

SUSY2016 @ Melbourne

Constraints on non-universal gaugino mass scenario using the latest LHC data

PTEP 2013 (2013) 013B02

PRD 93 (2016) no.5, 055019

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with
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Please see also:

J.K. and Y.Omura

“Diphoton excess at 750 GeV and LHC constraints in models with vectorlike particles”
Phys.Rev. D93 (2016) no.11, 115011

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Y.Omura (KMI, Nagoya Univ.)

little hierarchy problem

Higgs mass and SUSY search indicate high-scale SUSY

→ hierarchy between SUSY scale and EW scale

□ minimization condition of the Higgs potential

$$m_Z^2 \simeq -2 |\mu|^2 + 2|m_{h_u}^2|$$

EW scale

SUSY scale

- ✓ fine-tuning is required if $m_Z \ll \mu, m_{h_u}$
- ✓ at least $\mu(m_Z)$ must be small since it's unique SUSY parameter
- ✓ small $\mu(m_Z)$ means small $m_{H_u}(m_Z)$

Higgs mass vs little hierarchy

little hierarchy problem relates to the Higgs mass

□ RGE of m_{h_u}

$$16\pi^2 \frac{dm_{h_u}^2}{dt} \simeq 6y_t^2(m_{h_u}^2 + m_{\tilde{t}_L}^2 + m_{\tilde{t}_R}^2 + A_t^2) - 6g_2^2|M_2|^2 - \frac{6}{5}g_1^2|M_1|^2$$

- top squark parameters $m_{\tilde{t}_L}^2, m_{\tilde{t}_R}^2, A_t$ appear
- heavier top squark leads severer fine-tuning
- top squark mass is crucial for the Higgs mass

✓ 10 TeV top squark forces $10^{-3}\%$ tuning

maximal mixing

□ MSSM Higgs boson mass

$$m_h^2 \simeq m_Z^2 \cos^2 2\beta + \frac{3m_t^4}{8\pi^2 v_u^2} \left[\log \frac{M_{stop}^2}{m_t^2} + \frac{2A_t^2}{M_{stop}^2} \left(1 - \frac{A_t^2}{12M_{stop}^2} \right) \right]$$

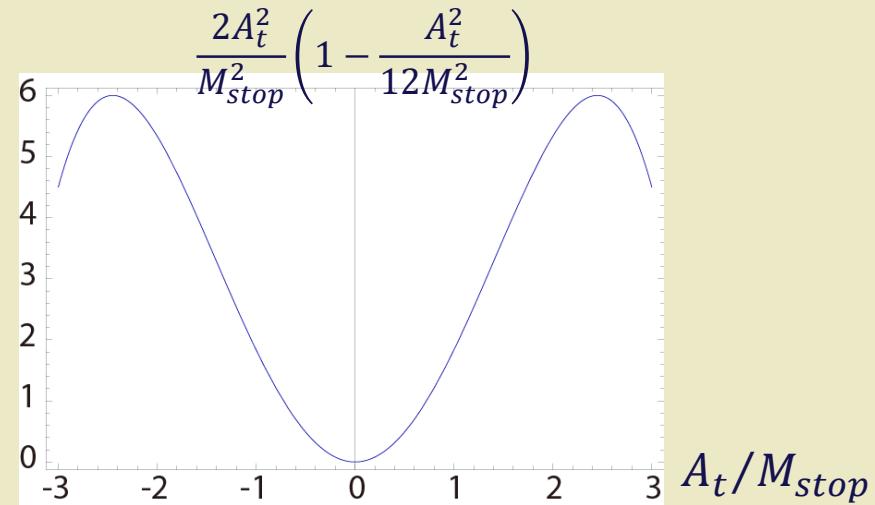
- $M_{stop} \simeq 10 \text{ TeV}$ if $A_t/M_{stop} \ll 1$
- Higgs mass requires severer bound than direct SUSY search

□ maximal mixing scenario

the last term is maximized at

$$A_t/M_{stop} \sim \sqrt{6}$$

“maximal mixing”



Higgs boson mass in NUGM

we assume universal soft masses m_0 and A-term A_0 at the GUT scale

□ top squark parameters

$$m_{\tilde{t}_L}^2(m_Z) \simeq +0.38M_2^2 + 5.63M_3^2 + 0.58m_0^2$$

$$m_{\tilde{t}_R}^2(m_Z) \simeq -0.21M_2^2 + 4.61M_3^2 + 0.19m_0^2 \quad \text{GUT scale}$$

$$A_t(m_Z) \simeq -0.21M_2 - 1.90M_3 + 0.18A_0$$

□ universal gaugino masses

$$M_{stop} \equiv \sqrt{m_{\tilde{t}_R} m_{\tilde{t}_L}}$$

$$M_2 = M_3 \gg m_0 \quad \rightarrow \quad \frac{A_t}{M_{stop}} \simeq \frac{2.11^2 \times M_3^2}{\sqrt{6.01 \cdot 4.40} \times M_3^2} \simeq 0.87$$

✓ 126 GeV Higgs boson requires heavy top squark \gtrsim sub TeV

Higgs boson mass in NUGM

□ top squark parameters

$$m_{\tilde{t}_L}^2(m_Z) \simeq +0.38M_2^2 + 5.63M_3^2 + 0.58m_0^2$$

$$m_{\tilde{t}_R}^2(m_Z) \simeq -0.21M_2^2 + 4.61M_3^2 + 0.19m_0^2 \quad \text{GUT scale}$$

$$A_t(m_Z) \simeq -0.21M_2 - 1.90M_3 + 0.18A_0$$

□ Non-Universal Gaugino Masses (NUGM)

