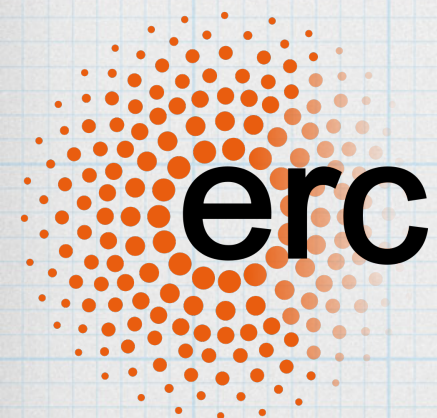


Vacuum Stability Bounds on Lepton Flavor Violating Tri-linear Soft-terms in the General MSSM

Debtosh Chowdhury
INFN, Rome

in collaboration w/ O. Eberhardt, A. Paul, and L. Silvestrini
[arXiv:1607.0XXXX]

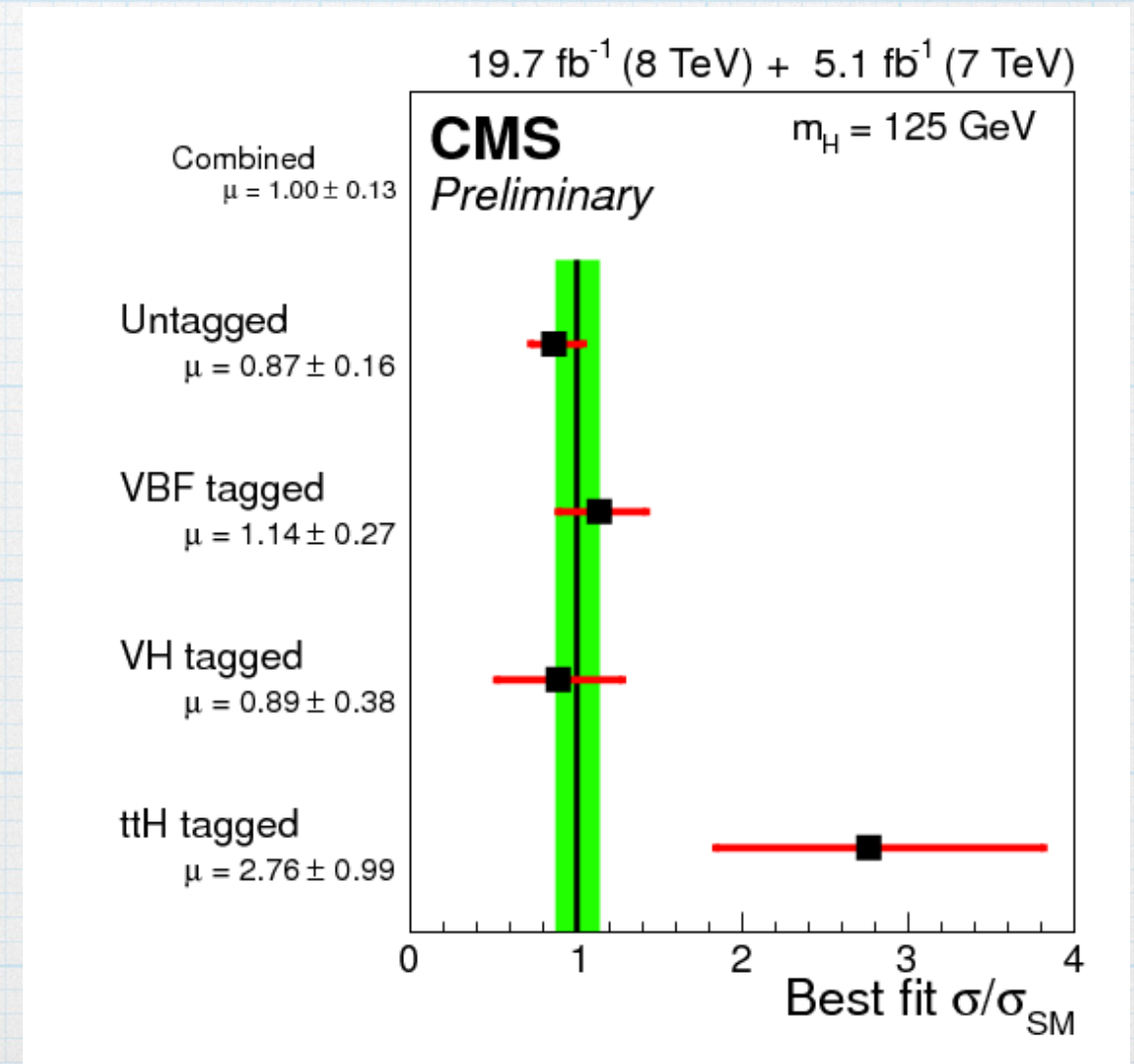
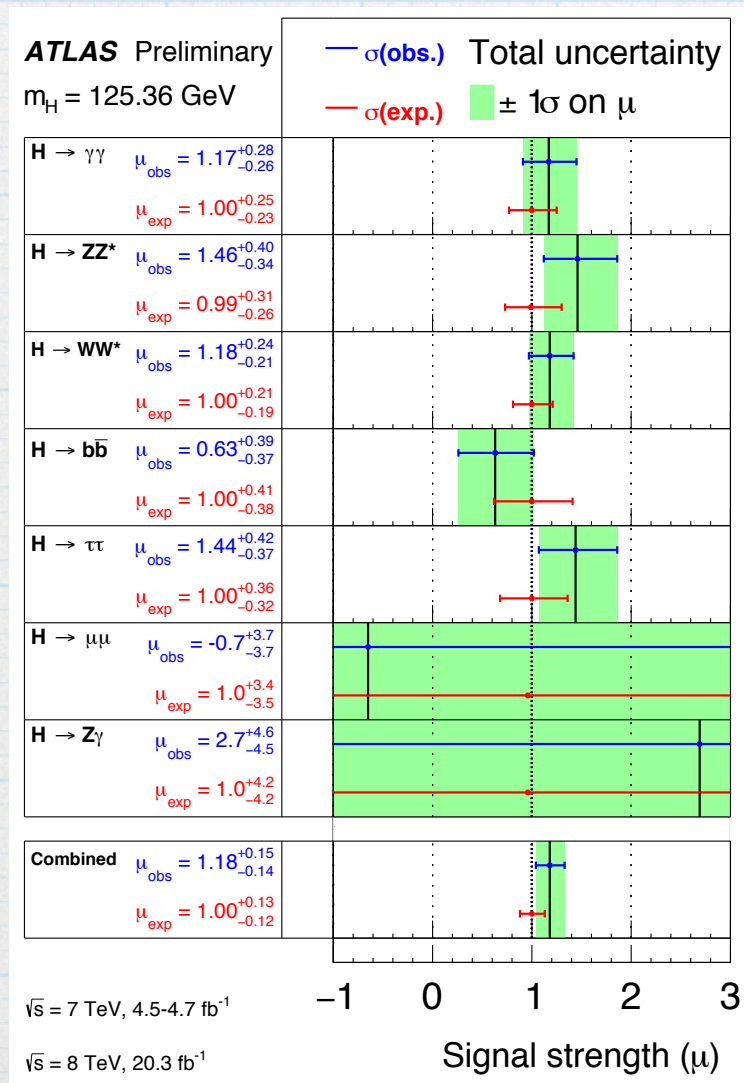


