

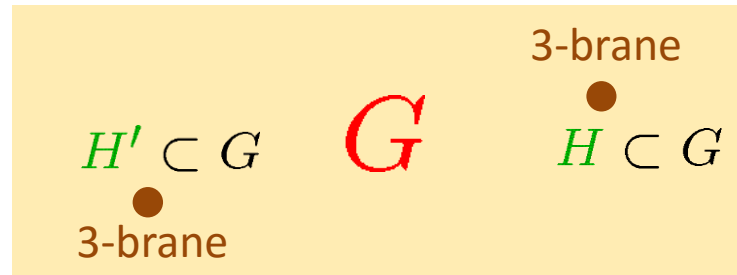


Spectrum in the presence of a brane-localized mass in six dimensions

Yutaka Sakamura
(KEK, SOKENDAI)

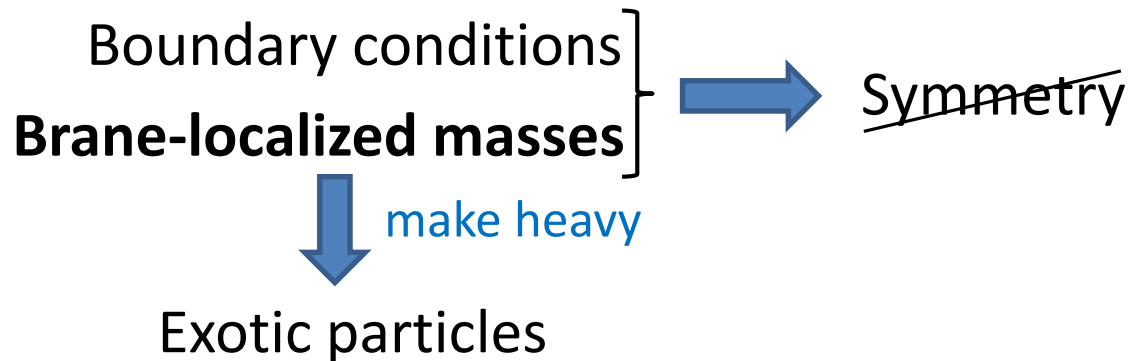
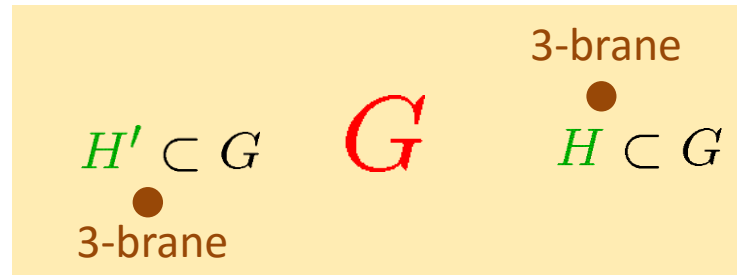
Introduction

Symmetry in the extra-dimensions



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Introduction

We study effects of the brane-mass on spectrum.

Questions

- Can the brane-mass make m_0 heavy enough?
- How much does it deform the mode functions?
- How much does m_0 depend on Λ_{cut} ?

lightest mass eigenvalue

cut-off scale

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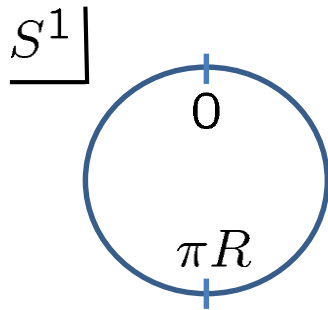
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5D theory

$$\mathcal{L} = -\partial^\mu \phi^* \partial_\mu \phi - \partial_y \phi^* \partial_y \phi$$

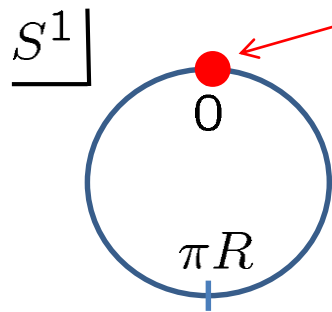


$$m_n^{(0)} = \frac{|n|}{R}$$

5D theory

$$\mathcal{L} = -\partial^\mu \phi^* \partial_\mu \phi - \partial_y \phi^* \partial_y \phi - \mu^2 |\phi|^2 \delta(y)$$

dimension 1
↓



$$m_n^{(0)} = \frac{|n|}{R}$$

$$m_n = \frac{\mu^2}{2} \cot(m_n \pi R)$$

$$\left(\lim_{\mu \rightarrow \infty} m_n = \frac{|n + \frac{1}{2}|}{R} \right)$$

