

From cold Fermi fluids to the hot QGP

Toward better understanding strongly coupled quantum fluids

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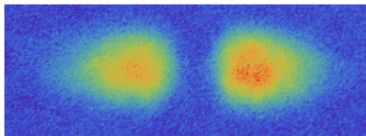
Critical Point and Onset of Deconfinement 2016 – Wrocław; June 2, 2016

Prelude

- ▶ ultracold Fermi gases share features with other strongly coupled quantum systems
- ▶ unique experimental playground to study fermionic many-body problem: widely variable densities, temperatures and interatomic interaction strengths
- ▶ Feshbach-resonant scattering → investigate crossover between BCS-type superfluid and BEC of strongly bound molecules
- ▶ ^{40}K , ^6Li : fermionic atoms with electron spin $S = 1/2$, nuclear spin $I = 1$
- ▶ cooled to μK -regime: s -wave scattering (length a) dominant

Prelude

Fermi gas cloud collision:



evolution within first 12 ms

- ▶ repulsive optical potential divides trapped gas into two clouds
 - ▶ clouds collide once repulsive potential switched off
- ⇒ formation of propagating shock fronts (fluid dynamical behavior)

Prelude

Fermi gas cloud collision:

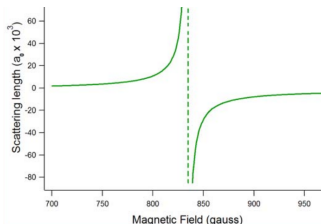
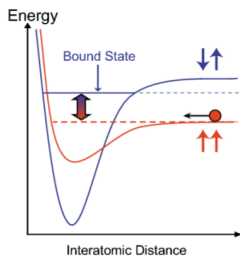
evolution within first 12 ms

- ▶ repulsive optical potential divides trapped gas into two clouds
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⇒ formation of propagating shock fronts (fluid dynamical behavior)

Tuning the interaction strength in Fermi gases

consider spin-balanced system of two (lowest) macroscopically occupied hyperfine eigenstates $|m_f = \pm \frac{1}{2}\rangle$ of ${}^6\text{Li}$



- ▶ scattering between two atoms predominantly in **triplet state** (Open Channel) of interaction potential
- ▶ at low incident energy: **singlet (bound) state** (Closed Channel) with a different magnetic moment

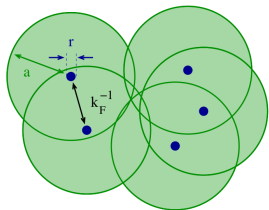
tune B -field \rightarrow CC couples resonantly with OC scattering continuum

\Rightarrow temporary capture in quasi-bound states: Fano-Feshbach resonance

Unitary Fermi gas

unitary limit $|a| \rightarrow \infty$:

- ▶ scale invariance, conformal invariance
- ▶ universal (*s*-wave) collision cross-section: $\sigma_{\text{coll}} = 4\pi/q^2$
- ▶ properties of the gas are universal functions of n and T ,
e.g. $\eta = \hbar n \cdot \alpha(z)$ with fugacity $z \sim n\lambda^3 \sim n/(mT)^{3/2}$



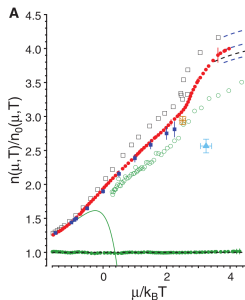
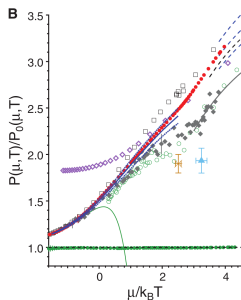
→ dilute regime: $r \cdot n^{1/3} \ll 1$

- ▶ extremely strong coupling/correlations reachable if a exceeds average interparticle spacing, i.e. $a \cdot n^{1/3} \gg 1$
- ▶ in high- T limit: $z \ll 1$, $\sigma_{\text{coll}} \sim 1/T$
→ kinetic theory applies

Equilibrium properties of unitary Fermi gases

universal thermodynamics:

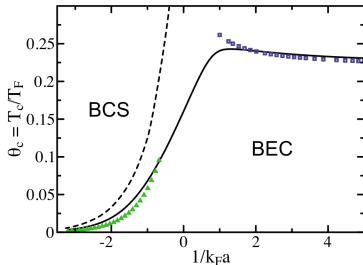
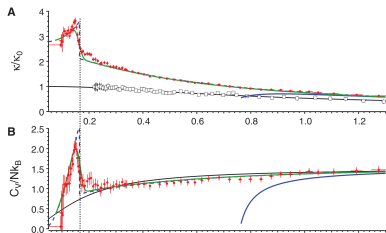
- ▶ $P = \frac{2}{3} \mathcal{E}_0$
- ▶ $P = T \lambda^{-3} f_P(z)$, $n = \lambda^{-3} f_n(z)$ with universal functions f_P and f_n



measurement across
superfluid transition

- ▶ change in behavior of $n(T, \mu)$ around $\mu/T \gtrsim 2.5$

Equilibrium properties of unitary Fermi gases



- ▶ rise in compressibility κ
 - ▶ λ -shape structure in specific heat c_V as functions of T/T_F
- \Rightarrow 2nd-order phase transition with $T_C/T_F = 0.167(13)$

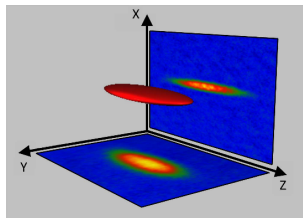
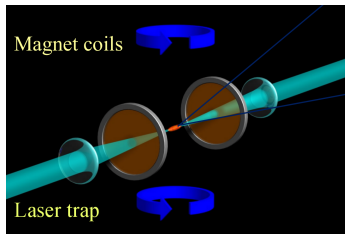
- ▶ self-consistent T -matrix, FRG and QMC approaches:
 $\rightarrow T_C/T_F \simeq 0.16$

Transport properties of unitary Fermi gases

1) How can they be measured?

- ▶ analysis of the gas cloud *flow dynamics* after release from a deformed trap or of the *damping* of the transverse breathing mode after recapturing the gas

→ **elliptic flow measurement:**



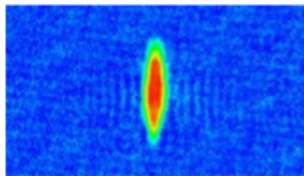
study expansion after release

Transport properties of unitary Fermi gases

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→ **elliptic flow measurement:**



100 μs ... 2 ms

- ▶ shows fluid dynamical behavior
- ▶ analyze time-evolution of gas cloud geometry

K.M. O'Hara et al., *Science* **298** (2002) 2179; C. Cao et al., *Science* **331** (2011) 58; J. Kinast et al., *PRA* **70** (2004) 051401 (R); M. Bartenstein et al., *PRL* **92** (2004) 203201; A. Turlapov et al., *J. Low Temp. Phys.* **150** (2008) 567; C. Cao et al., *New J. Phys.* **13** (2011) 075007; E. Elliott et al., *PRL* **112** (2014) 040405

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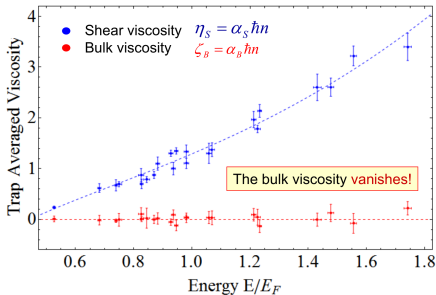
- ▶ shows fluid dynamical behavior
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100 μ s ... 2 ms

Transport properties of unitary Fermi gases

2) Experimental results:

- ▶ dependence on initial energy per particle E



in normal phase

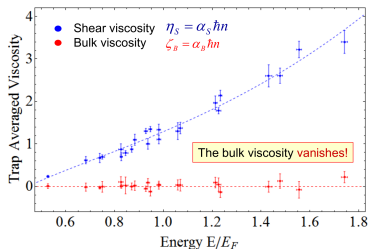
$$\rightarrow \langle \eta \rangle / \langle s \rangle \geq 0.4$$

\Rightarrow nearly perfect fluidity

\rightarrow only cloud-averaged quantities determined: $\alpha_S \sim \int \eta dV$

