

# Dark Matter Annihilation into Fermions and a Photon

Debtosh Chowdhury  
INFN, Rome

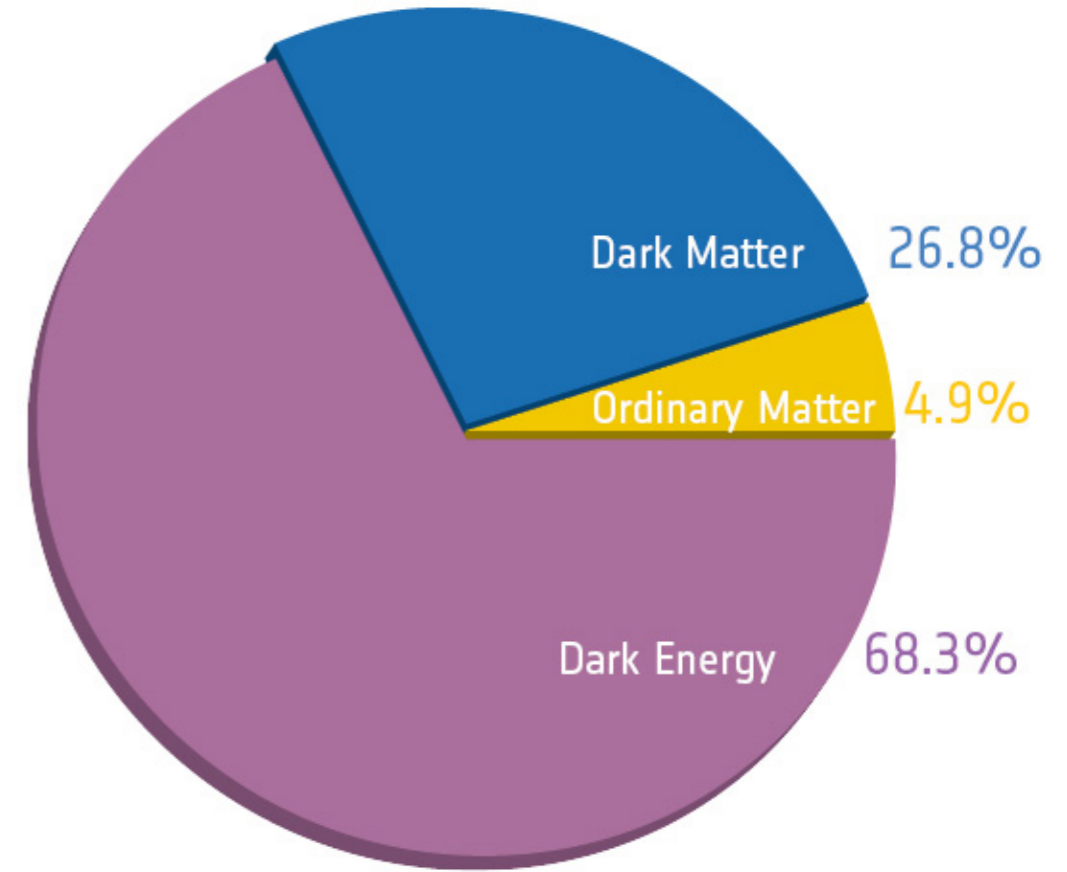
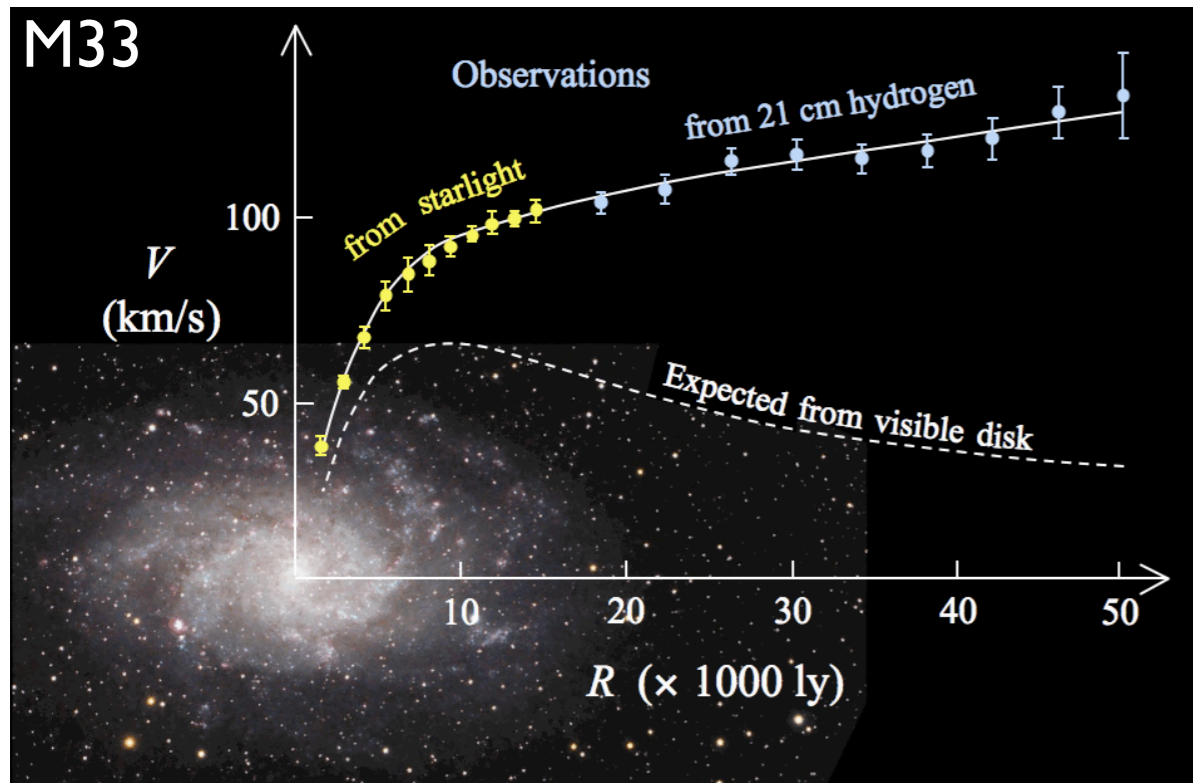
in collaboration w/ A. Iyer and R. Laha  
[arXiv:1601.06140]



