# **Muon System Performance**

FCC Hadron Detector Meeting Dec., 9<sup>th</sup> 2015

W. Riegler

## **Dates for Next Meetings**

Next hadron detector meetings:

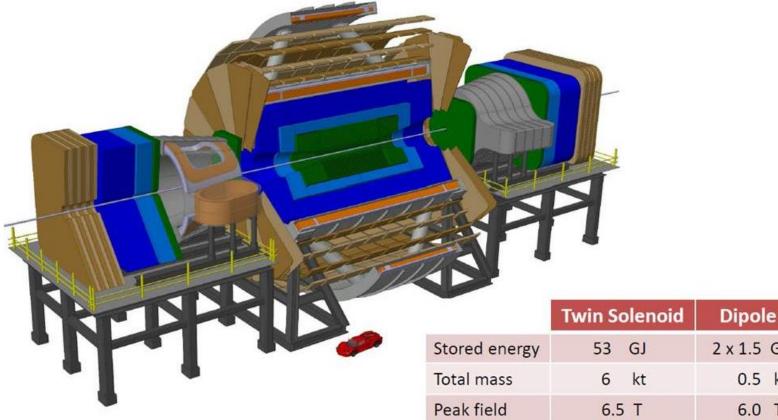
Jan. 21, 2016 Mar. 03, 2016 Apr. 06, 2016

FCC week Rome Apr. 11-15 2016

https://indico.cern.ch/category/6069/

## **Twin Solenoid + Dipole Magnet System**

Matthias Mentink, Alexey Dudarev, Helder Filipe Pais Da Silva, Christophe Paul Berriaud, Gabriella Rolando, Rosalinde Pots, Benoit Cure, Andrea Gaddi, Vyacheslav Klyukhin, Hubert Gerwig, Udo Wagner, and Herman ten Kate

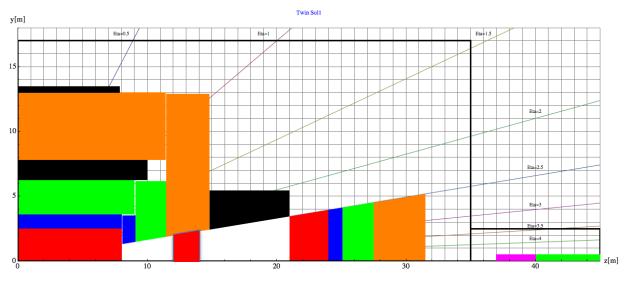


FCC Air core Twin solenoid and Dipoles

State of the art high stress / low mass design.

	Turn Solenoid	Dipole
Stored energy	53 GJ	2 x 1.5 GJ
Total mass	6 kt	0.5 kt
Peak field	6.5 T	6.0 T
Current	80 kA	20 kA
Conductor	102 km	2 x 37 km
Bore x Length	12 m x 20 m	6 m x 6 m

## **Baseline Geometry, Twin Solenoid**



## **Barrel:**

Tracker available space: R=2.1cm to R=2.5m, L=8m

EMCAL available space: R=2.5m to R=  $3.6m \rightarrow dR= 1.1m$ 

HCAL available space: R= 3.6m to R=6.0m → dR=2.4m

Coil+Cryostat: R= 6m to R= 7.825 → dR = 1.575m, L=10.1m

Muon available space: R= 7.825m to R= 13m  $\rightarrow$  dR = 5.175m

Coil2: R=13m to R=13.47m → dR=0.475m, L=7.6m

## Endcap:

EMCAL available space: z=8m to z=  $9.1m \rightarrow dz = 1.1m$ 

HCAL available space: z= 9.1m to z=11.5m  $\rightarrow$  dz=2.4m

Muon available space: z= 11.5m to z= 14.8m  $\rightarrow$  dz = 3.3m

## Forward:

Dipole: z= 14.8m to z= 21m  $\rightarrow$  dz=6.2m

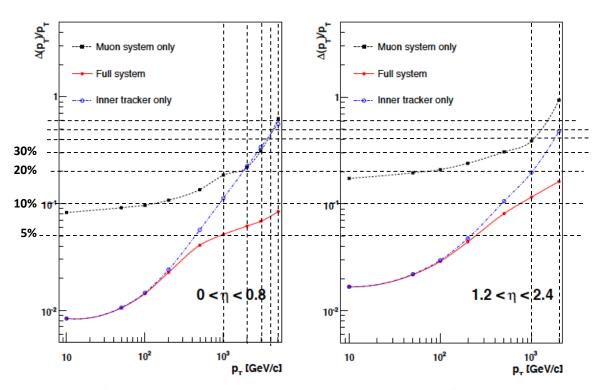
FTracker available space: z=21m to R=24m, L=3m

FEMCAL available space: Z=24m to z=  $25.1m \rightarrow dz= 1.1m$ 

FHCAL available space: z= 25.1m to z=27.5m  $\rightarrow$  dz=2.4m

FMuon available space: z= 27.5m to z=31.5m → dz=4m

#### **CMS Muon Performance**



e the first muon station, the offline bout 9% for small values of  $\eta$  and pdalone momentum resolution varies tum fit using also the inner tracker

**Figure 1.2**: The muon transverse-momentum resolution as a function of the transverse-momentum  $(p_T)$  using the muon system only, the inner tracking only, and both. Left panel:  $|\eta| < 0.8$ , right panel:  $1.2 < |\eta| < 2.4$ .