- ✓ $m_{\tilde{t}_R}(m_Z)$ decreases, $|A_t(m_Z)|$ increases as M_2 increases

$$\rightarrow A_t/M_{stop} \lesssim \sqrt{6}$$

- ✓ upper bound is $M_2/M_3 \lesssim 5$ for $m_{\tilde{t}_R}^2(m_Z) > 0$
- ✓ In other words, A_t/M_{stop} is maximized at $M_2/M_3 \sim 5$

brief summary

- ✓ large wino mass enhances the Higgs boson mass
- ✓ A_t/M_{stop} is maximized at $M_2/M_3 \simeq 5$

What happen for naturalness ?

naturalness in NUGM

□ RG-running of m_{H_u}

$$m_{h_u}^2(m_Z) \simeq +0.17M_2^2 - 0.20M_2M_3 - 3.09M_3^2 - 0.23m_0^2 \quad \text{GUT scale}$$

$$\rightarrow M_2 \simeq 5 \times M_3 \rightarrow m_{h_u}^2(m_Z) \simeq 0$$

- ✓ μ -parameter is minimized at $M_2/M_3 \simeq 5$
- ✓ Higgs mass is also maximized at $M_2/M_3 \simeq 5$!

suitably large wino reconcile the Higgs mass and naturalness

our tuning measure

□ our tuning measure (BG-type)

$$\Delta_\mu \equiv \left| \frac{d \ln m_Z^2}{d \ln \mu(\Lambda_{GUT})^2} \right|$$

□ minimization condition of the Higgs potential

$$m_Z^2 \simeq -2 |\mu|^2 + 2|m_{h_u}^2|$$

SUSY ~~SUSY~~

- ✓ we focus on tuning between μ -parameter and SUSY breaking parameters
- ✓ μ -parameter is unique SUSY dimensionful parameter in MSSM
- ✓ we expect some relations (e.g. gaugino mass ratio) among SUSY breaking parameters since these have same origin i.e. SUSY breaking mediation

NUGM from SUSY breaking mediation

□ possibilities of NUGM (large wino)

- mixed moduli/anomaly mediation (mirage mediation)

$$M_a = \frac{F^T}{T + \bar{T}} + \frac{g_0^2}{16\pi^2} b_a \frac{F^C}{C} \quad \text{where, } b_a = \left(\frac{33}{5}, 1, -3 \right) \\ a = U(1)_Y, SU(2)_L, SU(3)_C$$

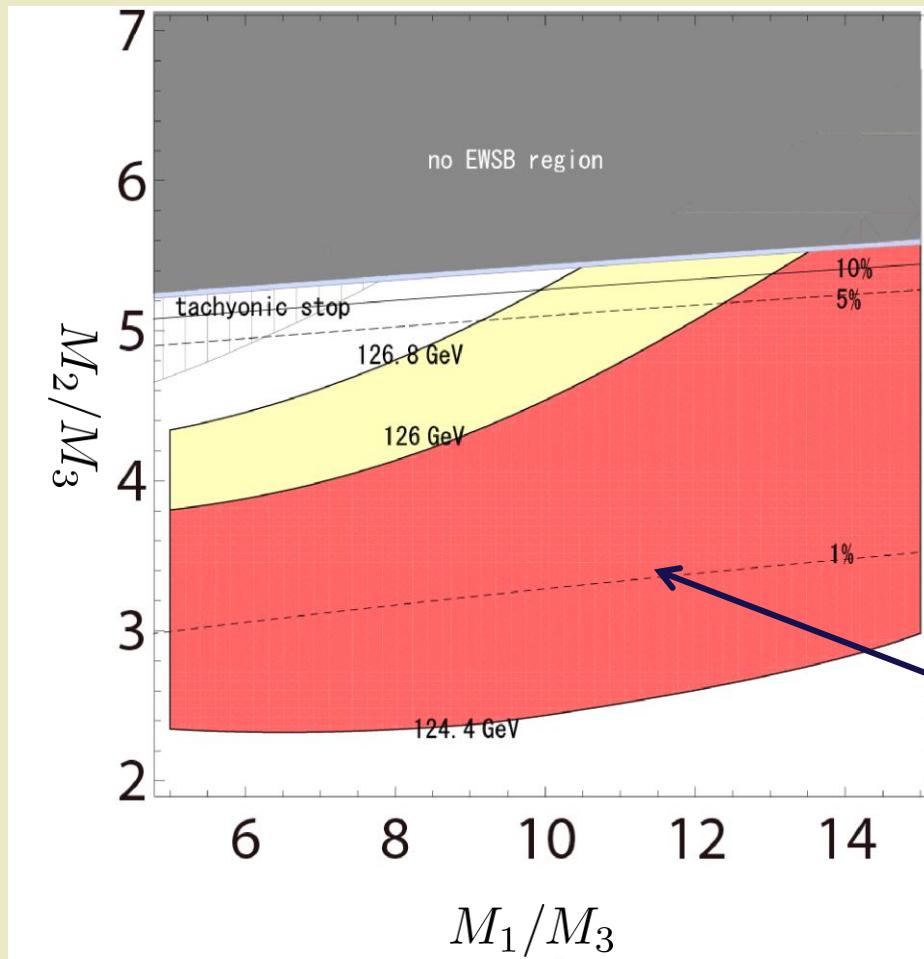
- moduli-mixing gauge kinetic function

$$\mathcal{L} \ni \int d^2\theta f_a(T) W^a W^a \ni f_a(T) F^{a\mu\nu} F_{\mu\nu}^a$$

$$f_a(T) = k_a^i T_i \quad \rightarrow \quad M_a = k_a^i \frac{F^{T_i}}{T_i + \bar{T}_i}$$

