



SUSY 2016

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The University of Melbourne

Radiative b -hadron decays at LHCb

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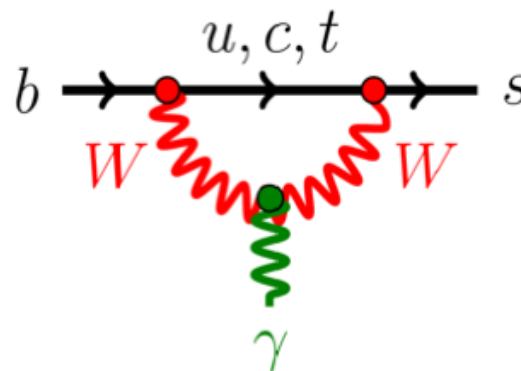
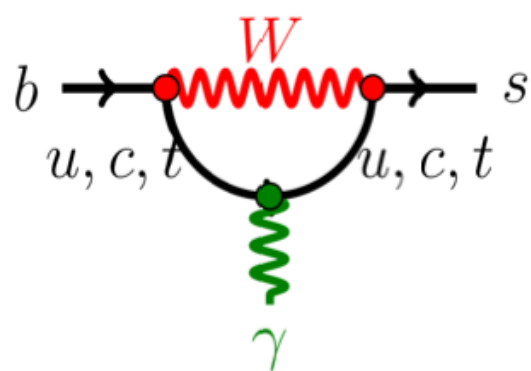
UBP/IN2P3/CNRS

On behalf of the LHCb collaboration

- Introduction
- Run1 achievements
 - BR & A_{CP} in $B \rightarrow V\gamma$ modes
 - Photon polarisation
- Prospective
- Conclusion

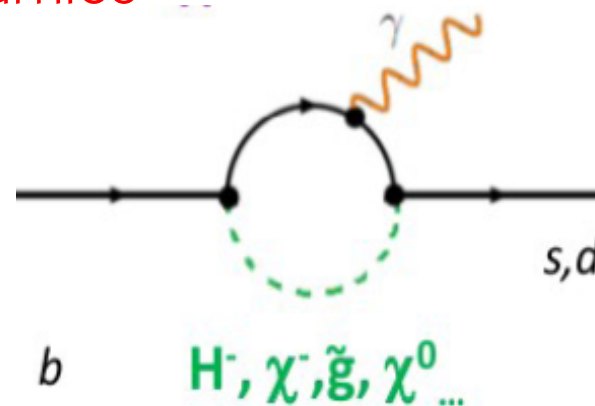
Radiative b transitions

$b \rightarrow q \gamma$ ($q=d,s$) transition : FCNC electro-magnetic penguin



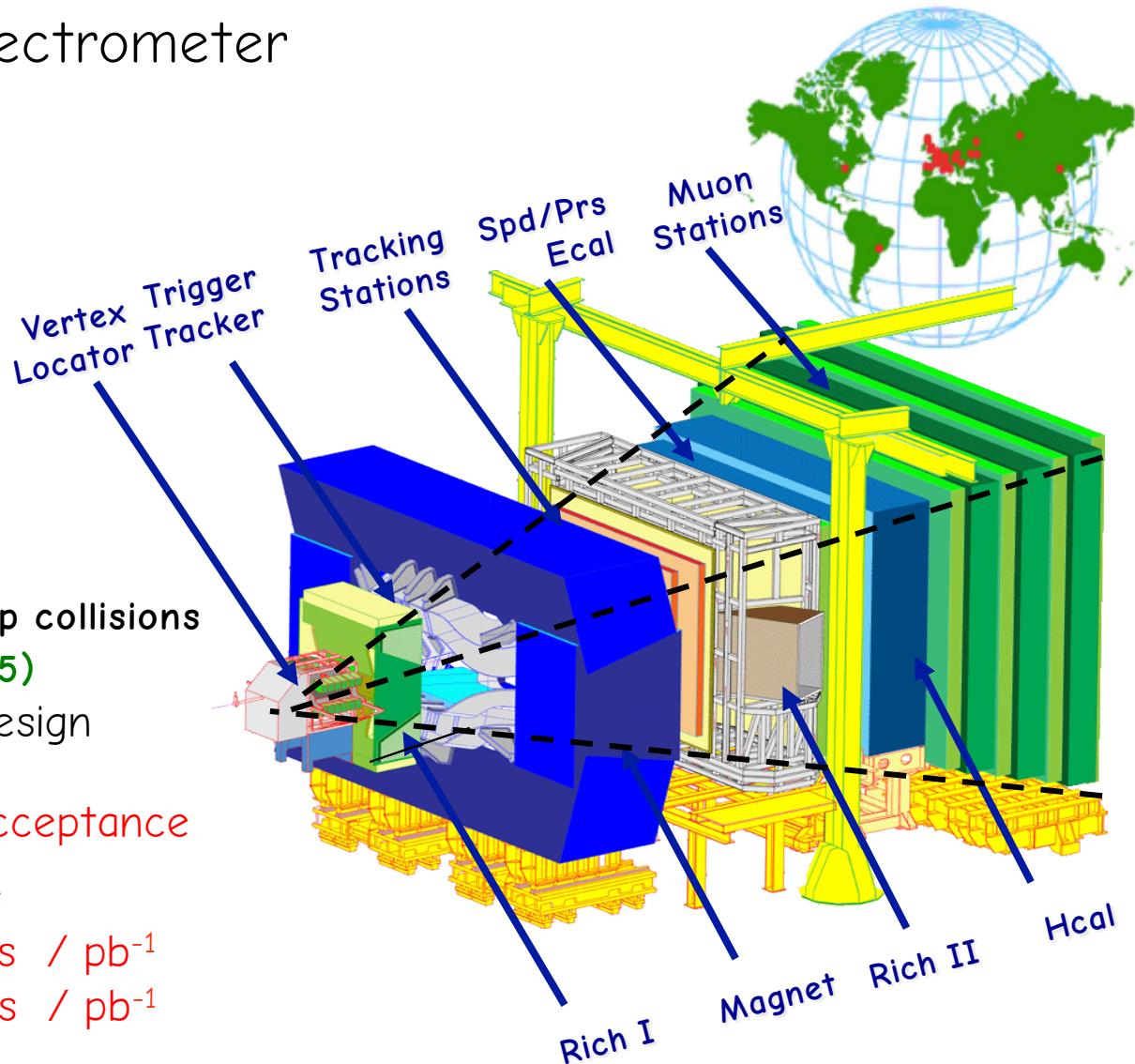
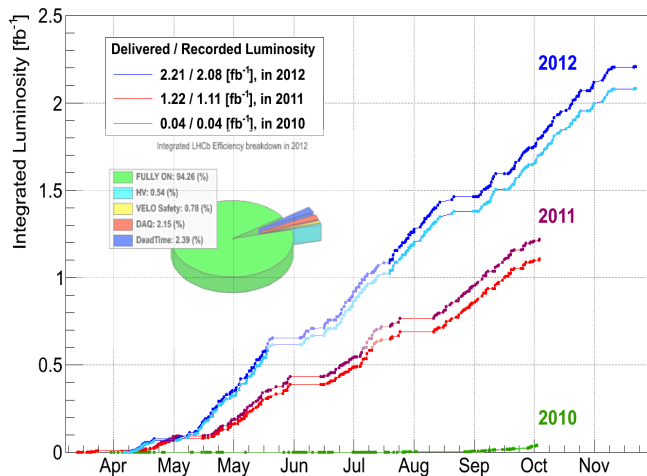
New physics affects the transition dynamics

