

# Experimental SUSY searches at the LHC

## Lecture 1

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# Bibliography

- [www.physik.uzh.ch/lectures/empp/15/lectures/empp15\\_LC\\_Trigger.pdf](http://www.physik.uzh.ch/lectures/empp/15/lectures/empp15_LC_Trigger.pdf)
- [https://www2.warwick.ac.uk/fac/sci/physics/staff/.../chill\\_warwick\\_lhc\\_2010\\_2.pdf](https://www2.warwick.ac.uk/fac/sci/physics/staff/.../chill_warwick_lhc_2010_2.pdf)
- <https://www-conf.slac.stanford.edu/ssi/2012/Presentations/Pralavorio-1.pdf>
- [ias.ust.hk/program/shared\\_doc/2016/.../20160118\\_Xuai\\_Zhuang\\_2042.pdf](http://ias.ust.hk/program/shared_doc/2016/.../20160118_Xuai_Zhuang_2042.pdf)
- [www.desy.de/~csander/Talks/131001\\_SUSYSummary.pdf](http://www.desy.de/~csander/Talks/131001_SUSYSummary.pdf)
- [www.hep.ph.imperial.ac.uk/~tapper/lecture/susy-lecture-1.pdf](http://www.hep.ph.imperial.ac.uk/~tapper/lecture/susy-lecture-1.pdf)
- [www.hep.ph.imperial.ac.uk/~tapper/lecture/susy-lecture-2.pdf](http://www.hep.ph.imperial.ac.uk/~tapper/lecture/susy-lecture-2.pdf)
- [moriond.in2p3.fr/QCD/2014/WednesdayMorning/Sekmen.pdf](http://moriond.in2p3.fr/QCD/2014/WednesdayMorning/Sekmen.pdf)
- <https://indico.in2p3.fr/event/12279/session/12/contribution/194/material/slides/>

# Where we're headed

Short Title of Paper	Date	$\sqrt{s}$ (TeV)	L (fb <sup>-1</sup> )	Document	Plots+Aux. Material	Journal
Stop to bs (RPV)	1/2016	8	17.4	<a href="#">1601.07453</a>	<a href="#">Link</a>	Accepted by JHEP
Stop to b $\tau$ au+gravitino	9/2015	8	20	<a href="#">1509.04976</a>	<a href="#">Link (+data)</a>	EPJ C (2016) 76:81
1/2 photon(s) + jets / b-jets / lepton + Etmis [GGM, Strong/EW production]	7/2015	8	20.3	<a href="#">1507.05493</a>	<a href="#">Link (+data)</a>	Phys. Rev. D 92 (2015) 072001
LLP with pixel dE/dx [split-SUSY, AMSB]	6/2015	8	18.4	<a href="#">1506.05332</a>	<a href="#">Link (+data)</a>	Eur. Phys. J. C (2015) 75:407
Displaced vertex [RPV, GMSB, split-SUSY]	4/2015	8	20.3	<a href="#">1504.05162</a>	<a href="#">Link (+data)</a>	Phys. Rev. D 92, 072004 (2015)
Heavy resonance to e $\mu$ , e $\tau$ , $\mu\tau$ [RPV-LFV, Z-LFV]	3/2015	8	20.3	<a href="#">1503.04430</a>	<a href="#">Link (+data)</a>	Phys. Rev. Lett. 115, 031801 (2015)
2 leptons (on/off-Z) + jets + Etmis [incl. squarks & gluinos, GGM]	3/2015	8	20.3	<a href="#">1503.03290</a>	<a href="#">Link (+data)</a>	Eur. Phys. J. C75 (2015) 318
Multijets [RPV]	2/2015	8	20.3	<a href="#">1502.05686</a>	<a href="#">Link (+data)</a>	Phys. Rev. D 91, 112016 (2015)
Monojet + Etmis (gravitino-squark, gravitino-gluino)	2/2015	8	20.3	<a href="#">1502.01518</a>	<a href="#">Link (+data)</a>	Eur. Phys. J. C (2015) 75:299
Chargino neutralino decaying via Higgs	1/2015	8	20.3	<a href="#">1501.07110</a>	<a href="#">Link (+data)</a>	Eur. Phys. J. C (2015) 75:208
1/2 leptons + jets + Etmis [incl. squarks & gluinos, mUED]	1/2015	8	20.3	<a href="#">1501.03555</a>	<a href="#">Link (+data)</a>	JHEP 04 (2015) 116
Search for scalar charm	1/2015	8	20.3	<a href="#">1501.01325</a>	<a href="#">Link (+data)</a>	Phys. Rev. Lett. 114, 161801 (2015)
Top spin correlations (stealth stop)	12/2014	8	20.3	<a href="#">1412.4742</a>	<a href="#">Link</a>	Phys. Rev. Lett. 114, 142001 (2015)

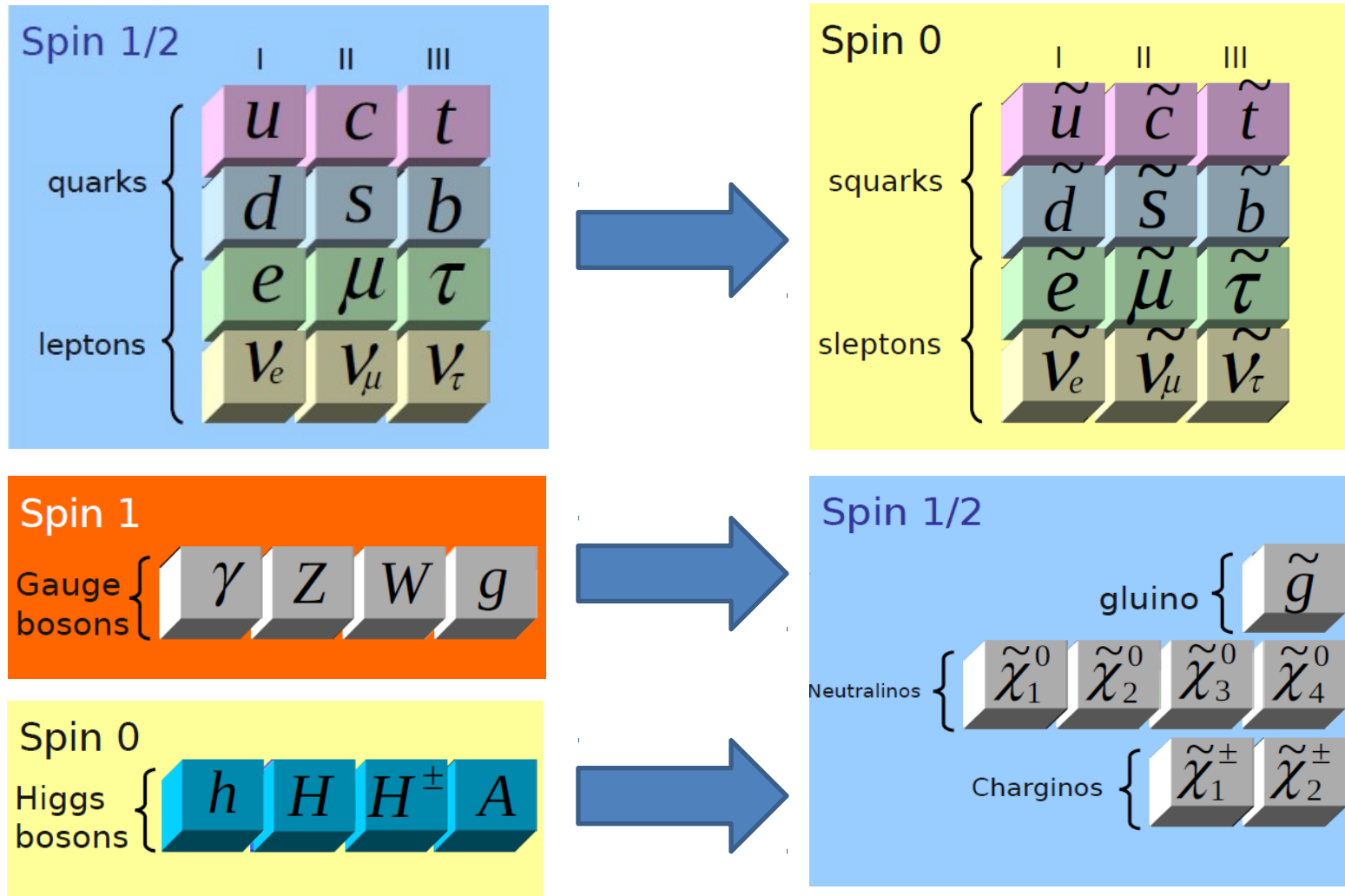
Short Title of Paper	Recent: Supersymmetry Preliminary Results	Date
Long-lived particles (sleptons, charginos, R-hadrons)	<a href="#">CMS-PAS-SUS-14-030</a>	June 2016
Mono-photon + Etmis [Degenerate squark/neutralino]	<a href="#">CMS-SUS-14-008</a>	29 May 2016
Non-pointing, delayed photons [LLP, GMSB]	<a href="#">CMS-PAS-SUS-13-012</a>	April 2016
0 leptons + mono-jet/c-jets + Etmis [Stop in charm+LSP]	<a href="#">CMS-PAS-SUS-13-007</a>	March 2016
1 lepton + 4(1 b-)jets + Etmis [Medium / heavy stop]	<a href="#">CMS-PAS-SUS-16-001</a>	March 2016
1-2 taus + 0-1 leptons + jets + Etmis [GMSB]	<a href="#">CMS-PAS-SUS-13-006</a>	March 2016
1-2 taus + 0-1 leptons + jets + Etmis [GMSB]	<a href="#">CMS-PAS-SUS-13-003</a>	March 2016
0-1 leptons + >=3 b-jets + Etmis [3rd gen. squarks]	<a href="#">CMS-PAS-SUS-16-004</a>	March 2016
2 taus + Etmis [EW production]	<a href="#">CMS-PAS-SUS-14-022</a>	March 2016
Stop constraints from precise t $\bar{b}$ cross-section [Light stop]	<a href="#">CMS-SUS-14-018</a>	29 February 2016
0 lepton + 6 (2 b-)jets + Etmis [Heavy stop]	<a href="#">CMS-SUS-14-013</a>	10 February 2016
0 leptons + 2-6 jets + Etmis [Incl. squarks & gluinos]	<a href="#">CMS-SUS-14-007</a>	9 February 2016
4 leptons + 2-6 jets + Etmis [Incl. squarks & gluinos]	<a href="#">CMS-PAS-SUS-13-002</a>	December 2015
1 lepton + 4 jets + Etmis [EW production, RPV]	<a href="#">CMS-PAS-SUS-13-002</a>	December 2015
2 same-sign / 3 -leptons + 0-3 b-jets + Etmis [Incl. squarks & gluinos]	<a href="#">CMS-PAS-SUS-13-004</a>	December 2015
2 leptons (e, $\mu$ ) + Etmis [chargino/neutralino/slepton]	<a href="#">CMS-PAS-SUS-13-007</a>	December 2015
Z + b-jet + jets + Etmis [Stop in GMSB, stop2]	<a href="#">CMS-PAS-SUS-13-008</a>	December 2015
2 leptons + (b)jets + Etmis [stop]	<a href="#">CMS-PAS-SUS-13-011</a>	December 2015
3 leptons (e, $\mu$ ,tau) + Etmis [chargino/neutralino]	<a href="#">CMS-PAS-SUS-13-010</a>	October 2015
Long-lived stopped gluino or squark R-hadrons [Split-SUSY]	<a href="#">CMS-PAS-SUS-14-018</a>	October 2015
Disappearing track + jets + Etmis [Direct long-lived chargino]	<a href="#">CMS-PAS-SUS-14-003</a>	September 2015
0 leptons + 2 b-jets + Etmis [Sbottom/stop]	<a href="#">CMS-PAS-SUS-14-017</a>	August 2015
0 leptons + >=7-10 jets + Etmis [Incl. squarks & gluinos]	<a href="#">CMS-PAS-SUS-14-019</a>	August 2015
	<a href="#">CMS-PAS-SUS-13-022</a>	August 2015
	<a href="#">CMS-SUS-14-013</a>	5 August 2015
	<a href="#">CMS-PAS-SUS-14-021</a>	April 2015
	<a href="#">CMS-SUS-14-001</a>	27 March 2015
	<a href="#">CMS-PAS-SUS-14-005</a>	March 2015
	<a href="#">CMS-SUS-14-011</a>	21 February 2015
	<a href="#">CMS-PAS-SUS-14-012</a>	January 2015
	<a href="#">CMS-SUS-14-010</a>	13 December 2014
	<a href="#">CMS-SUS-14-009</a>	26 November 2014
	<a href="#">CMS-PAS-SUS-14-004</a>	2014
	<a href="#">CMS-PAS-SUS-13-019</a>	2014
	<a href="#">CMS-PAS-SUS-14-006</a>	2014
	<a href="#">CMS-PAS-SUS-14-011</a>	2014
	<a href="#">CMS-PAS-SUS-14-002</a>	2014
	<a href="#">CMS-SUS-12-003</a>	15 May 2014
	<a href="#">CMS-PAS-SUS-13-018</a>	2014
	<a href="#">CMS-PAS-SUS-14-000</a>	2014

# Very rough outline

- We won't go through 200 different null SUSY searches at high speed with no details
  - you can get that at conferences (including this one)
  - I can't be bothered
  - **neither can you**
- What we will do
  - start with the basics, understand what we can measure in the LHC detectors
  - think about SUSY production processes at the LHC
  - walk through one analysis *in detail* so you can see how analyses are designed
  - sketch the main features of other analyses
  - learn how to tell if your favourite model is excluded from LHC data
- What I'd like us to discuss and learn:
  - what is good about LHC searches?
  - what is bad about LHC searches?
  - is SUSY dead, dying, stable or in good health?

- Lecture 1: the basics
  - what do I mean by SUSY?
  - what happens in an LHC collision?
  - how do the LHC detectors work?
  - which objects can we use to hunt for SUSY?
  - what (roughly) does SUSY look like at the LHC?
- Lecture 2 (mostly for experimentalists)
  - how are the LHC SUSY analyses designed?
  - discriminating variables
  - background measurements
  - detailed worked example (0 lepton search)
  - summary of different approaches
- Lecture 3 (mostly for theorists)
  - what have we learned about SUSY from the LHC (and LEP, etc...)
  - how can a theorist constrain a new model of SUSY (or other BSM physics?)

# Supersymmetry

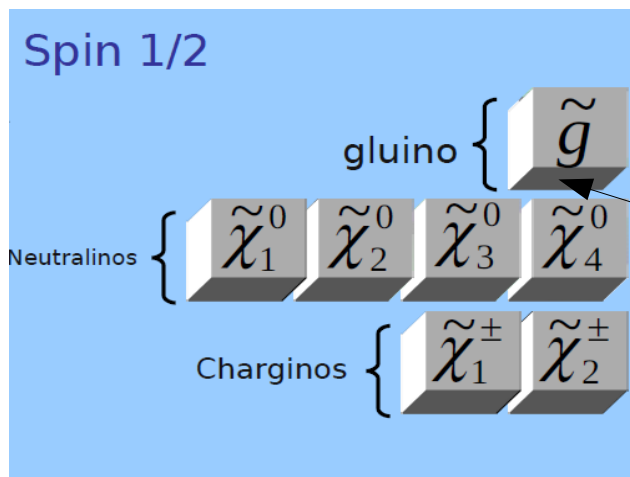


• “Supersymmetry” is a hypothesised symmetry connecting fermions and bosons

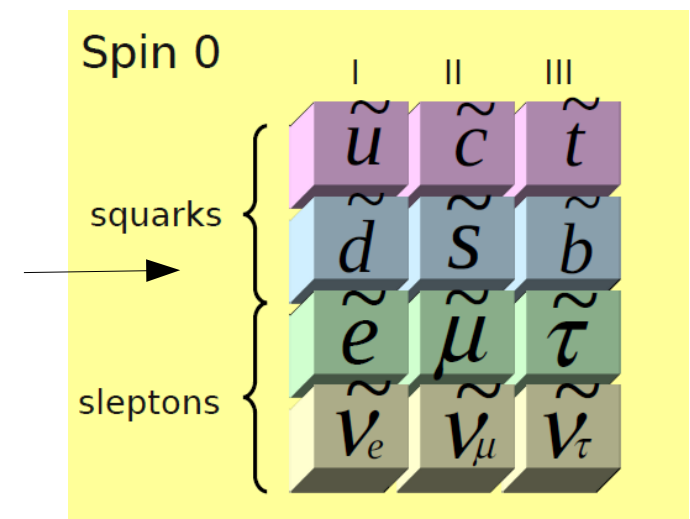
- SM particles get partners that cancel the corrections to the Higgs mass
- we also get viable dark matter candidates
- we also find that the SM forces unify at high energies

# What do I mean by SUSY?

- Exact SUSY would require identical SM & SUSY partner masses
  - SUSY is broken
- Breaking mechanism is *a priori* unknown
  - lots of GUT scale models (mSUGRA, GMSB, AMSB)
  - can use GUT or weak scale effective parameterisations (e.g. pMSSM)
  - lots of exotic SUSY models (e.g. NMSSM, E6SSM)
- For our purposes, SUSY means this:
  - any model that is supersymmetric and has some squarks, some sleptons, some extra Higgs bosons and some sort of EW gaugino sector



