

Measuring the Higgs Trilinear and Quartic Couplings at the LHC

Chung Kao
University of Oklahoma

[†]Presented at the 24th International Conference on Supersymmetry and
Unification of Fundamental Interactions (SUSY 2016),
Melbourne, Australia, July 4—8, 2016.

Measuring the Higgs Trilinear and Quartic Higgs Couplings at the LHC

Dicus, Kao, and Willenbrock, Phys. Lett. B203 (1988).

Dawson, Kao, and Wang, Phys.Rev. D75 (2007); Phys.Rev. D77 (2008).

Dicus, Kao, and Repko, Phys. Rev. D92 (2015); Phys. Rev. D93 (2016).

- Introduction
- Higgs Pair Production: Gluon Fusion and Bottom Fusion
- The Trilinear Higgs Coupling(s)
- The Discovery Potential of Higgs Pairs at the LHC
- Associated Production of Higgs Pairs and Triple Higgs
- Conclusions

Introduction

- Thus far the results from the LHC indicate that the couplings of the Higgs boson to other particles are consistent with the Standard Model.
- But the ultimate test as to whether this particle is the SM Higgs boson will be the trilinear Higgs coupling that appears in Higgs pair production.
- There are uncertainties in the factorization and renormalization scales as well as variations in the parton distribution functions.

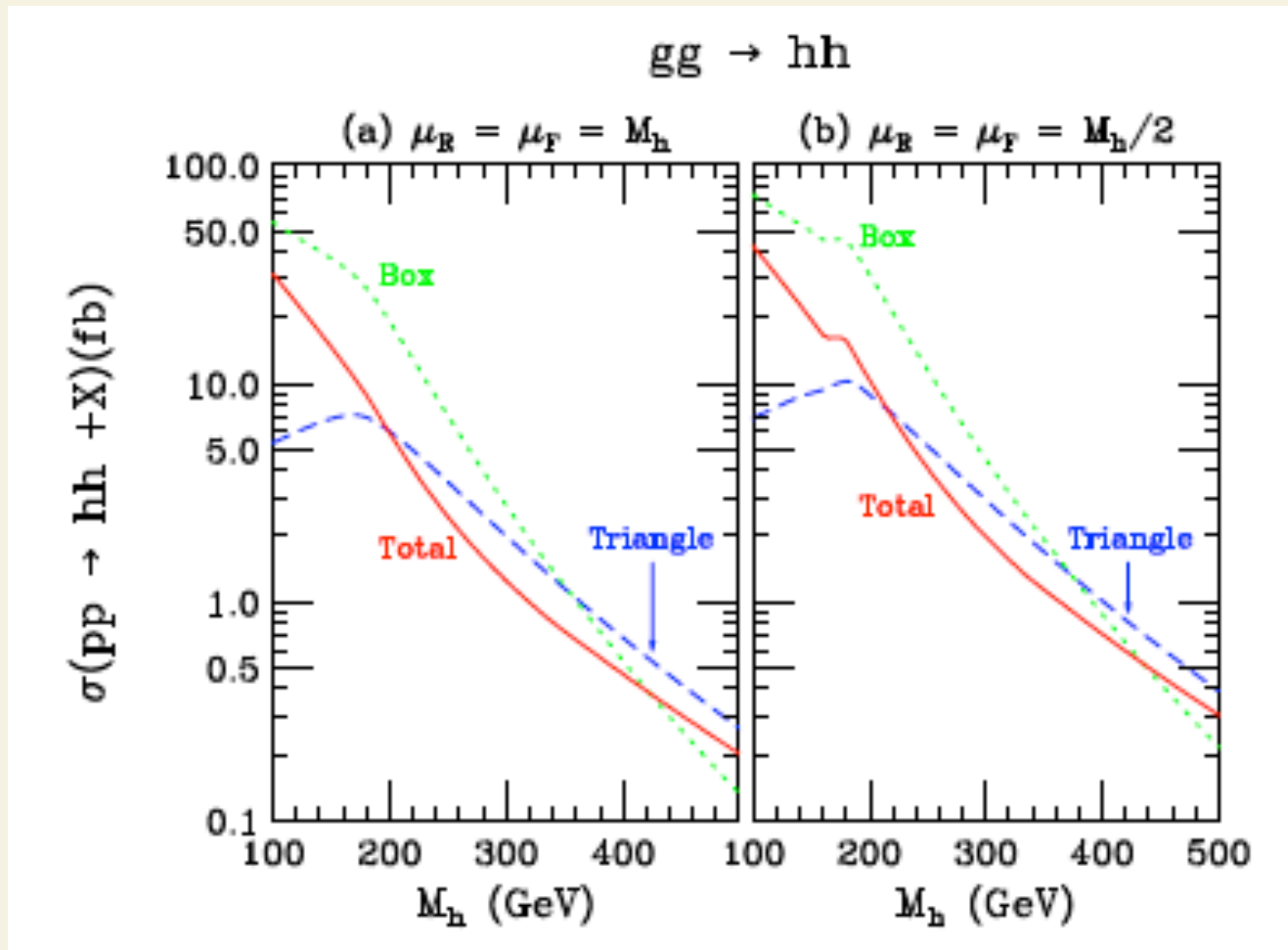
Higgs Pairs Production from Gluon Fusion

Dicus, Kao, and Willenbrock, Phys. Lett. **B203** (1988) 457;

Glover and van der Bij, Nucl. Phys. **B309** (1988) 282.

- For a light Higgs boson with $M_H < 500$ GeV, the dominant source of Higgs boson pair production is gluon fusion through both triangle and box diagrams.
- The triangle diagram involves the Higgs self-coupling while the box diagrams don't.
- For a heavy Higgs boson with $M_H \sim 1$ TeV, vector boson fusion can become significant.

Higgs Pairs Production from Gluon Fusion



Higgs Pair Production from Gluon Fusion

- Li and Voloshin, PRD 89, 013012 (2014); Dawson, Ismail and Low, PRD 91, no. 11, 115008 (2015).
- NLO: Plehn, Spira, and Zerwas, NPB 479 (1996) 46; NPB531 (1998) 655(E); Dawson, Dittmaier, and Spira, PRD 58 (1998) 115012; Jin, Li, Li, Liu and Oakes, PRD 71 (2005) 095004; Baglio, Djouadi, Gröber, Mühlleitner, Quevillon, and Spira, JHEP 1304 (2013) 151; Grigo, Hoff, Melnikov and Steinhauser, NPB 875 (2013) 1.
- NNLO: de Florian and Mazzitelli, Phys. Rev. Lett. 111 (2013) 201801.
- MSSM: Belyaev, Drees, Eboli, Mizukoshi and Novaes, PRD 60 (1999) 075008; Bendezu and Kniehl, PRD 64 (2001) 035006.
- A General Potential and a Non-perturbative Higgs Model: Haba, Kaneta, Mimura and Tsedenbalji, PRD 89 (2014) no.1, 015018.