SUSY 2016  
Melbourne, Australia  
July 4, 2016

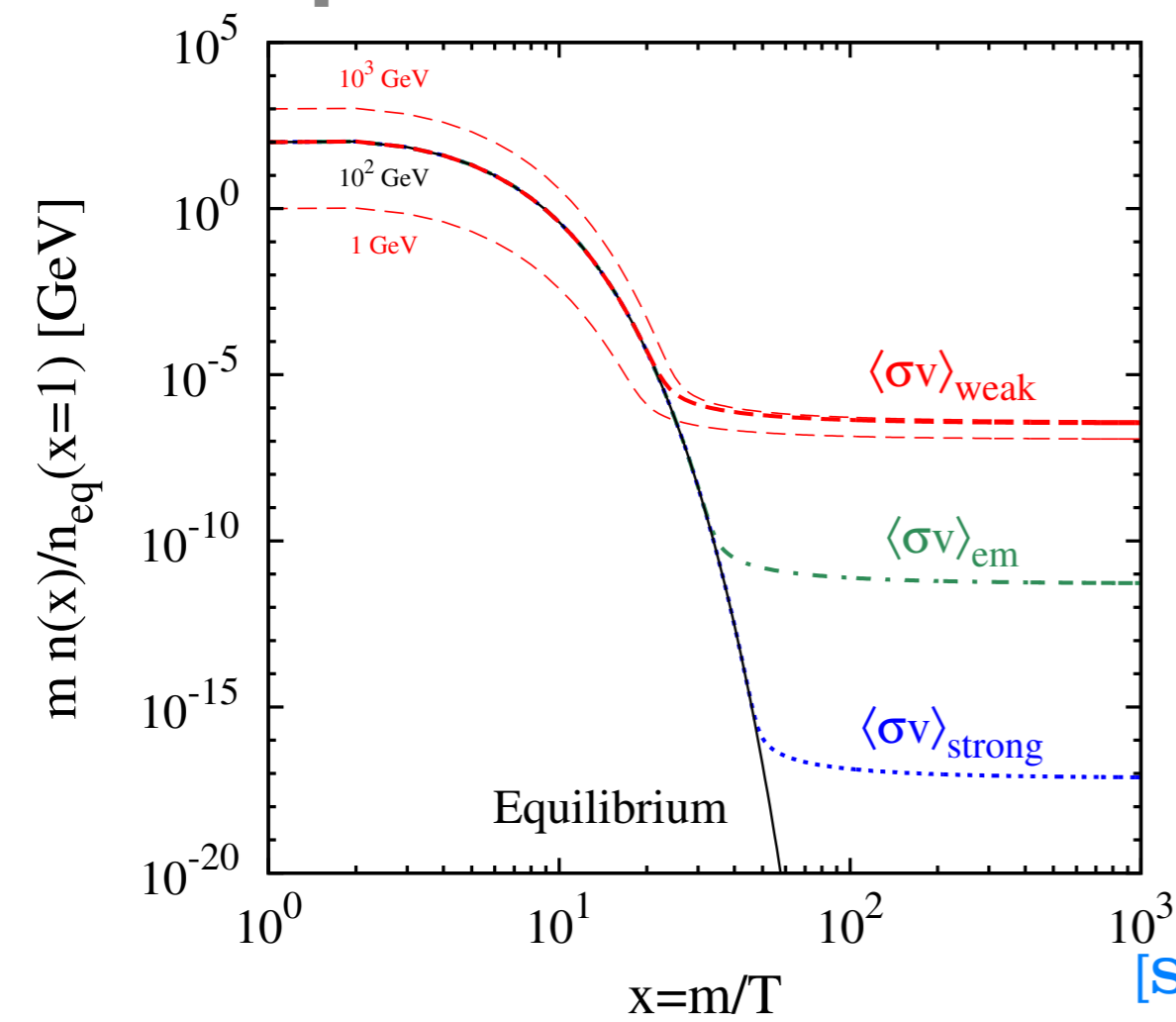
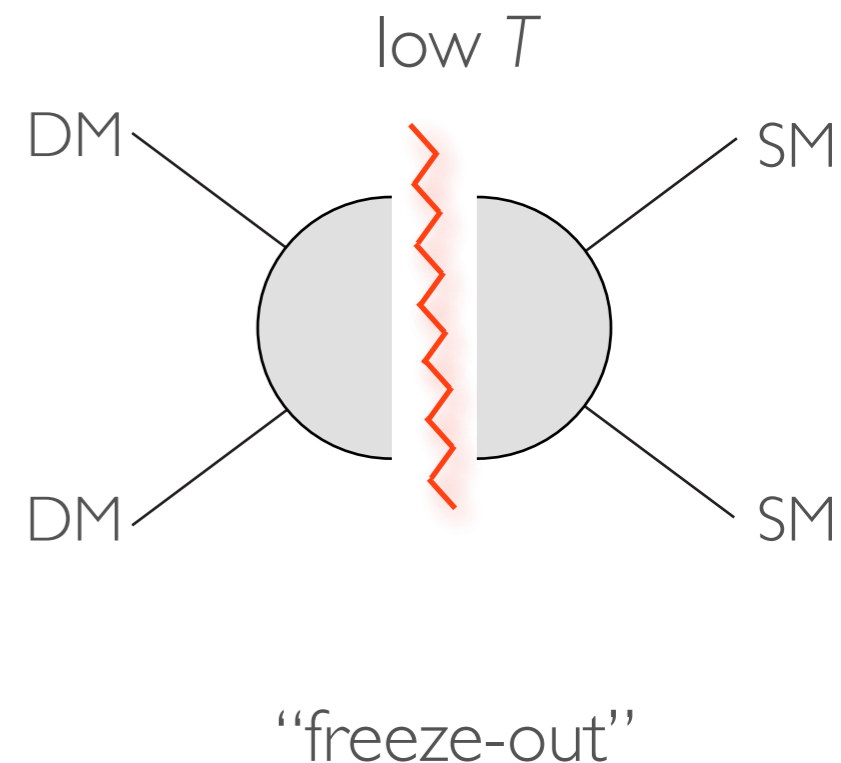
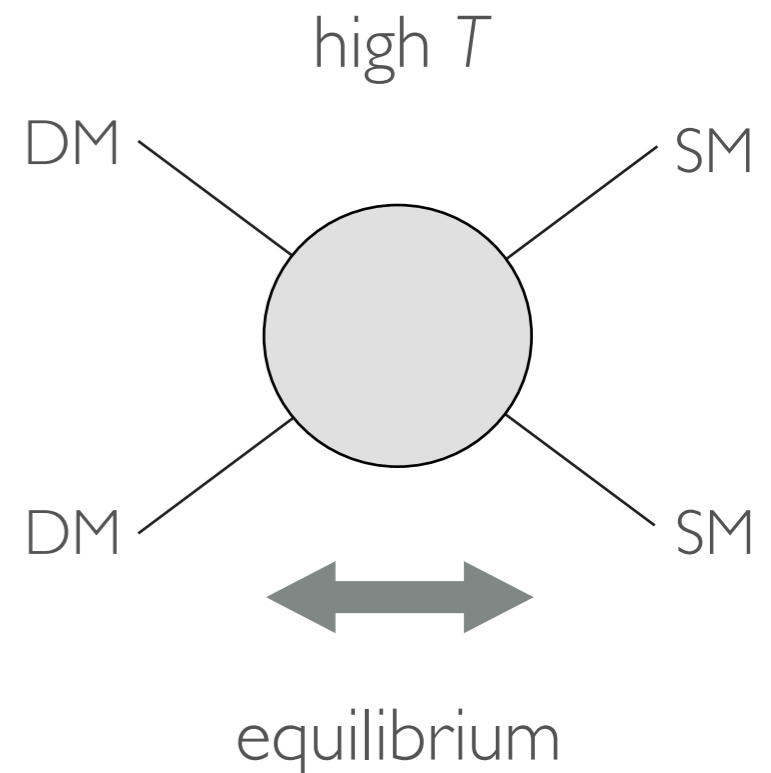


# Evidence of Dark Matter



# Thermal Relic

- \* In early Universe  $DM DM \rightleftharpoons SM SM$ .
- \* Universe cools down  $DM DM \Rightarrow SM SM$ .
- \* After further cooling, the interactions stop. DM **“freezes out”**.

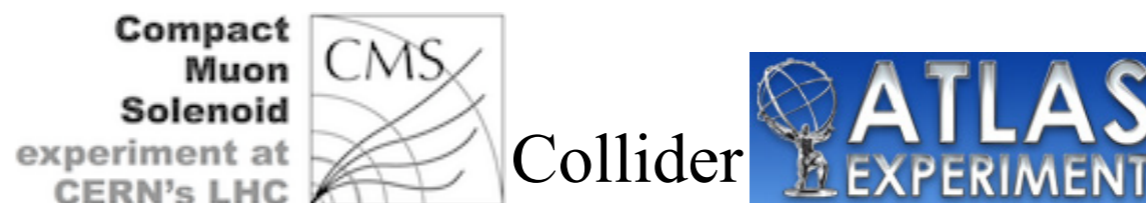
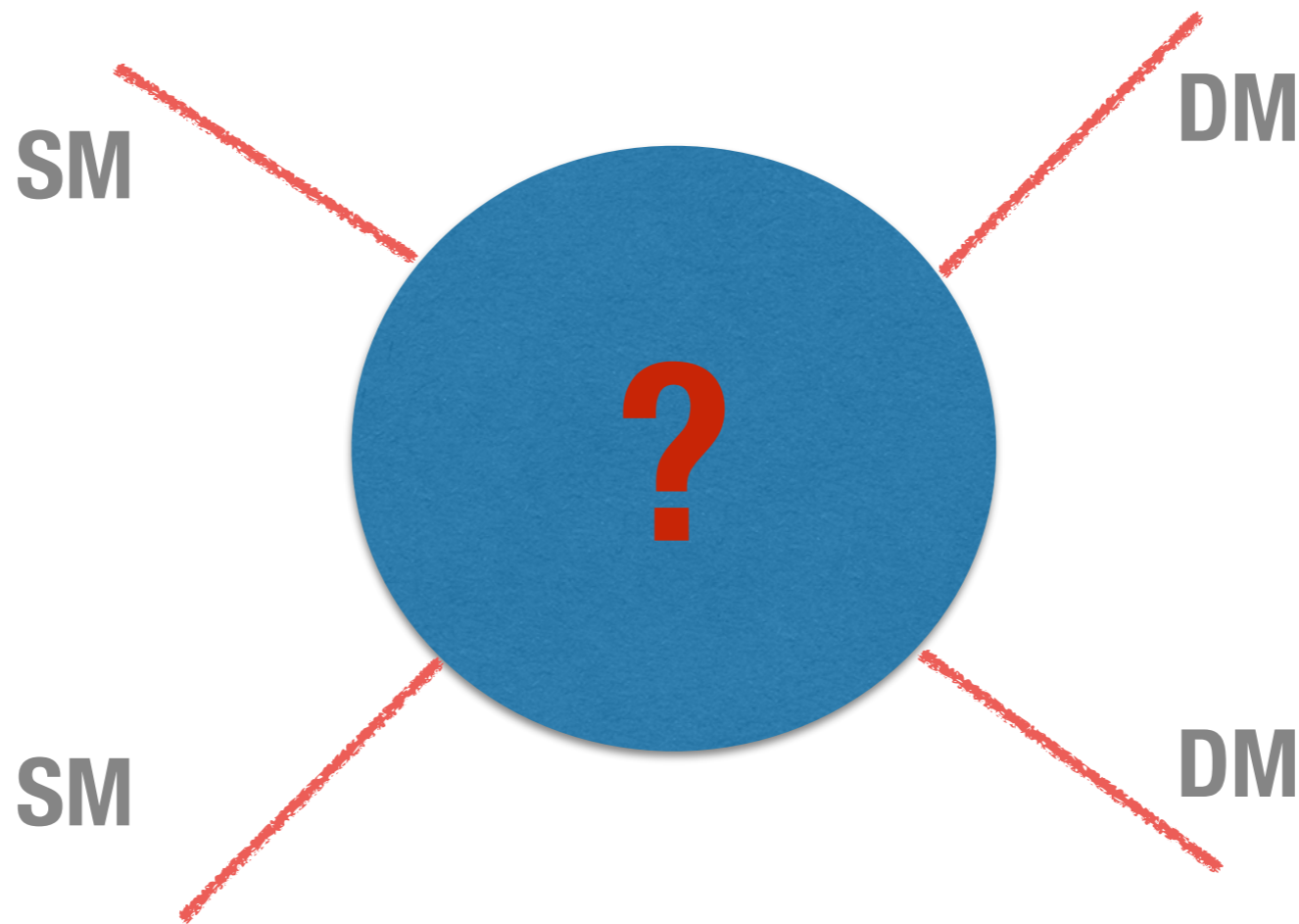
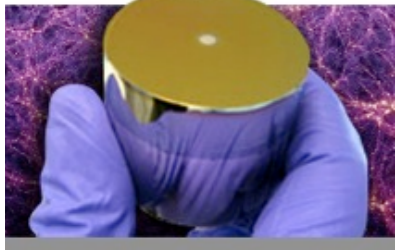