Higgs boson mass in NUGM

□ the Higgs boson mass and degree of tuning



$$M_3 = 385\text{GeV}$$

$$A_0 = -400\text{GeV}$$

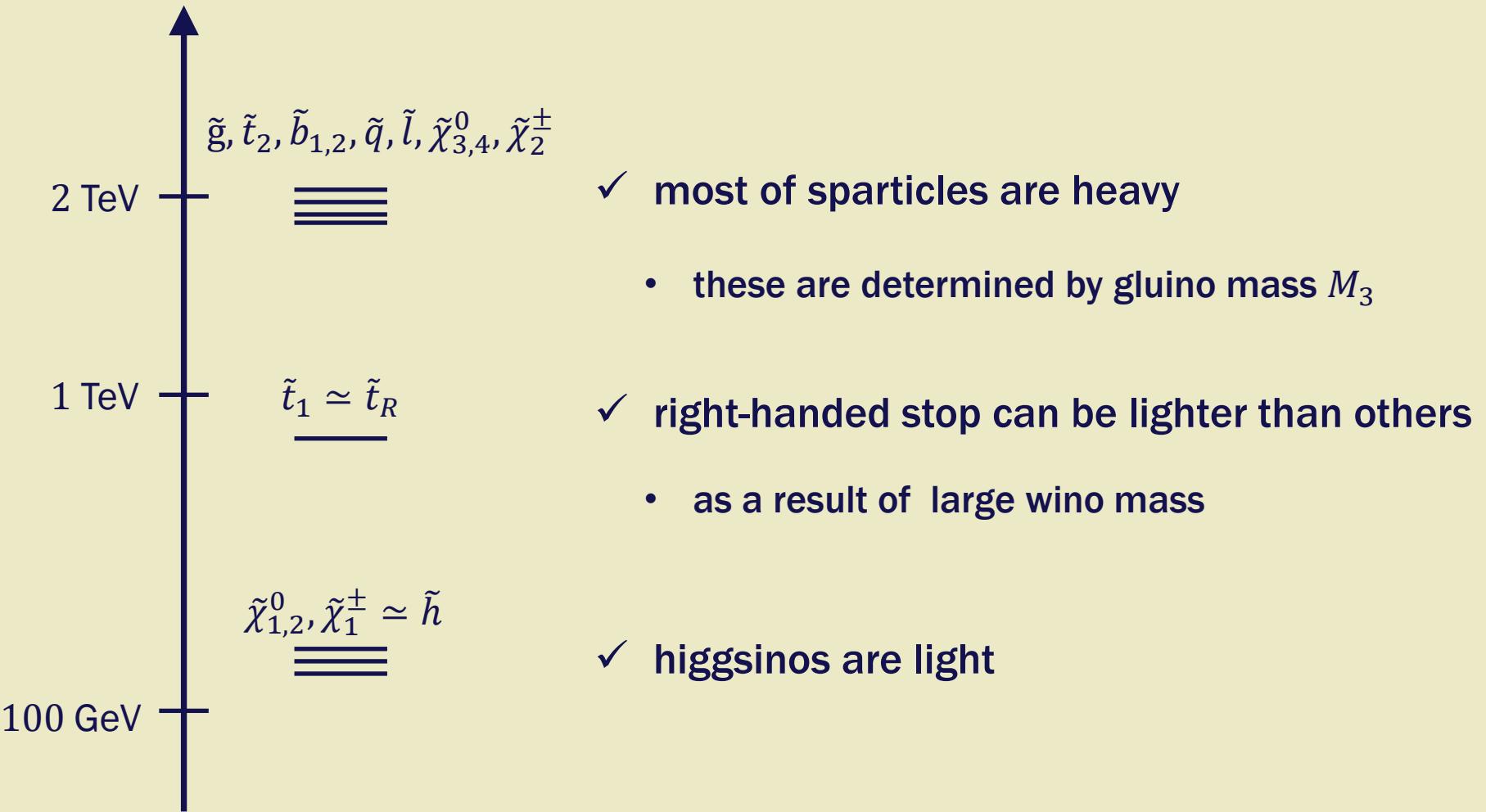
$$(m_0)_{\text{3rd}} = 200\text{GeV}$$

brief summary

- ✓ large wino mass enhances the Higgs boson mass
- ✓ A_t/M_{stop} is maximized at $M_2/M_3 \simeq 5$
- ✓ μ -parameter is also minimized at $M_2/M_3 \simeq 5$
- ✓ Higgs mass can reach 126 GeV even when $\Delta_\mu \lesssim 10$