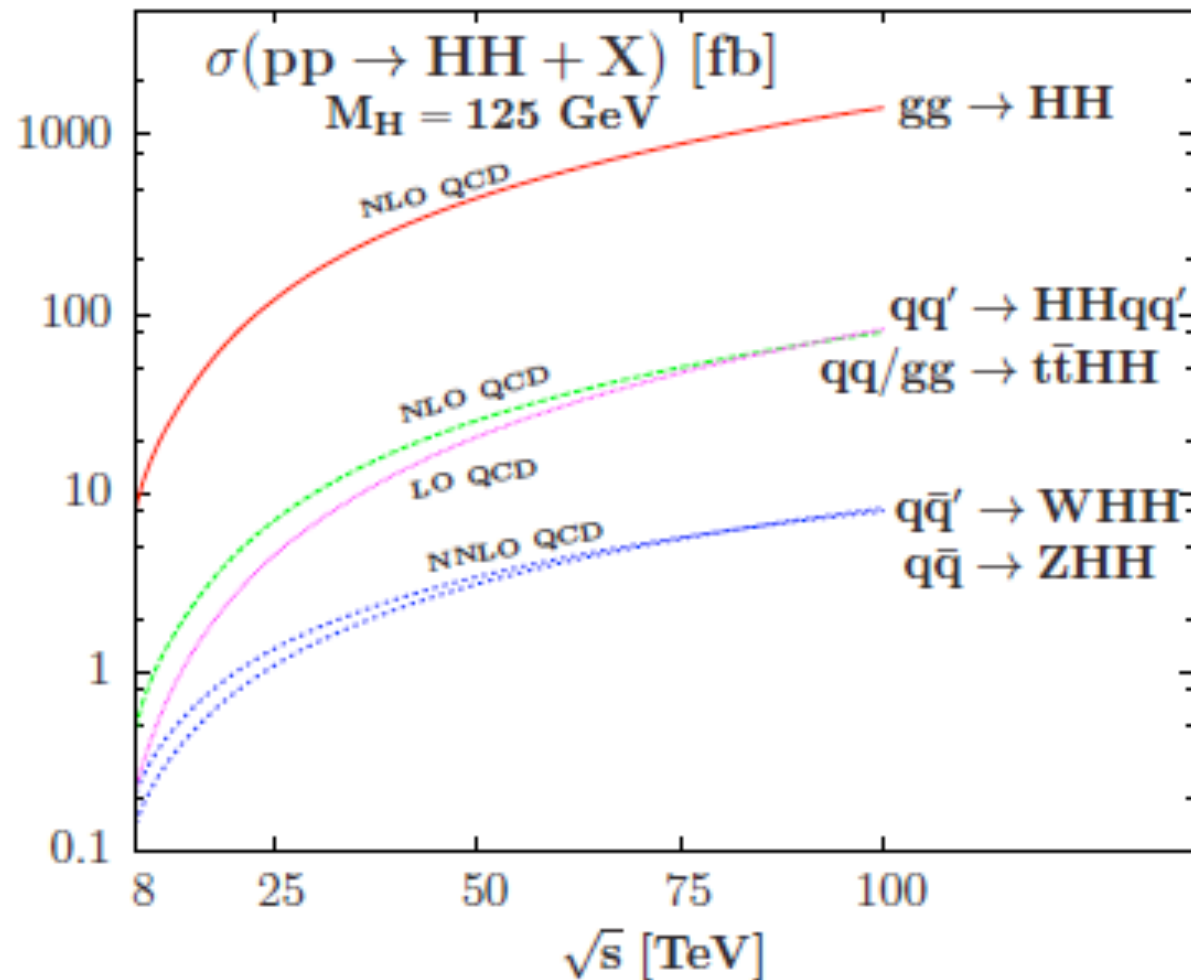
Discovery Channels of Higgs Pairs

- $bb\tau\tau$, $bb\gamma\gamma$, $bbWW$, $bbbb$:

U. Baur, T. Plehn, and D. L. Rainwater, PRD 67 (2003) 033003; PRD 69 (2004) 053004; Dib, Rosenfeld, Zerwekh, JHEP 0605 (2006) 074; Grober and Muhlleitner, JHEP 1106 (2011) 020; Dolan, Englert, Spannowsky JHEP 1210 (2012) 112; Barr, Dolan, Englert, Spannowsky PLB 728 (2014) 308; Papaefstathiou, Yang and Zurita, PRD 87 (2013) 011301; Goertz, Papaefstathiou, Yang and Zurita, JHEP 1306 (2013) 016.

Higgs Pair Production in Hadron Collisions

Baglio, Djouadi, Gröber, Mühlleitner, Quevillon, Spira, JHEP **1304** (2013) 151.



NNLO Higgs Pair Production at Hadron Colliders

de Florian and Mazzitelli, Phys. Rev. Lett. **111** (2013) 201801.

$$\begin{aligned}\sigma_{\text{LO}} &= 17.8_{-3.8}^{+5.3} \text{ fb}, & \sigma_{\text{NLO}} &= 33.2_{-4.9}^{+5.9} \text{ fb}, \\ \sigma_{\text{NNLO}} &= 40.2_{-3.5}^{+3.2} \text{ fb},\end{aligned}\tag{18}$$

$E_{\text{c.m.}}$	8 TeV	14 TeV	33 TeV	100 TeV
σ_{NNLO}	9.76 fb	40.2 fb	243 fb	1638 fb
Scale [%]	+9.0 – 9.8	+8.0 – 8.7	+7.0 – 7.4	+5.9 – 5.8
PDF [%]	+6.0 – 6.1	+4.0 – 4.0	+2.5 – 2.6	+2.3 – 2.6
PDF + α_S [%]	+9.3 – 8.8	+7.2 – 7.1	+6.0 – 6.0	+5.8 – 6.0

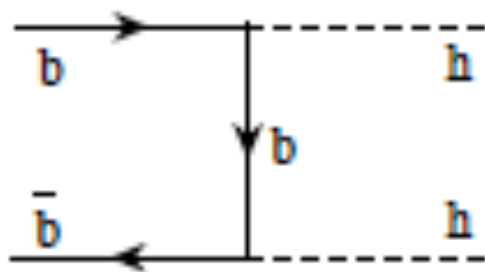
Higgs Pair Production via Bottom Quark Fusion

- In the Standard Model, bottom quark fusion is almost negligible for Higgs pair production.
- In two Higgs doublet models with Type II Yukawa interactions, the Hbb coupling is enhanced by a large value of $\tan\beta$. Thus for $\tan\beta > 7$, bottom quark fusion makes dominant contribution.
- The physical process is gg to $bbHH$.
- However, it is a good approximation to calculate bb to HH if no associate b quarks are tagged.