5D theory

$$m_0 < m_1 \leq m_2 \leq \cdots \leq m_N \leq \Lambda_{\text{cut}}$$

cut-off scale
||
 Λ_{cut}

Mass matrix

$$M_{ab}^2 = m_a^{(0)2} \delta_{ab}$$

5D theory

$$m_0 < m_1 \leq m_2 \leq \dots \leq m_N \leq \overset{\text{cut-off scale}}{\parallel} \Lambda_{\text{cut}}$$

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$$M_{ab}^2 = m_a^{(0)2} \delta_{ab} + \frac{\mu^2}{2\pi R}$$

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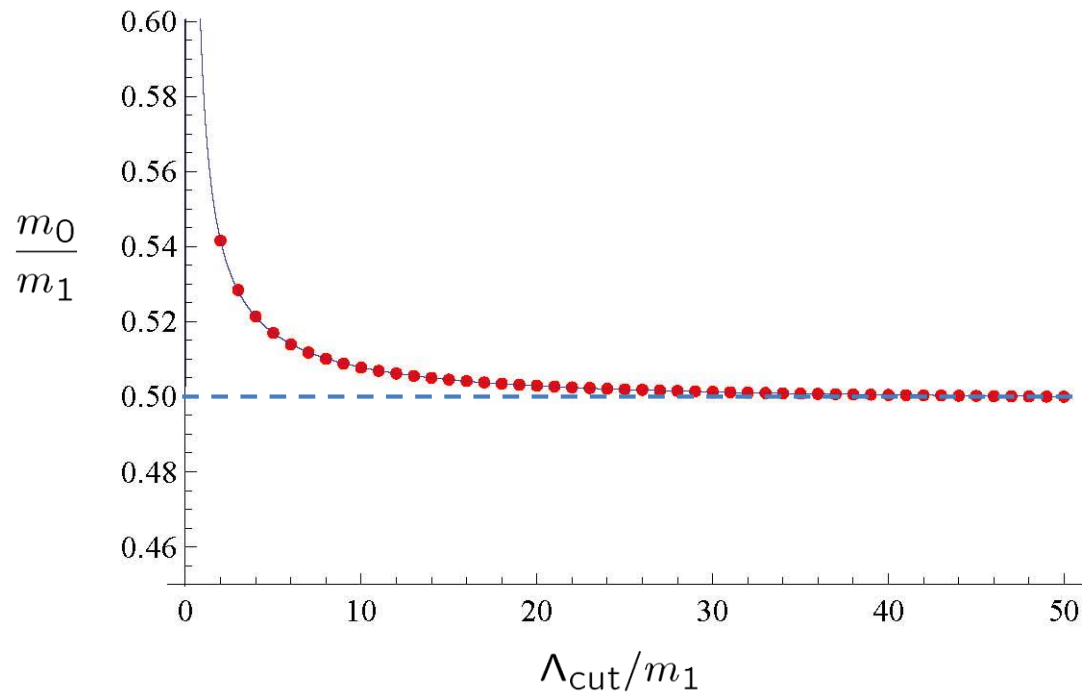


diagonalize

$$\begin{pmatrix} m_0(\mu) & & & \\ & m_1(\mu) & & \\ & & \dots & \\ & & & m_N(\mu) \end{pmatrix}$$

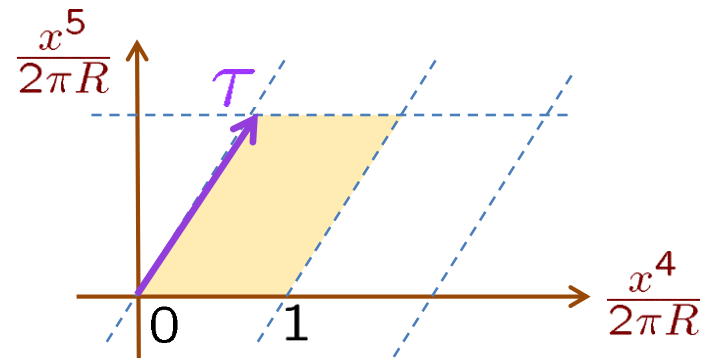
5D theory

Cut-off dependence of m_0



6D theory

T^2



dimensionless

$$\mathcal{L} = -\partial^\mu \phi^* \partial_\mu \phi - \sum_{m=4,5} \partial^m \phi^* \partial_m \phi - \mu^2 |\phi|^2 \delta^{(2)}(x^m)$$