P <sub>T</sub> =1TeV/c, 0 <eta 0.8:<="" <="" th=""><th>20% muon standalone (angle)</th><th>P<sub>T</sub>=1TeV/c, eta 0<eta<2.4:< th=""><th>40% muon standalone (angle)</th></eta<2.4:<></th></eta>	20% muon standalone (angle)	P <sub>T</sub> =1TeV/c, eta 0 <eta<2.4:< th=""><th>40% muon standalone (angle)</th></eta<2.4:<>	40% muon standalone (angle)
	10% inner tracker only		20% inner tracker only
	5% combined		10% combined

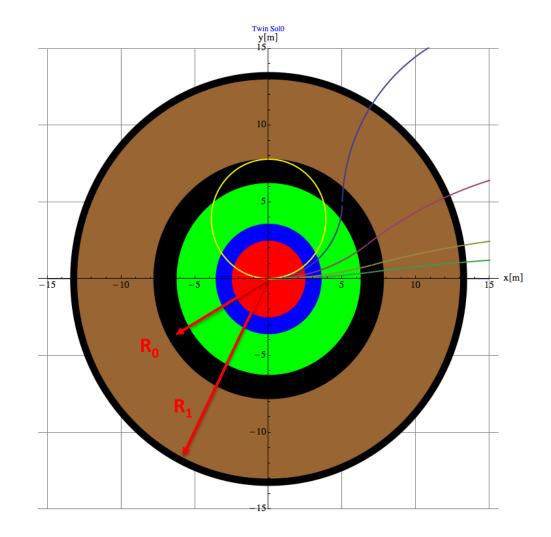
#### Muon Momentum can be measured by

1) The inner tracker

2) The track angle at the entrance of the muon system  $\rightarrow$  Trigger

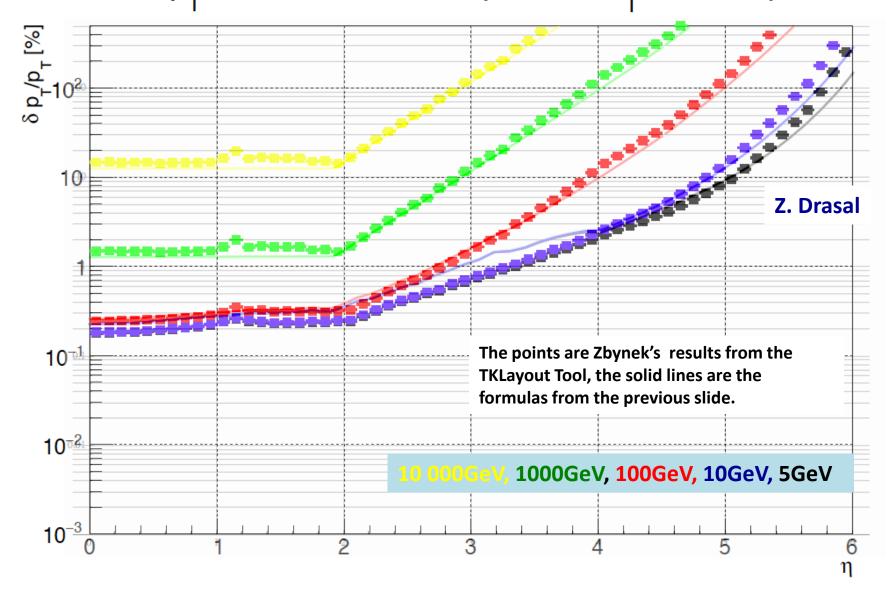
3) The combined fit of inner tracker and outer layers of the muon system.

4) A sagitta measurement in the muon system

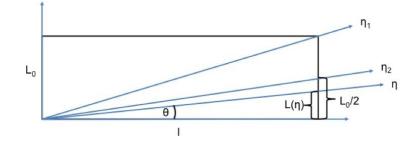


## Tracker

# $p_{\!_{\rm T}}$ resolution versus $\eta$ - const $P_{\!_{\rm T}}$ across $\eta$



Tracker



$$\eta_1 = -\ln an \left(rac{1}{2} \arctan rac{L_0}{l}
ight) \qquad \eta_2 = -\ln an \left(rac{1}{2} \arctan rac{L_0}{2l}
ight)$$