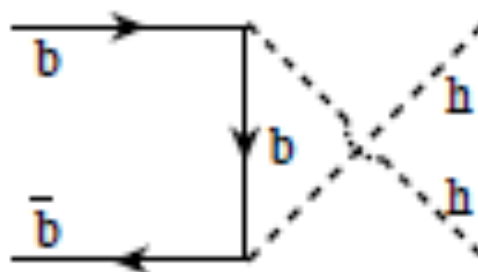
Higgs Pair Production via Bottom Quark Fusion

Dawson, Kao, Wang and Williams, Phys. Rev. D **75** (2007) 013007;

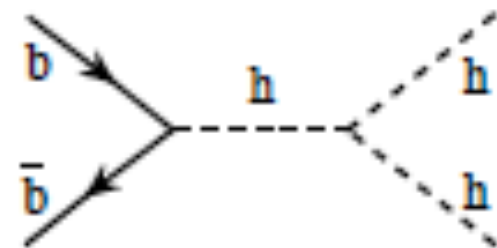
Dawson, Kao and Wang, Phys. Rev. D **77** (2008) 113005.



(1)



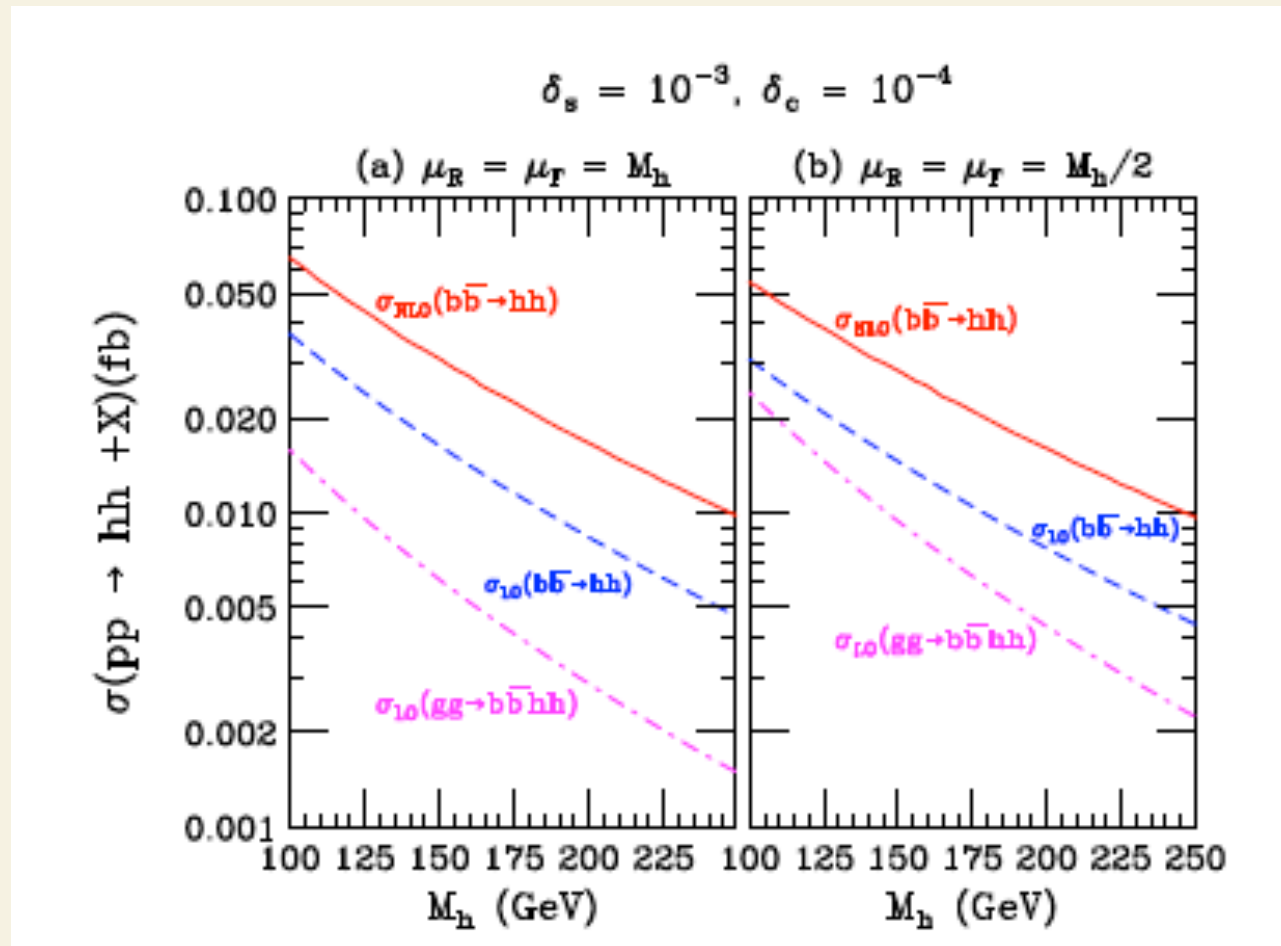
(2)



(3)

Higgs Pair Production via Bottom Quark Fusion

Dawson, Kao, Wang and Williams, Phys. Rev. **D75** (2007) 013007.



