

# Updated measurement of gamma-ray anisotropies

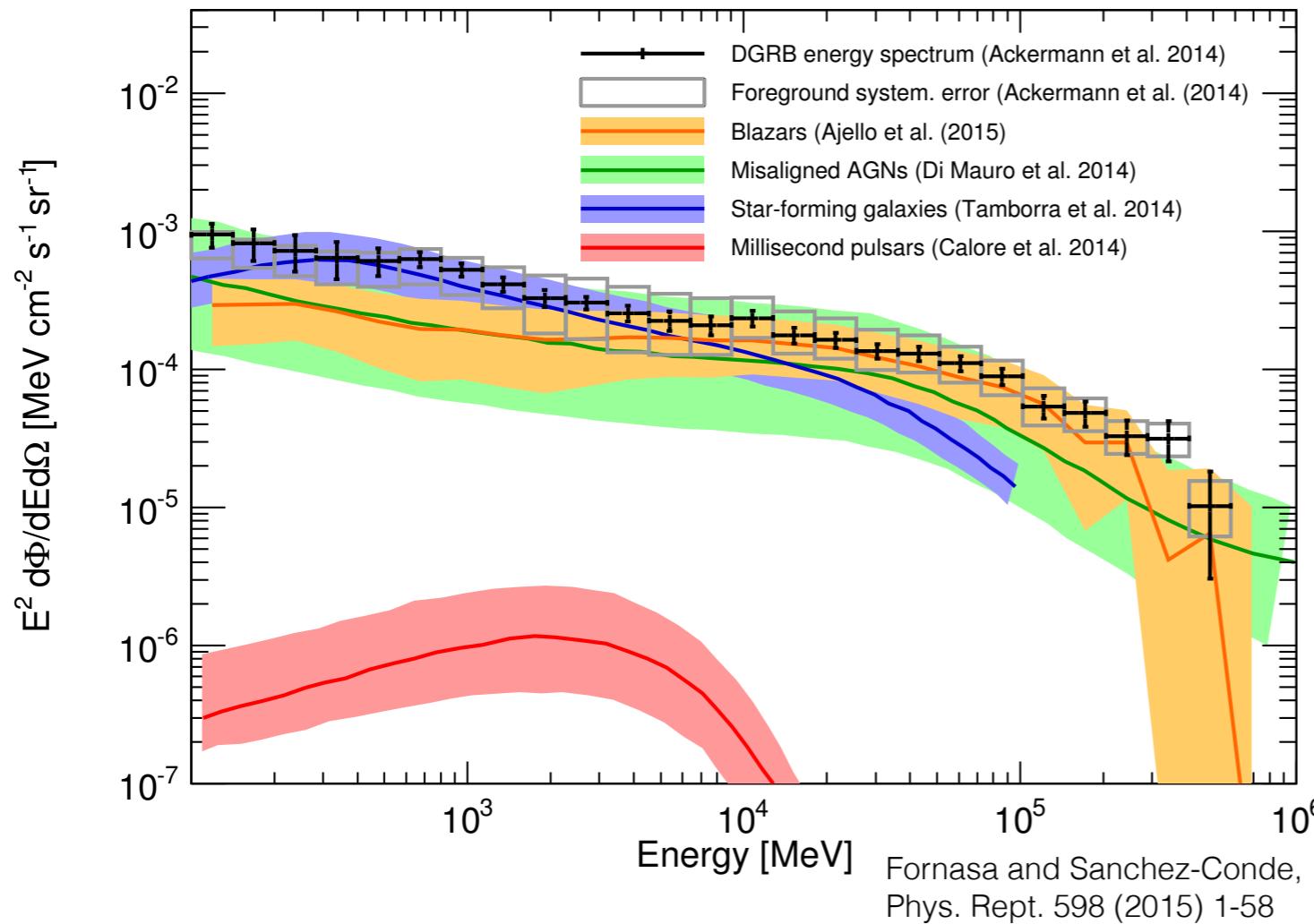
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In collaboration with A. Cuoco, J. Zavala, J. Gaskins, M. A. Sanchez-Conde,  
G. Gomez-Vargas, E. Komatsu, T. Linden, F. Prada, F. Zandanel and A. Morselli



# Anisotropies in the Diffuse Gamma-Ray Background



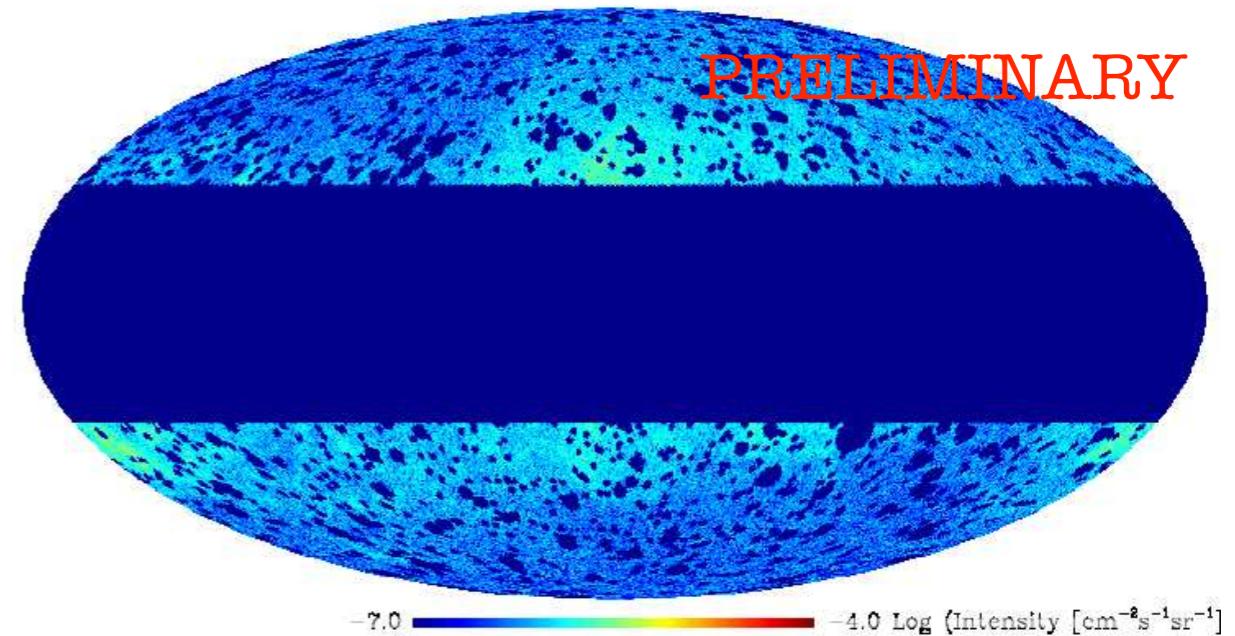
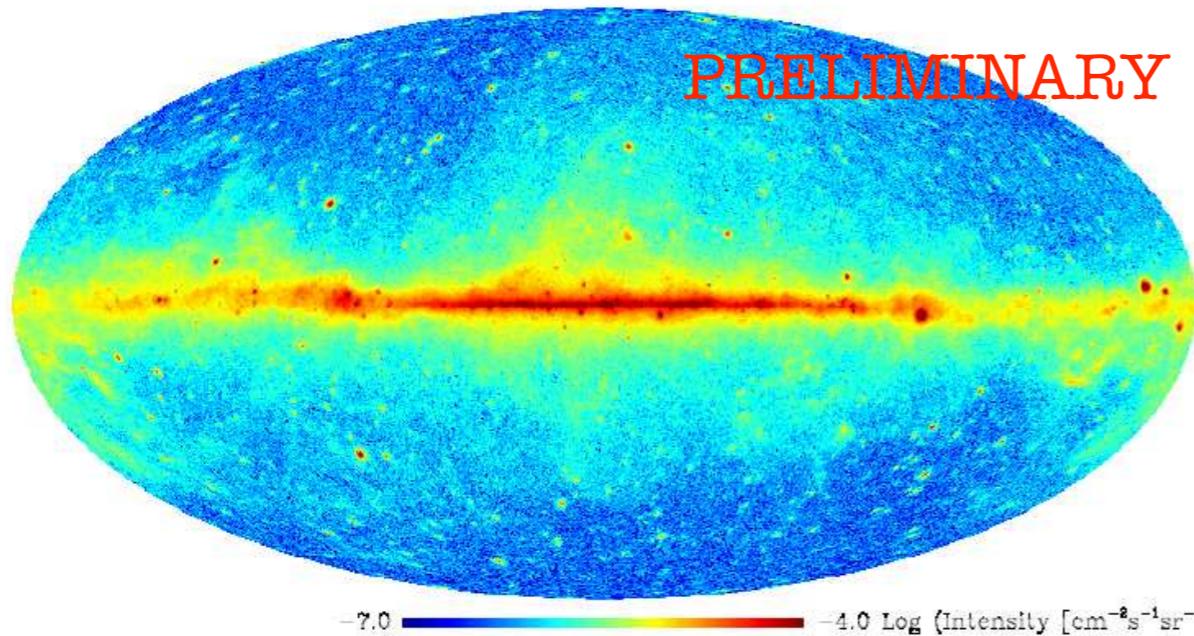
- cumulative emission of unresolved sources
- guaranteed components from unresolved astrophysical sources
- constraints on additional contributors (Dark Matter)

$$I(\psi) = \sum_{\ell m} a_{\ell,m} Y_{\ell,m}(\psi)$$

$$C_\ell = \frac{1}{2\ell + 1} \sum_{m=-\ell}^{\ell} |a_{\ell m}|^2$$

- measure  $C_\ell$  (update the 2012 detection by Fermi-LAT)
- develop a model of  $C_\ell$  in terms of astrophysical sources to fit the data

# New APS measurement



New measurement

81 months

Pass 7 reprocessed  
(ULTRACLEAN\_v15) front

13 energy bins  
between 0.5-500 GeV

masking sources in 3FGL

Ackermann et al. (2012)

22 months

Pass 6 (DIFFUSE\_v3) front and  
back

4 energy bins  
between 1-50 GeV

masking sources in 1FGL

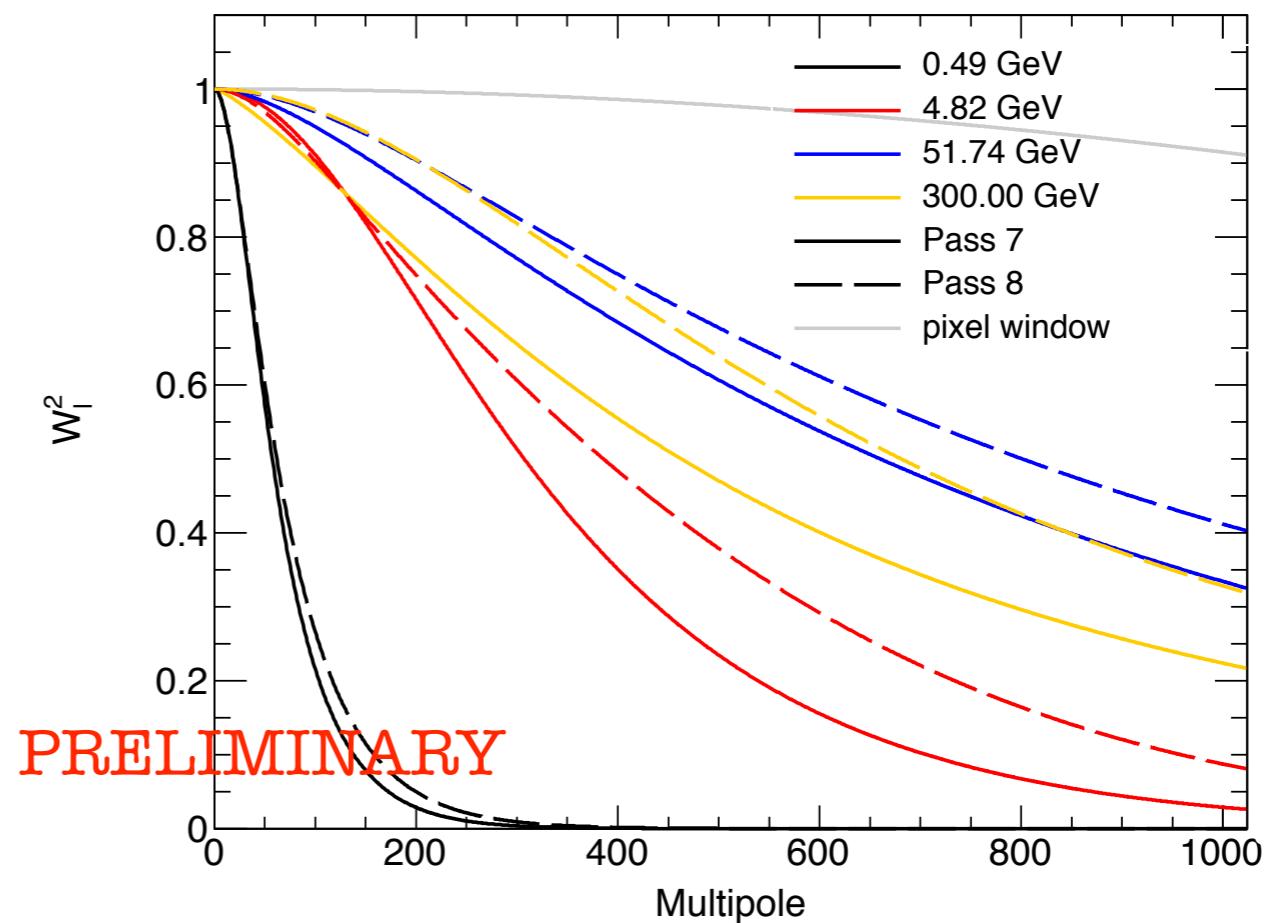
# APS estimator

output of the decomposition in spherical harmonics (already corrected for the effect of the mask)

window beam function  
(it corrects for the experimental PSF)

