

# Higgs Couplings and EFT Fits

Tilman Plehn

Universität Heidelberg

Santander, September 2015

## Standard Model operators [SFitter: Gonzalez-Fraile, Klute, TP, Rauch, Zerwas]

- assume: narrow CP-even scalar
- couplings from production & decay rates
- test Lagrangian [essentially non-linear sigma model: Buchalla etal]

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \Delta_W g m_W H W^\mu W_\mu + \Delta_Z \frac{g}{2c_W} m_Z H Z^\mu Z_\mu - \sum_{\tau, b, t} \Delta_f \frac{m_f}{v} H (\bar{f}_R f_L + \text{h.c.})$$

$$+ \Delta_g F_G \frac{H}{v} G_{\mu\nu} G^{\mu\nu} + \Delta_\gamma F_A \frac{H}{v} A_{\mu\nu} A^{\mu\nu} + \text{invisible decays}$$

- electroweak renormalizability through UV completion
- QCD renormalizability not an issue [ask Spirix]
- frequentist likelihood everywhere
- one key issue: theory uncertainties
- **total rates only**

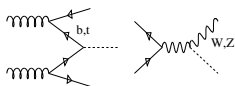
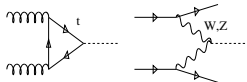
$$\begin{array}{l} gg \rightarrow H \\ qq \rightarrow qqH \\ gg \rightarrow t\bar{t}H \\ qq' \rightarrow VH \end{array}$$

 $\longleftrightarrow$ 

$$g_{HXX} = g_{HXX}^{\text{SM}} (1 + \Delta_X)$$

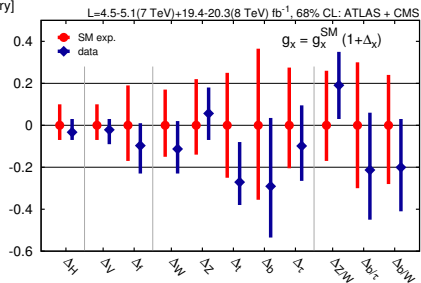
 $\longleftrightarrow$ 

$$\begin{array}{l} H \rightarrow ZZ \\ H \rightarrow WW \\ H \rightarrow b\bar{b} \\ H \rightarrow \tau^+ \tau^- \\ H \rightarrow \gamma\gamma \\ H \rightarrow \tilde{\nu}\tilde{\nu} \end{array}$$



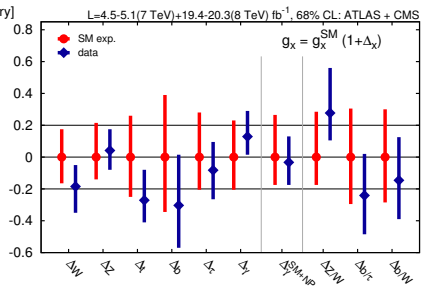
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- assume SM-like [secondary solutions secondary]
- ex: extract  $\Delta_H$  from general fit



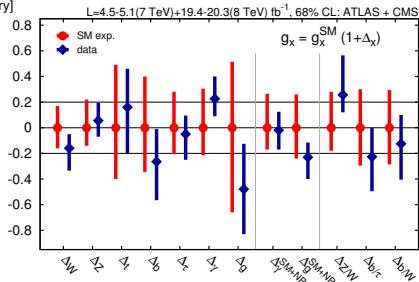
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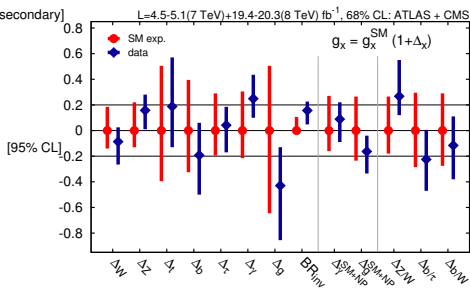
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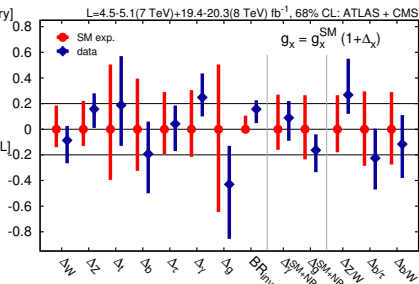
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  - 8 couplings best we can do
- ⇒ **Standard Model within 25%**



