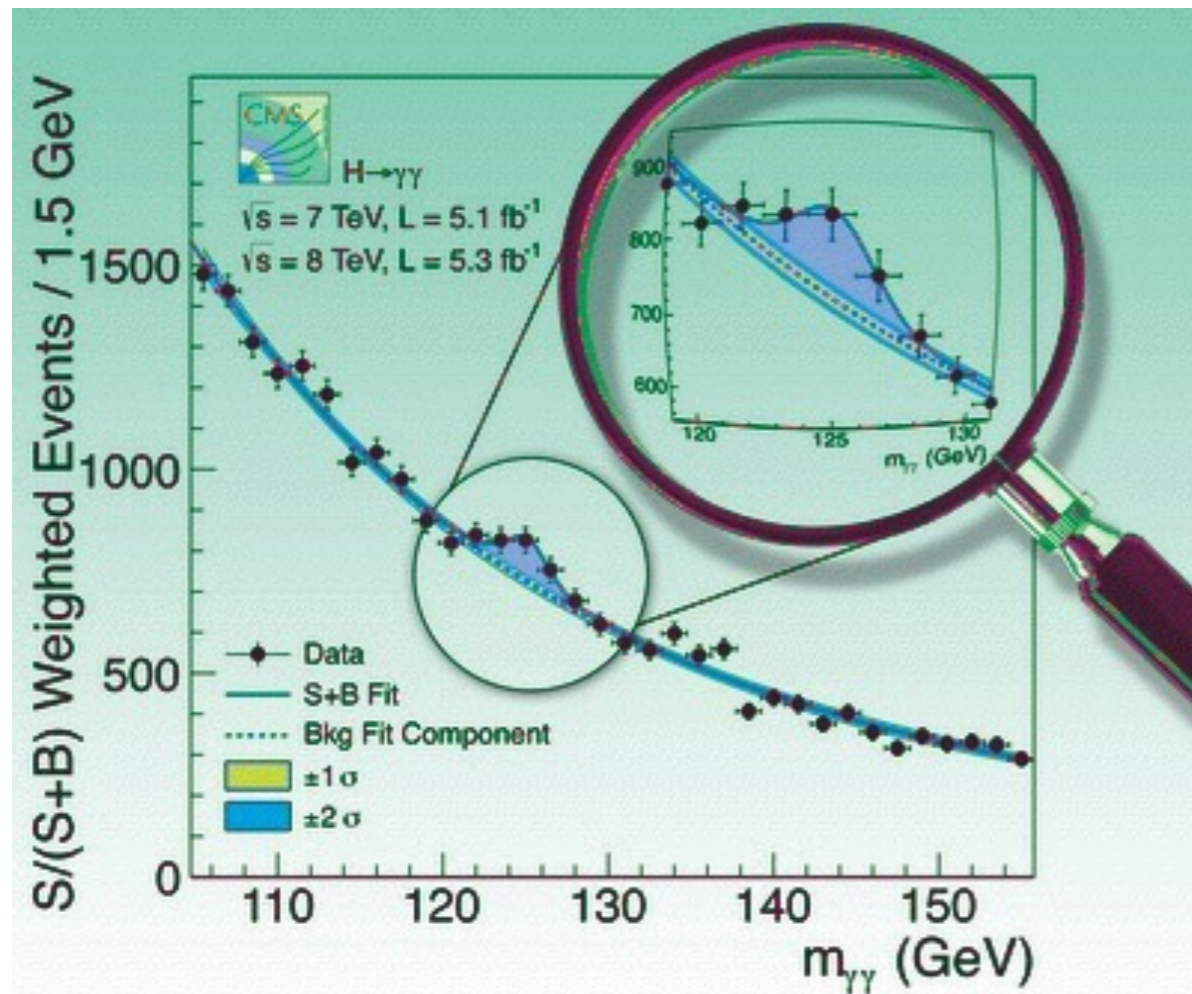


Strongly coupled gauge theories: What can lattice calculations teach us?

**Anna Hasenfratz
University of Colorado Boulder
CERN, Nov 25, 2015**

Higgs era of particle physics

The 2012 discovery of the Higgs boson “completed” the Standard Model



Yet many questions remain open

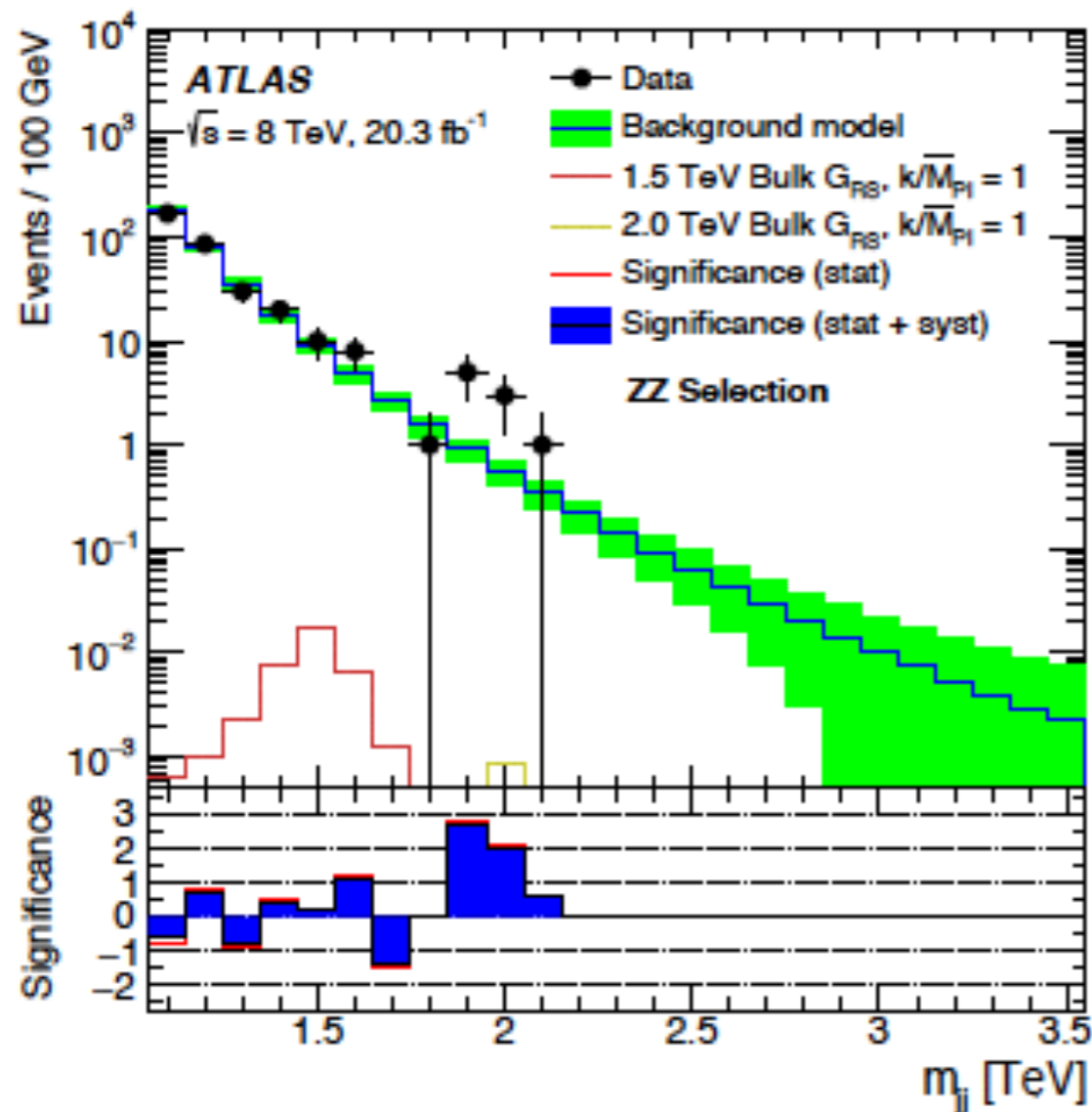
- The Standard Model is not UV complete
- What is the Higgs boson
- What is the nature of electroweak symmetry breaking?

No answers — yet

Higgs era of particle physics

Tantalizing possibility from ATLAS : 3.5σ excess
at 2TeV suggesting a vector resonance

(1506.00962)



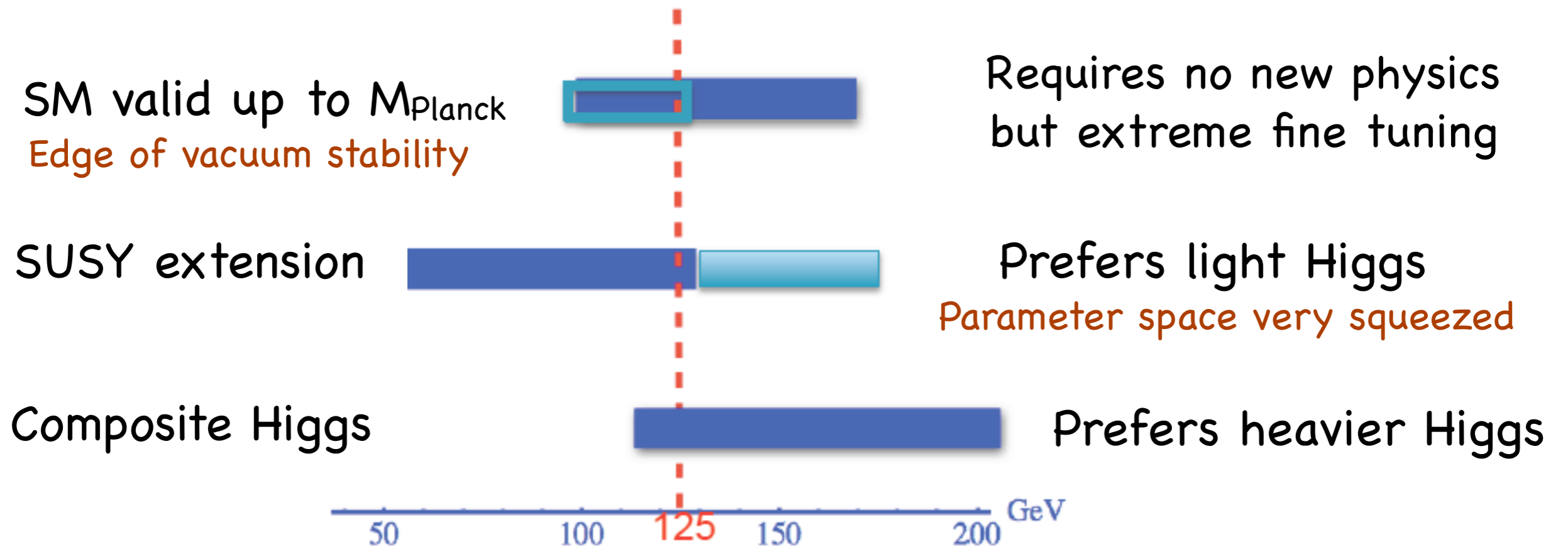
- Does the ATLAS result point to new physics?(if it is indeed there)
- Which BSM models can describe it?
- What other predictions do those models have?

Possible BSM models

a few popular options ...

blue bands represent typical ranges for Higgs mass

(from M. Carena)



New symmetries are needed to stabilize a scalar against (quadratic) UV divergences (below the Planck mass).

Composite Higgs - strong dynamics

Assume a new interaction:

N_f fermions, $SU(N_c)$ gauge fields, coupled to the Standard Model:

if spontaneously chirally broken

- ▶ 3 Goldstone pions become the longitudinal W, Z
—> break EW symmetry
- ▶ The Higgs could be a $\bar{q}q$ (possibly qq) bound state
- ▶ Tower of additional hadronic states appear in experiments

What keeps the Higgs light?

Spontaneously broken symmetry —> massless Goldstone bosons

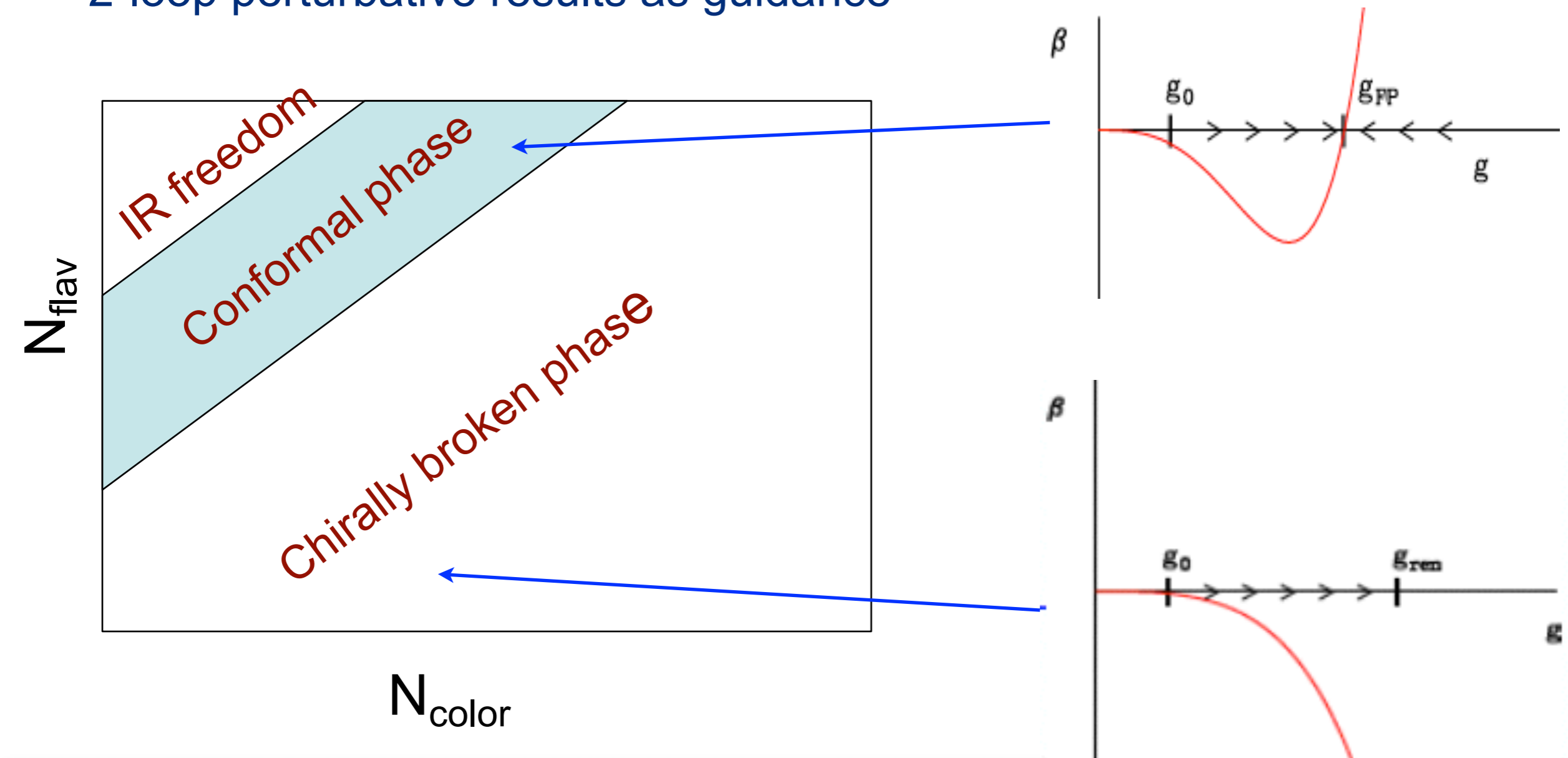
- **Flavor symmetry**: SSB leads to massless “pions”
- **Scale symmetry**: SSB leads to dilaton: **near-conformal models**

This is not the 80's technicolor!

Theory space:

$SU(N_{\text{color}} \geq 2)$ gauge fields + N_{flavor} fermions in some representation

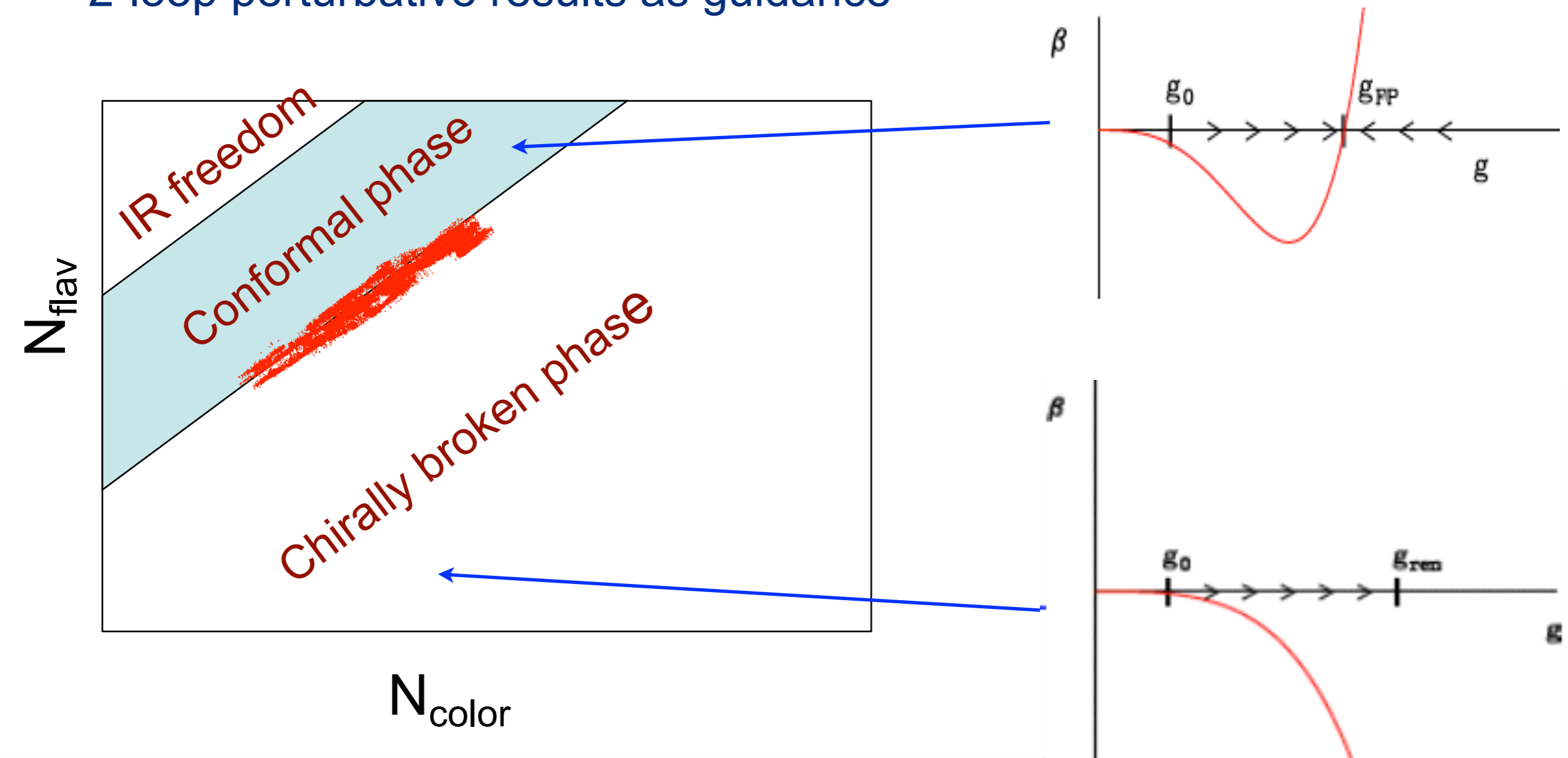
2-loop perturbative results as guidance



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2-loop perturbative results as guidance



Candidates

Any candidate model must be

- chirally broken - below the conformal window
- has 3 massless Goldstone pions
- most likely strongly coupled
- light scalar, enhanced condensate, etc

SU(2) gauge, adjoint fermions: $N_f = 2$: conformal ;

SU(2) gauge, fundamental fermions:

$N_f = 4$: chirally broken ; $N_f = 8$: conformal ; $N_f = 6$??