6D theory

Mass matrix $m_{a=(n,l)}^{(0)} = \frac{|n + l\tau|}{R \text{Im } \tau} \quad (n, l \in \mathbb{Z})$

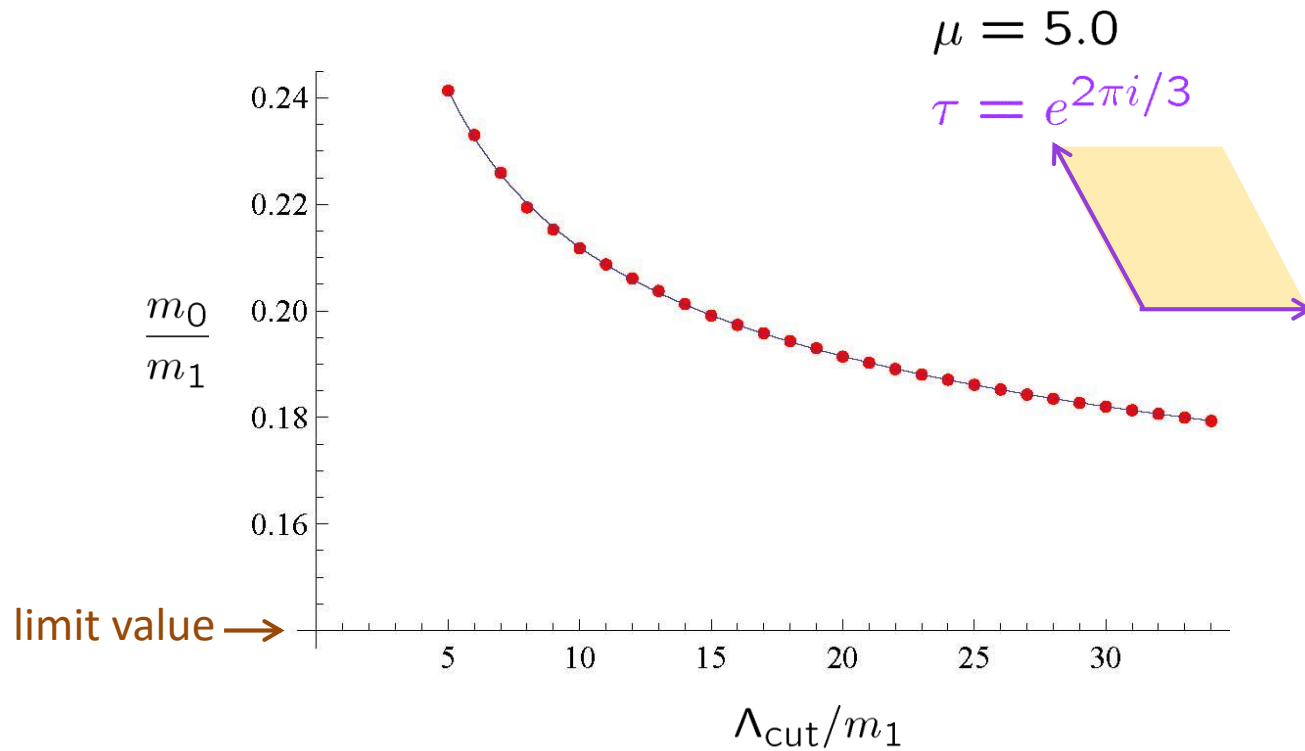
$$M_{ab}^2 = m_a^{(0)2} \delta_{ab} + \frac{\mu^2}{4(\pi R)^2 \text{Im } \tau}$$