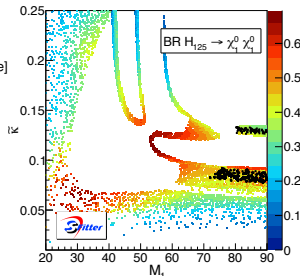
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## Model for invisible Higgs: Hooperon [SFitter: Butter et al]

- NMSSM with singlino dark matter [Ellwanger, ask Maggie]
  - simplified model: pseudo-scalar mediator Majorana dark matter
  - motivated by Fermi galactic center excess
  - different LHC signatures [Cao, Zurek,...]
- ⇒  **$BR_{inv}$  up to 40%**



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## Executive summary

- **couplings fit works great** [experimentally]
  - offers perfect th-ex interface [Cranmer, Kreiss, Lopez-Val, TP]
- (1) has issues with electroweak renormalization
  - (2) only describes total rate changes [theory-defined categories]
  - (3) does not easily replace model fits [correlations]
- ⇒ obvious answer: fit extended Higgs sectors... [that's for Sven]



## Higgs sector effective field theory [following Corbett, Eboli, Gonzalez-Fraile, Goncales-Garcia]

– set of Higgs-gauge operators

$$\mathcal{O}_{GG} = \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} \qquad \mathcal{O}_{WW} = \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \qquad \mathcal{O}_{BB} = \dots$$

$$\mathcal{O}_{BW} = \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \qquad \mathcal{O}_W = (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) \qquad \mathcal{O}_B = \dots$$

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- relevant part after equation of motion, etc

$$\mathcal{L}^{HVV} = - \frac{\alpha_s v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{BB}}{\Lambda^2} \mathcal{O}_{BB} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_{\Phi,2}}{\Lambda^2} \mathcal{O}_{\Phi,2}$$

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- Higgs couplings to SM particles

$$\begin{aligned} \mathcal{L}^{HVV} &= g_g H G_{\mu\nu}^a G^{a\mu\nu} + g_\gamma H A_{\mu\nu} A^{\mu\nu} \\ &+ g_Z^{(1)} Z_{\mu\nu} Z^\mu \partial^\nu H + g_Z^{(2)} H Z_{\mu\nu} Z^{\mu\nu} + g_Z^{(3)} H Z_\mu Z^\mu \\ &+ g_W^{(1)} \left( W_{\mu\nu}^+ W^{-\mu} \partial^\nu H + \text{h.c.} \right) + g_W^{(2)} H W_{\mu\nu}^+ W^{-\mu\nu} + g_W^{(3)} H W_\mu^+ W^{-\mu} + \dots \end{aligned}$$

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- plus Yukawa structure  $f_{\tau,b,t}$

⇒ 9 operators for Run I data

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- 11 Higgs couplings from 9 operators

$$g_g = \frac{f_{GG} v}{\Lambda^2} \equiv -\frac{\alpha_s}{8\pi} \frac{f_g v}{\Lambda^2} \quad g_\gamma = -\frac{g^2 v s_w^2}{2\Lambda^2} \frac{f_{BB} + f_{WW}}{2}$$

$$g_Z^{(1)} = \frac{g^2 v}{2\Lambda^2} \frac{c_w^2 f_{WW} + s_w^2 f_{BB}}{2c_w^2} \quad g_W^{(1)} = \frac{g^2 v}{2\Lambda^2} \frac{f_W}{2}$$

$$g_Z^{(2)} = -\frac{g^2 v}{2\Lambda^2} \frac{s_w^4 f_{BB} + c_w^4 f_{WW}}{2c_w^2} \quad g_W^{(2)} = -\frac{g^2 v}{2\Lambda^2} f_{WW}$$

$$g_Z^{(3)} = M_Z^2 (\sqrt{2} G_F)^{1/2} \left( 1 - \frac{v^2}{2\Lambda^2} f_{\Phi,2} \right) \quad g_W^{(3)} = M_W^2 (\sqrt{2} G_F)^{1/2} \left( 1 - \frac{v^2}{2\Lambda^2} f_{\Phi,2} \right)$$

$$g_f = -\frac{m_f}{v} \left( 1 - \frac{v^2}{2\Lambda^2} f_{\Phi,2} \right) + \frac{v^2}{\sqrt{2}\Lambda^2} f_f$$