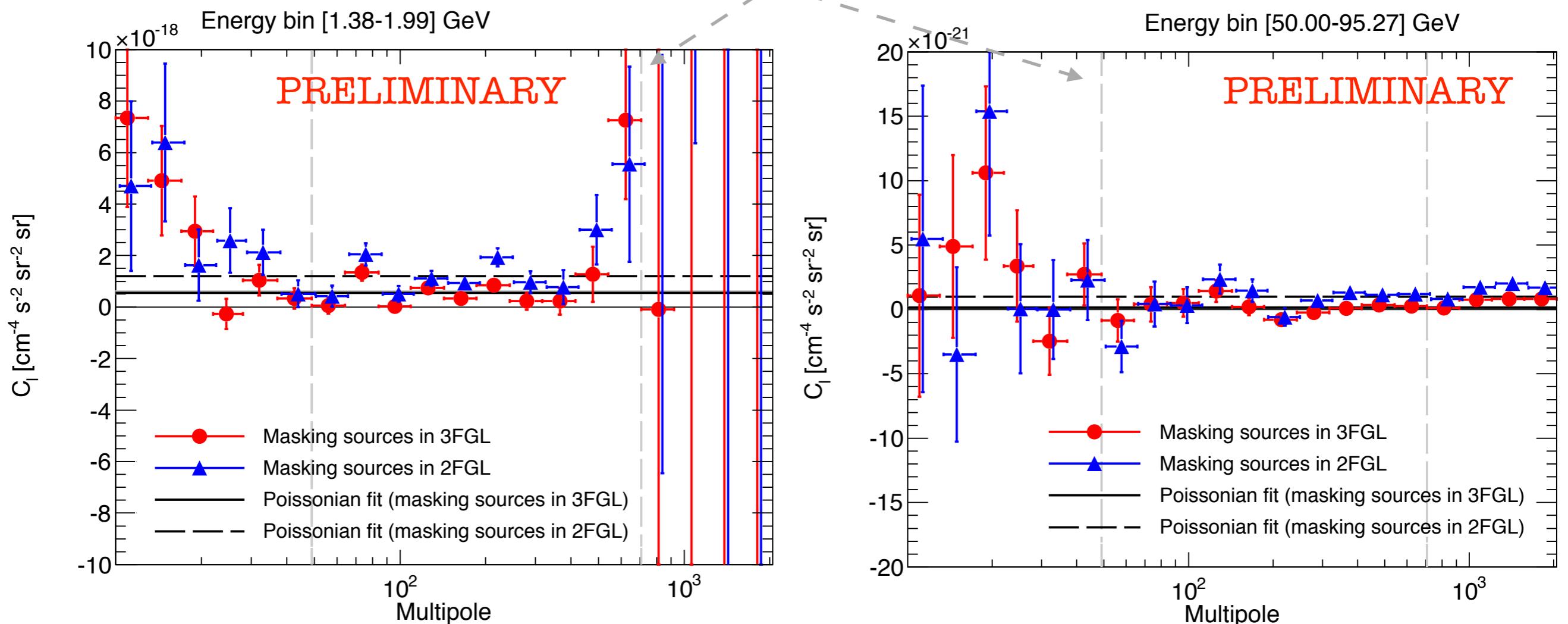
$$C_{\ell}^{\text{signal},ij} = \frac{C_{\ell}^{\text{Pol},ij} - C_N}{(W_{\ell}^{\text{beam},i} W_{\ell}^{\text{beam},j}) (W_{\ell}^{\text{pix}})^2}$$

photon noise  
(inversely proportional to the number of detected photons)



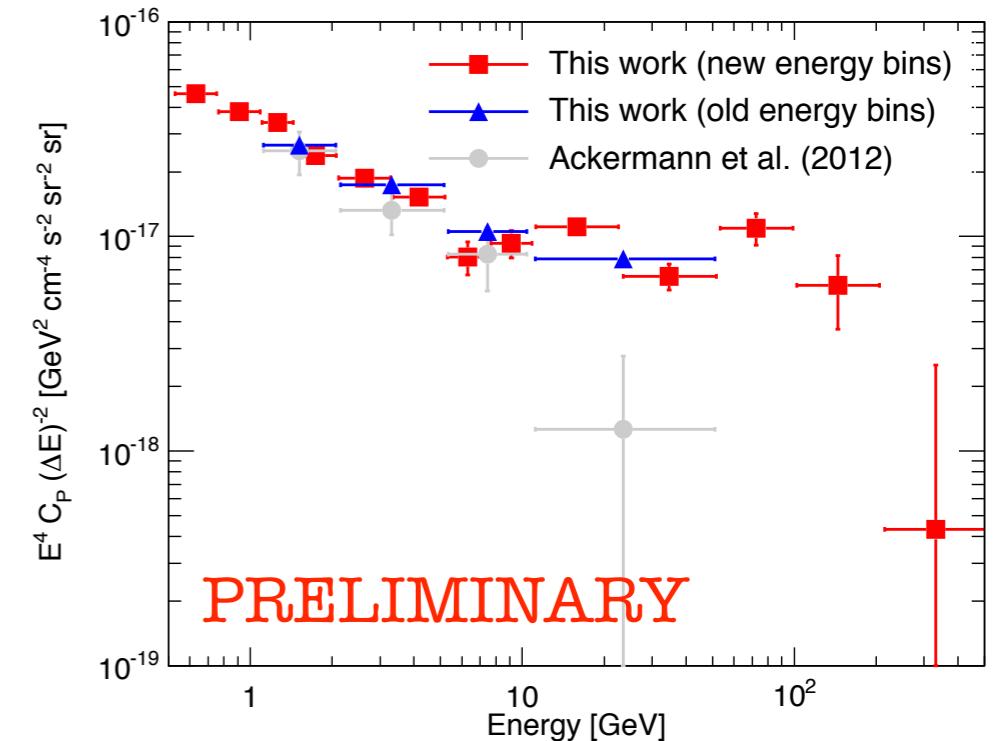
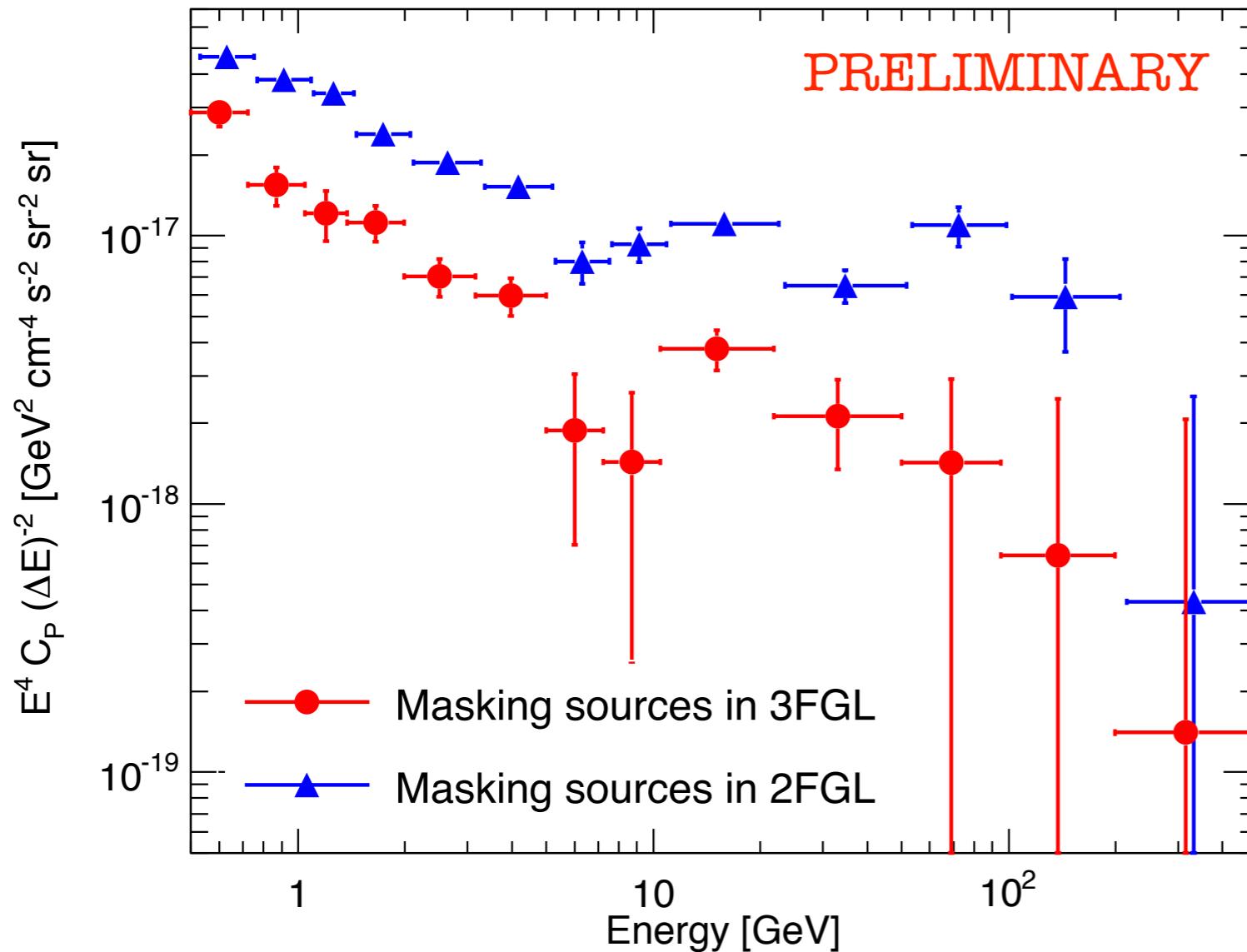
# Binned APS measurement

signal region between  $\ell=49$  and 706



- contamination of Galactic foreground at low  $\ell$  and effect of the beam window function at large  $\ell$
- fitting the data with a Poissonian APS:  $\chi^2/\text{dof} = 1.01$ ,  $p\text{-value}=0.61$
- fits with  $A(\ell/\ell_0)^\alpha$  and  $C_P + A(\ell/\ell_0)^\alpha$  have also been considered

# Anisotropy energy spectrum



$$I(\psi, E_i) = \sum_{\ell,m} a_{\ell,m}^i Y_{\ell,m}(\psi)$$

$$C_{\text{P}}^i = \sum_{\alpha} C_{i\text{P},\alpha} = \sum_{\alpha} I_{\alpha}^2(E_i) \tilde{C}_{\text{P},\alpha}$$

- anisotropy energy spectrum traces the intensity energy spectrum of sources
- features in the anisotropy energy spectrum hint at multiple components

# Cross-correlation APS

$$C_{\ell}^{ij} = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} a_{\ell m}^i a_{\ell m}^{j*}$$