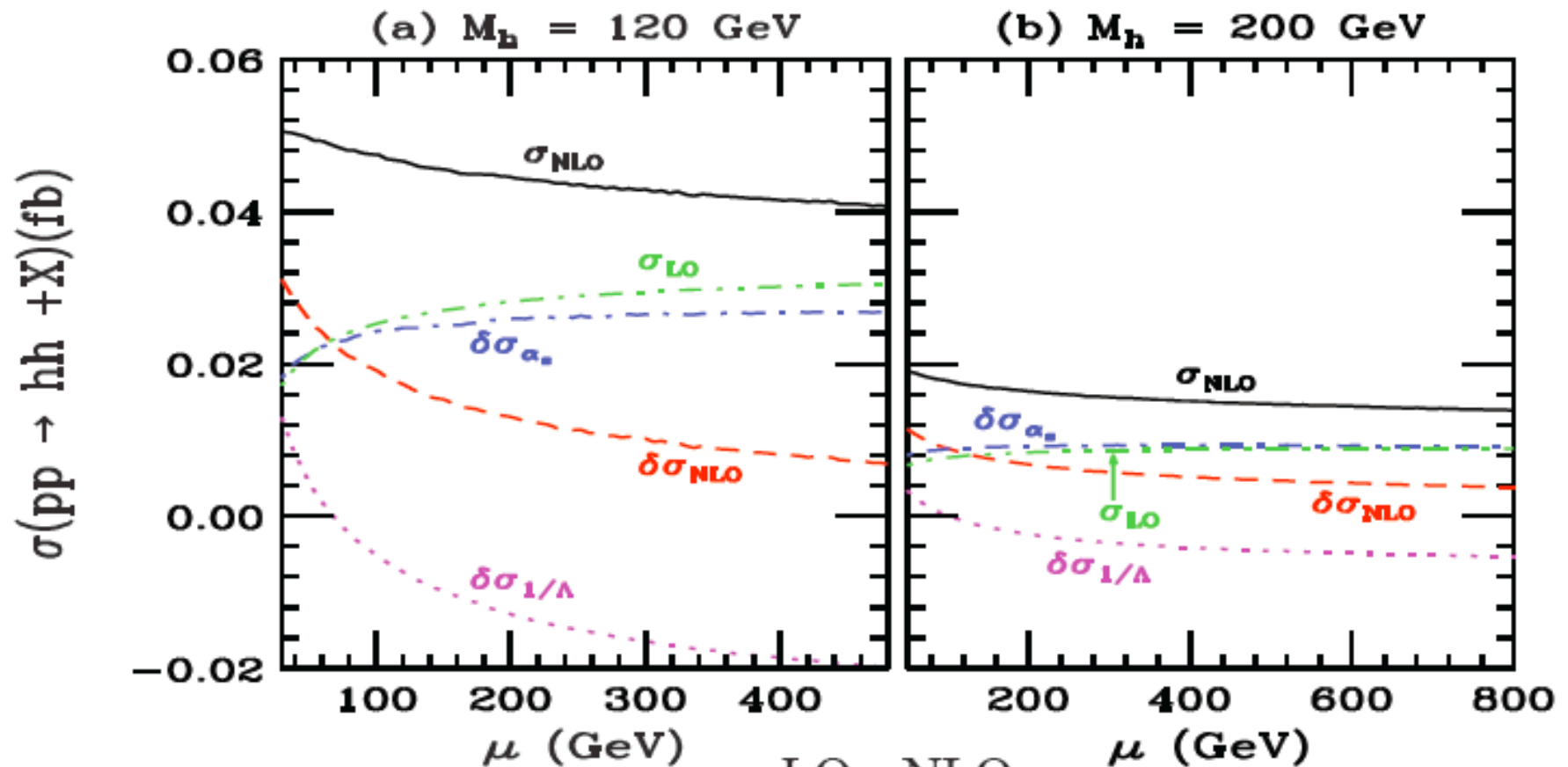
QCD CORRECTIONS TO $bb \rightarrow hh$

Dawson, Kao, Wang and Williams, Phys. Rev. **D75** (2007) 013007.

- **Next-to-Leading Order Corrections**
 - ▶ α_s **Corrections: Real Emission, $bb \rightarrow hhg$**
 - ▶ α_s **Corrections: Virtual Correction**
 - ▶ $1/\Lambda$ **Corrections: $bg \rightarrow bhh$ [$\Lambda = \ln(m_h/m_b)$]**
 - ▶ **$gg \rightarrow bbhh$ Cross Section ($1/\Lambda^2$)**

NLO Corrections to $bb \rightarrow hh$

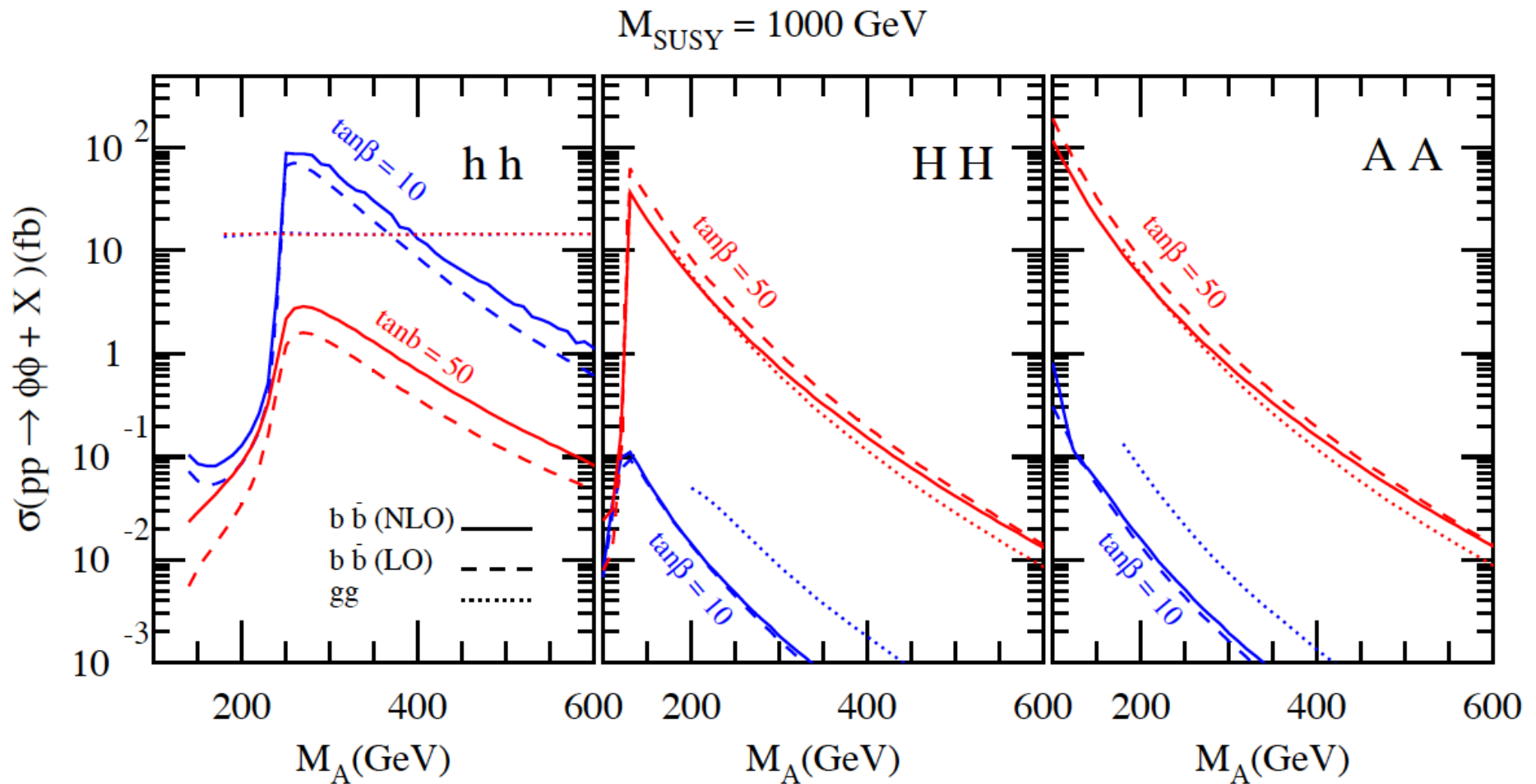
$$\delta_s = 10^{-3}, \delta_c = 10^{-4}$$



$$K = \frac{\text{LO} + \text{NLO}}{\text{LO}} \sim 1.5$$

NLO Corrections to $bb \rightarrow \phi\phi$ in MSSM

Dawson, Kao, and Wang, Phys.Rev. D77 (2008).