Transport properties of unitary Fermi gases

3) Theory results:

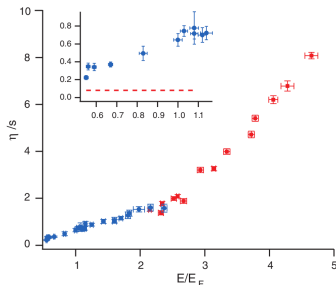


bulk viscosity $\zeta = 0$

D.T. Son, PRL **98** (2007) 020604; P. Massignan et al., PRA **71** (2005) 033607; G.M. Bruun, H. Smith, PRA **72** (2005) 043605; *ibid.* PRA **76** (2007) 045602; G. Rupak, T. Schäfer, PRA **76** (2007) 053607; T. Enss et al., Ann. Phys. **326** (2011) 770; H. Guo et al., New J. Phys. **13** (2011) 075011; *ibid.* PRL **107** (2011) 020403; G. Wlazłowski et al., PRL **109** (2012) 020406; *ibid.* PRA **88** (2013) 013639; Y. He, K. Levin, PRB **89** (2014) 035106; T. Schäfer, PRA **90** (2014) 043633

Transport properties of unitary Fermi gases

3) Theory results:



shear viscosity η

- ▶ in high- T limit

$$\rightarrow \eta = \frac{15}{32\sqrt{\pi}} (mT)^{3/2}$$

- ▶ at $z \simeq 12$ superfluid transition

$$\rightarrow \eta \sim 1/T^5 \text{ for phonons}$$

\Rightarrow suggests minimum near T_c

- ▶ for $z \gtrsim 1$

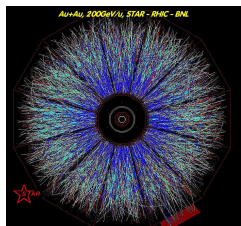
$$\rightarrow T\text{-matrix: } \eta/s \simeq 0.5$$

$$\rightarrow \text{QMC: } \eta/s \simeq 0.2$$

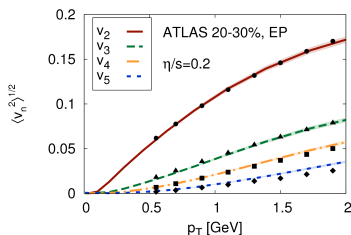
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Transport properties of the QGP

information on QGP transport properties come from observables associated with fluid dynamical flow in HIC



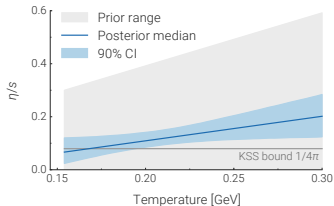
- ▶ in non-central HIC find anisotropies in azimuthal particle distribution



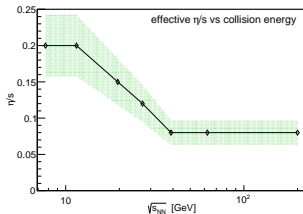
- ▶ analysis of flow harmonics:
 $\eta/s \simeq 0.12 \dots 0.24$
 \Rightarrow **nearly perfect fluidity**

Local determination of η in the QGP

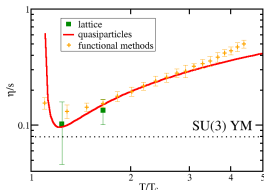
recent efforts aim at determining dependence of η/s on T and n



see talk by S. Bass



see talk by Iu. Karpenko



see talk by J.M. Pawłowski

MB, B. Kämpfer, K. Redlich, PRC **84** (2011) 025201

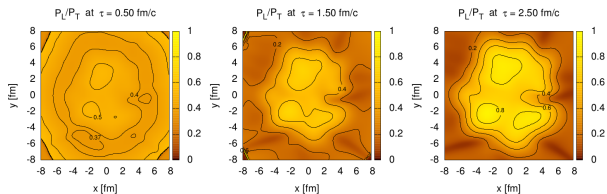
H. Niemi et al., PRC **86** (2012) 014909; Iu.A. Karpenko et al., PRC **91** (2015) 064901; J.E. Bernhard et al., 1605.03954; MB et al., PLB **709** (2012) 77; N. Christiansen et al., PRL **115** (2015) 112002; H. Berrehrah et al., Int. J. Mod. Phys. E **25** (2016) 1642003