For a geometry with  $L_0 = 2.4m$  and l = 8m we have  $\eta_1 = 1.9$  and  $\eta_2 = 2.6$ 

$$L(\eta) = L_0 \quad \eta < \eta_1 \qquad \qquad L(\eta) = L_0 \frac{\sinh \eta_1}{\sinh \eta} \quad \eta > \eta_1$$

$$\frac{\Delta p_T}{p_T}|_{reso.} = \frac{\sigma \, p_T}{0.3BL(\eta)^2} \sqrt{\frac{720}{N(\eta)+4}} \qquad \qquad \frac{\Delta p_T}{p_T}|_{m.s.} = \frac{0.0136}{0.3BL(\eta)} \sqrt{\frac{x}{X_0}(\eta)}$$

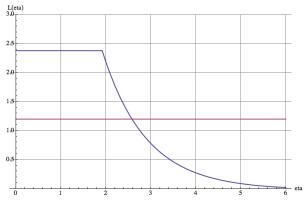
$$\frac{\Delta p_T}{p_T} = \sqrt{\left(\frac{\Delta p_T}{p_T}|_{reso.}\right)^2 + \left(\frac{\Delta p_T}{p_T}|_{m.s.}\right)^2}$$

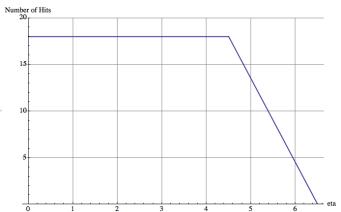
ln[60]:= L0 = 2.4;

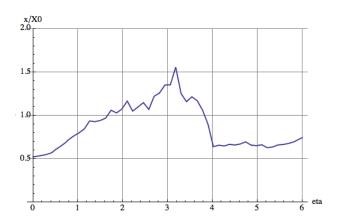
ln[61]:= 1 = 8;

In[62]:= B = 6;

 $\ln[63]:=$  sig = 23  $\pm$  10  $^{(-6)}$ ;







#### 2) Muon measurement by track angle at the entrance of the muon system

$$rac{\Delta p_T}{p_T} = \Delta heta \, \sqrt{\left(rac{2p_T}{0.3B_0R_0}
ight)^2 - 1} \quad pprox \quad rac{2p_T}{0.3B_0R_0} \, \Delta heta \quad ext{for a large } p_T$$

If we want 10% at 10TeV for B<sub>0</sub>=6T, R<sub>0</sub>=6m We need  $\Delta\theta$ =50µRad

→ 2 stations at 1.5m distance with 50um position resolution

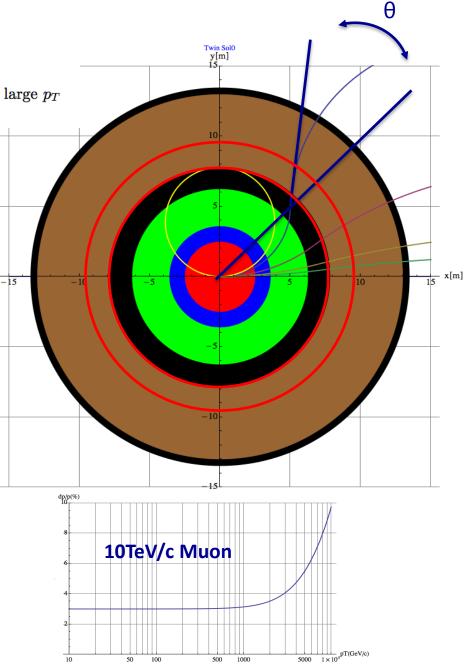
For low momentum, limit due to multiple scattering in the calorimeters and coil:

```
Calorimeter+Cryostat: 35X_0
HCAL: 110X_0
Coil: 5X_0
\rightarrow x_{tot}/X_0 \approx 150
```

 $\Delta \theta = 0.0136/p_{T}^{*}Sqrt(x_{tot}/X_{0})$ 

$\Delta p_T$	$2\times 0.0136$	$x_{tot}$	
$p_T$ –	$0.3B_0R_0$	$\sqrt{X_0}$	

 $B_0=6T$ ,  $R_0=6m → dp_T/p_T ≈ 3\%$ This is approximate for an infinitely thin coil and for η=0.



#### 3) Muon Measurement by combined measurement of inner tracker and muon system

If the full flux is returned trough the muon system, the muon trajectory at the exit of the system points exactly to the IP !

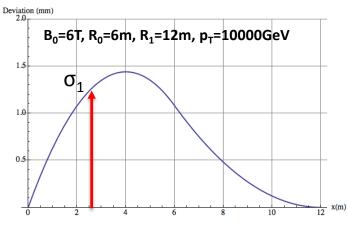
$$y_t(x) \;\;=\;\; rac{0.3B_0}{2p_T} \left( x^2 - rac{2R_0R_1}{R_0+R_1} x 
ight) \Theta(R_0-x) - rac{0.3B_0}{2p_T} rac{R_0^2}{(R_1-R_1)^2} \left( rac{R_0^2}{R_0^2} + r$$