SU(3) gauge, fundamental fermions:

$N_f = 6$: chirally broken ; $N_f = 12$: conformal ; $N_f = 8$??

SU(3) gauge, sextet (sym rep) fermions: $N_f = 2$??

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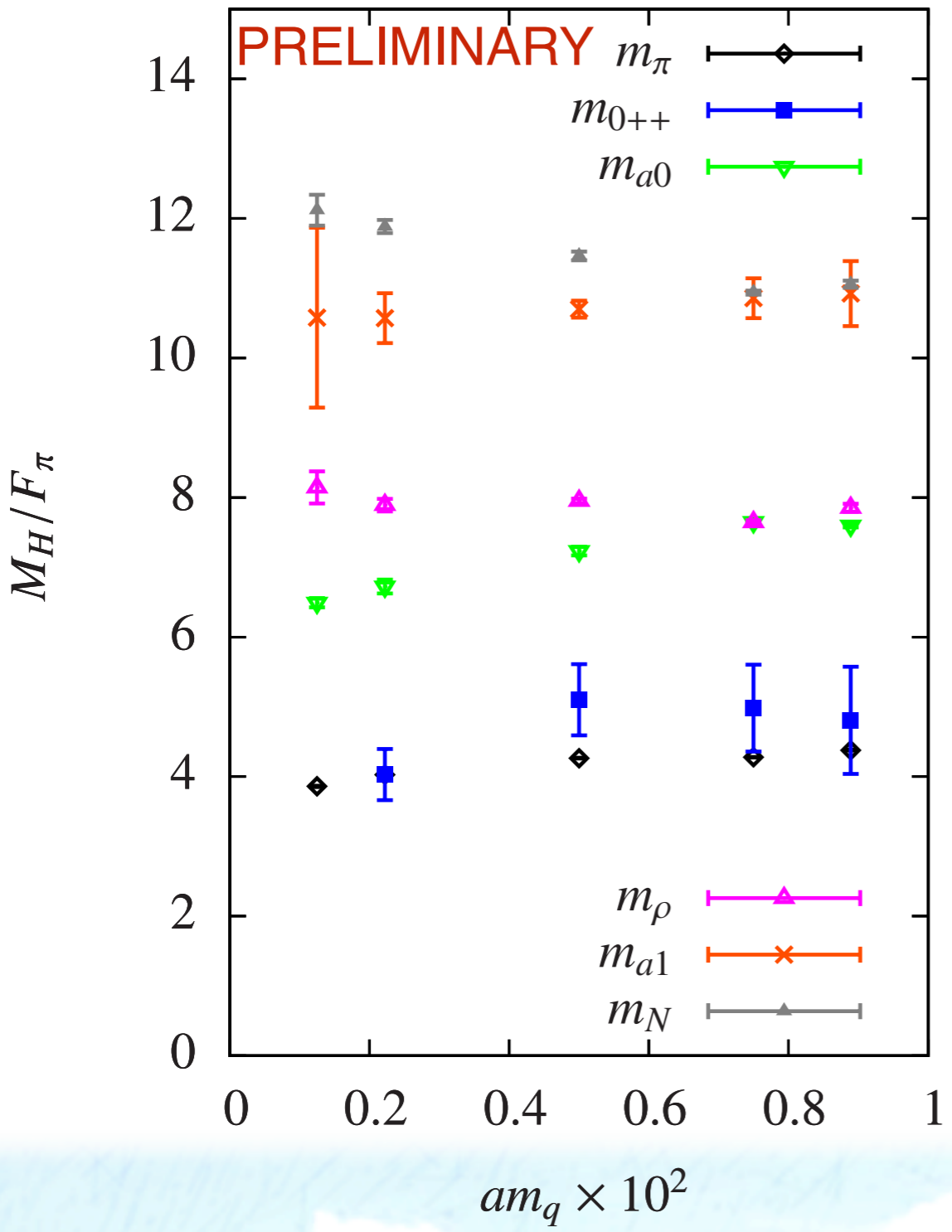
SU(3) gauge, sextet (sym rep) fermions: $N_f = 2$??

SU(3) with 8 fundamental fermions

8-flavor fundamental : old time favorite (?)

- close to the conformal window
(finite T studies inconclusive, AH, D. Schaich)
- if chirally broken, it has way too many Goldstones
- extensive studies by LatKMI collaboration and others
- new large scale study by LSD Coll. (1512.xxxx)

SU(3) with 8 fundamental fermions



$$M_H / F_\pi \text{ vs } m_f$$

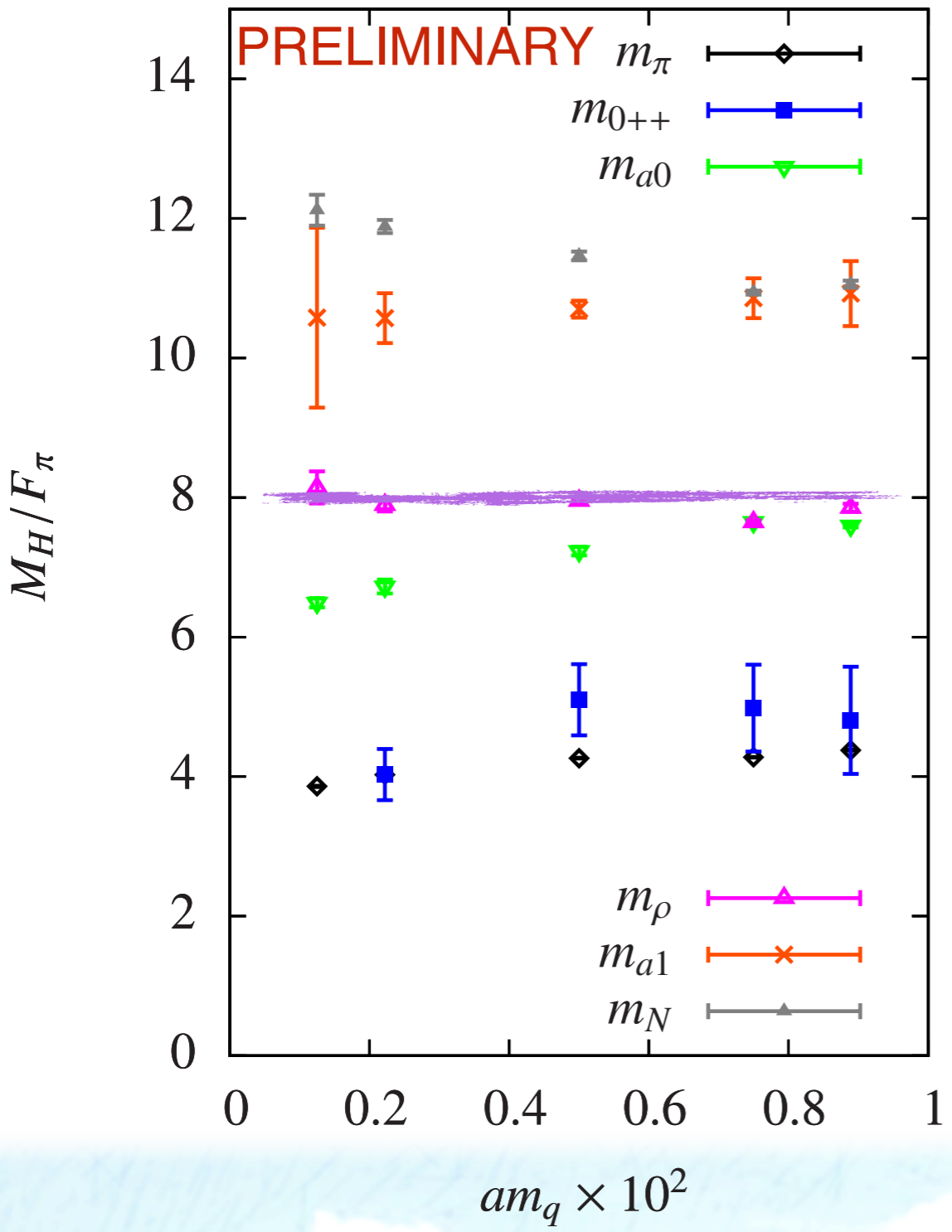
pion, rho, a0, a1, nucleon and 0^{++} scalar

In BSM scenario

$$F_\pi \sim vev = 250 GeV$$

With $M_\rho / F_\pi \approx 8$ this model predicts a 2TeV vector resonance!

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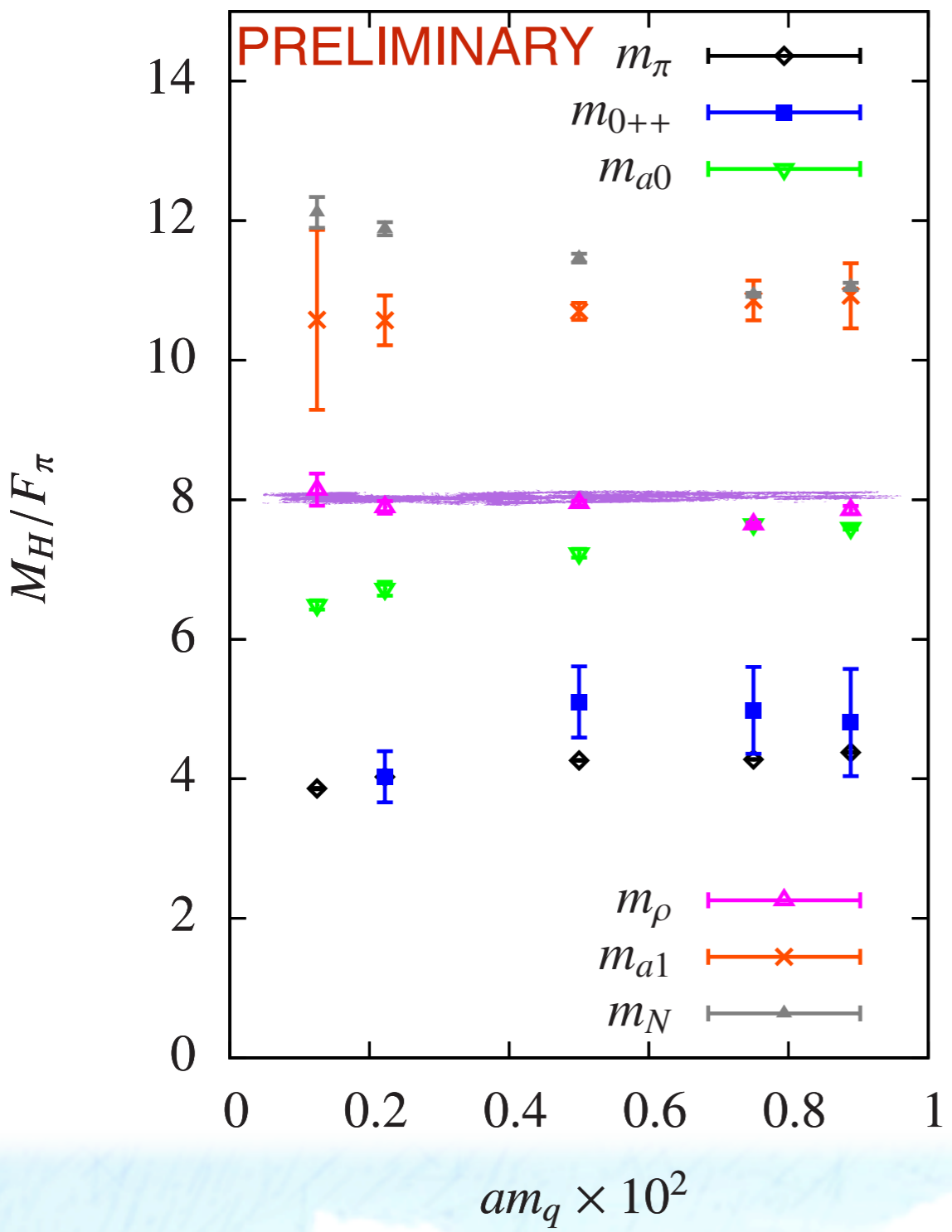
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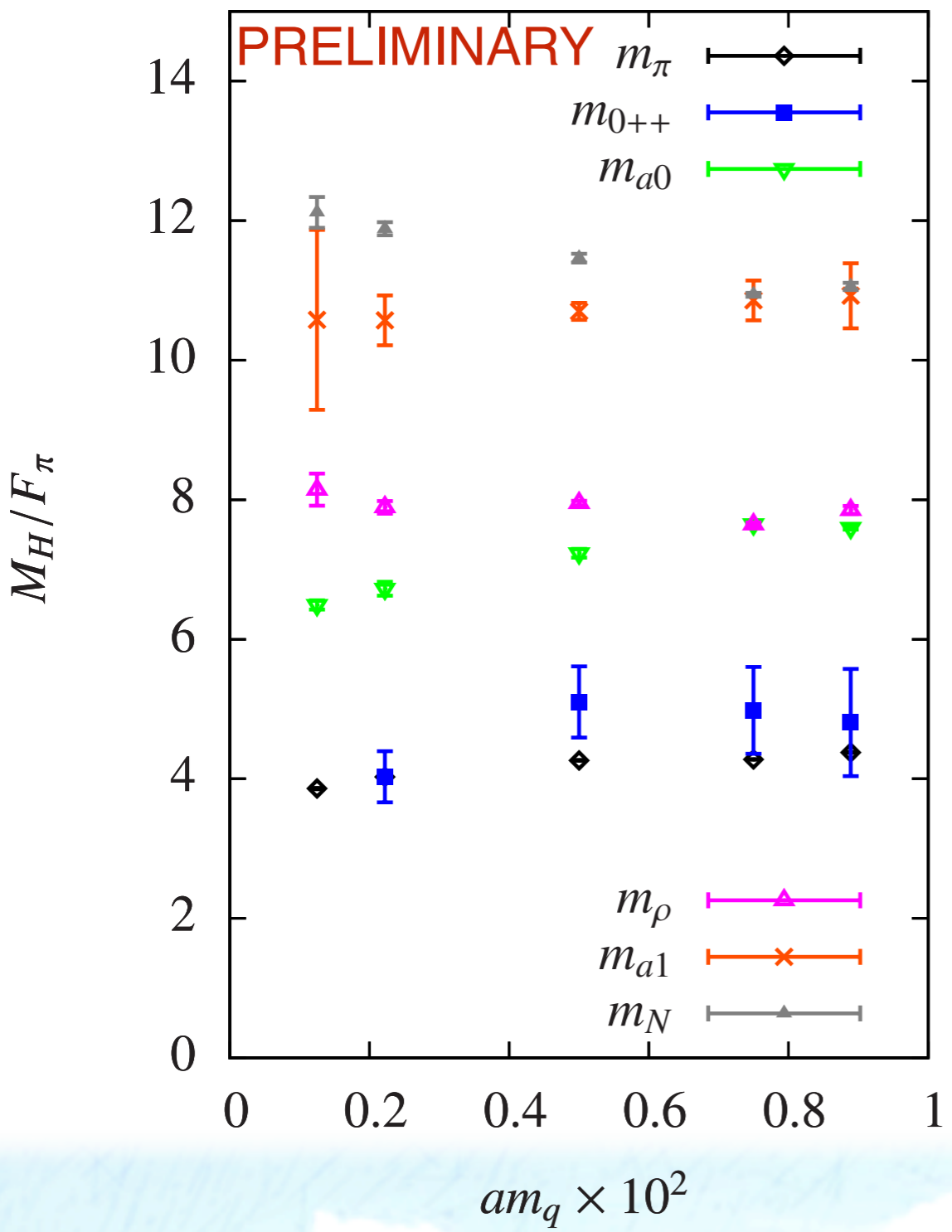


2TeV vector resonance
is very exciting

- But $M_\rho / F_\pi \approx 8$ in
- QCD
 - 4-flavor SU(3) fund.
 - 8 flavor SU(3) fund.
 - 2 flavor sextet
 - 12 flavor SU(3) fund.

Why?

SU(3) with 8 fundamental fermions



Several resonances are in the 2-3 TeV range

- 0^{++} scalar:
- it is degenerate with or lighter than the pion
 - in a chirally broken model it must “peel off” from the pion

How to take the chiral limit?

Near-conformal models with degenerate fermions

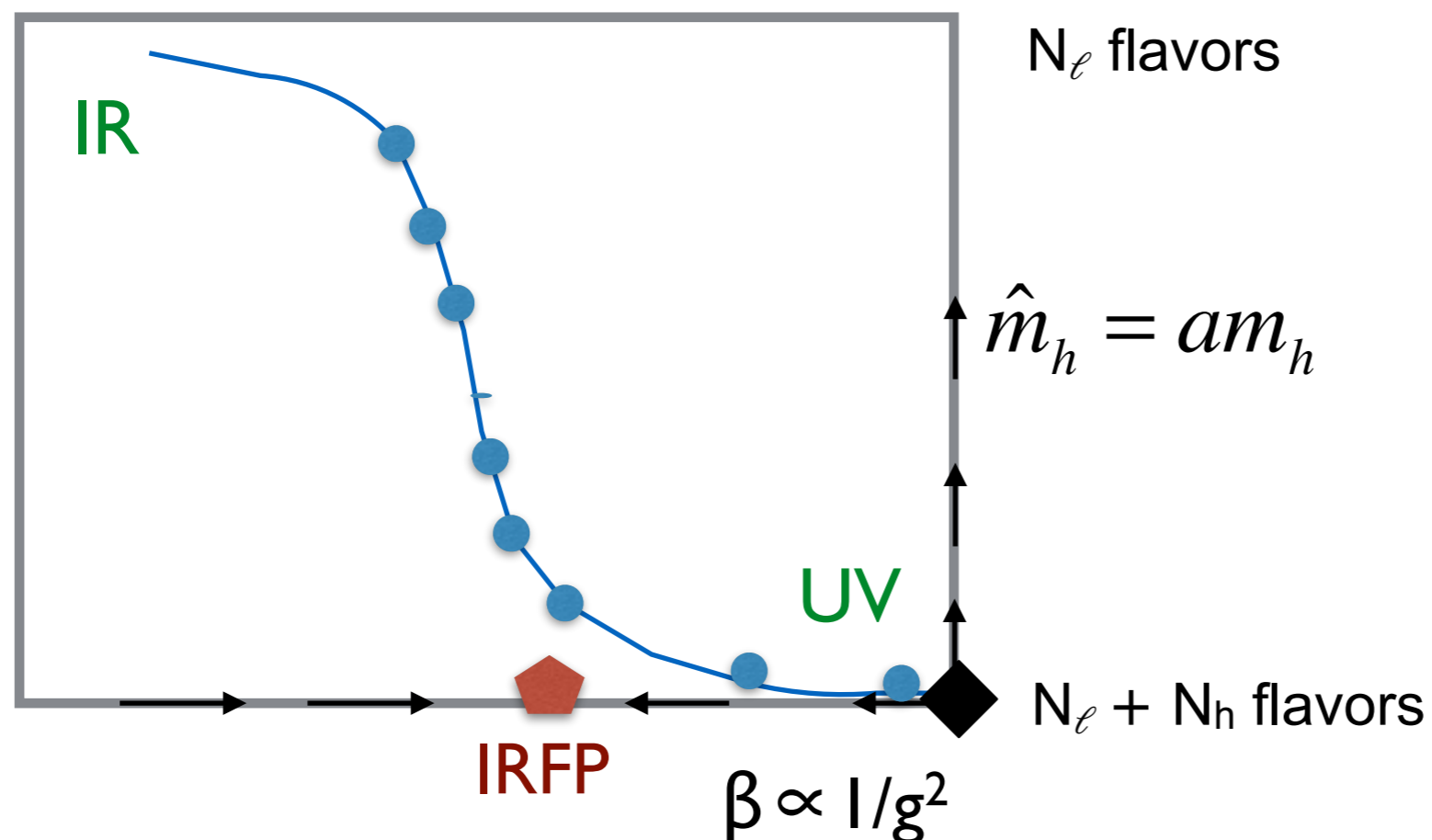
Issues, Questions:

- Is there a system with N_c color, N_f flavor in given rep that is close enough to the conformal window?
- Will the spontaneously broken conformal symmetry lead to a light dilaton? Or is a light scalar accidental?
- What happens to the extra Goldstone bosons?
- Is there walking? Are the EW constraints satisfied?