BR, A_{CP} , Isospin asymmetry,
helicity structure of the photon



The LHCb experiment

Single-arm forward spectrometer



- RUN 1 (2010-2013): 7/8 TeV pp collisions**

Visible pp interaction/crossing: **0(1.5)**

factor 4 beyond the design

Integrated luminosity: **3fb⁻¹**

$\sim 2 \times 10^{11}$ bb in LHCb acceptance

Radiative decay reconstruction rate

$B^0 \rightarrow K^{*0} \gamma$: ~ 7 events / pb^{-1}

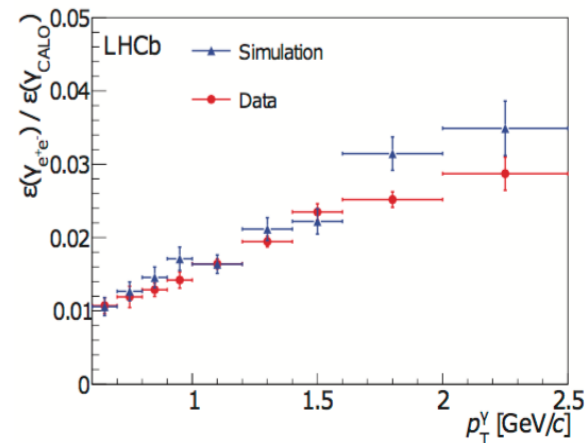
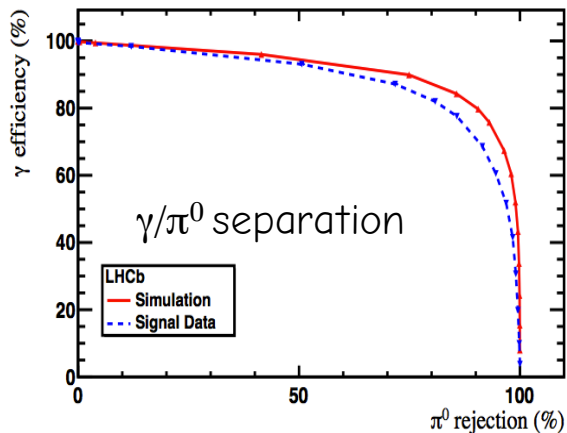
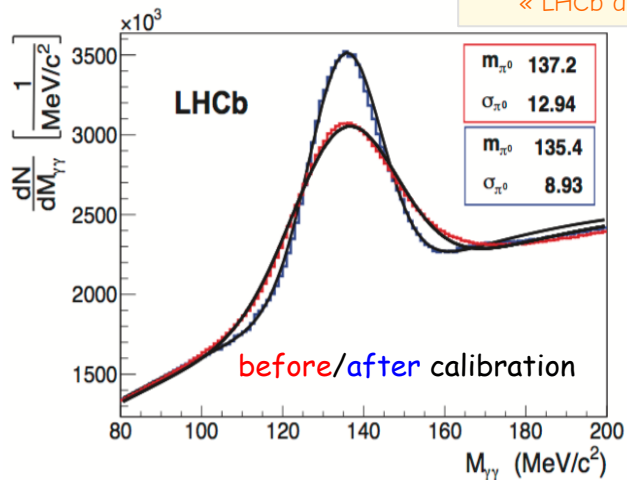
$B_s \rightarrow \phi \gamma$: ~ 1 events / pb^{-1}

- RUN 2 (2015-2018): 13 TeV pp collisions**

Photon reconstruction @ LHCb

- Calorimetric photons: **unconverted photons or conversion after magnet**
=> from calorimeters deposit
- Di-electron photons: **conversion before magnet**
=> from tracking system
- Large calorimeter occupancy : **large combinatorial background**
=> neutralID to separate neutral EM showers from hadronic and electrons deposits
- Above $p_T \sim 2.5 \text{ GeV}/c$ π^0 likely produced a single Ecal cluster
=> those π^0 represents an important background to high energy photons
=> γ/π^0 separation multivariate

« LHCb detector performances », IJMP A. Vol.30, No 7 (2015) 1530022



Radiative decay anatomy

Due to trigger constraint and large combinatorics the radiative decays mostly rely on high p_T photons

L0 threshold in 2011(2012) : $E_T(\gamma) > 2.5$ (3.0) GeV

Typical trigger efficiency on radiative modes ~ 30 -40%

For comparison : (di)muon channel $\epsilon_{\text{trg}} \sim 80$ -90%

Mass resolution driven by calorimeter resolution :

$$\sigma_M(B \rightarrow X \gamma) \sim 90 \text{ MeV}/c^2$$

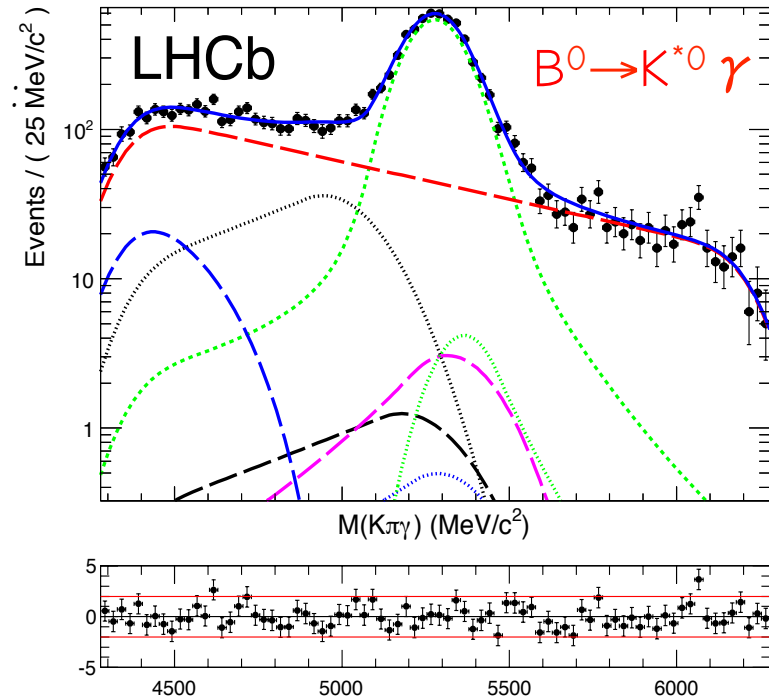
For comparison : $\sigma_M(B \rightarrow hh) \sim 25 \text{ MeV}/c^2$

$$\sigma_M(B \rightarrow J/\psi X) \leq 10 \text{ MeV}/c^2$$

No constraint on vertexing from γ / large photon multiplicity / limited mass resolution:

Large combinatorial background
partially rec'ed and peaking backgrounds

Tight selections are applied



- Generic background contamination :
 - Combinatorial background
 - Partially reconstructed $b \rightarrow s \gamma$ decays
 - Partially reconstructed $b \rightarrow c$ ($X + h h \pi^0$)
- Specific peaking backgrounds :
 - Charmless $B_{d,s} \rightarrow h^+ h^- \pi^0$
 - Irreducible $b \rightarrow d \gamma$: $B_s \rightarrow K^{*0} \gamma$
 - b-baryons cross-feed $\Lambda_b \rightarrow \Lambda^*(K-p) \gamma$



Run 1 achievements

- Do checklist
 - Post Checklist
 - Cross off Checklist
 - Something Else
-

B → V γ measurements

$B_s \rightarrow \phi \gamma$ branching fraction

$$\text{BR}(B^0 \rightarrow K^{*0} \gamma) = (4.33 \pm 0.15) \cdot 10^{-5} \quad [\text{Belle, Babar, Cleo}]$$

$$\text{BR}(B_s \rightarrow \phi \gamma) = (5.7^{+2.1}_{-1.8}) \cdot 10^{-5} \quad [\text{Belle}]$$

SM-predictions

large hadronic uncertainty mostly canceling in the ratio :

[Ali, Pecjak, Greub, 2008]

$$\text{BR}(B^0 \rightarrow K^{*0} \gamma) / \text{BR}(B_s \rightarrow \phi \gamma) = 1.0 \pm 0.2$$

LHCb result (1.0 fb⁻¹ - 2011 data)

