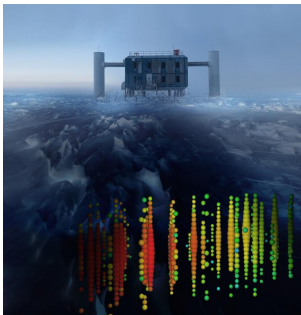


Neutrinos

Jenni Adams

University of Canterbury
New Zealand

PreSusy School 2016



What's so interesting about neutrinos?

Neutrinos are a key to understanding a range of physics.

- Neutrino sector expected to give clues to beyond the standard model physics and grand unification theories
- Neutrino nature is related to lepton number violation, which may be important for generating the matter/antimatter asymmetry in the early Universe
- There is a cosmic neutrino background (like the CMB) of $336 \nu/\text{cm}^3$ which affects the Universe's evolution and large scale structure formation
- Neutrinos are a unique cosmic messenger, able to escape from dense regions and unaffected by magnetic fields
- ...

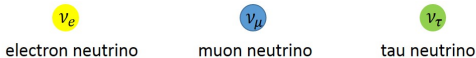
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Introducing neutrinos...

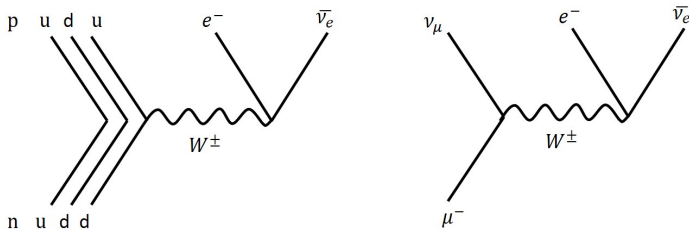
- Three neutrino flavours corresponding to the three charged leptons



- And three flavours of anti-neutrinos

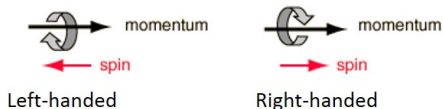


- Neutrino flavour and antiparticle/particle distinction determined by the interaction vertex and lepton flavour conservation

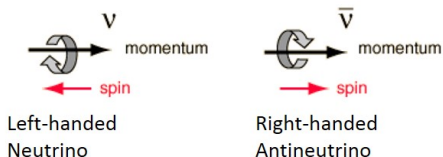


Introducing neutrinos... parity violation

- Parity refers to the relative orientation of the spin and momentum vectors of a particle

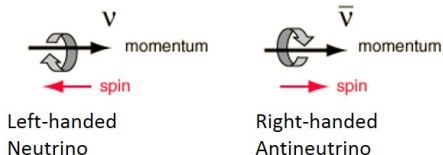


- All neutrinos are left-handed and all anti-neutrinos are right-handed

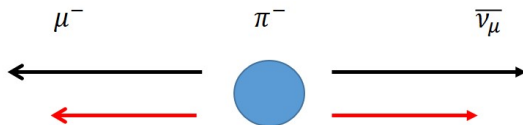


Introducing neutrinos... parity violation

- All neutrinos are left-handed and all anti-neutrinos are right-handed



- Pion decay - pion has spin 0, in rest frame of pion, μ and ν_{μ} emitted back to back eg



- Mirror process not allowed.

Introducing neutrinos... sterile neutrinos

- Three **active** neutrino flavours corresponding to the three charged leptons

 ν_e

electron neutrino

 ν_μ

muon neutrino

 ν_τ

tau neutrino

- It is possible that there are sterile neutrinos - neutrinos which do not interact by any force other than gravity
- Sterile neutrinos could oscillate with the active neutrinos - in that case the PMNS matrix would not be unitary (more later...)

Neutrino - Majorana or Dirac particle?

What is a Majorana particle?