We want to try and measure these



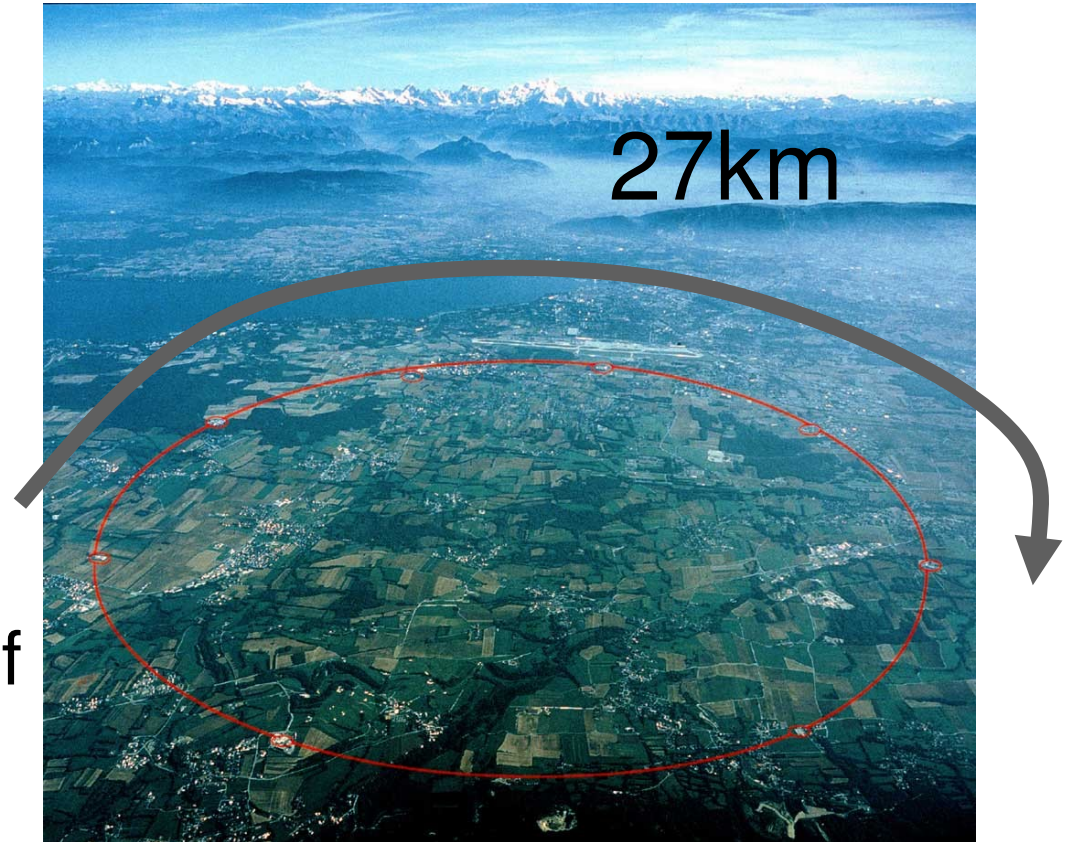
# Three sorts of paper on experimental searches

- Pre-LHC (up to 2007 or so)
  - how do we use the LHC to measure sparticle properties after discovery?
- Early LHC (2008 – 2011 or so)
  - how can we define new variables to better discover SUSY
- Now (2012+): Why the f\*\*k haven't we seen SUSY and what does it mean?
  - lots *and lots* of LHC exclusions
  - new baroque SUSY models that *are just around the corner*
  - why fine tuning doesn't matter, or why adding X@\$% relieves fine tuning

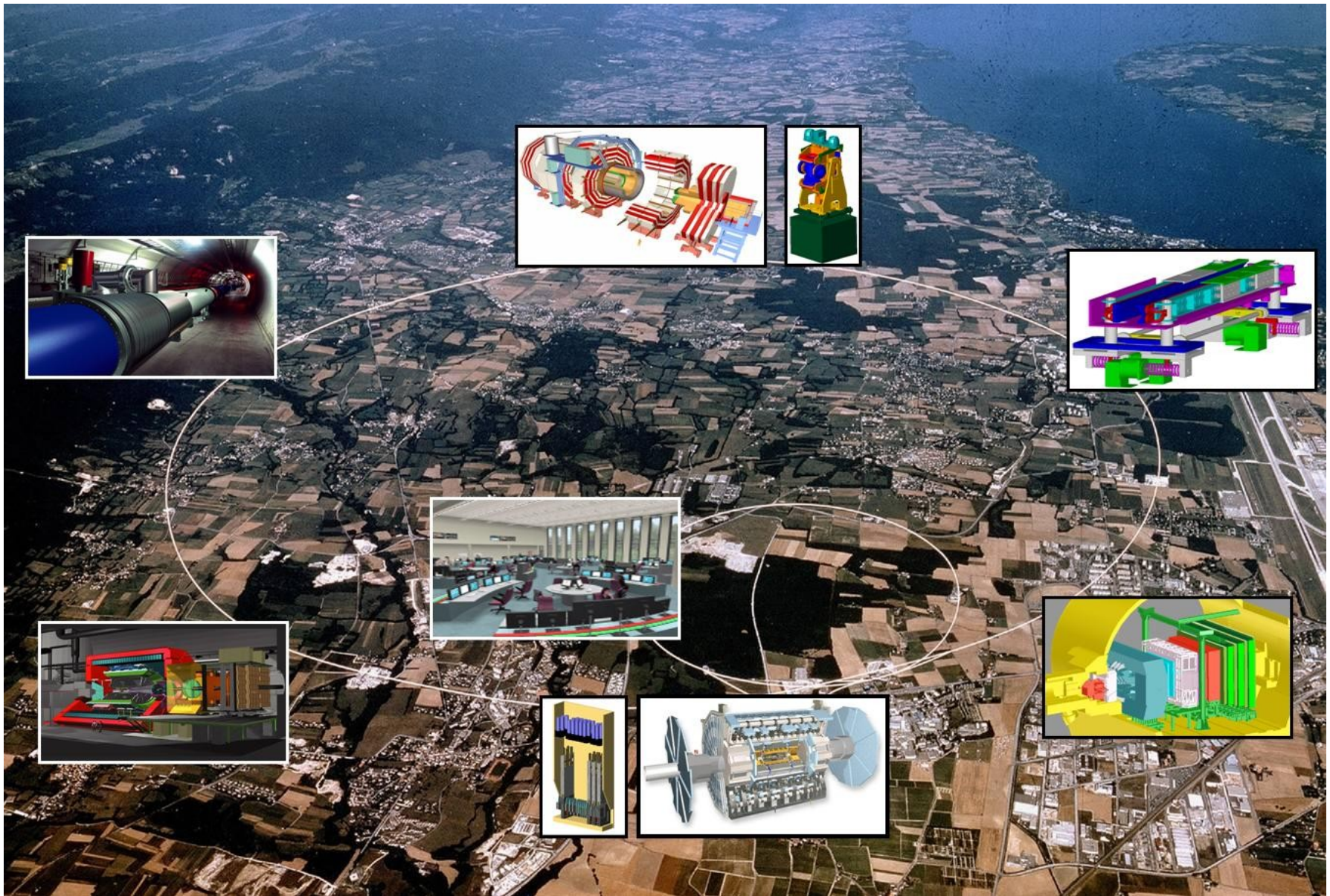
We won't spend much time on (1), but we will think a lot about (2) and (3)



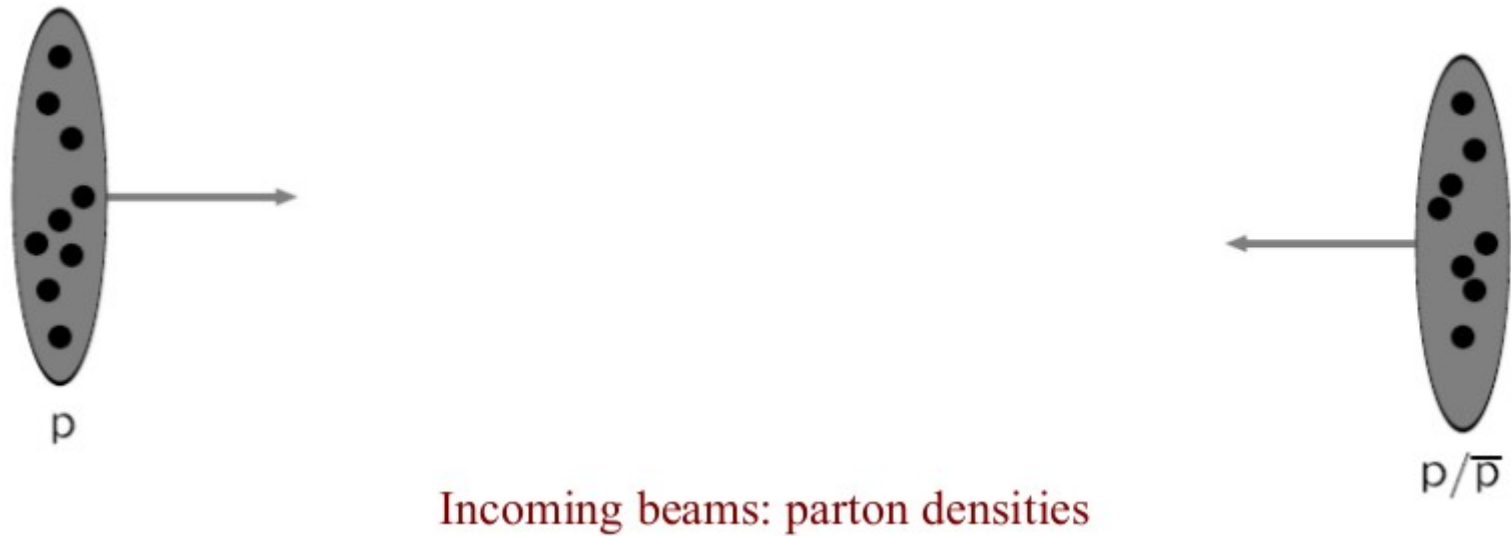
- Located at CERN
- Franco-Swiss border
- 27km circumference ring
- 100m underground
- Accelerates protons to *nearly* the speed of light
- Two counter rotating beams.



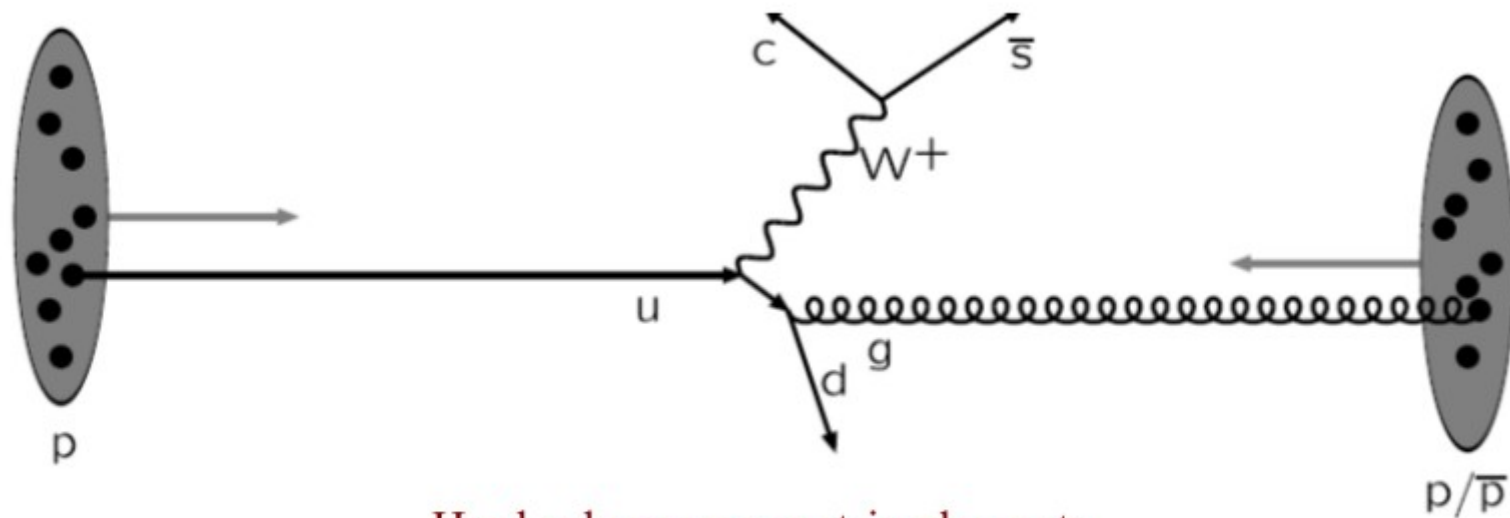
# The LHC



# What happens in a proton-proton collision

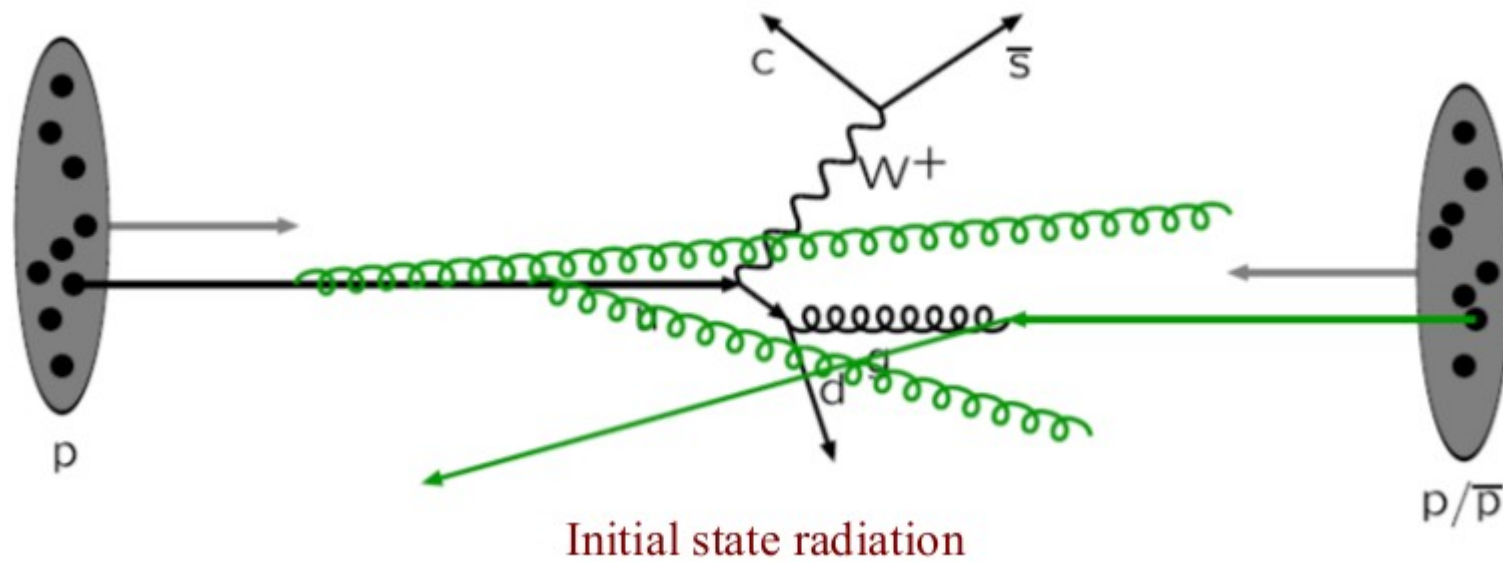


# What happens in a proton-proton collision

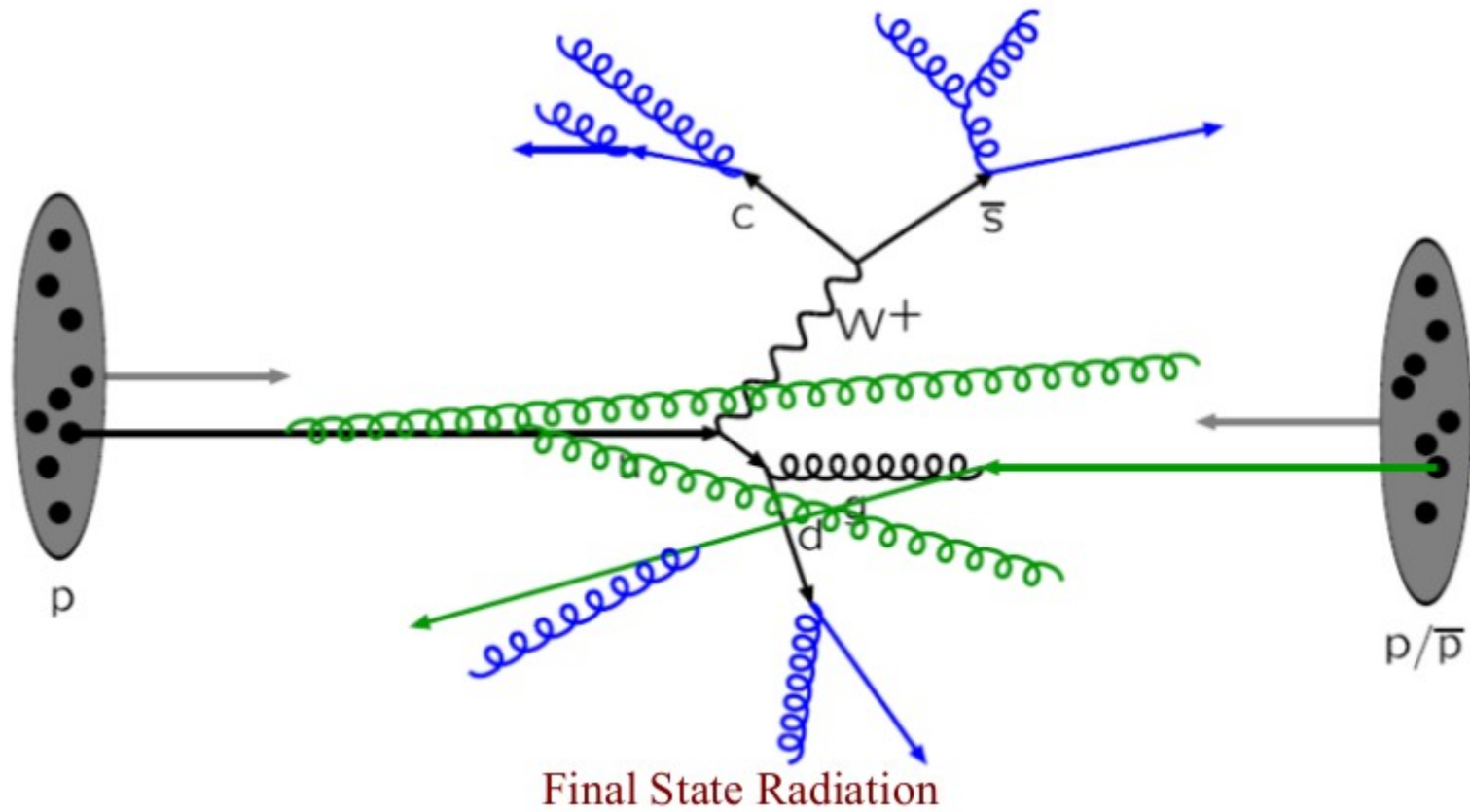


Hard subprocess: matrix elements,  
Resonance decays (correlated to hard process)