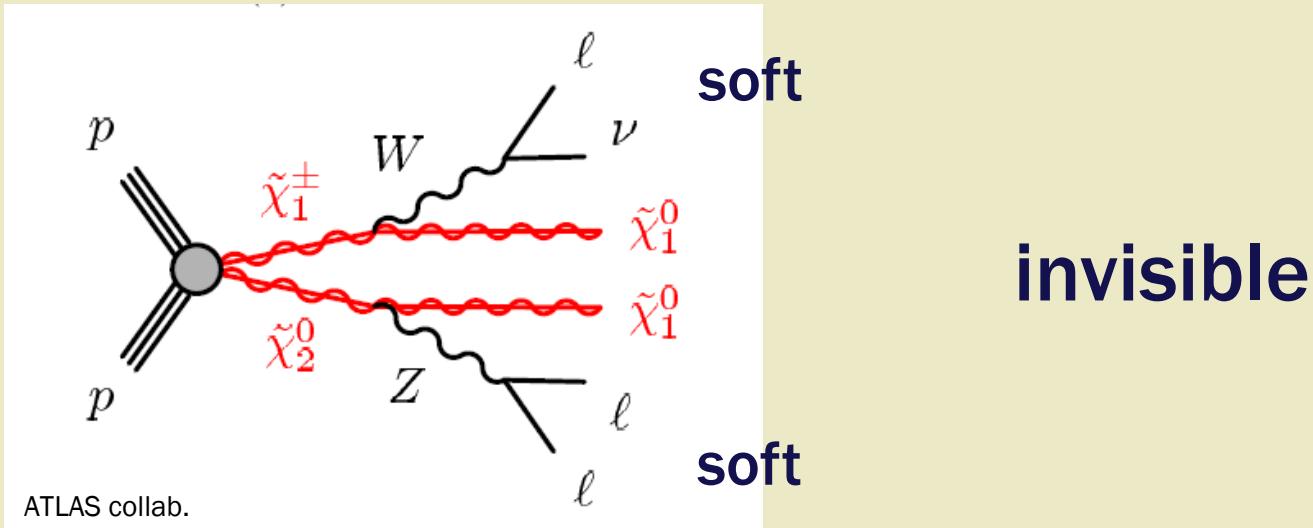
How to probe NUGM ?

typical mass spectrum



decays of higgsinos

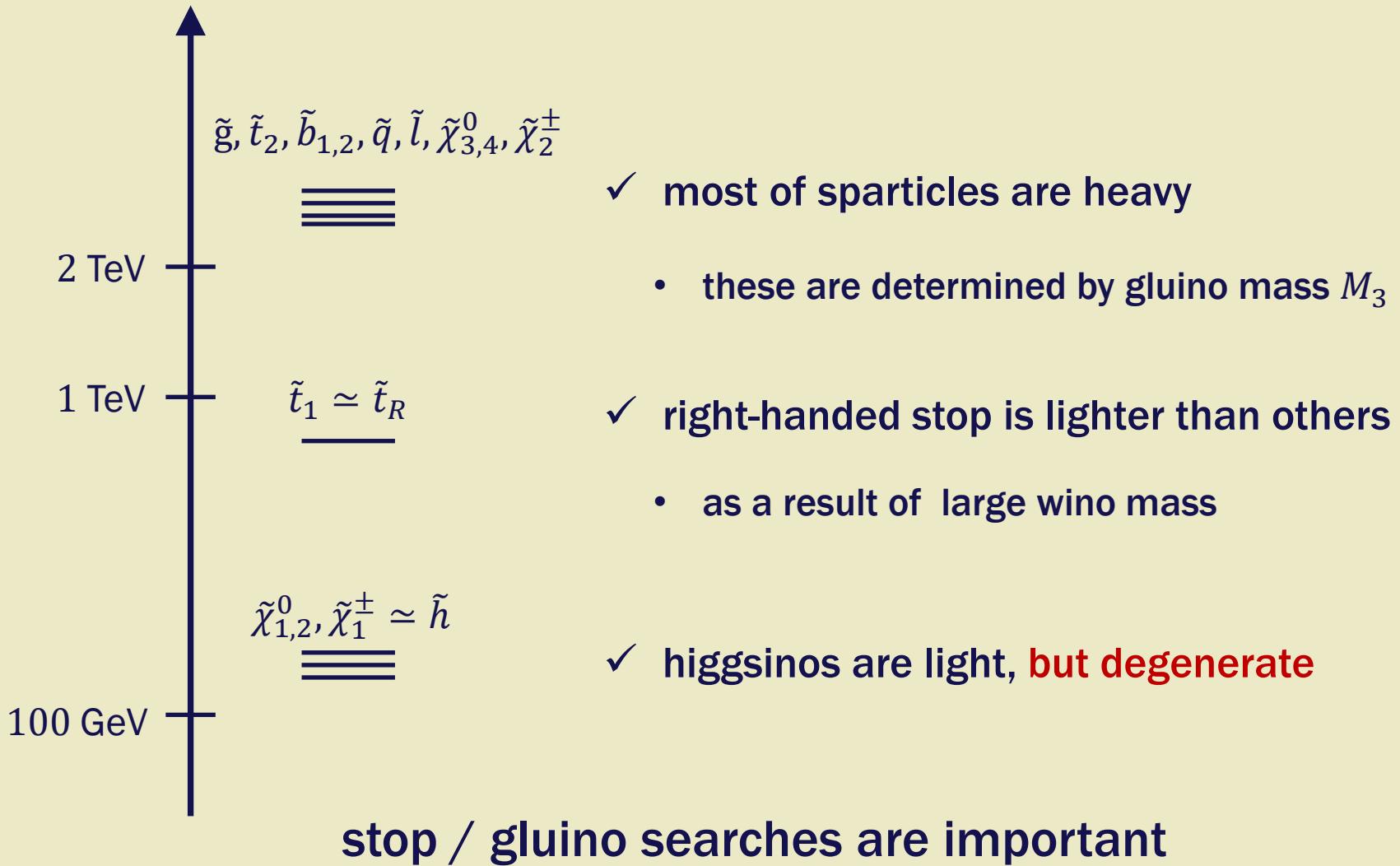
- higgsinos are light and degenerate $\Delta m_{\tilde{\chi}} \lesssim 2.0 \text{ GeV}$