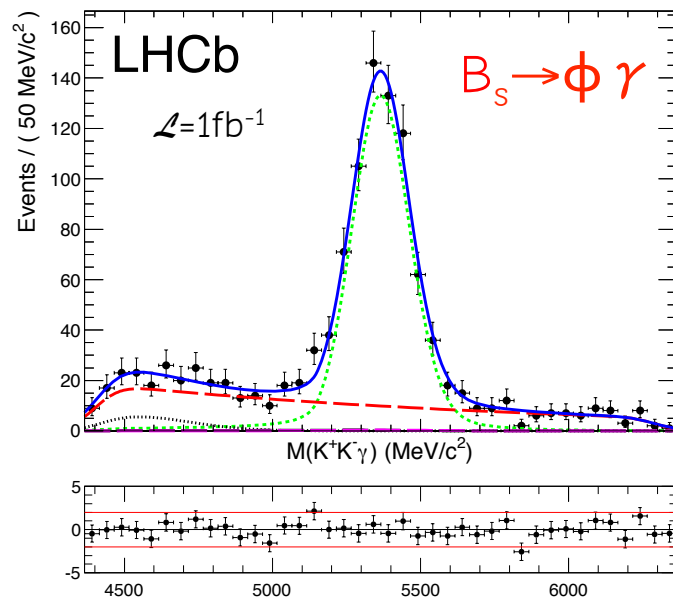
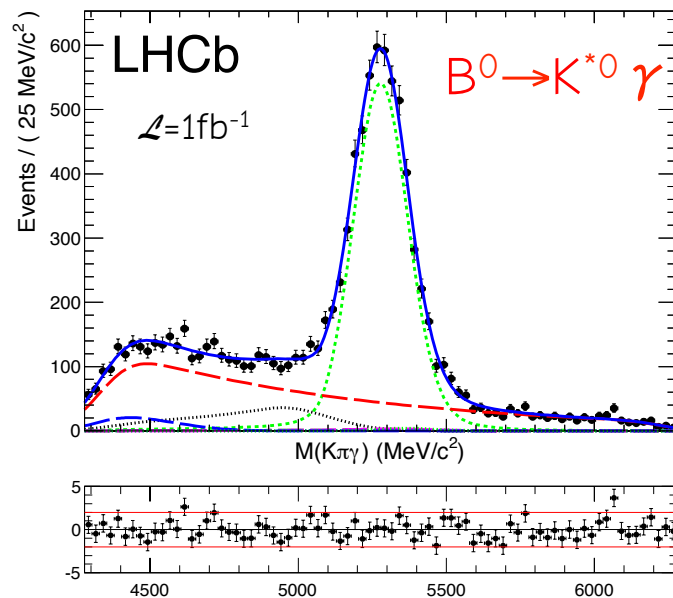
[Nuclear Physics B, 867, 1-18 (2013)]

$$\frac{\mathcal{B}(B^0 \rightarrow K^{*0} \gamma)}{\mathcal{B}(B_s^0 \rightarrow \phi \gamma)} = 1.23 \pm 0.06 \text{ (stat.)} \pm 0.04 \text{ (syst.)} \pm 0.10 \text{ (} f_s/f_d \text{)}$$

$$\mathcal{B}(B_s^0 \rightarrow \phi \gamma) = (3.5 \pm 0.4) \times 10^{-5}$$

Main systematics

- dominated by f_s/f_d : 8%
- reconstruction & selection : 2%
- background model : 2%



B → V γ measurements

Direct CP asymmetry in $B^0 \rightarrow K^{*0} \gamma$

SM-prediction :

$$A_{CP} = -0.0061 \pm 0.0043$$

Phys. Rev. D72 (2005) 014013

A_{CP} enhanced in NP scenarii

B-factory measurement

$$A_{CP} = -0.016 \pm 0.022 \pm 0.007$$

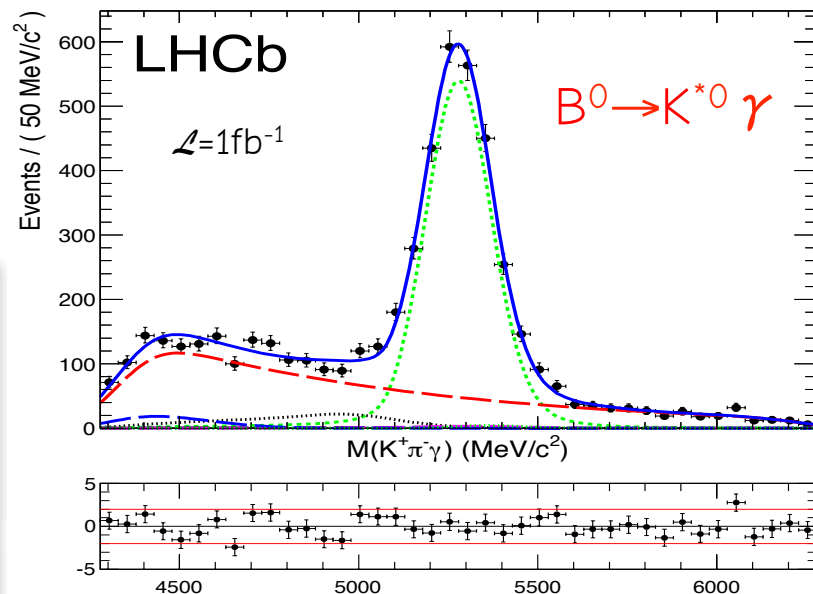
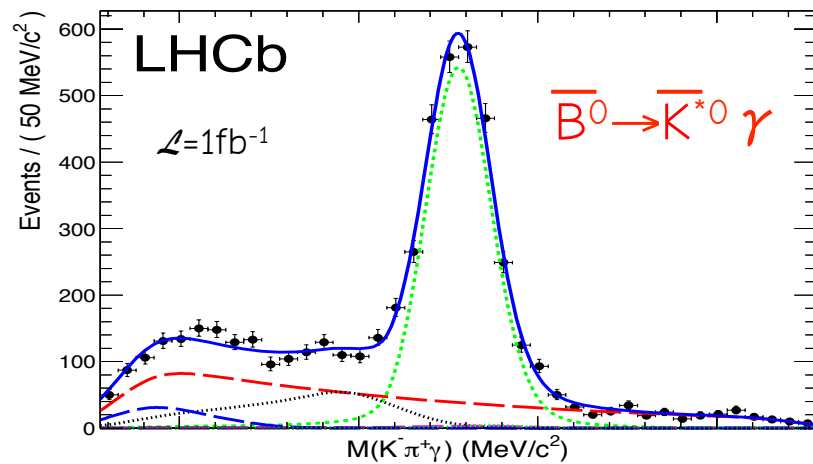
BABAR, Phys. Rev. Lett. 84, 5283-5287

LHCb result (1.0 fb⁻¹ - 2011 data) :

$$N_{B^0} + N_{\bar{B}^0} = 5300 \pm 100$$

Nuclear Physics B, 867, 1-18 (2013)

$$A_{CP}(B^0 \rightarrow K^{*0} \gamma) = 0.008 \pm 0.017(\text{stat}) \pm 0.009(\text{syst})$$