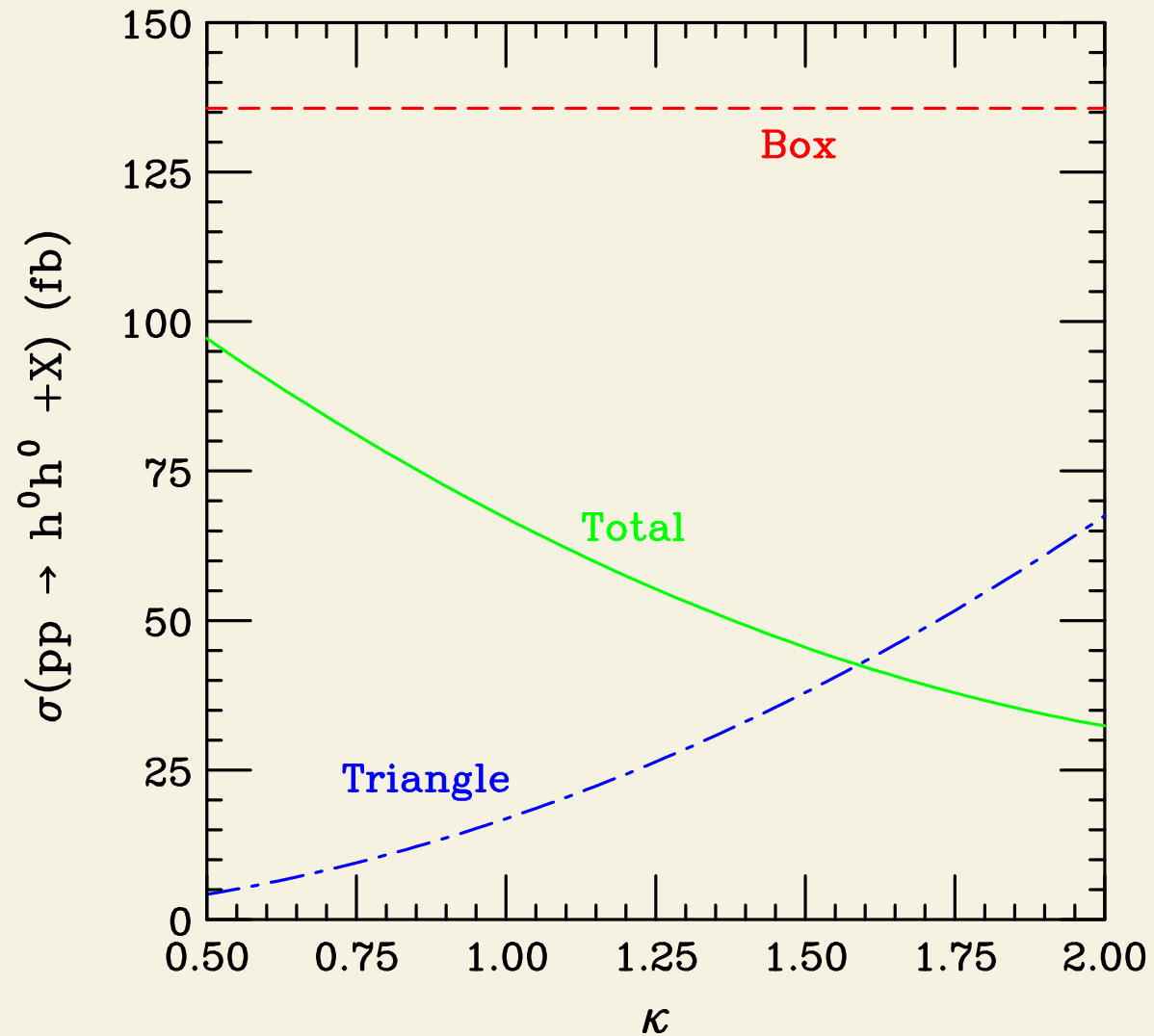
The Trilinear Higgs Coupling(s)

- Higgs pair production from gluon fusion involves $t\bar{t}H$ and HHH couplings.
- The box and triangle diagrams are separately gauge invariant so we can vary the two couplings independently by introducing parameters κ_t and κ or κ_H ,

$$\begin{aligned} t\bar{t}H &: -\frac{m_t}{v} k_t \\ HHH &: -\frac{3M_H^2}{v} \kappa \end{aligned}$$