- 91 independent combination of en. bins: 91 Poissonian  $C_P^{i,j}$

- cross correction coefficients

$$C_P^{i,j} / \sqrt{C_P^{i,i} C_P^{j,j}}$$

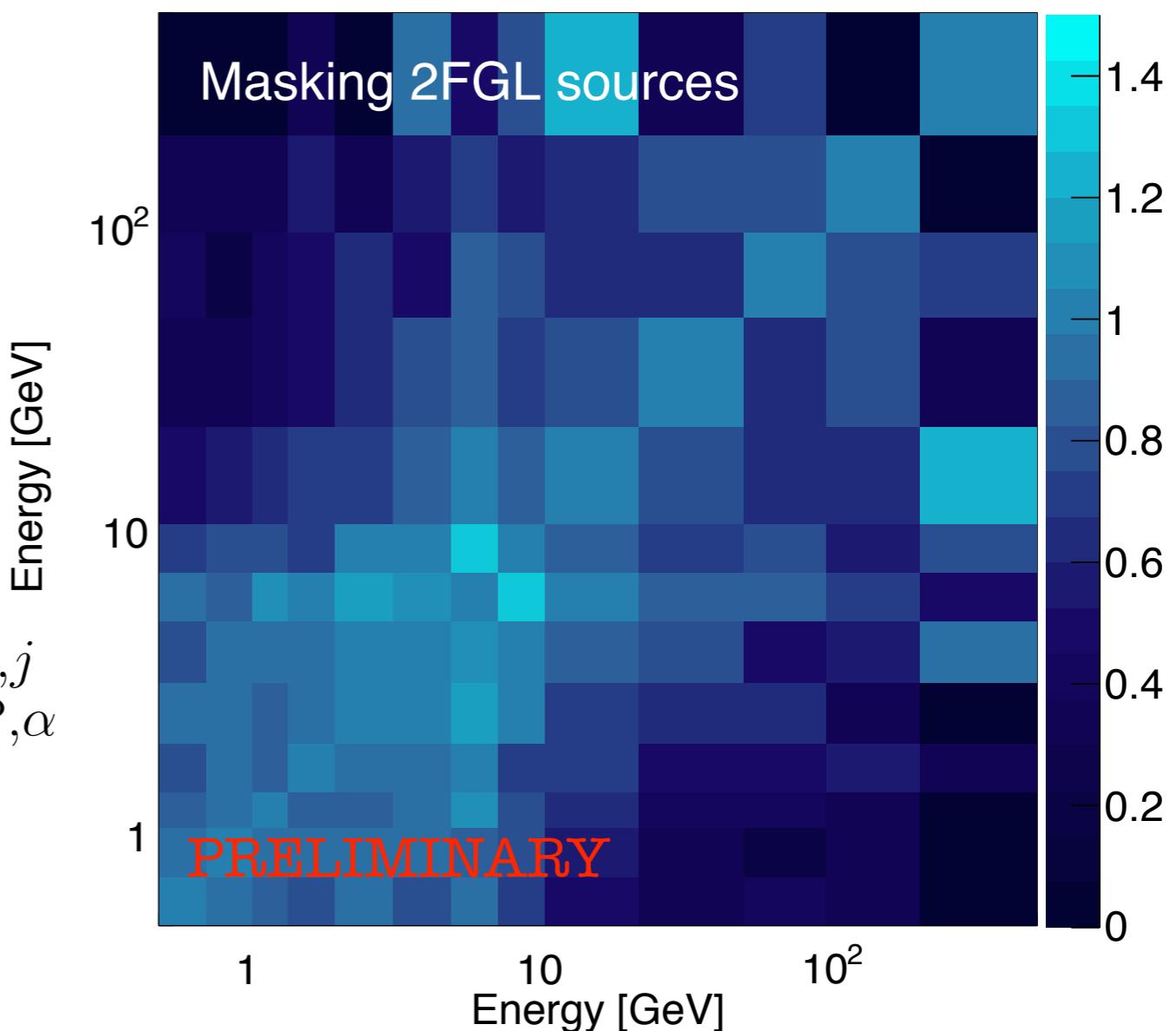
- one source class:

$$C_P^{i,j} = I(E_i)I(E_j)\tilde{C}_P$$

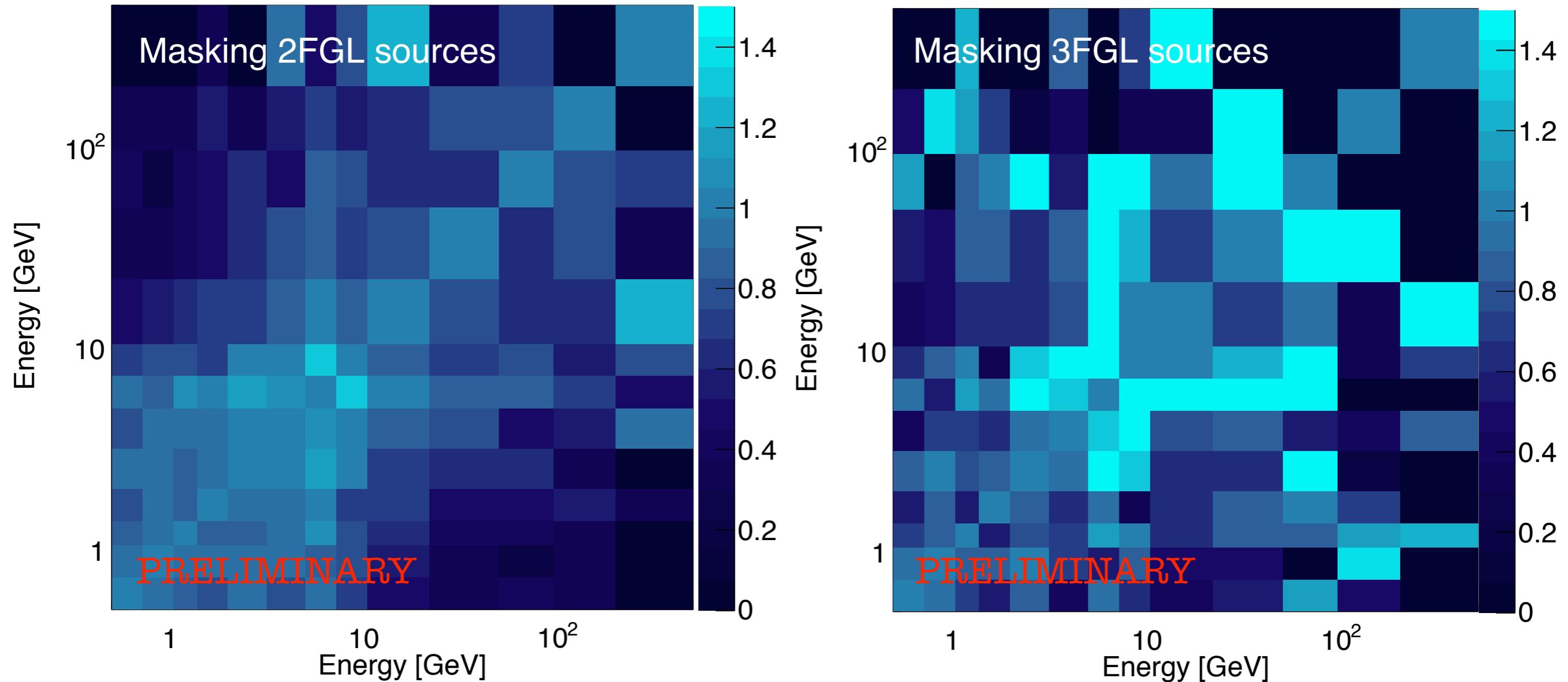
- multiple source classes:

$$C_P^{i,j} = \sum_{\alpha} C_{P,\alpha}^{i,j} = \sum_{\alpha} I(E_i)I(E_j)\tilde{C}_{P,\alpha}^{i,j}$$

- cross-correlation coefficients different than 1.0 hint at multiple components



# Cross-correlation APS

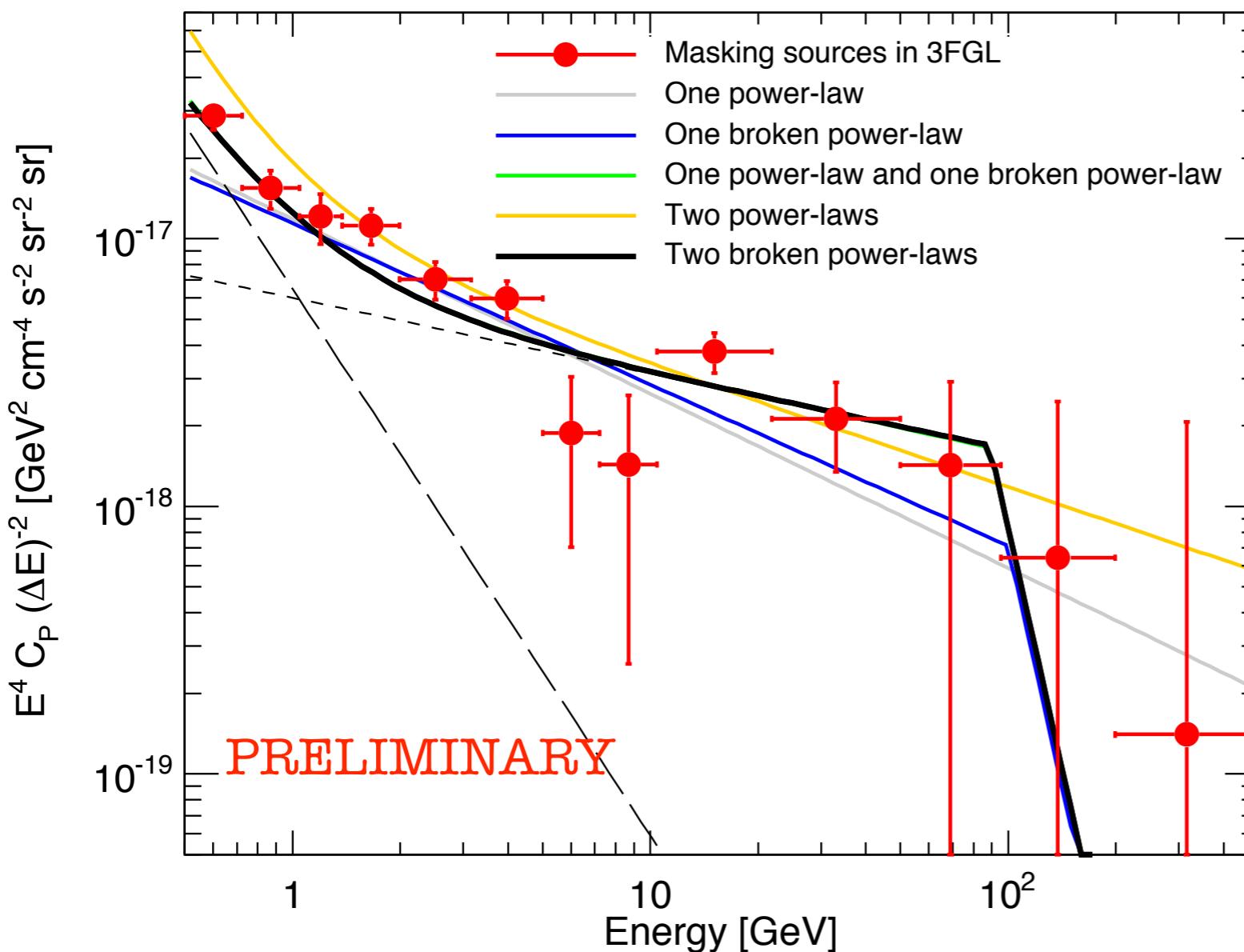


# Interpretation in terms of multiple populations

Fitting the data with one or more populations, assuming specific energy spectra:

$$I(E) \propto E^{-\alpha}$$

$$I(E) \propto \begin{cases} (E/E_0)^{-\alpha} & \text{if } E \leq E_b \\ (E_0/E_b)^{-\alpha+\beta}(E/E_0)^{-\beta} & \text{otherwise} \end{cases}$$



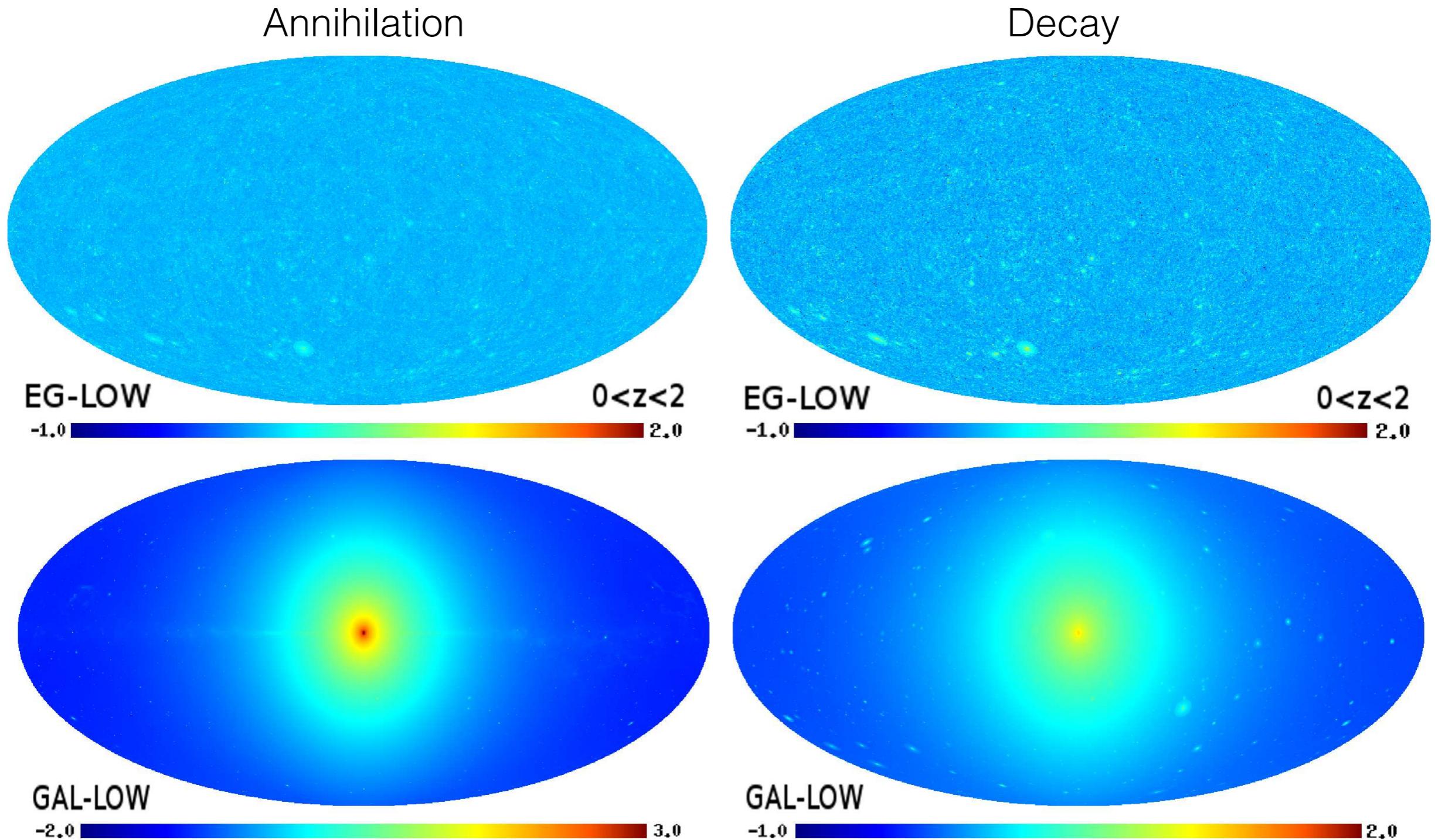
Best-fit model has two contributions both emitting as broken power laws:

- $E_b = (88.9_{-14.4}^{+9.6}) \text{ GeV}$ ,  $\alpha = 2.15 \pm 0.05$ ,  $\beta > 3.9$

- $E_b > 79 \text{ GeV}$ ,  
 $\alpha = 3.0_{-0.2}^{+0.3}$ ,  
 $\beta = 0.88_{-0.15}^{+0.09}$

$\chi^2/\text{dof} = 1.21$ ,  $p\text{-value} = 0.16$

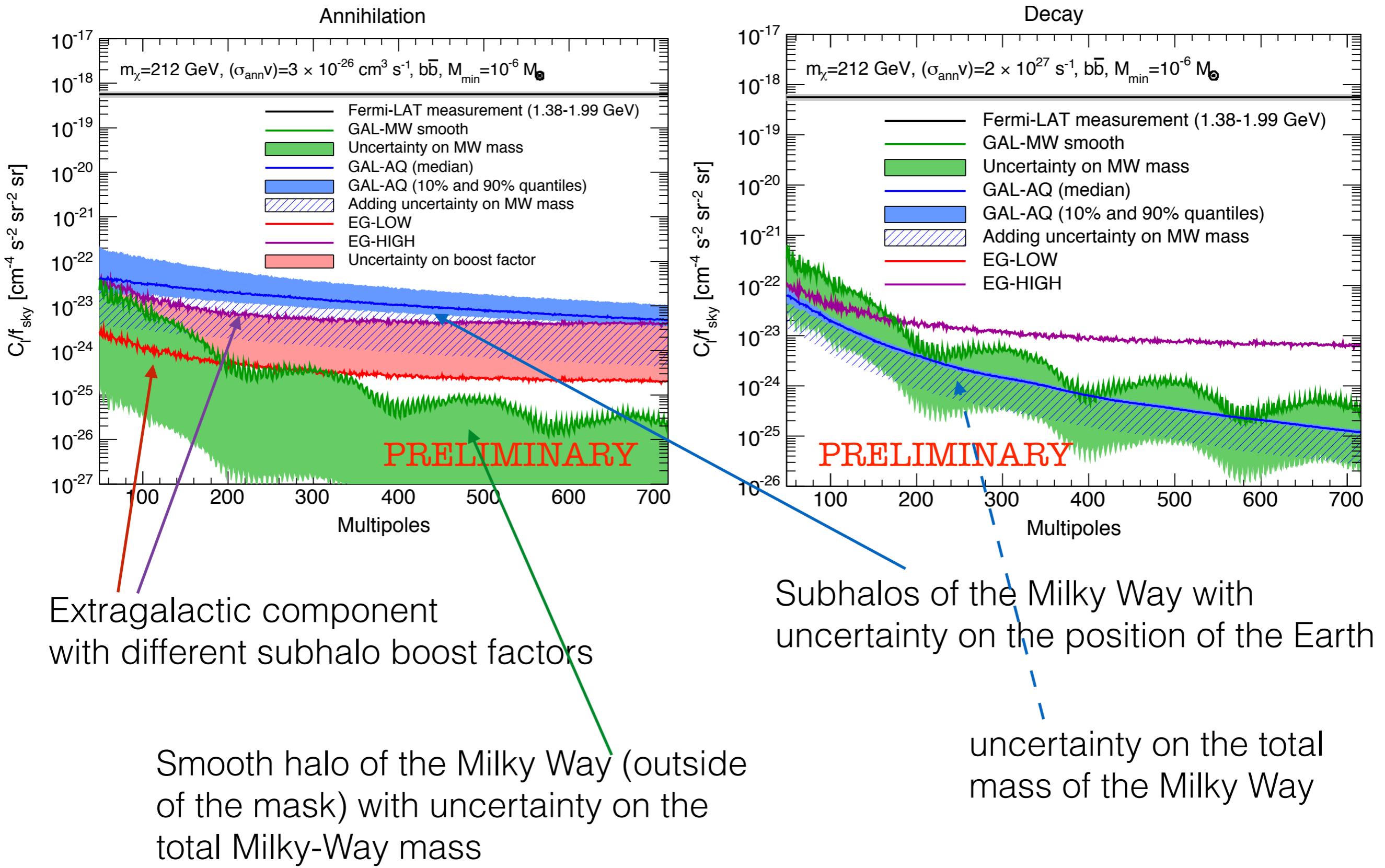
# Gamma-ray anisotropies from Dark Matter



$E=4 \text{ GeV}$ ,  $M_{\min}=10^{-6} M_{\odot}$ ,  $b$  quarks

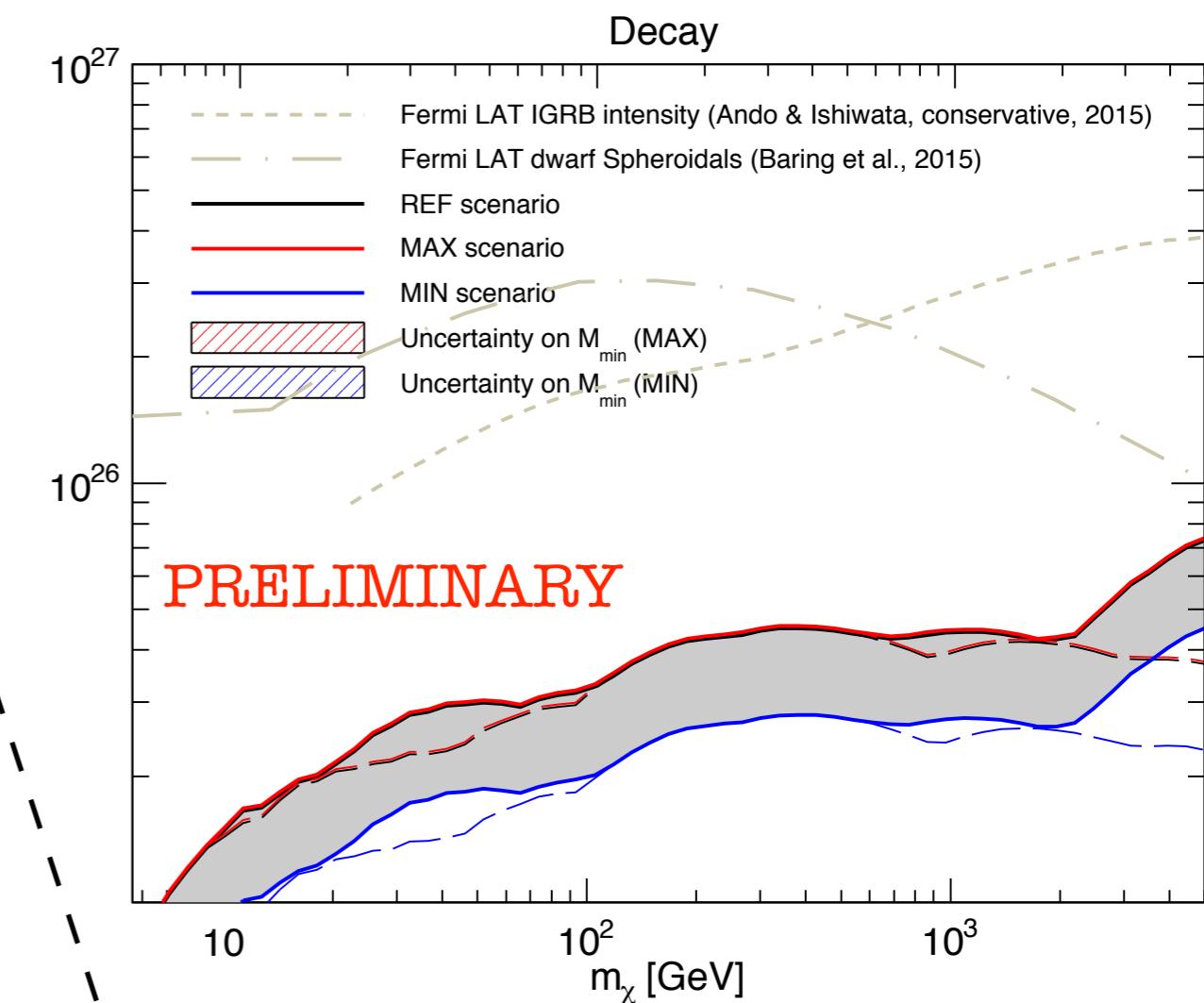
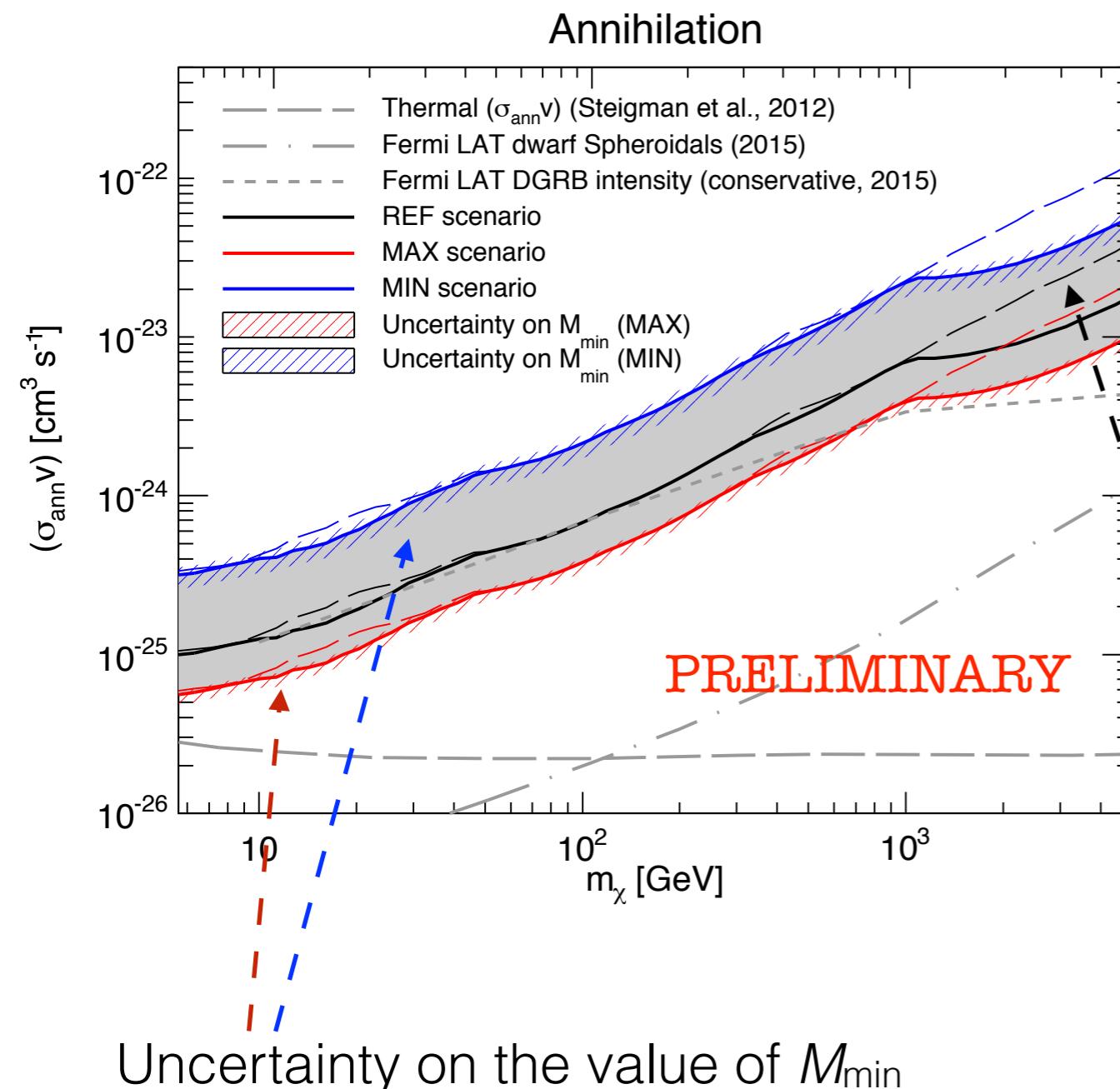
$m_{\chi}=200 \text{ GeV}$ ,  $\sigma v=3 \times 10^{-26} \text{ cm}^3 \text{s}^{-1}$  (annihilation),  $m_{\chi}=2 \text{ TeV}$ ,  $\tau=2 \times 10^{27} \text{ s}$  (decay)