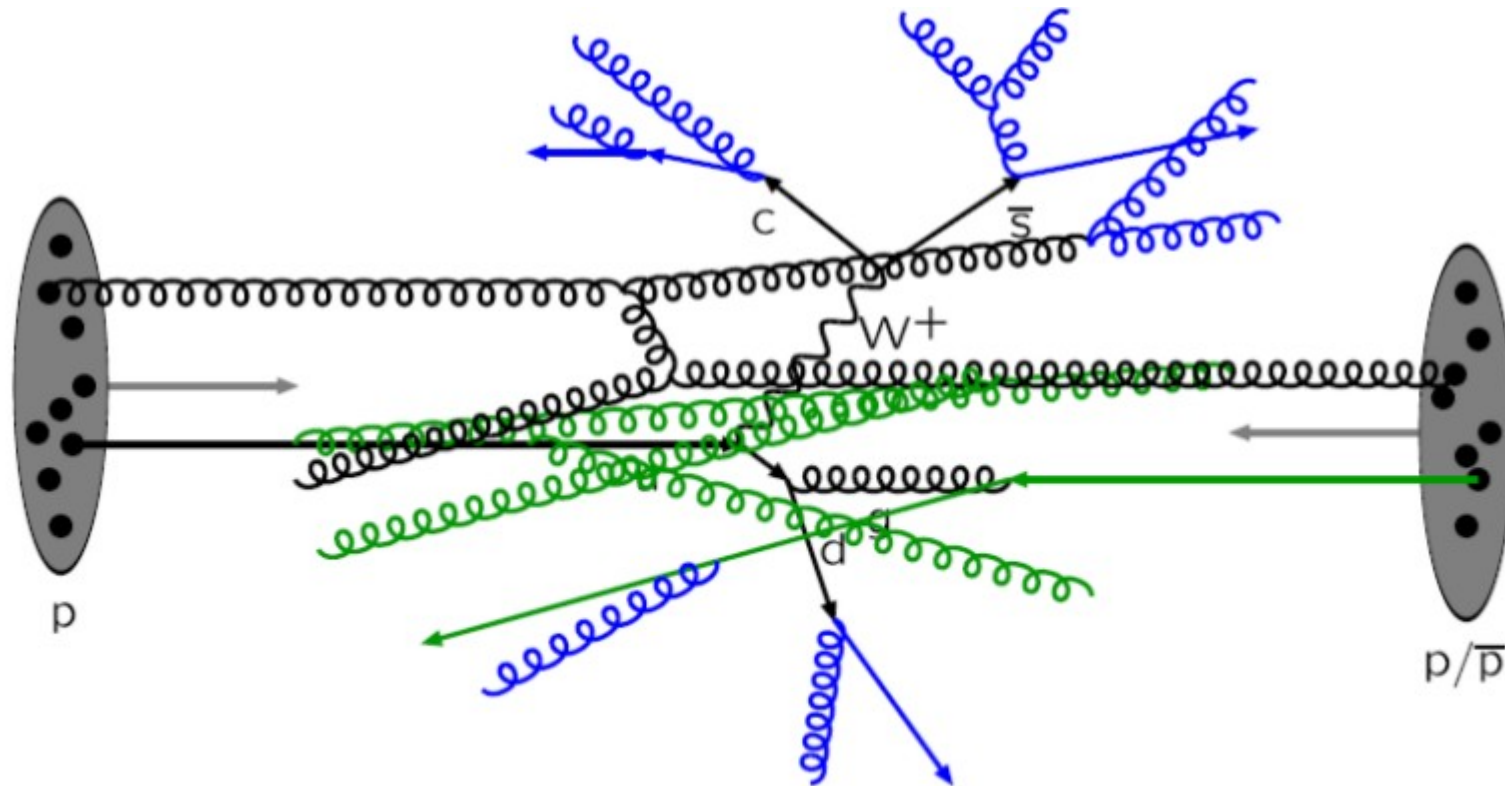
# What happens in a proton-proton collision



# What happens in a proton-proton collision

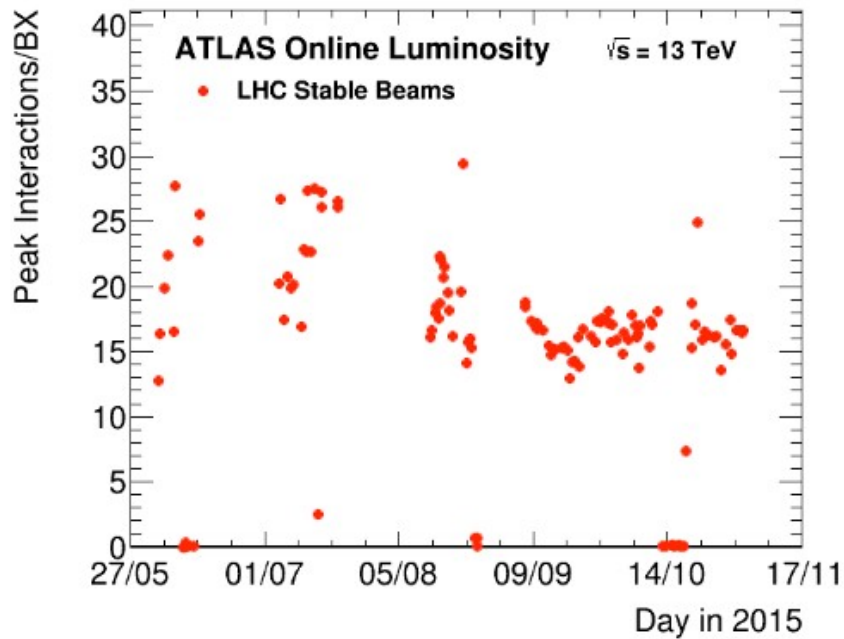
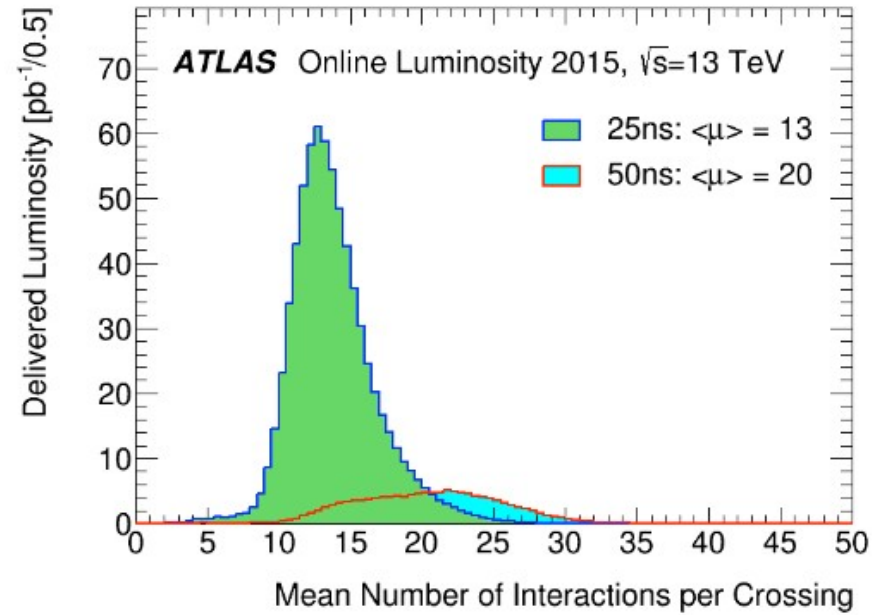
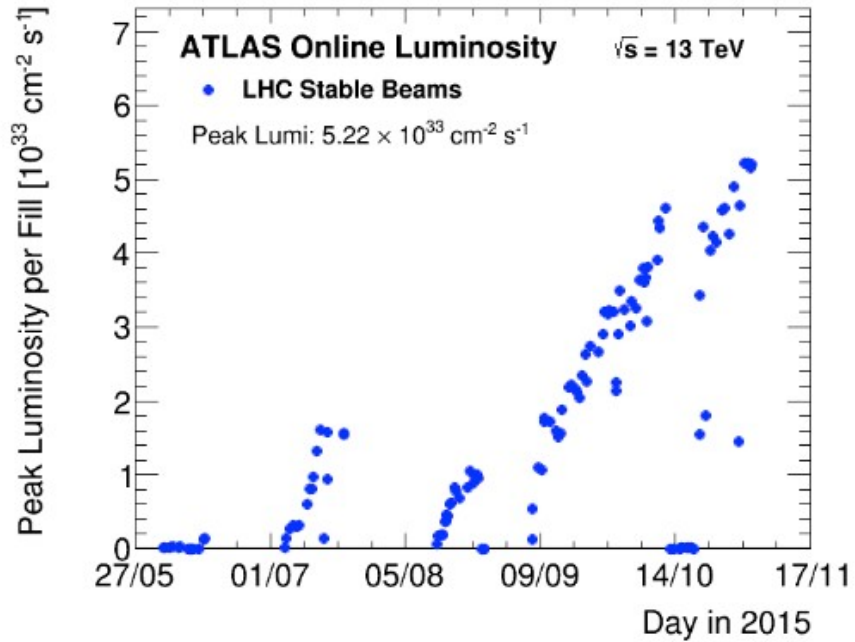


# What happens in a proton-proton collision



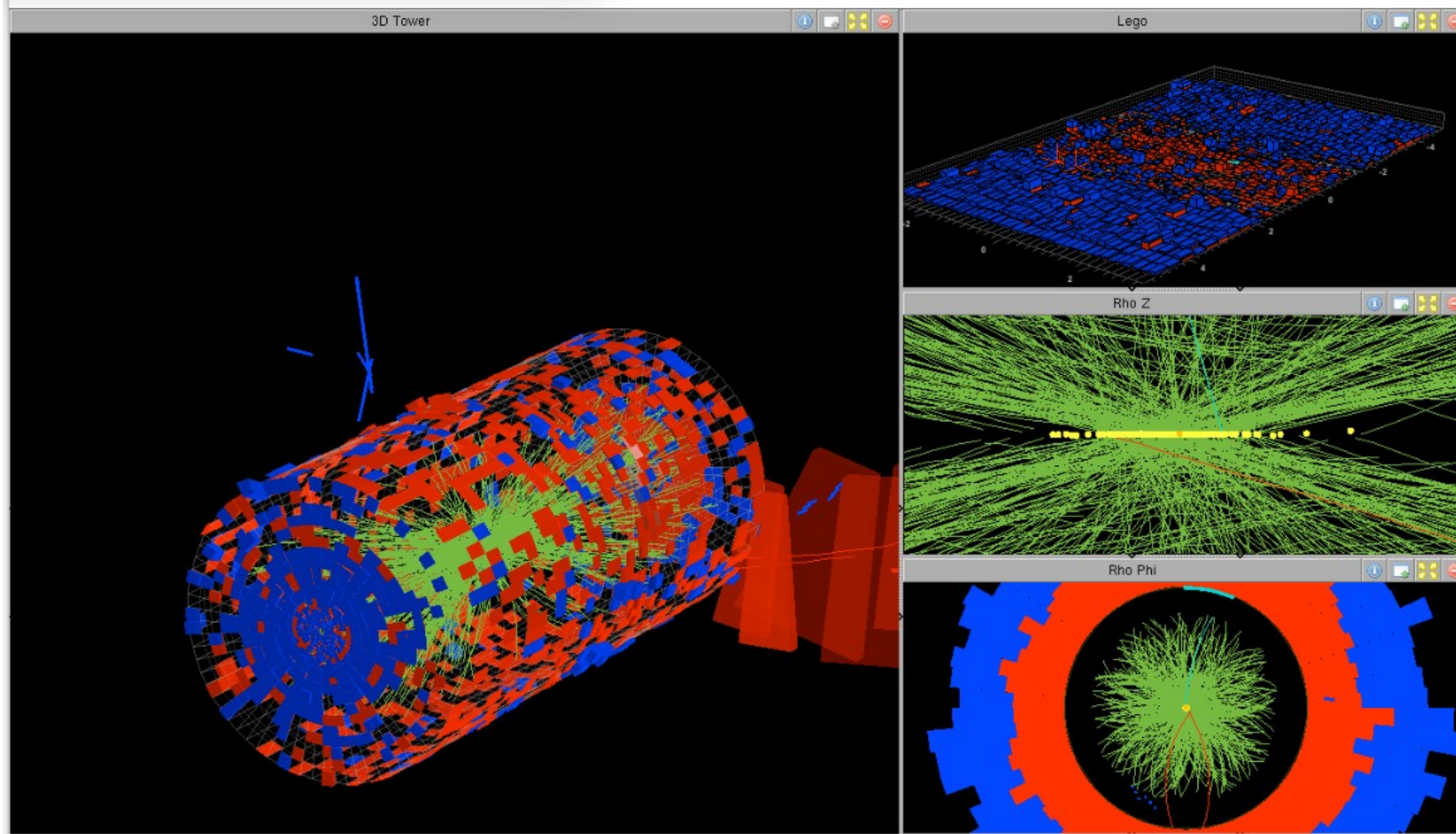
Multiple parton-parton interactions...  
...with their ISR/FSR

# LHC collisions are messy





# 8 TeV CMS collision



- ATLAS and CMS were designed as *general purpose* experiments
  - should be able to find any new physics at the Terascale
  - should also be able to cross-check each others results (different systematics)
- Need to be able to cope with intense environment of LHC collisions (for many years)
  - need radiation hard technology *or*
  - ability to replace components

# A word about coordinates

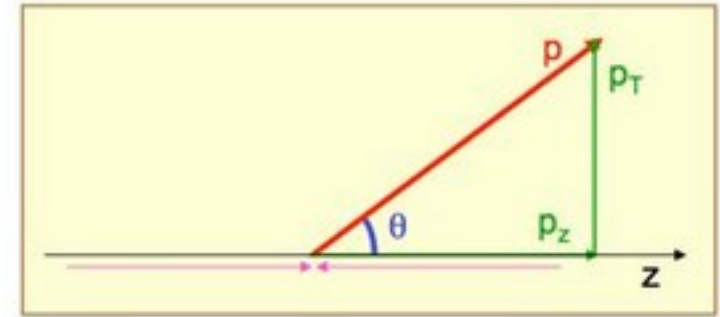
- The overall longitudinal boost is *not known* at a hadron collider

- thus cannot know  $p_z$  for any four vector

- can only work with transverse momentum  $p_T^2 = p_x^2 + p_y^2$

- The polar angle  $\theta$  is not Lorentz invariant

- Instead use rapidity (usually use pseudo-rapidity):



$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z} \approx \eta = -\ln \left( \tan \frac{\vartheta}{2} \right)$$

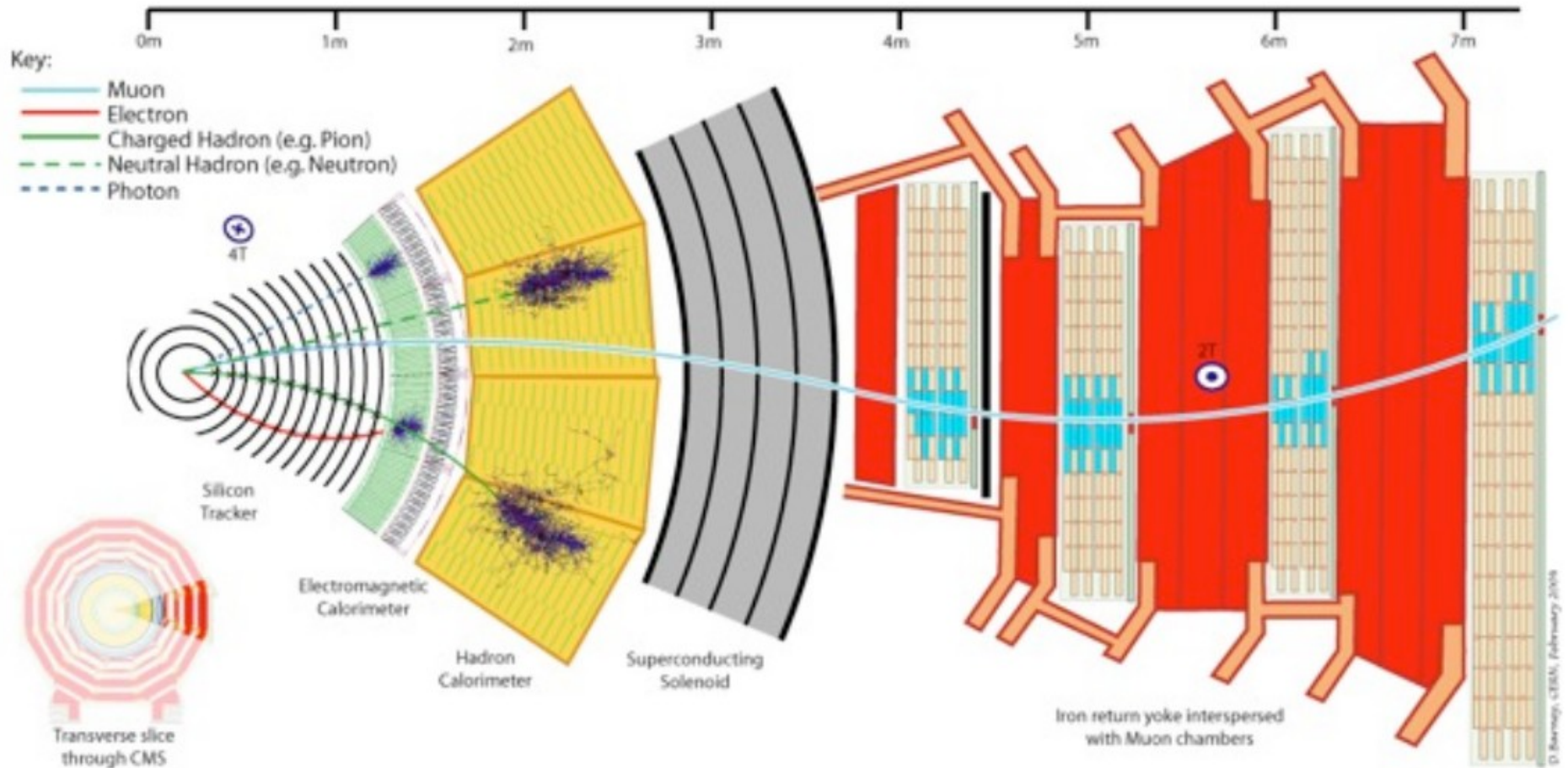
rapidity  $\approx$  pseudorapidity (|) in massless approximation