SUSY 2016
Melbourne, Australia
July 4, 2016



Higgs Discovery at LHC

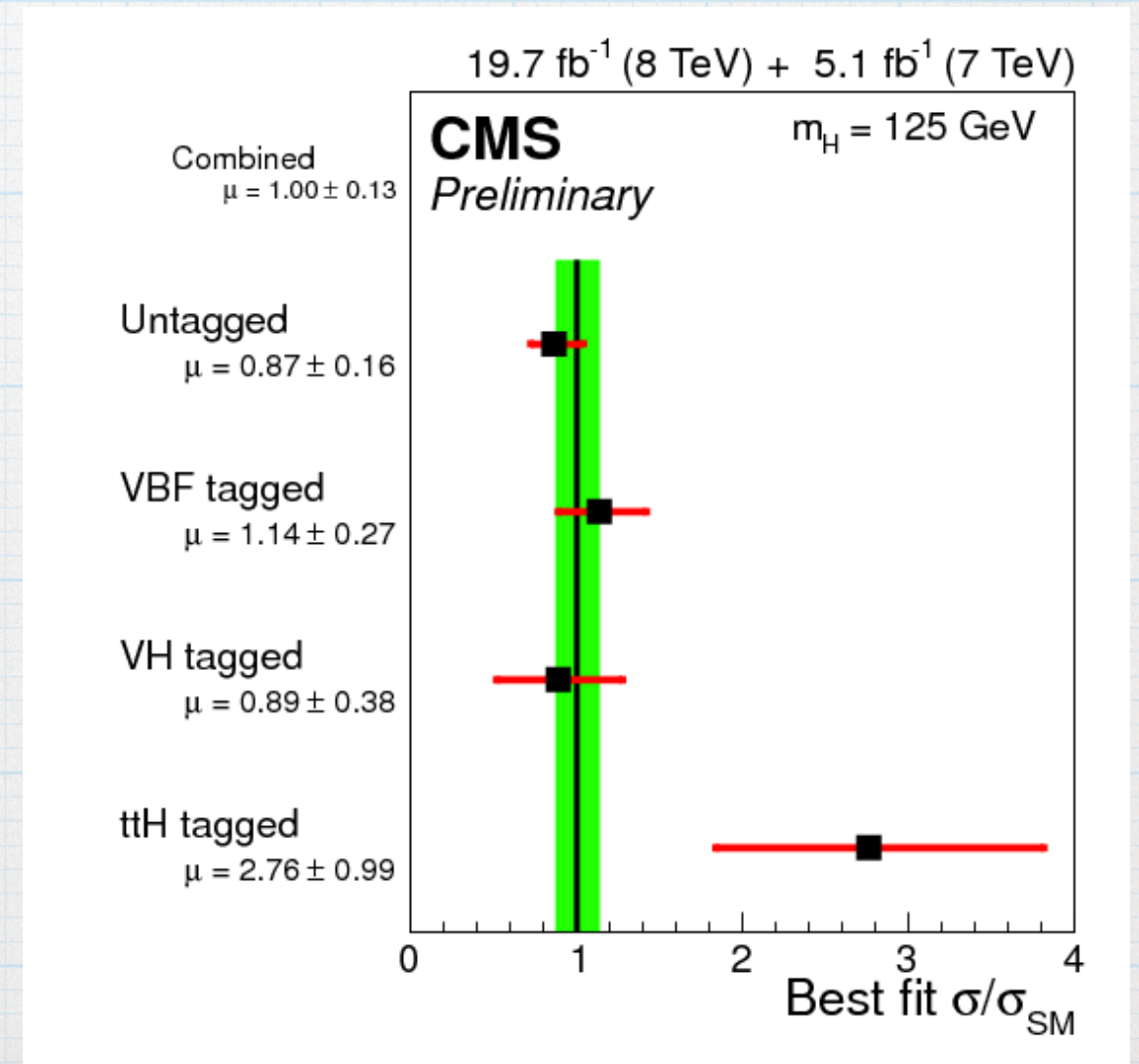
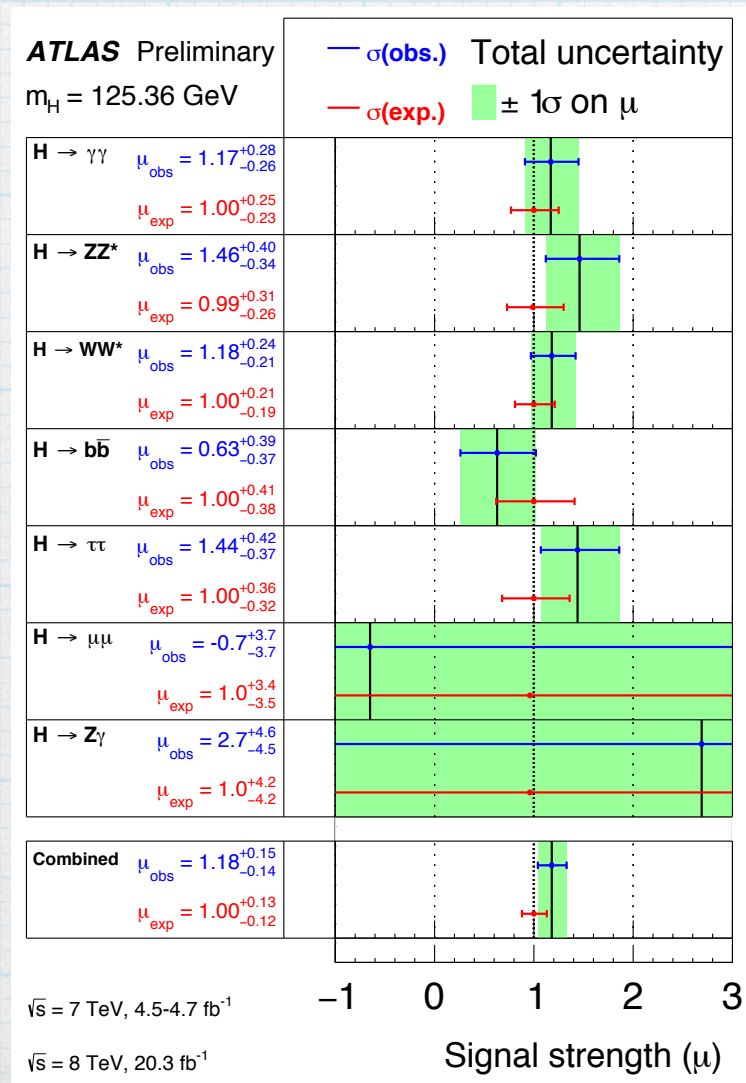
- * Mass of the Higgs is around 125 GeV
- * Higgs is parity even and spin-0
- * Higgs properties in Run-1 look “SM-like”



Higgs Discovery at LHC

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Discovery of 125 GeV Higgs at LHC puts constraints on the MSSM parameters.

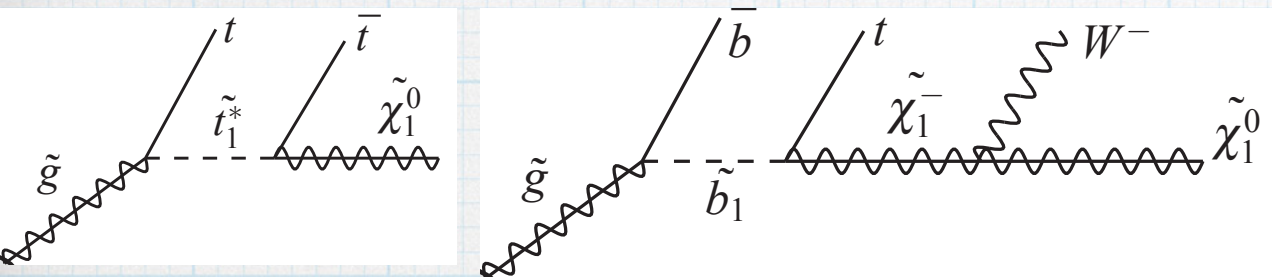


Testing SUSY

Testing SUSY

Direct Search

producing sparticles
at the colliders like LHC, ILC



insists constraints on
sparticle masses like

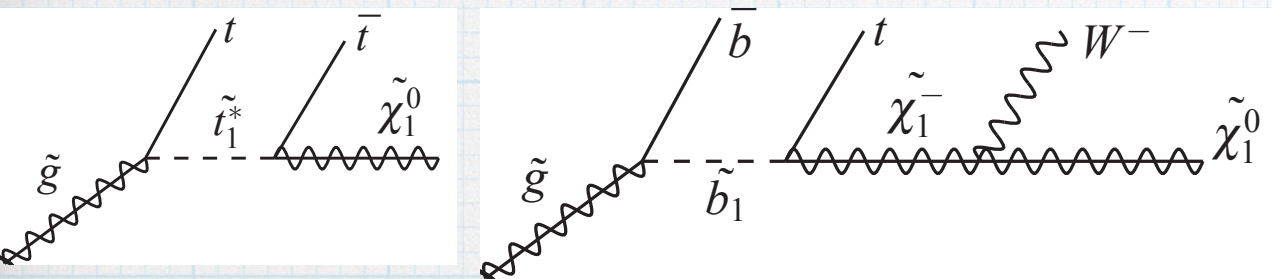
- gluinos are ruled out up to masses 1–1.25 TeV
- stops, sbottoms are > 600 GeV
- first two generations of squarks are > 0.9 TeV

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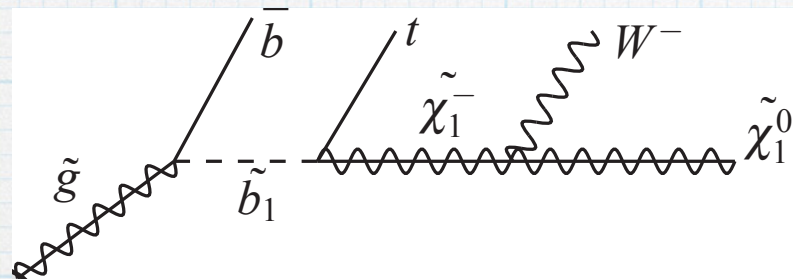
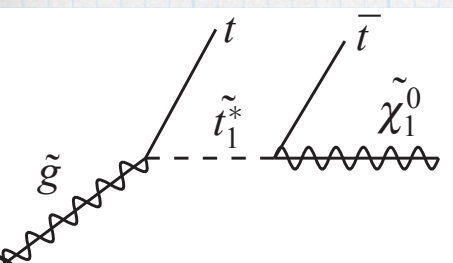
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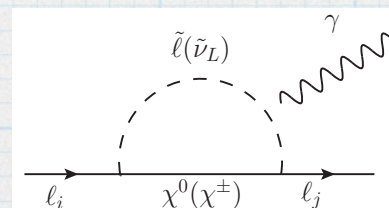
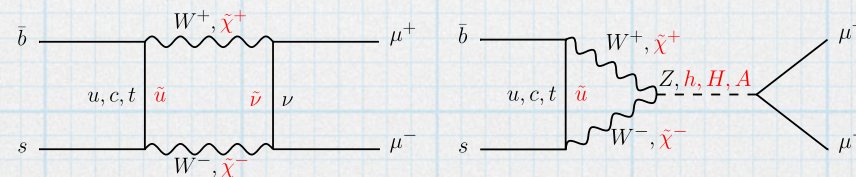
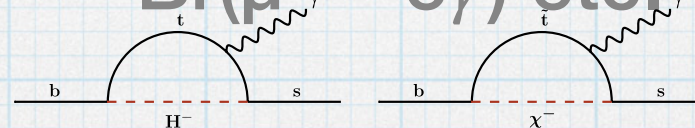
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Flavor

sparticles in the
loops
of the SM flavor
processes

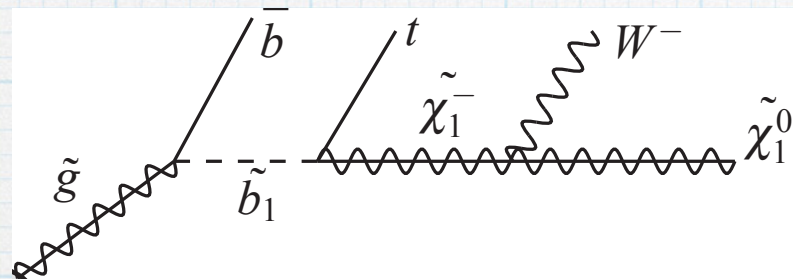
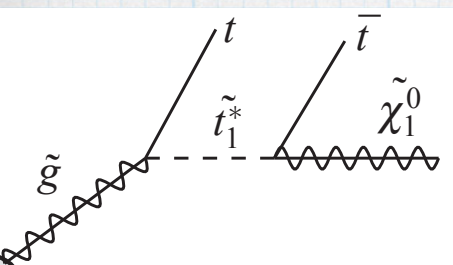
like, $\text{Br}(b \rightarrow s\gamma)$,
 $\text{Br}(B_s \rightarrow \mu\mu)$,
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Testing SUSY