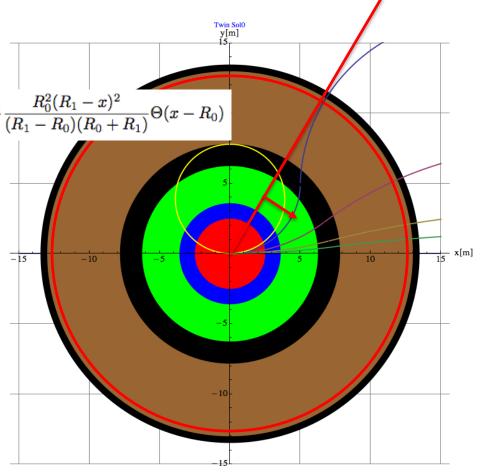
The maximum excursion  $y_t(x_0)$  is always at the same radial distance of  $x_0$ 

1

$$x_0 = rac{R_0 R_1}{R_0 + R_1} \qquad y_t(x_0) = -rac{0.3 B_0}{2 p_T} x_0^2 = -rac{0.3 B_0}{2 p_T} \left( rac{R_0 R_1}{R_0 + R_1} 
ight)^2$$

For values below:  $x_0=4m$ ,  $y_t(x0)=1.44mm$ Ideal measurement point is at the peak, but  $y_t(2.4m)=1.24mm$  still good !





 $\sigma^2 = \sigma_1^2 + (x/R_1\sigma_2)^2$ 

x=2.4m,R<sub>1</sub>=12m,  $\sigma_1$ =50 $\mu$ m,  $\sigma_1$ =250 $\mu$ m,  $\sigma$ =64 $\mu$ m, dp<sub>T</sub>/p<sub>T</sub>=5% at 10TeV !

Measuring just in the last tracker layer and in the outermost muon station already beats the full inner tracker performance (14 layers, 23um).

## 4) Muon measurement by sagitta measurement in the muon system

The return field is 2.45T

Measuring over the 5m lever arm with stations of sig=50um resolution we have

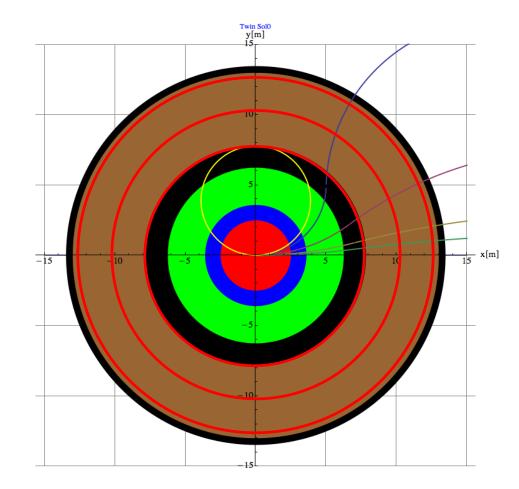
dp<sub>T</sub>/p<sub>T</sub>= sig\*p<sub>T</sub>/(0.3\*B\*L<sup>2</sup>)\*8 = 20% @ 10TeV

with possibly excellent performance at low  ${\bf p}_{\rm T}$  due to the absence of iron (vs. CMS) .

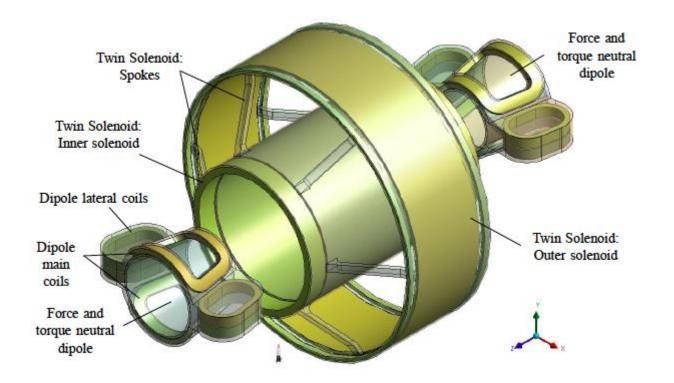
but very hard to beat the angular measurement at high  $p_T$  and the inner tracker at low  $p_T$ .

**Surface > 5000 m<sup>2</sup>** 

→ Seems not competitive



Muon performance versus η using the twin solenoid fieldmap and the baseline detector layout described in the previous meetings.



#### **Muon Trajectories in B-fields**

If we use the units  $[x] = [s] = m, [B] = T, [p] = GeV/c, q = e_0$  we have

$$\frac{d^2 \vec{x}(s)}{ds^2} = \frac{0.3}{p} \frac{d\vec{x}(s)}{ds} \times \vec{B}(\vec{x}(s))$$
(157)

or explicitly

$$x''(s) = \frac{0.3}{p} \left[ y'(s) B_z[x(s), y(s), z(s)] - z'(s) B_y[x(s), y(s), z(s)] \right]$$
(158)

$$y''(s) = \frac{0.3}{p} \left[ z'(s) B_x[x(s), y(s), z(s)] - x'(s) B_z[x(s), y(s), z(s)] \right]$$
(159)

$$z''(s) = \frac{0.3}{p} \left[ x'(s) B_y[x(s), y(s), z(s)] - y'(s) B_x[x(s), y(s), z(s)] \right]$$
(160)