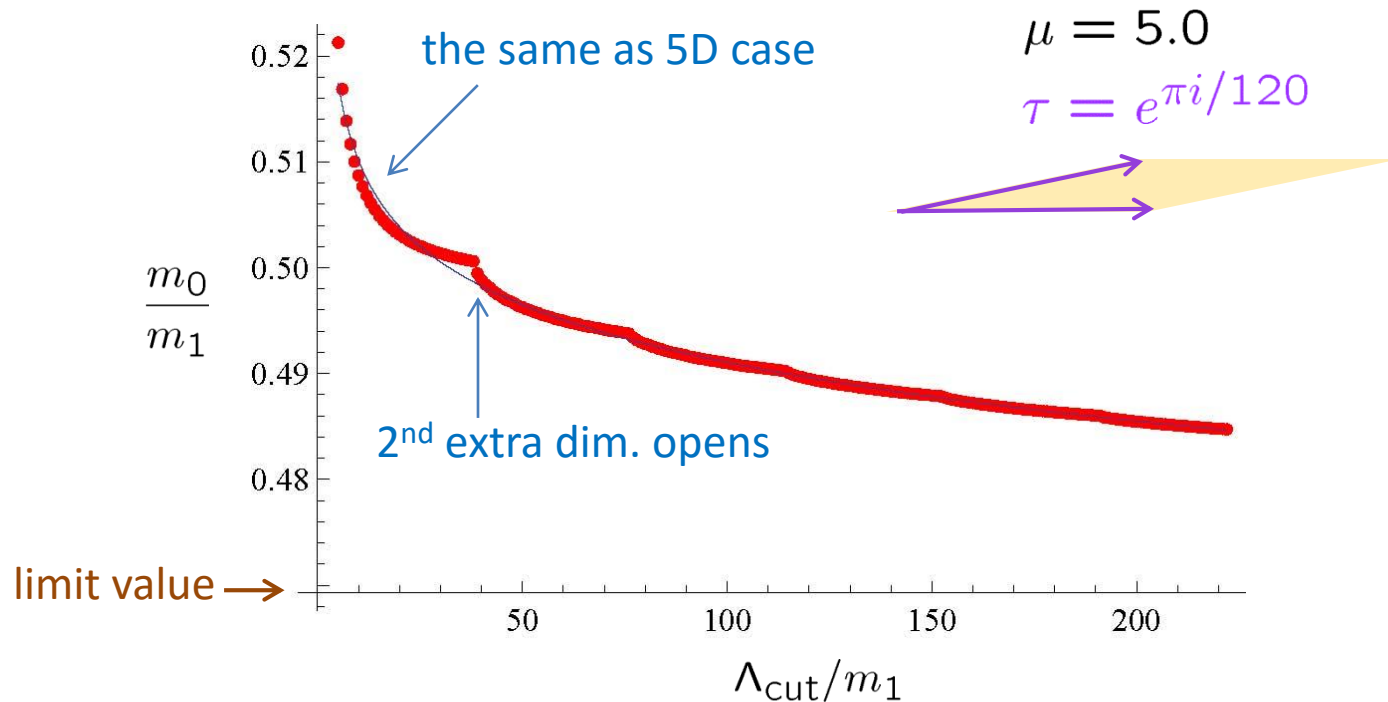
diagonalize

$$\begin{pmatrix} m_0(\mu) & & & \\ & m_1(\mu) & & \\ & & \dots & \\ & & & m_N(\mu) \end{pmatrix}$$