$$\Omega h^2 \propto \frac{1}{\langle \sigma v \rangle}$$

# Dark Matter Detection



Direct  
Detection





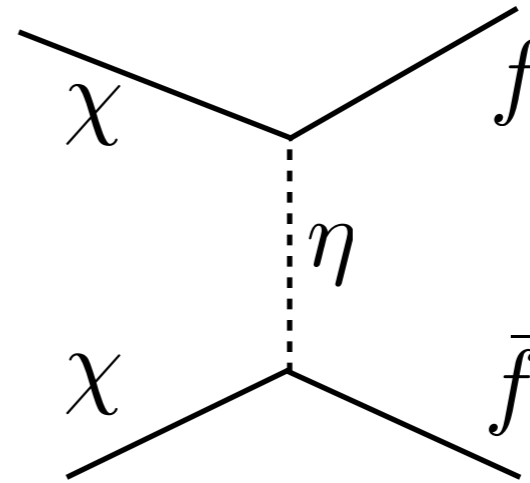
# DM Annihilation into Fermions

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_\chi + \mathcal{L}_\eta + \mathcal{L}_{\text{int}}$$

$$\mathcal{L}_\chi = \frac{1}{2} \bar{\chi}^c i \not{\partial} \chi - \frac{1}{2} m_\chi \bar{\chi}^c \chi,$$

$$\mathcal{L}_\eta = (D_\mu \eta)^\dagger (D^\mu \eta) - m_\eta^2 \eta^\dagger \eta$$

$$\mathcal{L}_{\text{int}} = -y \bar{\chi} \Psi_R \eta + \text{h.c.}$$



$$\langle \sigma v \rangle = a + bv^2 + \mathcal{O}(v^4)$$

$$\mu \equiv (m_\eta / m_\chi)^2$$

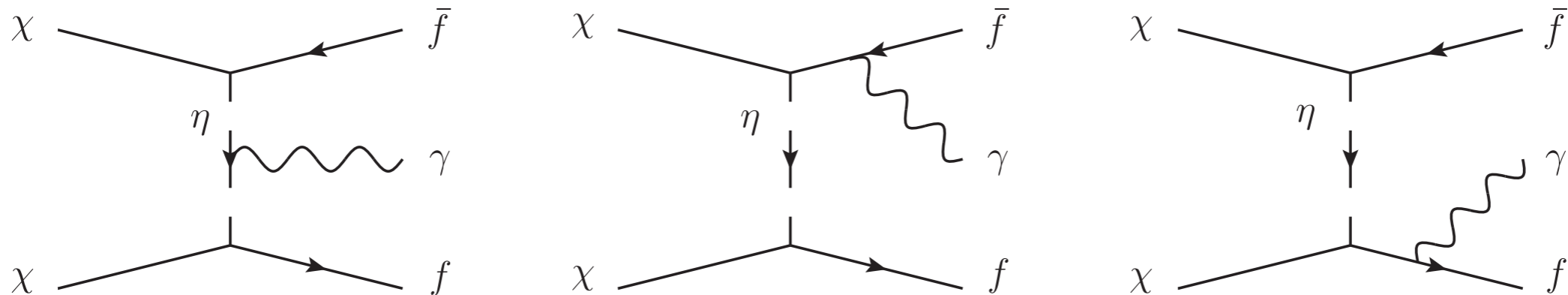
$$(\sigma v)_{2\text{-body}}^{\text{s-wave}} = \frac{y^4 N_c}{32\pi m_\chi^2} \frac{m_f^2}{m_\chi^2} \frac{1}{(1 + \mu)^2}$$

$$(\sigma v)_{2\text{-body}}^{\text{p-wave}} = v^2 \frac{y^4 N_c}{48\pi m_\chi^2} \frac{1 + \mu^2}{(1 + \mu)^4}$$

$$v \sim 10^{-3}$$

[Bergstrom '89, ...]

# Virtual Internal Bremsstrahlung



$$(\sigma v)_{3\text{-body}} \simeq \frac{\alpha_{\text{em}} y^4 N_c Q_f^2}{64\pi^2 m_\chi^2} \left\{ (\mu + 1) \left[ \frac{\pi^2}{6} - \ln^2 \left( \frac{\mu + 1}{2\mu} \right) - 2\text{Li}_2 \left( \frac{\mu + 1}{2\mu} \right) \right] \right. \\
 \left. + \frac{4\mu + 3}{\mu + 1} + \frac{4\mu^2 - 3\mu - 1}{2\mu} \ln \left( \frac{\mu - 1}{\mu + 1} \right) \right\}$$

**VIB lifts the helicity suppression of DM annihilation into pair of fermions**

[Bergstrom '89, ... ]

# VIB in Supersymmetry

Admixture of Bino/Wino/Higgsinos:

$$\chi = \alpha \tilde{B} + \beta \tilde{W} + \gamma \tilde{H}_d + \delta \tilde{H}_u$$

$SU(2)_L$     *singlet*    *triplet*    *doublets*

$$\mathcal{L}_{\text{int}}^{\chi \tilde{f} f} = y_L \bar{\chi} f_L \tilde{f}_L + y_R \bar{\chi} f_R \tilde{f}_R + \text{h.c.}$$

$$y_R = \sqrt{2} Q_f g \tan \theta_W N_{11}$$

$$y_L = -\frac{2Q_f \mp 1}{\sqrt{2}} g \tan \theta_W N_{11} \mp \frac{g}{\sqrt{2}} N_{12}$$

# VIB in Effective Operator Model

$$\mathcal{L}_{d=6} = \frac{1}{\Lambda^2} (\bar{\chi} \gamma^5 \gamma^\mu \chi) (\bar{f} \gamma_\mu f)$$

$$\mathcal{L}_{d=8} = \frac{1}{\Lambda^4} (\bar{\chi} \gamma^5 \gamma^\mu \chi) \left[ \left( \bar{f}_L \overleftarrow{D}_\rho \right) \gamma_\mu \left( \overrightarrow{D}^\rho f_L \right) + \left( \bar{f}_R \overleftarrow{D}_\rho \right) \gamma_\mu \left( \overrightarrow{D}^\rho f_R \right) \right]$$

$$\bar{f}_L \overleftarrow{D}_\mu = (\partial_\mu \bar{f}_L) - ig \frac{\sigma^i}{2} W_\mu^i \bar{f}_L - ig' Y_f B_\mu \bar{f}_L,$$

$$\overrightarrow{D}_\mu f_L = (\partial_\mu f_L) + ig \frac{\sigma^i}{2} W_\mu^i f_L + ig' Y_f B_\mu f_L,$$