Mass split model : a different approach

Build a model on a conformal IRFP :

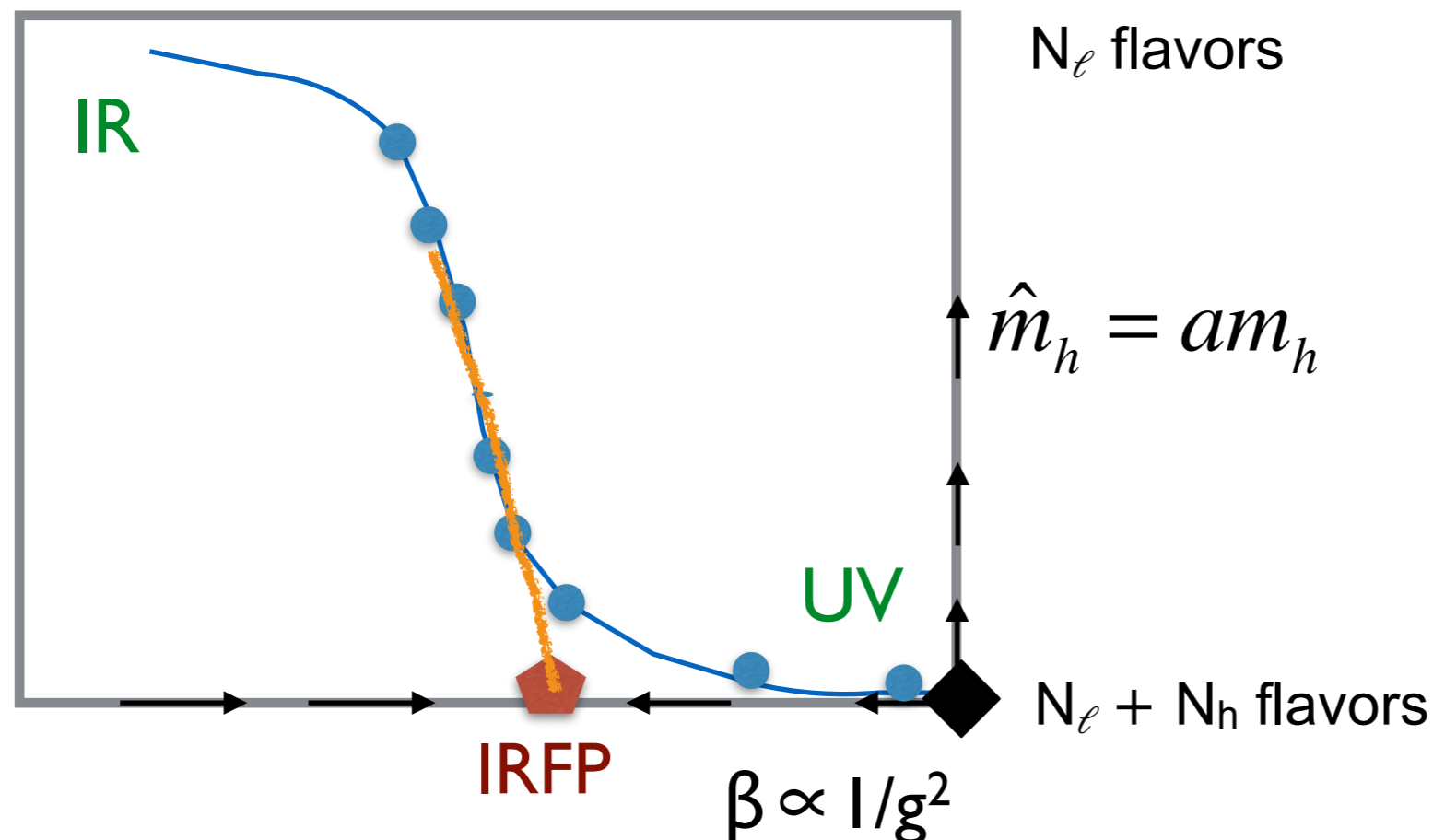
- Take N_f above the conformal window
- Split the masses: $N_f = N_\ell + N_h$
 - N_ℓ flavors are massless, $m_\ell = 0 \longrightarrow$ chirally broken
 - N_h flavors are massive, m_h varies \longrightarrow decouple in the IR



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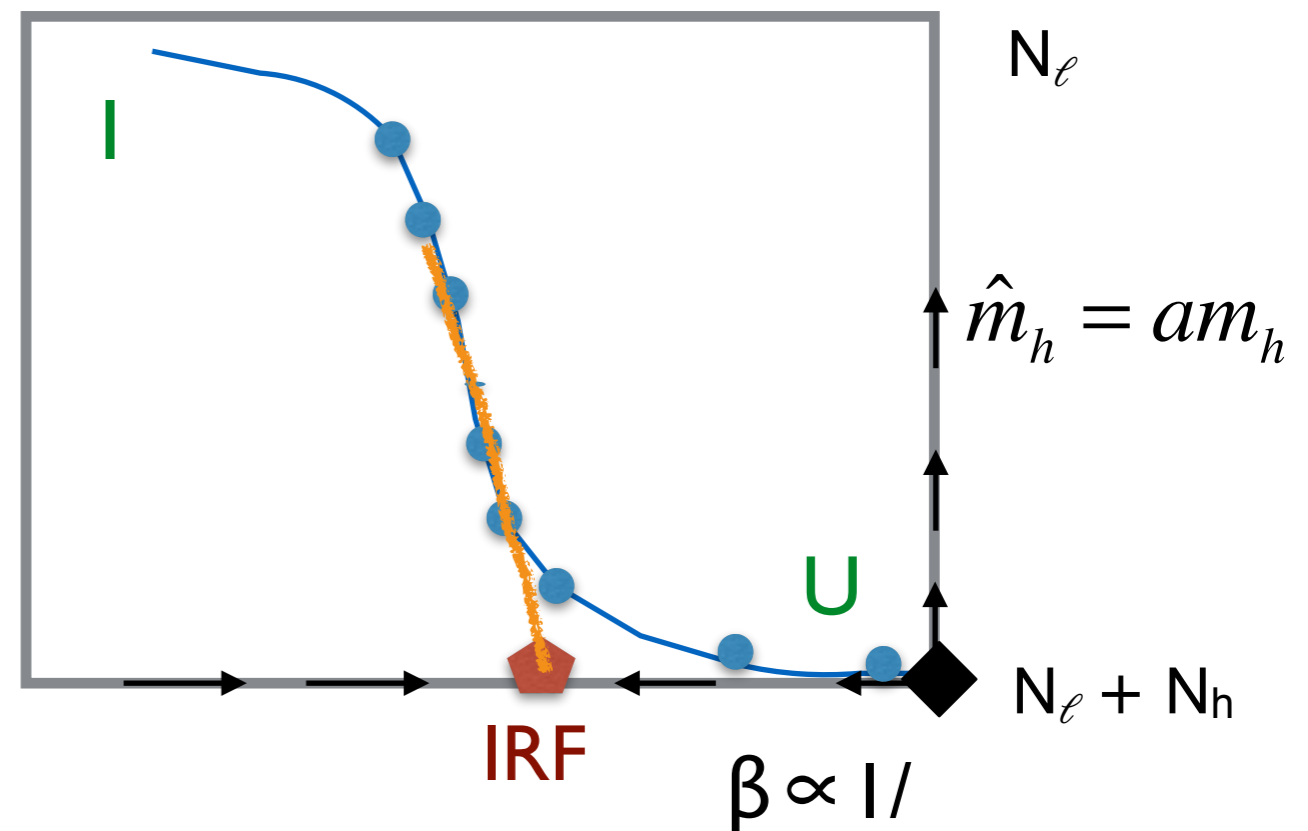
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Wilson renormalization group description

When discussing a conformal model it is convenient to use Wilson RG description:

- use bare parameters at the UV scale
 - define a continuum (infinite cut-off) system by tuning to criticality
-
- Critical surface is at $m_h = 0$;
 - There is a fixed point (IRFP) somewhere on this infinite dimensional surface —
 - The location of the IRFP is scheme dependent



Spectrum: Hyperscaling

In **conformal** systems Wilson RG considerations predict m_h dependence:

If the scale changes as $\mu \rightarrow \mu' = \mu / b$, $b > 1$
the couplings run as

$$\hat{m}(\mu) \rightarrow \hat{m}(\mu') = b^{y_m} \hat{m}(\mu) \quad (\text{increases})$$
$$g \rightarrow g^*$$

Any 2-point correlator at large b behaves

$$C_H(t; g_i, \hat{m}_i, \mu) = b^{-2y_H} C_H(t/b; g^*, b^{y_m} \hat{m}, \mu)$$

since $C_H(t) \propto e^{-M_H t}$,

$$aM_H \propto (\hat{m})^{1/y_m} \quad (\text{hyperscaling})$$

Amplitudes (F_π) also show hyperscaling

Hyper scaling in mass split systems

Nothing changes in the Wilson RG arguments if some of the masses remain massless:

mass split systems show the same hyperscaling in the $m_\ell = 0$ limit

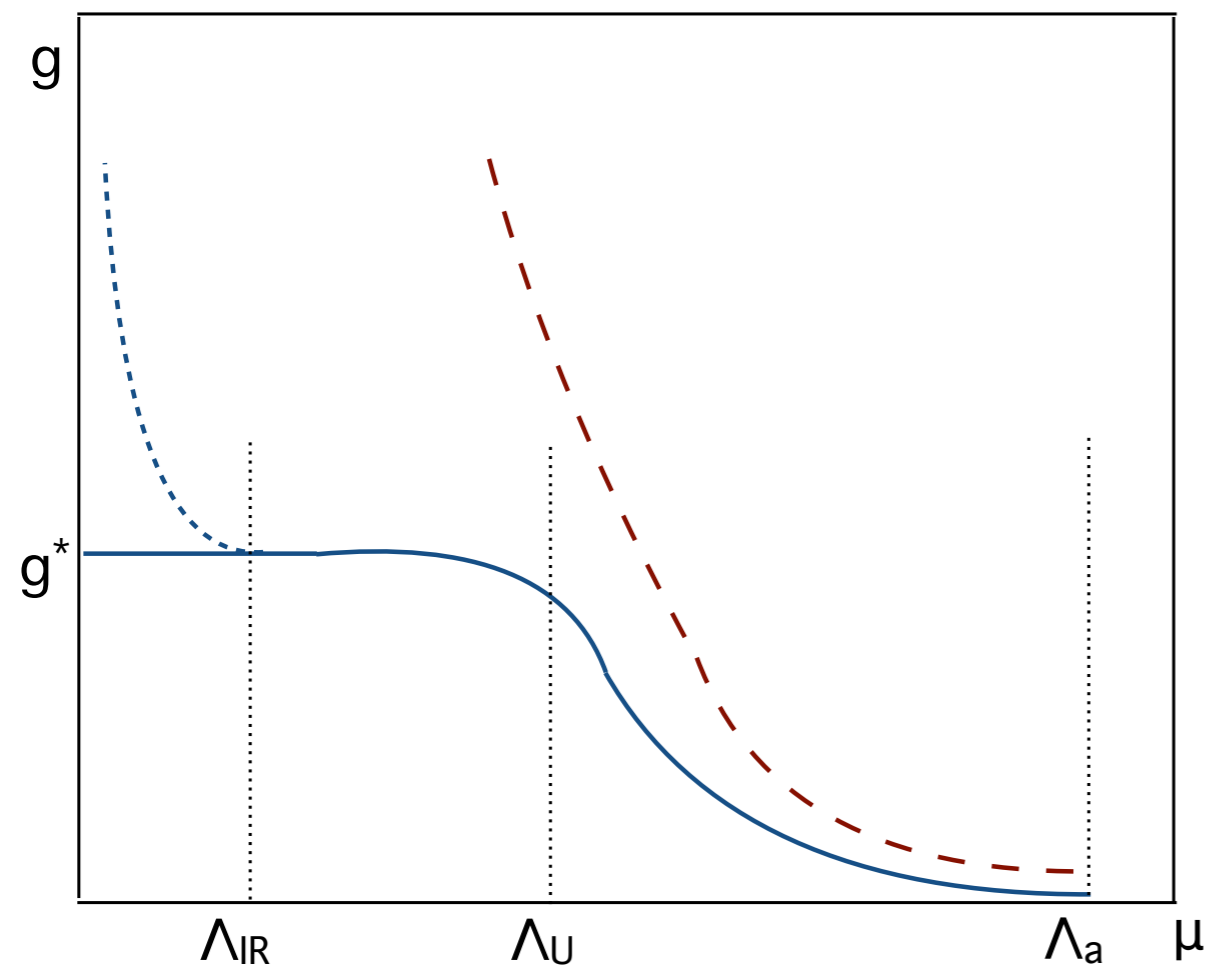
$$aM_H \propto (\hat{m})^{1/y_m}$$

(M_H can be all light, heavy or mixed fermion hadrons)

—> Ratios like M_H / F_π are independent of m_h

Mass split model - running coupling

RG flows predict the running coupling:

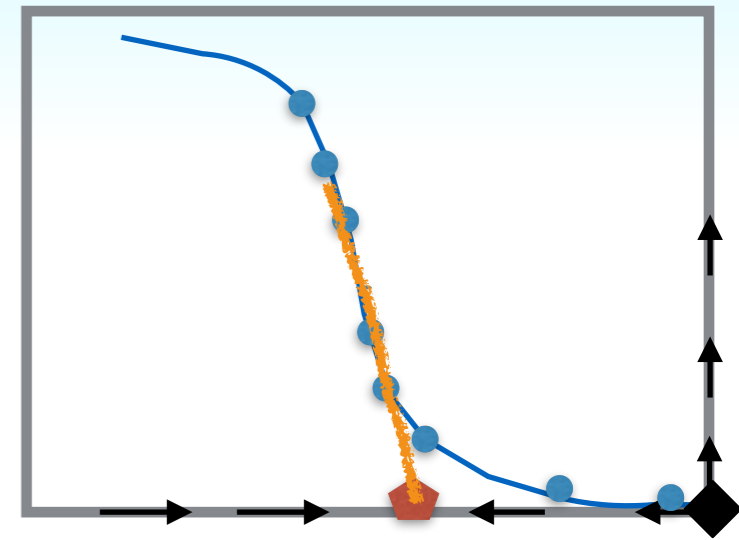


3 regions:

- UV :
from cut-off to $g \sim g^*$
- walking: m_h small, $g \sim g^*$
- IR :
heavy flavors decouple,
 N_ℓ light flavors are
chirally broken

walking can be tuned by

$$m_h \rightarrow 0$$



$N_\ell + N_h$ mass split systems

ideally

$N_\ell + N_h = 2 + 8$: $N_f = 10$ is close to the conformal window

(or

$N_\ell + N_h = 2 + 1 + 1 + 1 + 1 \dots$: cascade from $N_f = 16$ to $N_f = 2$)

Pilot study:

$N_\ell + N_h = 4 + 8$: conformal in the UV, $N_f=4$ flavor in the IR

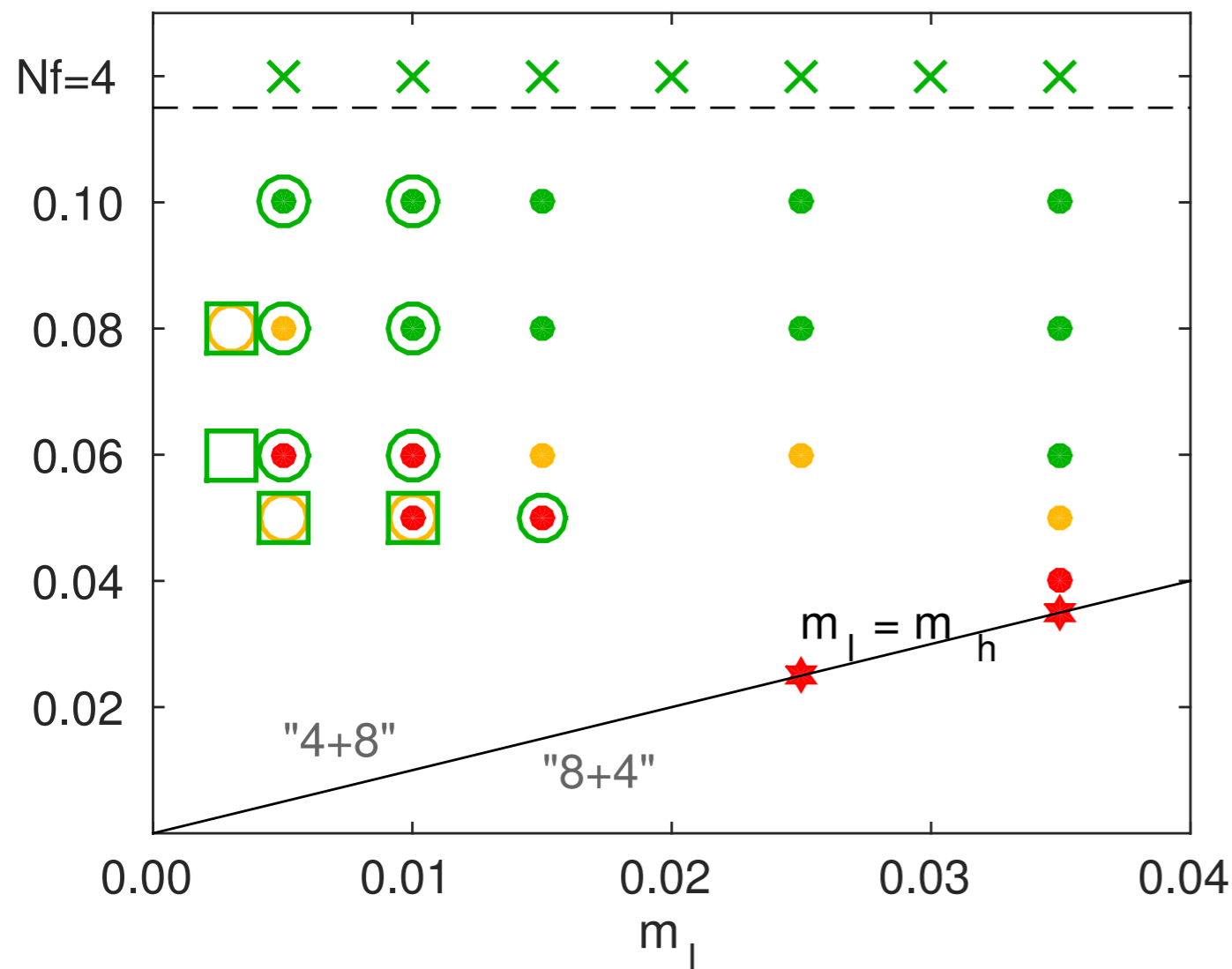
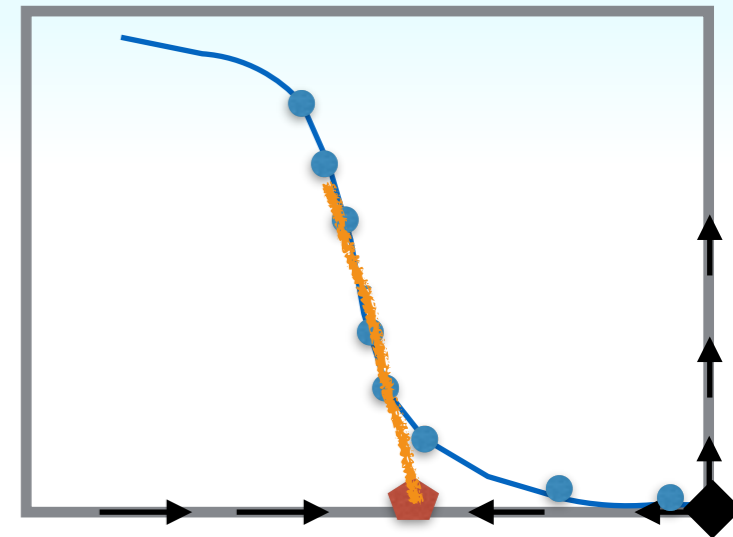
in collaboration with R. Brower, C. Rebbi, E. Weinberg, O. Witzel

arXiv:1411.3243, 1511.xxx

Why $4+8$? We use staggered fermions:
4 and 8 flavors do not require rooting

$N_\ell + N_h = 4 + 8$: Parameter space

- $\beta=4.0$ (close to the 12-flavor IRFP)
- $m_h = 0.100, 0.080, 0.060, 0.050$
- $m_\ell = 0.003, 0.005, 0.010, 0.015, 0.025, 0.035$



Volumes :

$24^3 \times 48$, (dots)

$32^3 \times 64$ (circle), $36^3 \times 64$

$48^3 \times 96$ (square)

Color: volume OK / marginal / squeezed

20-40,000 MDTU

Running coupling

Gradient flow transformation defines a renormalized coupling

arXiv:1006.4518

$$g_{GF}^2(\mu = \frac{1}{\sqrt{8t}}) = \frac{1}{\mathcal{N}} t^2 \langle E(t) \rangle$$

t: flow time;

E(t):energy density

g_{GF}^2 is used for scale setting as

$$g_{GF}^2(t = t_0) = \frac{0.3}{\mathcal{N}}$$

It is appropriate to determine the renormalized running coupling

- on large enough volumes
- at large enough flow time
- in the continuum limit

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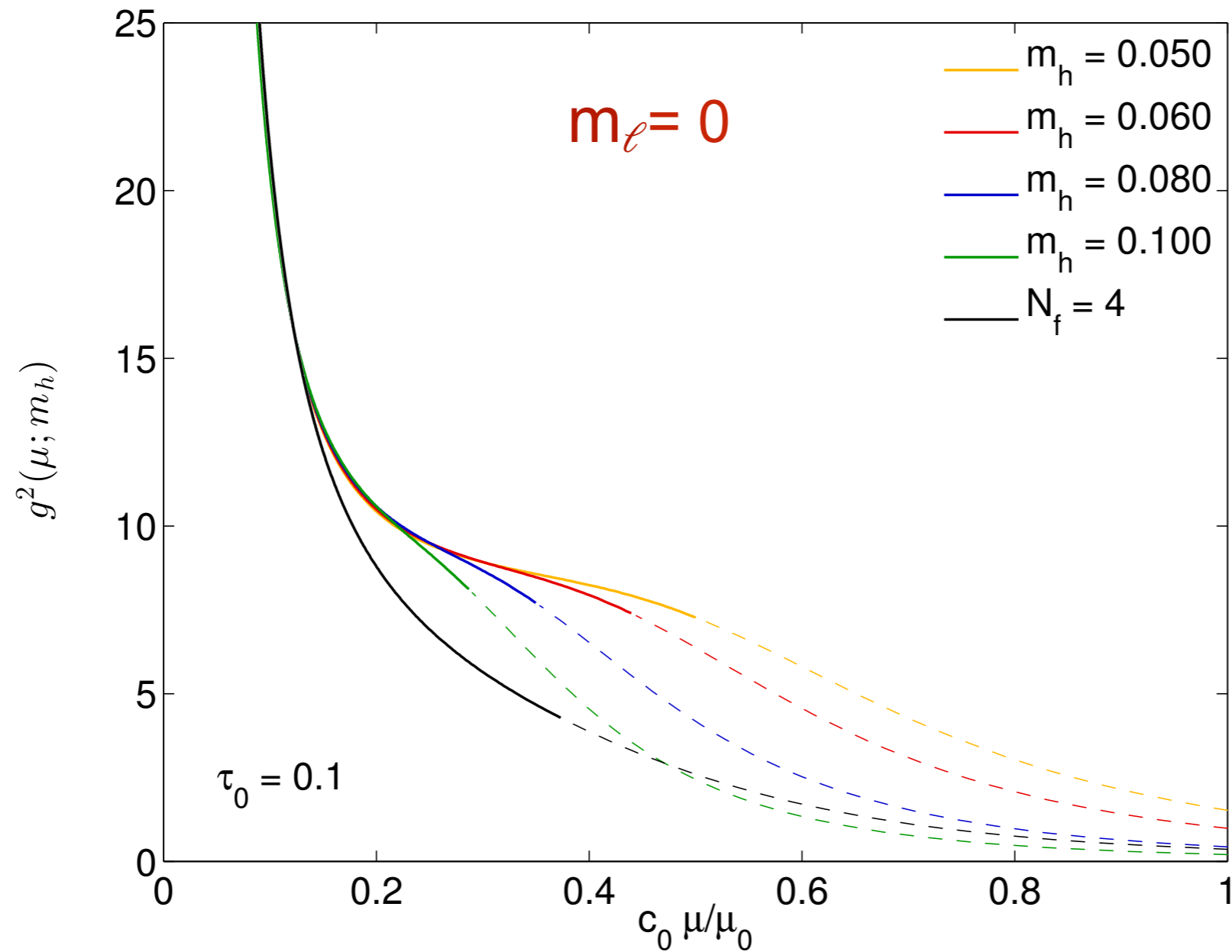
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} use t-shift improved coupling

Improved running coupling : 4+8 flavors



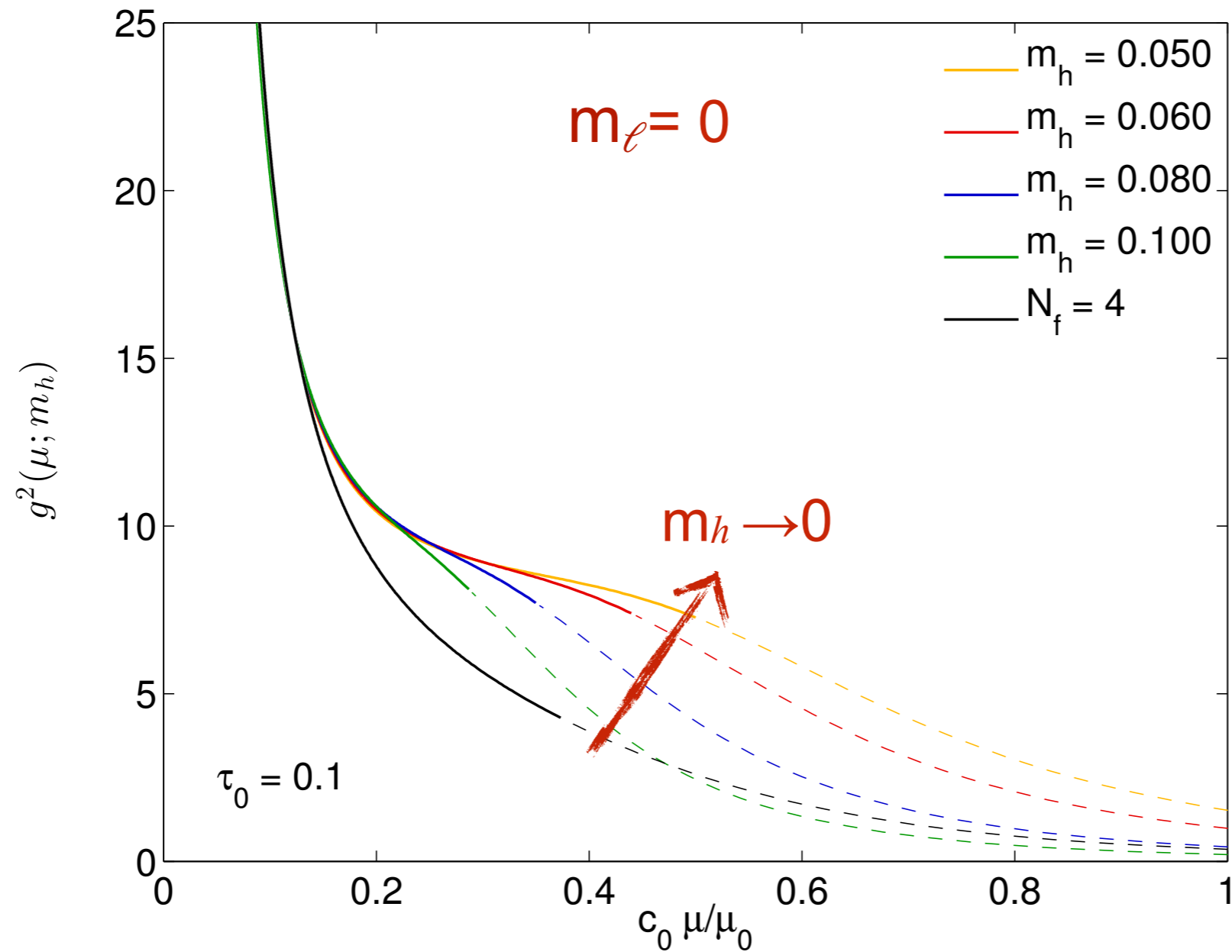
There are error bars on this plot!

$N_f=4$: running fast

$g_{GF}^2(\mu)$ develops a “shoulder” as $m_h \rightarrow 0$: this is walking !

Walking range can be tuned arbitrarily with m_h

Improved running coupling : 4+8 flavors



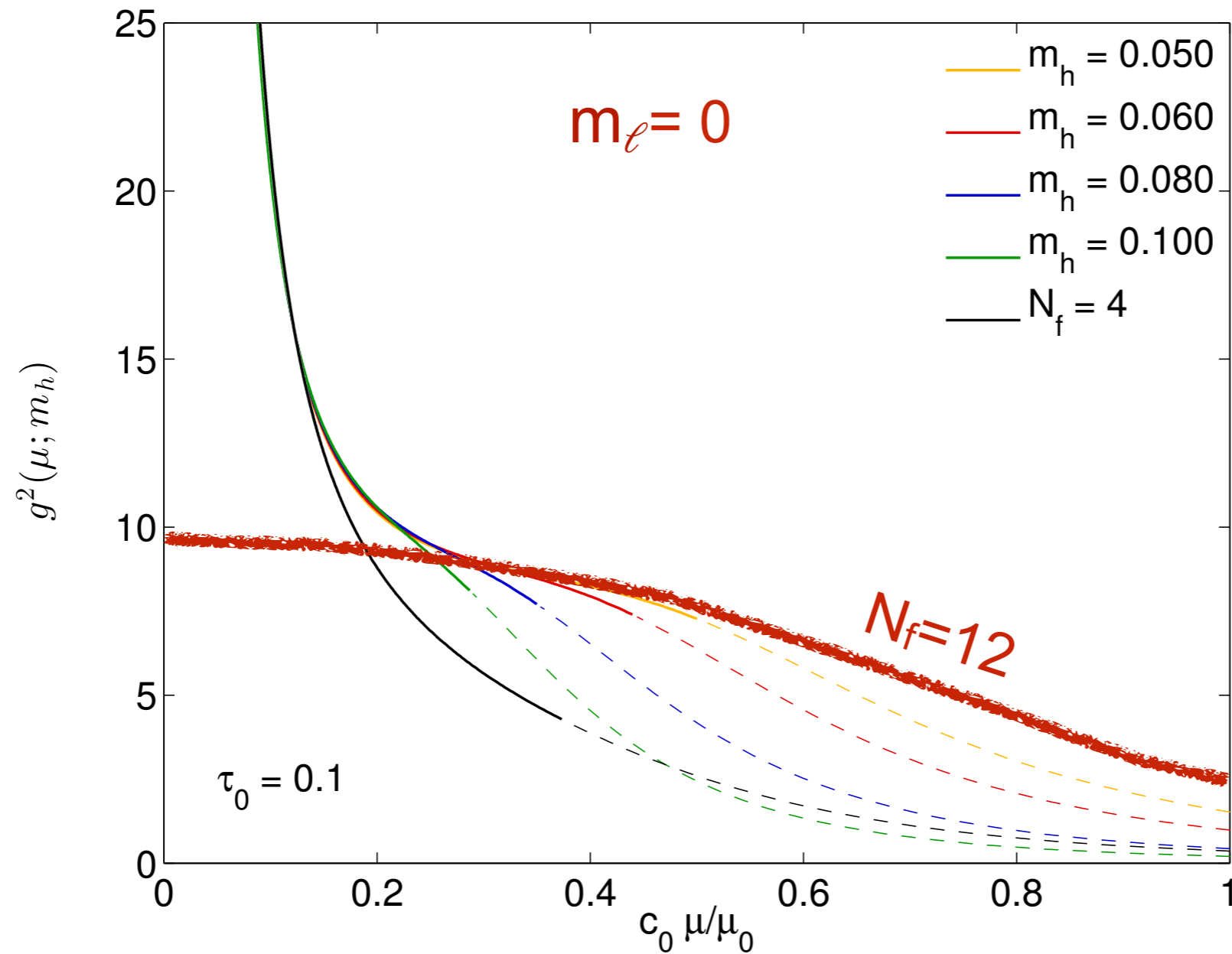
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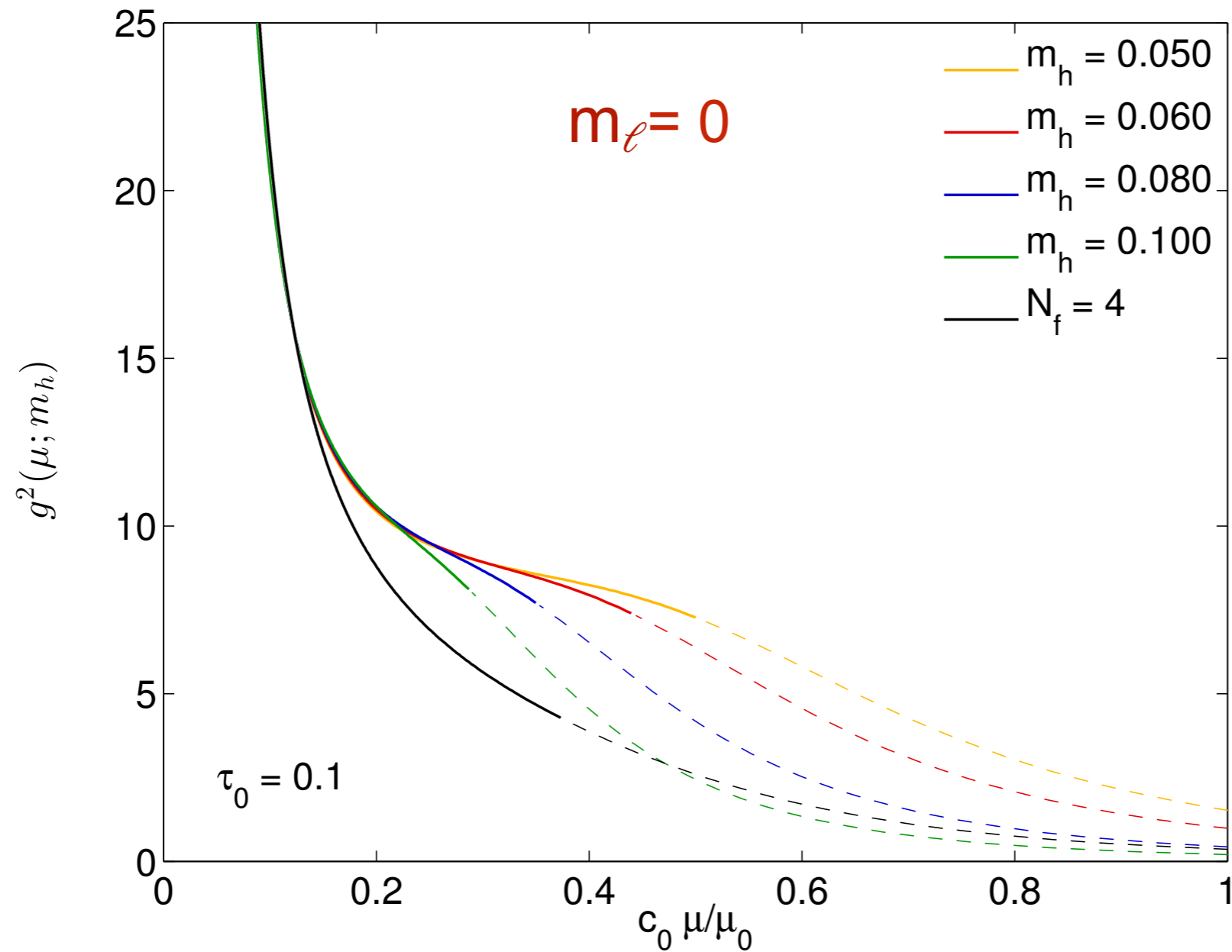
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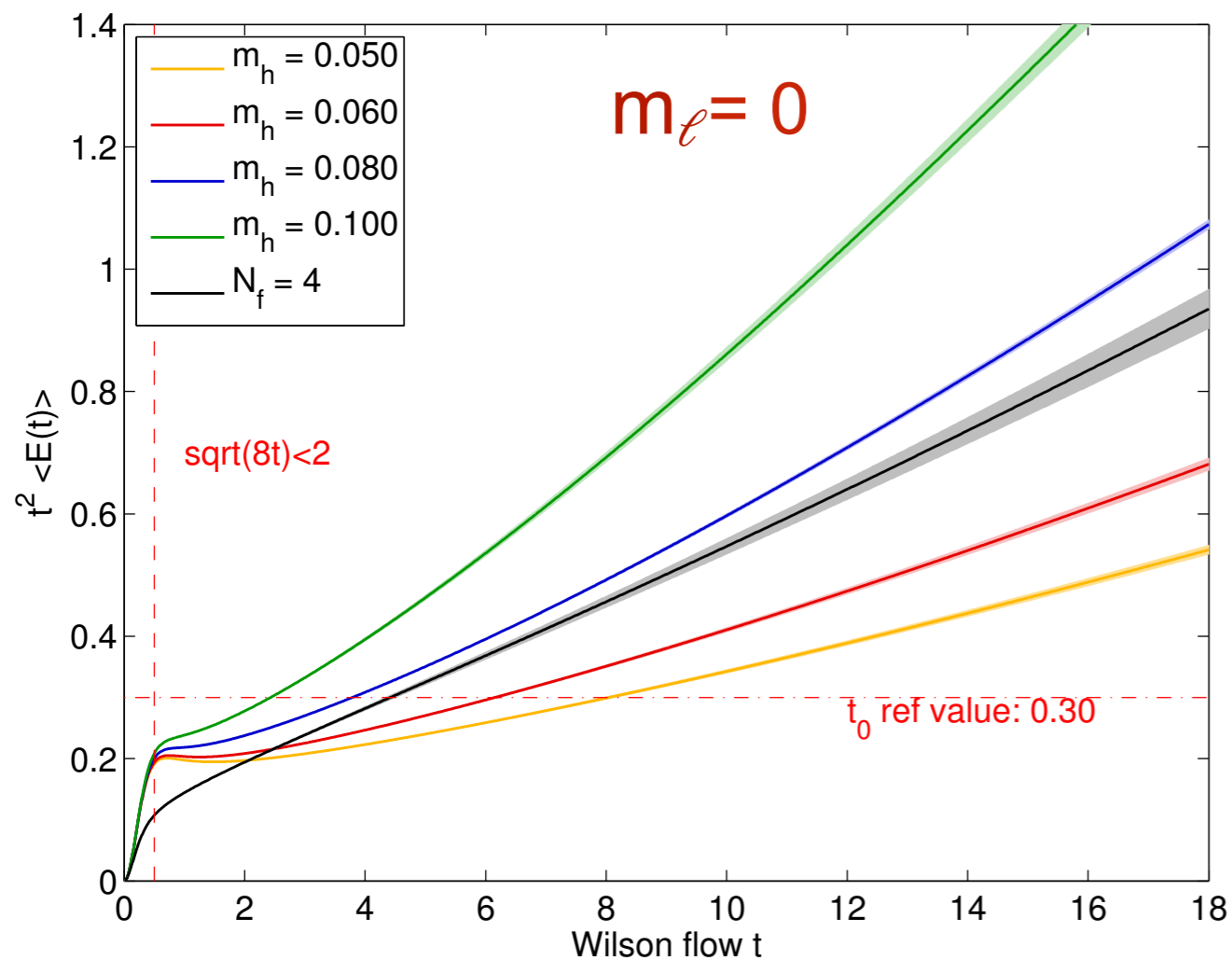
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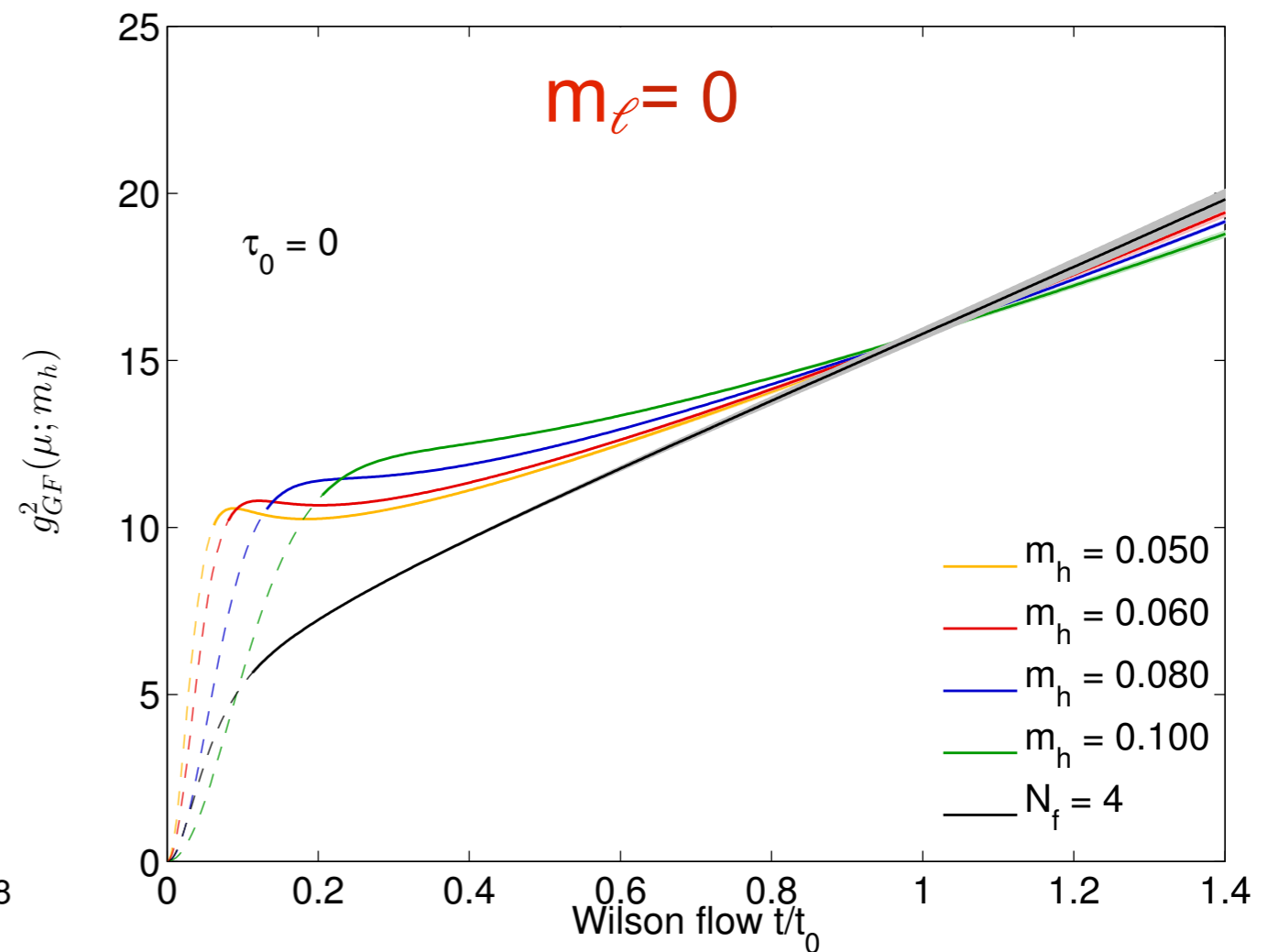
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Running coupling