# DM-induced APS



# Conservative exclusion limits

$$\langle C_{\ell, \text{DM}}^{i,j} \rangle < C_{\text{P}}^{i,j} + 1.64 \sigma_{C_{\text{P}}^{i,j}}$$

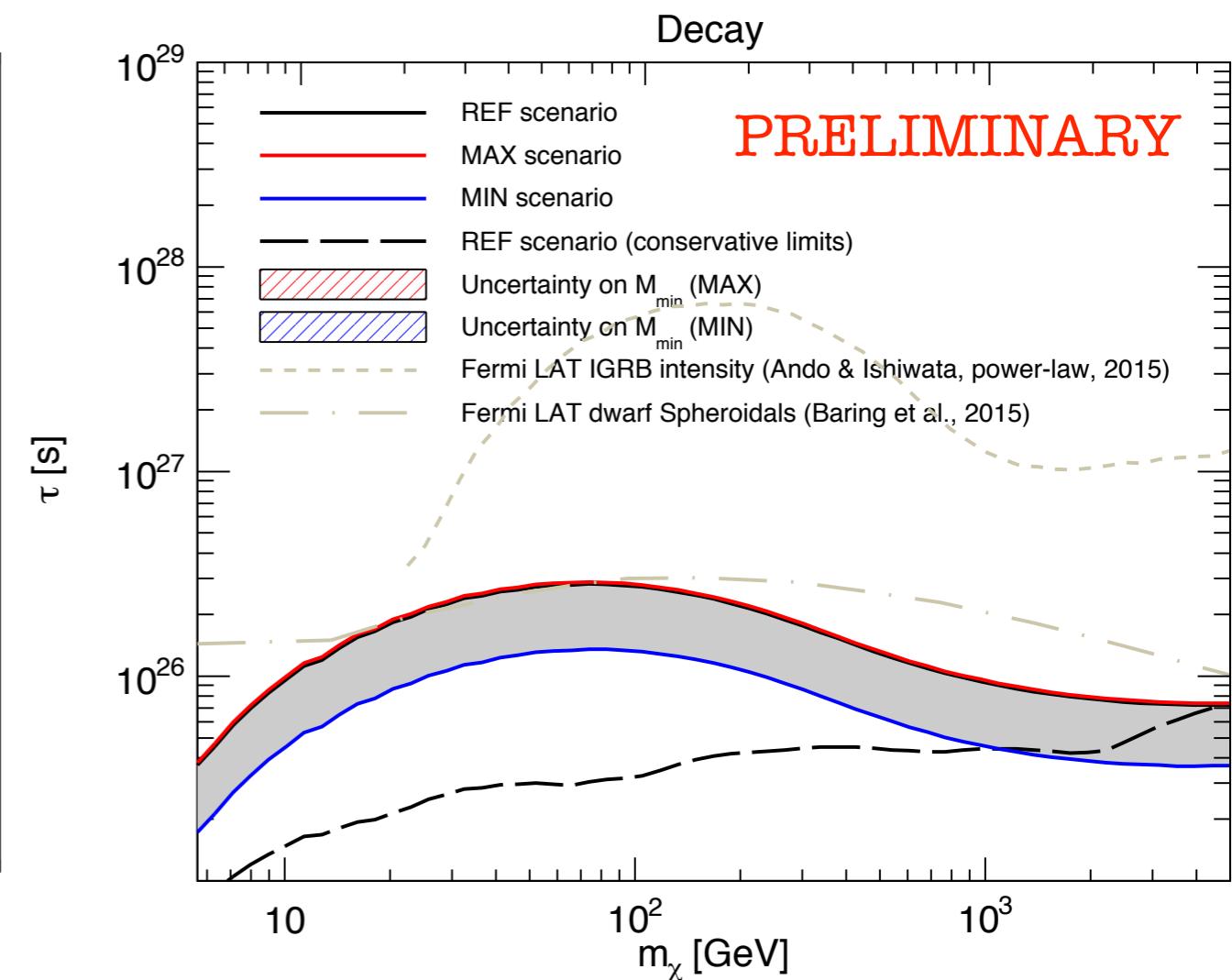
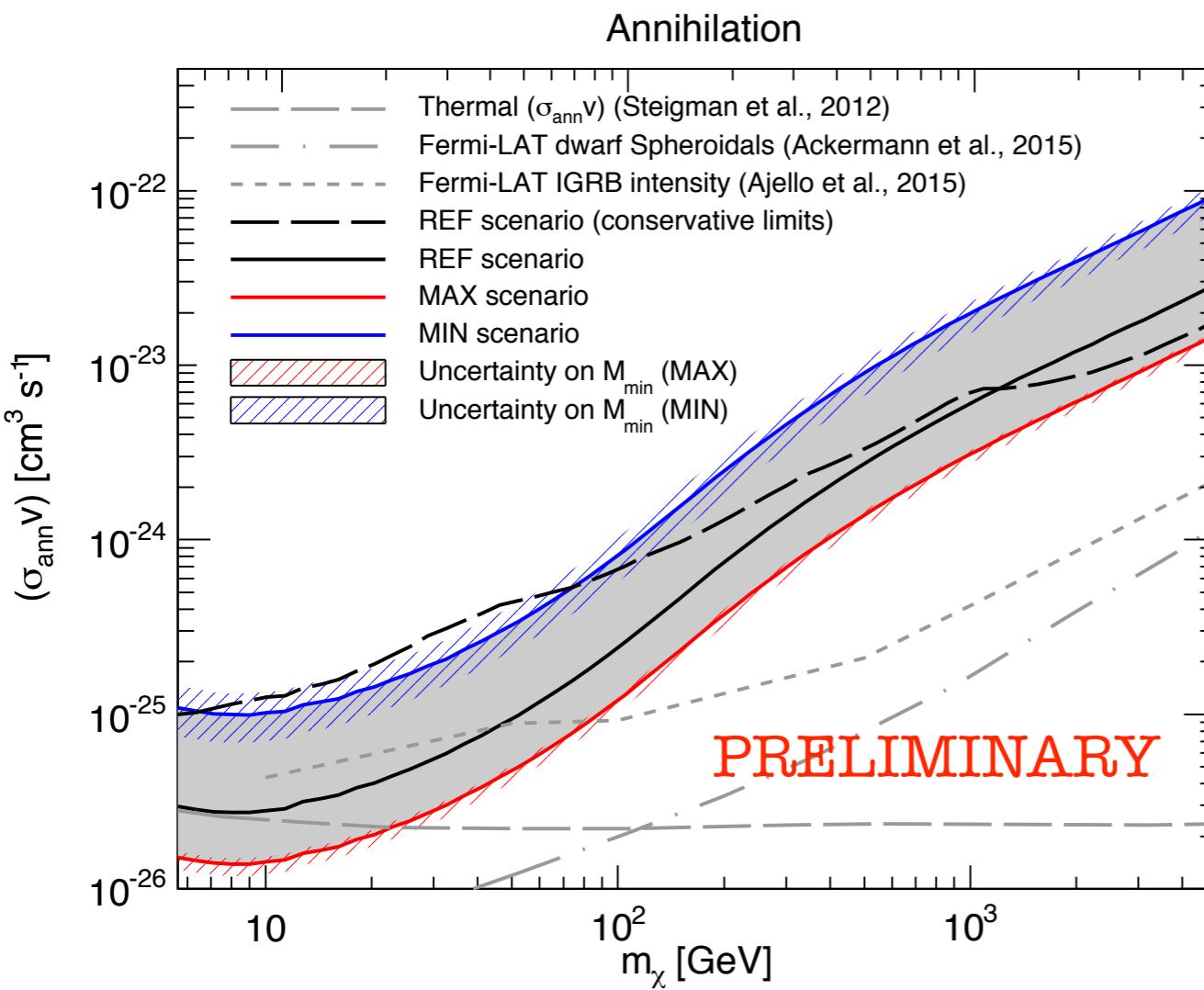


How the exclusion limits would look like if only the auto-correlation APS were used

# 2-component fit to the binned APS

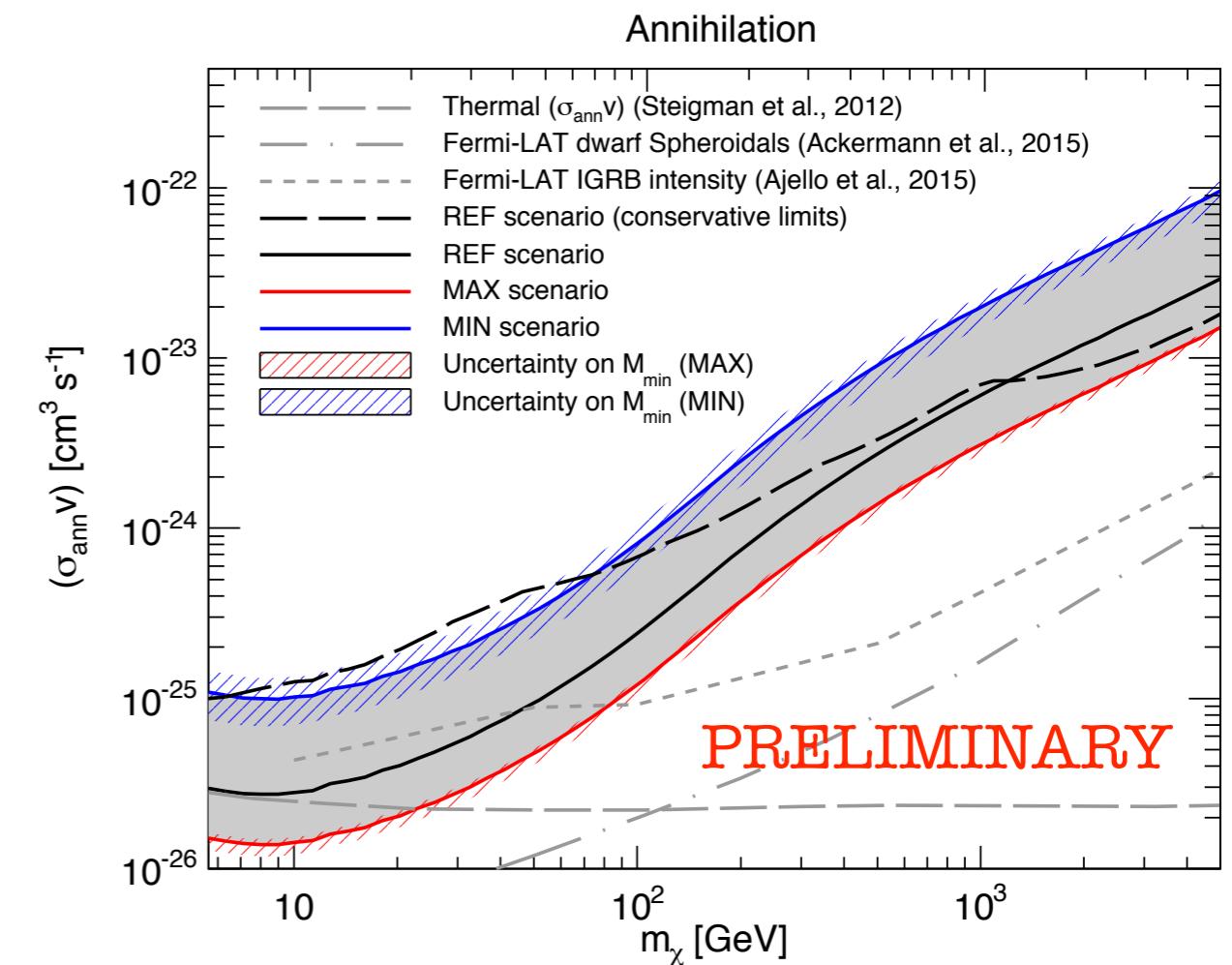
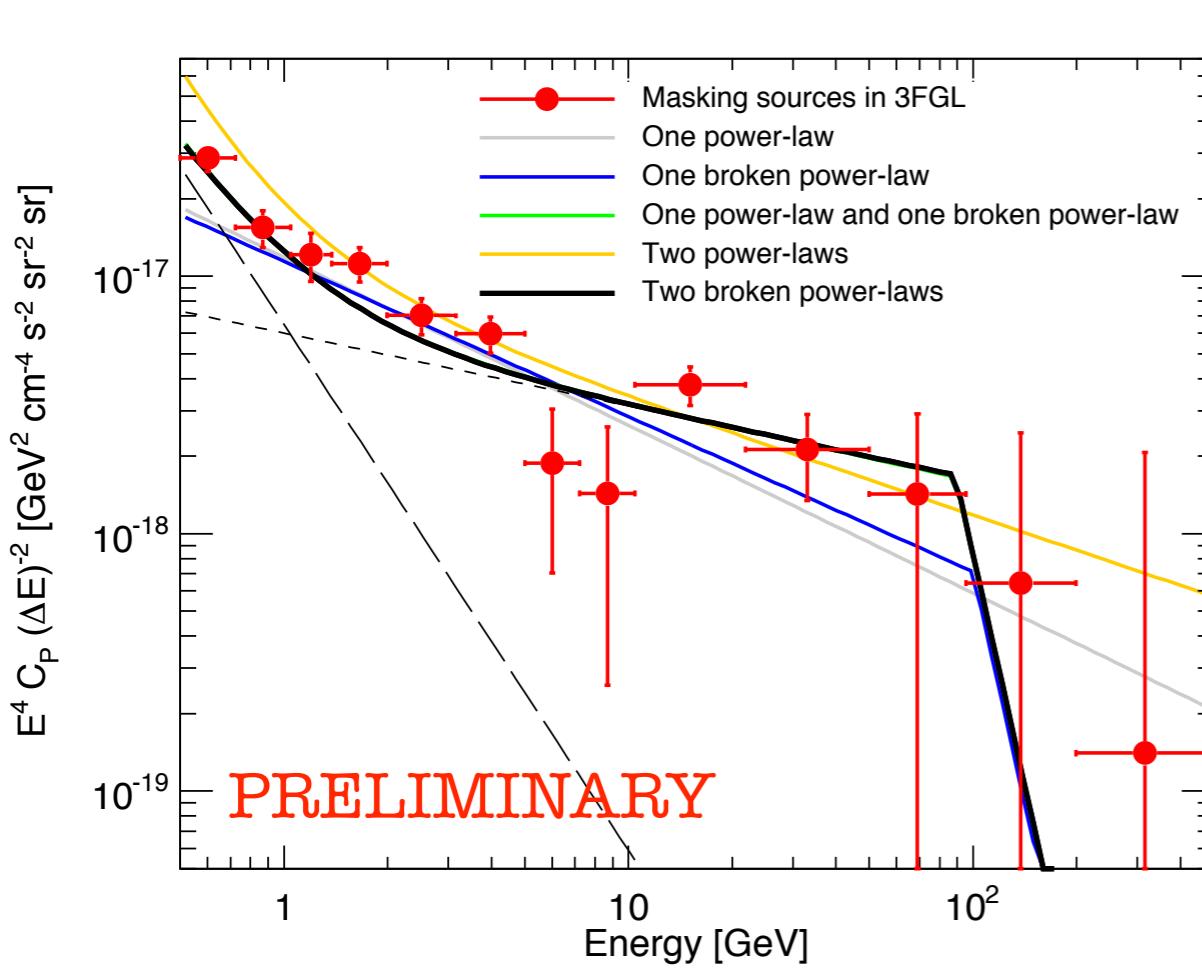
$$\chi^2 = - \sum_{i,j,\ell} \frac{[C_\ell^{i,j} - C_{\ell, \text{DM}}^{i,j} - C_{\text{P}}^{i,j}]^2}{\sigma_{C_\ell^{i,j}}^2}$$