The curvature is then given by

$$k = \frac{0.3}{p} \left| \frac{d\vec{x}(s)}{ds} \times \vec{B}(\vec{x}(s)) \right| = \frac{0.3}{p} B_T(\vec{x}(s))$$
(161)

Let's assume a straight line along a given  $\eta$  direction and a transverse B-field  $B_T(s)$  along this line. Assuming rotational symmetry of the magnetic field,  $B_T$  is inside one plane. A high momentum particle of momentum p starting from the origin i.e. y(0) = 0. The trajectory y(s) is defined by

$$k = \frac{y''(s)}{[1 + y'(s)^2]^{3/2}} = \frac{0.3}{p} B_T(s)$$
(162)

For small deflections,  $y'(s) \ll 1$  and the equation becomes

$$y''(s) = \frac{0.3}{p} B_T(s) \tag{163}$$

Since it holds that

$$\oint \vec{B}(\vec{x})d\vec{A} = 0 \tag{164}$$

and in case the magnetic field is zero outside the shielding coil and at large values of z, it holds with  $s=z/\cos\theta$ 

$$\int_0^a s B_T(s) ds = 0 \tag{165}$$

With this consition we have

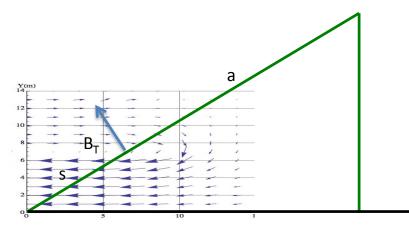
$$\int_{0}^{a} sy''(s)ds = 0 (166)$$

$$sy'(s)|_0^a - \int_0^a y'(s)ds = 0$$
(167)

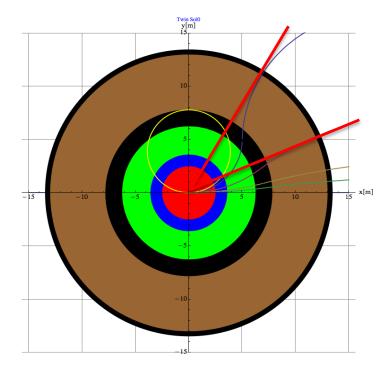
$$ay'(a) - y(a) + y(0) = 0 (168)$$

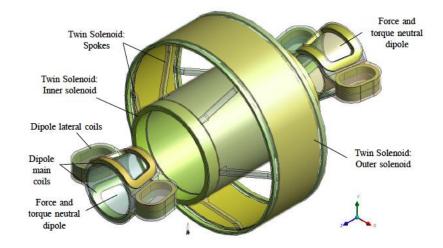
$$y'(a) = \frac{y(a)}{a} \tag{169}$$

so at the point x = a the particle trajectory points back to the origin.



#### Flux Return and Muon Trajectories





In[156]:= tab = Import["/Users/riegler/Documents/X/FCC/Bfield/Bfield\_may\_2015.txt", "table"];

In[157]:= (\*tab=Import["/Users/riegler/Documents/X/FCC/Bfield/Bfield\_test.txt","table"];\*)

ln[158] = n = Length[tab]

Out[158]= 624777

 $\ln[1756]:= eql = x''[s] = 0.3/p \star (y'[s] \star Bz[x[s], y[s], z[s]] - z'[s] \star By[x[s], y[s], z[s]]);$ 

 $\ln[1757]:= eq2 = y''[s] = 0.3/p \star (z'[s] \star Bx[x[s], y[s], z[s]] - x'[s] \star Bz[x[s], y[s], z[s]]);$ 

 $\ln[1758] := \mathbf{eq3} = \mathbf{z}''[\mathbf{s}] = \mathbf{0.3} / \mathbf{p} \star (\mathbf{x}'[\mathbf{s}] \star \mathbf{By}[\mathbf{x}[\mathbf{s}], \mathbf{y}[\mathbf{s}], \mathbf{z}[\mathbf{s}]] - \mathbf{y}'[\mathbf{s}] \star \mathbf{Bx}[\mathbf{x}[\mathbf{s}], \mathbf{y}[\mathbf{s}], \mathbf{z}[\mathbf{s}]]);$ 

 $ln[1759] \coloneqq (* \ X[s], \ Y[s], \ Z[s] \ --> \ parametrization \ of \ the \ muon \ trajectory \ *)$ 

ln[1760]:= {X, Y, Z} = NDSolveValue[{eq1, eq2, eq3, x'[0] == Sin[theta0] \* Cos[phi], x[0] == 0, y'[0] == Sin[theta0] \* Sin[phi], y[0] == 0, z'[0] == Cos[theta0], z[0] == 0},
{x, y, z}, {s, sd[5, theta0]}]

 $Out[1760]= \{InterpolatingFunction[\{\{0., 26.4379\}\}, <>\}, InterpolatingFunction[\{\{0., 26.4379\}\}, <>\}\}$ 