Cut-off dependence of m_0

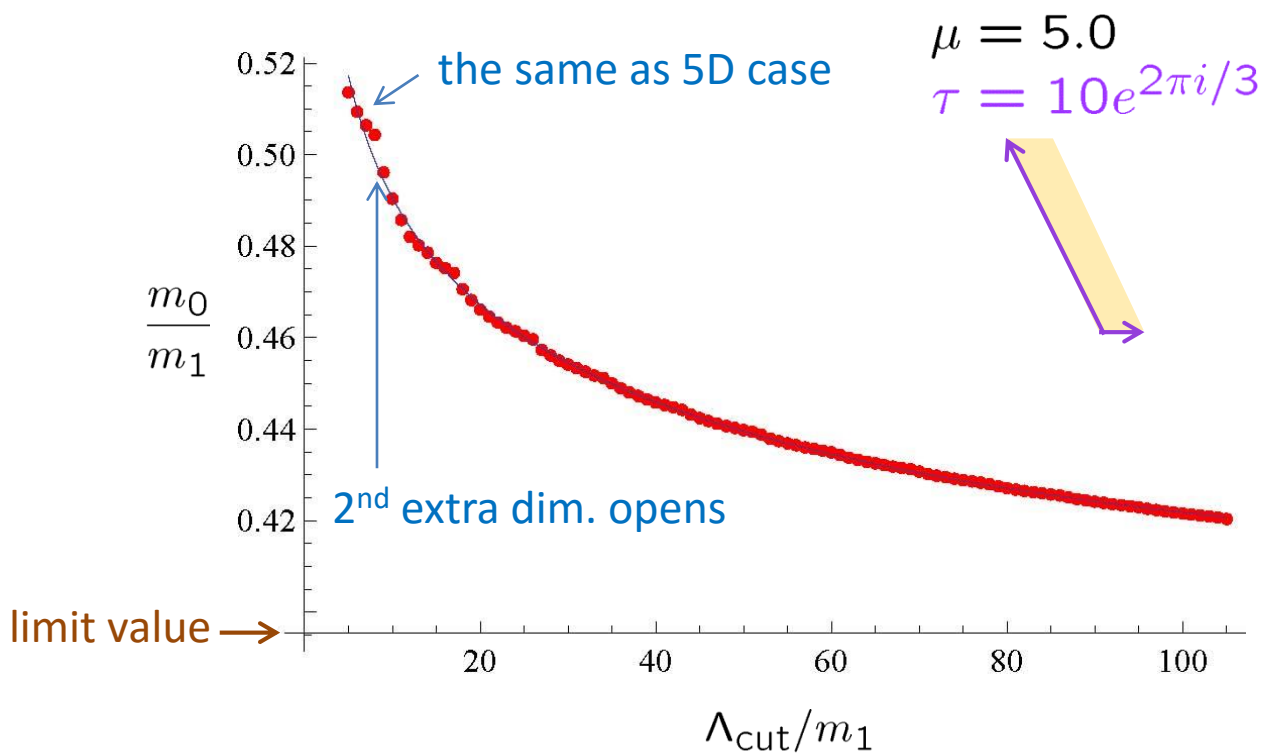


Cut-off dependence of m_0

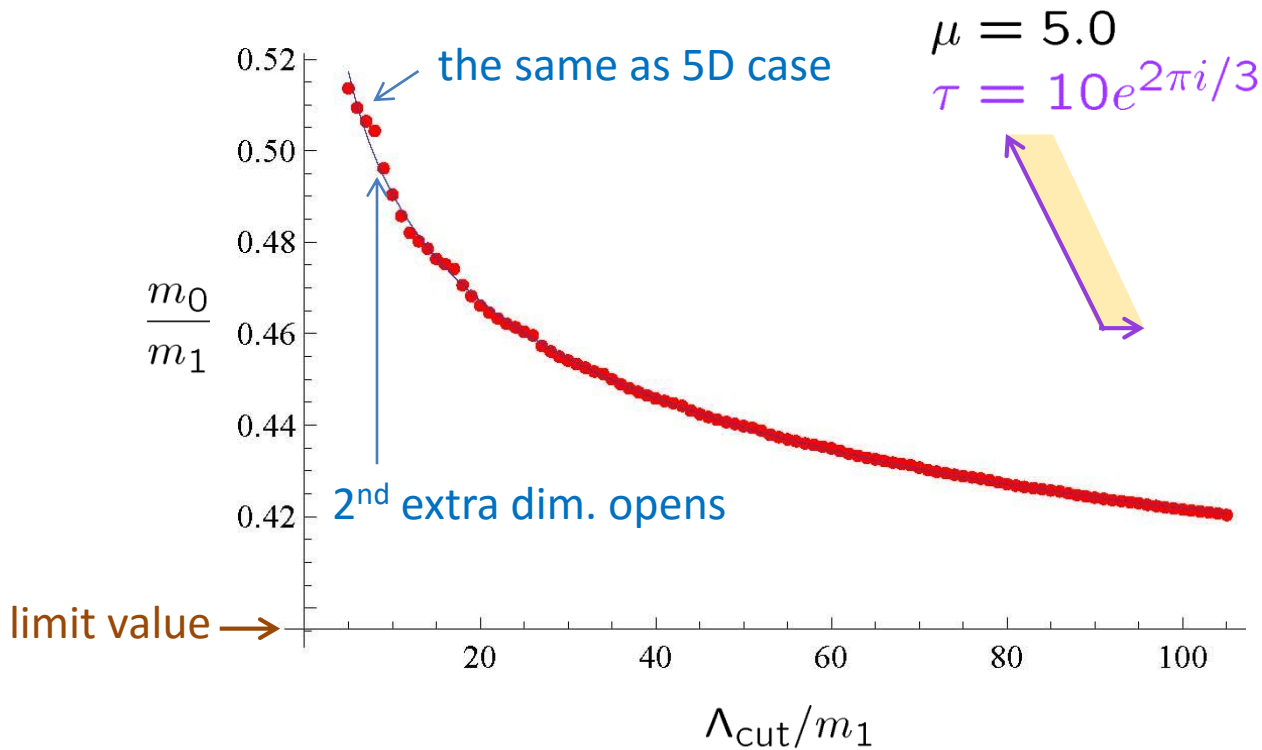


$$\text{At } \Lambda_{\text{cut}} = 15m_1, \quad \frac{m_0}{m_1} \simeq \lim_{\Lambda_{\text{cut}} \rightarrow \infty} \frac{m_0}{m_1} \times \begin{cases} 1.40 & (\tau = e^{2\pi i/3}) \\ 1.07 & (\tau = e^{\pi i/120}) \end{cases}$$

Cut-off dependence of m_0



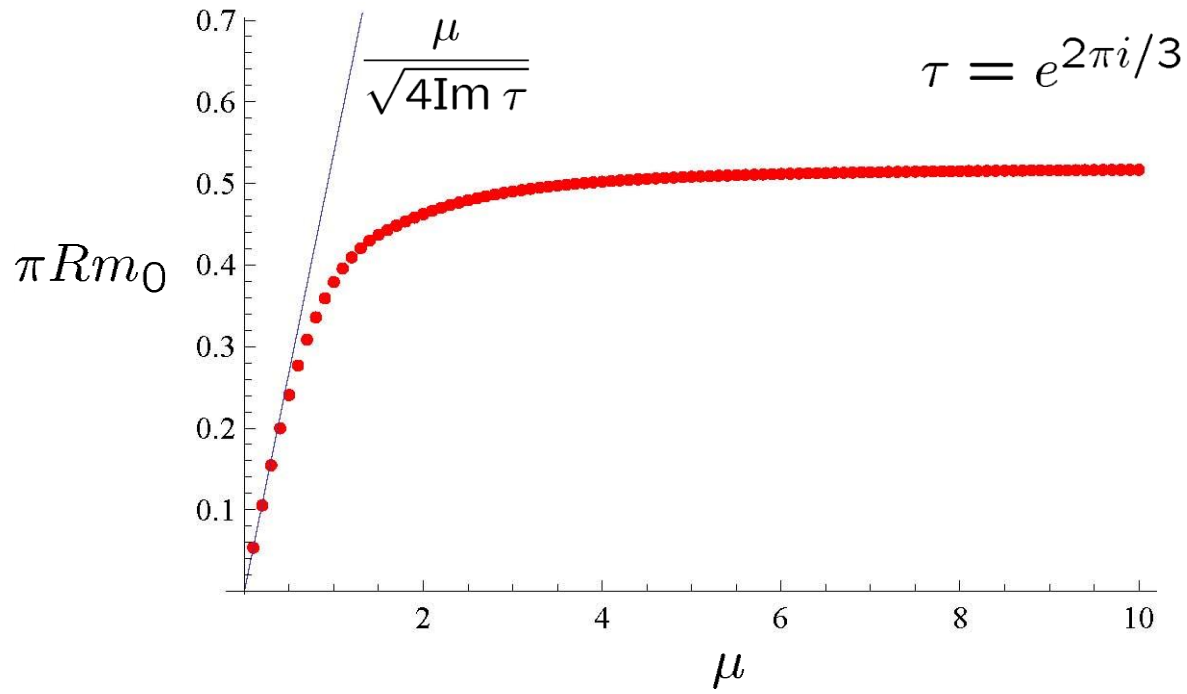
Cut-off dependence of m_0



In the following, we take a limit $\Lambda_{\text{cut}} \rightarrow \infty$.