95% CL exclusion limit when Test Statistics  $\Delta\chi^2=3.84$



# Conclusions

- updated measurement of anisotropy angular power spectrum
- new features, possible indication of multiple components
- impact on understanding of unresolved astrophysical sources
- limits on DM competitive with those coming from the overall intensity



# How to bin the APS

- produce 100 Monte Carlo realisations of the gamma-ray sky with a fixed nominal  $C_P$
- PolSpice computes  $C_\ell$  and estimates errors and covariances
- analytical expression for the error is

$$\sigma_\ell = \sqrt{2/(2\ell + 1)} \left( C_\ell + \frac{C_N}{W_\ell^2} \right)$$

- to bin  $C_\ell$  in one multipole bin, you can compute:
  - A. unweighted average**
  - weighted average with weight =  $1/\sigma_\ell$
  - weighted average with weight =  $1/\sigma_\ell$  and only photon noise
- Monte Carlo simulations prove that method B underestimates the APS
- method B was used in Ackermann et al. (2012)

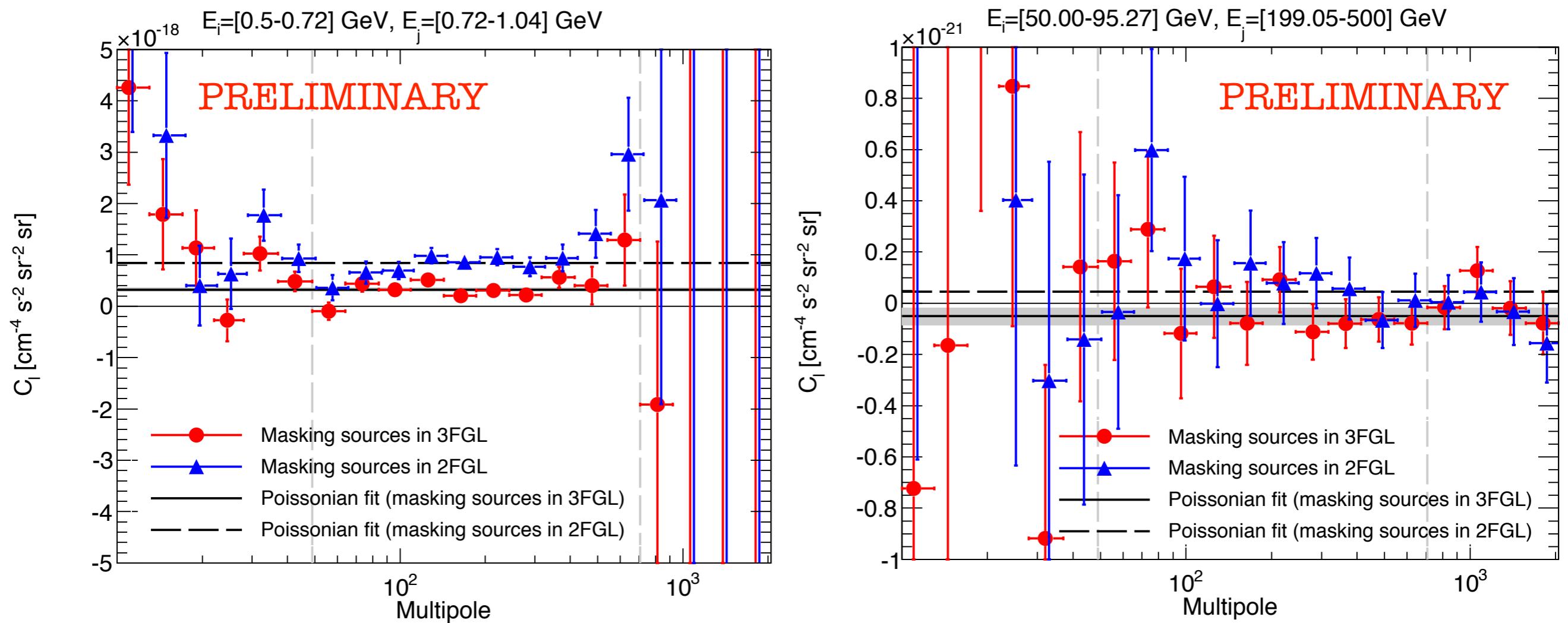
# How to estimate the error of the binned APS

- method A: average of the analytical expression for the error  $\sigma_\ell$

$$\sigma_\ell = \sqrt{2/(2\ell + 1)} \left( C_\ell + \frac{C_N}{W_\ell^2} \right)$$

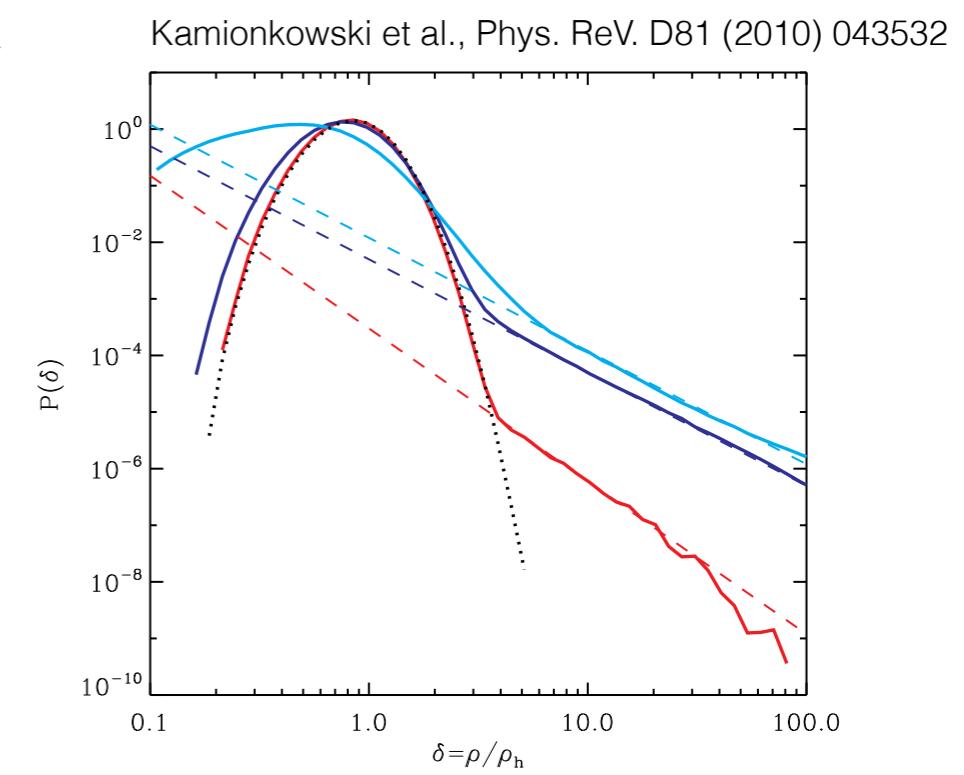
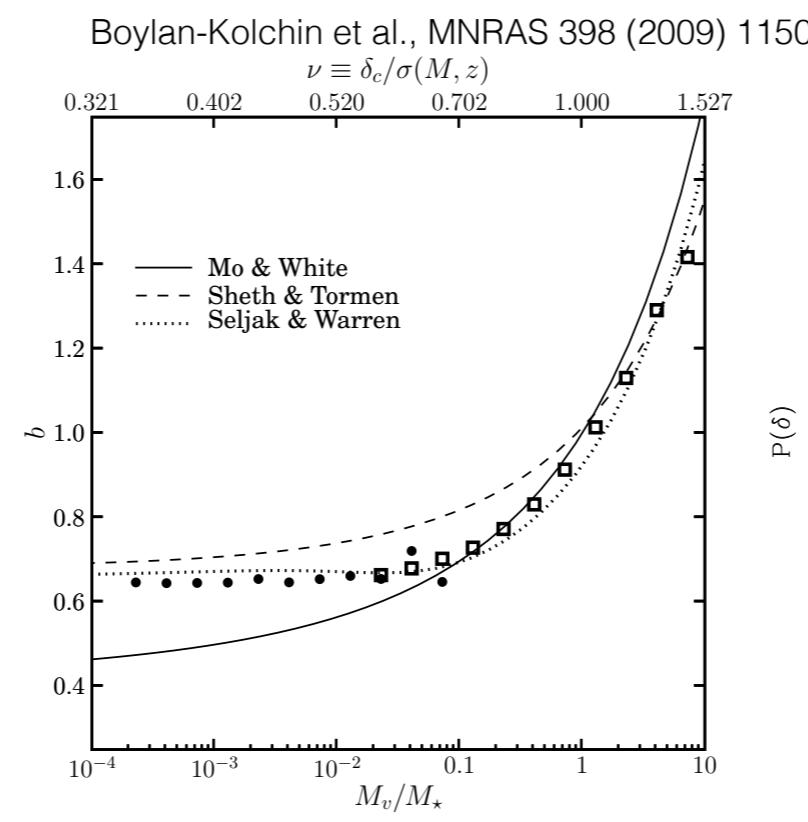
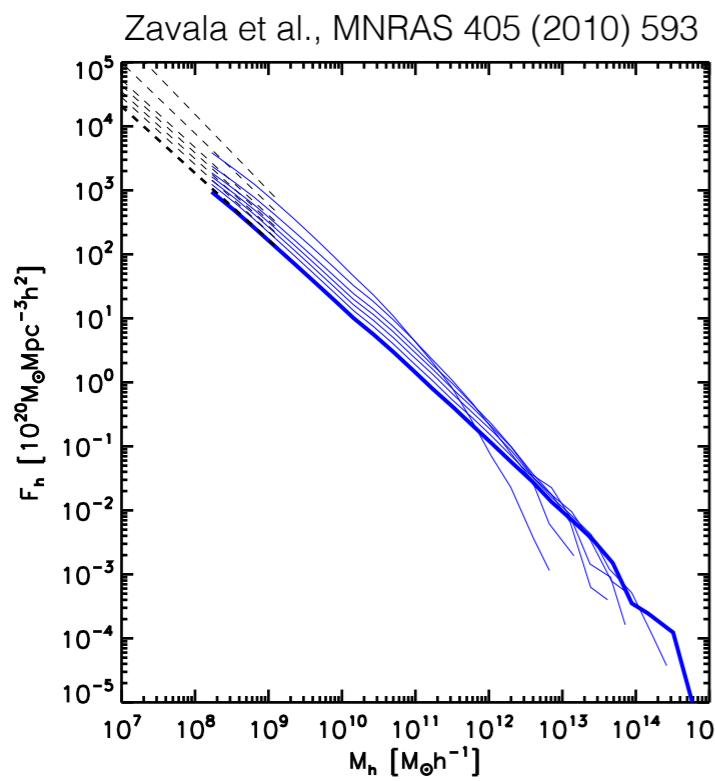
- method B: average of the variances and covariances computed by PolSpice
- the two methods agree
- the estimated error describes well the distribution of the binned  $C_\ell$  from the 100 Monte Carlo realisations