$t^2 \langle E(t) \rangle$ in the chiral limit
at various m_h values



$g_{GF}^2(t/t_0)$ rescaled by t_0
at various m_h values



Rescaling forces the renormalized couplings to agree at t_0
Fan-out before and after are due to cut-off lattice artifacts

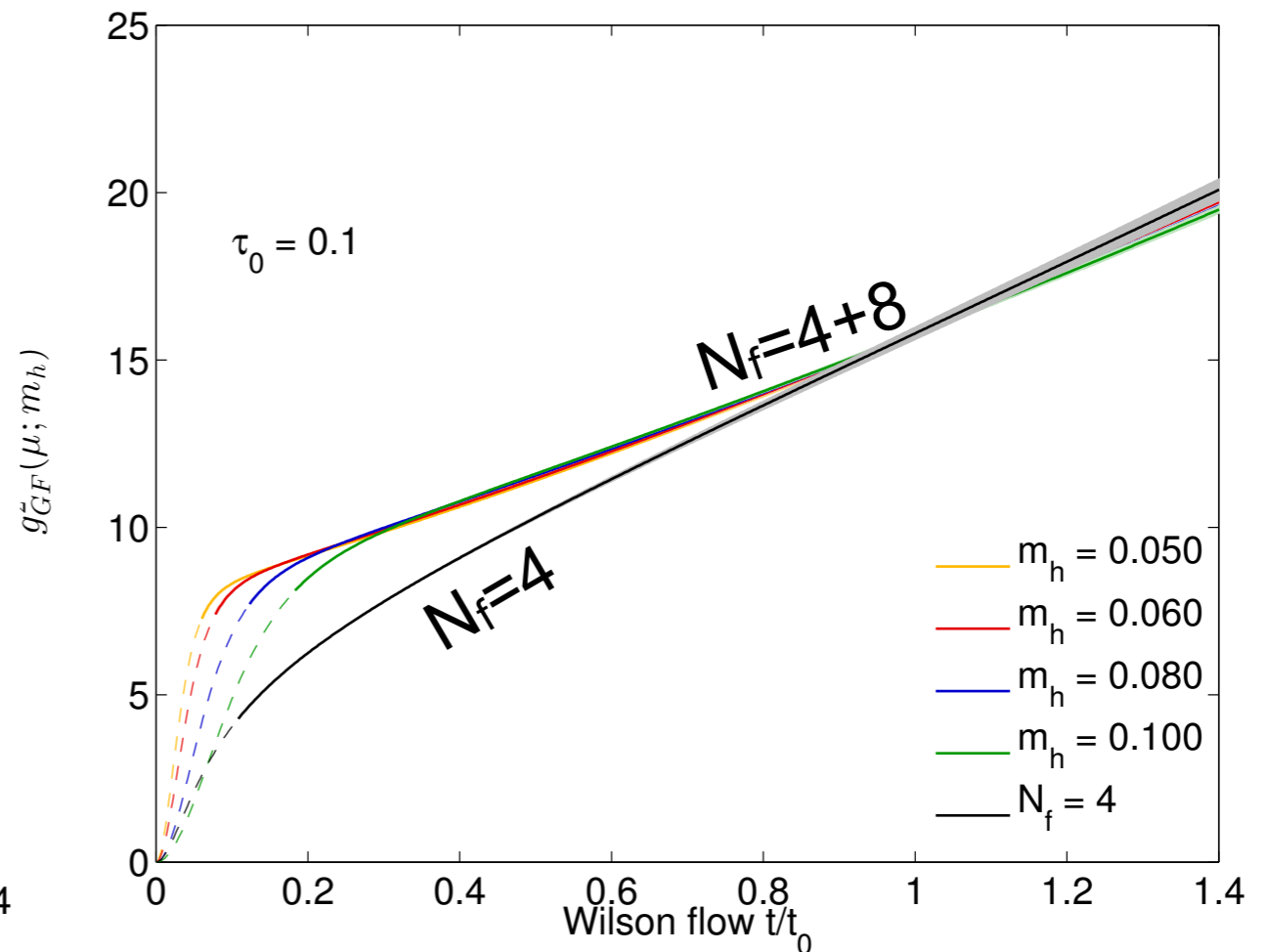
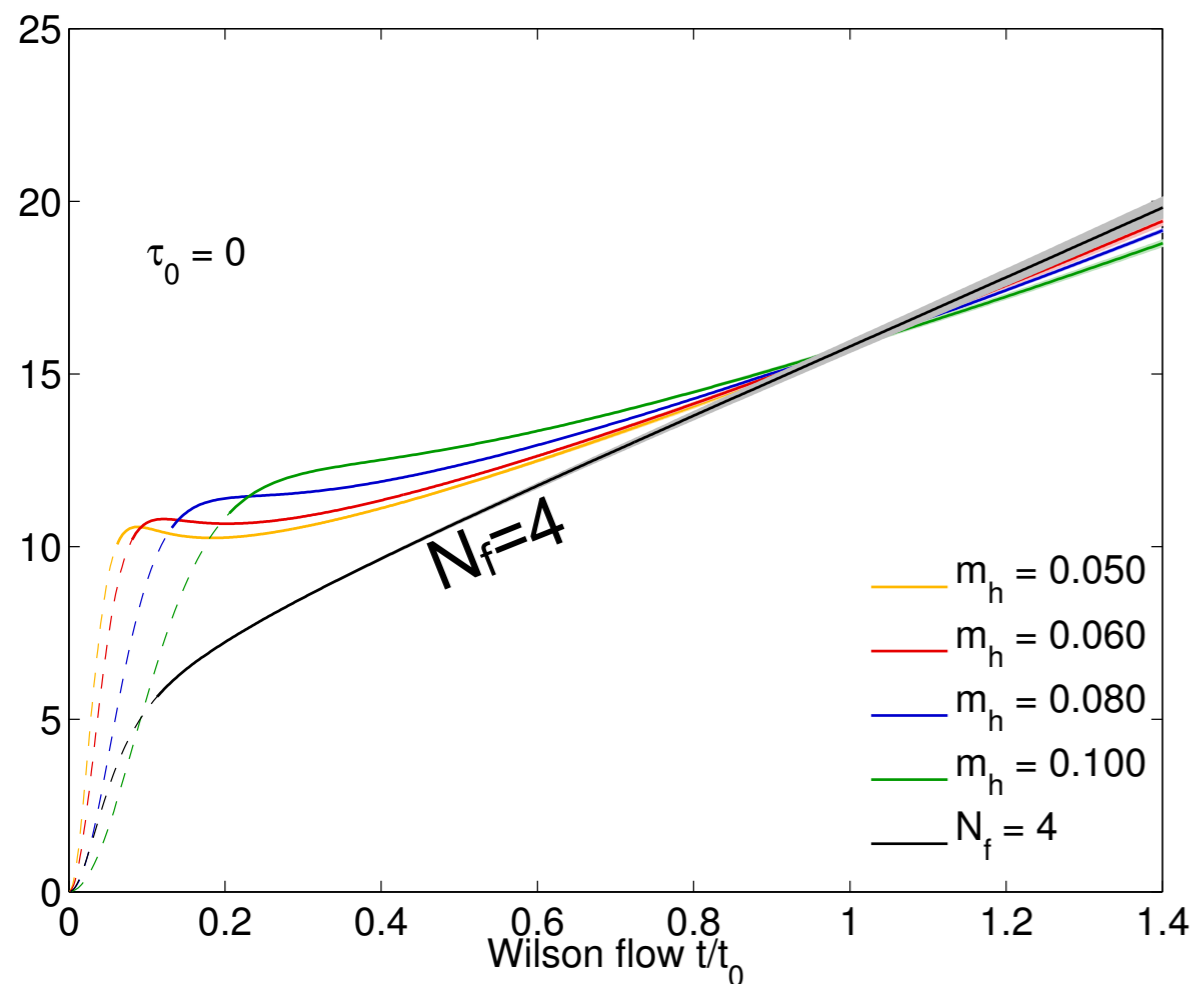
Improved running coupling

t-shift improved running coupling

$$\tilde{g}_{GF}^2(\mu = \frac{1}{\sqrt{8t}}) = \frac{1}{\mathcal{N}} t^2 \langle E(t + \tau_0) \rangle$$

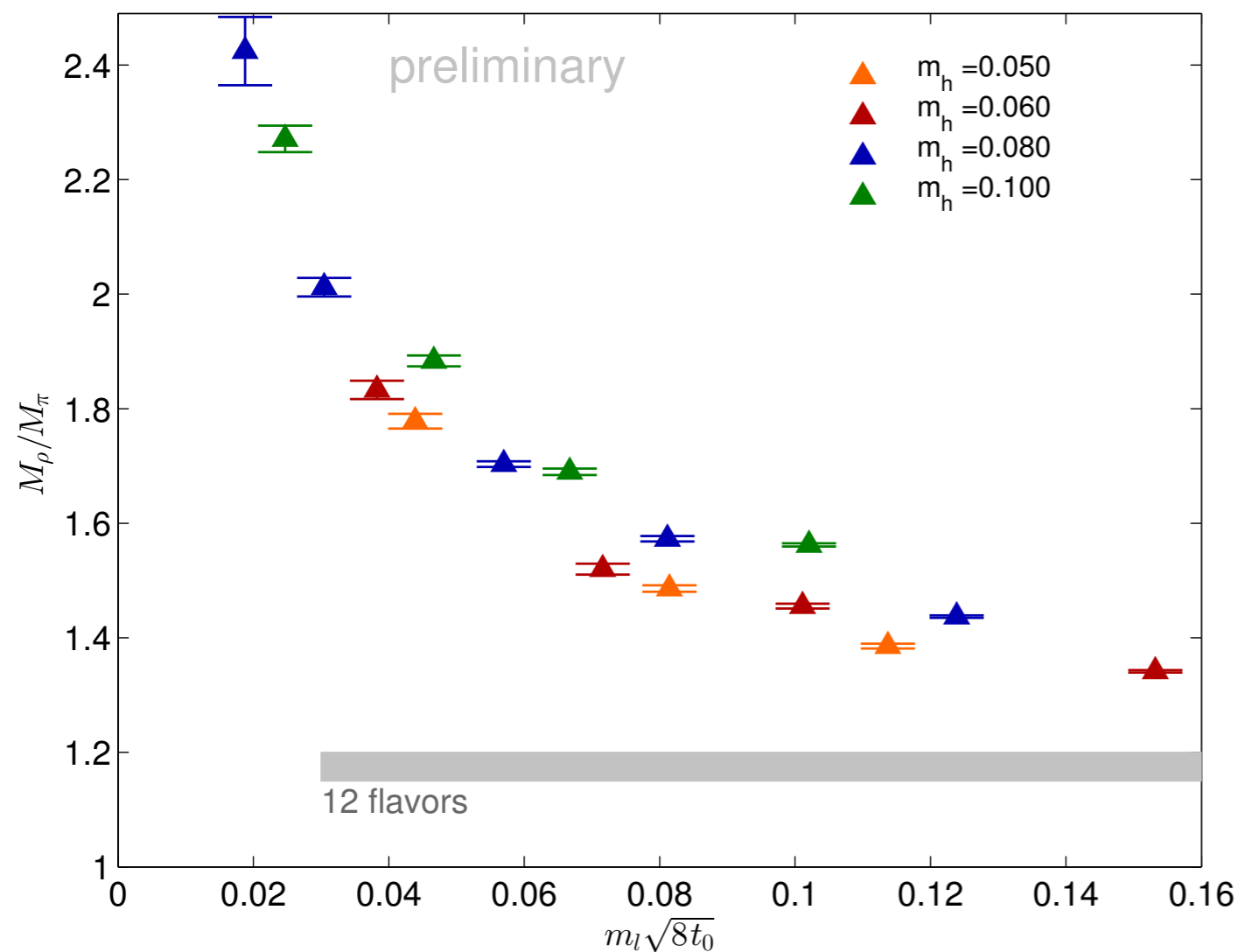
by adjusting τ_0 most cut-off effects can be removed

(1404.0984, 1501.07848)



Is the system chirally broken ?

