735 GeVB mass755 GeVQ mass690 GeVT _{5/3} mass840 GeV	1505.04306 1505.04306 1505.04306 1409.5500 1509.04261 1503.05425
Excited fermions	Excited quark $q^* \rightarrow q\gamma$ Excited quark $q^* \rightarrow qg$ Excited quark $b^* \rightarrow bg$ Excited quark $b^* \rightarrow Wt$ Excited lepton ℓ^* Excited lepton γ^*	1 γ - - 1 or 2 e, μ 3 e, μ 3 e, μ, τ	1 j 2 j 1 b, 1 j 1 b, 2-0 j –	- - Yes -	3.2 3.6 3.2 20.3 20.3 20.3	q* mass 4.4 TeV only u^* and d^* , $\Lambda = m(q^*)$ q* mass 5.2 TeV only u^* and d^* , $\Lambda = m(q^*)$ b* mass 2.1 TeV $f_g = f_L = f_R = 1$ d* mass 3.0 TeV $\Lambda = 3.0$ TeV v* mass 1.6 TeV $\Lambda = 1.6$ TeV	1512.05910 1512.01530 Preliminary 1510.02664 1411.2921 1411.2921
Other	LSTC $a_T \rightarrow W\gamma$ LRSM Majorana ν Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$ Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$ Monotop (non-res prod) Multi-charged particles Magnetic monopoles	1 e, μ, 1 γ 2 e, μ 2 e, μ (SS) 3 e, μ, τ 1 e, μ - - -	_ 2 j - 1 b - -	Yes - - Yes - - 3 TeV	20.3 20.3 20.3 20.3 20.3 20.3 20.3 7.0	a_T mass960 GeV N^0 mass2.0 TeV $H^{\pm\pm}$ mass551 GeV $H^{\pm\pm}$ mass551 GeV $H^{\pm\pm}$ mass400 GeVspin-1 invisible particle mass657 GeVmulti-charged particle mass785 GeVmonopole mass1.34 TeV $I = 1$ 1	1407.8150 1506.06020 1412.0237 1411.2921 1410.5404 1504.04188 1509.08059
						Mass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. Lower bounds are specified only when explicitly not excluded. *†Small-radius (large-radius) jets are denoted by the letter j (J).*











B2G: new physics searches with heavy SM particles

Resonances to dibosons



Excited quarks



+model-independent





+model-independent

backup



		excluded mass regions [TeV]		
signal	μ + jets observed (expected)	<i>e</i> + jets observed (expected)	combination observed (expected)	8 TeV combination
Z' (1% width)	0.5 - 1.8 (0.6 - 1.9)	1.0 - 1.1, 1.3 - 2.2 (0.9 - 1.7)	0.6 - 2.3 (0.6 - 2.1)	2.4 (2.4)
Z' (10% width)	0.5 - 3.2 (0.5 - 3.3)	0.5 - 3.2 (0.5 - 3.2)	0.5 - 3.4 (0.5 - 3.5)	2.9 (2.8)
Z' (30% width)	0.5 - 3.9(0.5 - 4.0)	0.5 - 3.8 (0.5 - 3.8)	$0.5 - 4.0 \ (0.5 - 4.0)$	
KK gluon	0.5 - 2.7 (0.5 - 2.8)	0.6 - 2.7 (0.6 - 2.5)	0.5 - 2.9 (0.5 - 2.9)	2.8 (2.7)

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dilepton resonances

• muons



Dimuon Invariant Mass [GeV]

e-mu resonances

[CMS-PAS-EXO-16-001]



tt resonances: semi-leptonic



what are VLQ:

- •they are quarks: coloured, charged, spin 1/2 particles
- •no difference between chiralities: they couple to left- and righthanded charged currents (in the same way)
- VL quarks can have mass terms without violating gauge invariance!
- not constrained through Higgs discovery
- new motivation through di-photon excess
- decay modes:
 T→ tH, tZ, bW
 B→bH, bZ, tW



vector-like quarks

[JHEP 08 (2015) 105]



vector-like quarks: single production



higher cross sections at high masses

allows to set limits on model parameters

vector-like quarks: single production $T \rightarrow tH_{[CMS-PAS-B2G-15-008]}$

background estimation

- entirely data-driven (shapes and normalisation)
- background region: no forward jet, subjet b-tag veto



vector-like quarks: pair production

[ATLAS-CONF-2016-13]