# Cross-correlation



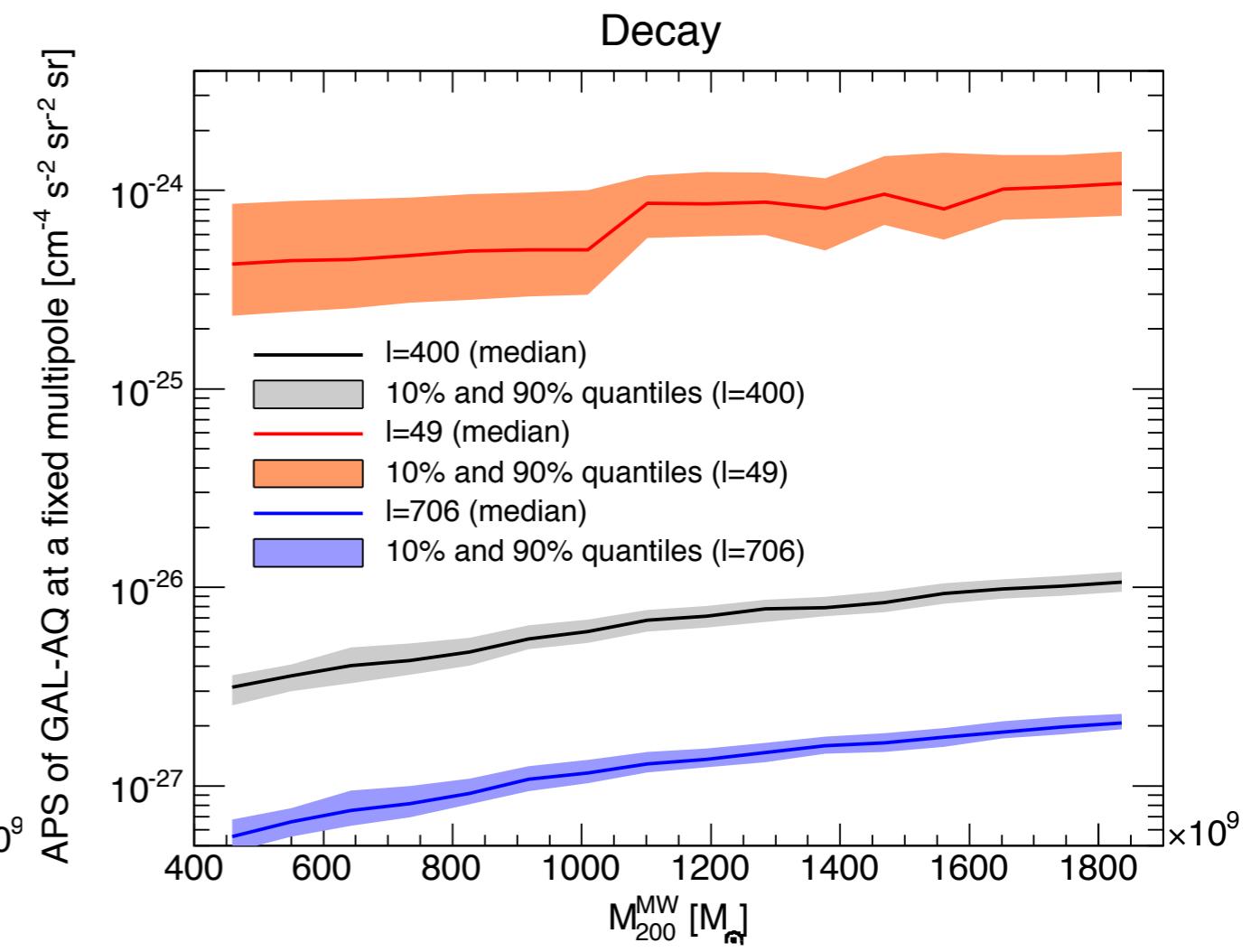
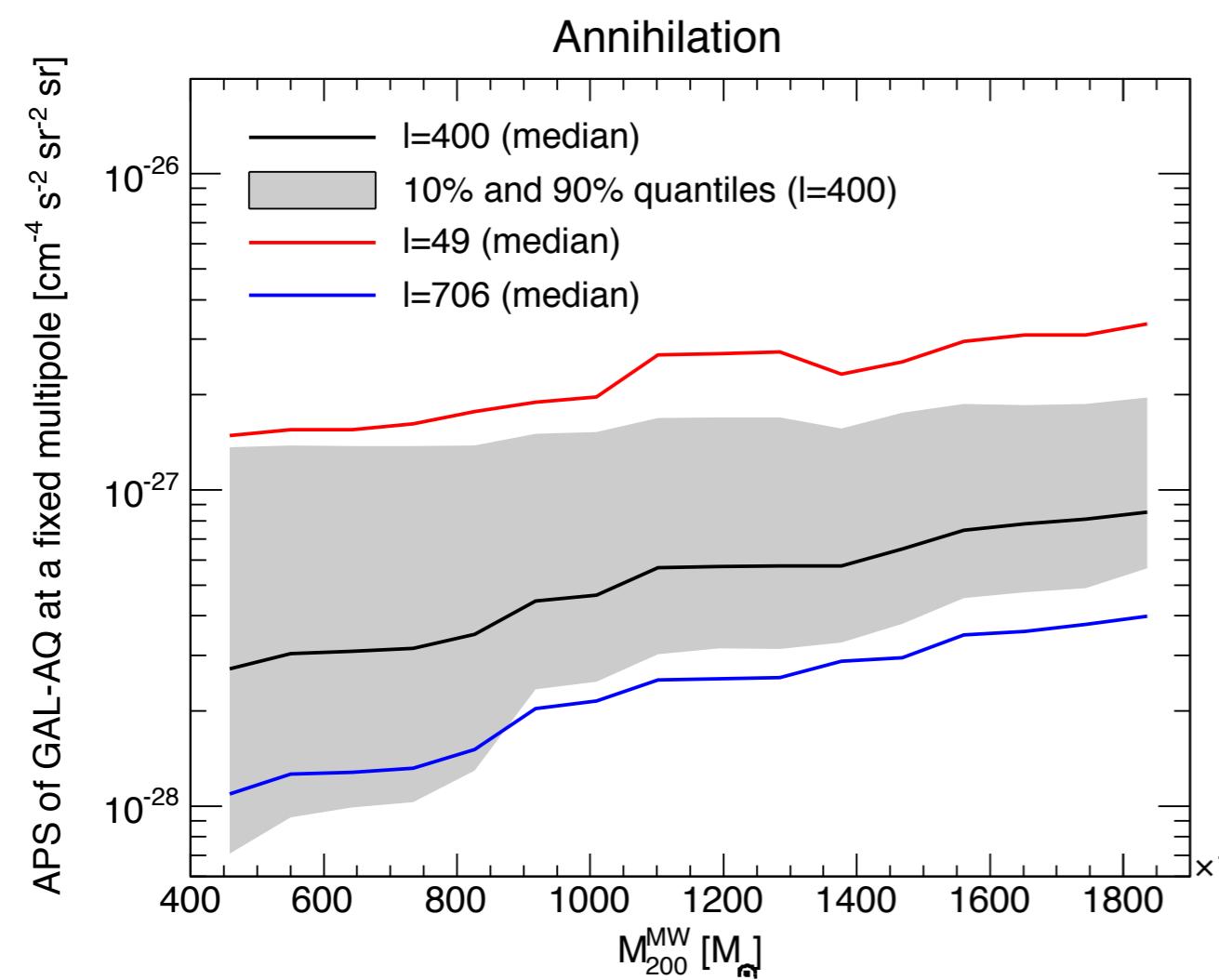
# DM-induced emission

- repetition of the Millennium-II simulation box to cover a large portion of the Universe
- extrapolation below the mass resolution of the Millennium-II (assuming low-mass halos trace the smallest halos in Millennium-II)
- unresolved subhalos accounted for through an analytic fit to  $P(\rho, r)$
- Milky Way smooth halo and Galactic subhalos from Aquarius (carved in the centre)



# Effect of an uncertain MW mass on GAL-AQ

- uncertainty of a factor 4 on the mass of the Milky Way (MW)
- 16 bins in  $M_{\text{MW}}$  accounting for a correspondent depletion in the amount of Galactic subhalos
- including uncertainty on the position of the observer