$B \rightarrow V \gamma$ measurements

Search for $B^0 \rightarrow J/\psi \gamma$ & $B_s \rightarrow J/\psi \gamma$

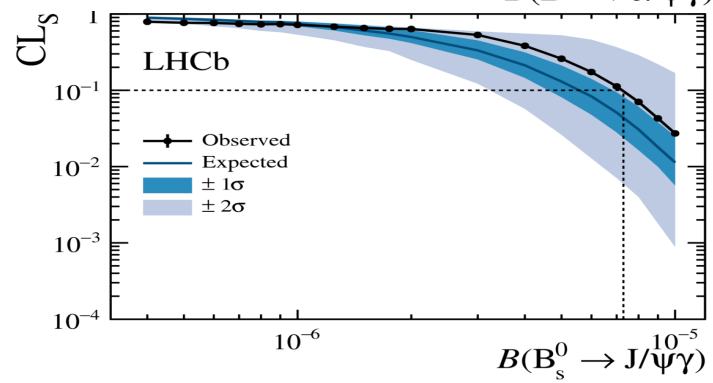
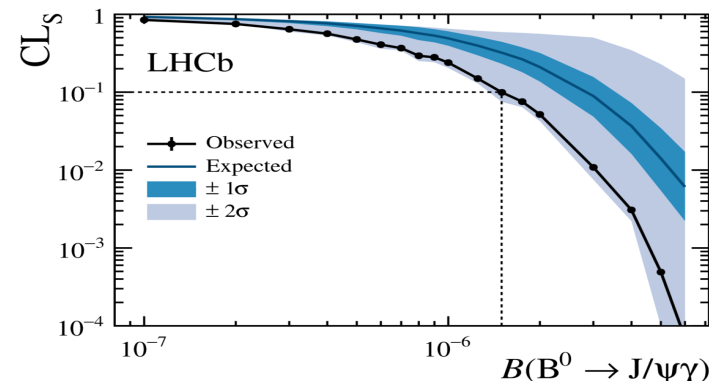
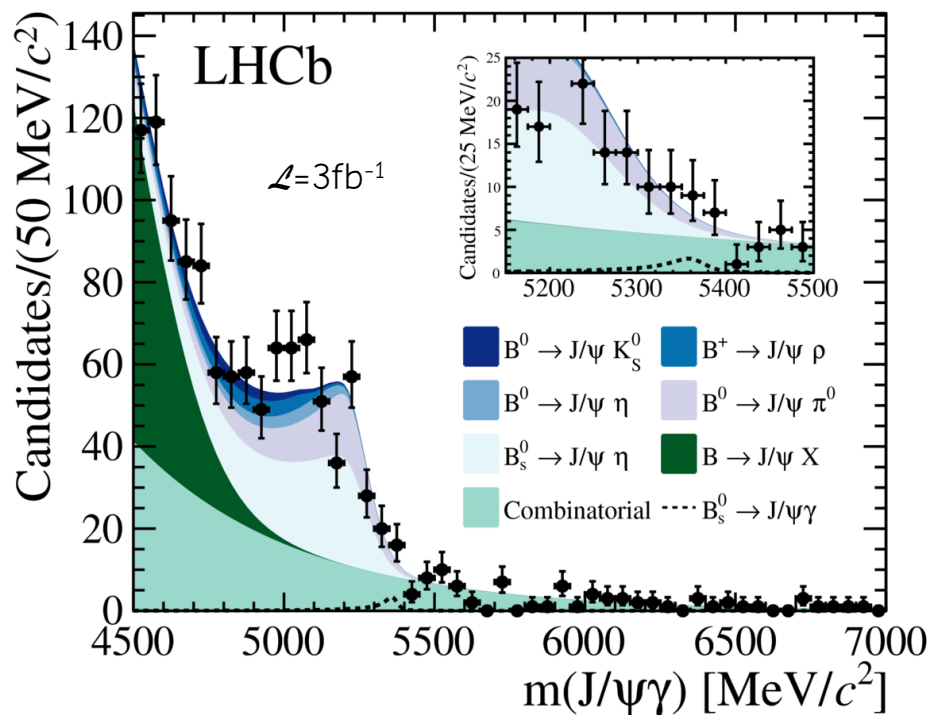
Not a radiative penguin transitions but share the same final-state problematics

PHYS. REV. D92 (2015) 112002

$$\text{BR}(B^0 \rightarrow J/\psi \gamma) < 1.7 \times 10^{-6}$$

$$\text{BR}(B_s \rightarrow J/\psi \gamma) < 7.4 \times 10^{-6}$$

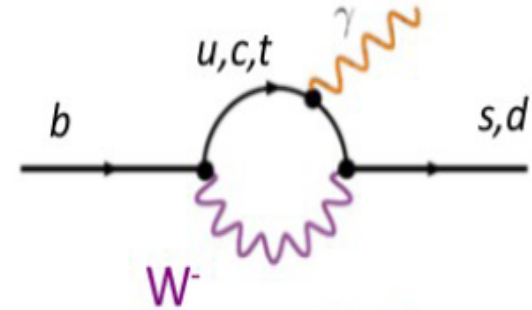
@ CL=90%



Photon polarisation

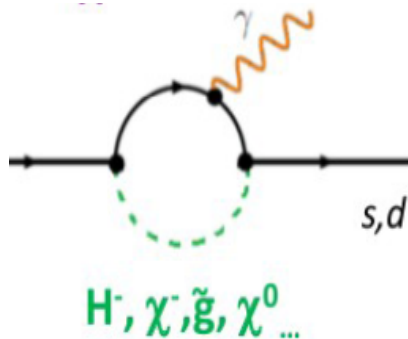
Due to the V-A structure of the electroweak interaction the photon is mostly left-handed in the radiative transition of the b-quark

Leading (EM dipole) operator in the effective Hamiltonian approach :



$$O_7 \propto m_b \bar{s} \sigma^{\mu\nu} F_{\mu\nu} (1 + \gamma_5) b + m_s \bar{s} \sigma^{\mu\nu} F_{\mu\nu} (1 - \gamma_5) b$$

$$\tan \psi = \left| \frac{A_L(b_L \rightarrow s_R \gamma_R)}{A_R(b_R \rightarrow s_L \gamma_L)} \right| \approx m_s / m_b$$

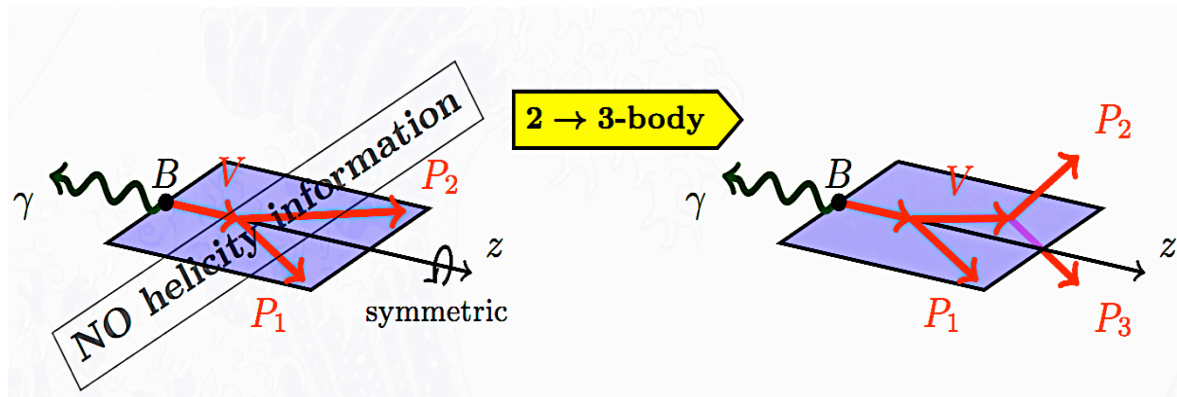


- Right-handed component could be enhanced in NP models

Photon polarisation

Experimentally, the photon polarization can be extracted from ...

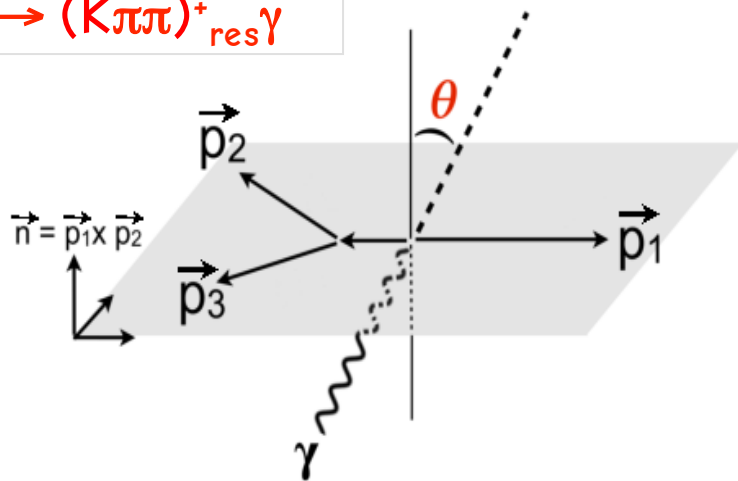
- angular analysis of the recoil 3-body in the $B \rightarrow \gamma + (hhh)_{res.}$ decay mode



- time-dependent analysis of the $B \rightarrow \gamma + \Phi_{CP}$ decay modes
- di-lepton angular analysis at low q^2 of the (virtual) photon decay in $B \rightarrow V e^+ e^-$
- angular analysis in the radiative transition of b-baryons