- A Majorana particle is its own antiparticle, ie it is self charge-conjugate $\psi = \psi^c \equiv C\bar{\psi}^T$ where C is the charge conjugate matrix which has the property $C\gamma_\lambda C^{-1} = -\gamma^{\lambda T}$
- A Majorana neutrino is described by the Lagrangian

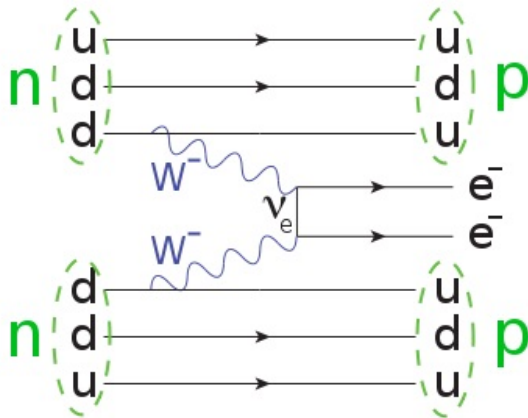
$$\mathcal{L}^M = \bar{\nu}_L i \not{\partial} \nu_L - \frac{m}{2} (\bar{\nu}_L^c \nu_L + \bar{\nu}_L \nu_L^c)$$

compare to the Dirac Lagrangian

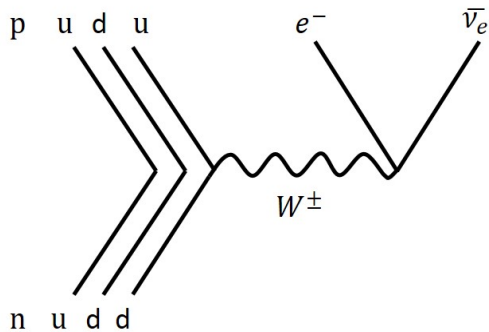
$$\mathcal{L}^D = \bar{\nu}_L i \not{\partial} \nu_L + \bar{\nu}_R i \not{\partial} \nu_R - m (\bar{\nu}_R \nu_L + \bar{\nu}_L \nu_R)$$

- Majorana particles imply lepton number violation

Majorana neutrinos \implies lepton violating
neutrinoless double beta decay $\beta\beta 0\nu$

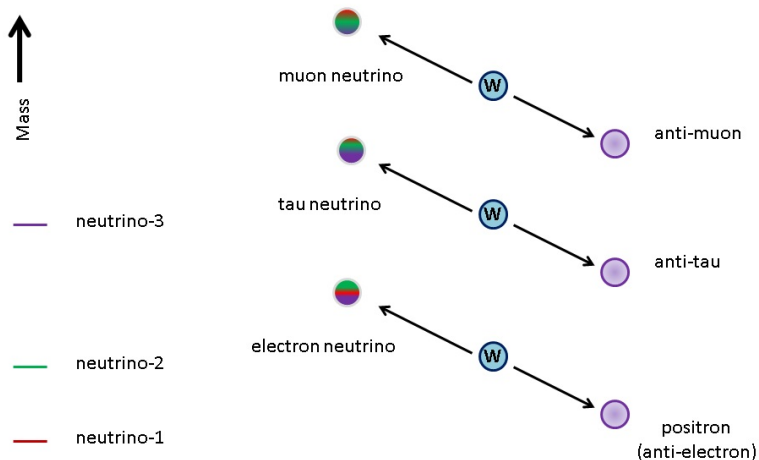


Compare beta decay



Neutrino masses

- Neutrino mass eigenstates are not neutrino flavour eigenstates



M. Strassler 2011

Neutrino masses

Pontecorvo-Maki-Nakagawa-Sakata Matrix relating flavour and mass eigenstates

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

ν_1  ν_2  ν_3 

Neutrino masses

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

This 3×3 unitary mixing matrix can be expressed in terms of four physical parameters conventionally chosen as **three mixing angles** $\theta_{12}, \theta_{23}, \theta_{13}$ (like 3 Euler angles describing rotation in 3D space) and **one phase** δ_{13}

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{13}} \\ 0 & 1 & 0 \\ -s_{13} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$c_{ab} \equiv \cos \theta_{ab} \quad s_{ab} \equiv \sin \theta_{ab} \quad 0 \leq \theta_{ab} \leq \frac{\pi}{2} \quad 0 \leq \delta_{13} \leq 2\pi$$