Direct Search

producing sparticles
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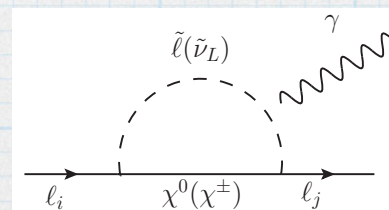
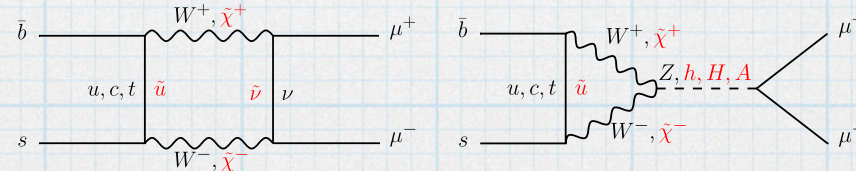


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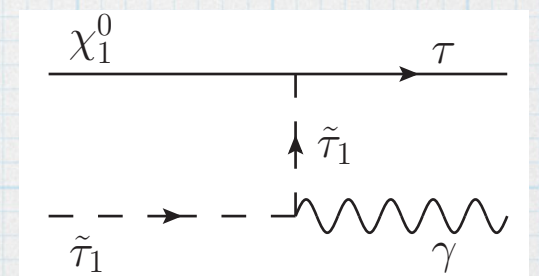
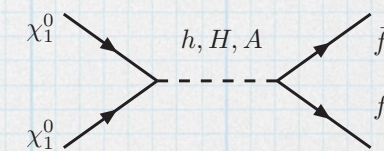
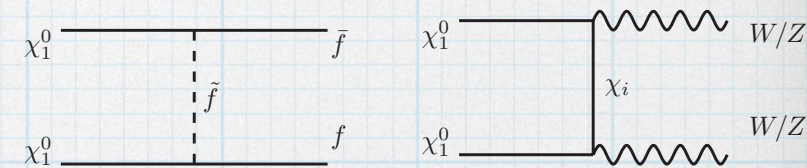
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Indirect Search

Flavor

Dark Matter

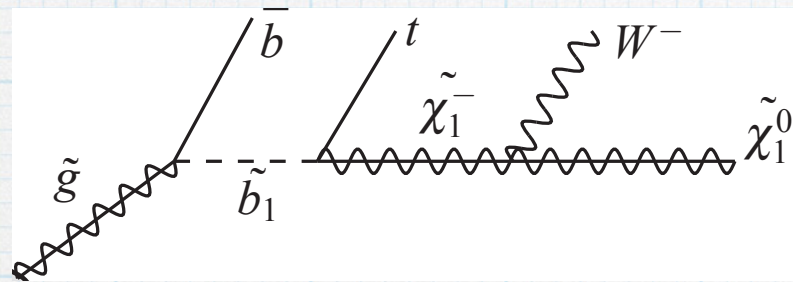
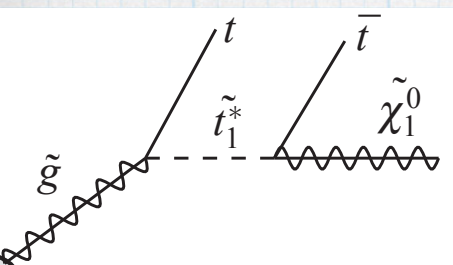
presence of new
vertices \rightarrow efficient
annihilation



Testing SUSY

Direct Search

producing sparticles
at the colliders like LHC, ILC

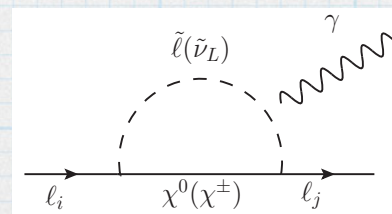
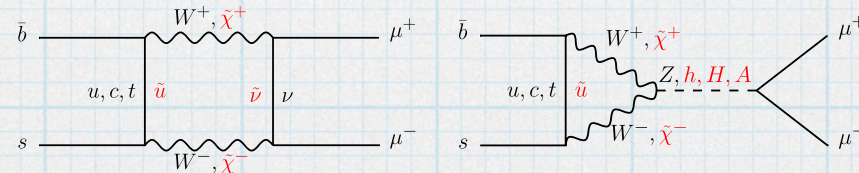


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Indirect Search

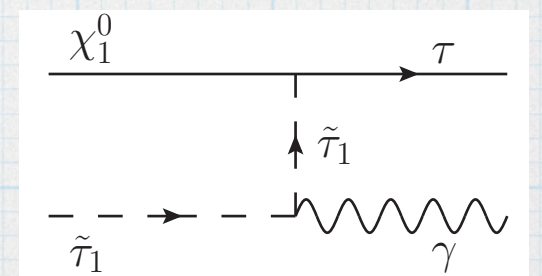
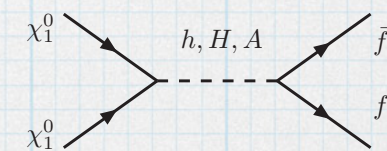
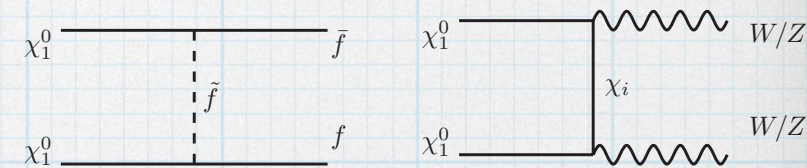
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puts constraints on
masses and couplings

SUSY Flavor Violation

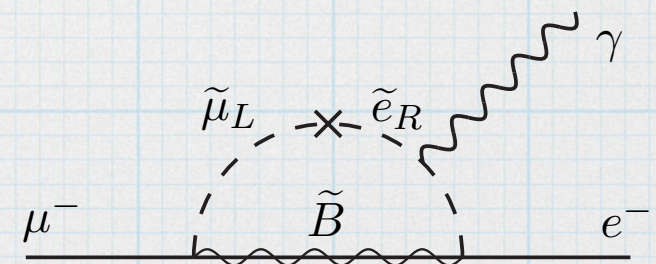
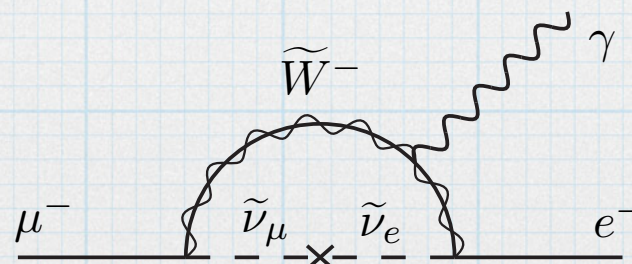
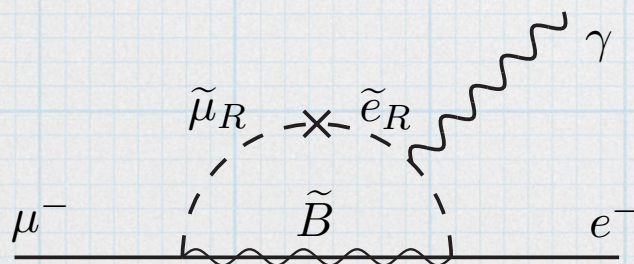
- In exact SUSY, equal masses for particles and its super-partners.

$$\mathcal{L}_{\text{soft}}^{\text{MSSM}} = -\frac{1}{2} \left(M_3 \tilde{g} \tilde{g} + M_2 \tilde{W} \tilde{W} + M_1 \tilde{B} \tilde{B} + \text{c.c.} \right) \\ - (\tilde{u} \mathbf{a}^u \tilde{Q} H_u - \tilde{d} \mathbf{a}^d \tilde{Q} H_d - \tilde{e} \mathbf{a}^l \tilde{L} H_d + \text{c.c.}) \\ - \tilde{Q}^\dagger \mathbf{m}_Q^2 \tilde{Q} - \tilde{L}^\dagger \mathbf{m}_L^2 \tilde{L} - \tilde{u} \mathbf{m}_u^2 \tilde{u}^\dagger - \tilde{d} \mathbf{m}_d^2 \tilde{d}^\dagger - \tilde{e} \mathbf{m}_e^2 \tilde{e}^\dagger \\ - m_{H_u}^2 H_u^* H_u - m_{H_d}^2 H_d^* H_d - (b H_u H_d + \text{c.c.}) .$$

- To realize in nature, SUSY must be broken.

- 3 kinds of soft-breaking terms introduced in the Lagrangian.