μ -dependence of m_0

$$M_{ab}^2 = m_a^{(0)2} + \frac{\mu^2}{4(\pi R)^2 \text{Im } \tau}$$



Approximate expression of m_0

- In the limit of $\left\{ \begin{array}{l} \arg \tau \rightarrow 0, \pi \text{ (squashed torus)} \\ \text{or} \\ |\tau| \rightarrow 0, \infty \text{ (long, thin torus)} \end{array} \right.$,

the spacetime approaches to 5D.

Namely,

$$\lim_{\mu \rightarrow \infty} m_0 \rightarrow \lim_{\mu \rightarrow \infty} \frac{m_1}{2} \equiv \frac{\bar{m}_1}{2}$$

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- Since m_a is a function on the torus,

$$m_a \left(\mu; -\frac{1}{\tau} \right) = |\tau| m_a(\mu; \tau)$$

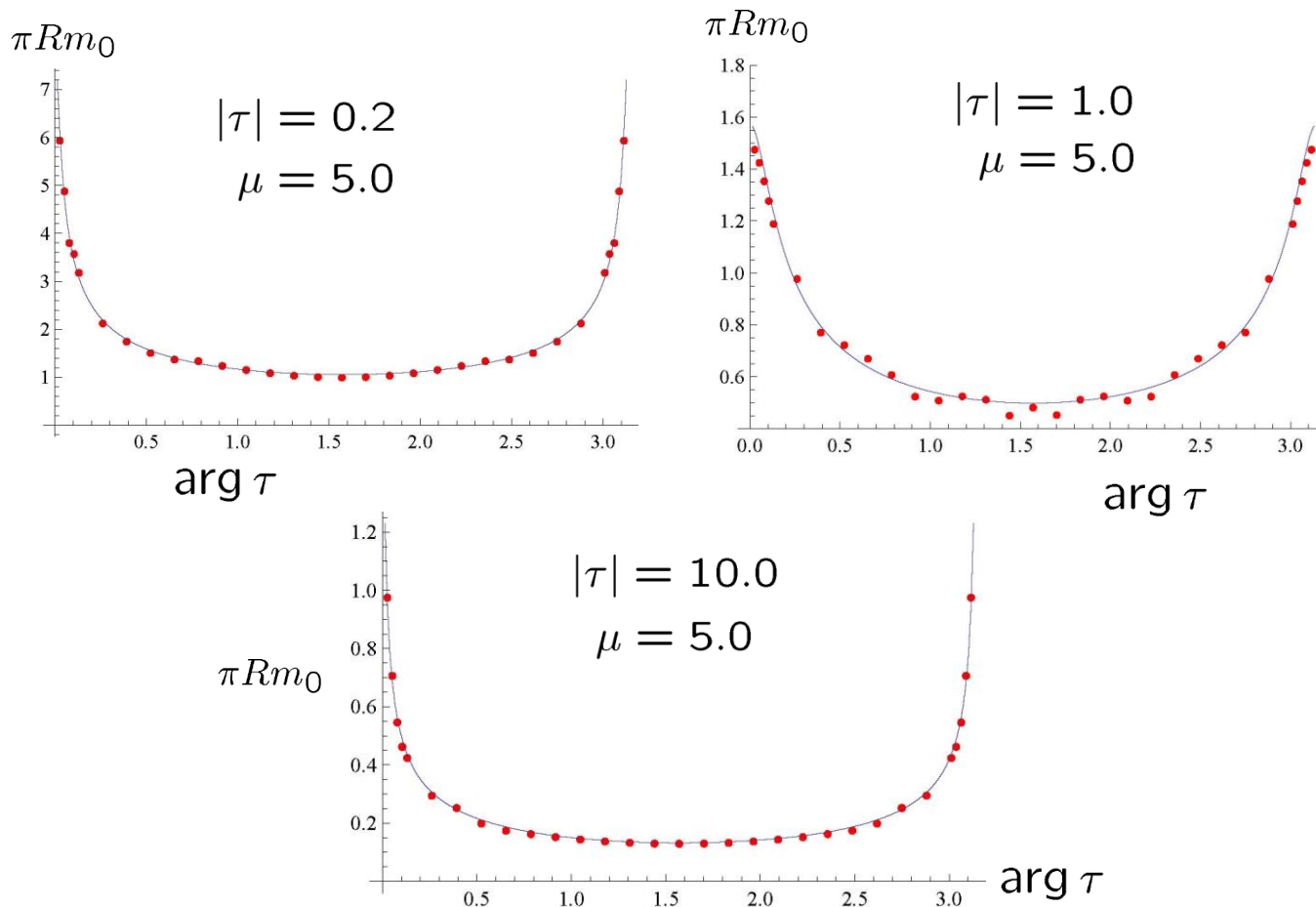
Approximate expression of m_0

$$\left\{ \begin{array}{l} \lim_{\mu \rightarrow \infty} m_0 = \frac{\bar{m}_1}{2} \\ m_a\left(\mu; -\frac{1}{\tau}\right) = |\tau| m_a(\mu; \tau) \end{array} \right.$$