M_ρ/M_π shows that we approach the chiral regime

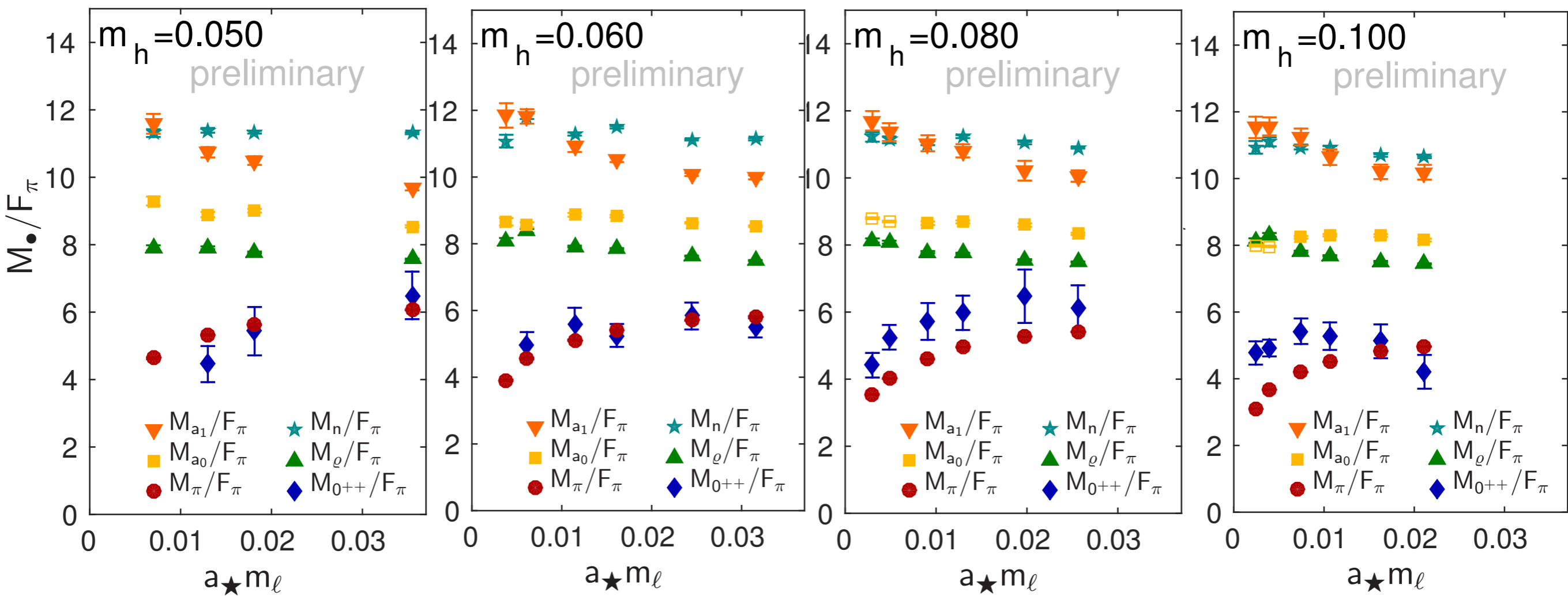


$< N_f=12$ predicts an almost constant ratio (as should be in a conformal system)

(arXiv:1401.0195)

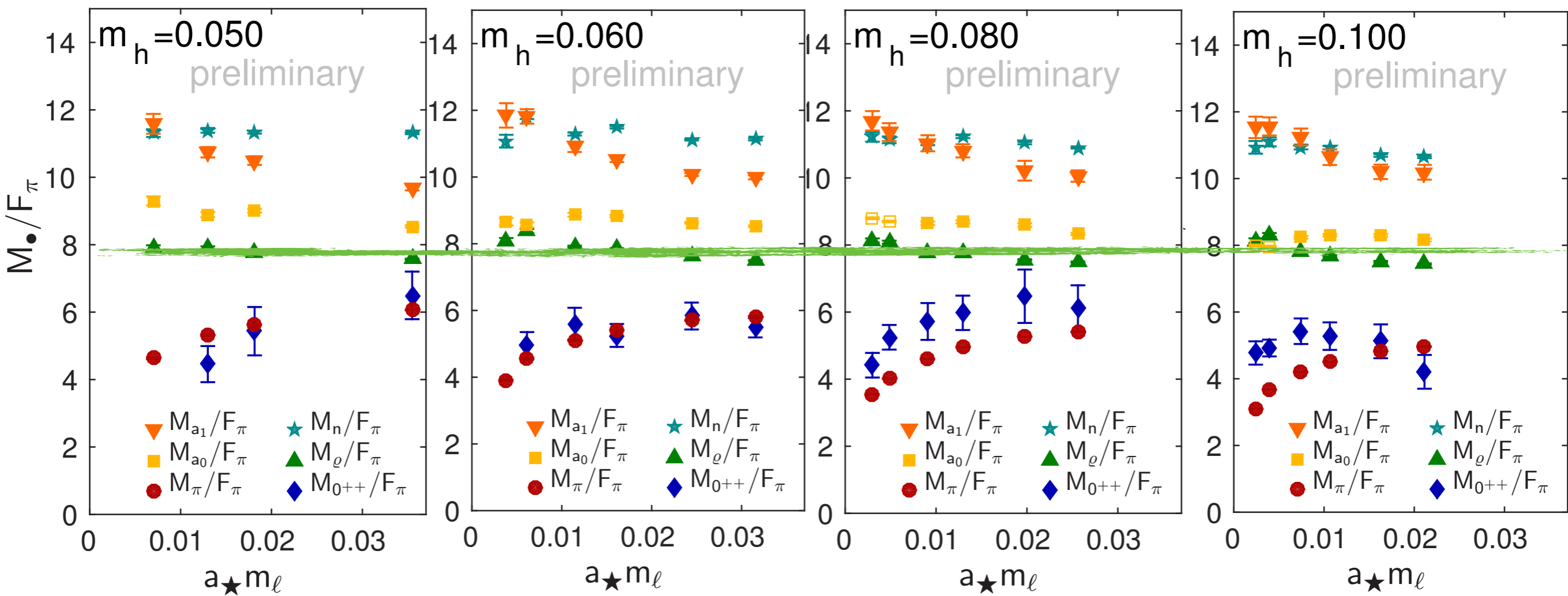
Spectrum, 4+8 flavors

Ratios M_H / F_π appear independent of m_h

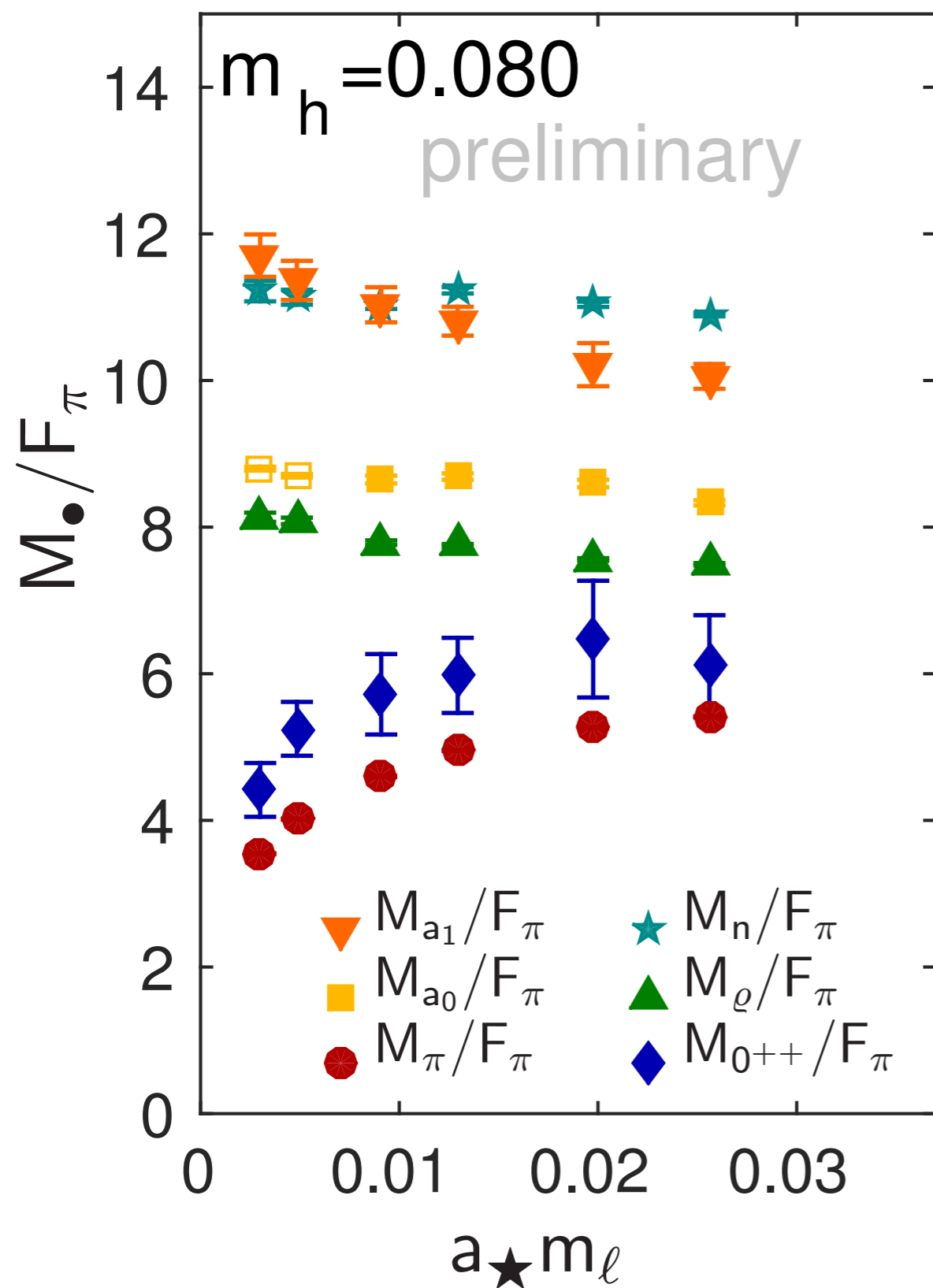


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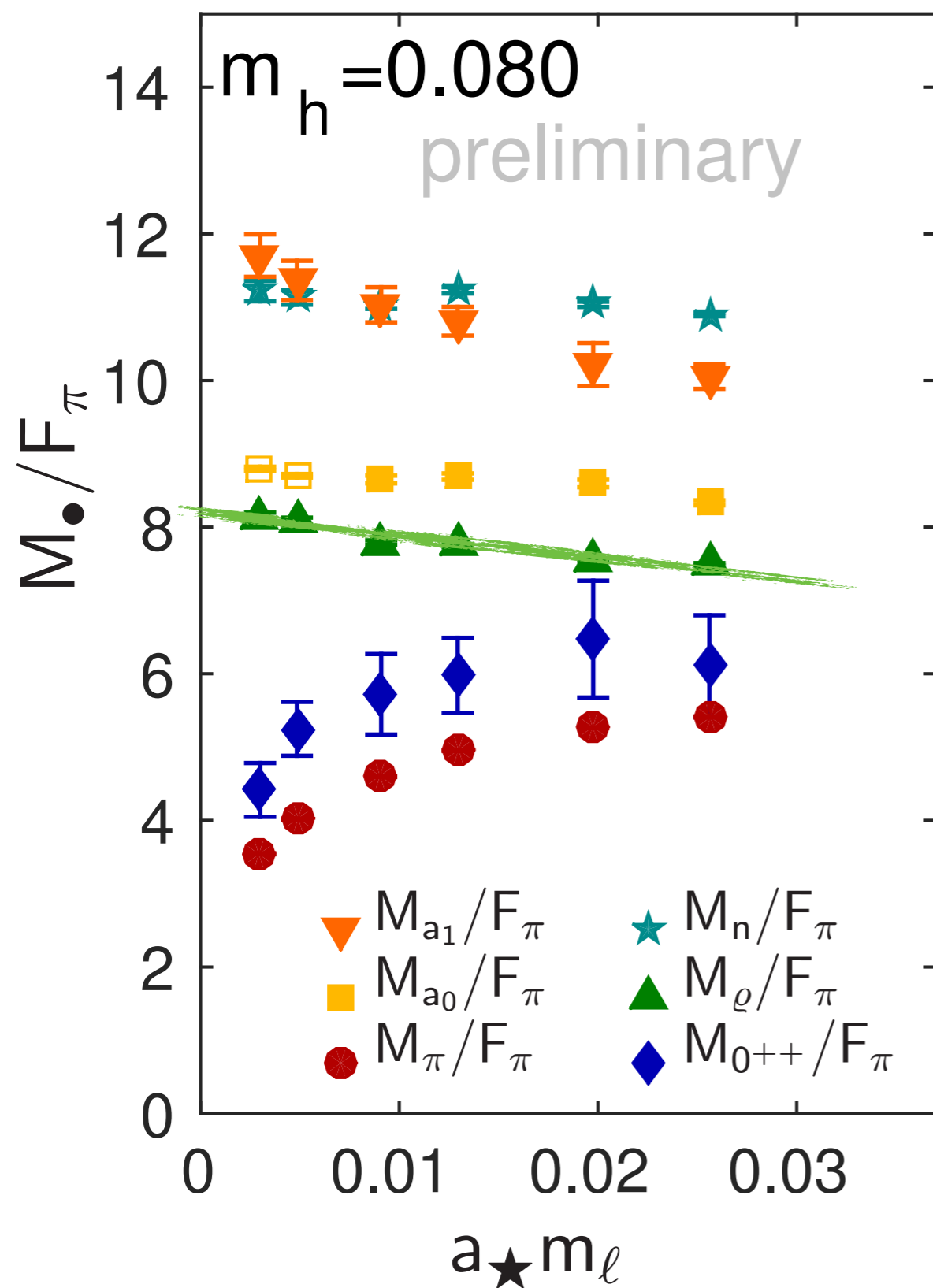
pion, rho, a0, a1, nucleon
and 0^{++} scalar

The ratios are very similar to
 $N_f=8$ and sextet

$$M_{\rho} / F_{\pi} \approx 8$$

0^{++} is just above, closely
following the pion -
chiral limit????