- soft-masses and tri-linear couplings => New sources of flavor violation. Flavor-changing neutral currents (e.g. $b \rightarrow s\gamma$, $\mu \rightarrow e\gamma$)



Charge breaking constraint

Soft-breaking tri-linear term in the Lagrangian

$$\mathcal{L}_{\text{soft}} = -\tilde{\bar{e}} a^l \tilde{L} H_d + \text{c.c.}$$

In the presence of H_d , $\tilde{\tau}_L$ and $\tilde{\mu}_R$

$$\begin{aligned} \mathcal{V} = & (m_{H_d}^2 + \mu^2) |H_d|^2 + m_{\mu_R}^2 |\tilde{\mu}_R|^2 + m_{\tau_L}^2 |\tilde{\tau}_L|^2 \\ & - (A_{23}^l H_d \tilde{\mu}_R^* \tilde{\tau}_L + \text{c.c.}) + Y_\tau^2 |H_d \tilde{\tau}_L|^2 + Y_\mu^2 |H_d \tilde{\mu}_R|^2 \\ & + \frac{g_1^2}{8} (2|\tilde{\mu}_R|^2 - |\tilde{\tau}_L|^2 - |H_d|^2)^2 + \frac{g_2^2}{8} (|\tilde{\tau}_L|^2 - |H_d|^2)^2 \end{aligned}$$

the CCB minima of the potential around the D-flat direction

$$|H_d| = |\tilde{\tau}_L| = |\tilde{\mu}_R| \equiv a$$

$$\mathcal{V} = (m_{H_d}^2 + \mu^2 + m_{\mu_R}^2 + m_{\tau_L}^2) a^2 - 2|A_{23}^l| a^3 + (Y_\mu^2 + Y_\tau^2) a^4$$

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$$a \sim 2 \frac{|A_{23}^l|}{Y_\tau^2}$$

absolute stability

$$|A_{23}^l|^2 \leq Y_\tau^2 (m_{H_d}^2 + \mu^2 + m_{\mu_R}^2 + m_{\tau_L}^2)$$

$$\frac{\Gamma}{V} = A e^{-S[\bar{\phi}]}$$

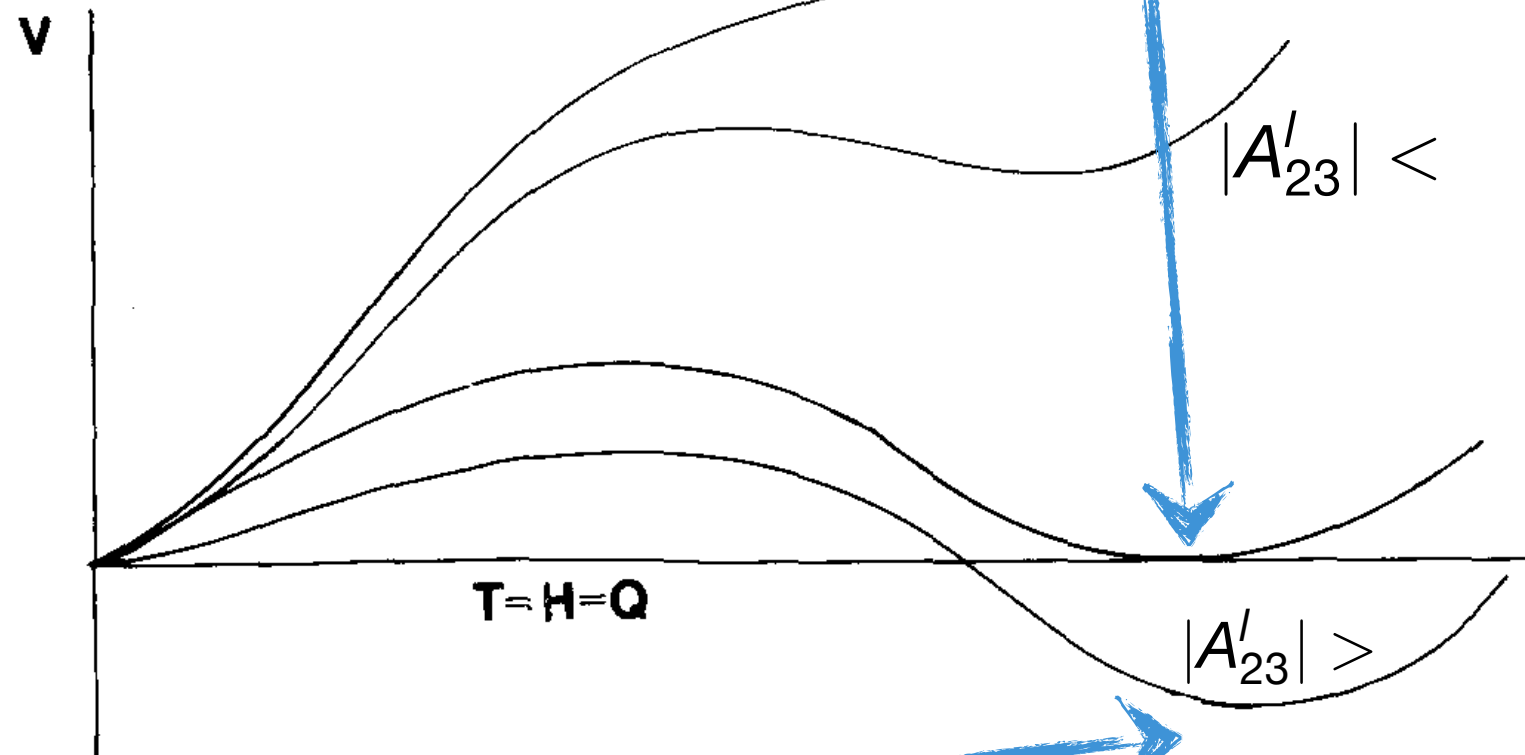
meta-stability: $\Gamma/V \sim H_0^4$

$$S[\bar{\phi}] \gtrsim 400$$

Coleman et al. '77

meta-stability

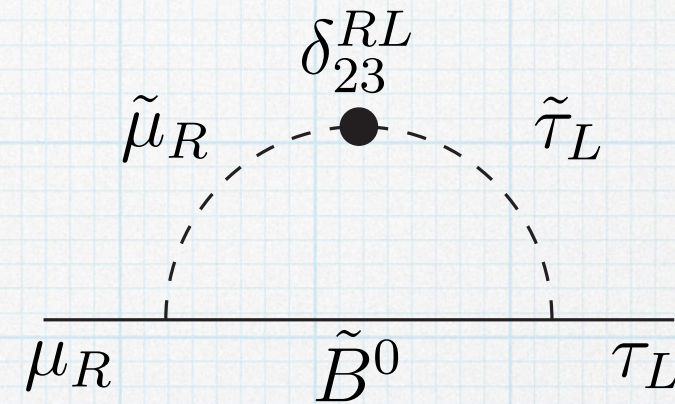
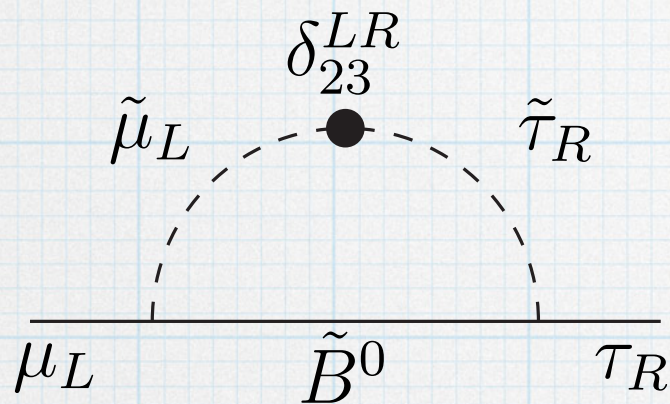
Claudson et al. '83



Quiros '94, Kusenko et al. '96, Carena et al. '96

$$\text{BR}(\tau \rightarrow \mu \gamma)$$

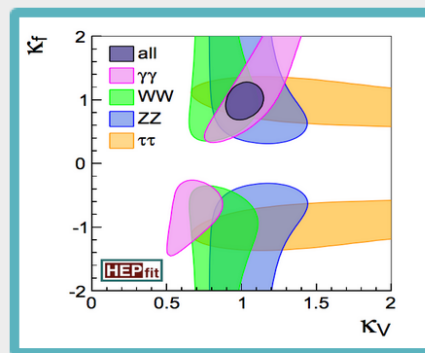
In the presence of A_{23}^l and A_{32}^l the branching fraction is



$$\text{BR}(\tau \rightarrow \mu \gamma) \propto \frac{M_1^2 \tan^2 \beta}{|\mu|^2} \left[\left(\delta_{23}^{LR} \right)^2 \left| \frac{m_R m_L}{m_\tau \mu \tan \beta} I_B \right|^2 + \left(\delta_{23}^{RL} \right)^2 \left| \frac{m_R m_L}{m_\tau \mu^* \tan \beta} I_B \right|^2 \right]$$

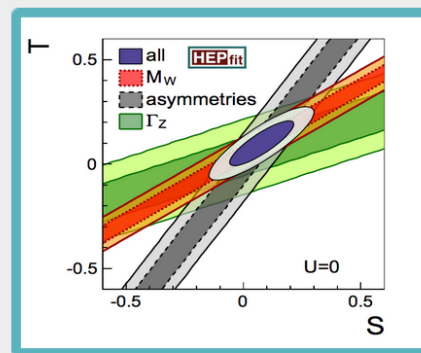
$$\delta_{23}^{RL} \equiv \frac{A_{23}^l \langle H_d^0 \rangle}{m_{\tilde{l}}^2} \quad m_{\tilde{l}} \equiv \sqrt{m_R^2 m_L^2}$$

HEPfit: a Code for the Combination of Indirect and Direct Constraints on High Energy Physics Models.



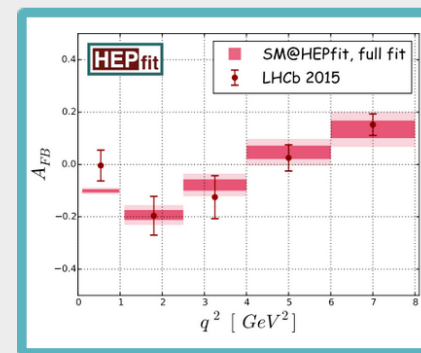
Higgs Physics

HEPfit can be used to study Higgs couplings and analyze data on signal strengths.



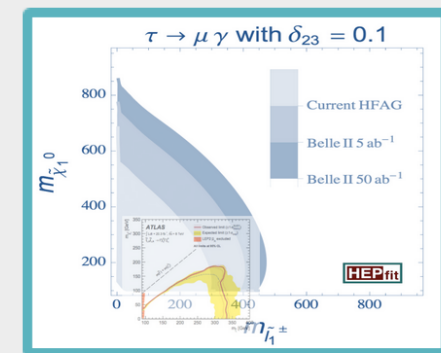
Precision Electroweak

Electroweak precision observables are included in HEPfit



Flavour Physics

The Flavour Physics menu in HEPfit includes both quark and lepton flavour dynamics.