We assume a high momentum particle originating from the IP in a given  $\theta$ ,  $\phi$  direction. The particle will deviate only by a small amount from a straight line along this direction and the deviation is given by the transverse magnetic field  $\vec{B_T}(s)$  along this direction. We have  $g''(s) = 0.3/pB_T(s) = kB_T(s)$ . We let the particle start from the origin and therefore have g(0) = 0 and also assume that this point is measured with infinite precision. Since we don't know the original direction of the particle we have to allow a small rotation and we have a template function f(s) given by

$$f(s) = k g(s) + \theta s \tag{170}$$

(172)

(173)

(174)

(175)

(176)

and we have to find the parameters k and  $\theta$  that minimize the difference to the measurement points

$$\sum_{n=1}^{N} w_n \left[ y_n - f(s_n) \right]^2 \quad \to \min$$
(171)

Defining  $g_n = g(s_n)$  we have

$$F(k,\theta) = \sum_{n=1}^{N} w_n \left[ y_n - kg_n - \theta s_n \right]^2 \qquad \frac{\partial F}{\partial k} = 0 \qquad \frac{\partial F}{\partial \theta} = 0$$

$$k\,F_{gg}+\theta\,F_{gs}=F_{gy}\qquad k\,F_{gs}+\theta\,F_{ss}=F_{sy}$$

$$F_{gg} = \sum_{n=1}^{N} w_n g_n^2$$
  $F_{gs} = \sum_{n=1}^{N} w_n s_n g_n$   $F_{ss} = \sum_{n=1}^{N} w_n s_n^2$ 

$$F_{sy}=\sum_{n=1}^N w_n s_n y_n \qquad F_{gy}=\sum_{n=1}^N w_n g_n y_n$$

$$k = \frac{F_{gs}F_{sy} - F_{gy}F_{ss}}{F_{gs}^2 - F_{gg}F_{ss}} \qquad \theta = \frac{F_{gs}F_{gy} - F_{gg}F_{sy}}{F_{gs}^2 - F_{gg}F_{ss}}$$

$$\Delta_k^2 = \sum_{n=1}^N \left(\frac{\partial k}{\partial y_n}\right)^2 \sigma_n^2 \tag{177}$$

$$\frac{\partial k}{\partial y_n} = \frac{w_n [F_{gs} s_n - F_{ss} g_n]}{F_{gs}^2 - F_{gg} F_{ss}}$$
(178)

and we have the general solution

$$\Delta_k^2 = \frac{F_{ss}^2 H_{gg} - 2F_{ss}F_{gs}H_{gs} + F_{sg}^2 H_{ss}}{(F_{gg}F_{ss} - F_{sg})^2} \tag{179}$$

with

$$H_{gg} = \sum_{n=1}^{N} \sigma_n^2 w_n^2 g_n^2 \qquad H_{gs} = \sum_{n=1}^{N} \sigma_n^2 w_n^2 g_n s_n \qquad H_{ss} = \sum_{n=1}^{N} \sigma_n^2 w_n^2 s_n^2$$
(180)

Knowing the position of the detection layers  $s_n$ , the resolution of these layers  $\sigma_n$  and the value of the template function at the detection layer  $g_n = g(s_n)$ , the variance  $\Delta_k^2$  is only a function of the relative weights of the measurement points and has to be minimized with respect to these weights.

$$\Delta_k(w_2, w_3, ..., w_N) \to \min \tag{181}$$

For the momentum resolution we then have

$$k = \frac{0.3}{p} \longrightarrow \frac{\sigma_p}{p} = \frac{0.3}{p} \Delta_k$$
 (182)

Analytic, back of the envelope, pseudo-precise determination of the muon momentum resolution using the tracker+muon system measurement.

Template function g(s) is calculated with MATHEMATICA based on the correct field map.

14 equally spaced layers of central tracker are assumed with a resolution of 23um per layer.

The layers in the muon system are then added with their respective resolution.

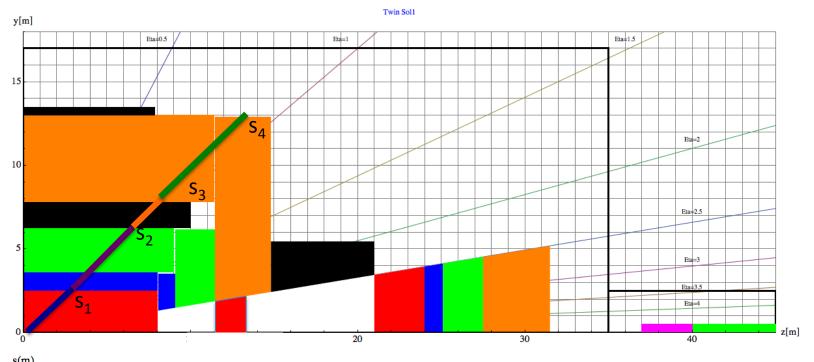
A global minimization of  $\Delta_k$  for the weights  $w_n$  is performed.