Recent development in fluid dynamics for HICs

at early times large momentum-space anisotropies within fireball from pre-equilibrium evolution

- ▶ viscous hydro and AdS/CFT estimates give $P_L/P_T \lesssim 0.3$:
→ anisotropy stronger for larger η/s and lower T

⇒ development of (viscous) anisotropic hydrodynamics for relativistic fluids



see talk by
W. Florkowski

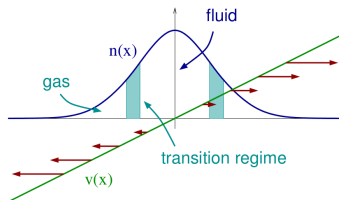
- ▶ reproduces known exact solutions of Boltzmann equation accurately

M.P. Heller et al., PRL **108** (2012) 201602; W. van der Schee et al., PRL **111** (2013) 222302; M. Martinez, M. Strickland, NPA **848** (2010) 183; W. Florkowski, R. Ryblewski, PRC **83** (2011) 034907; M. Martinez et al., PRC **85** (2012) 064913; W. Florkowski et al., NPA **916** (2013) 249; *ibid.* PRC **88** (2013) 024903; D. Bazow et al., PRC **90** (2014) 054910; L. Tinti, W. Florkowski, PRC **89** (2014) 034907; W. Florkowski et al., PRC **89** (2014) 054908; *ibid.* PRC **89** (2014) 054909; G.S. Denicol et al., PRL **113** (2014) 202301; *ibid.* PRD **90** (2014) 125026; M. Nopoush et al., PRD **91** (2015) 045007; W. Florkowski et al., PRC **92** (2015) 054912; M. Alqahtani et al., PRC **92** (2015) 054910; *ibid.* 1605.02101; E. Molnár et al., 1602.00573

Local η determination in unitary Fermi gases

Navier-Stokes theory breaks down in dilute corona \rightarrow paradoxical fluid dynamical behavior

- ▶ must either assume $\eta \sim n$ or introduce cut-off radius

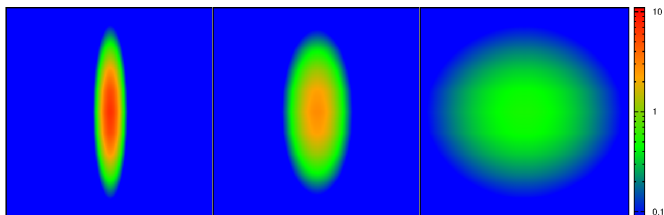


- ▶ reliable treatment of low density regime:
solve Boltzmann equation numerically; Lattice Boltzmann simulation; Quantum Monte Carlo simulation; combine fluid dynamics with kinetic transport approach; **Anisotropic fluid dynamics**

Anisotropic fluid dynamics

based on guided choice for f_p^{an} (exact solution of Boltzmann equation for harmonic potentials in ballistic limit)

- ▶ conservation laws + evolution equations for *non-hydrodynamic* d.o.f. \mathcal{E}_a with $\sum_a \mathcal{E}_a = \mathcal{E}$: depend on **relaxation time** $\tau = \eta / P$

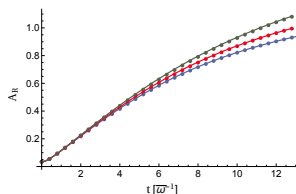


MB, T. Schäfer, PRA **92** (2015) 043602

- ▶ small τ : fast relaxation to Navier-Stokes theory
 - ▶ large τ : additional conservation laws \rightarrow ballistic expansion
- \Rightarrow **provides graceful exit for fluid dynamics in dilute regime**

Application of anisotropic fluid dynamics

baseline test:



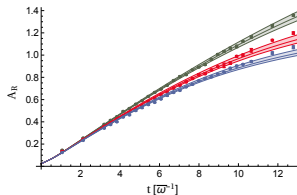
T/T_F

0.79

1.11

1.54

expansion data analysis:



MB, T. Schäfer, PRL **116** (2016) 115301

compare aspect ratio $A_R(t)$ with Boltzmann equation results (symbols)