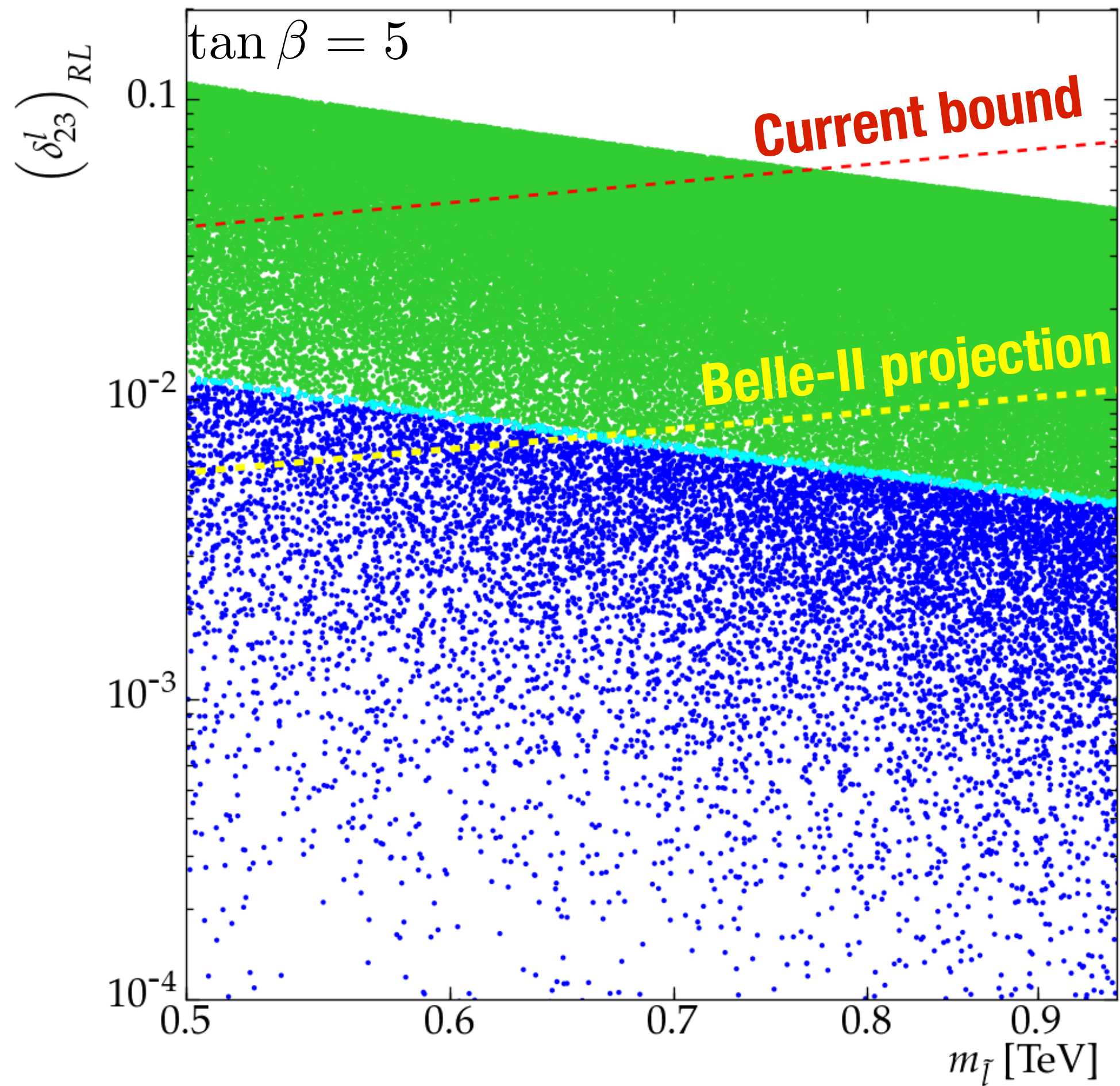


BSM Physics

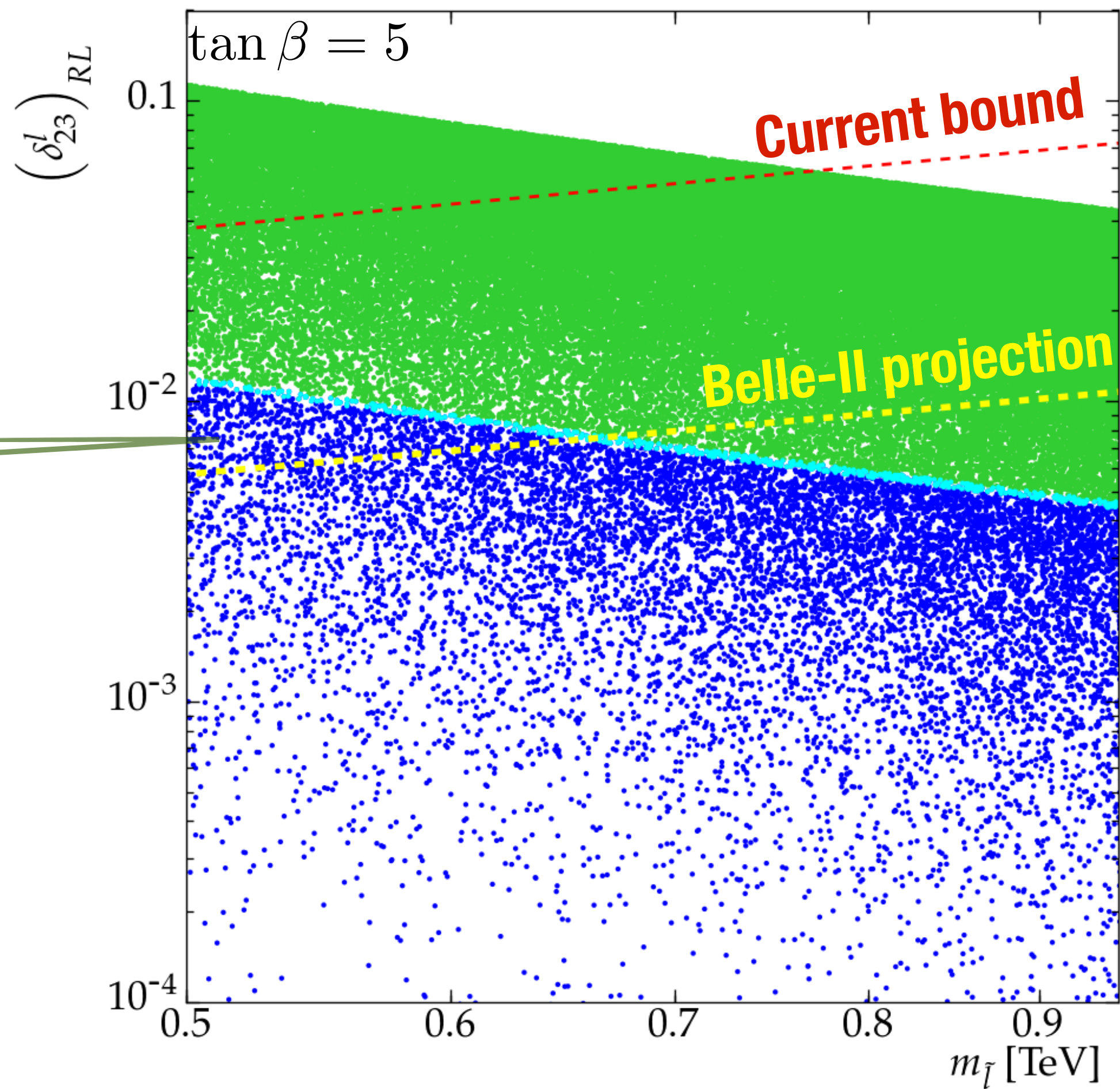
Dynamics beyond the Standard Model can be studied by adding models in HEPfit.

- **Flexible open-source C++ code to do calculations with various observables in the SM and beyond:**
 - Simple user-defined models and/or observables
 - Stand-alone or library modes to compute single observables.
 - Optional Bayesian fitting framework to do global statistical analyses (run-time optimized, parallelized; can be replaced by a different one)