- scalar sum of lepton, jets, missing energy
- used to discriminate signal and background





vector-like quarks: pair production

[ATLAS-CONF-2016-13]



X^{5/3} in same-sign di-lepton final state:

After requiring two tight, same-sign leptons with p_T greater than 30 GeV we impose the following requirements:

- Quarkonia Veto: *M*_{ll} > 20 GeV
- Associated Z-boson Veto: veto any event where either of the leptons in the samesign pair reconstructs to within 15 GeV of the mass of the Z-boson with any other lepton in the event not in the same-sign pair.
- Primary Z-boson Veto: invariant dilepton mass (M_{ll}) > 106.1 or < 76.1 GeV for dielectron channel only. If the muon charge is mismeasured, its momentum will also be mismeasured so a selected muon pair from a Zboson will not fall within this invariant mass range.
- Leading lepton p_T > 40 GeV
- Number of Constituents >= 5
- $H_{\rm T}^{\rm lep} > 900 \,{\rm GeV}$

X^{5/3} in lepton+jet final state:

We begin by defining event preselection criteria consistent with the signal topology. The following selection is applied: lepton $p_T > 50 \text{ GeV}$, $E_T^{\text{miss}} > 100 \text{ GeV}$, $N_{jets} \ge 3$ with jet $p_T > 30 \text{ GeV}$ and $|\eta| < 2.4$, p_T (leading jet) > 200 GeV, and p_T (sub leading jet) > 90 GeV. In addition at least one of the jets is required to be tagged as originating from a b quark [28]. In Fig. 4, we show the differences in signal and background as characterized by a few kinematic variables, the number of jets (N_{jets}), missing E_T (E_T^{miss}). Decay products of heavy particles such as $X_{5/3}$ can have large Lorentz boost, and their subsequent decay products can merge into a single jet. The "n-subjettiness" [29] algorithm measures the likelihood of a jet having *n* subjets (1, 2, 3, etc). Mass grooming techniques are used to remove soft jet constituents so that the mass of the hard constituents can be measured more clearly. The "pruning" [30] algorithm is currently used to identify boosted hadronic W boson decays. We require a W tagged jet to have $p_T > 200$ GeV, $|\eta| < 2.4$, pruned mass between 65 and 105 GeV, and the ratio of n-subjettiness variables $\tau_2/\tau_1 < 0.55$. Figure 4 shows the number of b tagged and W tagged jets (bottom row).

T^{2/3} in lepton+jets final state:

We select events with one electron or muon, usually from the decay of a W boson in the bW channel or from the t \rightarrow Wb decay in the tZ or tH channels. Electrons and muons must have $p_T > 40$ GeV and pass the tight identification and isolation requirements described previously. We discard events having extra loose electrons or muons with $p_T > 10$ GeV.

We select events with three or more jets, where the highest corrected jet p_T is greater than 150 GeV, and the second highest p_T is greater than 75 GeV. E_T^{miss} must be greater than 60 GeV to reduce multijet background events. The final selection criterion is a veto on boosted Higgs tagged jets, which creates a signal region orthogonal to that of an ongoing dedicated $T \rightarrow tH$ analysis. Figure 2 shows distributions of lepton p_T , jet p_T , E_T^{miss} along with the total number of jets, number of b tagged jets, and number of W tagged large-radius jets. These distributions include the jet p_T reweighting procedure described below.