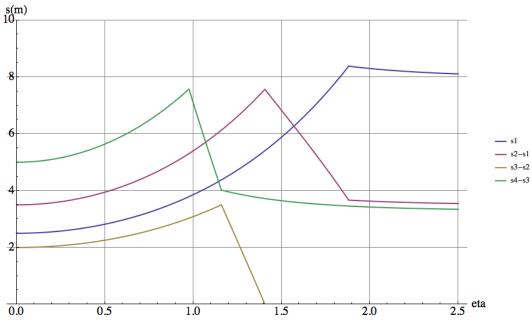
Before assuming specific resolution number for the muon system, a key question is: How is the position measurement in the muon system affected by the multiple scattering the in calorimeters ?

For N = 2 we have

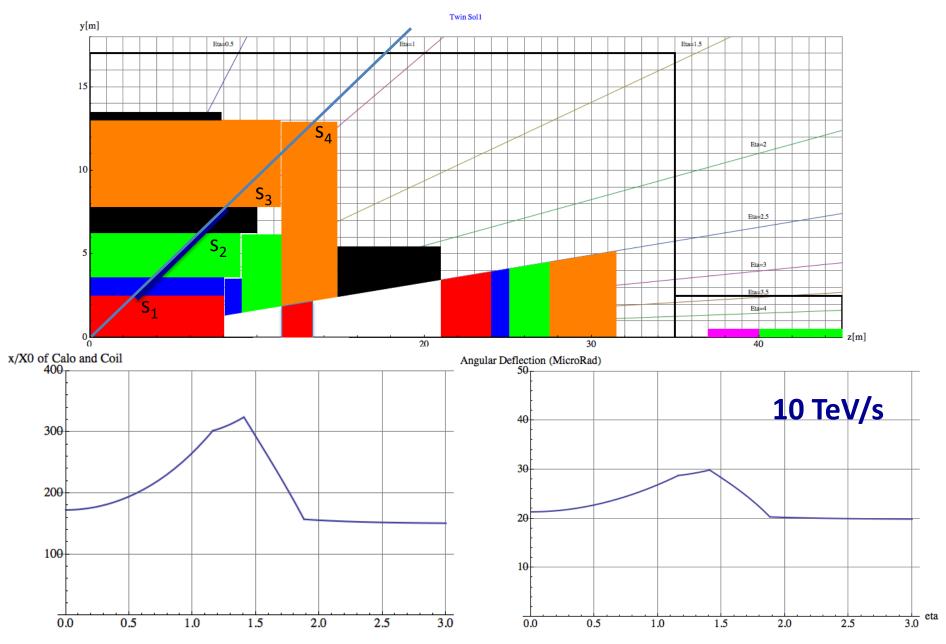
$$\Delta_k^2 = \frac{s_1^2 \sigma_2^2 + s_2^2 \sigma_1^2}{[s_1 g_2 - s_2 g_1]^2} \tag{183}$$

## Thickness of different subsystems seen by a high energy muon

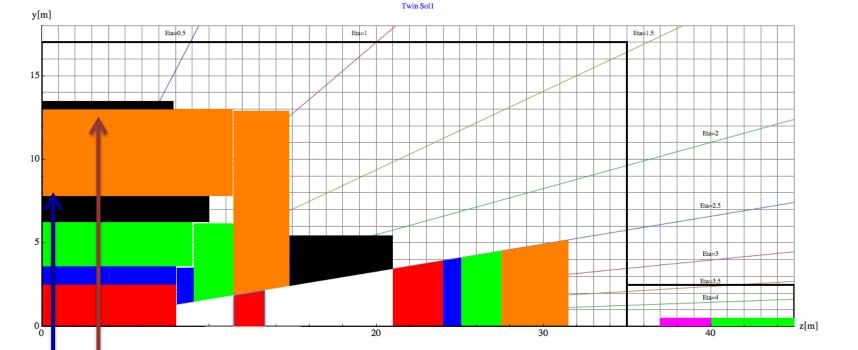




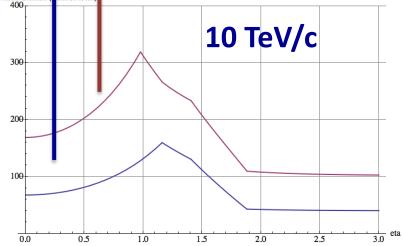
## Radiation length and angular deflection of muons in calo+coil



#### Measurement error in the muon system due to multiple scatering



MS Error at Muon Station ( ficroMeter)