For more details, look at the talk by J. de Blas this afternoon

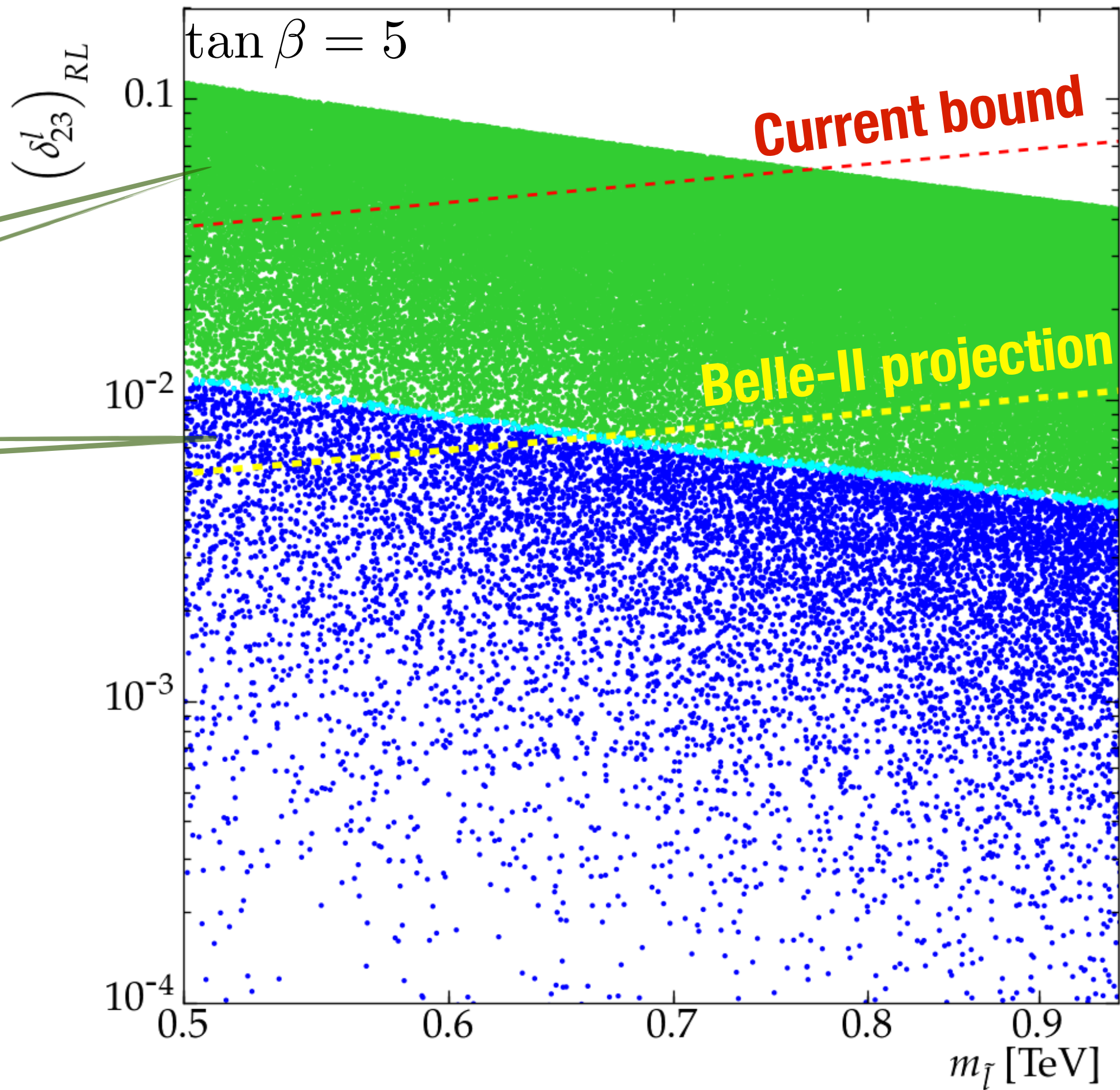


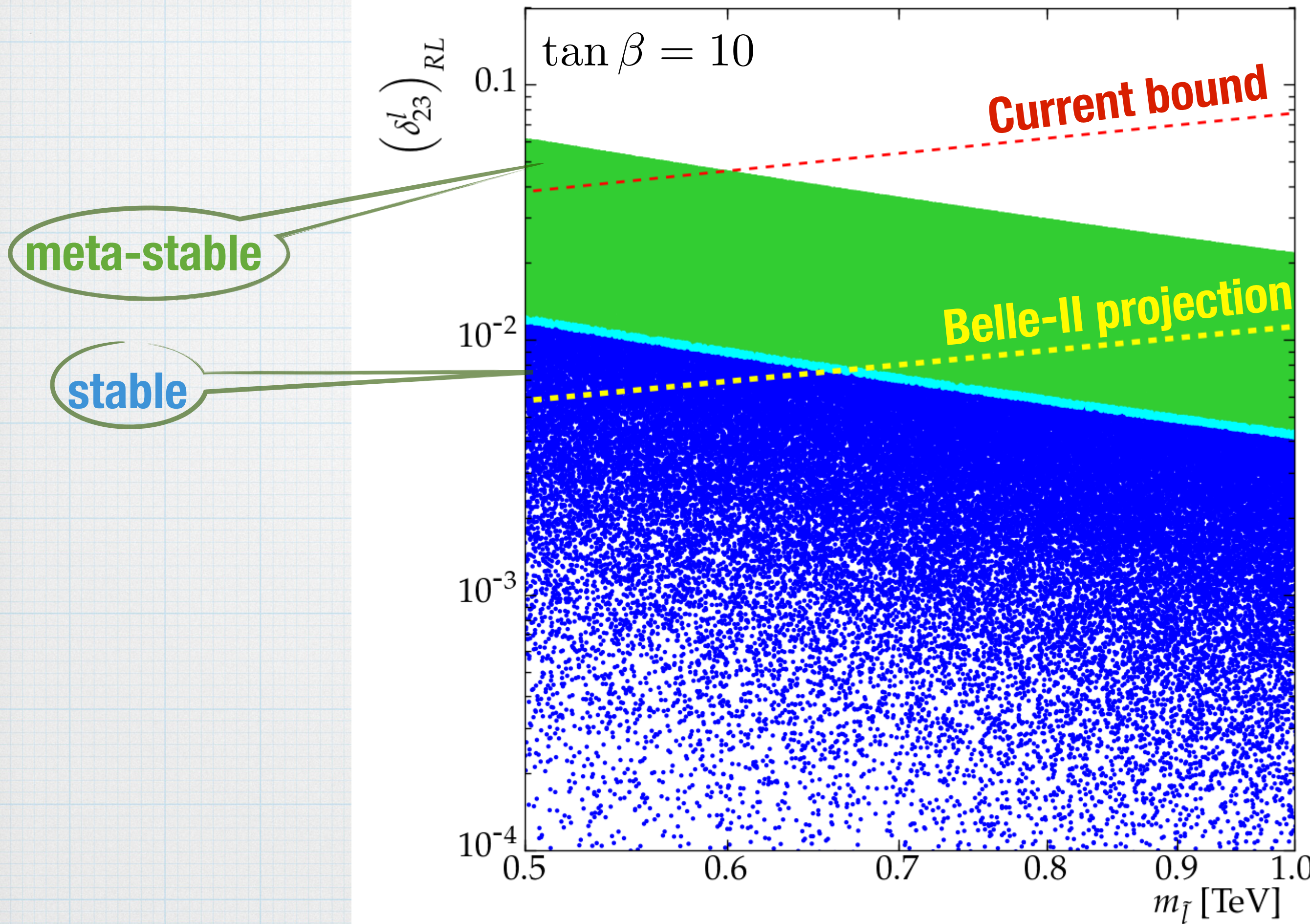
stable



meta-stable

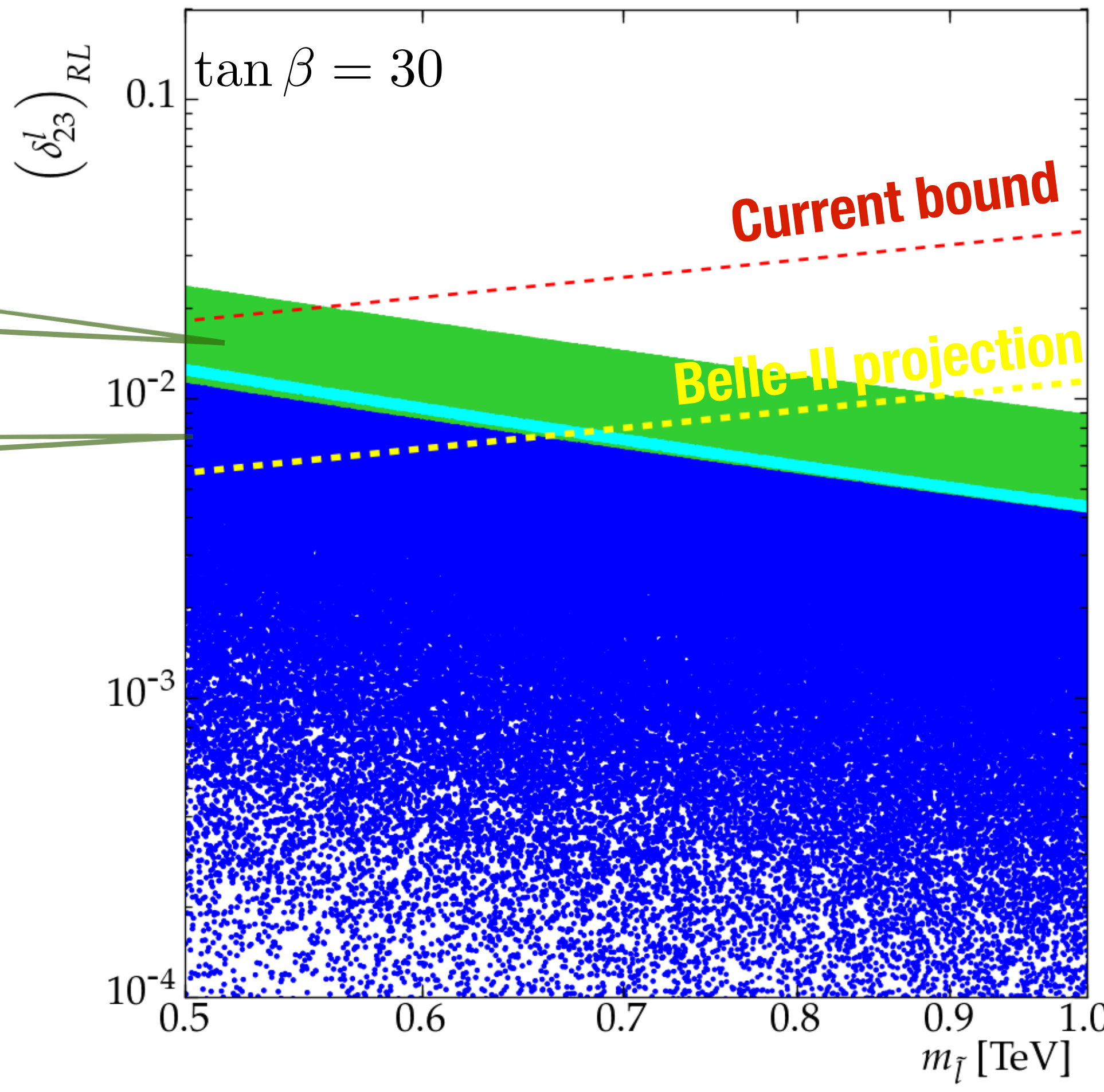
stable





meta-stable

stable



Conclusion

- * Unlike the FCNC bounds, the strength of the CCB and UFB bounds does not decrease as the scale of supersymmetry breaking increases.**
- * meta-stability relaxes the bounds on the flavor violating tri-linear couplings.**
- * Future flavor factories, e.g. Belle-II will be able to probe large fractions of the meta-stable region for A_{23}^I .**