Helicity structure in $B \rightarrow (\text{K}\pi\pi)_{\text{res}} \gamma$

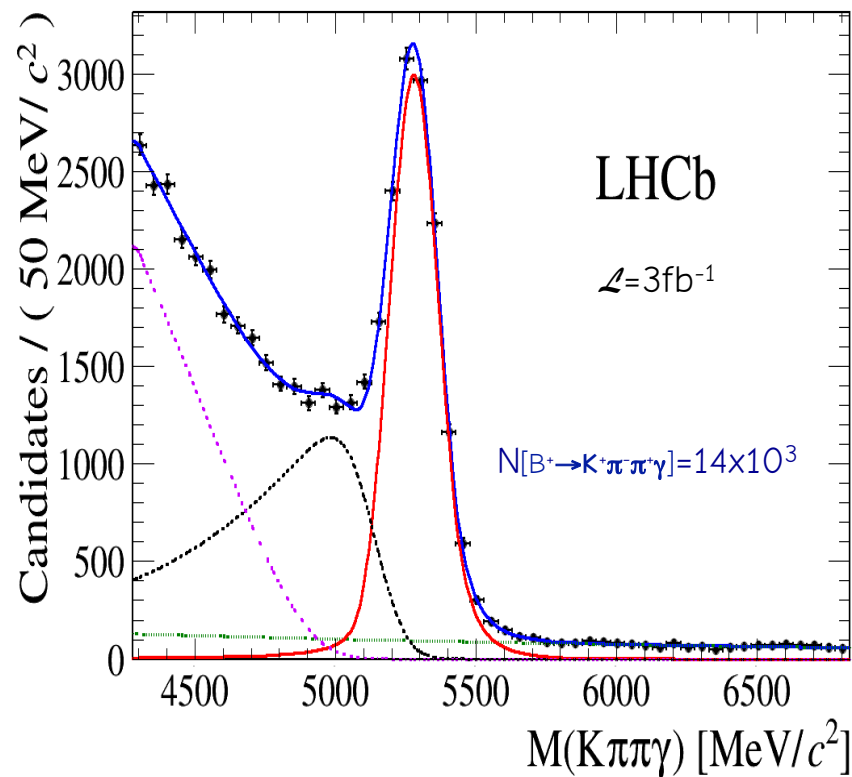
$$B^+ \rightarrow (\text{K}\pi\pi)^+_{\text{res}} \gamma$$



For a mixture of spin-parity $\text{K}_{\text{res}}(1^+, 2^+, 1^-)$:

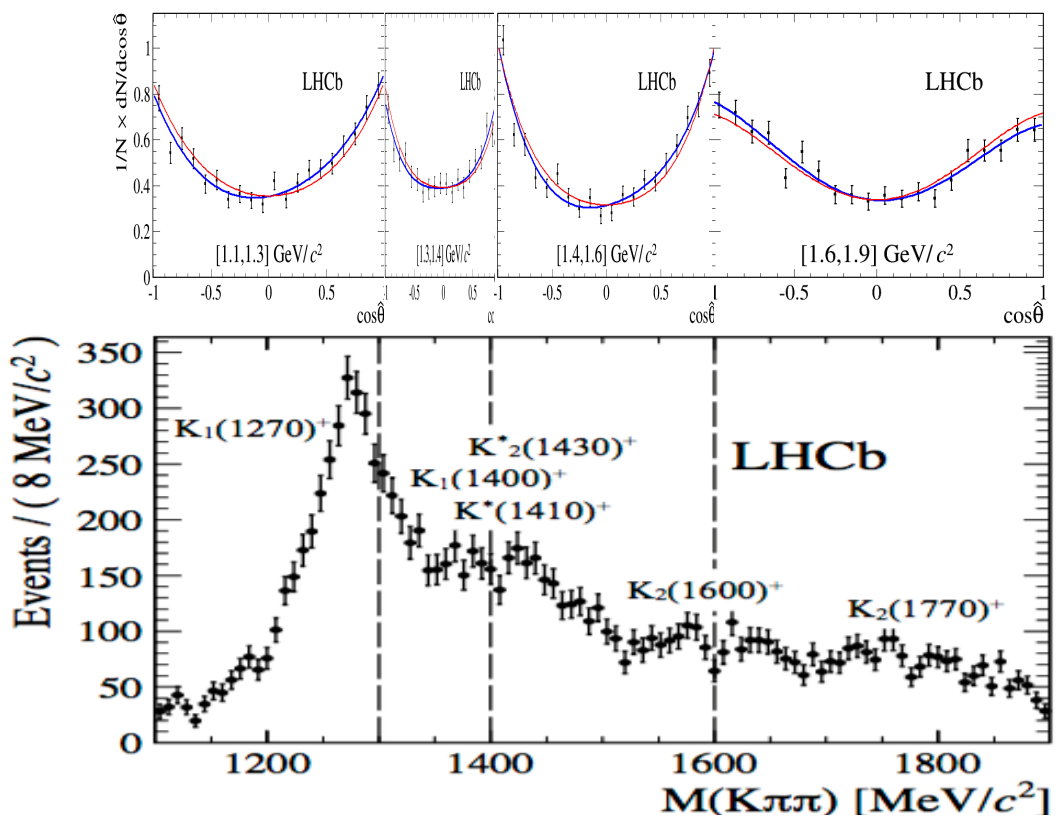
$$\frac{d\Gamma}{ds ds_{13} ds_{23} d \cos \theta} \propto$$

$$\sum_{i=0,2,4} a_i(s, s_{13}, s_{23}) \cos^i \theta + \lambda_\gamma \sum_{j=1,3} a_j(s, s_{13}, s_{23}) \cos^j \theta$$



Up-down photon asymmetry is proportional to the photon polarisation λ_γ
 Angular analysis of photon direction wrt to $(\text{K}\pi\pi)_{\text{res}}$ decay plane in different mass bins

Helicity structure in $B \rightarrow (K\pi\pi)_{res} \gamma$



[Phys. Rev. Lett. 112, 161801 (2014)]

Photons from radiative decays are polarized @ 5.2σ significance

First direct observation of photon polarization in $b \rightarrow s \gamma$ transition

as a by-product:

LHCb-CONF-2013-009

$$\mathcal{A}_{CP} = -0.007 \pm 0.015 \text{ (stat)}_{-0.011}^{+0.012} \text{ (syst)}$$

Measuring the λ_γ value from the up-down asymmetry require to separate the $(K\pi\pi)$ resonances & theoretical determination of the helicity amplitude

Phys. Rev. Lett. 88 (2002) 051802

e.g. for a single 1+ resonance

$$\frac{d\Gamma(B \rightarrow K\pi\pi\gamma)}{ds ds_{13} ds_{23} d\cos\theta} \propto \frac{1}{2} |\vec{\mathcal{J}}|^2 (1 + \cos^2\theta) + \lambda_\gamma \cos\theta \text{Im}[\vec{n} \cdot (\vec{\mathcal{J}} \times \vec{\mathcal{J}}^*)]$$

Virtual photon : $B^0 \rightarrow K^{*0} e^+ e^-$

$B^0 \rightarrow K^*(\gamma^* \rightarrow ee)$ in the low q^2 region

Branching fraction in $[30 ; 1000] \text{ MeV}/c^2$
 1.0 fb^{-1} - 2011 data :

J. High Energy Phys. 05 (2013) 159

$$\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)_{30-1000 \text{ MeV}/c^2} = (3.1^{+0.9}_{-0.8} \text{ }^{+0.2}_{-0.3} \pm 0.2) \times 10^{-7}$$

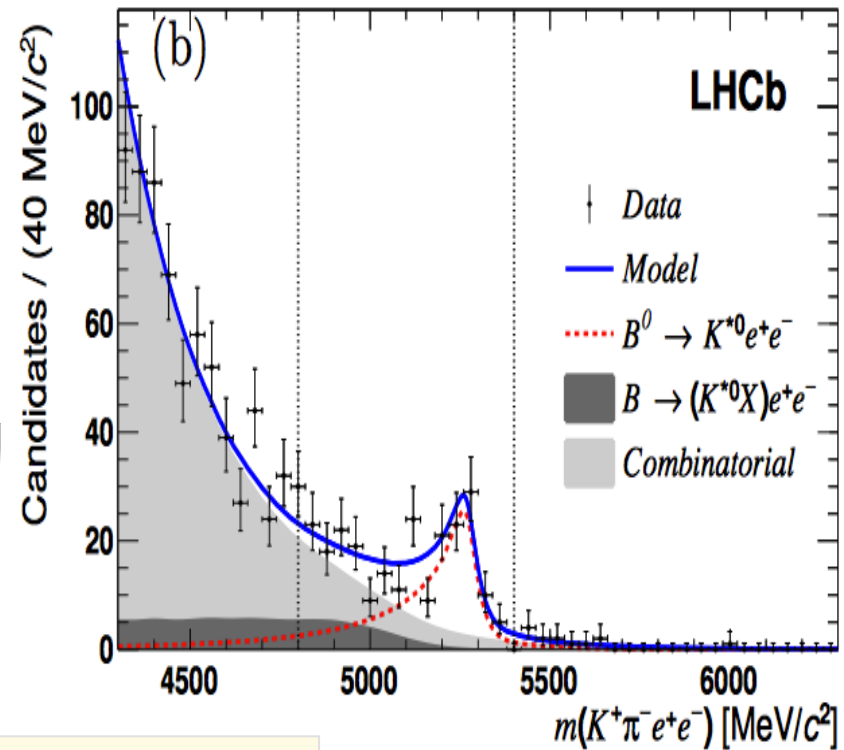
Full angular analysis in $[20; 1120] \text{ MeV}/c^2$
 3.0 fb^{-1} - 2011+2012 data :

