TELE script

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Final Focus meeting October 22, 2014

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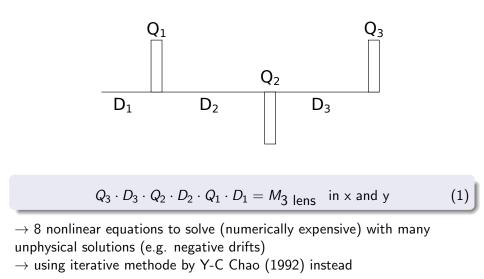
- Code to automatically produce 4 lens telescopes with given magnification
- used in FFADA program
- rewritten from FORTRAN to Python
- based on calculations for 3 thin lenses \rightarrow requires:
 - first drift length (e.g. L^*)
 - last drift length (little to no impact on result when kept short)
 - first quadrupole (roughly defined by L^* and maximum gradient)
- ullet \Rightarrow ideal for FFS design

$M_{tot} = D_4 \cdot Q_4 \cdot D_3 \cdot Q_3 \cdot D_2 \cdot Q_2 \cdot D_1 \cdot Q_1 \cdot D^*$

 L^* is given, first guess for Q_1 is defined by L^* and max. gradient, D_4 chosen freely (result depends only weakly on D_4), M_{tot} defined by magnification

$$\Rightarrow D_4^{-1} M_{tot} D^{*-1} Q_1^{-1} = Q_4 \cdot D_3 \cdot Q_3 \cdot D_2 \cdot Q_2 \cdot D_1 = M_3 \text{ lens}$$

 \rightarrow reduced 4 lens system to 3 lens system



Iterative methode

Step 1:

$$Q_{2} \cdot D_{2} = M^{1}$$

$$\Rightarrow \begin{bmatrix} 1 & L_{2} \\ K_{2} & 1 + K_{2}L_{2} \end{bmatrix} = M^{1}$$

$$\Rightarrow L_{2} = M_{1,2}^{1} \quad K_{2} = M_{2,1}^{1}$$

$$\Rightarrow \begin{bmatrix} 1 & M_{1,2}^{1} \\ M_{2,1}^{1} & 1 + M_{1,2}^{1}M_{2,1}^{1} \end{bmatrix} = M^{1}$$
(3)

Step 2:

$$M^1 = Q_2 \cdot D_2 = D_3^{-1} M^2 Q_1^{-1} \tag{4}$$

Image: Image:

Using Eq. 3 we get expressions for L_3, K_1, L_2, K_2 as functions of matrix M^2

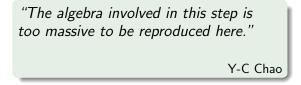
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Step 3:

$$M^2 = D_3 \cdot Q_2 \cdot D_2 \cdot Q_1 = Q_3^{-1} M^3 D_1^{-1}$$

similar to Step 2 just with larger equations



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REFERENCE

 O. Napoly, Thin lens telescopes for final focus systems, CERN/LEP-1H/99-09, CLIC Note 102.

APPENDIX & EXPLICIT FORMS OF X, Y, AND Z IN EQ. (3.18)

By careful manipulation of Eq. (1.2), taking advantage of the simplified forms for the sum and difference of its x-part and y-part, one arrives at two independent and relatively simple equations for L_1 and K_3 as follow:

 $[g^2 - 2(\overline{a} - 1)\overline{a}]\tilde{K}_3\tilde{L}_1 + (\overline{a} - 1)\tilde{B}\tilde{K}_3 + (\overline{a} - 1)\tilde{C}\tilde{L}_1 + g(\overline{a} - 1)^2 = 0,$

 $[(z-1)z^2 - g^2] \vec{K}_3 \vec{L}_1 - (z-1)z \vec{B} \vec{K}_3 \qquad (A.1)$

 $-(u - 1)u\hat{C}\hat{L}_{1} + i(u - 1)\hat{C}\hat{B} - a^{2}(u - 1)i = 0$

where

$\widetilde{K_2} = gK_2 - \overline{\epsilon}$,	$\widetilde{L_1} = gL_1 - k$
$\hat{B} = (g\bar{\delta} - \bar{a}\bar{g})$,	$\hat{C} = (g_2 - 4\hat{c})$,
$= \frac{M_{e11} + M_{g11}}{2}$	$a = \frac{M_{p11} - M_{p11}}{1 + M_{p11}}$
$= \frac{M_{x12} + M_{y12}}{2}$	$k = \frac{M_{F12} - M_{F12}}{2}$,
$=\frac{M_{e31}+M_{e31}}{2}$	$c = \frac{M_{e31} - M_{g31}}{2}$.

Elimination of say, \mathcal{K}_3 in Eq. (A.1) thus leads to the following equation:

 $A_1 E_1^* + B_1 E_1 + C_1 = 0$.