- This *is* Lorentz invariant, and tends to infinity at the beam axis

- An LHC detector needs:

- full azimuthal coverage (need to measure **missing transverse energy**)
- forward coverage (up to  $|\eta| = 5$ )
- good **lepton identification** with excellent  $\mathbf{p}_T$  resolution
- good **photon identification**
- good **jet energy resolution** (and **missing transverse energy** resolution)
- efficient **secondary vertex finding** (b-tagging, tau ID)

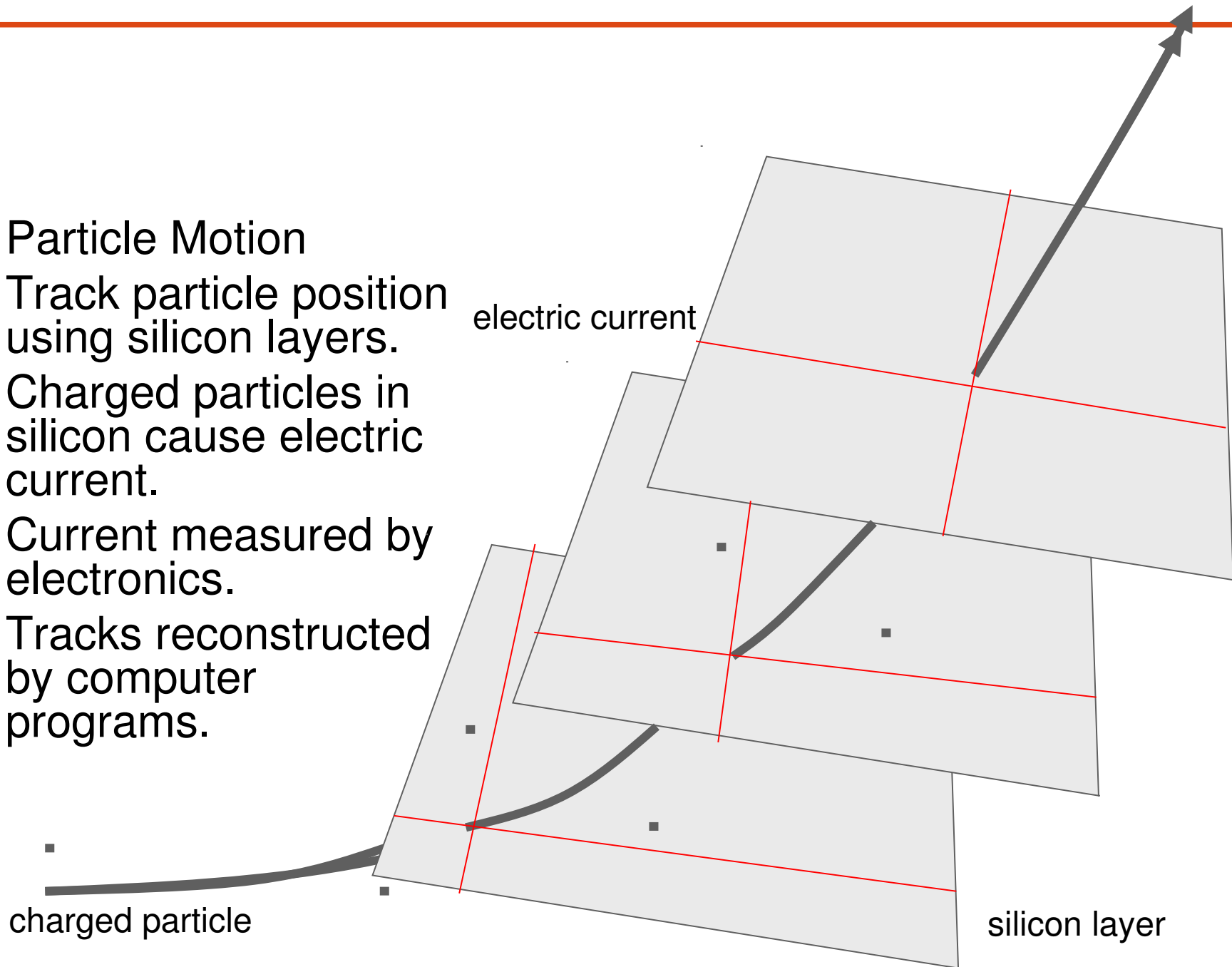
# How does a detector work?



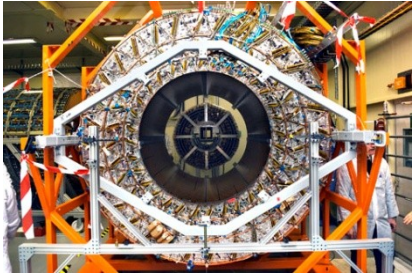
- Built around collision point → full azimuthal coverage
- Momentum: measure charged particle trajectories in inner part of detector (apply B field)
- Energy: slam particles into Pb or steel, stop them, measure the energy released (calorimeters)
- Muons: have special detectors on the outside of the detector

# Measuring position

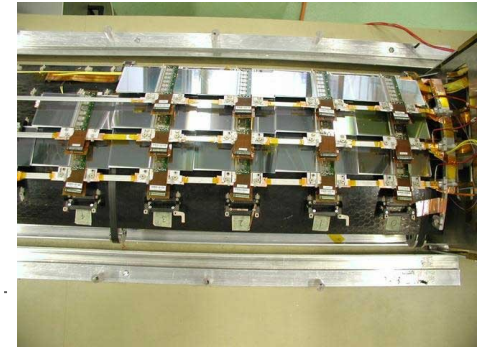
- Particle Motion
- Track particle position using silicon layers.
- Charged particles in silicon cause electric current.
- Current measured by electronics.
- Tracks reconstructed by computer programs.



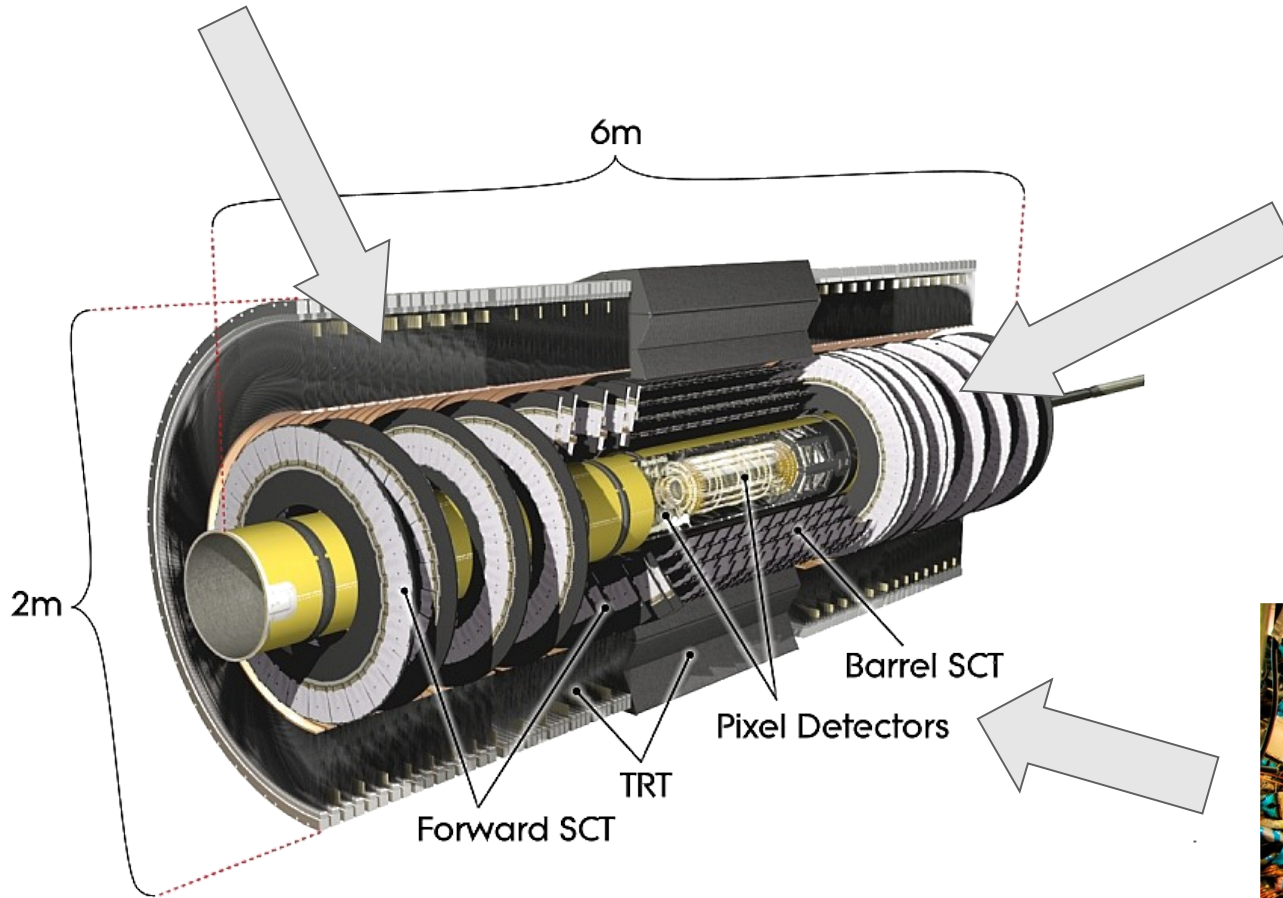
# ATLAS tracker



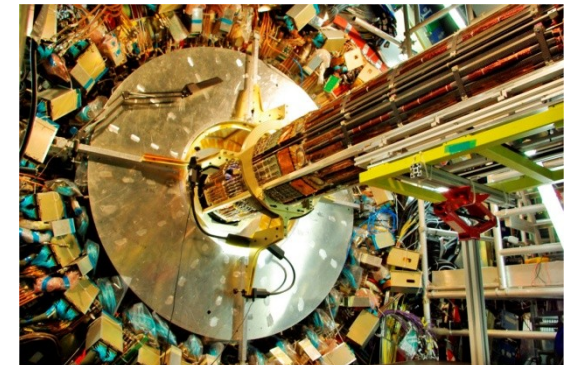
420,000 Xenon Straws



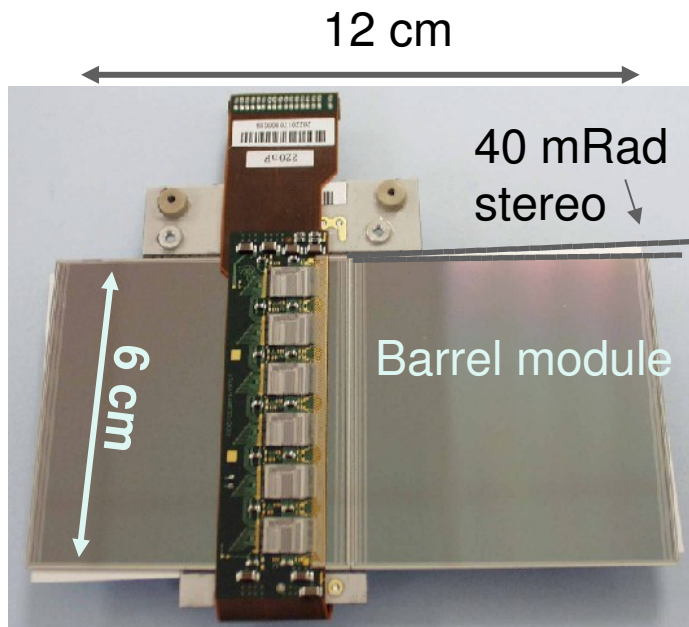
6 Million Silicon Strips



135 Million Pixels

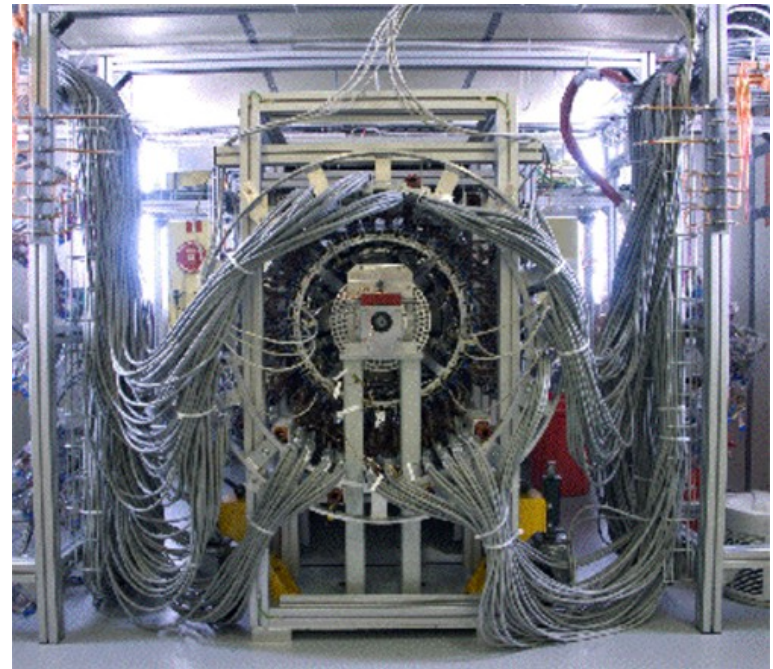


# ATLAS semiconductor tracker



- 4 barrel layers
- 9 endcap wheels
- 6 million silicon strips

- Need to configure and read-out each module.