$$F_L = 0.16 \pm 0.06 \pm 0.03$$

$$A_T^{(2)} = -0.23 \pm 0.23 \pm 0.05$$

$$A_T^{\text{Im}} = +0.14 \pm 0.22 \pm 0.05$$

$$A_T^{\text{Re}} = +0.10 \pm 0.18 \pm 0.05,$$



J. High Energy Phys. 04(2015) 064

Related to the
 photon polarisation

Time-dependent decay rate

Direct access to the polarization via the time-dependent decay rate of $B \rightarrow \Phi^{CP} \gamma$

$$\Gamma_{B(\bar{B})_{(s)}^0 \rightarrow \Phi^{CP} \gamma}(t) = |A|^2 e^{-\Gamma_{(s)} t} \left(\underbrace{\cosh(\Delta\Gamma_{(s)} t/2) + A_{\Delta} \sinh(\Delta\Gamma_{(s)} t/2)}_{\text{untagged}} \right) \bigg/ \left(\underbrace{\pm C_{CP} \cos(\Delta m_{(s)} t) \mp S_{CP} \sin(\Delta m_{(s)} t)}_{\text{tagged analysis required}} \right)$$

$$\begin{aligned} S_{CP} &\sim \sin 2\psi \sin \phi_{(s)} \\ A_{\Delta} &\sim \sin 2\psi \cos \phi_{(s)} \end{aligned} \quad \tan \psi = \left| \frac{A_R}{A_L} \right| \quad \phi_{(s)} : B_{(s)}^0 \text{ mixing phase}$$

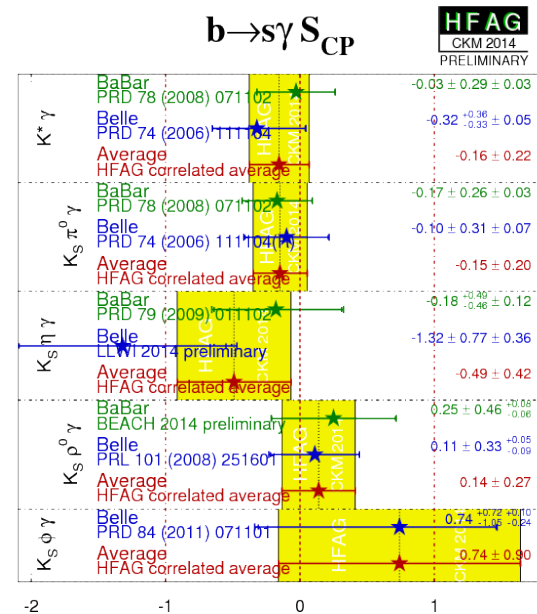
- B^0 decays : access to the polarisation through the TD asymmetry term S_{CP}
Out of LHCb reach for the main decay mode $B^0 \rightarrow K_S \pi^0 \gamma$

- B_s decays : access through the mixing term A_{Δ}

Muheim et al., PLB664(08)17

SM: $A_{\Delta} = 0.047 \pm 0.025 \pm 0.015$
 Left-Right Symmetric model: $A_{LRSM}^{\Delta} \sim 0.7$

LHCb : ongoing untagged analysis of the $B_s \rightarrow \phi \gamma$ decay rate
 Expected statistical resolution : $\sigma_{A_{\Delta}} \sim 0.4$
 Measurement statistically limited (3×10^3 rec'ed B_s)





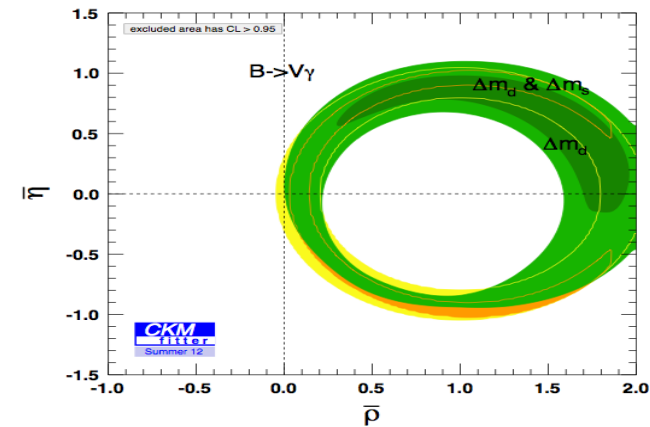
The expectedly large run2 statistics will allow to explore suppressed radiative modes

- e.g. V_{td} suppressed $b \rightarrow d\gamma$ penguin

Branching ratio & asymmetry of exclusive $b \rightarrow (d + s)\gamma$ modes provide a direct constraint on UT

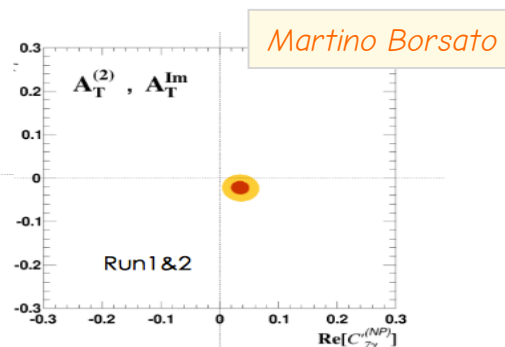
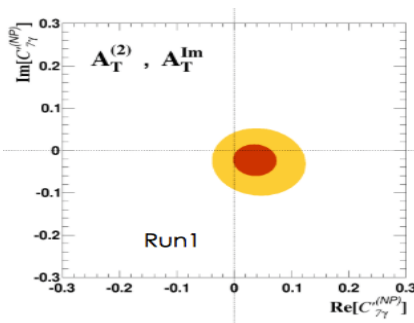
Such transition could be accessible in LHCb via $B^0 \rightarrow \rho^0/\omega \gamma$, $B^+ \rightarrow a_1^+ \gamma$

Could separate the $b \rightarrow d\gamma$ transition $B_s \rightarrow K^* \gamma$ from $b \rightarrow s\gamma$ in B^0 using converted photons



- Photon polarisation : reach < 10 % resolution

scenario II: $C_{7\gamma}^{(NP)} = 0$, $C_{7\gamma}'^{(NP)} \in \mathbb{C}$



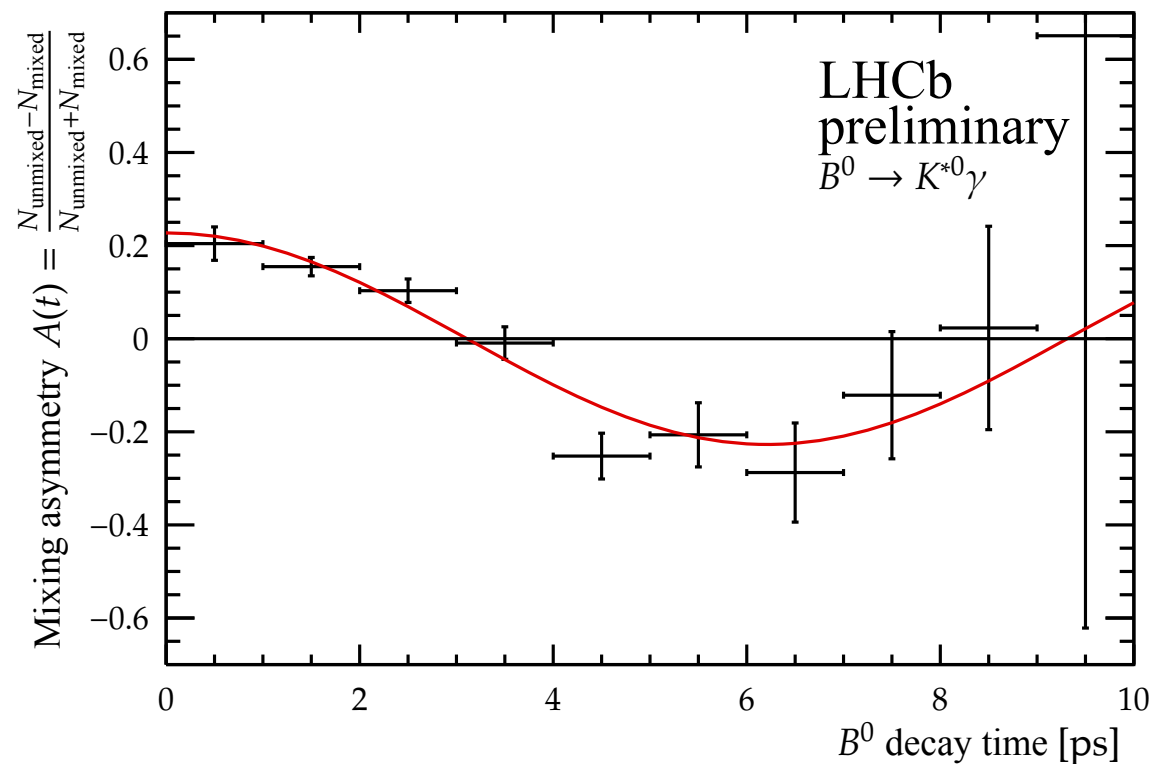
Prospective from $K^* ee$ angular analysis