- decay products are too soft to be reconstructed
- no charged tracks unlike pure wino

higgsino searches are not efficient

typical mass spectrum

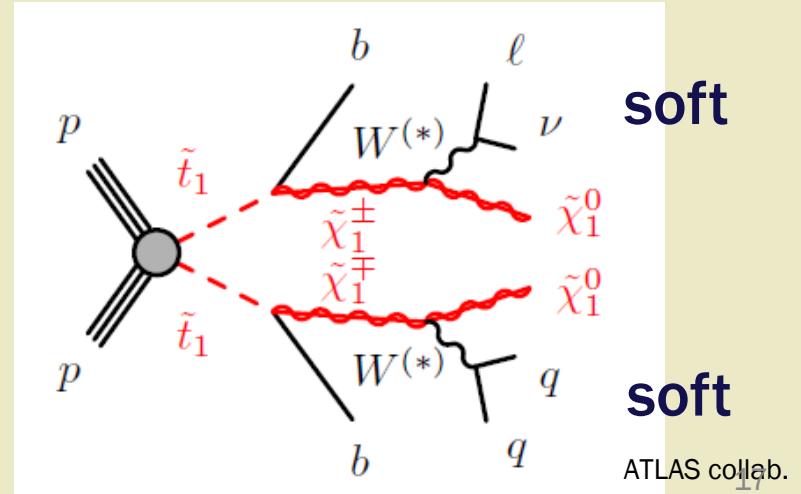
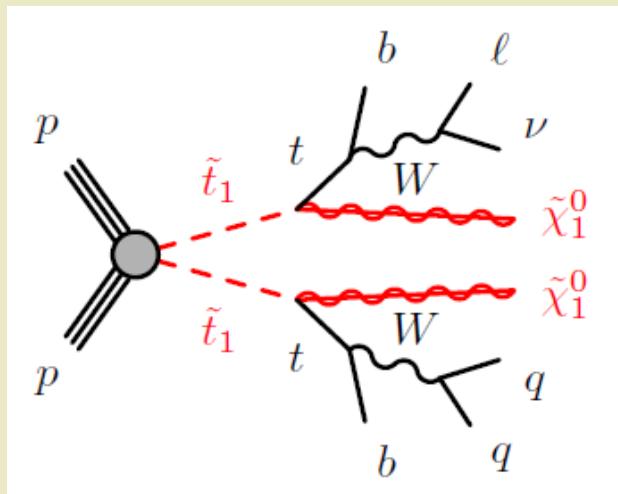


top squark decays

□ right-handed top squark is light in NUGM

$$W_{MSSM} \ni y_t (t_L \tilde{h}_u^0 - b_L \tilde{h}_u^+) \tilde{t}_R$$

- top squark decays to $t + \tilde{\chi}_1^0$ or $b + \tilde{\chi}_1^\pm$
- right-handed top squark couples to quark/higgsinos universally
- $\text{Br}(\tilde{t}_1 \rightarrow b \tilde{\chi}_1^\pm) = 1 - \text{Br}(\tilde{t}_1 \rightarrow t \tilde{\chi}_1^0) \simeq 0.5$ unless $m_{\tilde{t}_1} \simeq m_{\tilde{\chi}_1^\pm}$

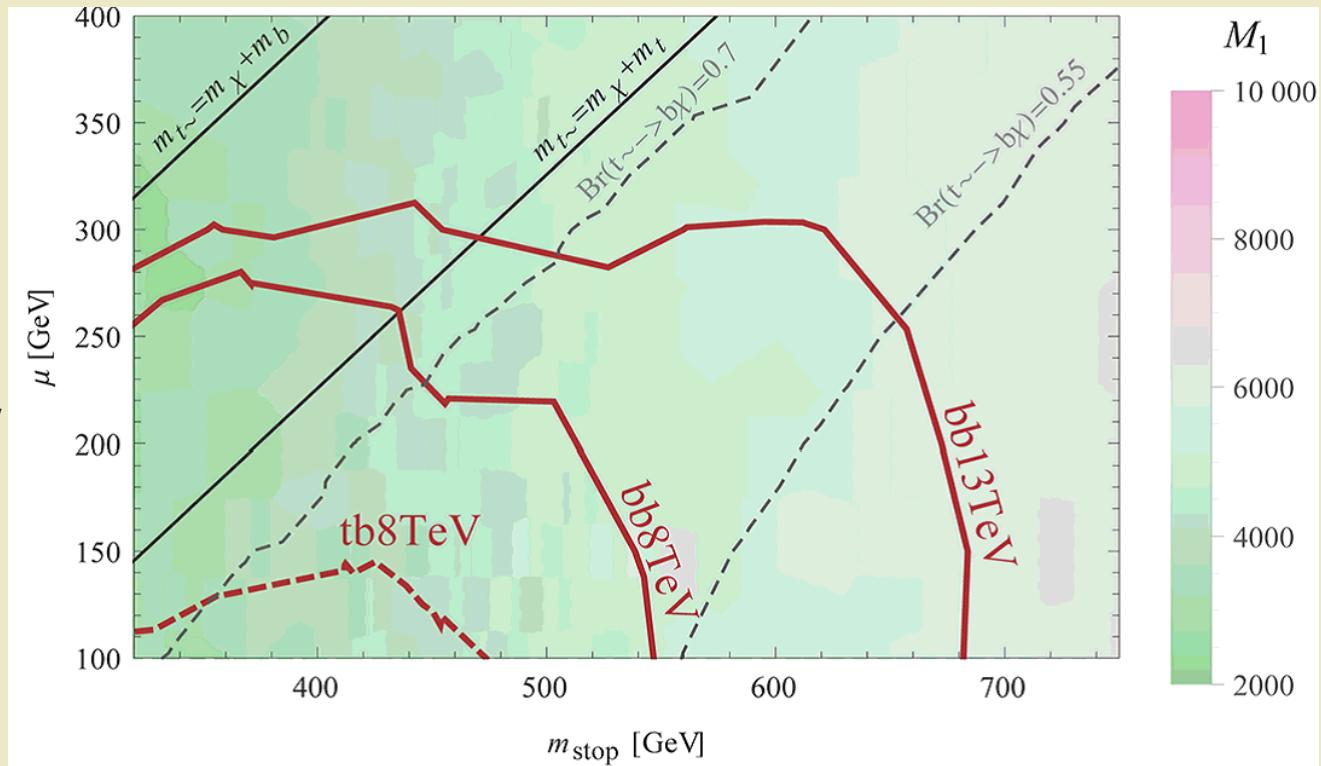


ATLAS collab.
17

top squark search

- ✓ signals are tt (25%) / tb (50%) / bb (25%) + MET
- ✓ bb+MET channel gives the severer bound than tt+MET in run-1 result [1]
- ✓ 13TeV data[2] has already given the severest bound