Neutrino masses - 2 flavour oscillations

$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

Amplitude for oscillation from flavour state α to state β

$$A(\nu_\alpha \rightarrow \nu_\beta) = \sum_i [A(\text{neutrino born flavour } \alpha \text{ is a } \nu_i) \\ \times A(\nu_i \text{ propagates}) \times A(\text{when } \nu_i \text{ interacts it makes flavour } \beta)]$$

In terms of time t and position L each mass eigenstate propagates as $e^{-i(E_i t - p_i L)}$ and where $t \cong L$ it becomes $e^{-i(E_i - p_i)L}$ where

$$p_i = \sqrt{E^2 - m_i^2} \cong E - m_i^2/2E \text{ so}$$

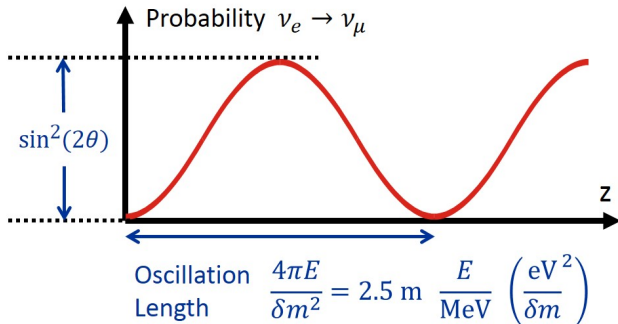
$$A(\nu_i \text{ propagates}) = e^{-i(m_i^2/2p)L}$$

and

$$P(\nu_\alpha \rightarrow \nu_\beta) = |A(\nu_\alpha \rightarrow \nu_\beta)|^2 = \sin^2(2\theta) \sin^2\left(\frac{\Delta m_{12}^2 L}{4E}\right)$$

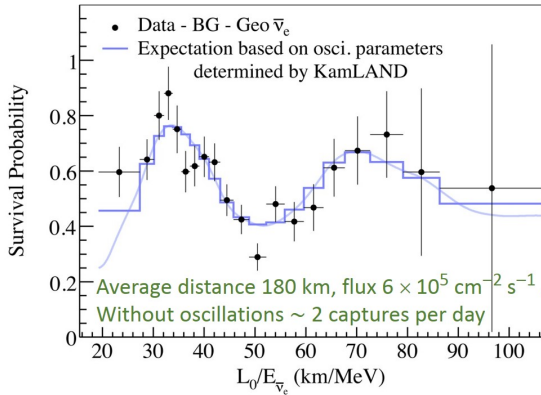
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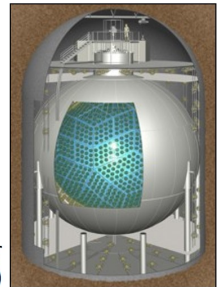


Oscillation of reactor neutrinos at KamLAND

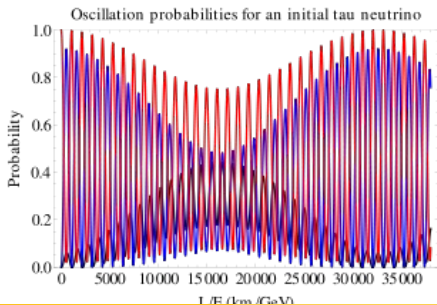
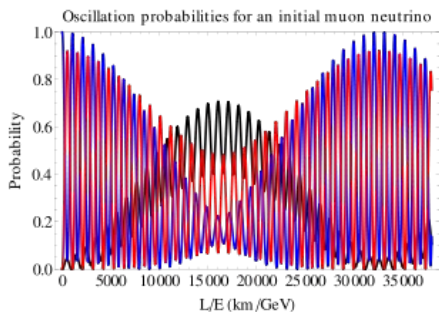
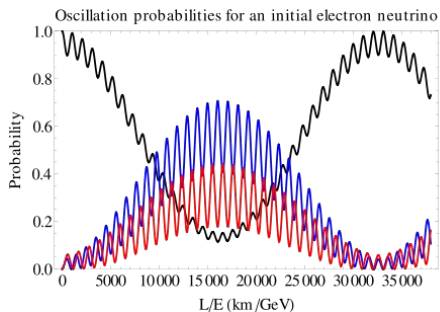
Oscillation pattern for anti-electron neutrinos from Japanese power reactors as a function of L/E