- ▶ essentially perfect agreement (for $\tau = \eta^{\text{CE}}/P$ with η^{CE} for unitary σ_{coll})
- ▶ accurate representation of kinetic theory with 2-body collisions

- ▶ fit result: $\eta = \eta_0(mT)^{3/2}$ with $\eta_0 = 0.282$ (compare $\eta_0^{\text{CE}} = 0.269$)
- ▶ **model-independent** benchmark result

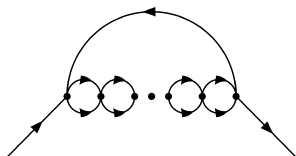
Bulk viscosity and conformal symmetry breaking

▶ in thermodynamics: $1 - \frac{2}{3} \frac{\mathcal{E}_0}{P} = \frac{\langle \mathcal{O}_C \rangle}{12\pi m a P} \sim z \frac{\lambda}{a}$

→ conformal symmetry breaking parameter

▶ impact on ζ : in high- T limit $\zeta = \frac{\lambda^{-3}}{24\sqrt{2}\pi} \left(z \frac{\lambda}{a} \right)^2$

→ physical origin:



$$\mathcal{L}_{\text{eff}} = \psi^\dagger \left(i\partial_t + \frac{\nabla^2}{2m} \right) \psi - \frac{C_0}{2} (\psi^\dagger \psi)^2; \quad C_0 = 4\pi a/m$$

→ well-defined fermionic quasiparticles with **scale-invariance breaking self-energy**

$$\text{Re}\Sigma(p) \sim \left(z \frac{\lambda}{a} \right) \sqrt{\frac{T^3}{E_p}}$$

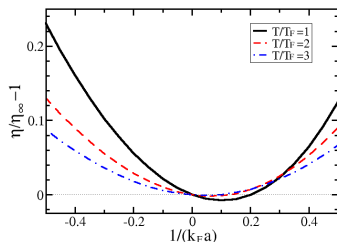
→ ratio $\zeta/\eta_\infty \sim \left(1 - \frac{2}{3} \frac{\mathcal{E}_0}{P} \right)^2$

Shear viscosity and conformal symmetry breaking

in high- T limit \rightarrow fugacity expansion:

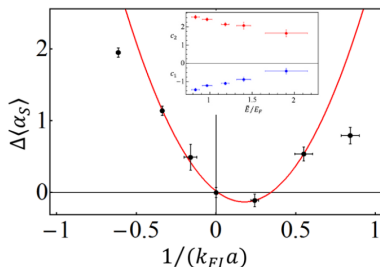
$$\eta = \eta_\infty \left(1 + \hat{c}_2 \left(\frac{\lambda}{a} \right)^2 + \hat{c}_1 \left(\frac{z\lambda}{a} \right) + \dots \right); \hat{c}_1 < 0$$

$\mathcal{O}(z\lambda/a)$ originates from in-medium effects in scattering amplitude $\mathcal{A} = C_0/(1 - \Pi C_0)$ and in quasiparticle properties



MB, T. Schäfer, PRA **90** (2014) 063615

\Rightarrow Pauli-blocking more efficient on BCS-side



\rightarrow minimum on BEC-side

\rightarrow with increasing T minimum shifted to unitarity

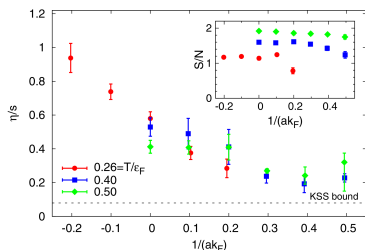
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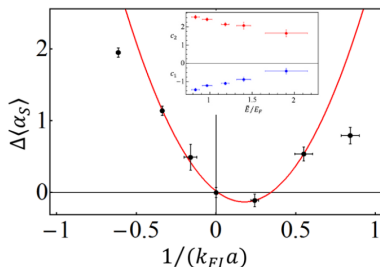
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QMC also qualitative agreement

$\rightarrow (\eta/s)_{\min} \simeq 0.2$ at $1/(k_F a) \simeq 0.4$

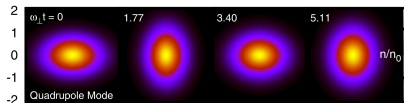


\rightarrow minimum on BEC-side

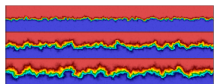
\rightarrow with increasing T minimum shifted to unitarity