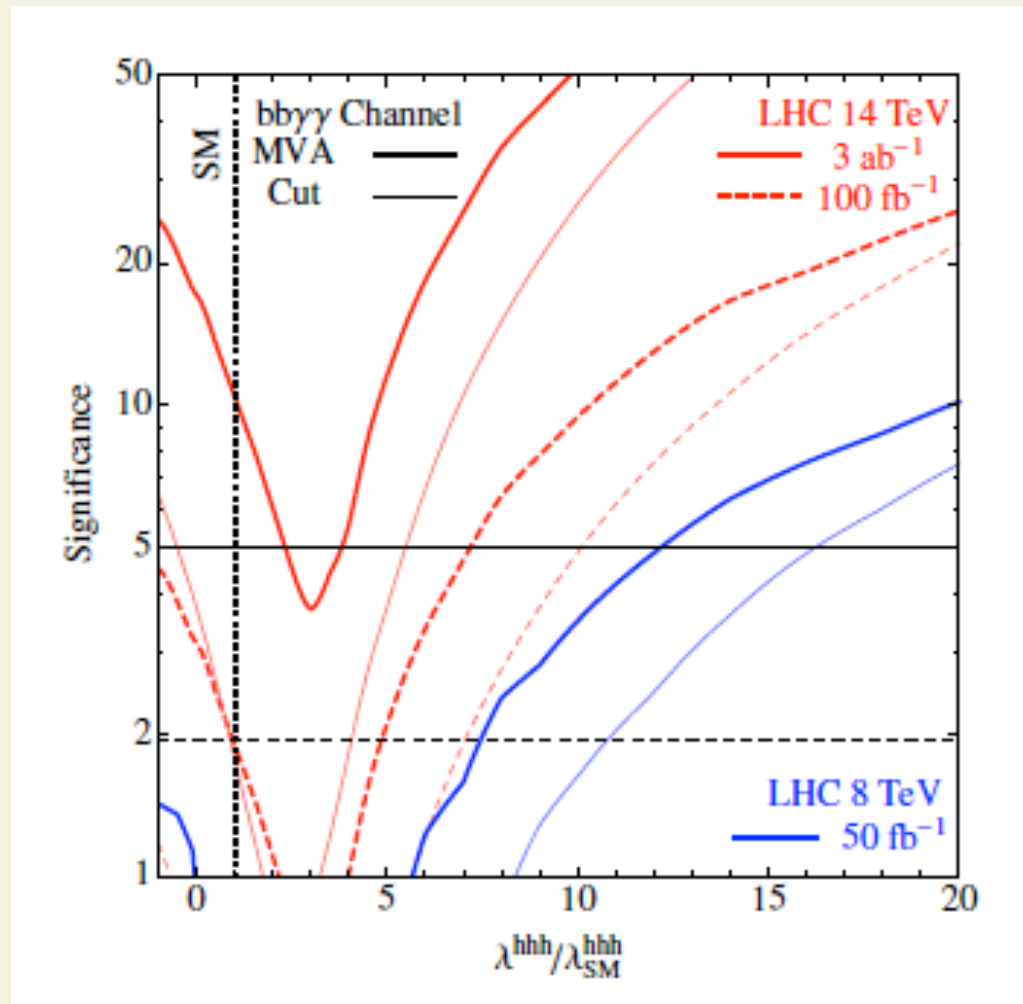
Effects of kappa with $\kappa > 0$

$$\sqrt{s} = 14 \text{ TeV}, K = 1.9$$



The Discovery Potential of Higgs Pairs

Barger, Everett, Jackson, Shaughnessy, PLB 728 (2014) 433.



Conclusions

Barger, Everett, Jackson, Shaughnessy, Phys. Lett. B728 (2014) 433;
Goertz, Papaefstathiou, Yang and Zurita, JHEP 1306 (2013) 016;
de Lima, Papaefstathiou, Spannowsky, JHEP 1408 (2014) 030.

- The $b\bar{b}\gamma\gamma$, $b\bar{b}\tau\tau$, $b\bar{b}WW$ and $b\bar{b}bb$ final states are promising channels to measure the Higgs trilinear coupling(s) at the LHC.
- LHC data at 7–8 TeV should probe large deviations of λ_{hhh} from the SM ($\kappa_h > 7.5$ at 95% C.L.).
- At the LHC with a CM energy of 14 TeV and an integrated luminosity of 3 ab^{-1} , ATLAS and CMS will be able to measure λ_{hhh} within 20%–40% uncertainty.

Interference in Higgs Pair Production

$$V(H) = \lambda v^2 H^2 + \lambda v H^3 + \frac{1}{4} \lambda H^4$$

$$\lambda = \frac{m_H^2}{2v^2}$$

$$\sigma_{TOT} = \sigma_T + \sigma_B = 2 \cos(\alpha_I) \sqrt{\sigma_T \sigma_B}$$

pp to HH +X

\sqrt{s} (TeV)	$\sigma_B(fb)$	$\sigma_T(fb)$	$\sigma_{TOT}(fb)$	$\cos(\alpha_I)$
8	9.06	1.34	4.11	-0.902
12	26.0	3.62	12.2	-0.899
13	31.6	4.36	14.9	-0.898
14	37.8	5.16	17.9	-0.898
33	243.	30.3	120.	-0.893
60	760.	89.8	383.	-0.893
100	1900.	212.	965.	-0.904

ud to ud HH

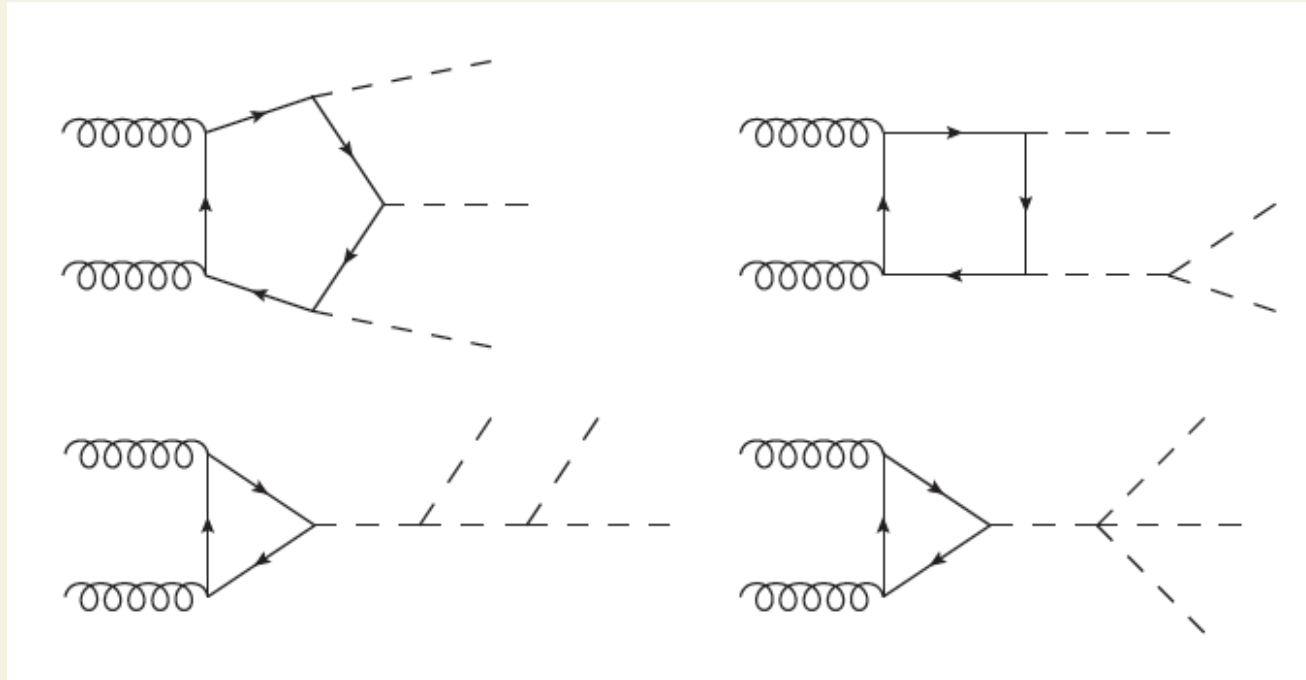
\sqrt{s} (TeV)	σ_B	σ_T	σ_{TOT}	$\cos(\alpha_I)$
8	0.404	0.128	0.141	-0.860
12	0.997	0.295	0.377	-0.844
13	1.18	0.344	0.452	-0.841
14	1.36	0.394	0.530	-0.836
33	6.24	1.60	2.76	-0.804
60	15.6	3.72	7.32	-0.788
100	32.2	7.36	15.8	-0.772