# Measuring momentum

- For a charged particle moving in a magnetic field:

$$\mathbf{p} = \mathbf{B} \times \mathbf{q} \times \mathbf{r}$$

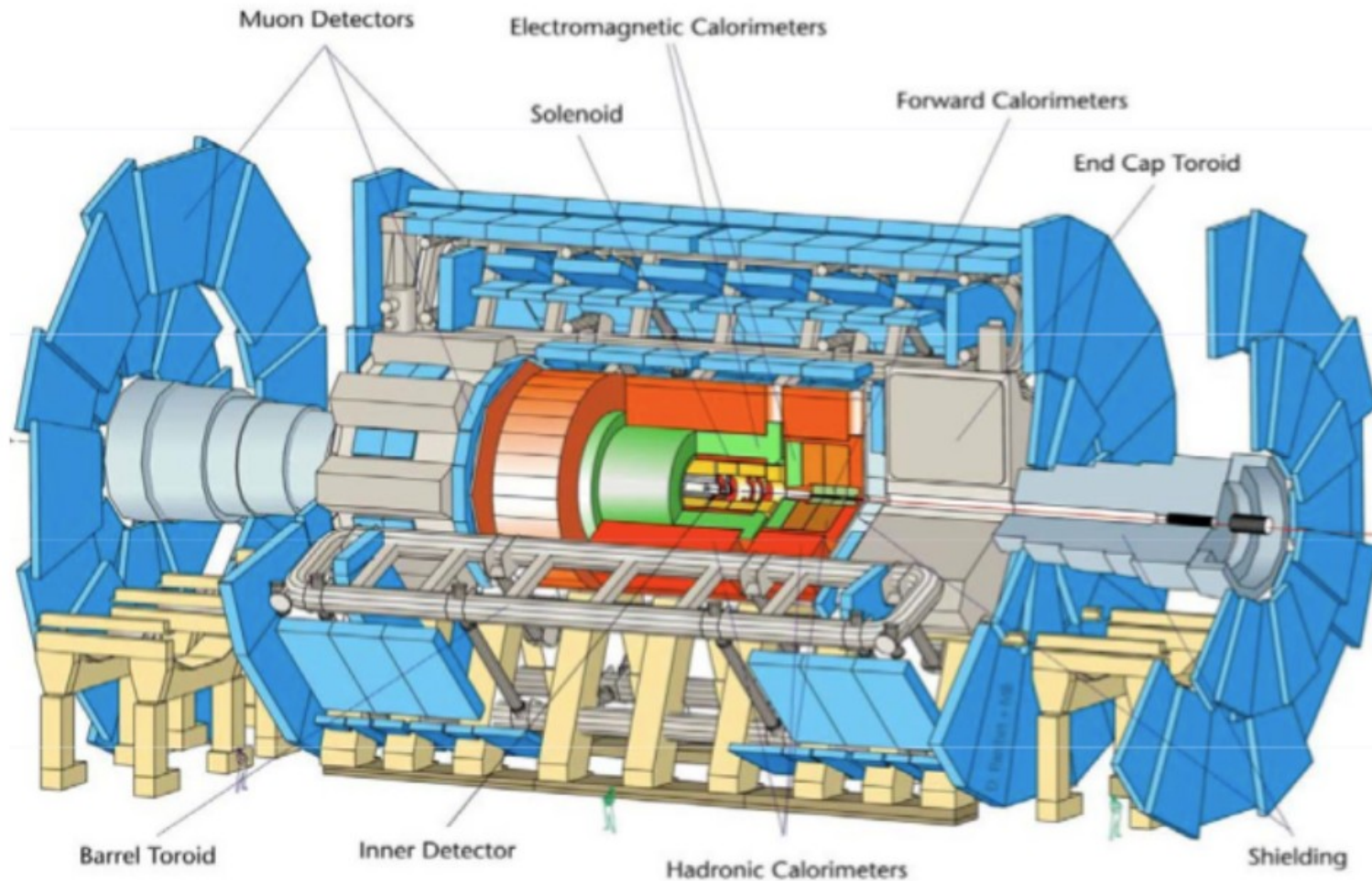
- The momentum resolution goes as:

$$\frac{dp}{p} \propto \frac{p}{B\ell^2}$$

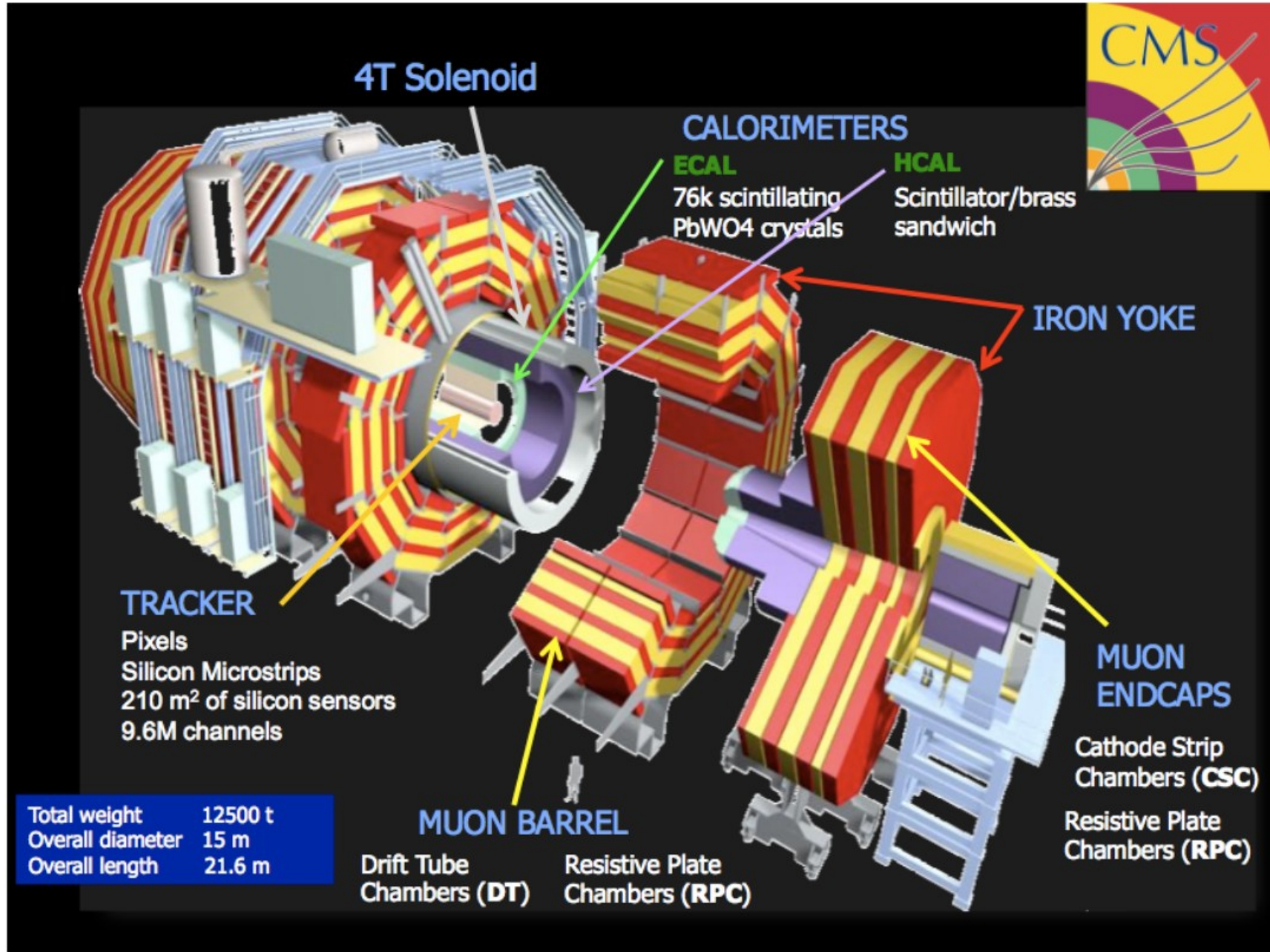
- Therefore need either:

- a large lever arm (this is basically what ATLAS did)
- a large magnetic field (this is basically what CMS did)

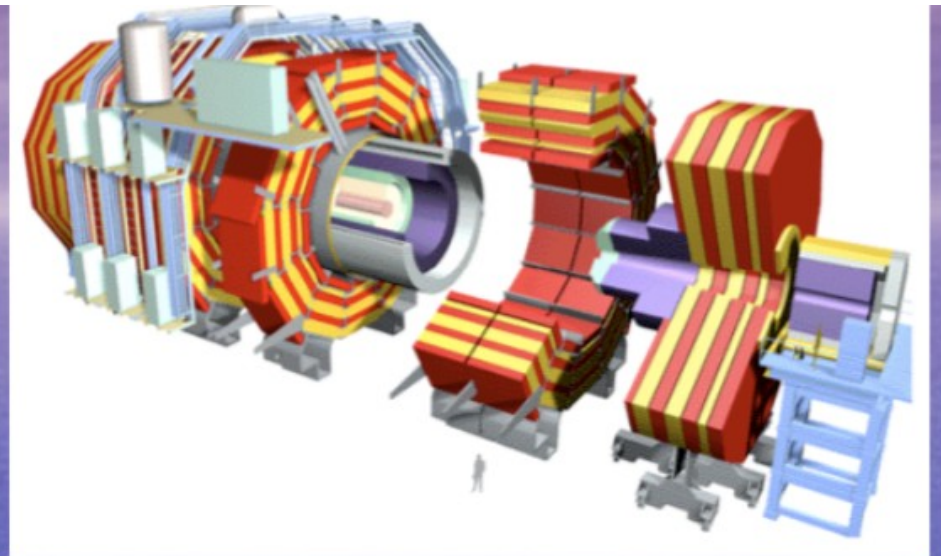
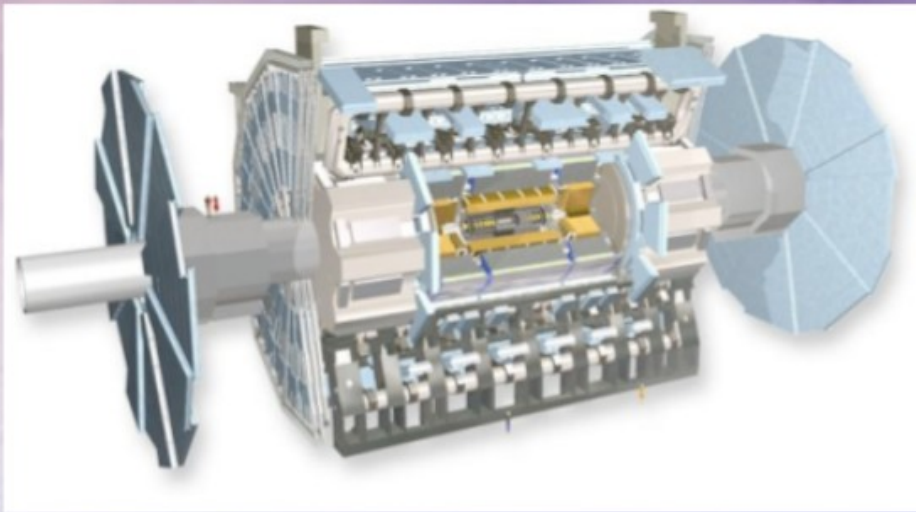
# ATLAS



Diameter	25 m
Barrel toroid length	26 m
End-cap end-wall chamber span	46 m
Overall weight	7000 Tons



# Important numbers



Tracker

$|\eta| < 2.5$  coverage

$$\sigma / p_T \approx 5 \cdot 10^{-5} p_T \oplus 0.01 [\text{GeV}]$$

$|\eta| < 2.6$  coverage

$$\sigma / p_T \approx 1.5 \cdot 10^{-5} p_T \oplus 0.005$$

EM Calorimeter

$|\eta| < 4.9$  coverage

$$\sigma / E \approx 10\% / \sqrt{E} [\text{GeV}]$$

$|\eta| < 4.9$  coverage

$$\sigma / E \approx 2 - 5\% / \sqrt{E}$$

HAD Calorimeter

$|\eta| < 4.9$  coverage

$$\sigma / E \approx 50\% / \sqrt{E} \oplus 0.03 [\text{GeV}]$$

$|\eta| < 4.9$  coverage

$$\sigma / E \approx 100\% / \sqrt{E} \oplus 0.05$$

Muon Spectrometer

$|\eta| < 2.7$  coverage:

$$\sigma / p_T \approx 0.07 \text{ (1 TeV muons)}$$

$|\eta| < 2.6$  coverage:

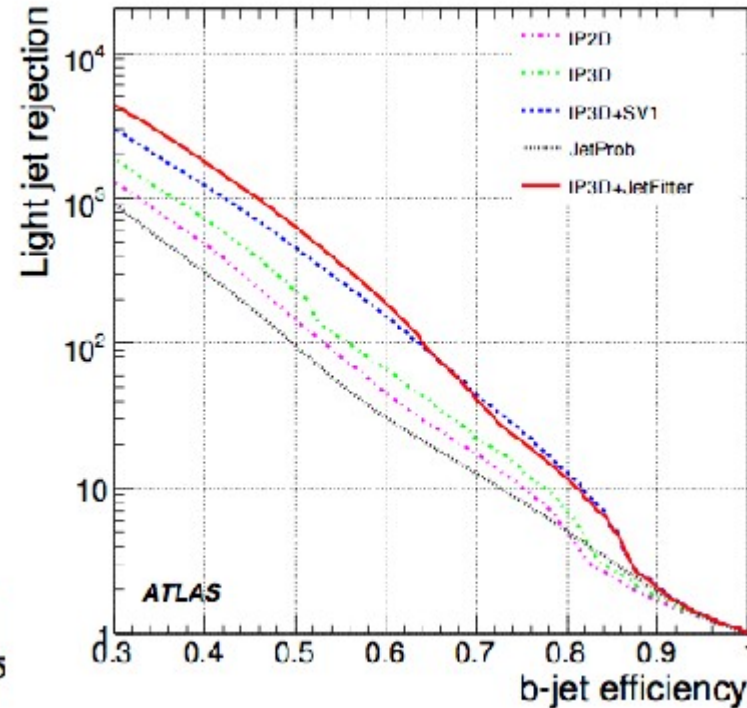
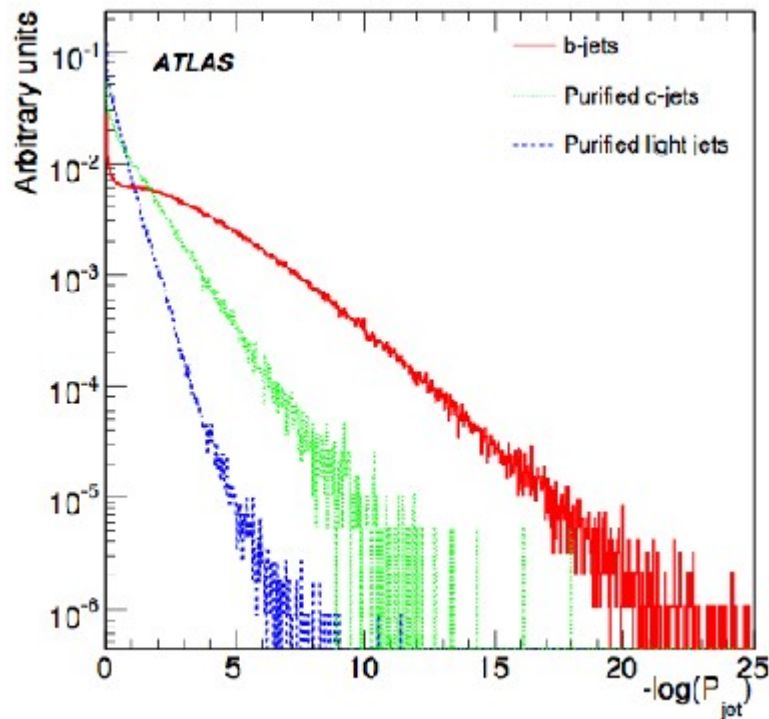
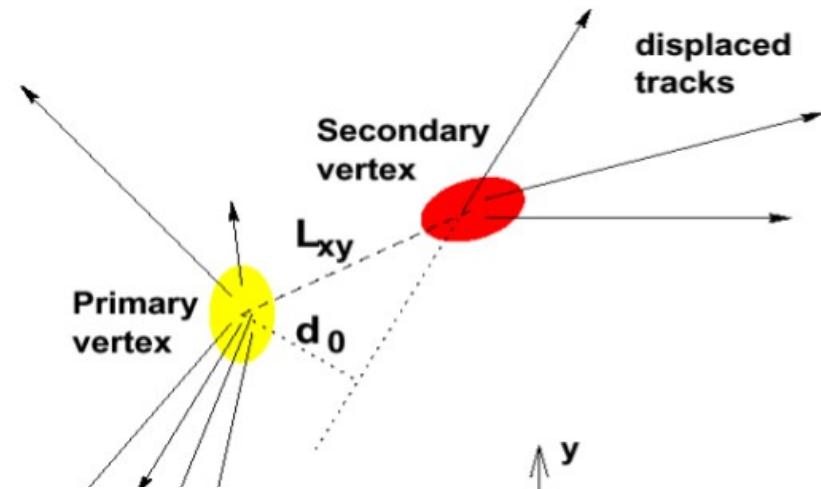
$$\sigma / p_T \approx 0.10 \text{ (1 TeV muons)}$$

- All quarks except the top quark form bound states at the LHC
- In the case of *b* quarks, a single *B* or *D* meson carries most of the energy
- When a *B* meson decays it will:
  - typically travel 450  $\mu\text{m}$  before decaying
  - typically decay to  $Dl\nu$  or  $D + \text{other hadrons}$
  - a *D* meson typically decays to kaons plus leptons or hadrons
- All of this makes *b* jets a bit special:
  - typically has a displaced vertex (or two)
  - often has more than 2 tracks missing the collision point by 200-1000  $\mu\text{m}$
  - has a muon embedded in it  $\sim 15\%$  of the time
- Charm quarks show similar behaviour, but are less likely to have displaced tracks and vertices

# *b*-tagging

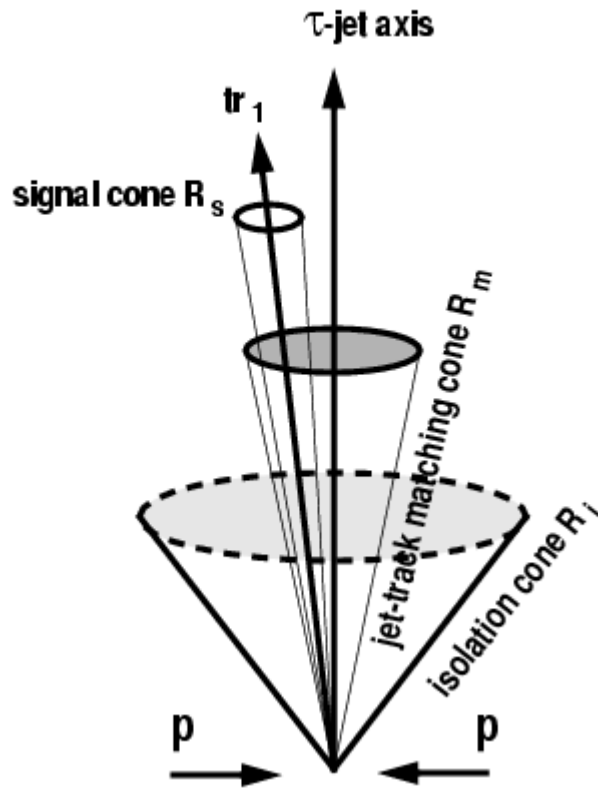
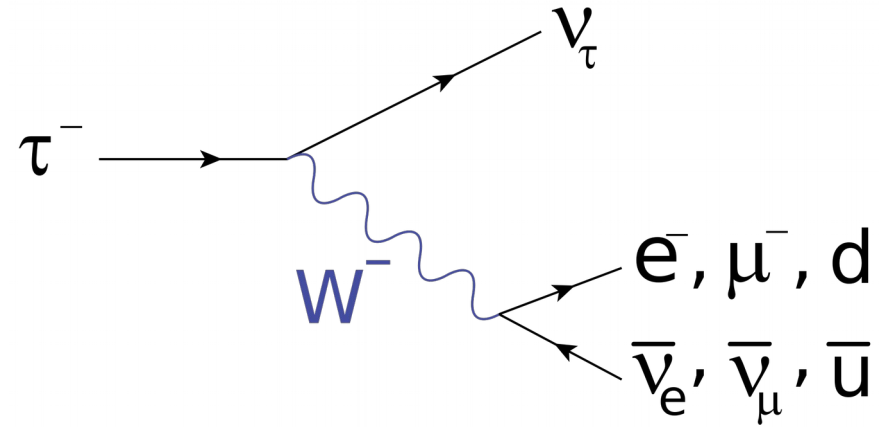
- Measure the decay length  $L_{xy}$  or impact parameter  $d_0$

- typically a modern *b* tag algorithm uses multivariate techniques



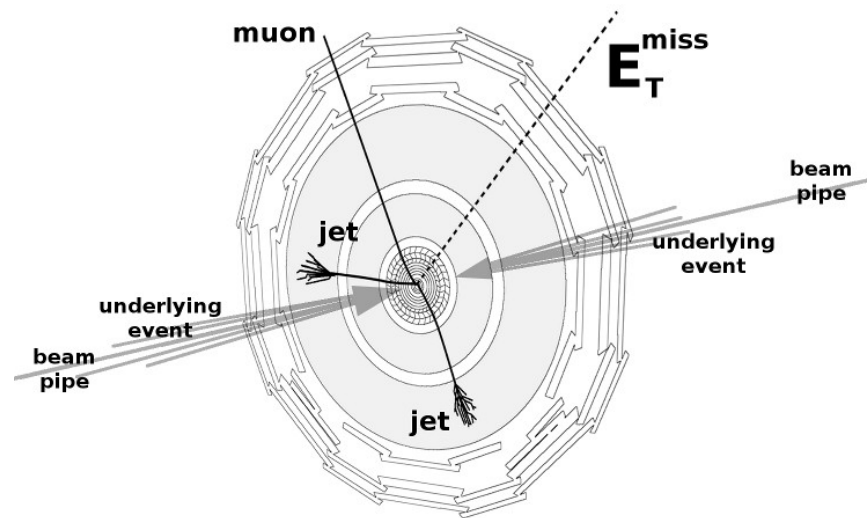
# $\tau$ -tagging

- 34% to  $e$  or  $\mu$
- 50% to one charged hadron (plus neutral hadrons)
- 15% to 3 charged hadrons (plus neutral hadrons)
- **Always** get a tau neutrino – can never measure  $E_\tau$ !