- 7 EFT couplings identical to  $\Delta_x$ , suppressed by  $v^2/\Lambda^2$
- 4 EFT couplings  $g_{W,Z}^{(1,2)}$  in addition, suppressed by  $\partial\partial/\Lambda^2$

## Higgs sector effective field theory [following Corbett, Eboli, Gonzalez-Fraile, Goncales-Garcia]

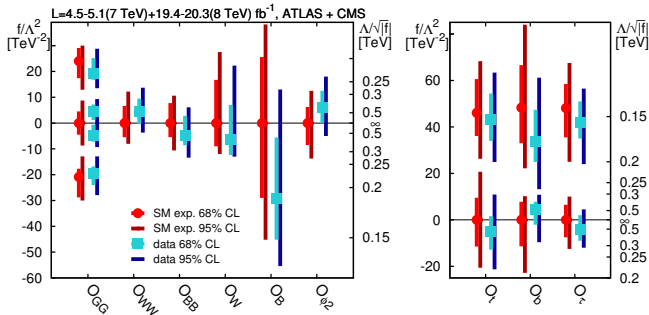
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## SFitter rate analysis

- setup and data identical to  $\Delta_\chi$  fit



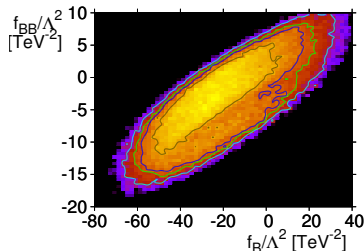
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- correlations through larger basis [problem for #3]
- diagonalization essentially means  $\Delta_x$
- price to pay for theory issue #1?



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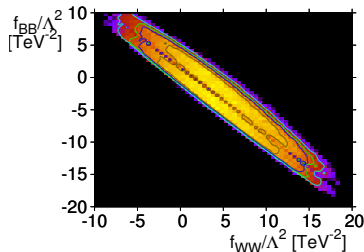
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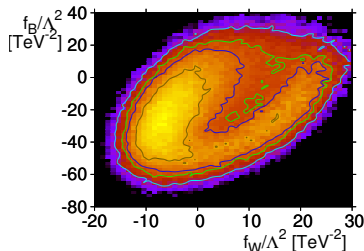
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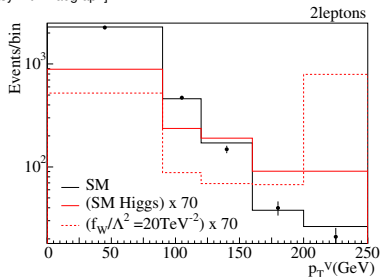
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- $\mathcal{O} \propto \partial\partial/\Lambda^2$  testing  $p_{T,V}$  or  $\Delta\Phi_{jj}$ , #2 [easy with Madgraph]



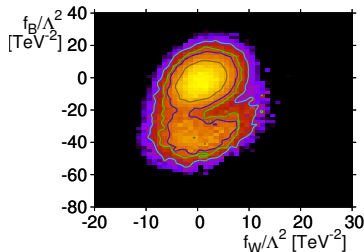
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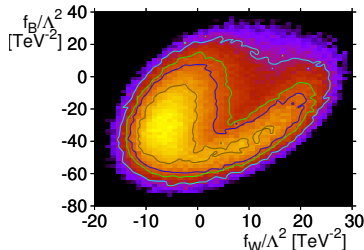
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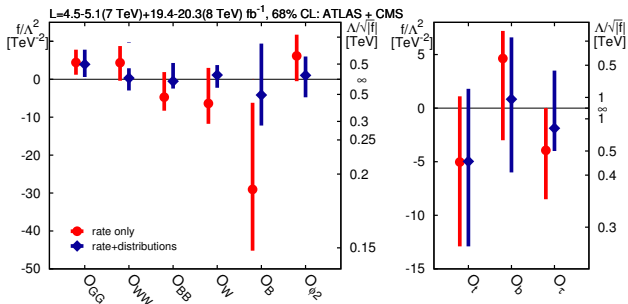
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⇒ Run I legacy