The r.m.s. (in plane) angular deflection  $\theta_0$  and spatial displacement  $\sigma_y$  of a particle passing through material of thickness x are given by

$$\theta_0 = \frac{0.0136}{\beta p [GeV/c]} \sqrt{\frac{x}{X_0}} \left( 1 + 0.038 \log \frac{x}{X_0} \right)$$
(184)

$$\sigma_y = \frac{1}{\sqrt{3}} x \,\theta_0 \tag{185}$$

For the proper correlated M.C. generation of these two quantities, one uses two independent Gaussian random variables  $z_1, z_2$  with mean of zero and sigma of one, and calculates

$$y = z_1 x \frac{\theta_0}{\sqrt{12}} + z_2 x \frac{\theta_0}{2}$$
(186)

$$\theta = z_2 \theta_0 \tag{187}$$

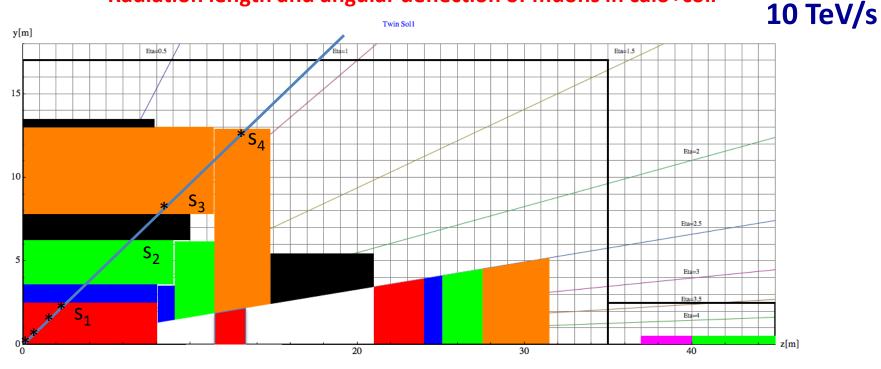
In case a muon passes the calorimeter of thickness  $L_1$  and then propagates through the (material less) muons system of length  $L_2$  the spatial deflection is given by

$$s = z_1 L_1 \frac{\theta_0}{\sqrt{12}} + z_2 L_1 \frac{\theta_0}{2} + z_2 \theta_0 L_2 \tag{188}$$

and therefore a position error of

$$\sigma_s^2 = \theta_0^2 \left[ \frac{L_1^2}{12} + \left( \frac{L_1}{2} + L_2 \right)^2 \right]$$
(189)

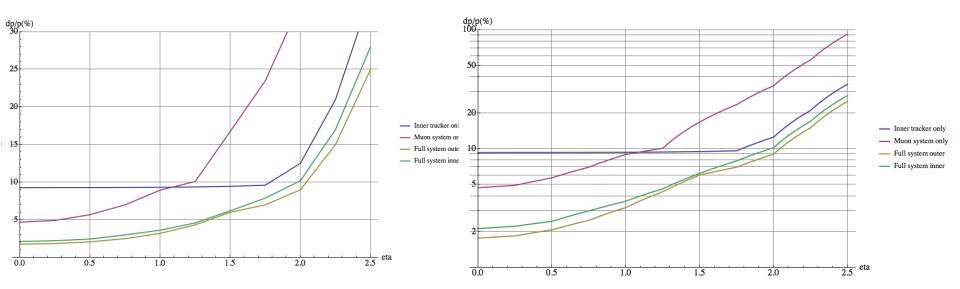
## Radiation length and angular deflection of muons in calo+coil



Fit of the template function to 14 measurement points of 23um in the inner tracker and 1 layer in the muon system with a resolution that is much better than the multiple scattering limit.

i.e. assuming the resolution at s4 to be better than 100um and at s3 to be better than 50um and the resolution of the angular measurement better than 20uRad.

#### Momentum resolution for a 10 TeV/s muon vs. eta



Twin Solenoid assuminginner tracker with baseline resolution curves and multiple scattering limit in the muons system.

 $P_{T}=10TeV/c$  eta = 0:

5% muon standalone (angle) 10% inner tracker only 2% combined

 $P_{\tau}$ =10TeV/c eta=2.: 35% muon standalone (angle) 12.5% inner tracker only 8% combined

#### **Compare to the CMS numbers:**

P <sub>T</sub> =1TeV/c, 0 <eta 0.8:<="" <="" th=""><th>20% muon standalone (angle)</th><th>P<sub>T</sub>=1TeV/c, eta 0<eta<2.4:< th=""><th>40% muon standalone (angle)</th></eta<2.4:<></th></eta>	20% muon standalone (angle)	P <sub>T</sub> =1TeV/c, eta 0 <eta<2.4:< th=""><th>40% muon standalone (angle)</th></eta<2.4:<>	40% muon standalone (angle)
	10% inner tracker only		20% inner tracker only
	5% combined		10% combined

# **Summary**

First shot at the global muon performance.

What CMS does for 1TeV/c Muons seems feasible for 10TeV/c at FCC.

For now, the 'global fit' assumed just one measurement point in the muon system in addition to the 14 points of the inner tracker.

We have to do proper evaluation of the radiation length of the Calo using FLUKA/GEANT.

We have to understand whether the scattering of high energy muons is really approximated well by the given formulas (see todays presentation on Muon scattering in GEANT).

We have to include energy loss.

Another attempt to include all this in the semi-analytical calculation will be done.

A proper track fit in a proper MC simulation will be needed soon to have a solid result.