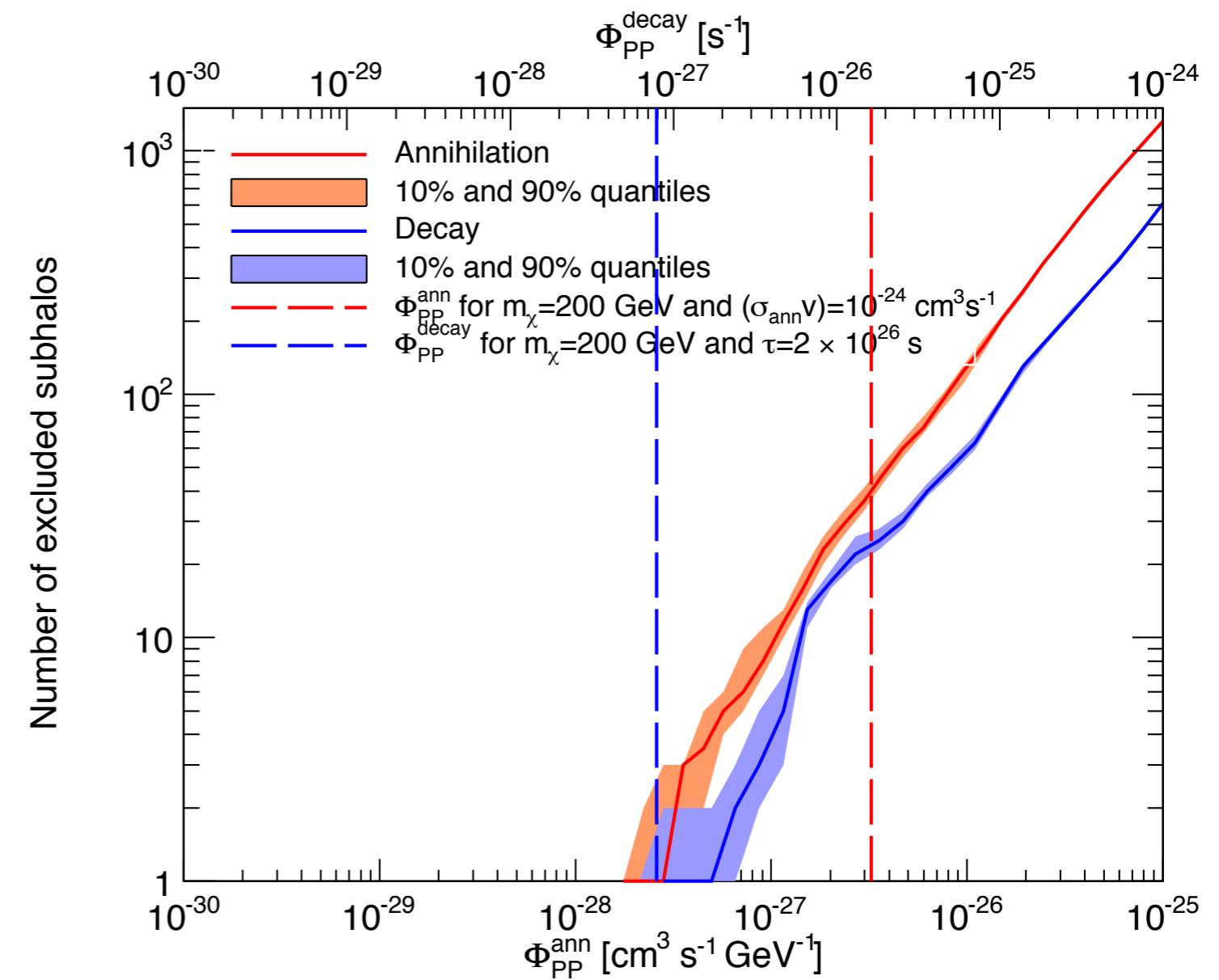


# Effect of an too-bright subhalos on GAL-AQ

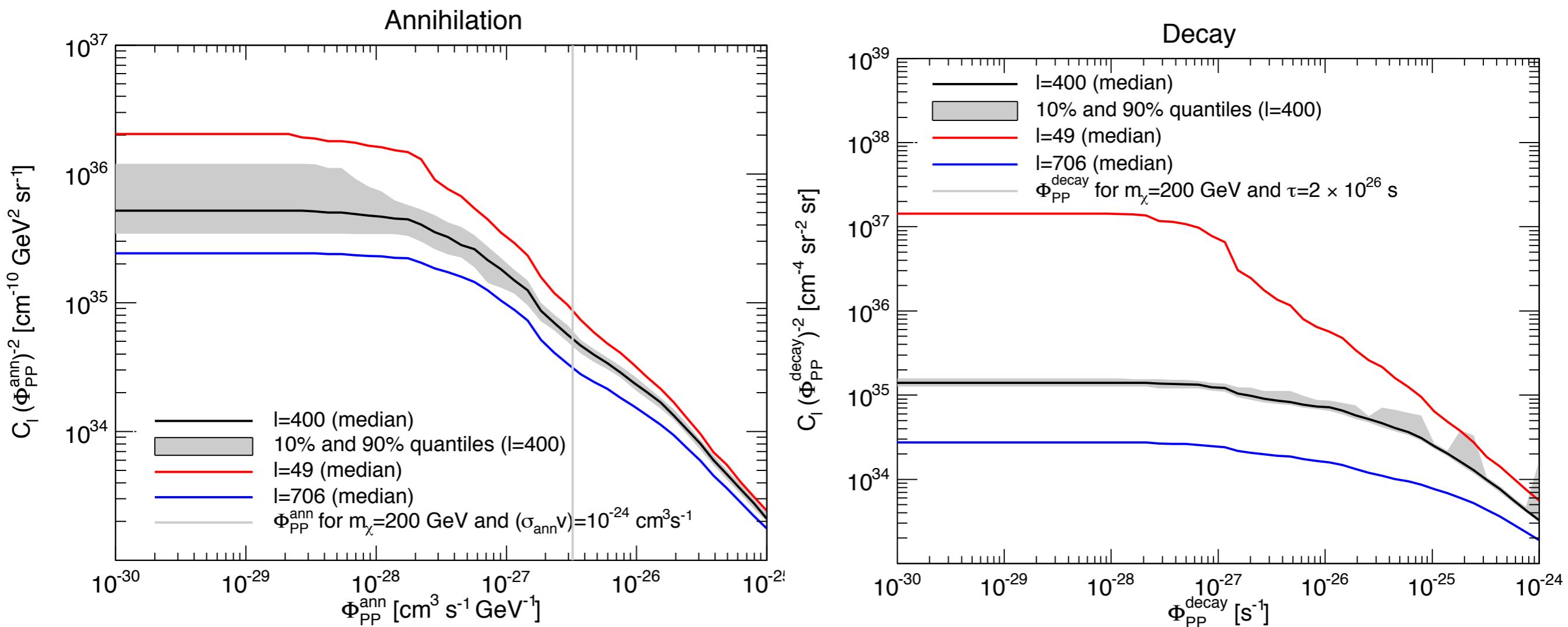
- for certain combination of  $(m_\chi, \sigma_{\text{ann}}v)$  and  $(m_\chi, \tau)$ , some subhalos are brighter than the 3FGL sensitivity
- those structures should be masked

$$\Phi_{\text{PP}}^{\text{ann}} = \frac{(\sigma_{\text{ann}}v)}{2m_\chi^2} \int_{\bar{E}} E \frac{dN_\gamma^{\text{ann}}}{dE} dE$$

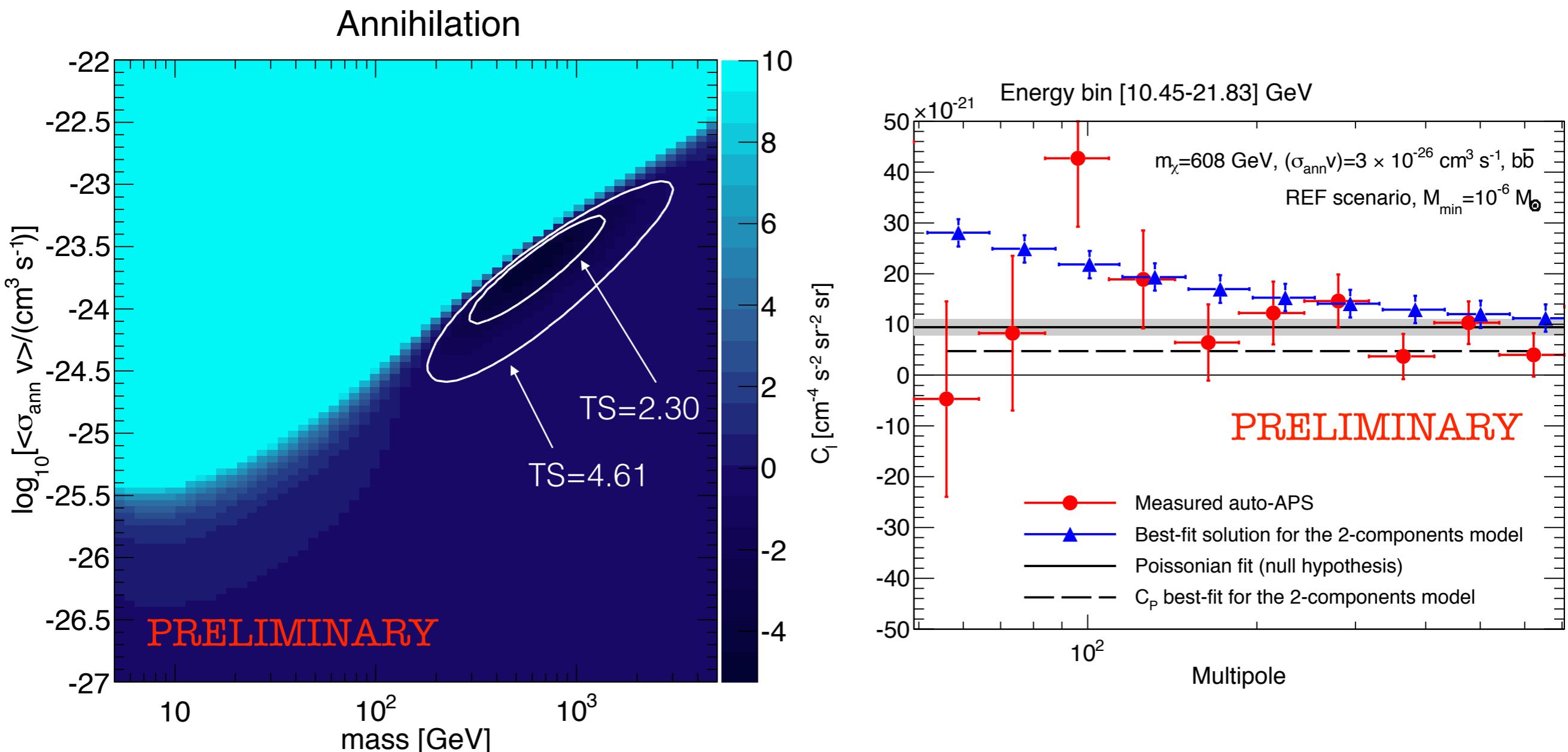
$$\Phi_{\text{PP}}^{\text{decay}} = \frac{1}{m_\chi \tau} \int_{\bar{E}} E \frac{dN_\gamma^{\text{decay}}}{dE} dE$$



# Effect of an too-bright subhalos on GAL-AQ

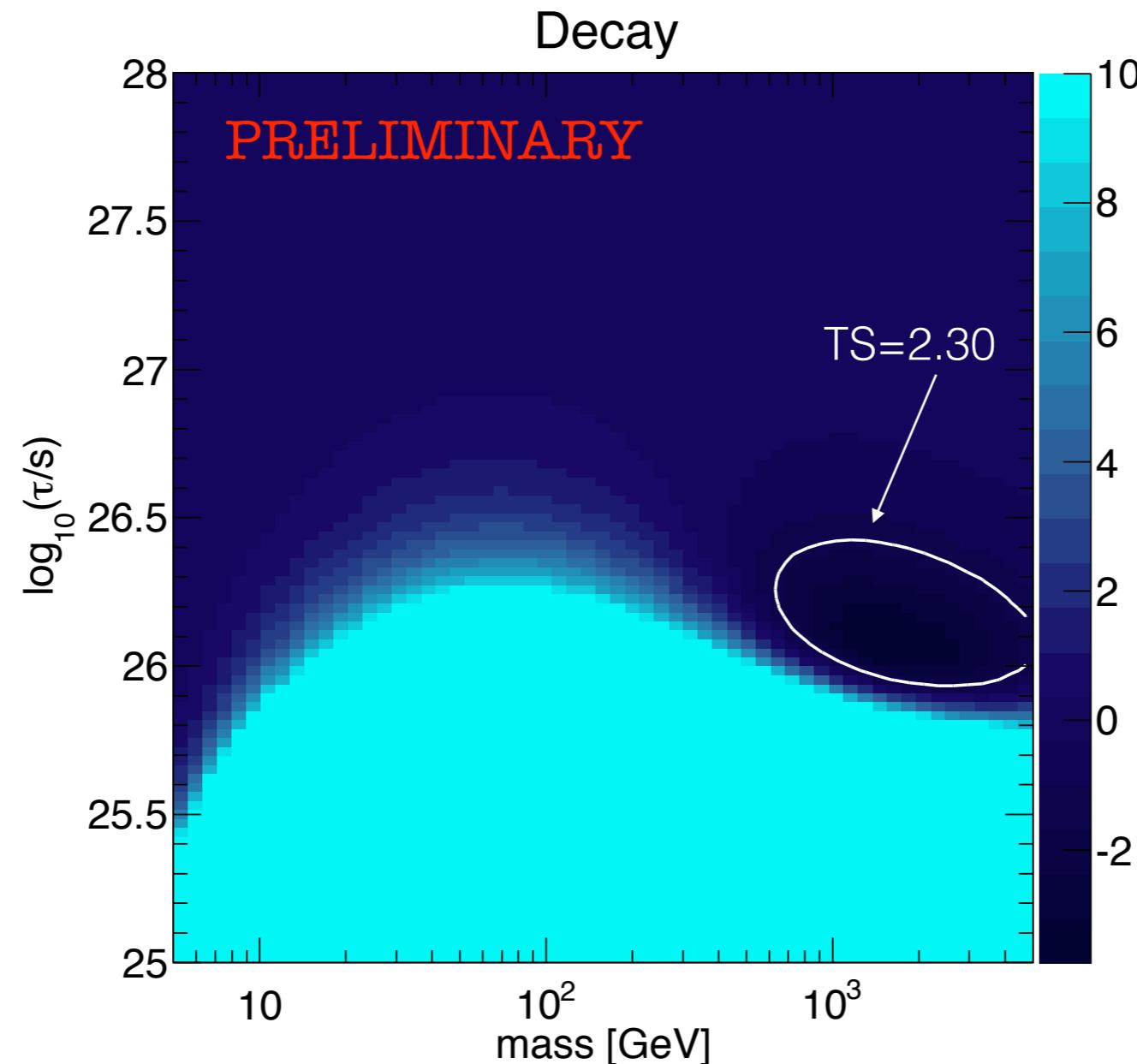


# 2-component fit to the binned APS



- $\text{TS} = -2 \ln[\chi^2(\text{no DM})] + 2 \ln[\chi^2(m_\chi, \sigma v)]$
- best-fit solution has  $\text{TS}=-4.5$ ,  $m_\chi=607 \text{ GeV}$ ,  $(\sigma_{\text{ann}} v)=2.2 \times 10^{-24} \text{ cm}^3 \text{s}^{-1}$

# 2-component fit to the binned APS



best-fit solution has  $m_\chi = 1743 \text{ GeV}$ ,  $\tau = 1.2 \times 10^{26} \text{ s}$