- Photon polarisation (cont')

Enhanced sensitivity in neutral B decays from tagged analysis giving access to TD asymmetries

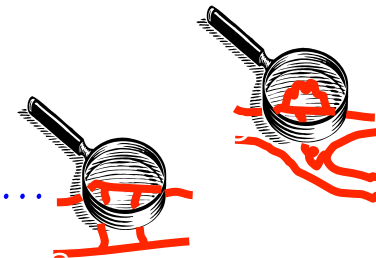
For illustration :

LHCb tagging performance applied for the first time on a radiative decay



LHCb provides an unique laboratory for precise measurements in radiative decay

Great harvest of result with 2011-2012 Run 1
World best measurements in radiative $V\gamma$ decays ...



... consistent with SM expectation



Many updated or new results expected soon
run 2 will allow to explore rarest radiative decays



$B_s \rightarrow \phi \gamma$ Branching Fraction

- Systematic uncertainty dominated by f_s/f_d ($\pm 8\%$)

[Phys. Rev. D 85 (2012) 032008]

from semi-leptonic $B_{u,d,s} \rightarrow D_{(s)} \mu \nu X$ and hadronic $B_{u,d,s} \rightarrow D_{(s)} h$

$$\frac{f_s}{f_d} = 0.267^{+0.021}_{-0.020}$$

- Background model ($\pm 2\%$)

Contamination level and shape

- Reconstruction and selection ($\pm 2\%$)

Trigger and selection efficiencies, Particle reconstruction & identification

Update with whole 3fb^{-1} sample ongoing

both statistical and systematical uncertainty will improve
(more precise f_s/f_d , improved background model ...)

- $K^+ \pi^- / K^- \pi^+$ detection asymmetry

From charm $D^0 \rightarrow K \pi$ large control sample

$$A_D(K\pi) = \frac{\varepsilon(K^- \pi^+) - \varepsilon(K^+ \pi^-)}{\varepsilon(K^- \pi^+) + \varepsilon(K^+ \pi^-)} = (-1.0 \pm 0.2)\%$$

- B production asymmetry

From large $B \rightarrow J/\psi K^$ sample*

$$A_p(B) = \frac{R(\bar{B}) - R(B)}{R(\bar{B}) + R(B)} = (1.0 \pm 1.3)\%$$

LHCb-CONF-2011-042.

- Background model

$$\Delta A_{CP} = (-0.2 \pm 0.7)\%$$

Contamination level, shape & CP asymmetry in various background components

Dominated by the unknown asymmetry from the misidentified $\Lambda_b \rightarrow (pK)\gamma$ contamination

- Detector non-uniformity

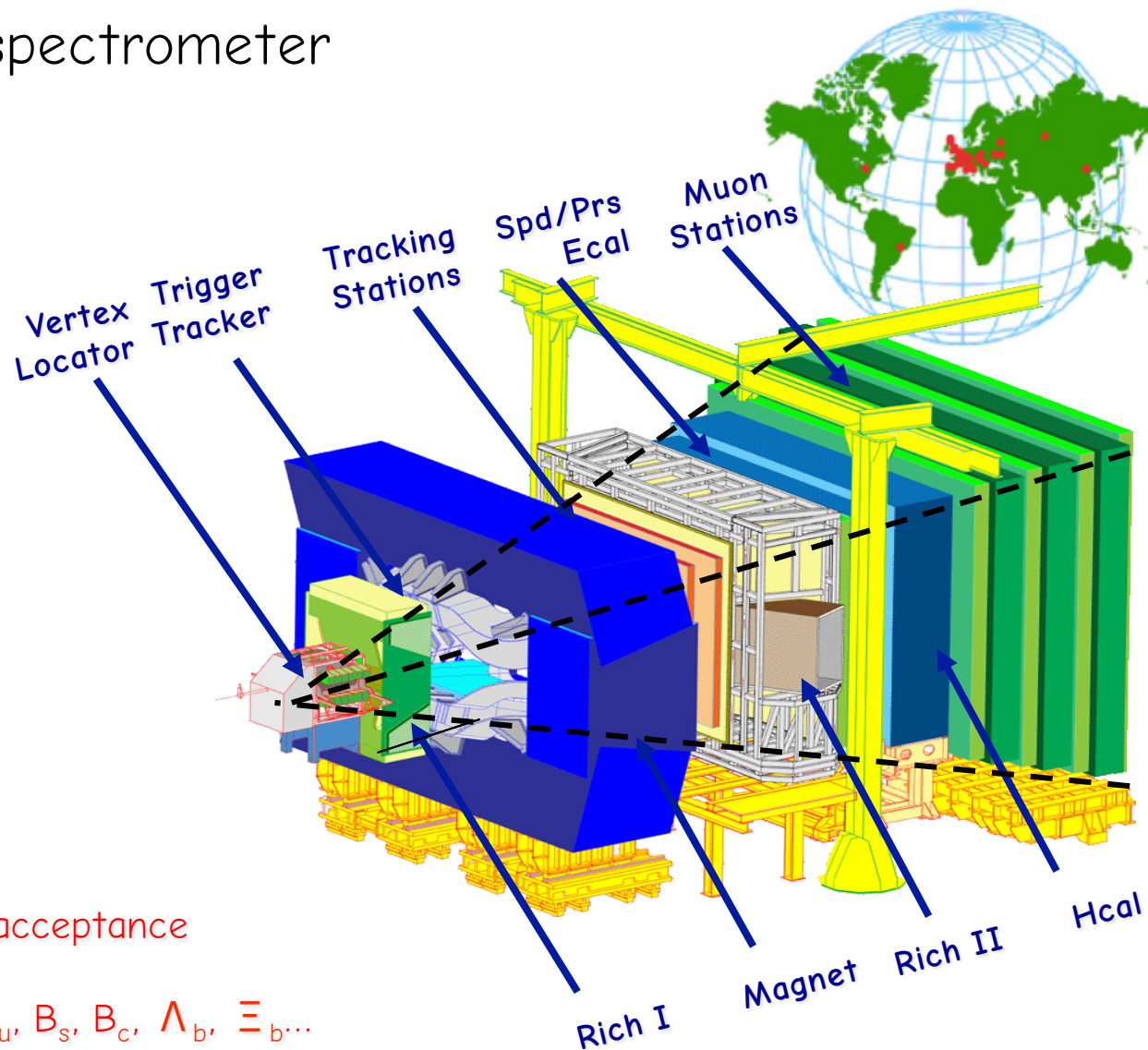
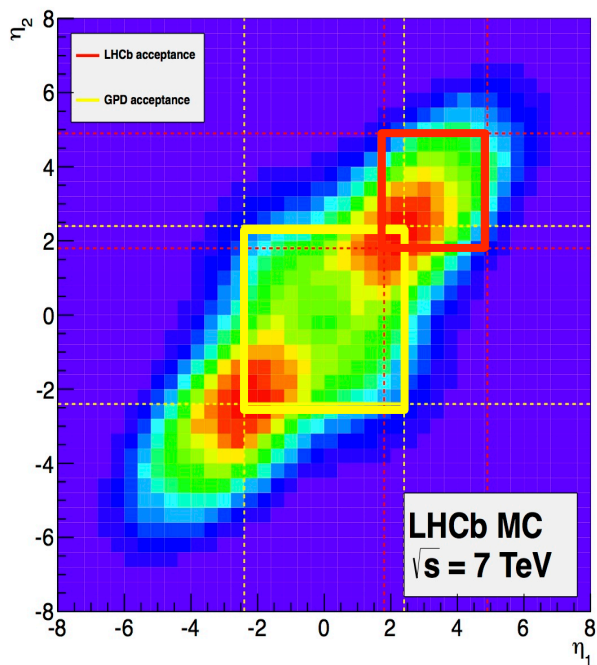
$$\Delta A_{CP} = (+0.1 \pm 0.2)\%$$