Spectrum, 4+8 flavors



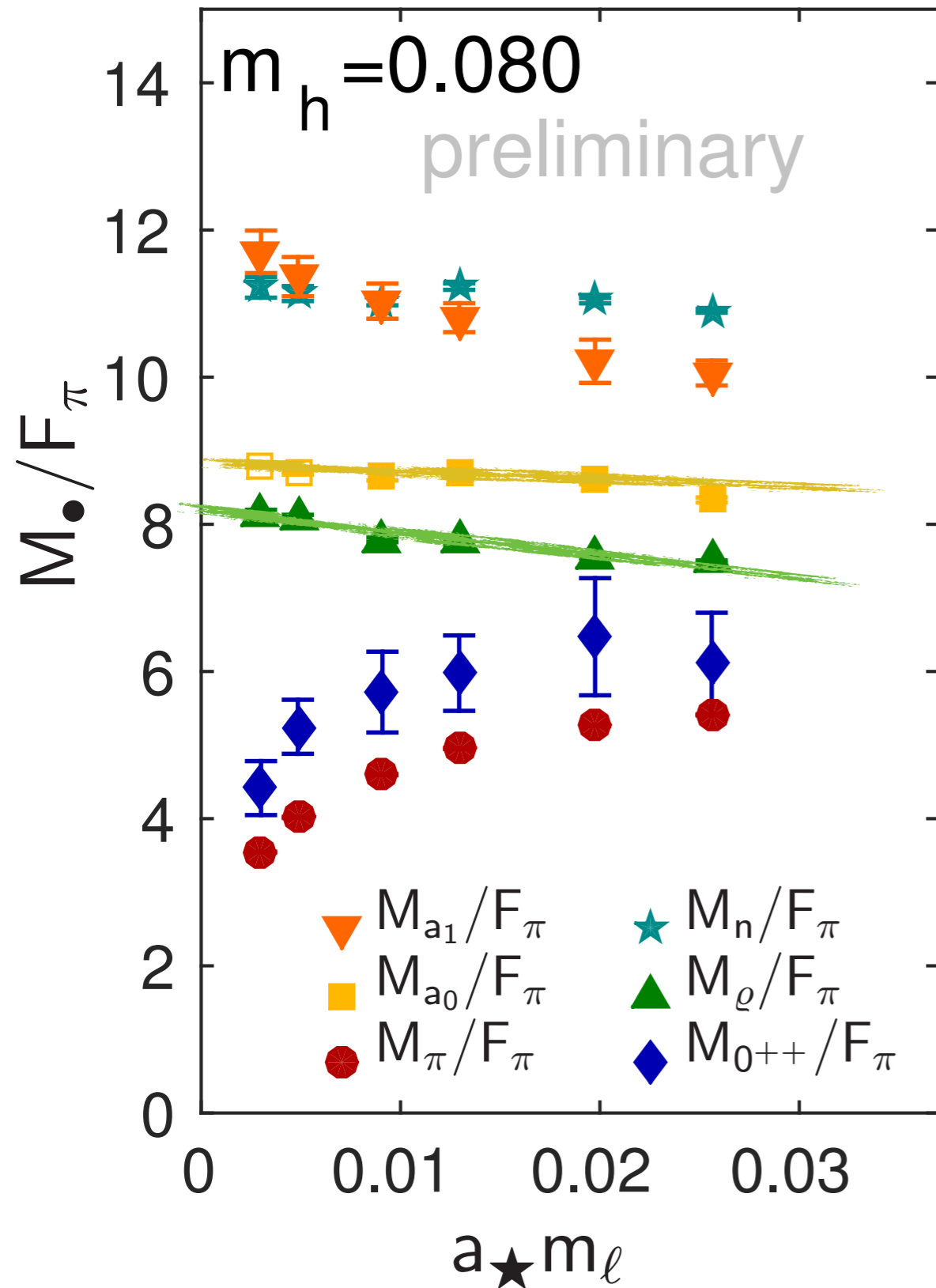
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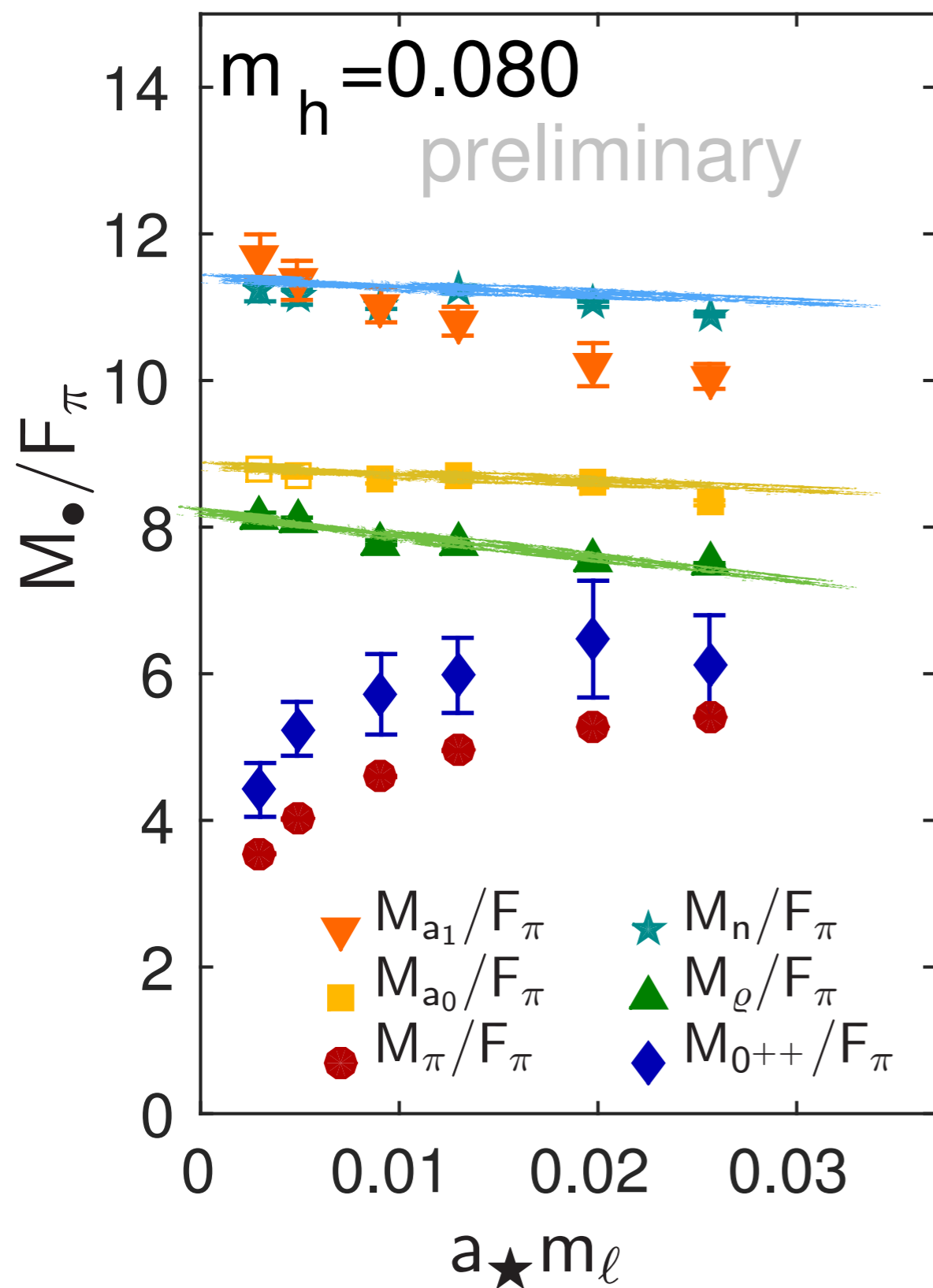
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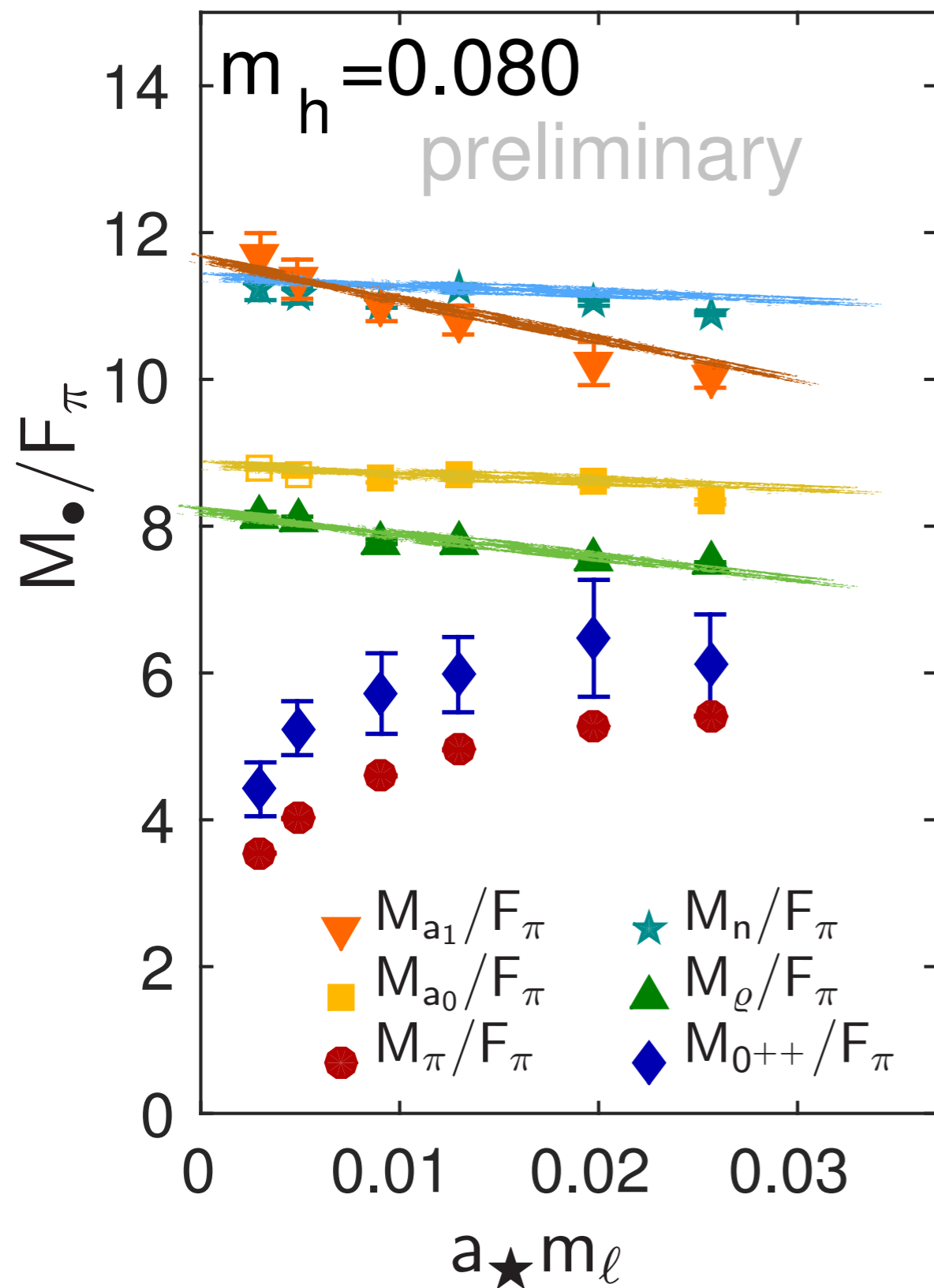
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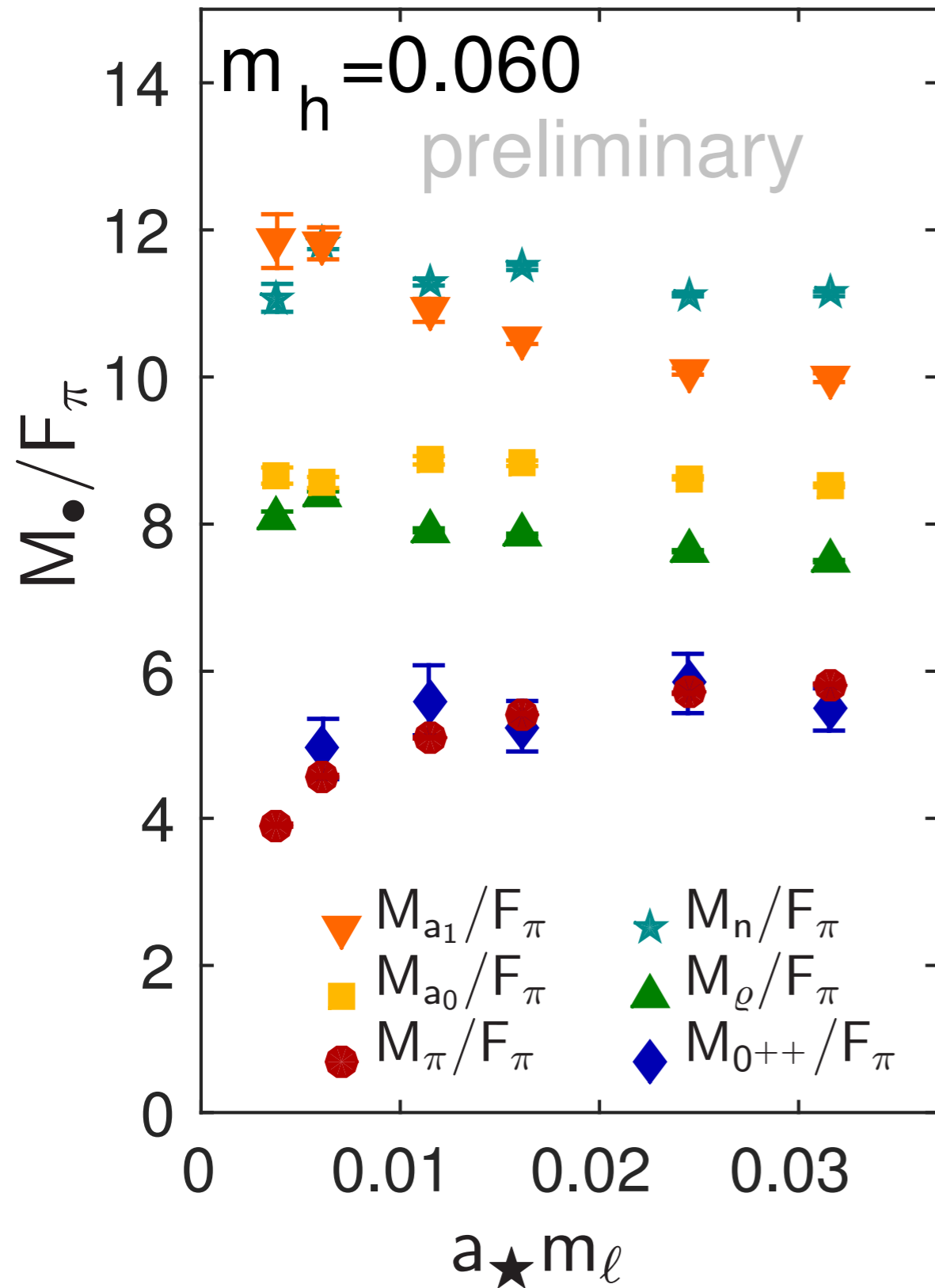
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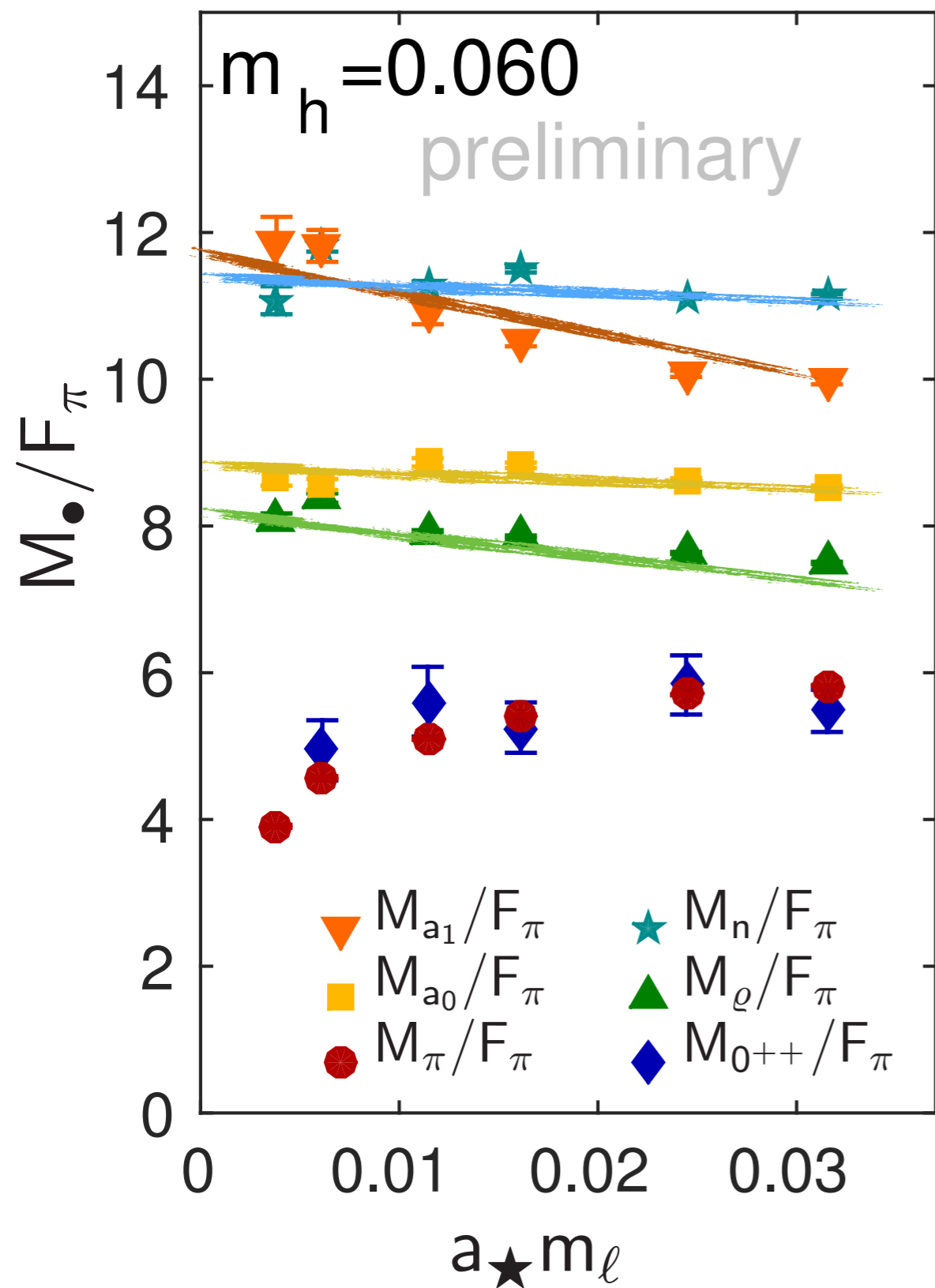
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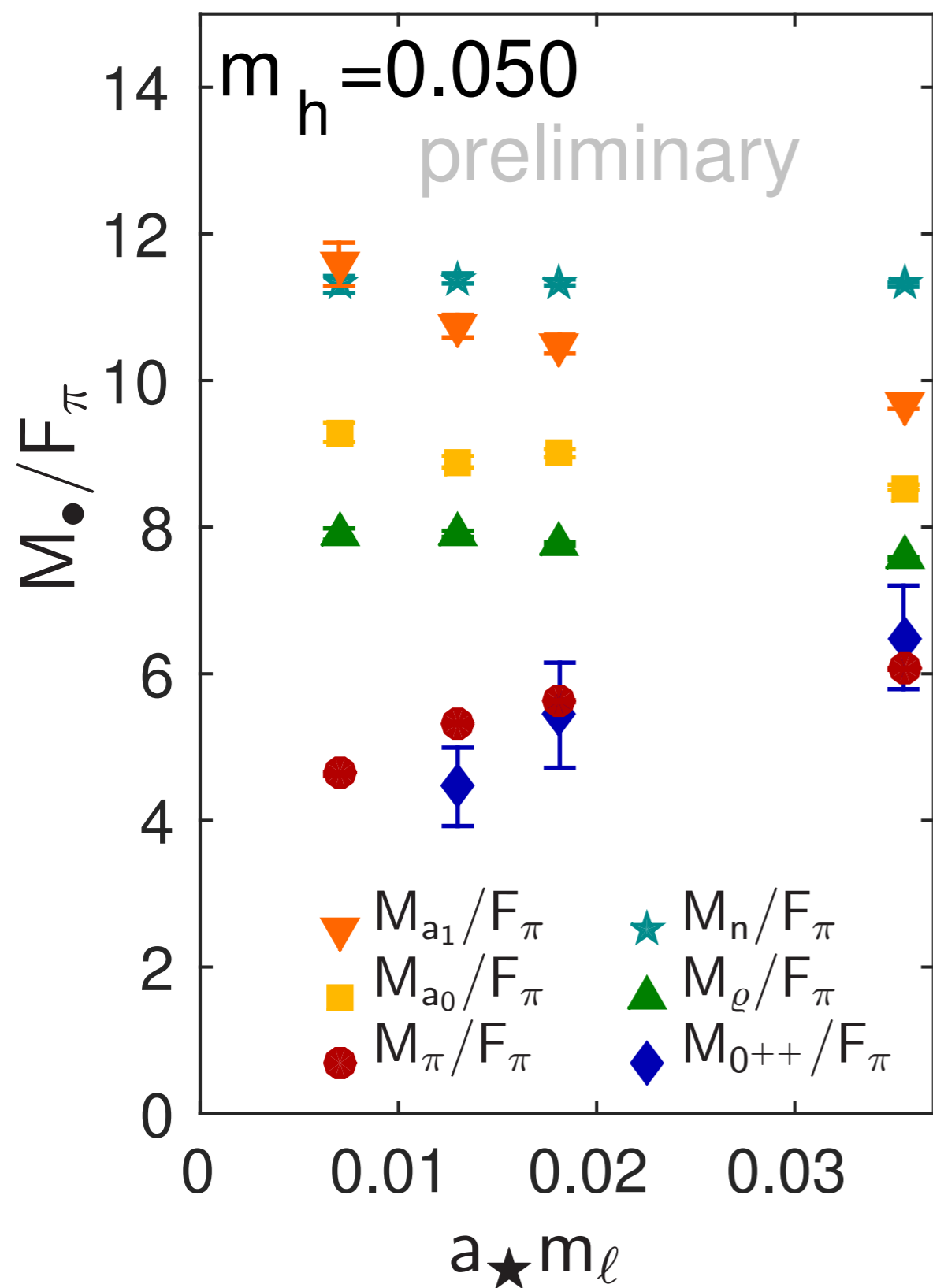
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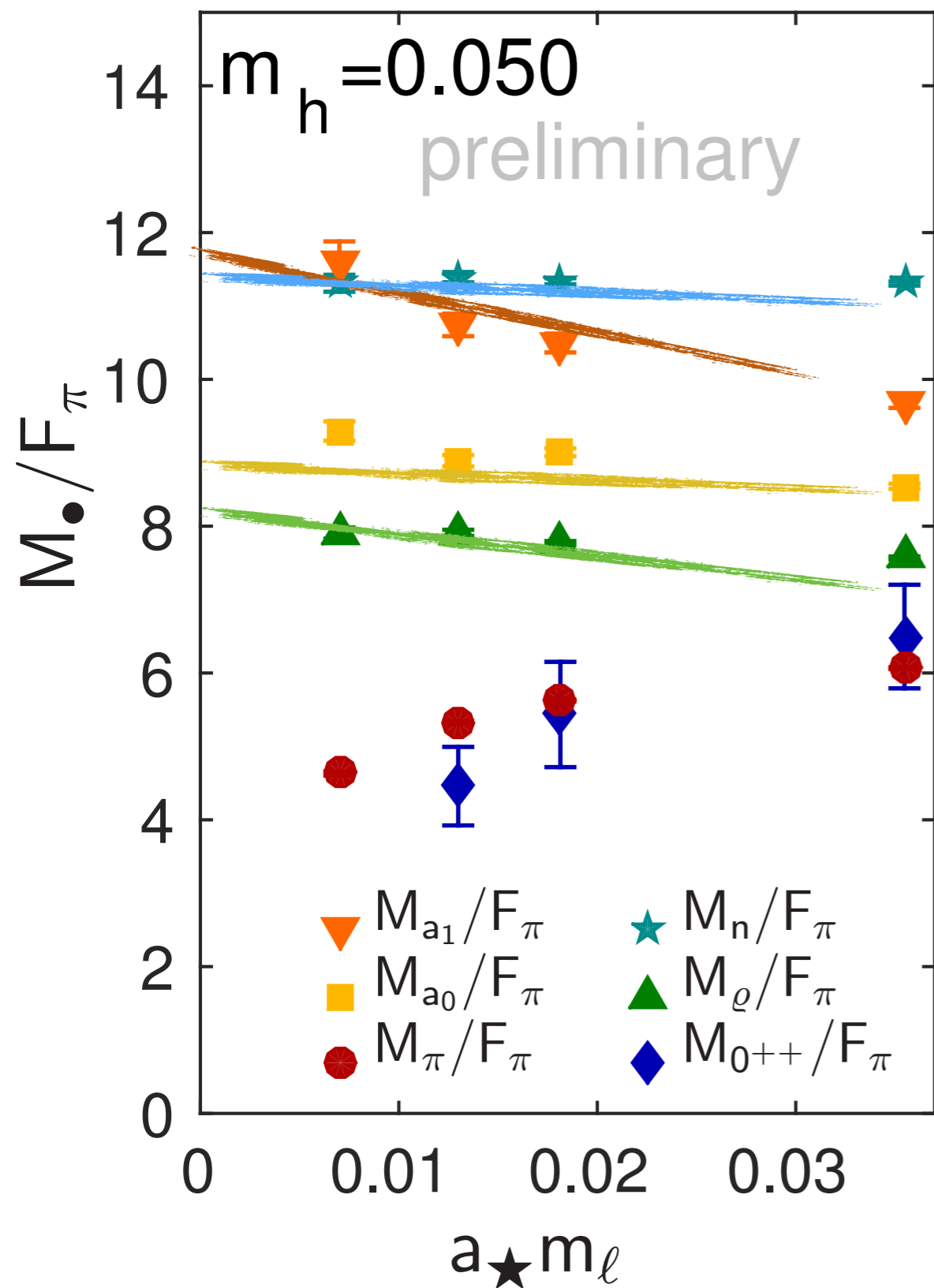
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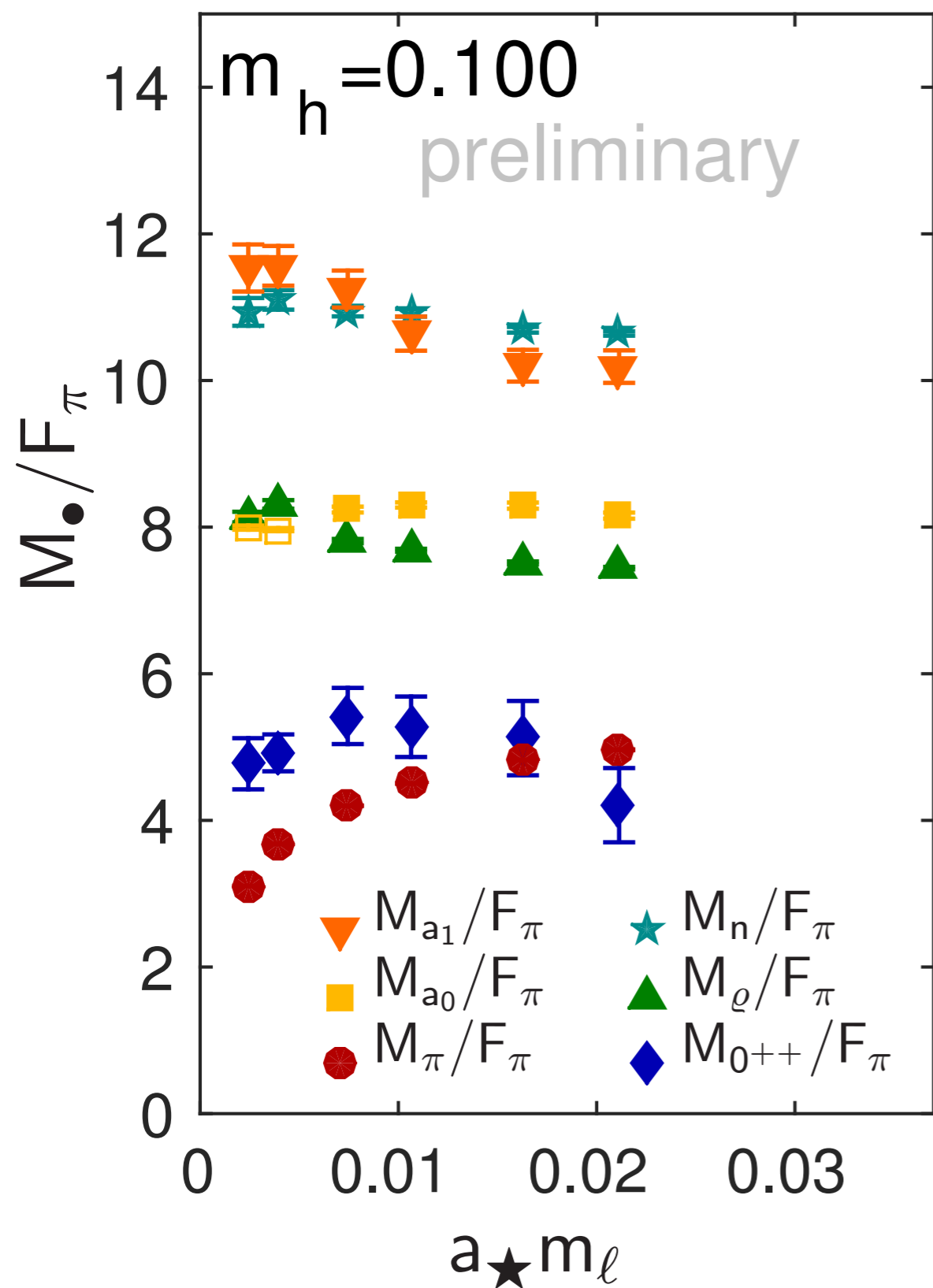
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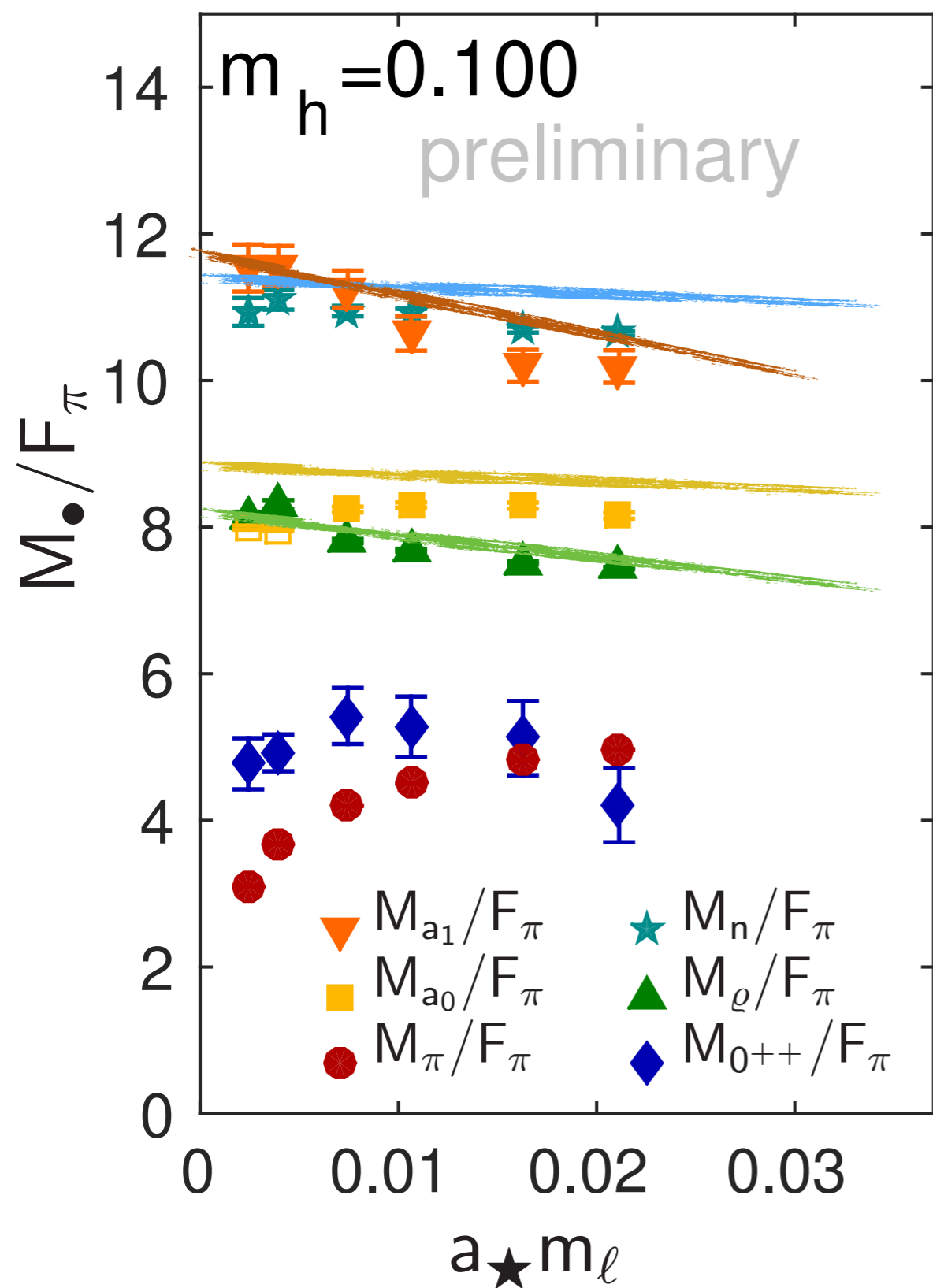
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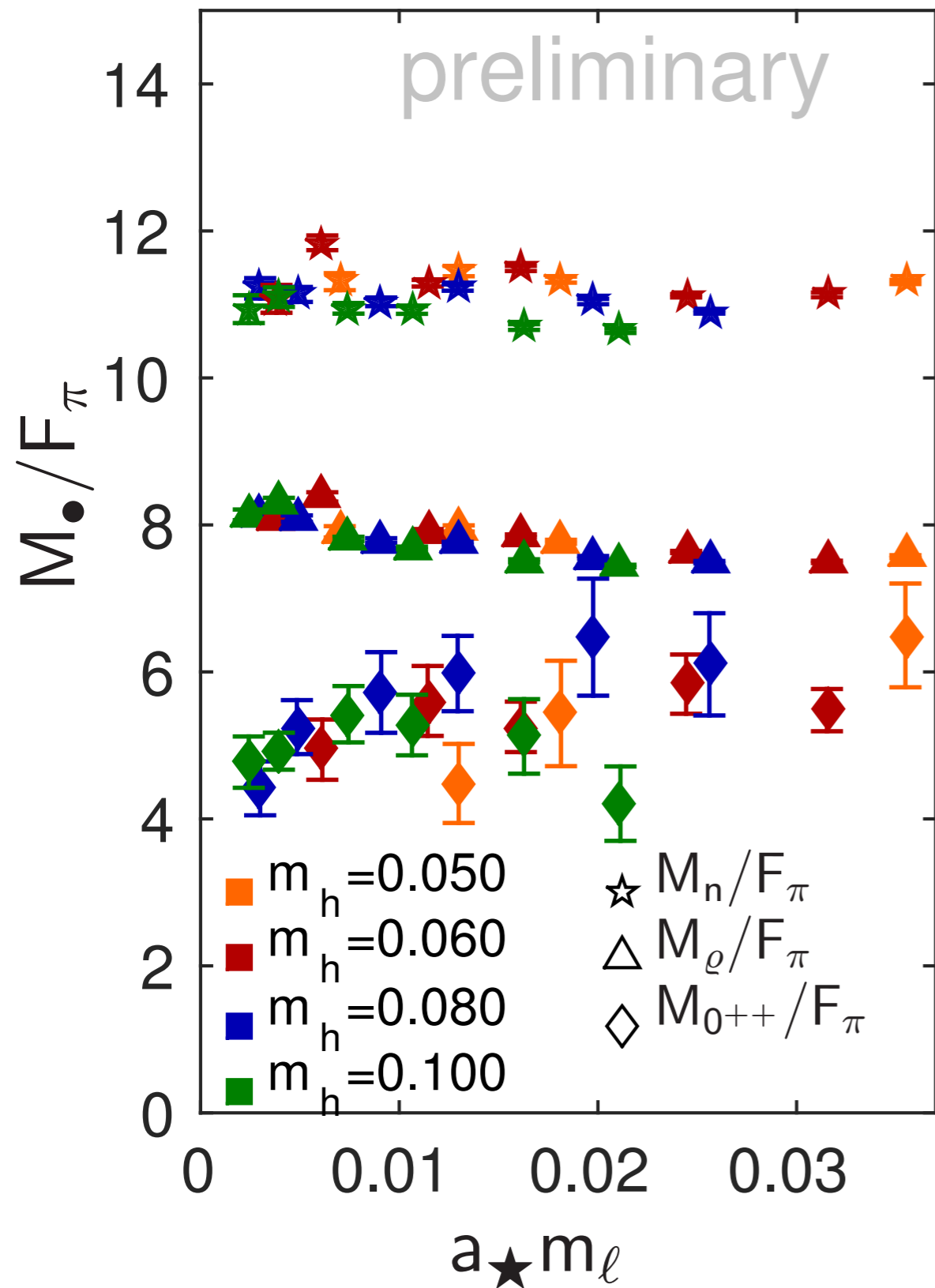
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Spectrum, 4+8 flavors: Put it all together