## Width measurements [Kauer & Passarino; Caola & Melnikov; Ellis & Williams]

- peak cross section vs off-shell interference in  $H \rightarrow ZZ$

$$\sigma_{\text{peak}} \sim \frac{g_g^2 g_Z^2}{(s - m^2)^2 + m^2 \Gamma^2} = \frac{g_g^2 g_Z^2}{m^2 \Gamma^2} \quad \sigma_{\text{off}}(g_g g_Z) \sim \sigma_{\text{cont}} - \frac{A_{\text{int}} g_g g_Z}{s - m^2} + \frac{A_H g_g^2 g_Z^2}{(s - m^2)^2}$$

- top–Higgs–gluon sector  $\Delta_t$  vs  $\Delta_g$  or  $f_t$  vs  $f_g$   $[m_{4\ell} \gg m_t > m_H]$

$$\mathcal{M}_{gg \rightarrow ZZ} \sim \pm \frac{m_t^2}{m_Z^2} \log^2 \frac{m_{4\ell}^2}{m_t^2}$$

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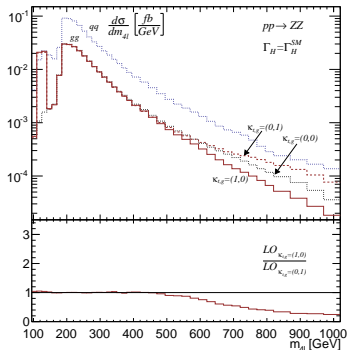
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## Measuring $\Delta_{t,g}$ from $m_{4\ell}$ distributions [Buschmann, Goncalves, Kuttimalai, Schönherr, Krauss, TP]

- simulation: MCFM
- sensitive region  $m_{4\ell} > 500$  GeV



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- peak cross section vs off-shell interference in  $H \rightarrow ZZ$

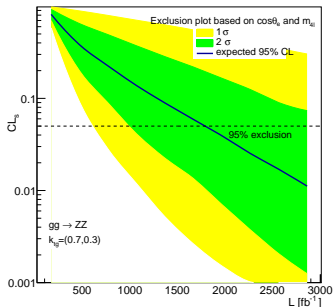
$$\sigma_{\text{peak}} \sim \frac{g_g^2 g_Z^2}{(s - m^2)^2 + m^2 \Gamma^2} = \frac{g_g^2 g_Z^2}{m^2 \Gamma^2} \quad \sigma_{\text{off}}(g_g g_Z) \sim \sigma_{\text{cont}} - \frac{A_{\text{int}} g_g g_Z}{s - m^2} + \frac{A_H g_g^2 g_Z^2}{(s - m^2)^2}$$

- top-Higgs-gluon sector  $\Delta_t$  vs  $\Delta_g$  or  $f_t$  vs  $f_g$   $[m_{4\ell} \gg m_t > m_H]$

$$\mathcal{M}_{gg \rightarrow ZZ} \sim \pm \frac{m_t^2}{m_Z^2} \log^2 \frac{m_{4\ell}^2}{m_t^2}$$

## Measuring $\Delta_{t,g}$ from $m_{4\ell}$ distributions [Buschmann, Goncalves, Kuttimalai, Schönherr, Krauss, TP]

- simulation: MCFM
  - sensitive region  $m_{4\ell} > 500$  GeV
  - most optimistic: statistics only  
 $H \rightarrow ee\mu\mu$  analysis  
 2D likelihood study of  $\cos \theta_e, m_{4\ell}$
- $\Rightarrow \Delta_t = -0.3$  to 95% CL with  $1700 \text{ fb}^{-1}$