$\tan\beta = 15$
 $m_0 = 1 \text{ TeV}$
 $M_3 = 1 \text{ TeV}$
 $m_h = 125.8 \pm 0.3 \text{ GeV}$
 M_2, A_t : tuned



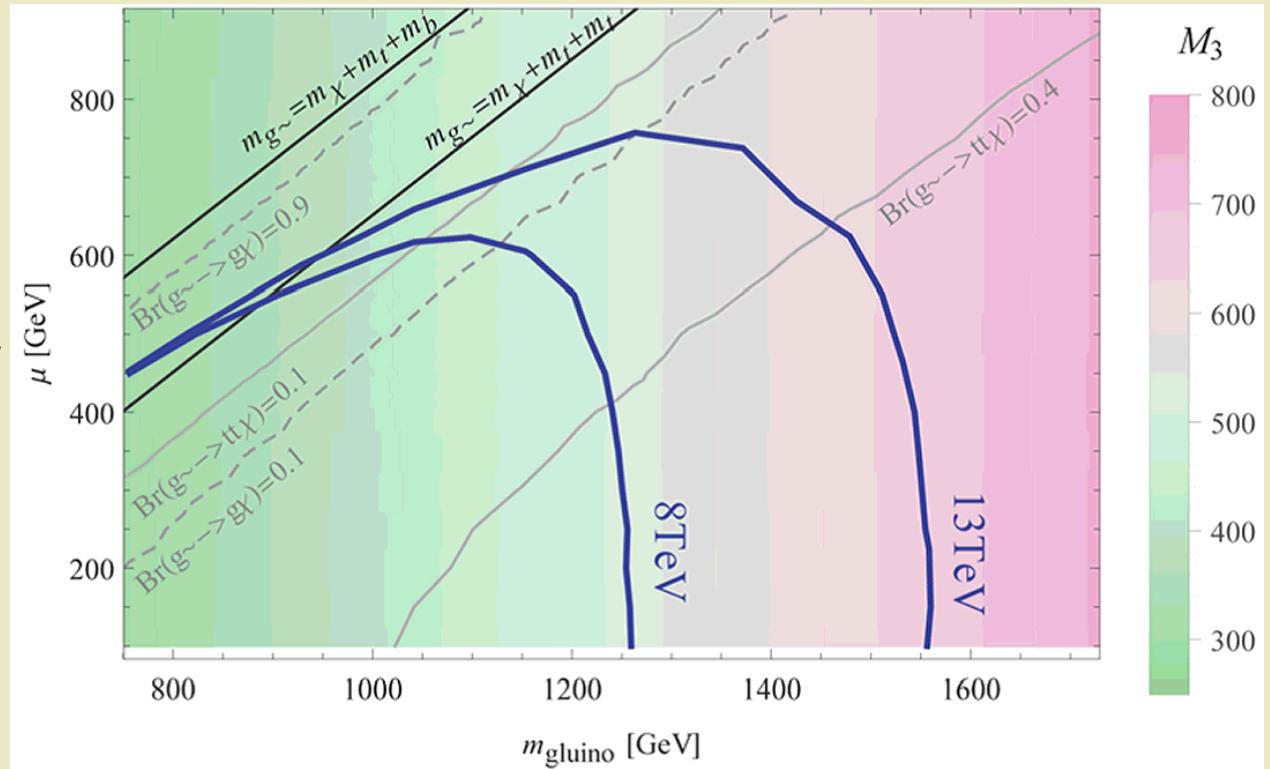
[1] ATLAS collab., JHEP 1310, 189 (2013), Eur.Phys.J.C 75, no.10, 510 (2015)

[2] TLAS-CONF-2015-066

gluino search

- ✓ gluino decays to top and stop: $\tilde{g} \rightarrow t \tilde{t}_1 \rightarrow t + t\tilde{\chi}_1^0/b \tilde{\chi}_1^+$
- ✓ signals are characterized by 4 bottoms and large MET
- ✓ 13TeV data [3] has already given the severest bound