KamLAND Scintillator detector (1000 t)



Neutrino masses - 3 flavour oscillations

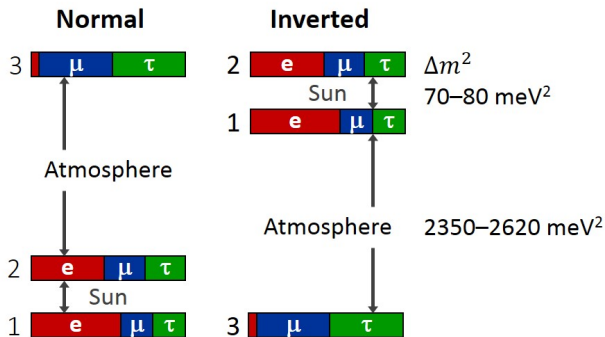


Neutrino masses - 3 flavour oscillations

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\substack{39^\circ < \theta_{23} < 53^\circ \\ \text{Atmospheric/LBL-Beams}}} \underbrace{\begin{pmatrix} c_{13} & 0 & e^{-i\delta} s_{13} \\ 0 & 1 & 0 \\ -e^{i\delta} s_{13} & 0 & c_{13} \end{pmatrix}}_{\substack{7^\circ < \theta_{13} < 9^\circ \\ \text{Reactor}}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\substack{31^\circ < \theta_{12} < 37^\circ \\ \text{Solar/KamLAND}}} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

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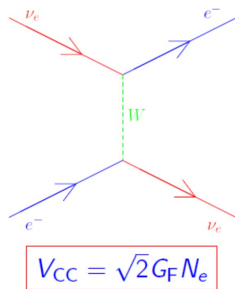


Ongoing

- Precision for all angles
- CP-violating phase δ ?
- Mass ordering? (normal vs inverted)

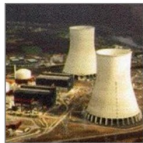
Oscillations in matter - MSW effect

- The Mikheyev Smirnov Wolfenstein effect or matter effect describes a change in oscillation parameters for neutrinos propagating in matter
- Charged current coherent forward scattering of the electron neutrinos off electrons in matter changes the neutrino energy levels and oscillation amplitudes



- Propagation amplitude determined by Hamiltonian which is modified
 $H = H_0 + H_I$

Sources of neutrinos



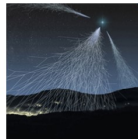
Reactors



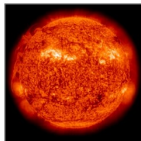
Particle accelerators



Geo-neutrinos



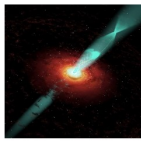
Atmospheric neutrinos –
interactions of cosmic rays in
the Earth's atmosphere