- Leptonic taus look just like  $e$  or  $\mu$
- Hadronic taus look like special jets
  - 1 or 3 tracks (typical jets  $\sim 5 - 20$  tracks)
  - invariant mass of decay products  $< m_\tau \sim 1.8$  GeV
- For, e.g. GMSB, tau searches are crucial
- Can tag taus with typical efficiency of 40%

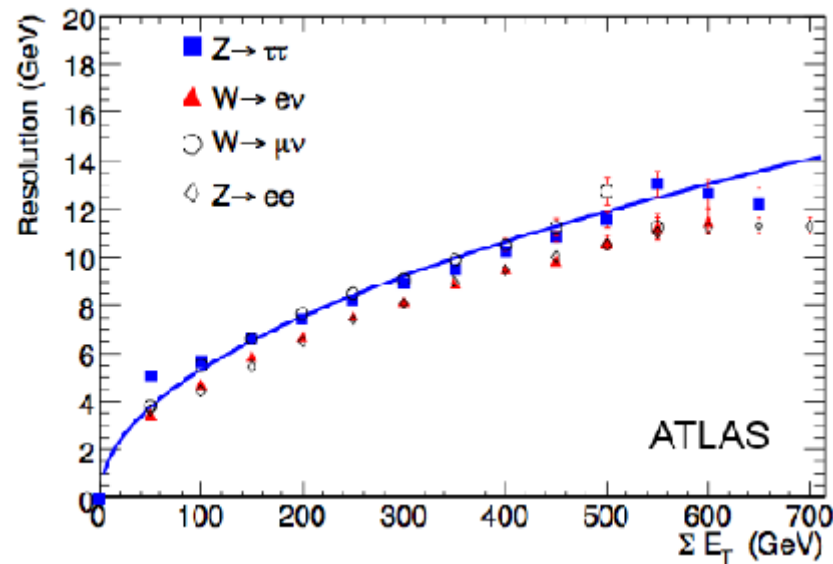
# Missing transverse energy



- We expect a certain amount of missing  $E_T$  in the SM
  - neutrino production
  - mismeasurement of visible objects
  - other fake sources (e.g. cosmic muons overlaid on collision data)
- Substantial missing  $E_T$  is a *generic* feature of models with dark matter candidates
  - these models are cooked to feature a Weakly Interacting Massive Particle (WIMP)
  - this is stable, neutral and invisible to ATLAS and CMS
  - e.g. SUSY, UED, Little Higgs with T parity, many more...



# Missing transverse energy



- Careful calibration required for accurate MET
- Use standard processes known to have specific MET distributions
  - important to calibrate at several points → extrapolation to higher values
- Rough MET resolution follows the calorimeter resolution

$$\sigma = 0.57 \sqrt{E \text{ (GeV)}}$$

# Summary of detector output

- LHC detector output (after reconstruction) consists of:
  - four vectors of jets, leptons and photons (plus particle identification)
  - tagging of  $b$  jets  $\sim 70\%$  of the time
  - tagging of  $\tau$  leptons  $\sim 40\%$  of the time
  - missing transverse energy
  - EXTRA: evidence for exotic objects (long-lived sparticles?)
- This is our complete LHC toolkit!

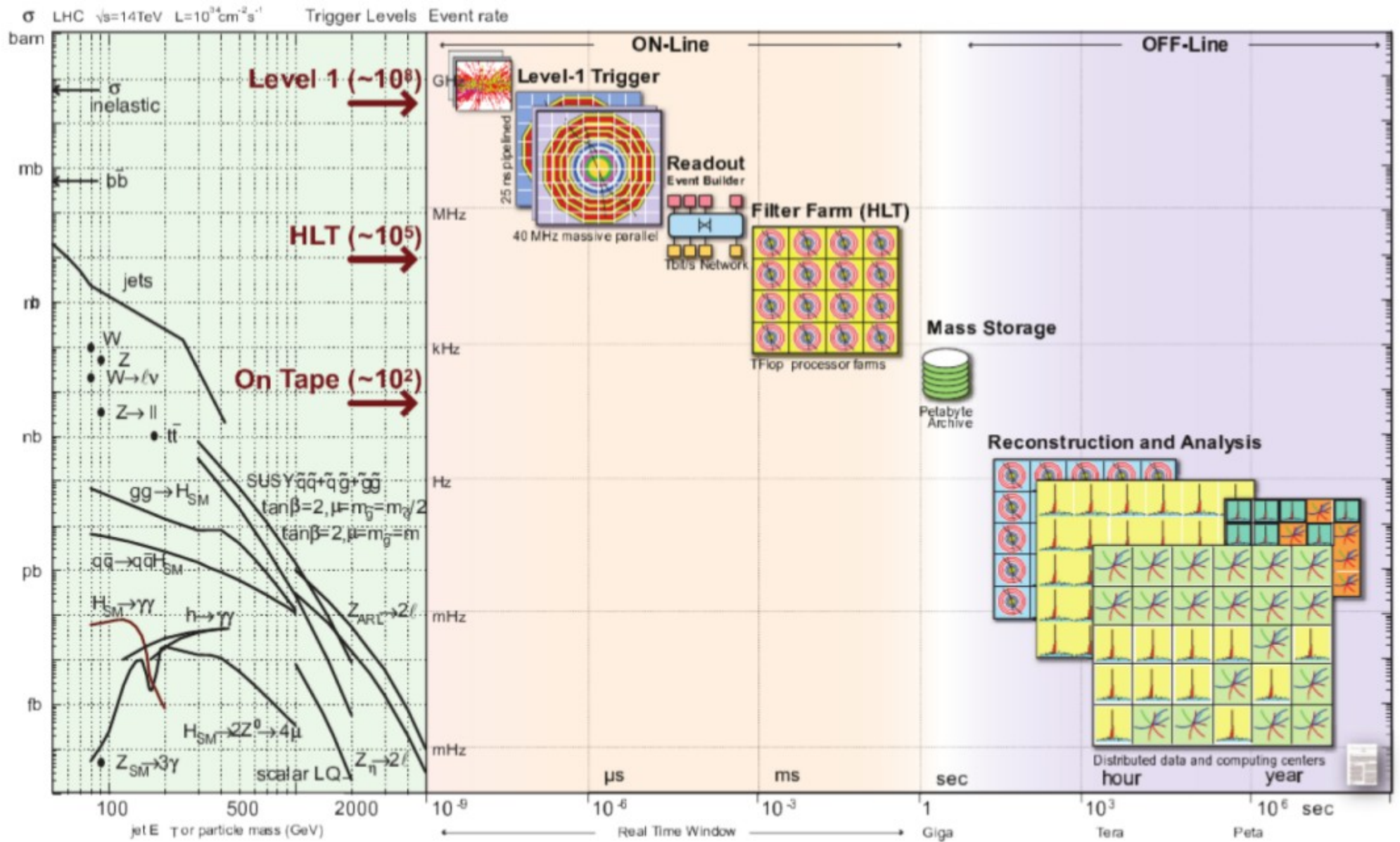
***The whole game of experimental SUSY searches is to use only this information to discover sparticle production, then measure sparticle properties in case of discovery***

# One last thing: triggers

---

- Proton-proton collision rate at the LHC is 40 MHz
- Pixel detector has  $\sim 66\text{M}$  readout channels,  $O(10^4)$  are hit per event
- Data size to be stored per pixel is 4 bytes
  - $\Rightarrow > 120,000$  TB of data per day!
- Need to reduce data volume online: use a **trigger**
  - hardware + software triggers reduce events in real time
  - can use a variety of detector information

# Interesting things << boring things



# Trigger design

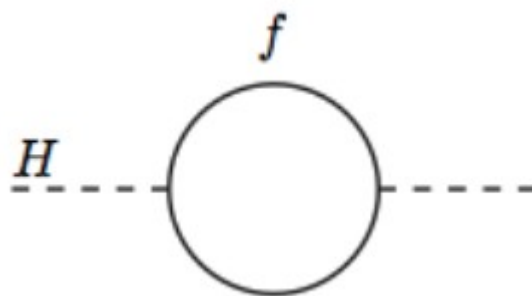
- Good knowledge of detector and signatures is needed to efficiently select interesting events
  - relevant detector parts and their performance
  - needed/desired measurement precision
  - physical properties of signal and background events (kinematics, multiplicities, etc)
- Trigger strategy needs careful planning and has scientific, technical and political elements
  - bandwidth will always be limited
- Theorists:
  - don't forget that untriggered events are *lost forever!*
  - gives you a “get out of jail free card” for your favourite difficult models
- Experimentalists:
  - trigger strategy plays a central role in designing effective analyses
  - do not be afraid of rewriting the rule book (e.g. razor triggers)

## A very simple trigger menu example

Object	Physics coverage	Object name
electrons	Higgs, new gauge bosons, extra dim., <b>SUSY</b> , W/Z, top	e25i, 2e15i, e60
Photons	Higgs, <b>SUSY</b> , extra dim.	$\gamma$ 60, 2 $\gamma$ 20i
Muons	Higgs, new gauge bosons, extra dim., <b>SUSY</b> , W/Z, top	$\mu$ 20i, 2 $\mu$ 10
Jets	<b>SUSY</b> , compositness, resonances	j400, 3j165, 4j110
Jets+missEt	<b>SUSY</b> , leptoquarks	j70+xE70
Tau+missEt	Extended Higgs models (e.g. MSSM), <b>SUSY</b>	$\tau$ 35i+xE45

- Note SUSY trigger requirements overlap with a lot of other SM & BSM scenarios

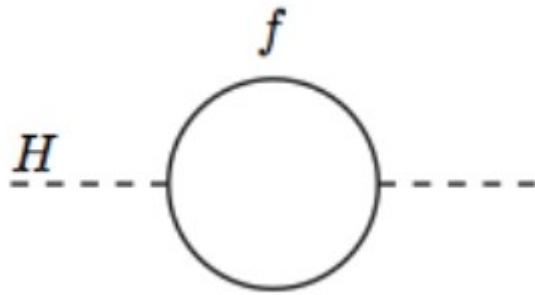
# Why bother with SUSY at the LHC at all?



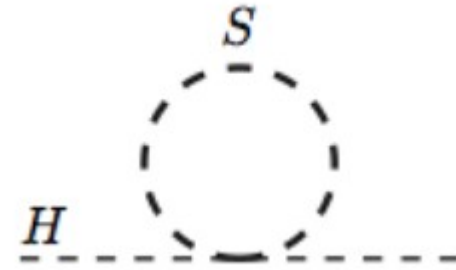
$$\Delta m_H^2 = \frac{\lambda_f^2}{8\pi^2} \left[ -\Lambda^2 + 6m_f^2 \ln \frac{\Lambda}{m_f} \right]$$

- The Higgs mass receives quantum corrections through Feynman diagrams as shown
- We know that  $m_h \sim 125$  GeV
  - naively this should be of the order of the Planck scale
  - this is the **hierarchy problem**

# The hierarchy problem solved



$$\Delta m_H^2 = \frac{\lambda_f^2}{8\pi^2} \left[ -\Lambda^2 + 6m_f^2 \ln \frac{\Lambda}{m_f} \right]$$



$$\Delta m_H^2 = \frac{\lambda_S}{16\pi^2} \left[ \Lambda^2 - 2m_S^2 \ln \frac{\Lambda}{m_S} \right]$$

- Scalar sparticles provide extra contributions to the Higgs quantum corrections

- these conspire to cancel the  $\Lambda^2$  term

- We are left with a *much nicer* logarithmic correction:

$$\Delta m_H^2 \approx \frac{\lambda_f^2}{4\pi^2} (m_S^2 - m_f^2) \ln \frac{\Lambda}{m_S}$$

- *But:* need  $m_S$  not too far from  $m_f \Rightarrow$  TeV scale SUSY?



# Naturalness gets a bit more complicated

## 1. Tree level: light higgsinos

- For large  $\tan \beta \rightarrow$  EWSB condition  $\frac{M_Z^2}{2} = \frac{(m_{H_d}^2 + \Sigma_d) - (m_{H_u}^2 + \Sigma_u) \tan^2 \beta}{\tan^2 \beta - 1} - |\mu|^2$

$$\rightarrow M_Z^2 = -2(m_{H_u}^2 + |\mu|^2) + \dots$$

- $|\mu| \lesssim 200$  GeV for  $\Delta \leq 10$  with  $\Delta[a_i] = \frac{\partial \ln m_Z^2}{\partial \ln a_i^2}$ ,  $\Delta = \max \Delta[a_i]$  and fundamental parameters  $a_i$

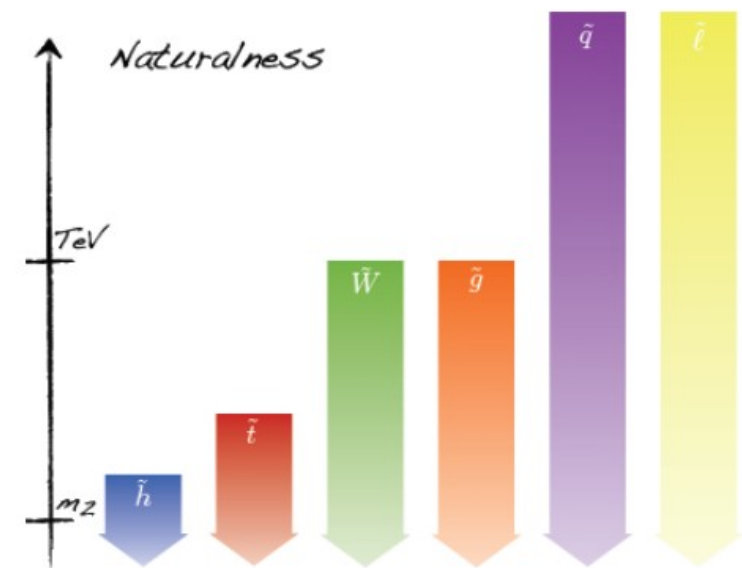
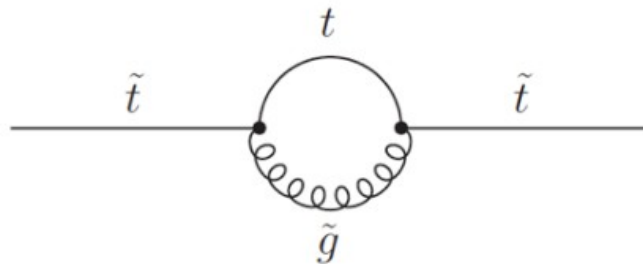
## 2. 1-loop: light stops and winos

- Stop masses below 400 GeV for  $\Delta \leq 10$
- Wino masses below 1 TeV

## 3. 2-loop: light gluinos

- Related to naturalness of other particles
- Stop mass gets corrections from gluino mass

$$\rightarrow m_{\tilde{g}} \lesssim 2m_{\tilde{t}}$$



SUSY mass scales motivated by electroweak naturalness

N. Craig, arXiv:1309.0528

# What about the 125 GeV Higgs?

- Things get more complicated still
- We actually need the radiative corrections to  $m_h$  to be *rather large* in the MSSM

$$M_h^2 \approx M_Z^2 \cos^2 2\beta + \frac{3m_t^4}{4\pi^2 v^2} \left[ \log \frac{\Delta_S^2}{m_t^2} + \frac{X_t^2}{\Delta_S^2} \left( 1 - \frac{X_t^2}{12\Delta_S^2} \right) \right] \quad \text{with} \quad \Delta_S^2 = m_{\tilde{t}_1} m_{\tilde{t}_2}$$

- Either stops are heavy ( $\gg 1$  TeV)  $\rightarrow$  Tension to “naturalness” ( $\rightarrow$  fine tuning)
- or large stop mass splitting (“Maximal mixing”)

$$\mathcal{M}_{\tilde{t}} = \begin{pmatrix} \tilde{m}_{t_L}^2 & \overbrace{m_t(A_t - \mu \cot \beta)}^{X_t} \\ m_t(A_t - \mu \cot \beta) & \tilde{m}_{t_R}^2 \end{pmatrix}$$

- So maybe we shouldn't be surprised that we don't see SUSY at the LHC?