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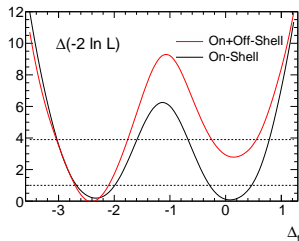
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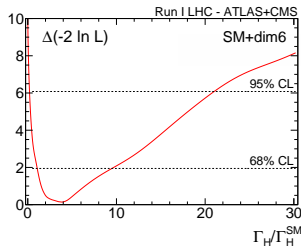
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  - ...either improving  $\Delta_t$  or  $f_t$  measurement
  - ...or measuring unobserved Higgs decays
- eventually a measured distribution



## Complete models vs EFT signatures [Brehmer, Freitas, Lopez-Val, TP]

- push **models** to visible deviations at 13 TeV  
Higgs portal, 2HDM, stops, vector triplet
- simulate distributions in full models  
 $H \rightarrow \gamma\gamma, 4\ell, \text{WBF}, VH, HH$
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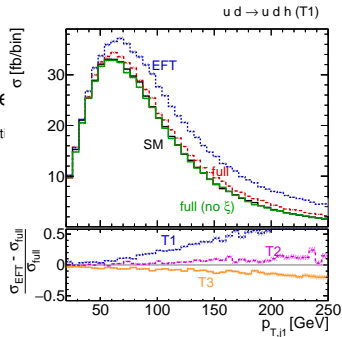
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**Tagging jet observables** [Brehmer, Jäckel, TP]

- polarization defined in Higgs frame
- transverse momenta

$$P_T(x, p_T) \sim \frac{1 + (1 - x)^2}{x} \frac{p_T^3}{((1 - x)m_W^2 + p_T^2)^2}$$

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**Tagging jet observables** [Brehmer, Jäckel, TP]

- polarization defined in Higgs frame
- transverse momenta
- azimuthal angle

$$A_\phi = \frac{\sigma(\Delta\phi_{jj} < \frac{\pi}{2}) - \sigma(\Delta\phi_{jj} > \frac{\pi}{2})}{\sigma(\Delta\phi_{jj} < \frac{\pi}{2}) + \sigma(\Delta\phi_{jj} > \frac{\pi}{2})}$$

# Longitudinal WW scattering

## WW scattering at high energies [Han et al; Dawson]

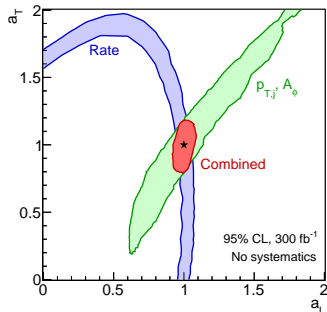
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## Tagging jet observables [Brehmer, Jäckel, TP]

- polarization defined in Higgs frame
  - transverse momenta
  - azimuthal angle
  - total rate  $\sigma \sim (A_L a_L^2 + A_T a_T^2)$
- ⇒ simple question, clear answer



## Higgs couplings

- **couplings fit works great** [experimentally]
  - offers perfect th-ex interface [Cranmer, Kreiss, Lopez-Val, TP]
- (1) has issues with electroweak renormalization
  - (2) only describes total rate changes [theory-defined categories]
  - (3) does not easily replace model fits [correlations]

## Higgs effective theory

- is harder than  $\Delta_x$  for  $v^2/\Lambda^2$   
describes distributions though  $\partial\partial/\Lambda^2$
- is easy to simulate through MC
- currently excludes D8 operators ex cathedra
- will hardly replace model fits [correlations and matching]
- **explains why nothing new happens with  $\Lambda < 400$  GeV**

Lectures on LHC Physics, Springer, arXiv:0910.4182 updated under [www.thphys.uni-heidelberg.de/~plehn/](http://www.thphys.uni-heidelberg.de/~plehn/)

Much of this work was funded by the BMBF Theorie-Verbund which is ideal for relevant LHC work



Higgs Fits

Tilman Plehn

Couplings

Operators

Off-shell

**Limitations**