Possible detector bias strongly reduced by switching regularly the magnet polarity

Update with whole 3fb^{-1} sample ongoing

both statistical and systematical uncertainty will improve :
 more precise detection and production asymmetry,
 CP asymmetry from background in particular $\Lambda_b \rightarrow (pK)\gamma$

Single-arm forward spectrometer



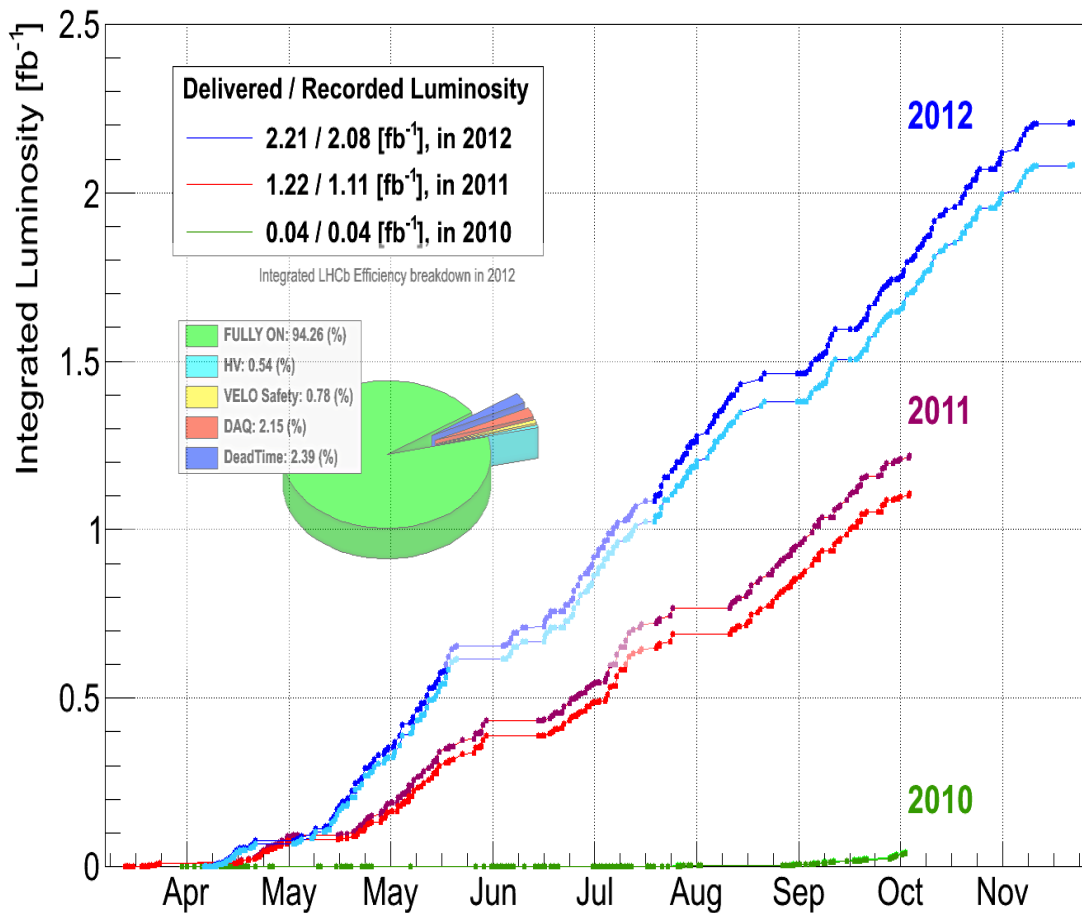
Large $b\bar{b}$ rate in the detector acceptance

Access to all b species : $B_d, B_u, B_s, B_c, \Lambda_b, \Xi_b, \dots$

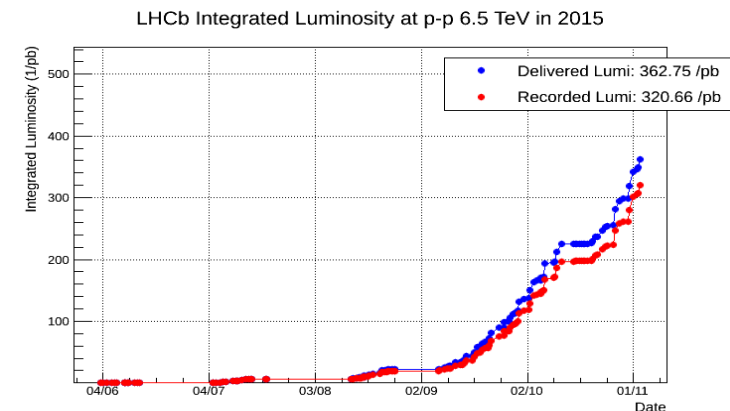
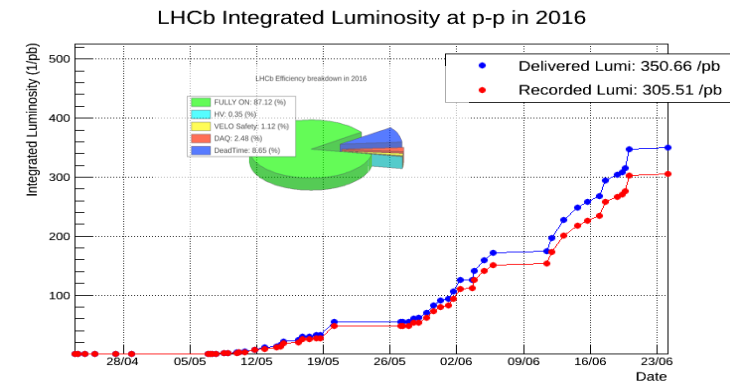
LHCb operations

- RUN 1 (2010-2013) : $\sqrt{s}=7$ TeV and $\sqrt{s}=8$ TeV pp collisions

Instantaneous luminosity : $4 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$ factor 2 beyond the design
 Visible pp interaction/crossing : $0(1.5)$ factor 4 beyond the design
 Integrated luminosity : 3 fb^{-1} $\sim 2 \times 10^{11} \text{ bb}$ in LHCb acceptance



- RUN 2 (2015-2018) $\sqrt{s}=13$ TeV



Because of trigger constraints and the hadronic environment, some radiative channels are out of reach for LHCb

- What LHCb can do (or has already done) :

- $(h^+h^-)\gamma$ channels from $b \rightarrow s\gamma$ transition of neutral b-hadrons : B^0, B_s, Λ_b
- $(h^+h^-h^+)\gamma$ channel from charged b-hadrons
- Same for suppressed $b \rightarrow d\gamma$ transitions when accumulating statistics

- What LHCb could probably do :

- $(h^+h^-\pi^0)\gamma, (h^+h^-K_s)\gamma$: e.g. $K_1^0\gamma$ maybe $B \rightarrow \omega\gamma$
- $B^+ \rightarrow K^{*+}(K_s\pi^+)\gamma$: e.g. isospin asymmetry in $K^*\gamma$
- $\Lambda_b \rightarrow \Lambda_{1115}\gamma$: very challenging

- What LHCb can't do

- $B^+ \rightarrow K^{*+}(K^+\pi^0)\gamma$ and a fortiori $B^+ \rightarrow \rho^+(\pi^+\pi^0)\gamma$
- $B^0 \rightarrow K^{*0}(K_s\pi^0)\gamma$: time-dependent asymmetry in radiative Bd
- inclusive $b \rightarrow X_s\gamma$ BR