- or maybe the MSSM is too simple?

***In any case we must continue to hunt!***

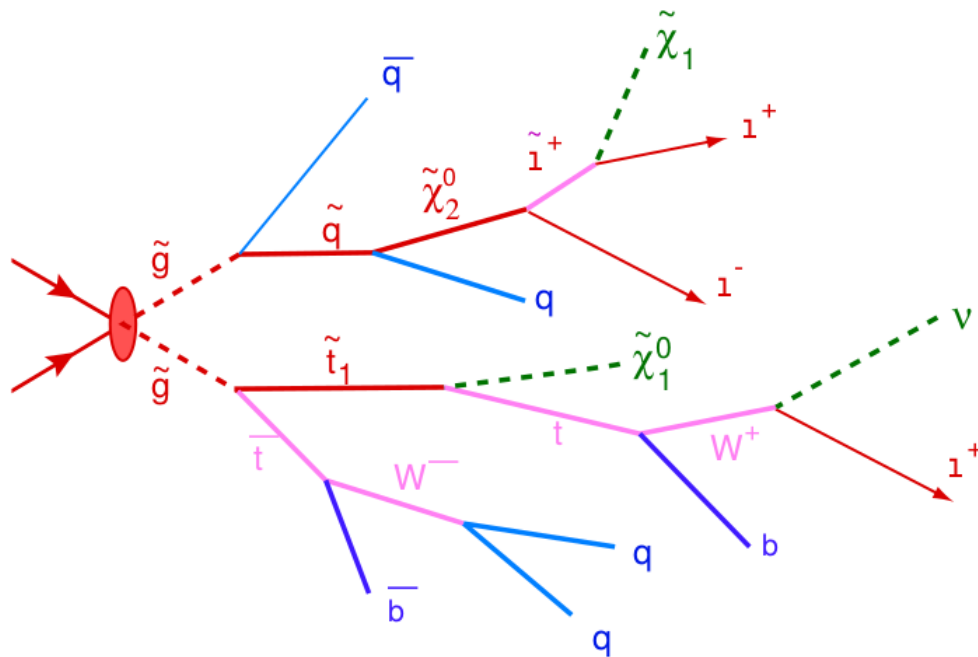
# R-parity conservation

- In the MSSM, baryon number and lepton number are no longer conserved
  - but we have precise limits on their conservation!
- Can introduce a  $Z_2$  symmetry called “R-parity” that prevents the extra couplings:

$$P_R = (-1)^{3(B-L)+2s}$$

- Two immediate consequences:
  - 1) The lightest supersymmetric particle is absolutely stable
  - 2) Sparticles must be *pair-produced* at the LHC
- We will mostly concentrate on R-parity conserving (RPC) searches

# A typical SUSY hard process (assuming R-parity conservation)

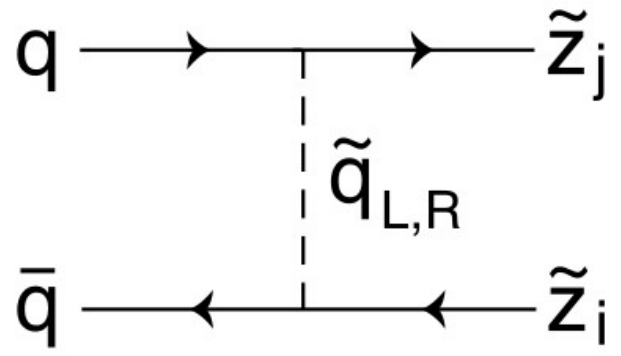
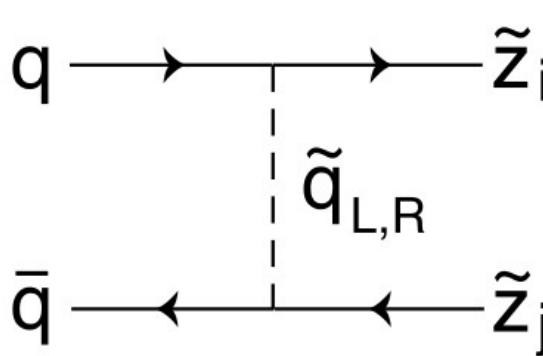
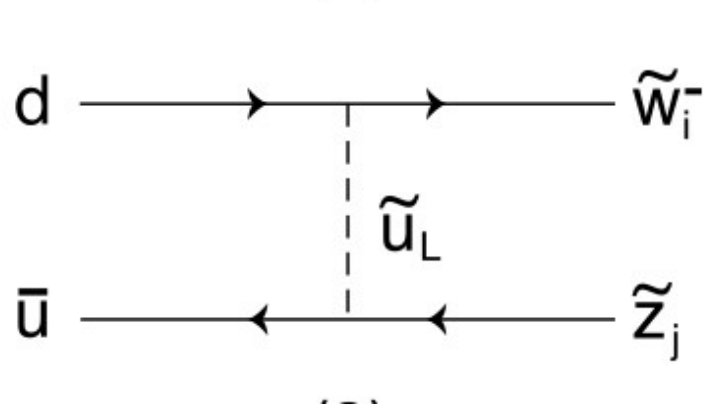
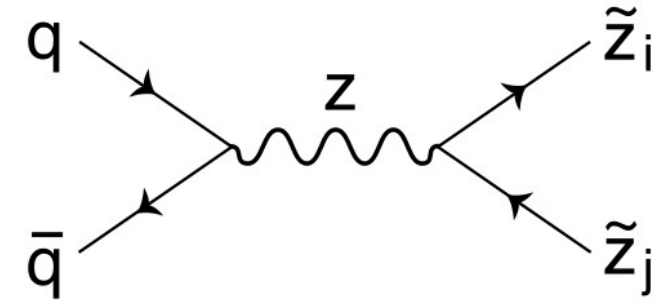
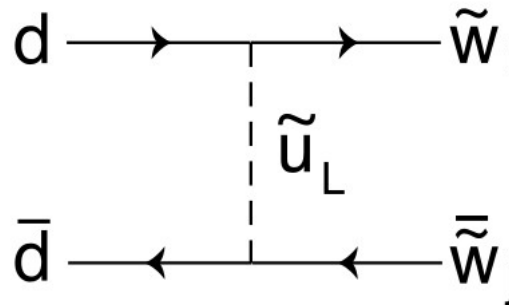
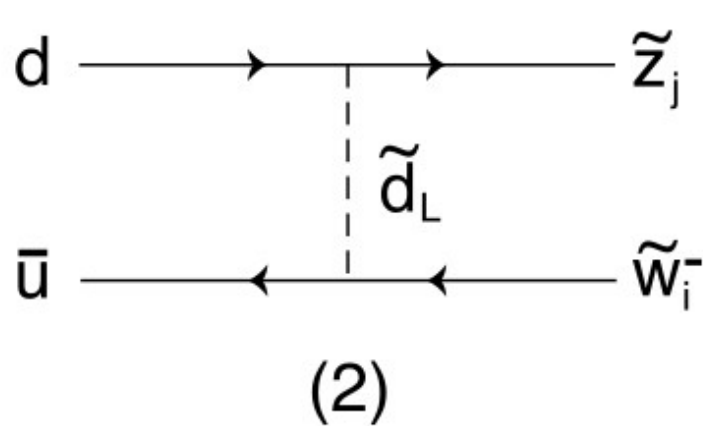
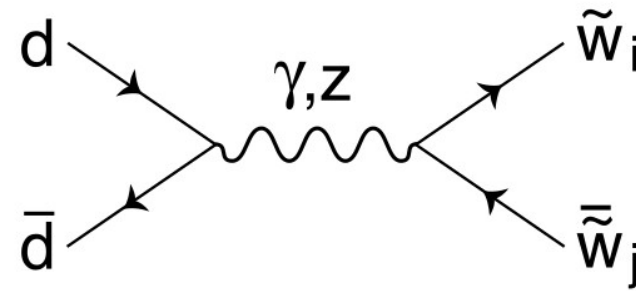
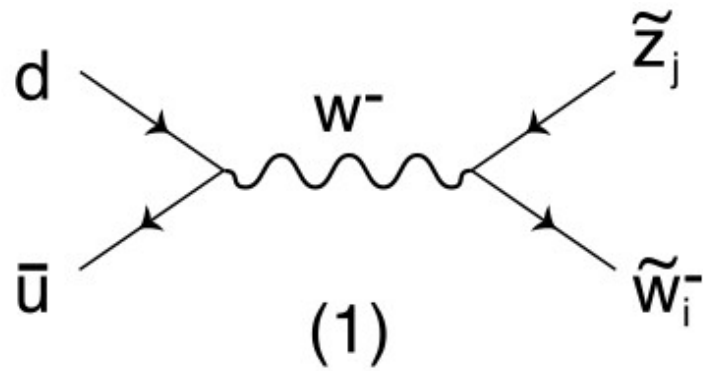


- $\cancel{E}_T$ : from LSP escaping detection
- High  $E_T$  jets: guaranteed if squarks/gluinos if unification of gaugino masses assumed.
- Multiple leptons ( $Z$ ): from decays of Charginos/neutralinos in cascade
- Multiple  $\tau$ -jets or  $b$ -jets ( $h$ ): Often abundant production of third generation sparticles

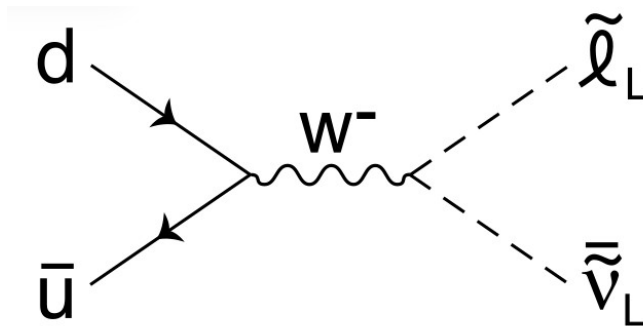
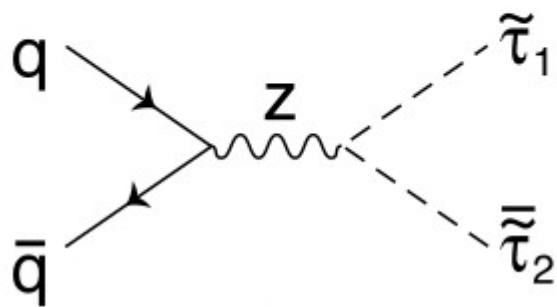
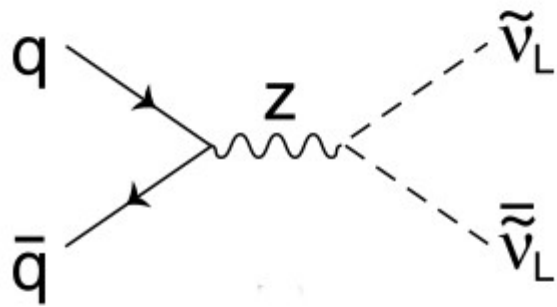
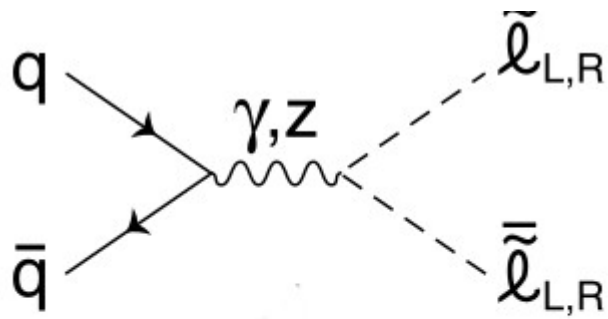
# What would actually happen?

- The answer is **completely model-dependent!**
- Depends on the answers to at least the following questions:
  - is the gluino light?
  - how many of the squarks are light?
  - what are the details of weak gaugino mixing?
  - are there degeneracies in the mass spectrum?
  - how heavy are the extra Higgs bosons, and what are their couplings?
  - is the MSSM true, or something more complicated?
- Please don't forget that **every event is different**
  - will come back to this when I tell you about simplified models

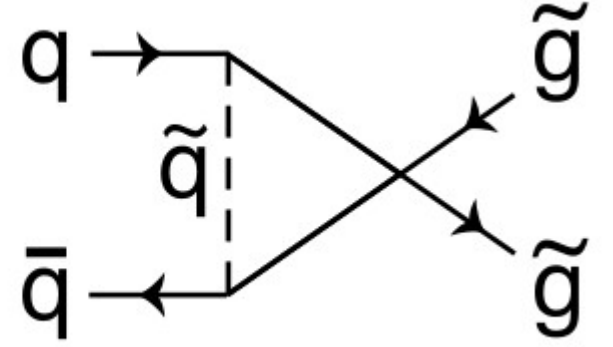
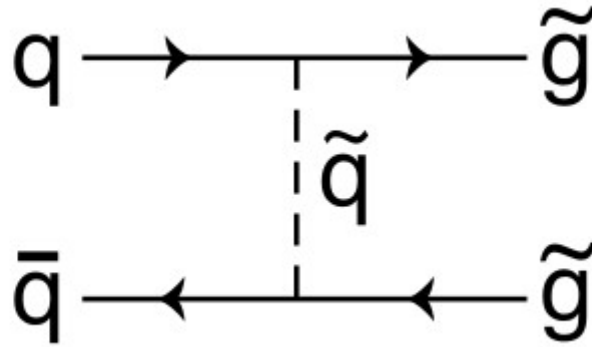
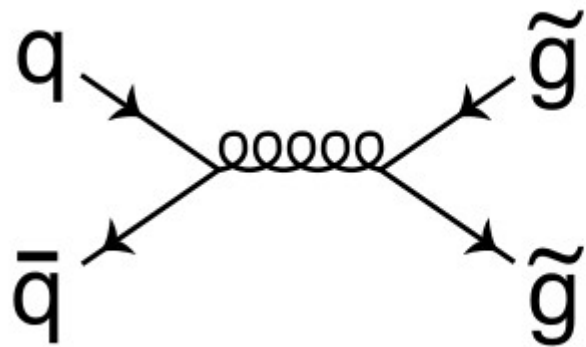
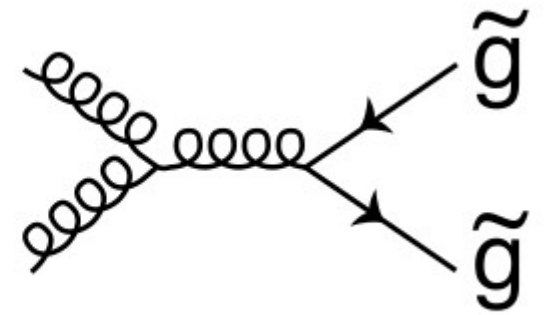
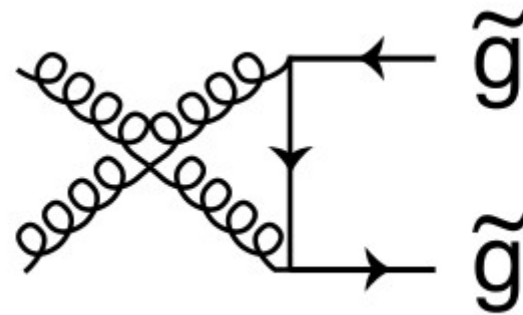
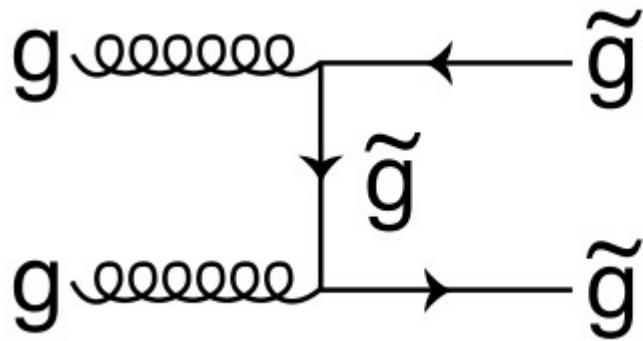
# Tree level SUSY production diagrams: EW gauginos