(A.2)

where

 $A_{c} = \langle \pi - 1 \rangle \langle \pi - g \rangle \langle \pi + g \rangle \tilde{C}^{\dagger}$.

 $\omega_{1} \ = \ -(aa^{*}-aa-g^{*})\,\hat{\omega}\,\hat{\psi} - g(a^{*}-aa^{*}-ag^{*}a^{*}+a^{*}+g^{*}a+g^{*}+g^{*})\;,$

 $C_1 = (\pi - 1) \hat{B} (\hat{B}\hat{C} + g\pi^2 - g\pi - g^3)$.

The above expression is really more symmetric in the matrix elements if one uses the explicit forms of s and π :

 $A_{1} = \frac{\dot{C}M_{g11}M_{g11}(M_{g11} + M_{g11} - 2)}{\dot{C}M_{g11}M_{g11} + M_{g11} - 2}$

$$\begin{split} B_{\ell} &= -\frac{\hat{B}\hat{U}(M_{211}^2 + \delta M_{211}M_{211} - 4M_{211} + M_{211}^2 - 4M_{211})}{4} \\ &= -\frac{(M_{211} - M_{211})}{18} - (M_{111}^2 - 8M_{111}^2 M_{211}^2 + 1M_{211}M_{111}^2) \\ &= -4M_{111}^2 + 7M_{111}^2 M_{111} + M_{211}^2 - 4M_{211}^2) \,, \end{split}$$

 $C_{1} = \frac{-B(M_{p11} + M_{p11} - 2)(2M_{p11}M_{p11}^2 - M_{p11}^2 - 2M_{p11}^2M_{p11} + M_{p11}^2 - 4BC)}{(2M_{p11}M_{p11}^2 - M_{p11}^2 - 4BC)}$

Finally X, Y, and Z in equation Eq. (3.18) are given by

 $X = 2A_l$, $Y = -B_l$, and $Z = B_l^2 - 4A_lC_l$. (A.3)

The complexity of Eqs. (A.2) and (3.19) does not pose any problem in terms of coding into a FORTRAN program and maintaining high degrees of accuracy and efficiency. Eq. (3.18) into equations obtained at earlier stages of the iteration with

fewer elements:

 $L_2 = \frac{L_{2D}}{D}$, $L_3 = \frac{L_{3D}}{D}$, (3.19)

where

 (r, σ_{in})

. ...

$$K_{1D} = -M_{x21} - M_{y21} + M_{x11}K_3 - M_{y11}K_3$$
,

$$_{2D} = -(M_{p12} - M_{p11}L_1)(M_{p22} - M_{p12}K_3 - (M_{p21} - M_{p11}K_3)L_1) +$$

 $(M_{p12} - M_{p11}L_1)(M_{p22} + M_{p12}K_3 - (M_{p21} + M_{p11}K_3)L_1),$

 $K_{2D} \ = \ (M_{g21} + M_{g11}K_3) \left(M_{g22} - M_{g12}K_3 - \left(M_{g21} - M_{g11}K_3 \right) L_1 \right)$

+ $(M_{x21} - M_{x11}K_3)(M_{y22} + M_{y12}K_3 - (M_{y21} + M_{y11}K_3)L_1)$,

 $L_{3D} = -M_{x12} + M_{y12} + M_{x11}L_1 - M_{y11}L_1 ,$

 $D = -M_{e22} + M_{g22} + M_{e12}K_8 + M_{g12}K_8 + M_{e21}L_1 -$

 $M_{\rm y21}L_1 - M_{\rm g11}K_{\rm S}L_1 - M_{\rm y11}K_{\rm S}L_1 \ . \label{eq:masses}$

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- no simple criterion to ensure physical solutions
- no proof the results are always real numbers
- \bullet explicit solution \rightarrow fast numerical calculation

• physical solution (i.e. positive drift lengths) not guaranteed \rightarrow introduction of *Figure Of Merit*

$$FOM = \frac{1}{\sum_{i} L_{i} \cdot \sum_{j} K_{j}}$$

- maximizing FOM = best balance between overall length and integrated Quadrupole strength (chromaticity)
- TELE increases/decreases length of first quadrupole (without changing gradient) till a local maximum of the FOM is found
- physical solution still not guaranteed but more likely
- optimum solution also not guaranteed due to search for local maximum

- O. Napoly, B. Dunham, 'FFADA, Computer Design of Final Focus Systems for Linear Colliders', EPAC94, London, England, June 1994, p. 698
- Y-C Chao, J. Irwin, 'Solution of a Three-Thin-Lens System', SLAC-PUB-5834 (1992)
- B. Zotter, 'Design Considerations for a Chromatically Corrected Final Focus System for TeV Colliders', CLIC-Note-64 (1988)
- O. Napoly, 'Thin Lens Telescopes for Final Focus Systems', CLIC Note 102 (1989)