Associated Production with Top Pair

\sqrt{s} (TeV)	σ_B	σ_T	σ_{TOT}	$\cos(\alpha_I)$
8	6.47×10^{-2}	1.96×10^{-3}	8.56×10^{-2}	0.841
12	0.178	5.47×10^{-3}	0.235	0.826
13	0.212	6.54×10^{-3}	0.280	0.825
14	0.249	7.68×10^{-3}	0.328	0.815
33	1.19	3.75×10^{-2}	1.56	0.788
60	2.98	9.51×10^{-2}	3.89	0.765
100	6.15	0.197	8.04	0.769

Production of Triple Higgs Bosons

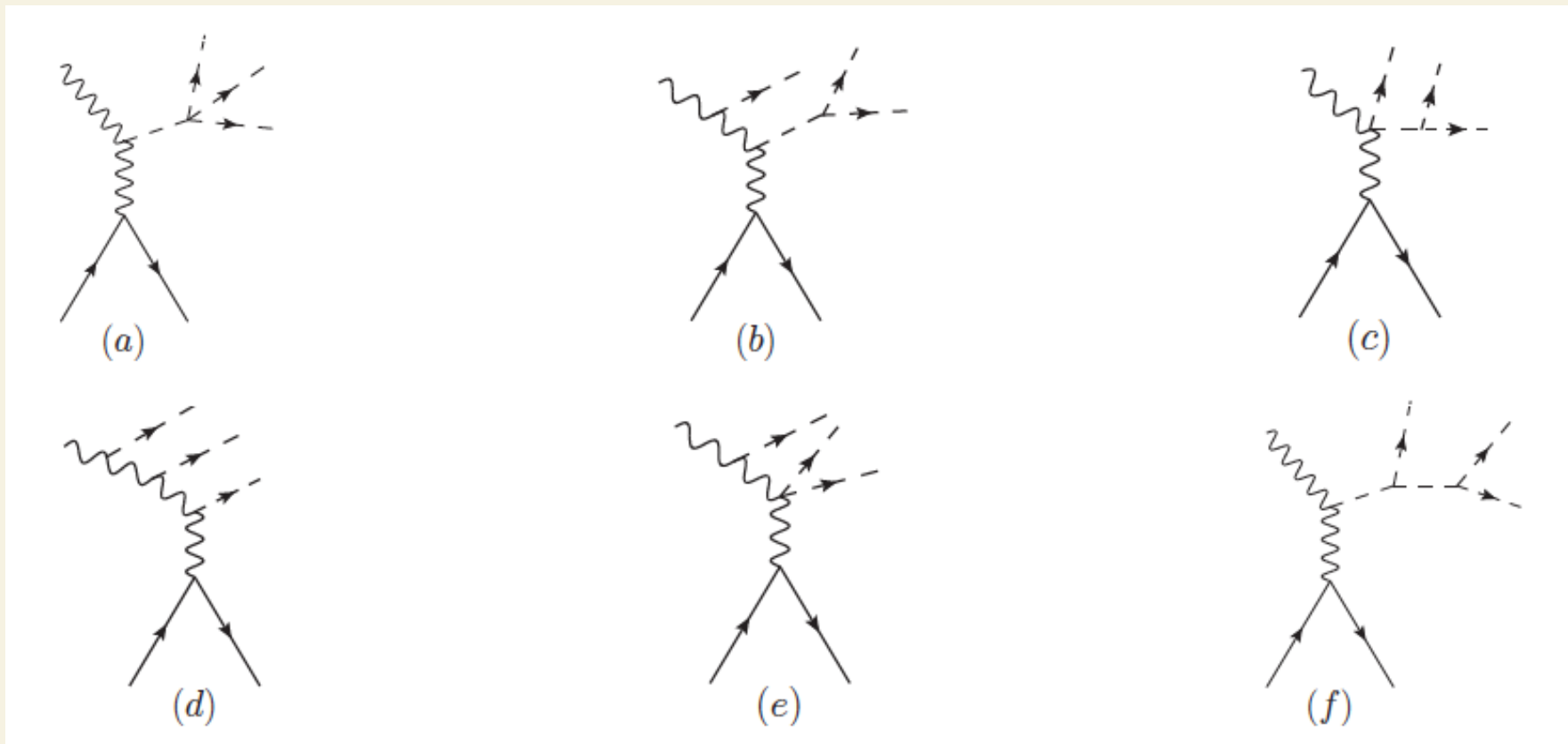
Maltoni, Vryonidou and Zaro, JHEP 1411 (2014) 079.



$\sigma(HHH)$ [fb]	$\sqrt{s} = 14$ TeV	$\sqrt{s} = 33$ TeV	$\sqrt{s} = 100$ TeV
LO FT	0.0557 $^{+34.5+2.5\%}_{-24.0-2.7\%}$	0.438 $^{+26.8+1.5\%}_{+20.0-2.0\%}$	3.78 $^{+24.1+0.9\%}_{-18.7-1.7\%}$
NLO FT _{approx}	0.0894 $^{+16.5+2.5\%}_{-14.6-3.2\%}$	0.677 $^{+14.5+1.4\%}_{-13.4-1.7\%}$	5.09 $^{+13.5+1.0\%}_{-12.7-1.3\%}$

Associated Triple Higgs Boson Production

Dicus, Repko and Kao, Phys.Rev. D93 (2016) no.11, 113003.



Cancellation of High Energy Behavior

- Let us consider the matrix element as

$$\mathcal{M} = f \bar{v}(p_2) \gamma_\mu (g_V + g_A \gamma_5) u(p_1) X^{\mu\nu} \epsilon_\nu(p)$$

- The spinor factor goes as E^1 .
- $X^{\mu\nu}$ has two or three Z propagators, some goes as E^0 .
- The longitudinal Z polarization vector goes as E^1 .
- We have explicitly implemented cancellation of energy dependence by proper grouping of terms.