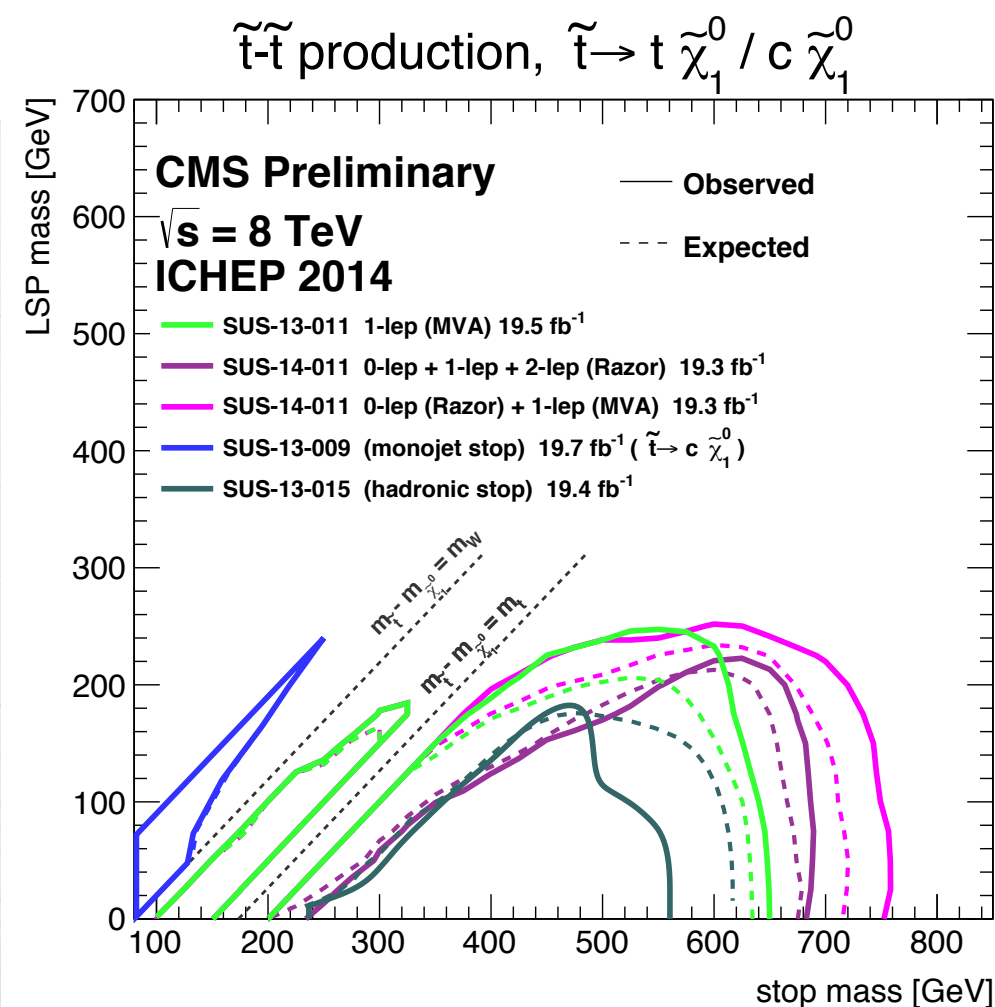
Thank You!!

Extras

ICHEP 2014

For decays with intermediate mass,

$$m_{\text{intermediate}} = x \cdot m_{\text{mother}} + (1-x) \cdot m_{\text{sp}}$$