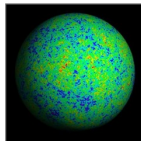
Sun



Supernova

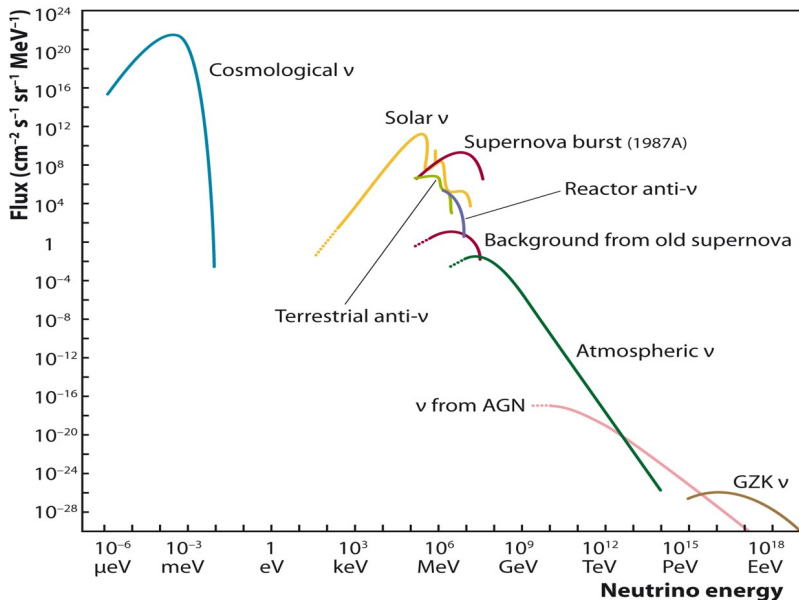


Astrophysical
accelerators

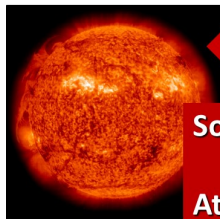


Cosmic neutrinos –
prediction $337 \nu/\text{cm}^3$

Sources of neutrinos



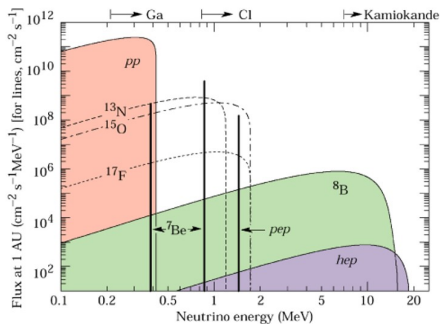
Solar neutrinos



8.3 light minutes

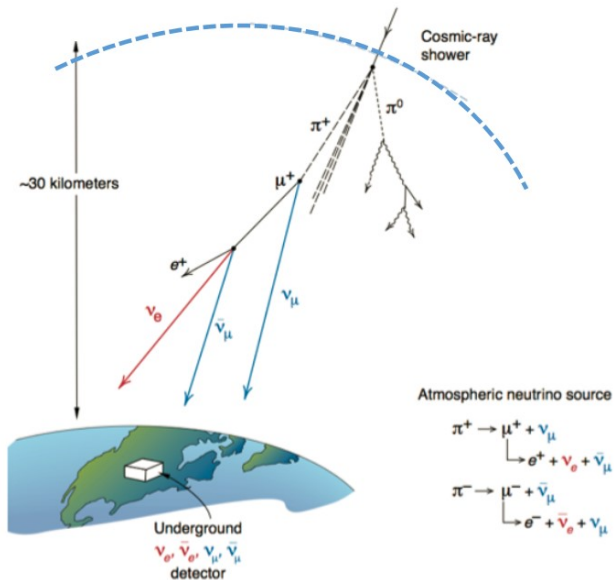


Solar radiation: 98 % light (photons)
2 % neutrinos
At Earth 66 billion neutrinos/cm² sec



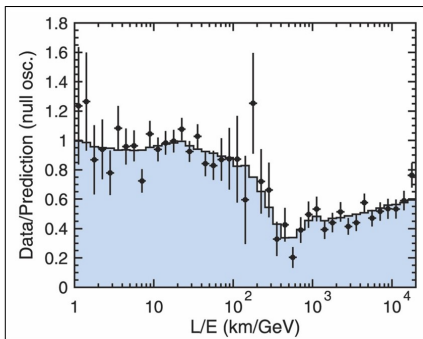
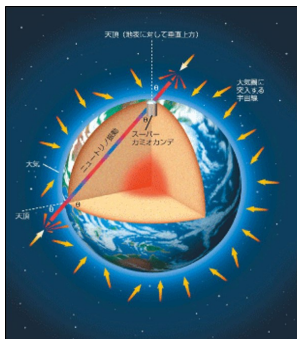
Several light years of lead
needed to shield solar
neutrinos

Atmospheric neutrinos



Atmospheric neutrino oscillations

Observation by SuperKamiokande - 2008



Atmospheric neutrino oscillations show characteristic L/E variation