$\sigma(\text{pp to ZHHH} + X)$ in fb

\sqrt{s}	σ_{44}	σ_{3333}	σ_{433}	σ_{40}	σ_{330}	σ_{43}	σ_0	σ_{333}	σ_{30}	σ_{33}	σ_{TOT}
8	$4.72 \cdot 10^{-7}$	$1.20 \cdot 10^{-6}$	$1.43 \cdot 10^{-6}$	$2.38 \cdot 10^{-6}$	$3.38 \cdot 10^{-6}$	$6.03 \cdot 10^{-6}$	$7.69 \cdot 10^{-6}$	$9.01 \cdot 10^{-6}$	$3.05 \cdot 10^{-5}$	$3.60 \cdot 10^{-5}$	$9.80 \cdot 10^{-5}$
13	$1.57 \cdot 10^{-6}$	$3.61 \cdot 10^{-6}$	$4.47 \cdot 10^{-6}$	$6.94 \cdot 10^{-6}$	$9.31 \cdot 10^{-6}$	$1.80 \cdot 10^{-5}$	$2.32 \cdot 10^{-5}$	$2.55 \cdot 10^{-5}$	$9.05 \cdot 10^{-5}$	$1.09 \cdot 10^{-4}$	$2.92 \cdot 10^{-4}$
14	$1.85 \cdot 10^{-6}$	$4.21 \cdot 10^{-6}$	$5.22 \cdot 10^{-6}$	$8.01 \cdot 10^{-6}$	$1.07 \cdot 10^{-5}$	$2.08 \cdot 10^{-5}$	$2.70 \cdot 10^{-5}$	$2.94 \cdot 10^{-5}$	$1.05 \cdot 10^{-4}$	$1.27 \cdot 10^{-4}$	$3.39 \cdot 10^{-4}$
33	$9.37 \cdot 10^{-6}$	$1.90 \cdot 10^{-5}$	$2.46 \cdot 10^{-5}$	$3.47 \cdot 10^{-5}$	$4.38 \cdot 10^{-5}$	$9.24 \cdot 10^{-5}$	$1.23 \cdot 10^{-4}$	$1.24 \cdot 10^{-4}$	$4.64 \cdot 10^{-4}$	$5.84 \cdot 10^{-4}$	$1.52 \cdot 10^{-3}$
60	$2.43 \cdot 10^{-5}$	$4.67 \cdot 10^{-5}$	$6.16 \cdot 10^{-5}$	$8.41 \cdot 10^{-5}$	$1.04 \cdot 10^{-4}$	$2.26 \cdot 10^{-4}$	$3.06 \cdot 10^{-4}$	$2.97 \cdot 10^{-4}$	$1.14 \cdot 10^{-3}$	$1.45 \cdot 10^{-3}$	$3.74 \cdot 10^{-3}$
100	$5.15 \cdot 10^{-5}$	$9.55 \cdot 10^{-5}$	$1.27 \cdot 10^{-4}$	$1.70 \cdot 10^{-4}$	$2.07 \cdot 10^{-4}$	$4.61 \cdot 10^{-4}$	$6.26 \cdot 10^{-4}$	$5.96 \cdot 10^{-4}$	$2.31 \cdot 10^{-3}$	$2.97 \cdot 10^{-3}$	$7.62 \cdot 10^{-3}$

$\sigma(\text{pp to } W\text{HHH} + X) \text{ in fb}$

\sqrt{s}	σ_{44}	σ_{3333}	σ_{433}	σ_{40}	σ_{330}	σ_{43}	σ_0	σ_{333}	σ_{30}	σ_{33}	σ_{TOT}
8	$6.58 \cdot 10^{-7}$	$1.63 \cdot 10^{-8}$	$1.96 \cdot 10^{-8}$	$2.27 \cdot 10^{-8}$	$3.17 \cdot 10^{-8}$	$7.26 \cdot 10^{-8}$	$6.14 \cdot 10^{-8}$	$1.07 \cdot 10^{-5}$	$2.87 \cdot 10^{-5}$	$4.16 \cdot 10^{-5}$	$1.04 \cdot 10^{-4}$
13	$2.03 \cdot 10^{-8}$	$4.53 \cdot 10^{-8}$	$5.65 \cdot 10^{-8}$	$6.00 \cdot 10^{-8}$	$7.96 \cdot 10^{-8}$	$1.99 \cdot 10^{-5}$	$1.70 \cdot 10^{-5}$	$2.78 \cdot 10^{-5}$	$7.79 \cdot 10^{-5}$	$1.17 \cdot 10^{-4}$	$2.85 \cdot 10^{-4}$
14	$2.36 \cdot 10^{-8}$	$5.19 \cdot 10^{-8}$	$6.52 \cdot 10^{-8}$	$6.58 \cdot 10^{-8}$	$9.02 \cdot 10^{-8}$	$2.28 \cdot 10^{-5}$	$1.96 \cdot 10^{-5}$	$3.16 \cdot 10^{-5}$	$8.93 \cdot 10^{-5}$	$1.34 \cdot 10^{-4}$	$3.27 \cdot 10^{-4}$
33	$1.08 \cdot 10^{-5}$	$2.12 \cdot 10^{-5}$	$2.78 \cdot 10^{-5}$	$2.66 \cdot 10^{-5}$	$3.34 \cdot 10^{-5}$	$9.15 \cdot 10^{-5}$	$8.08 \cdot 10^{-5}$	$1.21 \cdot 10^{-4}$	$3.55 \cdot 10^{-4}$	$5.54 \cdot 10^{-4}$	$1.32 \cdot 10^{-3}$
60	$2.66 \cdot 10^{-5}$	$5.01 \cdot 10^{-5}$	$6.65 \cdot 10^{-5}$	$6.13 \cdot 10^{-5}$	$7.56 \cdot 10^{-5}$	$2.13 \cdot 10^{-4}$	$1.91 \cdot 10^{-4}$	$2.77 \cdot 10^{-4}$	$8.29 \cdot 10^{-4}$	$1.31 \cdot 10^{-3}$	$3.10 \cdot 10^{-3}$
100	$5.43 \cdot 10^{-5}$	$9.95 \cdot 10^{-5}$	$1.33 \cdot 10^{-4}$	$1.20 \cdot 10^{-4}$	$1.64 \cdot 10^{-4}$	$4.21 \cdot 10^{-4}$	$3.81 \cdot 10^{-4}$	$5.40 \cdot 10^{-4}$	$1.63 \cdot 10^{-3}$	$2.62 \cdot 10^{-3}$	$6.14 \cdot 10^{-3}$

Summary for Triple Higgs Production

- In general the problem of determining κ_4 will be very difficult.
- Similar conclusions have been reached by Binoth, Karg, Kauer, and Ruckl (2006) and others for the gluon fusion process.

Bonus

Loop Integrals

- Scalar Integrals: 't Hooft and Veltman (1979)
- Tensor Integrals: Passarino and Veltman (1979)
- FF: G.J. van Oldenborgh (1990)
- LOOP: Dicus and Kao (1990)
- LoopTools: Hahn and Perez-Victoria (1999)
- BlackHat: Bern, Dixon, and Kosower et al. (2013)

Decoupling Limit of General 2HDMs