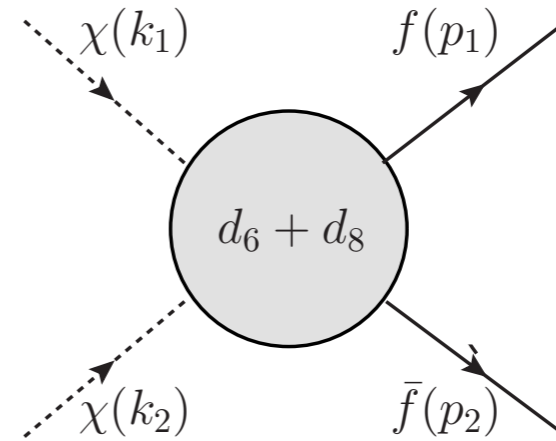
$$\bar{f}_R \overleftarrow{D}_\mu = (\partial_\mu \bar{f}_R) - ig' Y_f B_\mu \bar{f}_R,$$

$$\overrightarrow{D}_\mu f_R = (\partial_\mu f_R) + ig' Y_f B_\mu f_R,$$

$$\mathcal{L} = d_6 \mathcal{L}_{d=6} + d_8 \mathcal{L}_{d=8}$$



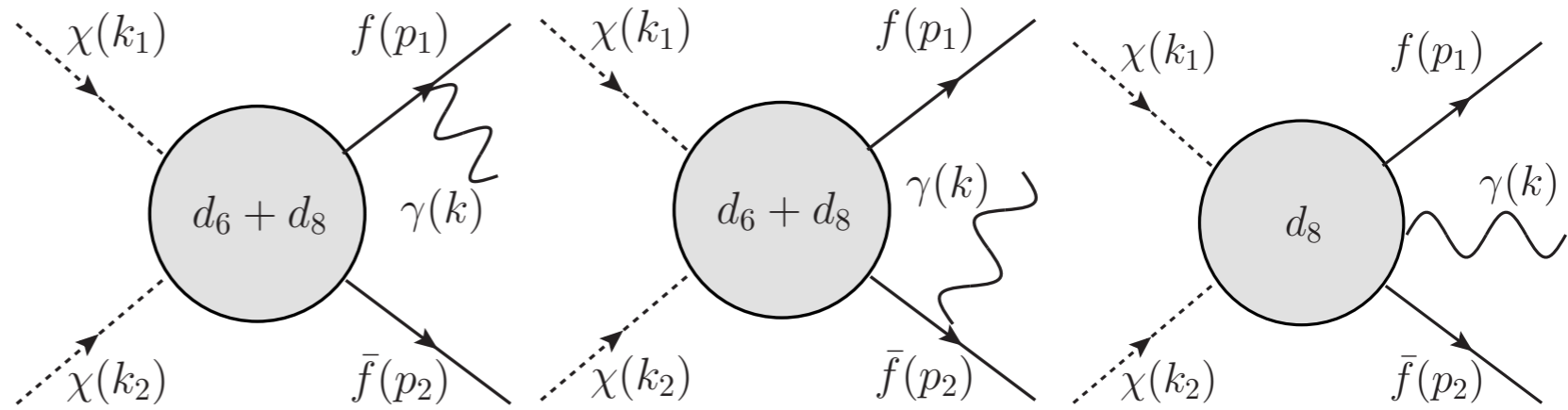
$$\chi\chi \rightarrow f\bar{f}$$



$$\mathcal{L} = \frac{d_6}{\Lambda^2} (\bar{\chi} \gamma^5 \gamma^\mu \chi) (\bar{f} \gamma_\mu f) + \frac{d_8}{\Lambda^4} (\bar{\chi} \gamma^5 \gamma^\nu \chi) \left[ (\partial_\rho \bar{f}) \gamma_\nu (\partial^\rho f) \right]$$

$$\sigma v = \frac{v^2}{3\pi\Lambda^8} \sqrt{1 - \frac{m_f^2}{m_\chi^2}} (2m_\chi^2 + m_f^2) \left\{ |d_6|^2 \Lambda^4 + |d_8|^2 (2m_\chi^2 + m_f^2)^2 - \Lambda^2 (d_6 d_8^* + d_6^* d_8) (2m_\chi^2 - m_f^2) \right\}$$

$$\chi\chi \rightarrow f\bar{f}\gamma$$

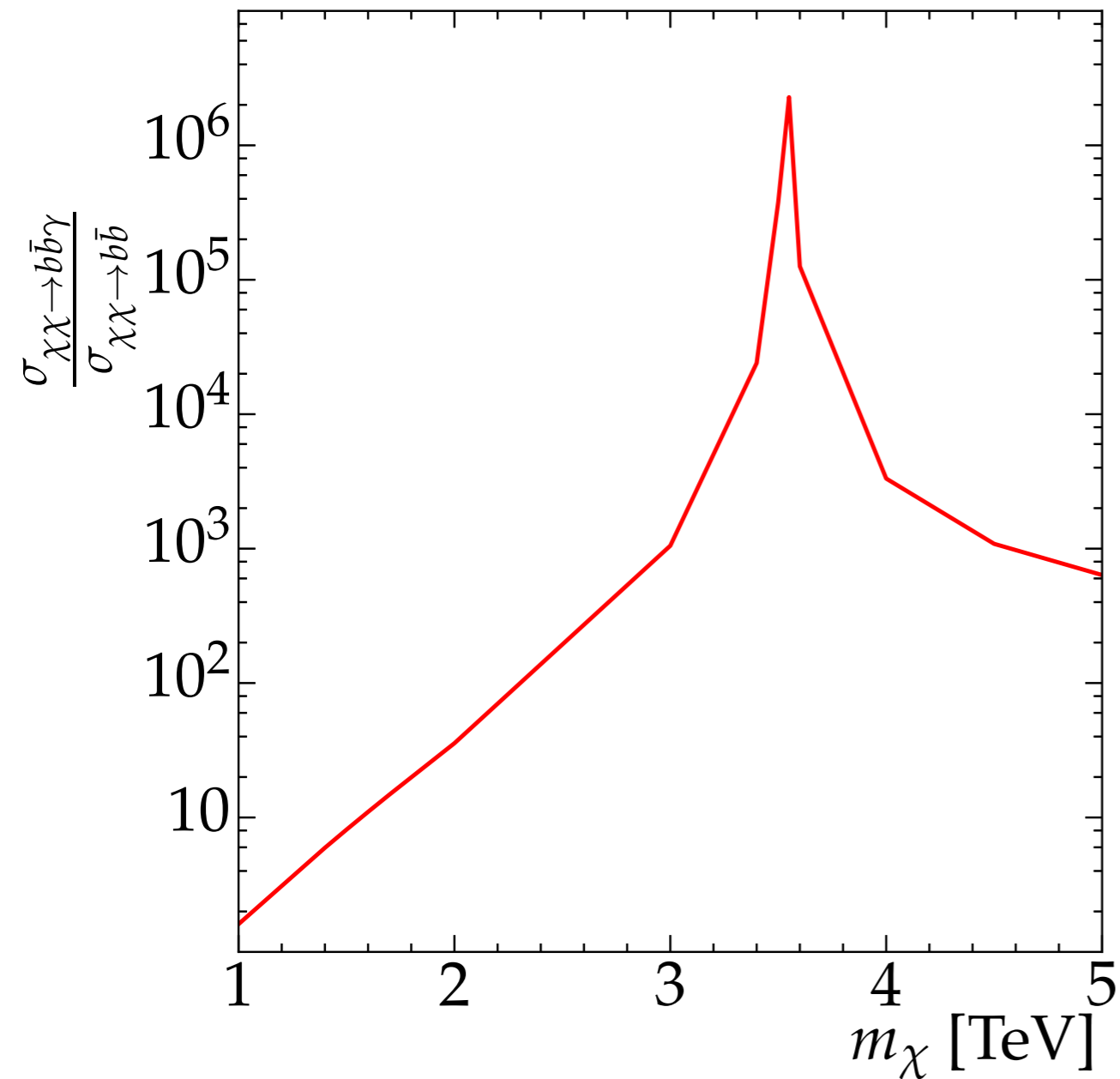


$$\begin{aligned} \mathcal{L}_{\text{eff}} = & \frac{d_6}{\Lambda^2} (\bar{\chi}\gamma^5\gamma^\mu\chi)(\bar{f}\gamma_\mu f) \\ & + \frac{d_8}{\Lambda^4} (\bar{\chi}\gamma^5\gamma^\nu\chi) \left[ (\partial_\rho\bar{f})\gamma_\nu(\partial^\rho f) \right. \\ & \left. + i e Q A_\rho \left\{ (\partial^\rho\bar{f})\gamma_\nu f - \bar{f}\gamma_\nu(\partial^\rho f) \right\} \right] \end{aligned}$$