Outlook and Conclusions

- ▶ extend anisotropic fluid dynamics studies to low- T to unfold full T and n dependence of η ; requires inclusion of superfluid phase
- ▶ study spin diffusion in comparison to experiments; test quantum limit near T_C , connection with momentum-diffusion, ...
- ▶ study emergence of non-hydrodynamic modes



- ▶ study fluid dynamical fluctuations



BEST
COLLABORATION

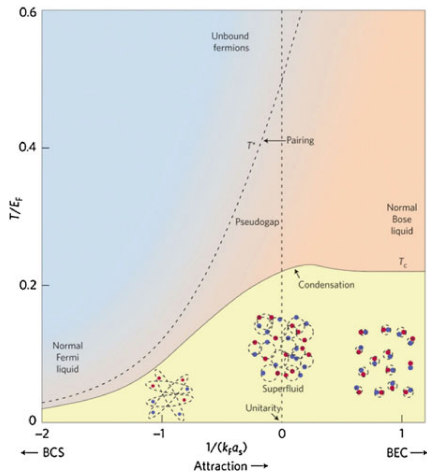
see talk by M. Nahrgang

- ▶ ...

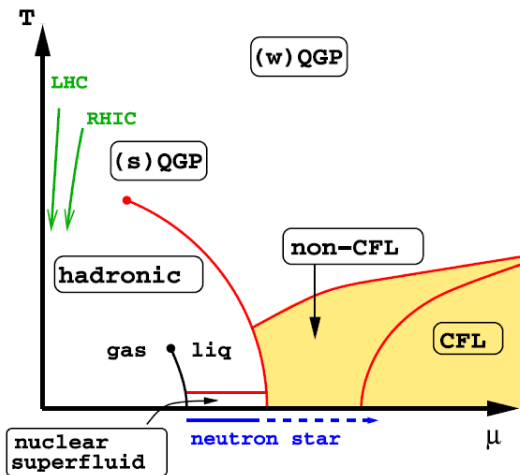
Outlook and Conclusions

- ▶ experimental and theoretical study of cold atoms spans a bridge across different areas in physics

BCS–BEC crossover in cold atoms

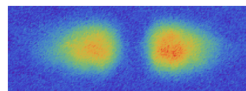


QCD phase diagram

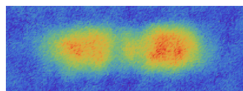


Fermi gas cloud collision

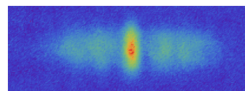
Fermi gas cloud collision:



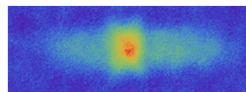
$t = 0$ ms



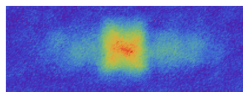
$t = 1$ ms



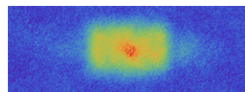
$t = 3$ ms



$t = 5$ ms



$t = 6$ ms



$t = 9$ ms

- ▶ repulsive optical potential divides trapped gas into two clouds
- ▶ clouds collide once repulsive potential switched off

⇒ formation of propagating shock fronts (fluid dynamical behavior)

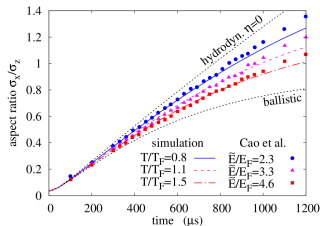
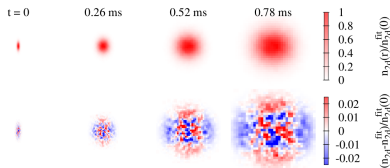
Numerical solution of the Boltzmann-equation

solve Boltzmann-equation numerically within *test-particle method*

$$(\partial_t + \vec{v} \cdot \vec{\nabla}_x - \vec{F} \cdot \vec{\nabla}_p) f_p(\vec{x}, t) = C[f_p]$$



L. Boltzmann (1844-1906)



→ not straightforward to extract dependence of η on n and T