In addition, a Universal Extra Dimensions (UED) model is considered involving new heavy particles. The UED model considered has two extra dimensions that are compactified using the geometry of the real projective plane (2UED/RPP) [41], leading to a discretisation of the momenta along their directions. A tier of Kaluza–Klein towers is labelled by two integers, k and ℓ , referred to as "tier (k, ℓ) ". Within a given tier, the squared masses of the particles are given at leading order by $m^2 = k^2/R_4^2 + \ell^2/R_5^2$, where πR_4 and πR_5 are the sizes of the two extra dimensions. The model is parameterised by R_4 and R_5 or, alternatively, by $m_{\rm KK} = 1/R_4$ and $\xi = R_4/R_5$. Four-top-quark production can arise from tier (1,1), where particles from this tier have to be pair produced because of symmetries of the model. Then they chaindecay to the lightest particle of this tier, the heavy photon $A^{(1,1)}$, by emitting SM particles (Fig. 3(c)). The branching ratios of $A^{(1,1)}$ into SM particles are not predicted by the model, although the decay into $t\bar{t}$ is expected to be dominant [42]. Four-top-quark events can also arise from tiers (2,0) and (0,2) via a similar mechanism. In this case the expected cross section for four-top-quark production is reduced compared to that from tier (1,1) since each state in tiers (2,0) and (0,2) can decay directly into a pair of SM particles or into a pair of states in tiers (1,0) or (0,1) via bulk interactions, resulting in smaller branching ratios for the decay into $t\bar{t}$ [42]. In the following, when considering four-top-quark production from a given tier, it is assumed that the A photon in that tier decays exclusively into $t\bar{t}$ while A photons from other tiers cannot decay into $t\bar{t}$. Within this model, observations of dark-matter relic abundance prefer values of $m_{\rm KK}$ between 600 GeV and 1200 GeV [43].

77 2.3 Trigger-Level Analysis in ATLAS

Partially-built events are collected by means of an additional TLA data stream. In this stream, all events 78 containing at least one L1 jet ROI with $E_T > 75$ GeV at the electromagnetic scale ² are recorded. For 79 each trigger jet recorded at the HLT with $p_T > 4$ GeV, the stream records the jet four-momentum and 80 a set of variables characterizing the detector information within the jet (such as variables needed for 81 jet identification [25], and limited information about the structure of the jet). The data format does not 82 include any readout of individual calorimeter elements, nor does it record information from the tracking 83 or muon detectors. The size of partially-built TLA events in 2015 data is less than 5% of the size of full 84 events, allowing for higher event rates to be recorded. TLA events are recorded with a rate of 2 kHz, in 85 addition to the fully-recorded higher $p_{\rm T}$ events. 86

The search uses a sample of events which contain at least one photon with $p_T > 130$ GeV as well as at least two jets. Jets are required to have $p_T > 25$ GeV and $|\eta| < 2.8$. In addition, half the difference in rapidity y of the leading two jets,

$$y^* = \frac{y_{j_1} - y_{j_2}}{2}$$

is required to satisfy $|y^*| < 0.8$ to reduce the backgrounds from processes that yield an ISR photon and a non-resonant jet pair from QCD processes, which tend to have higher-rapidity separation. To suppress such jets surviving the photon isolation criterion, and to further suppress background from fragmentation photons where the photon is inside or near a jet, the selected photon is required to be separated from the closest jet by $\Delta R > 0.85$. Figure 2 shows the m_{ij} distribution of events satisfying these criteria.



[CMS-PAS-EXO-16-008]









jet grooming



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jet shapes

energy patterns within a jet:

- characterise multi-prong properties
- many definitions and varieties

n-subjettiness:

$$\tau_{N} = \frac{\sum_{i=1}^{n_{\text{constituents}}} p_{\text{T},i} \min\{\Delta R_{1,i}, \Delta R_{2,i}, ..., \Delta R_{N,i}\}}{\sum_{i=1}^{n_{\text{constituents}}} p_{\text{T},i}R}$$

- how consistent is jet with having N subjets
- ratios discriminate hypotheses: τ_2/τ_1

many other on the market:

- energy correlation functions
- Q-jet volatility



13 TeV

performances:

 $^{\mathrm{B}}_{\mathrm{B}}$

 several variables scanned individually

background rejection at 30% efficiency