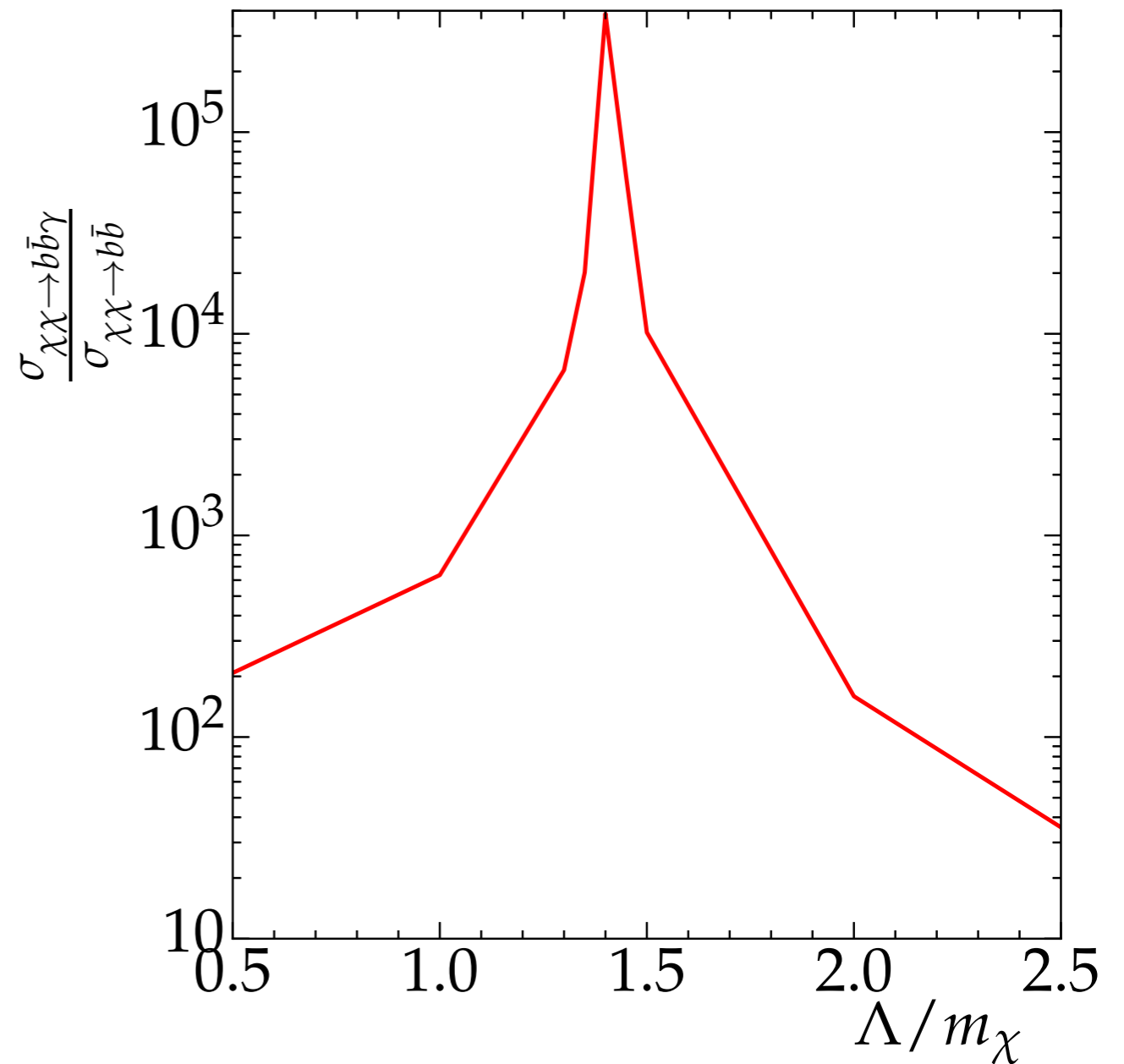
$$\begin{aligned} \mathcal{M} = & -i e Q \bar{u}(k_1)\gamma^5\gamma^\mu v(k_2) \left[ \left( \frac{d_6}{\Lambda^2} - \frac{d_8}{\Lambda^4} p_2 \cdot (p_1 + k) \right) \bar{u}(p_1) \not{\epsilon}(k) \frac{p_1 + \not{k} + m_f}{(p_1 + k)^2 - m_f^2 + i\epsilon} \gamma_\mu v(p_2) \right. \\ & + \left( \frac{d_6}{\Lambda^2} - \frac{d_8}{\Lambda^4} p_1 \cdot (p_2 + k) \right) \bar{u}(p_1) \gamma_\mu \frac{-p_2 - \not{k} + m_f}{(p_2 + k)^2 - m_f^2 + i\epsilon} \not{\epsilon}(k) v(p_2) \\ & \left. - \frac{d_8}{\Lambda^4} \bar{u}_{p_1} \epsilon_\rho(k) \gamma_\mu v_{p_2} (p_1^\rho - p_2^\rho) \right] \end{aligned}$$

# Ratio of the cross-sections

$$d_6 = 1, d_8 = 1, v = 10^{-3}$$

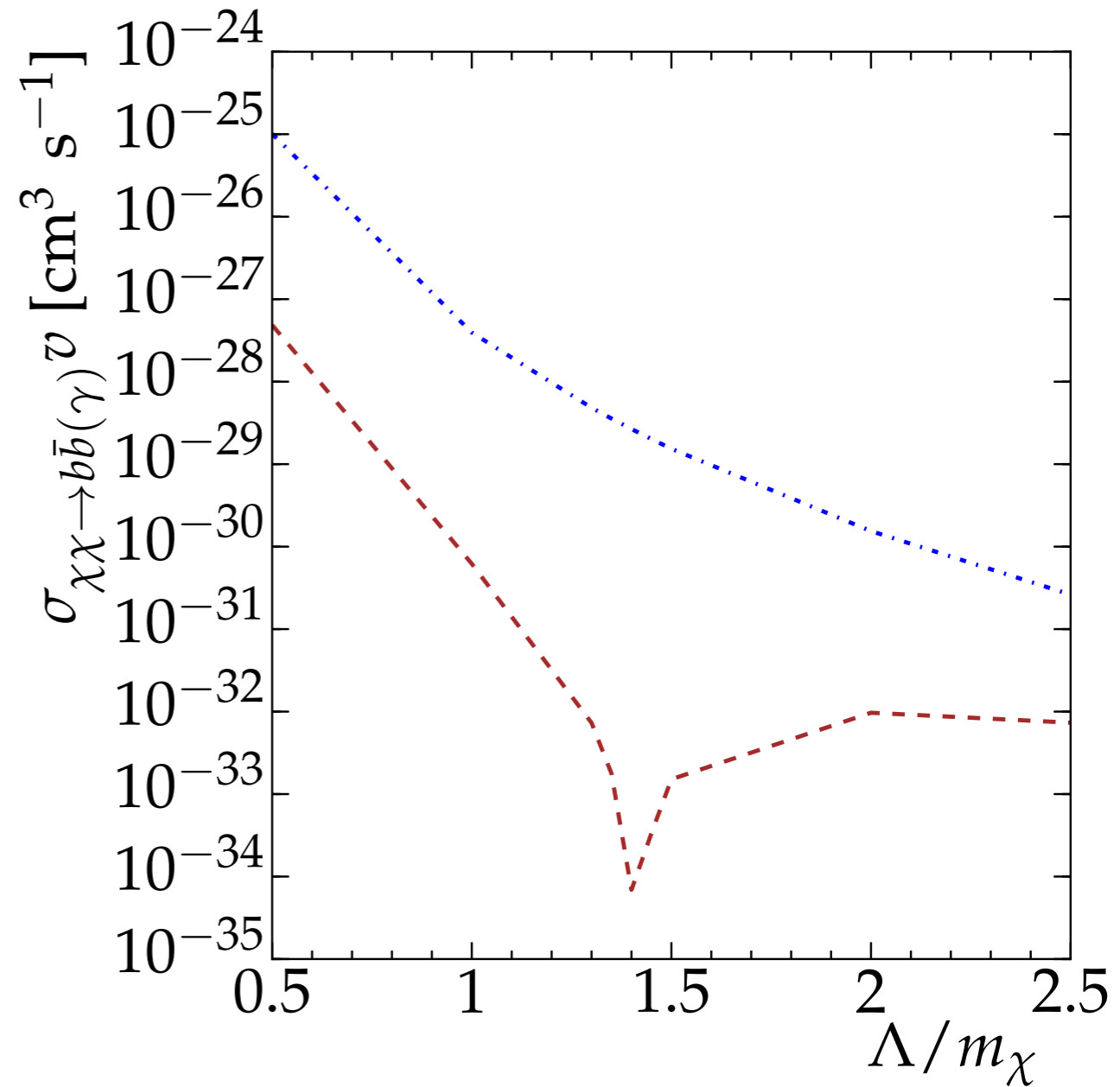
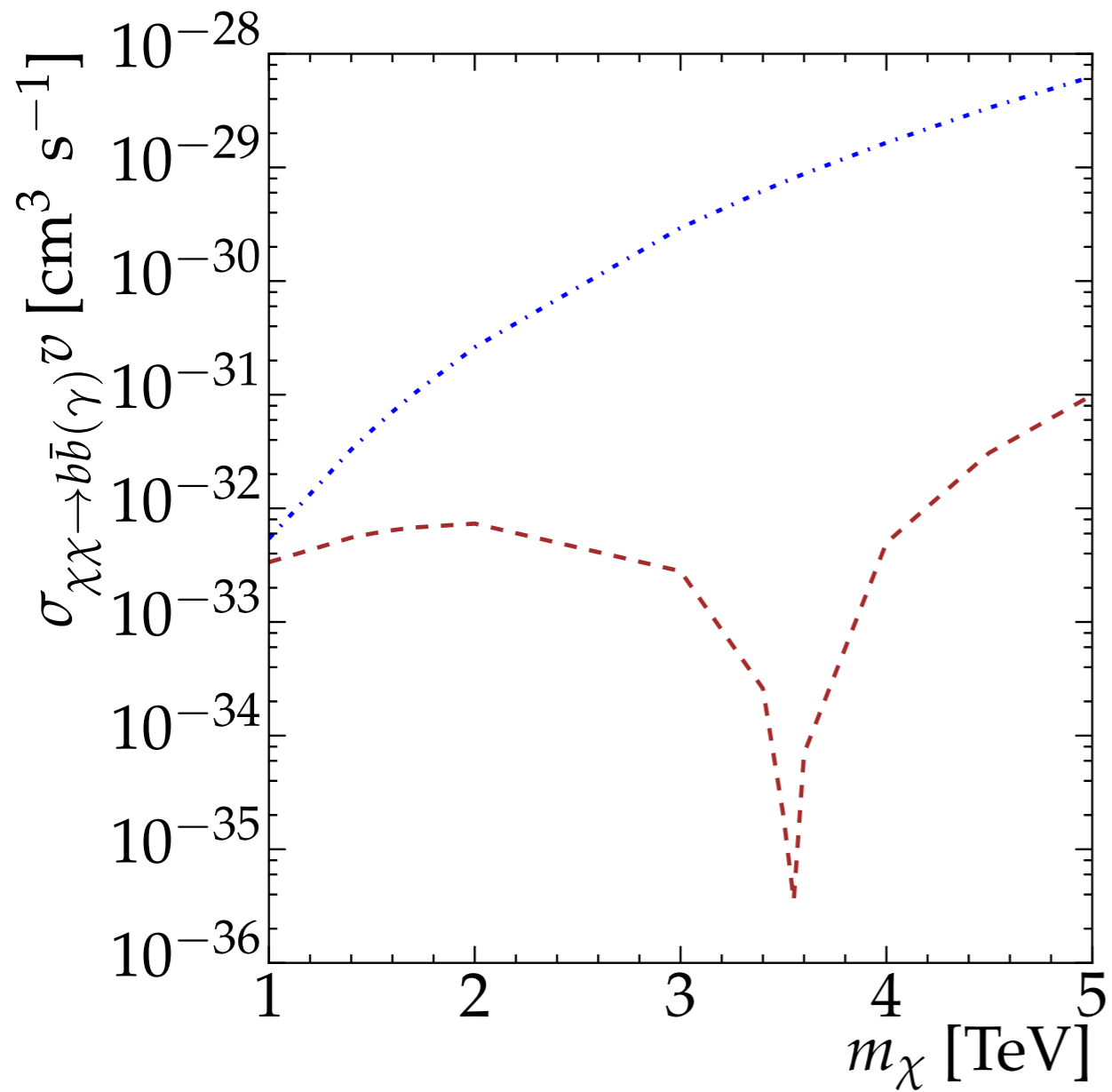


