### Study of $ee \rightarrow \gamma\gamma$

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# Signal process

- e<sup>+</sup>e<sup>-</sup>→γγ provides possibility to test various QED extensions, in particular to probe a finite size of the electron
- Need to measure absolute crosssection and angular spectrum of γγ



#### **Results from LEP**



## New physics with $e^+e^- \rightarrow \gamma \gamma$

- QED cut-off parameter (finite electron size)
  - The most sensitive way to search for non-zero electron size
  - The most precise limit from combination of the LEP data: 4.6x10<sup>-17</sup>cm (431 GeV)
- Compactified extra dimensions:  $4D \rightarrow (4+n)D$ 
  - The Plank mass is actually 1 or few TeV; the "observed" very large M<sub>PL</sub> is simply because we put wrong number of dimensions in our formula
- Excited electron e\* exchanged in the t-channel

#### Finite electron size

$$\left(\frac{d\sigma}{d\Omega}\right)_{\Lambda_{\pm}} = \left(\frac{d\sigma}{d\Omega}\right)_{\rm Born} \pm \frac{\alpha^2 s}{2\Lambda_{\pm}^4} (1 + \cos^2\theta)$$

#### Finite electron size



#### Extra dimensions and excited electrons

$$\left(\frac{d\sigma}{d\Omega}\right)_{\rm M_s} = \left(\frac{d\sigma}{d\Omega}\right)_{\rm Born} - \frac{\alpha s}{2\pi} \frac{\lambda}{M_s^4} \left(1 + \cos^2\theta\right) + \frac{s^3}{16\pi^2} \frac{\lambda^2}{M_s^8} \left(1 - \cos^4\theta\right)$$
$$\lambda = \pm 1$$

$$\begin{pmatrix} \frac{d\sigma}{d\Omega} \end{pmatrix}_{e^*} = \left( \frac{d\sigma}{d\Omega} \right)_{Born} + \frac{\alpha^2}{16} \frac{f_{\gamma}^4}{\Lambda^4} s \sin^2 \theta \left[ \frac{p^4}{(p^2 - M_{e^*}^2)^2} + \frac{q^4}{(q^2 - M_{e^*}^2)^2} \right] \\ - \frac{\alpha^2}{2s} \frac{f_{\gamma}^2}{\Lambda^2} \left[ \frac{p^4}{(p^2 - M_{e^*}^2)} + \frac{q^4}{(q^2 - M_{e^*}^2)} \right] ,$$

$$p^2 = -\frac{s}{2} (1 - \cos \theta) \text{ and } q^2 = -\frac{s}{2} (1 + \cos \theta)$$

$$\text{Igor Boyko}$$

#### Signal and main background



## **Selection cuts**

- Two photons, the most energetic above 1300 GeV, another above 1200 GeV
- No third photon above 50 GeV
- Back-to-back photons:  $\pm 10^{\circ}$  in  $\theta$ ,  $\pm 10^{\circ}$  in  $\phi$
- Track veto: no tracks with |p|>300 GeV/c within 20° from a photon candidate (even "bad tracks"!)

#### Track veto



## "Bad track" issue

- We study the absolutely extreme part of the kinematic space. Our signal is made of 1.5 TeV photons, background are 1.5 TeV electrons
- At 1.5 TeV/c the track momentum resolution approaches σ(p)/p=15% - the limit of "good track" definition.
- In the forward region nearly all electrons are reconstructed as good photons (no good track!)
- Hence, the "track veto" must reject also the "bad tracks"

### **Selection efficiency**



- Efficiency and angular acceptance much better than at LEP
- Angular coverage: down to 8-9° for SiD, down to 12-13° for ILD

### Selected collision energy



- Selection cuts ensure nearly nominal collision energy
- Average  $\sqrt{s}$ : 2878 GeV
- Reduced √s will somewhat reduce sensitivity

## Signal versus ee background



### Signal versus ee background



#### Finite electron size ±4 TeV



# Result: sensitivity to electron size versus precision on luminosity



- Sensitivity is improved by precise luminosity, but not dramatically.
- For σ(L)/L=0.5%, the 95% CL exclusion limit is Λ>6.33 TeV
- For comparison: at LEP the limit was Λ>431 GeV

# Summary

- CLIC is the best place for QED tests
- Preliminary result on sensitivity to electron size: 6.33 TeV or 3.1x10<sup>-18</sup>cm
  - LEP: 431 GeV or 4.6x10<sup>-17</sup>cm
- To do next:
  - Add (small) background from  $\gamma e \rightarrow \gamma e$
  - Estimate systematic errors (other than lumi)
  - Repeat the ILD study with SiD geometry
  - Test more models of New Physics

## Spare slides

# Track veto: Is selection efficiency still acceptable?



#### $\theta$ reconstruction