rho, nucleon
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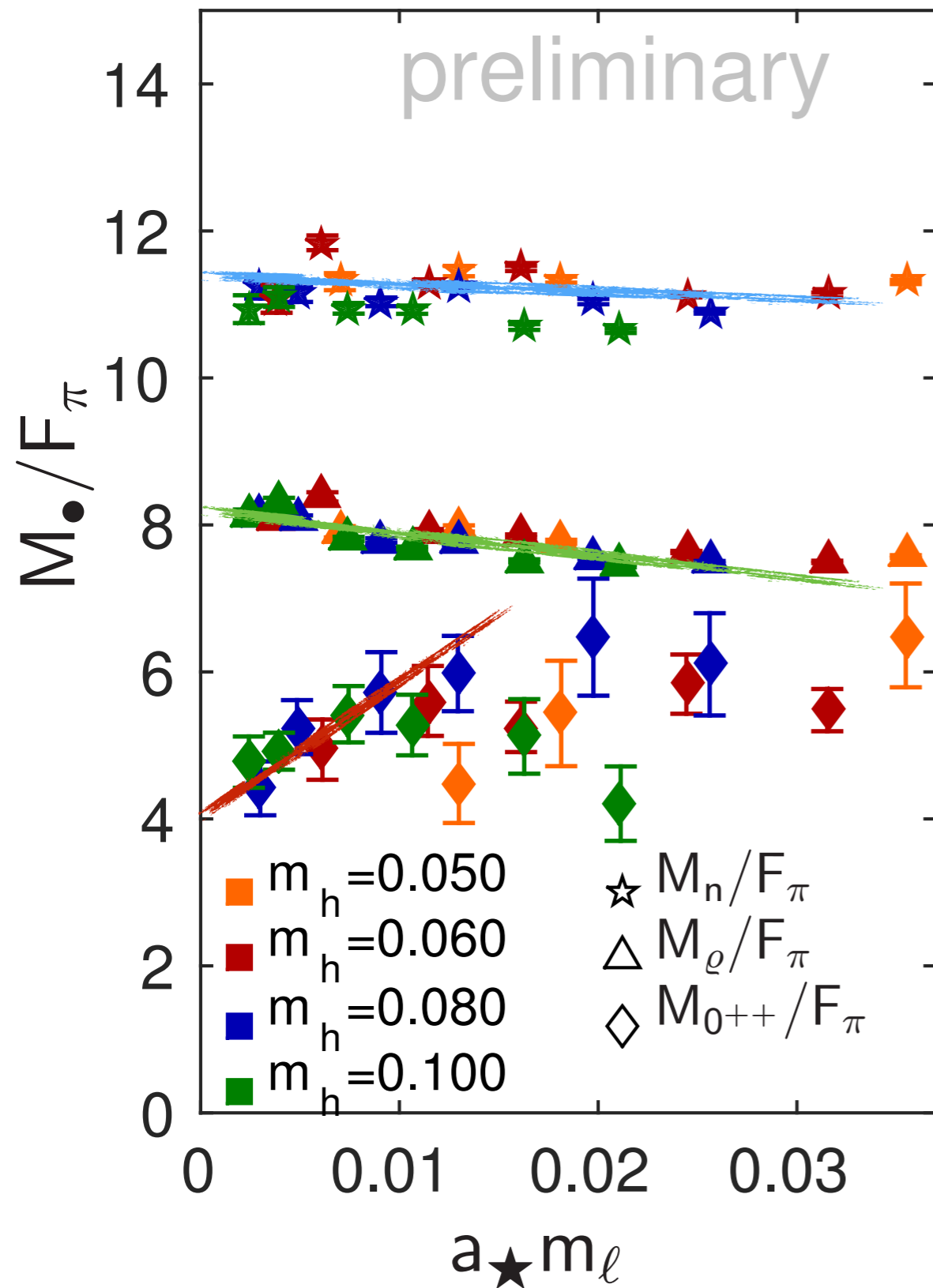
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$$M_N / F_{\pi} \approx 11$$

0^{++} : chiral limit is difficult but
well separated from the rho

Spectrum, 4+8 flavors: Put it all together



rho, nucleon
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$$M_{\rho} / F_{\pi} \approx 8$$

$$M_N / F_{\pi} \approx 11$$

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Mass split models

The 4+8 system is not ideal:

- we need 2 light flavors
- $N_f = 12$ is far above the conformal window with small anomalous dimension $\gamma_m \approx 0.25$

Yet it shows:

- hyperscaling,
- similar general properties as $N_f = 8$
- 0^{++} well separated from heavier excitations

How will the spectrum change if we change N_f or cascade the mass?

Conclusion & Summary

Lots of interesting possibilities

Lattice studies are needed to investigate strongly coupled systems
- individual and generic properties

There appear to be some very general features between
different models: $M_\rho / F_\pi \approx 8$
with other nearby resonances

Models with split fermion masses, built on a conformal IRFP,
offer new approach - yet still show similar general features

LHC Run-2 could verify / falsify many of the BSM models

EXTRA SLIDES

If the dilaton is not naturally light:

The heavy flavors can be decoupled sequentially, making the beta function walk across a wide coupling range