ATLAS SUSY Searches* - 95% CL Lower Limits

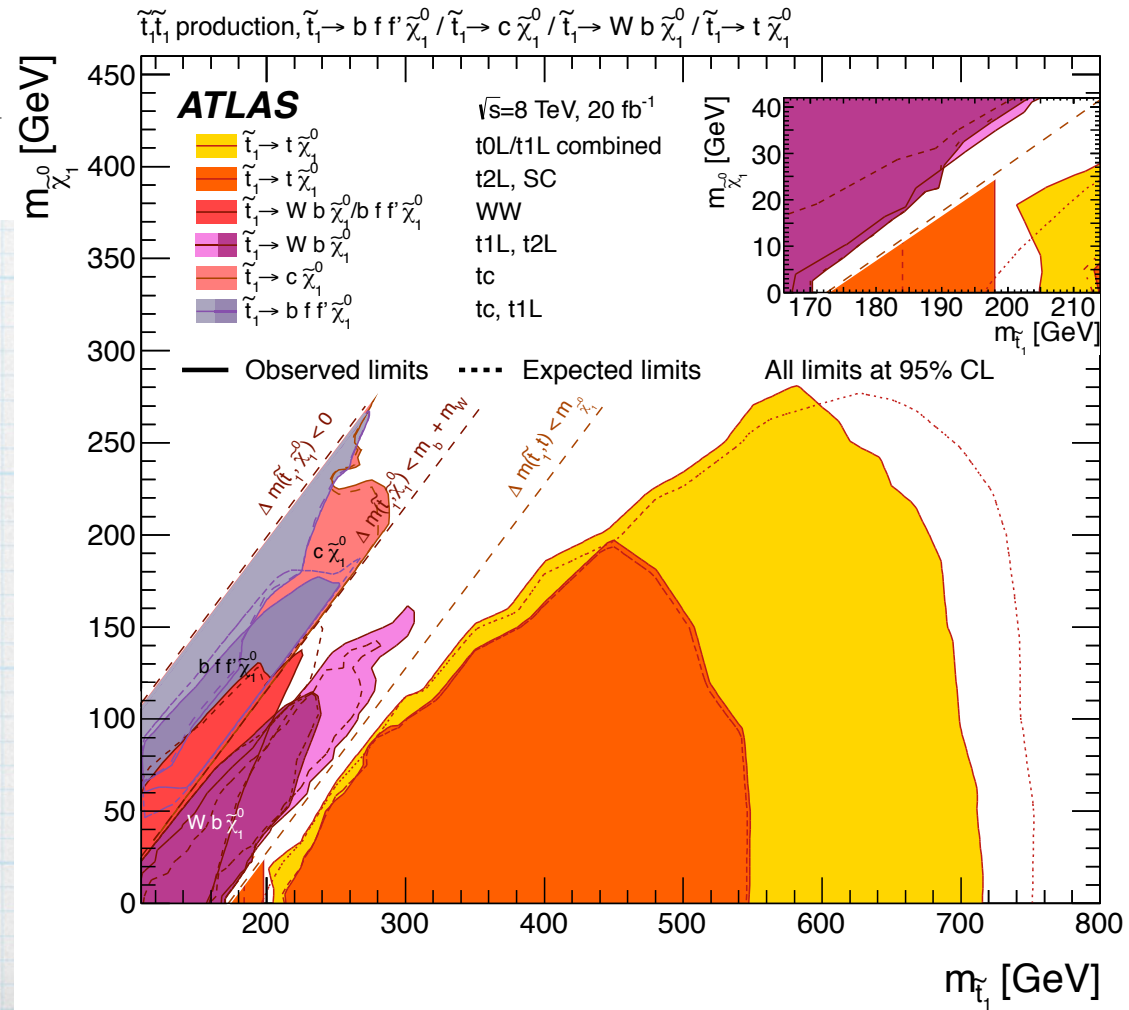
Status: July 2015

ATLAS Preliminary

$\sqrt{s} = 7, 8 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	$\sqrt{s} = 7 \text{ TeV}$	$\sqrt{s} = 8 \text{ TeV}$	Reference
Inclusive Searches	MSUGRA/CMSSM	0-3 e, μ /1-2 τ	2-10 jets/3 b	Yes	20.3	\tilde{q}, \tilde{g}	1.8 TeV	$m(\tilde{g})=m(\tilde{q})$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q}	850 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \tilde{q})=m(2^{\text{nd}} \text{ gen. } \tilde{q})$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed)	mono-jet	1-3 jets	Yes	20.3	\tilde{q}	100-440 GeV	$m(\tilde{g})=m(\tilde{\chi}_1^0)<10 \text{ GeV}$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q(\ell\ell)/(\nu\nu)\tilde{\chi}_1^0$	2 e, μ (off-Z)	2 jets	Yes	20.3	\tilde{q}	780 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g}	1.33 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0 \rightarrow q\tilde{q}W^\pm\tilde{\chi}_1^0$	0-1 e, μ	2-6 jets	Yes	20	\tilde{g}	1.26 TeV	$m(\tilde{\chi}_1^0)<300 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)/(\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20	\tilde{g}	1.32 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	GMSB ($\tilde{\ell}$ NLSP)	1-2 τ + 0-1 ℓ	0-2 jets	Yes	20.3	\tilde{g}	1.6 TeV	$\tan\beta > 20$
	GGM (bino NLSP)	2 γ	-	Yes	20.3	\tilde{g}	1.29 TeV	$c\tau(\text{NLSP})<0.1 \text{ mm}$
	GGM (higgsino-bino NLSP)	γ	1 b	Yes	20.3	\tilde{g}	1.3 TeV	$m(\tilde{\chi}_1^0)<900 \text{ GeV}, c\tau(\text{NLSP})<0.1 \text{ mm}, \mu<0$
3 rd gen. \tilde{g} & med.	GGM (higgsino-bino NLSP)	γ	2 jets	Yes	20.3	\tilde{g}	1.25 TeV	$m(\tilde{\chi}_1^0)<850 \text{ GeV}, c\tau(\text{NLSP})<0.1 \text{ mm}, \mu>0$
	GGM (higgsino NLSP)	2 e, μ (Z)	2 jets	Yes	20.3	\tilde{g}	850 GeV	$m(\text{NLSP})>430 \text{ GeV}$
	Gravitino LSP	0	mono-jet	Yes	20.3	$E^{1/2}$ scale	865 GeV	$m(\tilde{G})>1.8 \times 10^{-4} \text{ eV}, m(\tilde{g})=m(\tilde{q})=1.5 \text{ TeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g}	1.25 TeV	$m(\tilde{\chi}_1^0)<400 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g}	1.1 TeV	$m(\tilde{\chi}_1^0)<350 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	1.34 TeV	$m(\tilde{\chi}_1^0)<400 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	1.3 TeV	$m(\tilde{\chi}_1^0)<300 \text{ GeV}$
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1	100-620 GeV	$m(\tilde{\chi}_1^0)<90 \text{ GeV}$
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^0$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{b}_1	275-440 GeV	$m(\tilde{\chi}_1^0)=2 m(\tilde{\chi}_1^0)$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^0$	1-2 e, μ	1-2 b	Yes	4.7/20.3	\tilde{t}_1	110-167 GeV	$m(\tilde{\chi}_1^0) = 2m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=55 \text{ GeV}$
3 rd gen. squarks direct production	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{\nu}_t^0$	0-2 e, μ	0-2 jets/1-2 b	Yes	20.3	\tilde{t}_1	90-191 GeV	$m(\tilde{\chi}_1^0)=1 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/ c -tag	Yes	20.3	\tilde{t}_1	90-240 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0)<85 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_1	150-580 GeV	$m(\tilde{\chi}_1^0)<150 \text{ GeV}$
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_2	290-600 GeV	$m(\tilde{\chi}_1^0)<200 \text{ GeV}$
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$	90-325 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu(\tilde{\ell}\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$	140-465 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tilde{\tau}\bar{\nu})$	2 τ	-	Yes	20.3	$\tilde{\chi}_1^\pm$	100-350 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^0 \rightarrow \tilde{\ell}_1\tilde{\nu}_\ell\ell(\tilde{\nu}_\ell), \ell\tilde{\nu}_\ell\ell(\tilde{\nu}_\ell)$	3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^0$	700 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^\pm Z\tilde{\chi}_1^0$	2-3 e, μ	0-2 jets	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^0$	420 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^\pm h\tilde{\chi}_1^0, h \rightarrow b\tilde{b}/WW/\tau\tau/\gamma\gamma$	e, μ, γ	0-2 b	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^0$	250 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$
EW direct	$\tilde{\chi}_2^0\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$	4 e, μ	0	Yes	20.3	$\tilde{\chi}_2^0$	620 GeV	$m(\tilde{\chi}_2^0)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_2^0)+m(\tilde{\chi}_1^0))$
	GGM (wino NLSP) weak prod.	1 $e, \mu + \gamma$	-	Yes	20.3	\tilde{W}	124-361 GeV	$c\tau < 1 \text{ mm}$
	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$	270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)\sim 160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$
	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm$ prod., long-lived $\tilde{\chi}_1^\pm$	dE/dx trk	-	Yes	18.4	$\tilde{\chi}_1^\pm$	482 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)\sim 160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)<15 \text{ ns}$
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	27.9	\tilde{g}	832 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$
	Stable \tilde{g} R-hadron	trk	-	-	19.1	\tilde{g}	1.27 TeV	
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	1-2 μ	-	-	19.1	$\tilde{\chi}_1^0$	537 GeV	$10 < \tan\beta < 50$
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$	435 GeV	$2 < \tau(\tilde{\chi}_1^0) < 3 \text{ ns}, \text{SPS8 model}$
	$\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow e\bar{e}\nu/\mu\bar{\mu}\nu$	displ. $e\bar{e}/\mu\bar{\mu}$	-	-	20.3	$\tilde{\chi}_1^0$	1.0 TeV	$7 < c\tau(\tilde{\chi}_1^0) < 740 \text{ mm}, m(\tilde{g})=1.3 \text{ TeV}$
	GGM $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow Z\tilde{G}$	displ. vtx + jets	-	-	20.3	$\tilde{\chi}_1^0$	1.0 TeV	$6 < c\tau(\tilde{\chi}_1^0) < 480 \text{ mm}, m(\tilde{g})=1.1 \text{ TeV}$
Long-lived particles	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	-	-	20.3	$\tilde{\nu}_\tau$	1.7 TeV	$\lambda_{111}^{\tau\tau}=0.11, \lambda_{132/133/233}=0.07$
	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g}	1.35 TeV	$m(\tilde{g})=m(\tilde{q}), c\tau_{\text{LSP}}<1 \text{ mm}$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^\pm\tilde{\chi}_1^0 \rightarrow e\bar{e}\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^\pm$	750 GeV	$m(\tilde{\chi}_1^0)>0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{121}\neq 0$
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^\pm\tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^\pm$	450 GeV	$m(\tilde{\chi}_1^0)>0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{133}\neq 0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	6-7 jets	-	20.3	\tilde{g}	917 GeV	$\text{BR}(\tilde{g})-\text{BR}(\tilde{b})-\text{BR}(\tilde{c})=0\%$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\tilde{q}q$	0	6-7 jets	-	20.3	\tilde{g}	870 GeV	$m(\tilde{\chi}_1^0)=600 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\tilde{q}q$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{g}	850 GeV	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{s}$	0	2 jets + 2 b	-	20.3	\tilde{t}_1	100-308 GeV	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\ell}$	2 e, μ	2 b	-	20.3	\tilde{t}_1	0.4-1.0 TeV	$\text{BR}(\tilde{t}_1 \rightarrow b\ell/\mu\nu)>20\%$
	Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	2 c	Yes	20.3	\tilde{c}	490 GeV	$m(\tilde{\chi}_1^0)<200 \text{ GeV}$
Other								1501.01325

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.



- The Standard model has passed almost all the experimental tests over decades.

- In spite of these successes SM has quite a few shortcomings, like

- Neutrino mass
- Dark matter
- Baryon asymmetry of the Universe
- Fine tuning

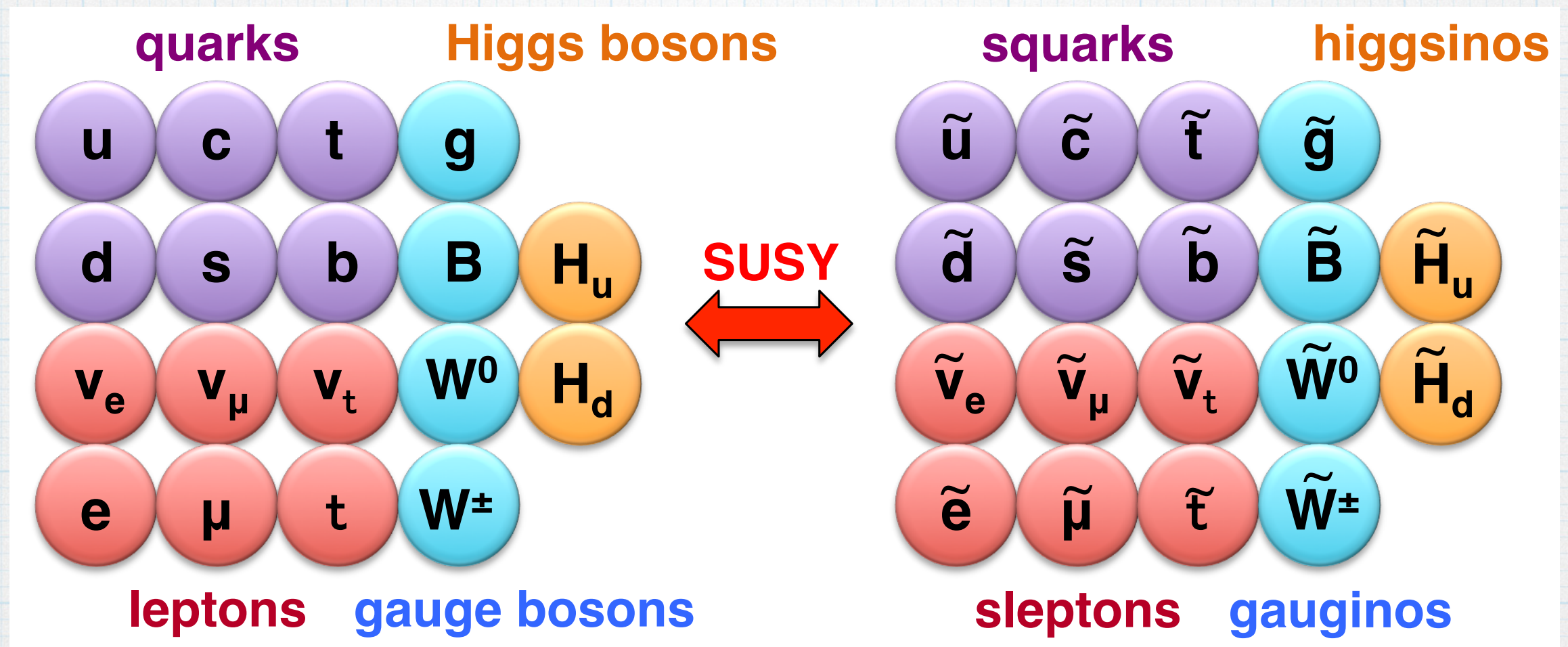
- To address these issues SM has to be extended.

- Many possible extensions are proposed, like SUSY, Composite Higgs, Extra Dimensions ...



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Supersymmetry



- Symmetry between fermions and bosons. One supersymmetric multiplet for every standard model particle.
- In conserved SUSY particles in a multiplet share the same couplings and masses.
- To realize in nature SUSY must be broken.

Advantages of SUSY

- * protects the Higgs mass by introducing a new physics scale.
- * calculable and thus in principle, predictable.
- * provides a viable Dark Matter candidate if R-parity is conserved.
- * lightest Higgs boson can be SM-like in regions of parameter space.
- * unifies the SM gauge couplings.