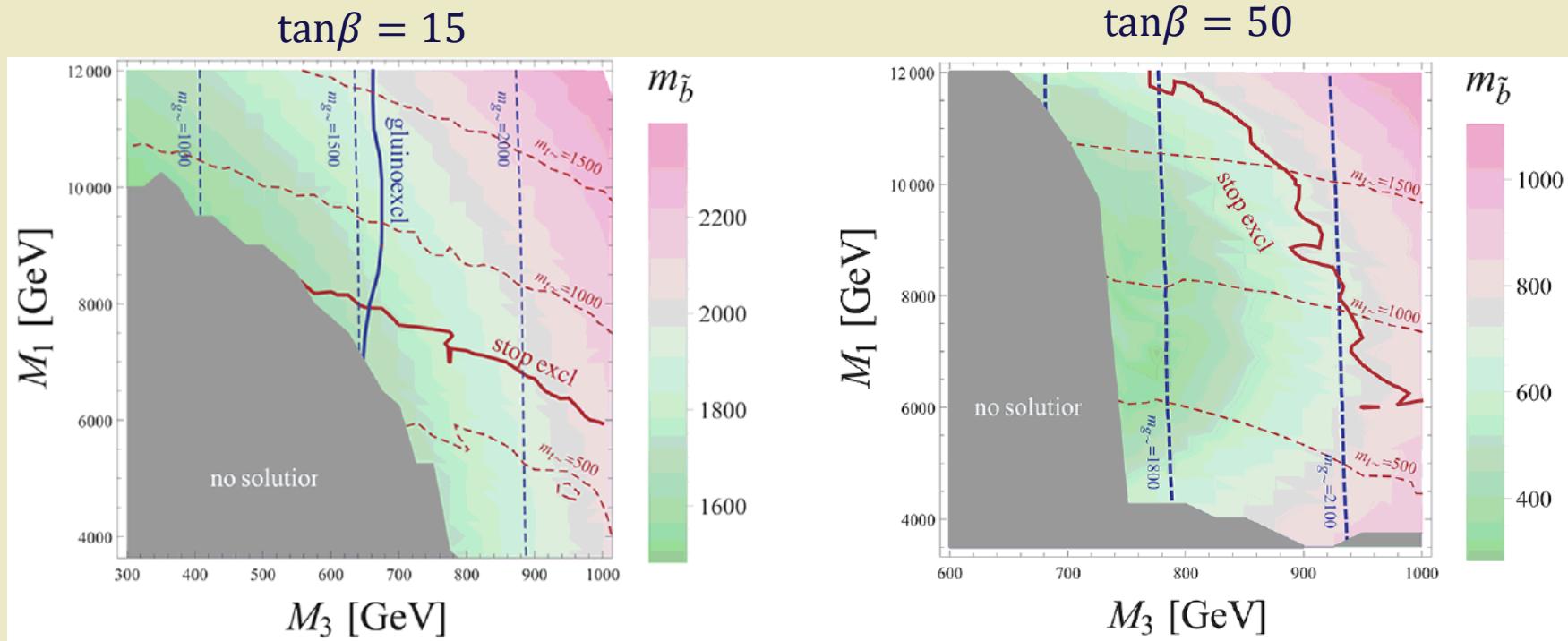
$\tan\beta = 15$
 $m_0 = 1 \text{ TeV}$
 $M_1 = 12 \text{ TeV}$
 $m_h = 125.8 \pm 0.3 \text{ GeV}$
 M_2, A_t : tuned



bounds on boundary conditions

- ✓ right-handed sbottom can also be light for large $\tan\beta$
- ✓ sbottom pair production gives same signal as stop
- ✓ stop tends to be tachyonic for small M_1 and M_3

$m_0 = 1 \text{ TeV}$
 $\mu = 150 \text{ GeV}$
 $m_h = 125.8 \pm 0.3 \text{ GeV}$
 M_2, A_t : tuned



conclusion

- NUGM can realize 126GeV Higgs and small μ -parameter
- right-handed top squark tends to be light
- stop/gluino search are important for NUGM scenario
- $m_{\tilde{t}_1} \lesssim 700$ GeV, $m_{\tilde{g}} \lesssim 1.6$ TeV is excluded by the latest data

thank you for your attention

backups

degenerate higgsinos

□ higgsinos are light and degenerate

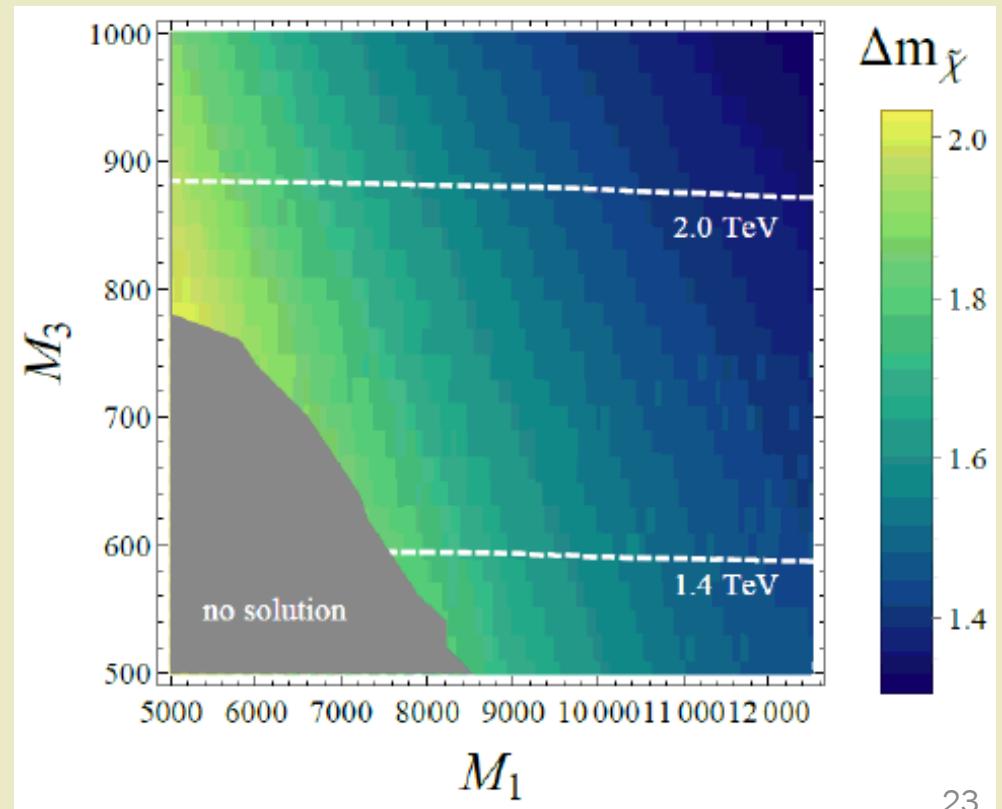
heavy bino, wino $M_{1,2} \gg \mu$ leads $\Delta m_{\tilde{\chi}} = O(1\text{GeV})$