# Tree level SUSY production diagrams: sleptons

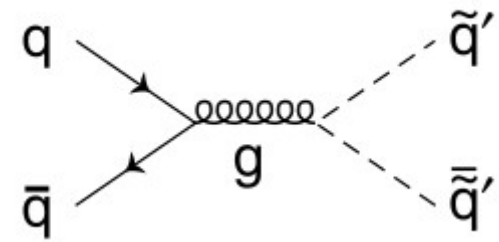
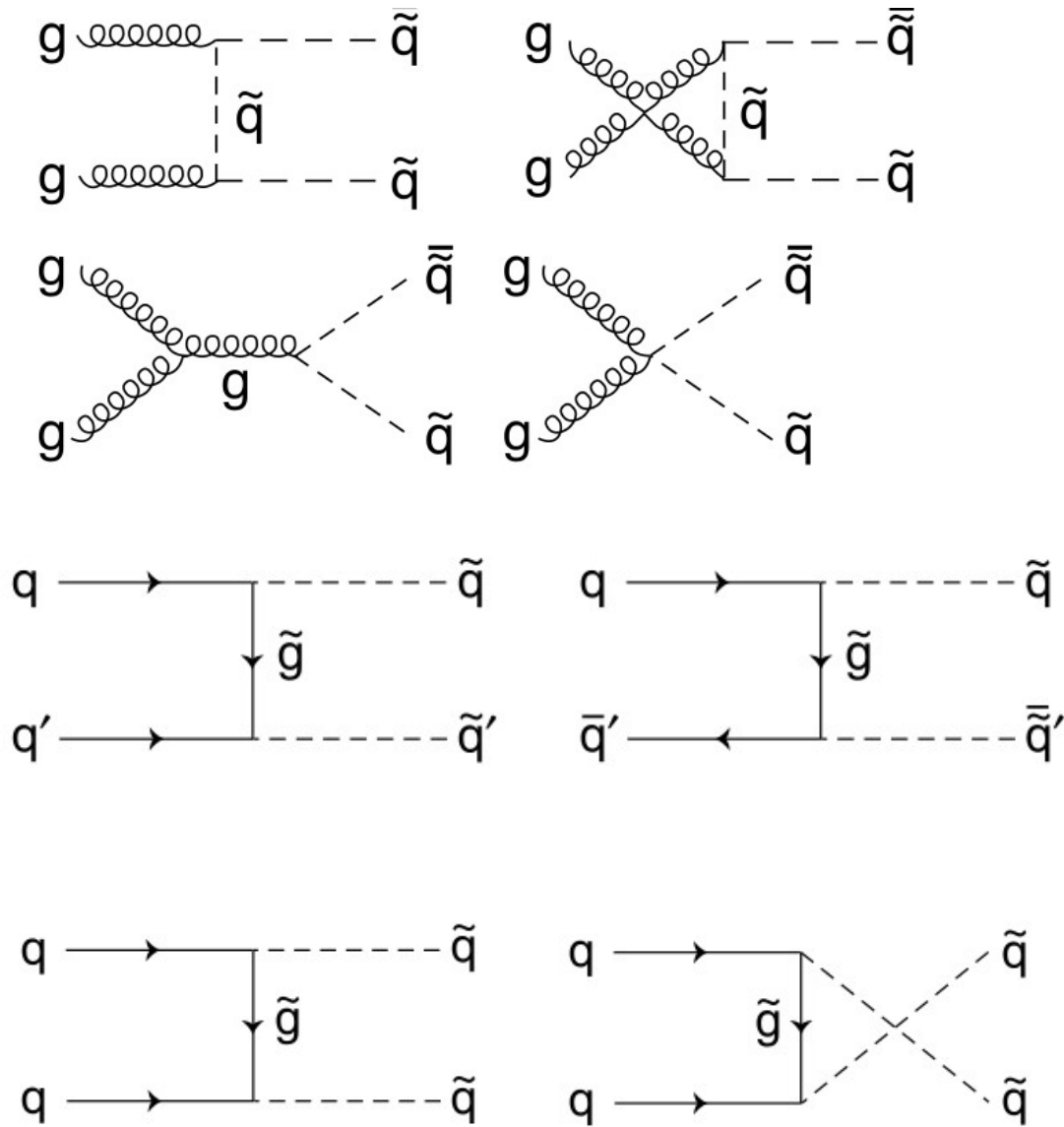


# Tree level SUSY production diagrams: gluino pair

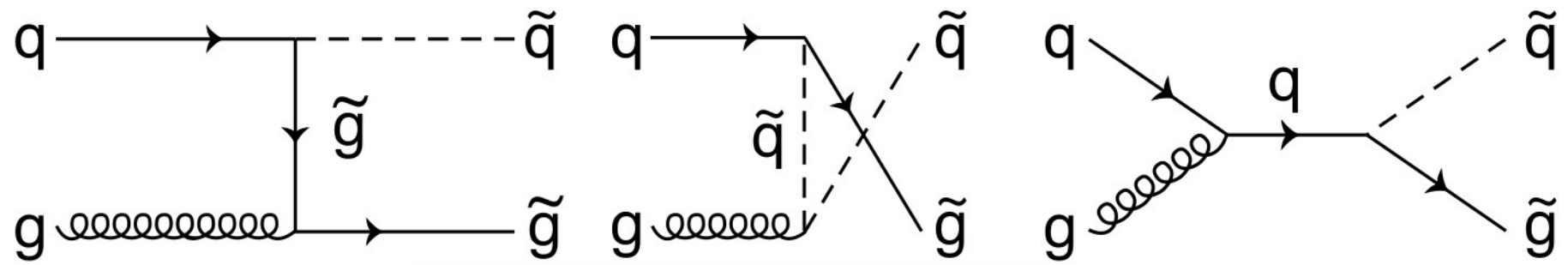




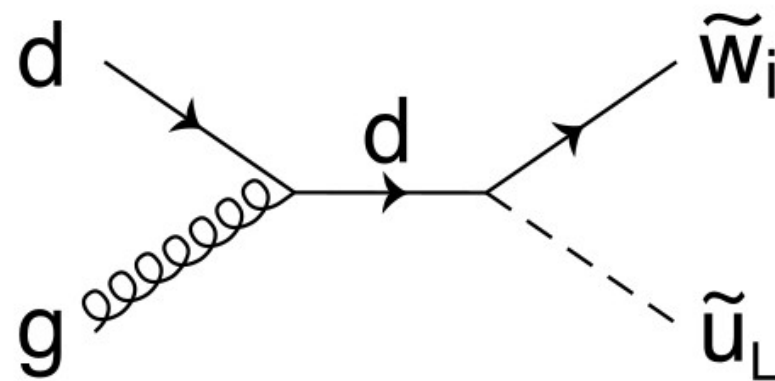
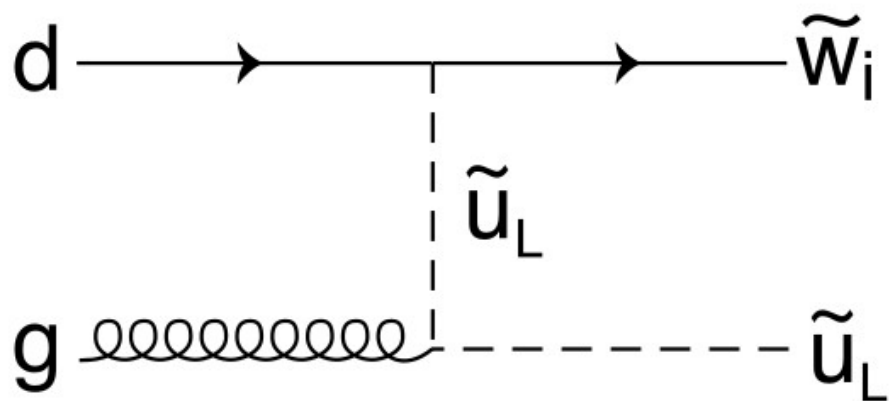
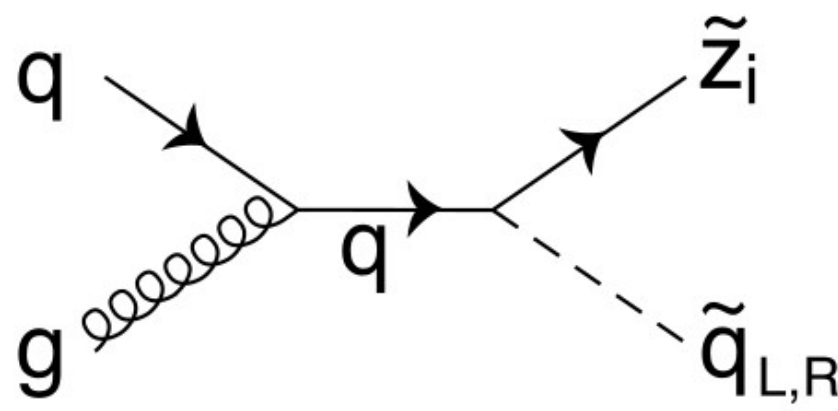
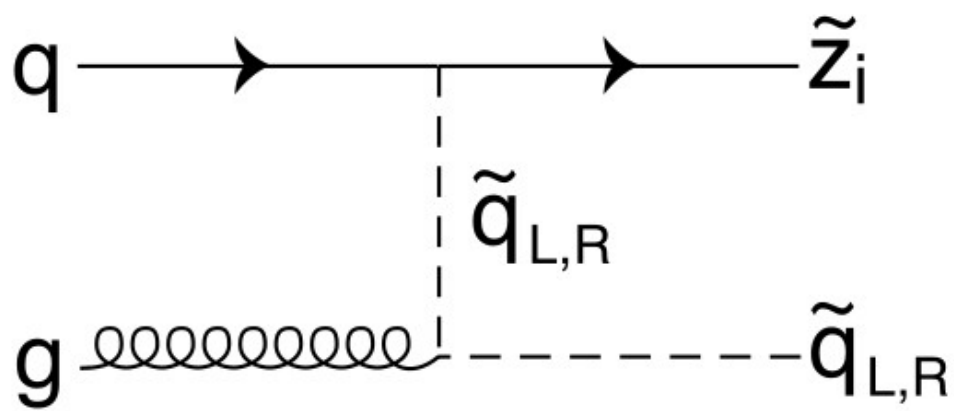
# Tree level SUSY production diagrams: squark pair



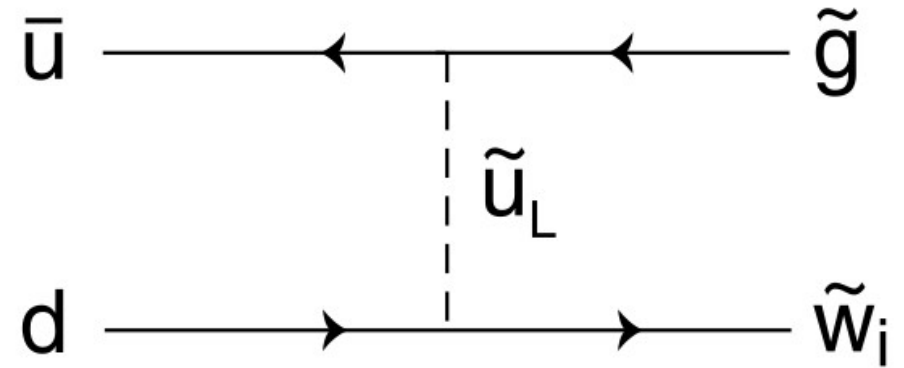
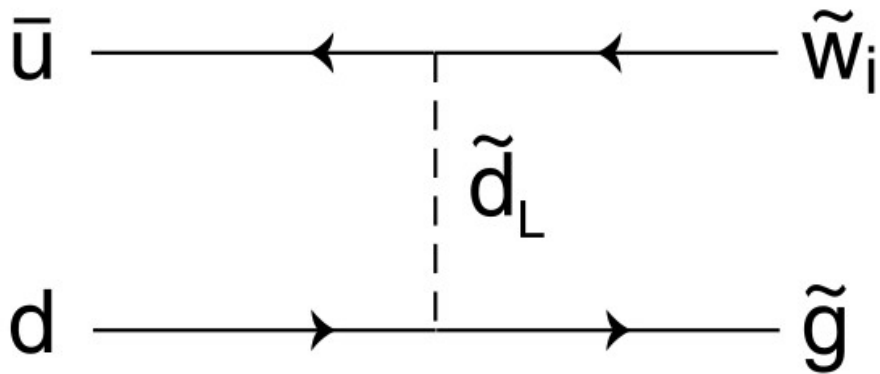
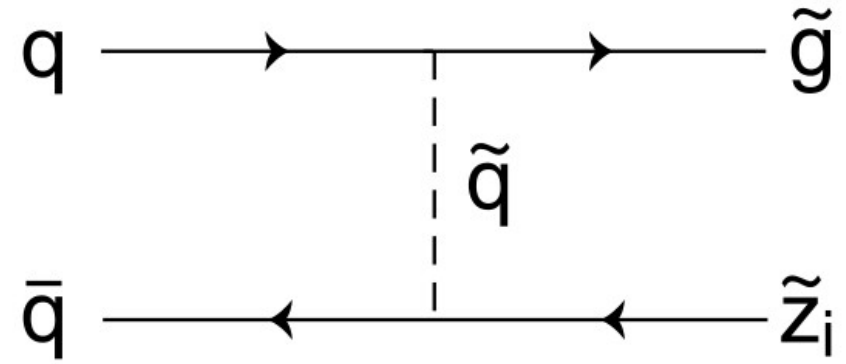
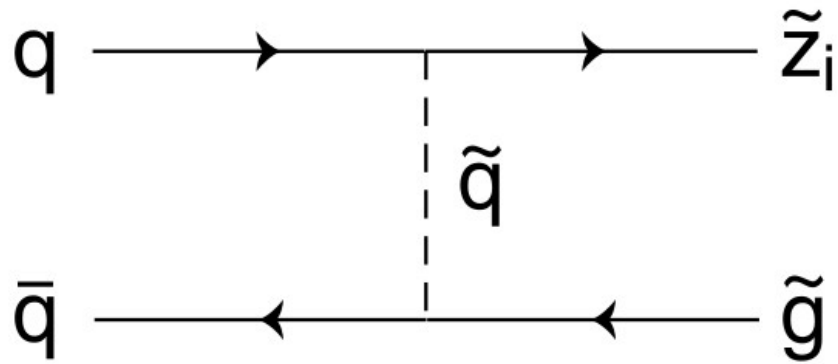
# Tree level SUSY production diagrams: gluino-squark

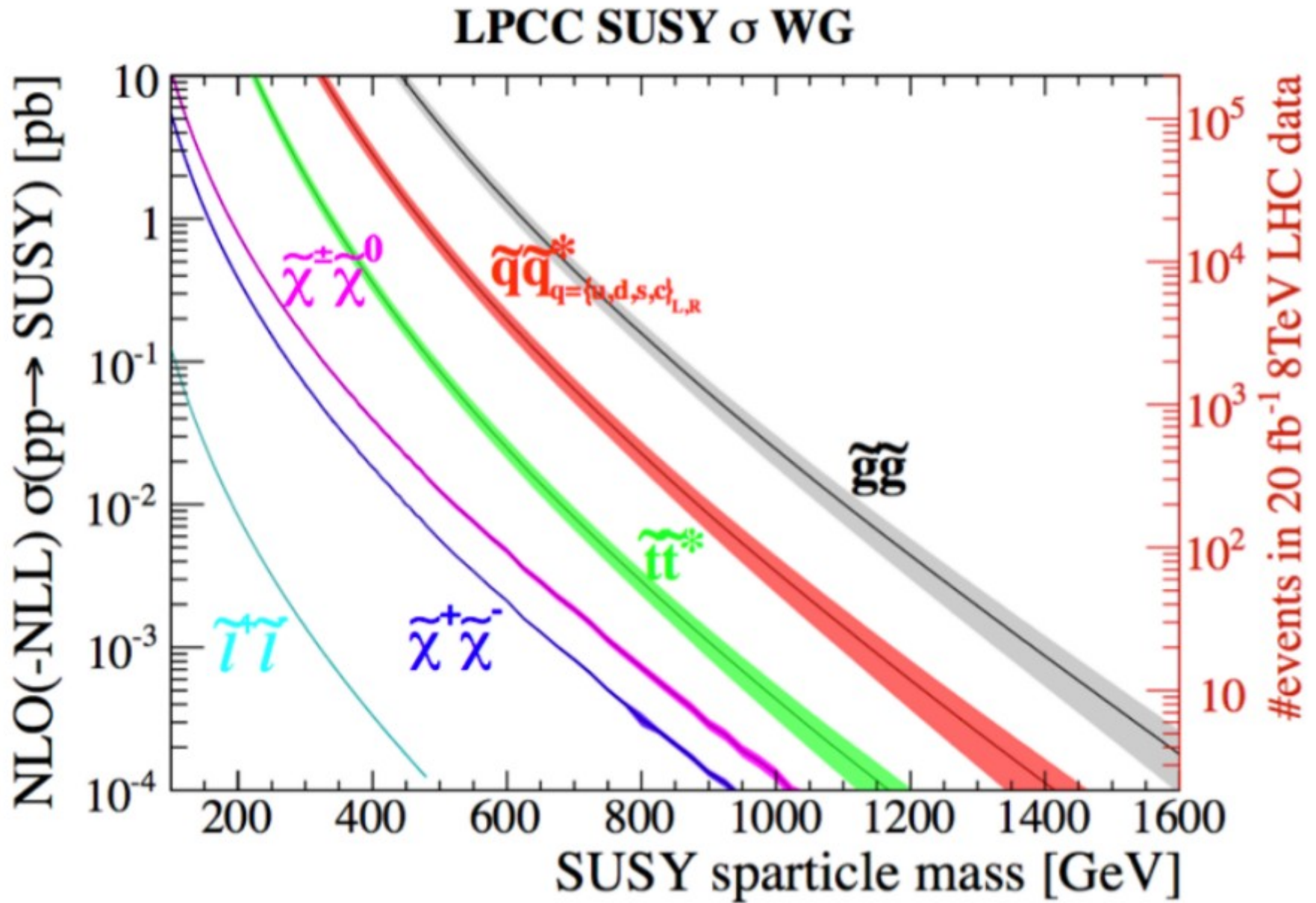


# Tree level SUSY production diagrams: squark-neutralino/chargino



# Tree level SUSY production diagrams: gluino-neutralino/chargino





# Summary of lecture 1

- We have no idea exactly what SUSY will look like at the LHC
  - naturalness is nice, but doesn't give us many clues
  - the Higgs boson discovery complicates matters
  - null observation so far is perfectly consistent with both SUSY and the Higgs mass
- Lots of potential options to look for:
  - high multiplicities
  - lots of missing energy
- LHC detectors give us:
  - particle ID and four vectors for jets, leptons and photons
  - tau and b tagging
  - missing transverse energy

***Next lecture: How do we use this information to discover SUSY?!***