$\Lambda = 5$  TeV



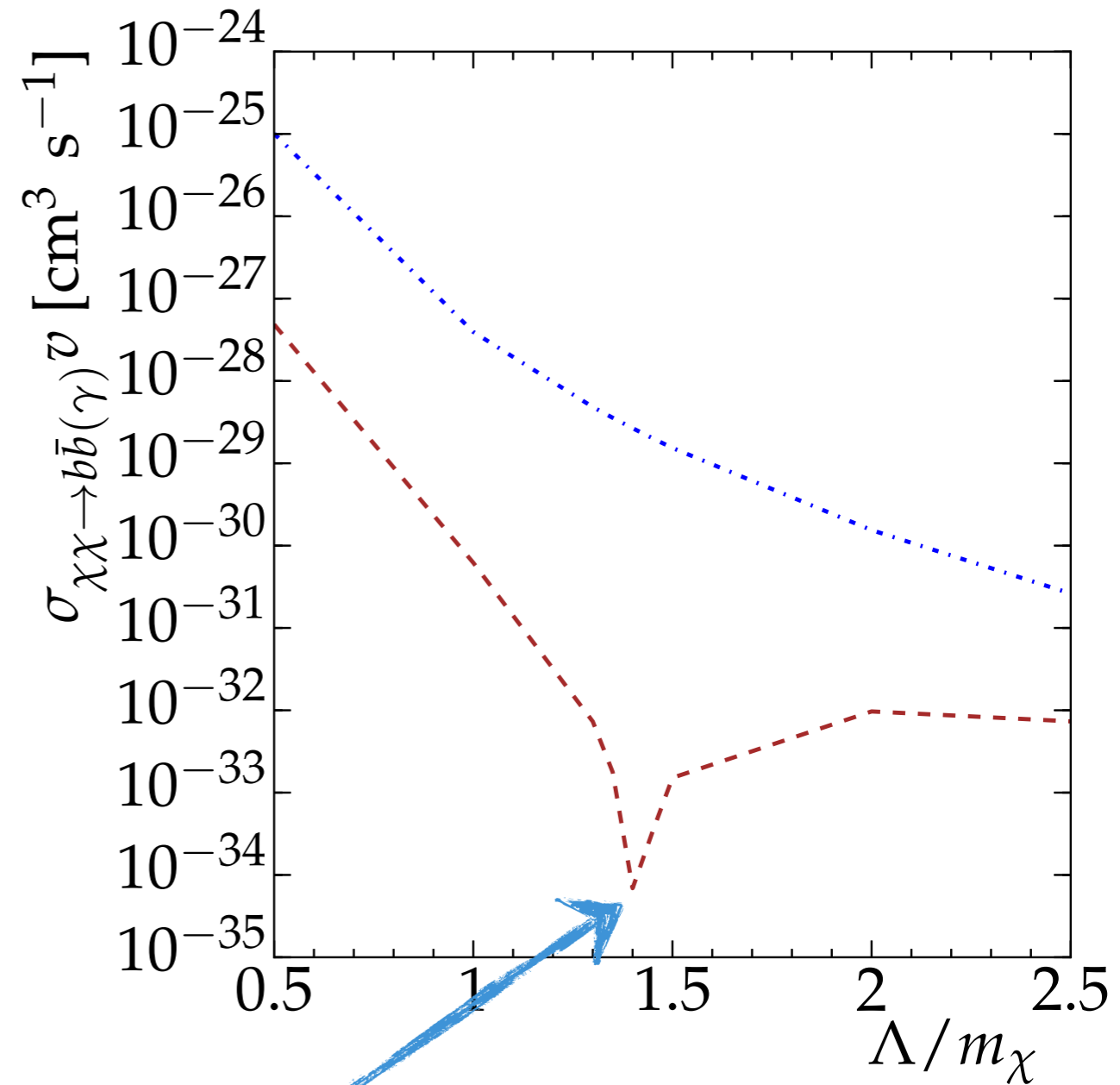
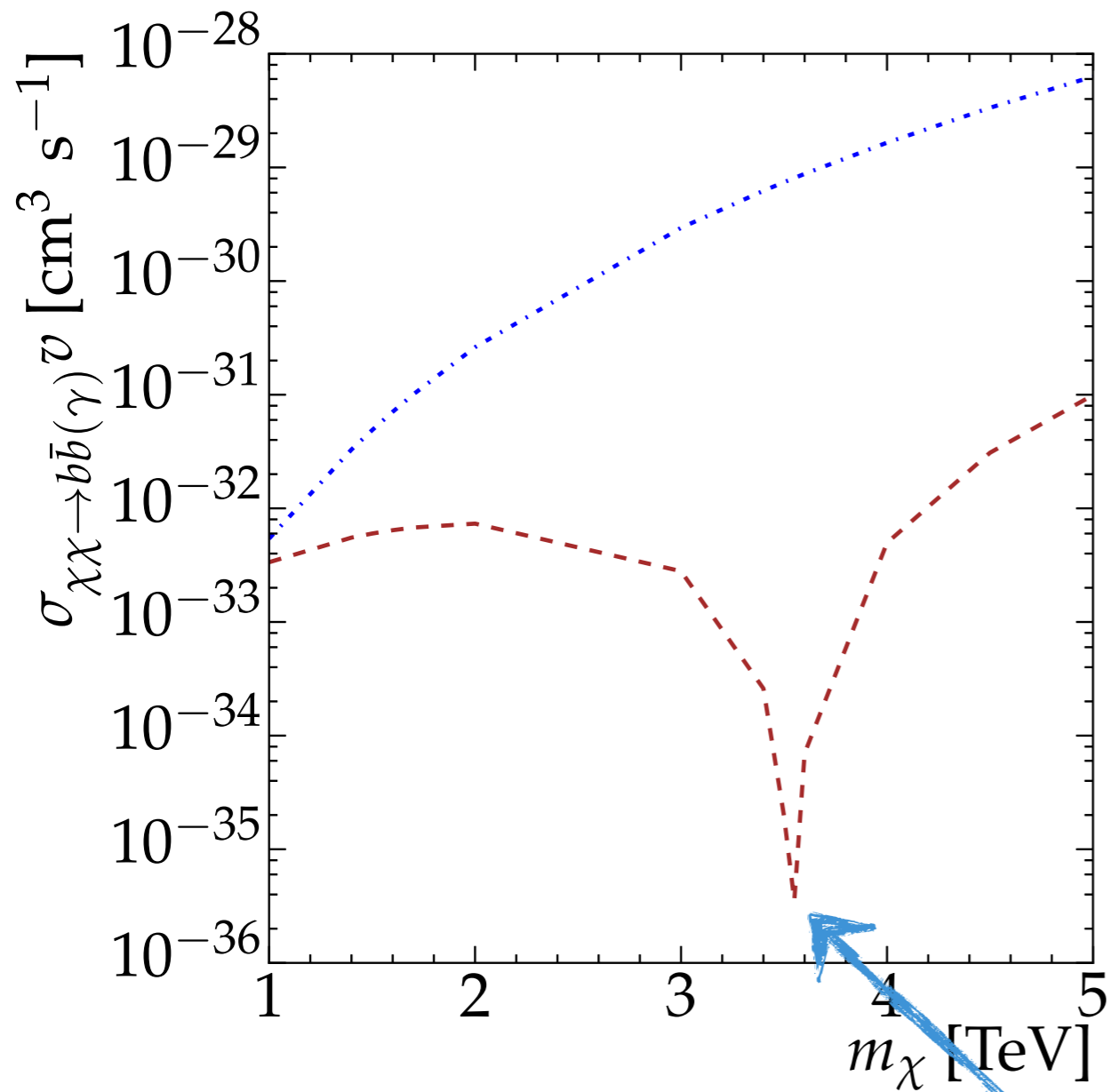
$m_\chi = 2$  TeV