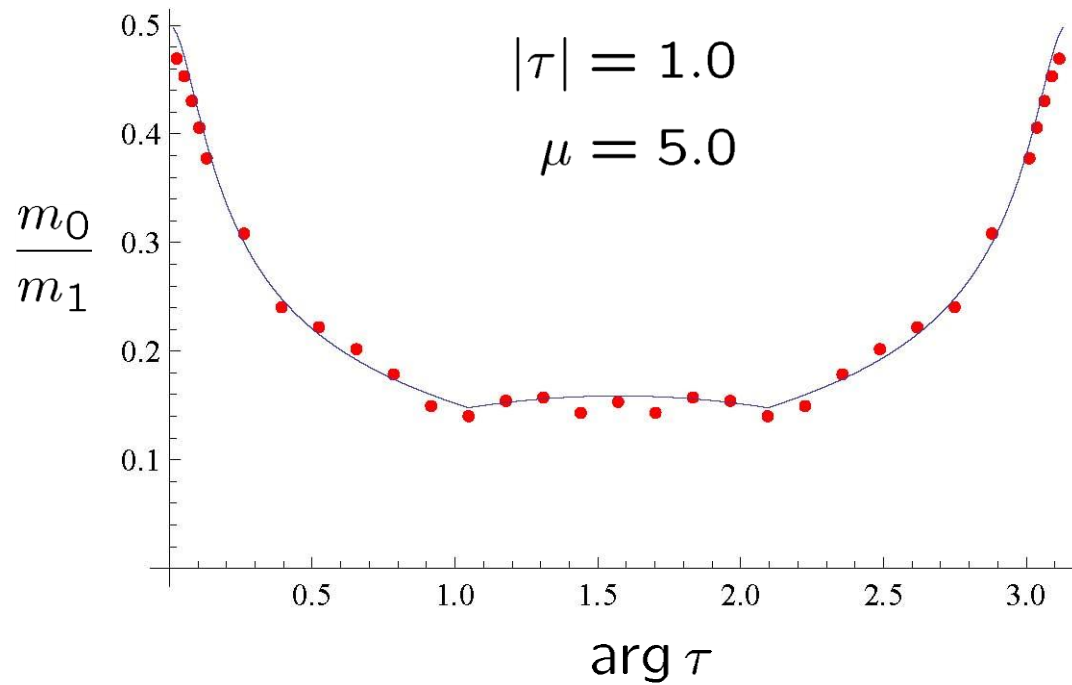
Approximate formula

$$\lim_{\mu \rightarrow \infty} m_0 \simeq \frac{\sqrt{\sin\{\arcsin((\pi R \bar{m}_1)^2 \text{Im } \tau)\}}}{2\pi R \sqrt{\text{Im } \tau}}$$

Approximate expression of m_0



Ratio to KK scale



Summary

- We investigate the spectrum of a 6D scalar theory in the presence of a **brane-localized mass term**.
- The cut-off dependence is non-negligible.
- An approximate expression of m_0 is found.
- m_0/m_{KK} is much smaller than the 5D case.

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- We investigate the spectrum of a 6D scalar theory in the presence of a **brane-localized mass term**.
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Future works

- Calculate the deformed mode functions
- Discuss more general set-ups
(more than one brane-masses, magnetic fluxes, ...)