$m_0 = 1 \text{ TeV}$
 $\tan\beta = 15$
 $\mu = 150 \text{ GeV}$
 M_2, A_t : tuned

$$\Delta m_{\tilde{\chi}} \equiv m_{\tilde{\chi}_1^\pm} - m_{\tilde{\chi}_1^0} \simeq \frac{m_Z^2}{M_2(m_Z)}$$

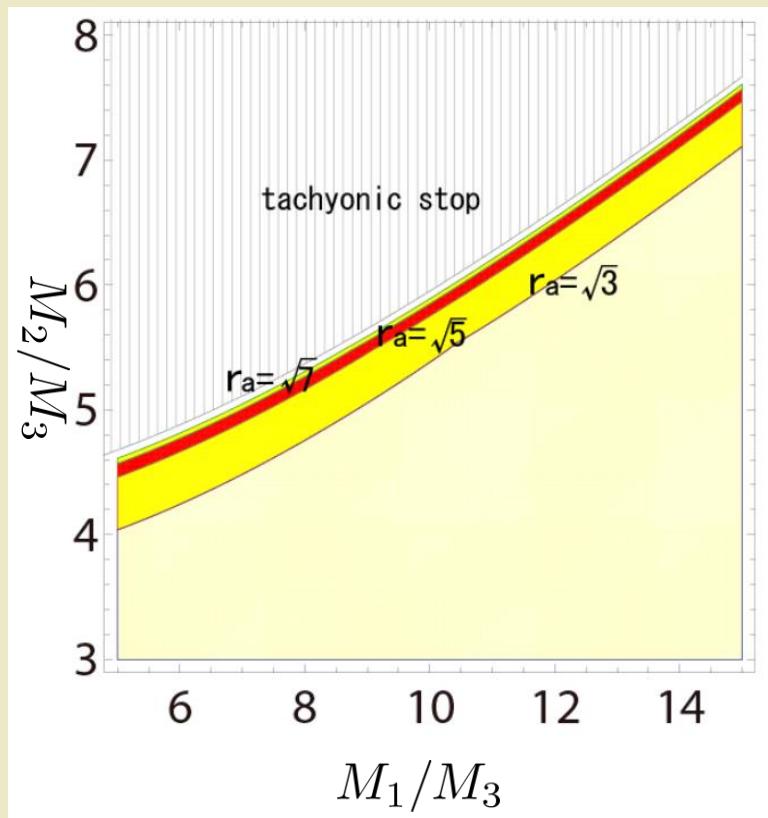
$$M_{1,2} \simeq 2 - 4 \text{ TeV}$$

$$\rightarrow \Delta m_{\tilde{\chi}} \lesssim 2.0 \text{ GeV}$$



Higgs boson mass in NUGM

$$r_a \equiv A_t / M_{st}$$

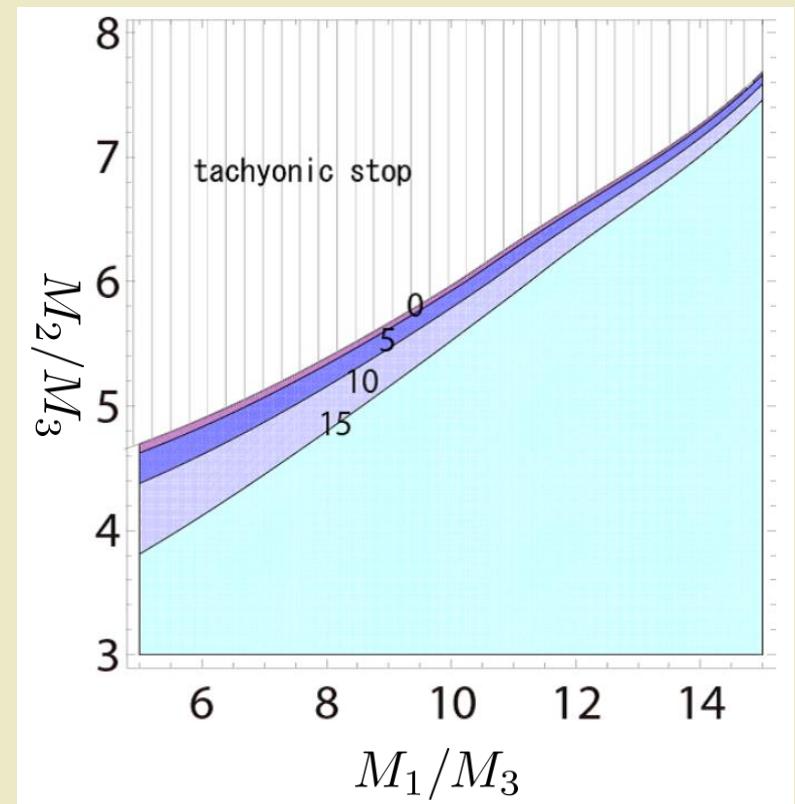


$$M_3 = 385 \text{ GeV}$$

$$A_0 = -400 \text{ GeV}$$

$$(m_0)_{\text{3rd}} = 200 \text{ GeV}$$

$$M_{stop} / m_{\tilde{t}_L}$$



$$M_{stop} \equiv \sqrt{m_{\tilde{t}_R} m_{\tilde{t}_L}}$$