# Ratio of the cross-sections





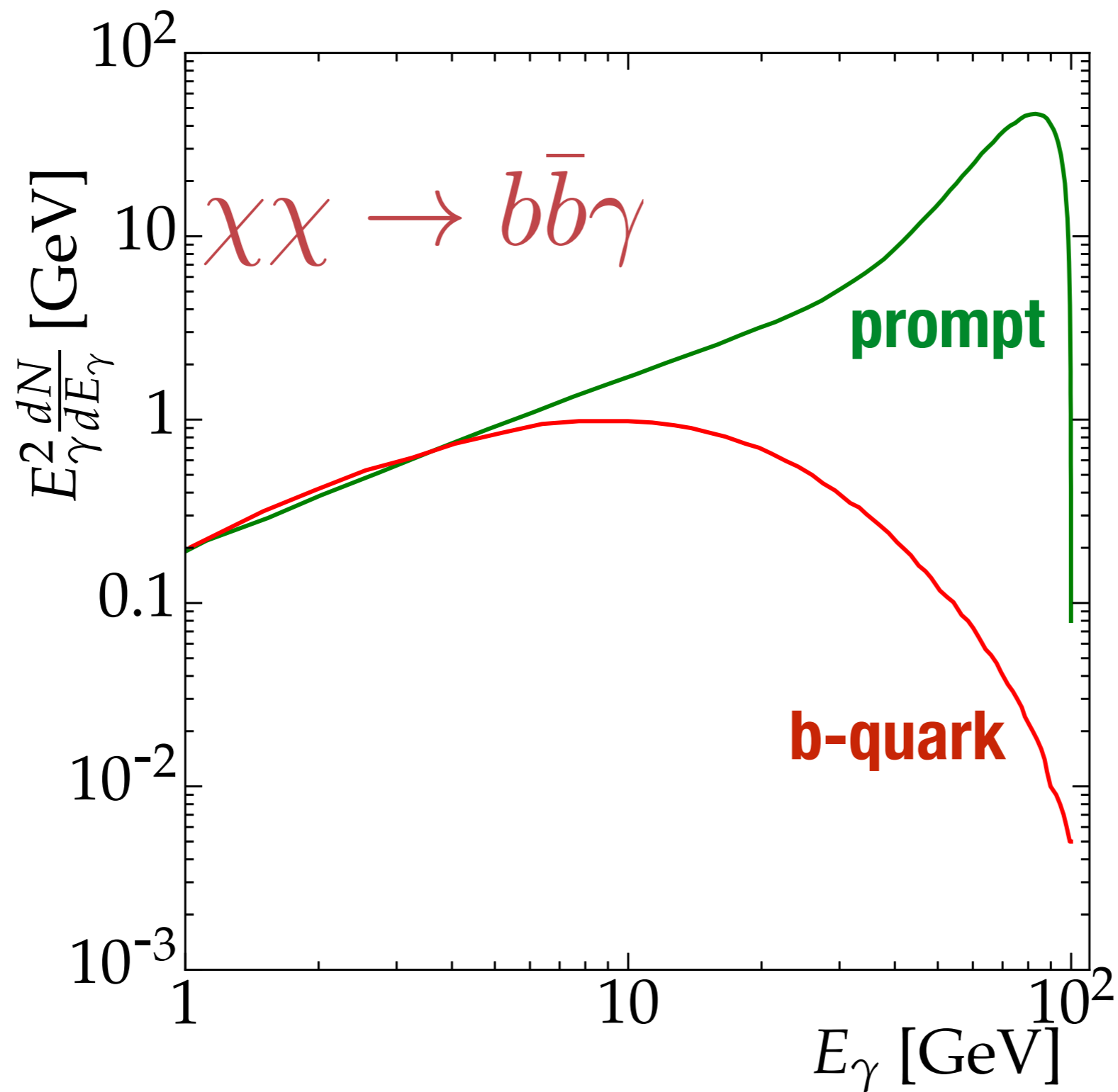
# Ratio of the cross-sections



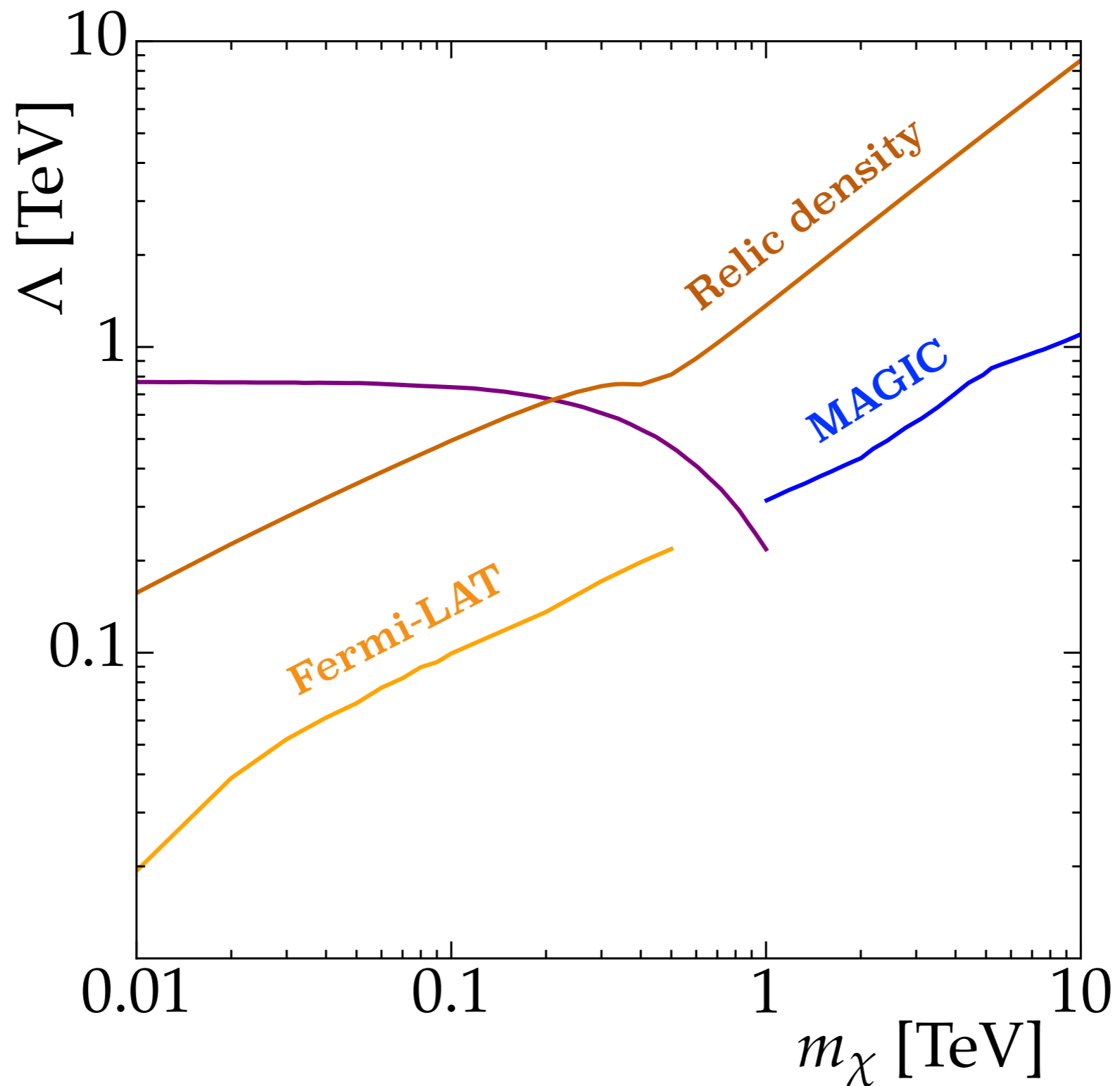
$$\Lambda/m_\chi \simeq 1.4$$

# Photon Energy Spectrum

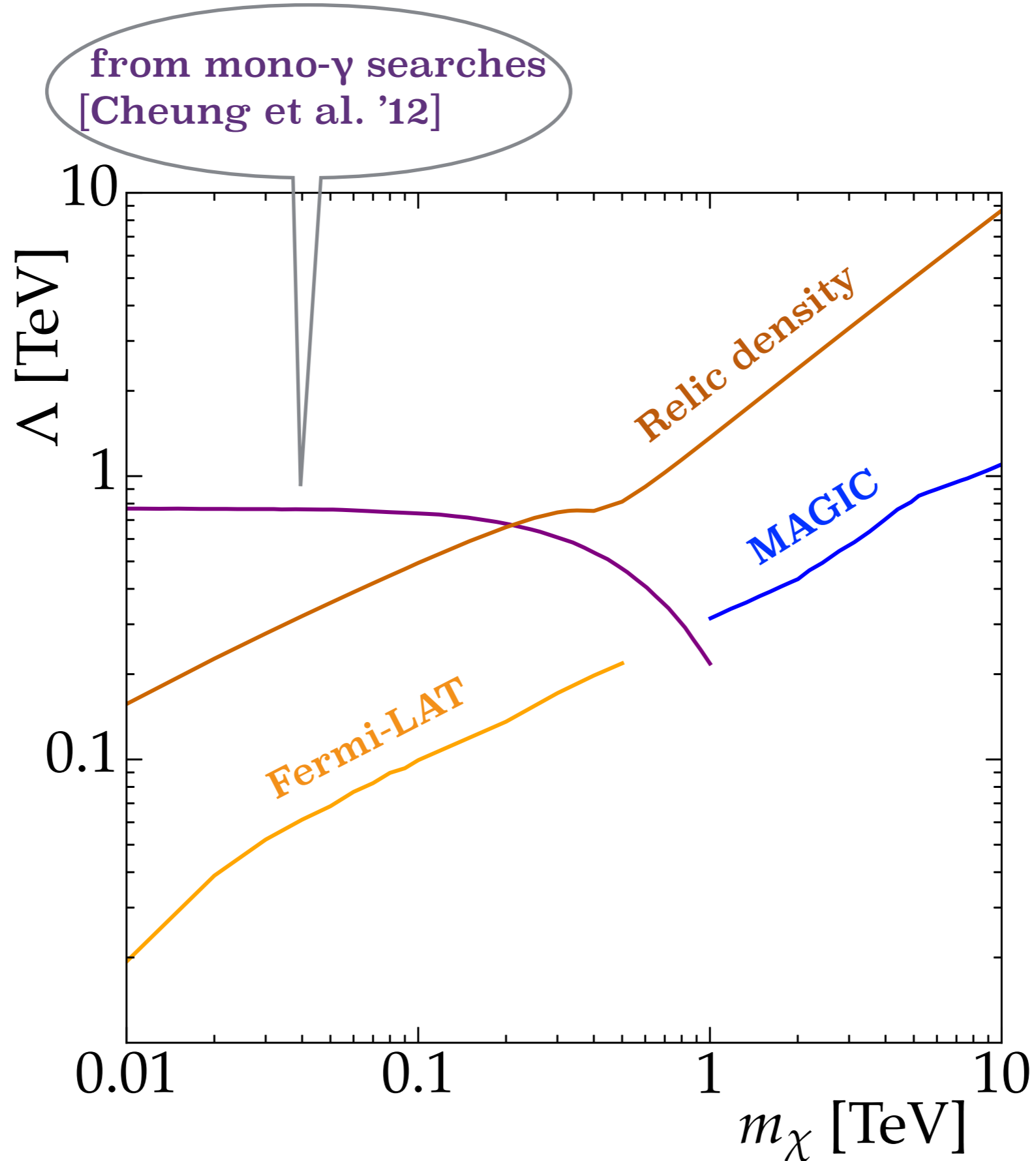
$d_6 = 1, d_8 = 1, v = 10^{-3}, M_{\text{DM}} = 100 \text{ GeV}, \Lambda = 1 \text{ TeV}$



# Bound on effective operator scale



# Bound on effective operator scale





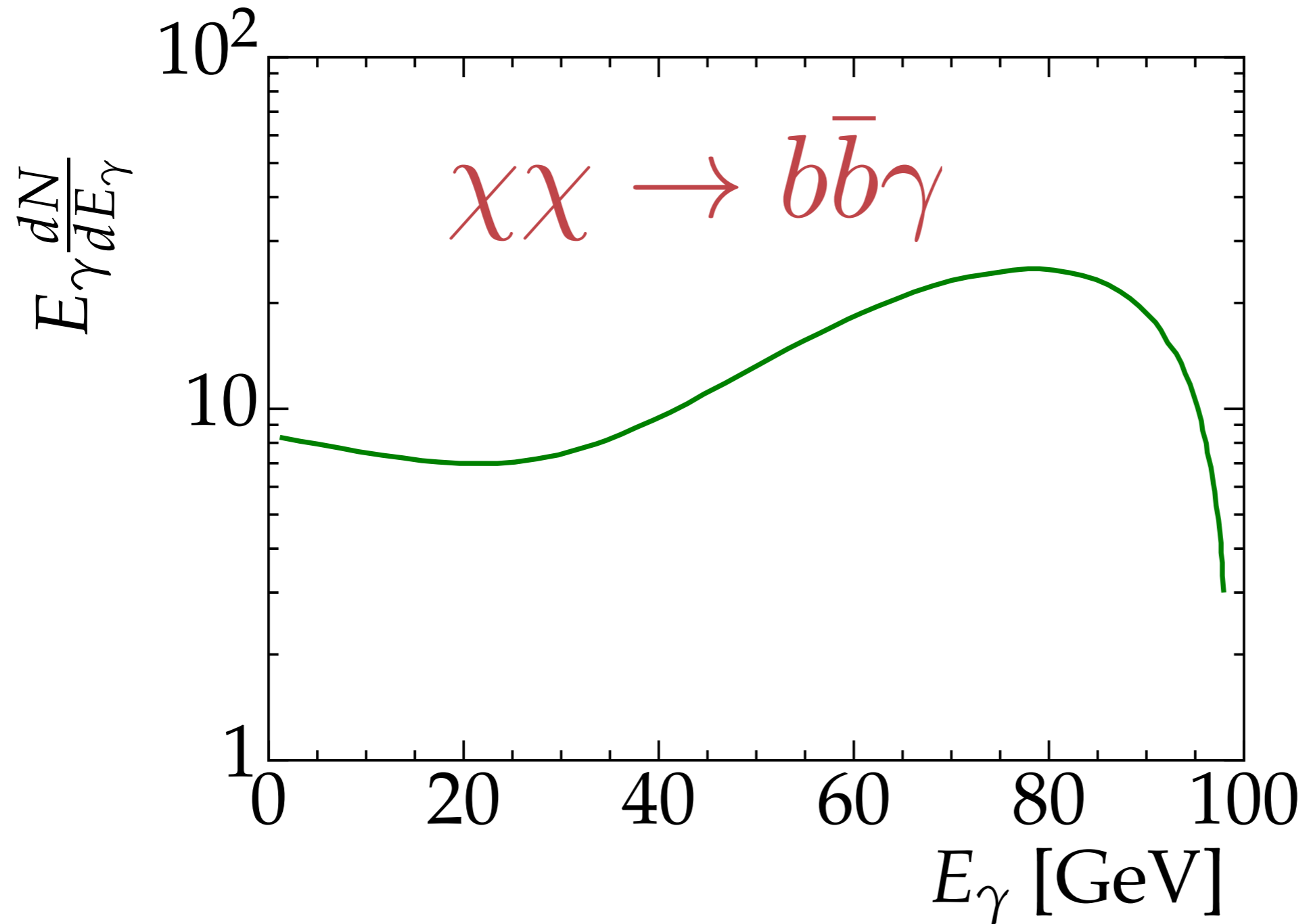
# Conclusions

- \* Electroweak bremsstrahlung lifts the helicity suppression of DM annihilation to fermions.
- \* dim-8 operator encodes this effect in EFT framework.
- \* In spite of higher dimensionality of dim-8 operator, it does not suffer any suppression from DM relative velocity.
- \* Annihilation cross-section from dim 8 operator to the dim 6 operator is always larger at all dark matter mass scales  $> 1$  TeV.
- \* Cancellation in the 2-body cross-section between dim 6 and dim 8 for  $\Lambda/m_\chi \simeq 1.4$
- \* Bounds on the  $\gamma$ -ray flux and relic density translates into stringent bound on the EFT scale.
- \* For high mass DM ( $> 200$  GeV), relic density provides the strongest bound on  $\Lambda$ .

*Thank You!!*

# Photon Energy Spectrum

$d_6 = 1, d_8 = 1, v = 10^{-3}, M_{\text{DM}} = 100 \text{ GeV}, \Lambda = 1 \text{ TeV}$



# Photon Energy Spectrum

$d_6 = 1, d_8 = 1, v = 10^{-3}, M_{\text{DM}} = 100 \text{ GeV}, \Lambda = 1 \text{ TeV}$

