

# Z'-induced Invisible Right-handed Sneutrino Decays at the LHC

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# Outline

## **1** Motivation

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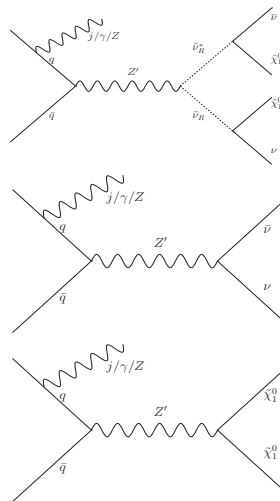
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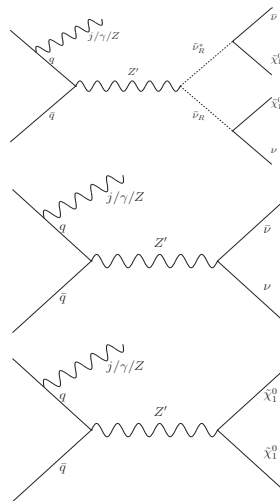
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- The advantage of the signal is that the final state objects carry a very large (transverse) missing energy, since the  $Z'$  is naturally massive and constrained by direct searches to be at least in the TeV scale region.
- Under these conditions the efficiency in accessing the invisible final state and rejecting the SM background is very high.
- Another special feature of this invisible BLSSM signal is its composition, which is often dominated by RH sneutrino decays.
- Sensitivity of the CERN machine can therefore help disentangling the BLSSM from more popular SUSY models.



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BLSSM 



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- For MC efficiency and in order to obtain reasonable statistics, we have applied a parton level (generation) cut of  $p_T(j_1) > 120$  GeV for all signals and backgrounds.
- According to the estimation of the QCD background in the SUSY mono-jets analysis, the multi-jet background can be reduced to a negligible level by requiring a large  $\cancel{E}_T$  cut, so we have applied another parton level cut of  $\cancel{E}_T > 100$  GeV for both signals and backgrounds.

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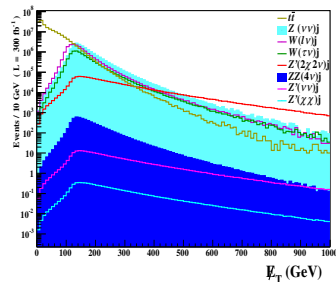
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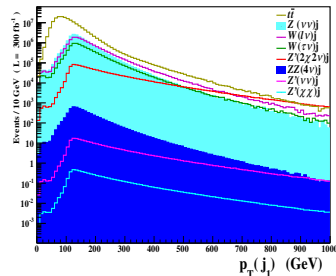
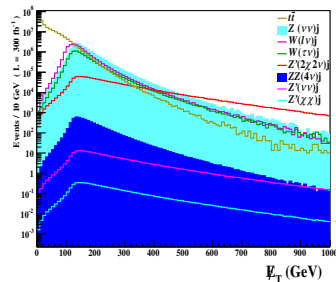
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- 5 the di-boson background  $pp \rightarrow ZZ(\rightarrow 2\nu 2\bar{\nu}) + j$ , which is generically suppressed due to its small cross section at production level but topologically mimic our signals rather well.

- As explained, owing to the large mass of the  $Z'$ , we can in the end afford a rather stiff cut in  $\cancel{E}_T$  in order to enhance our signals. This is done by setting  $\cancel{E}_T > 500$  GeV.



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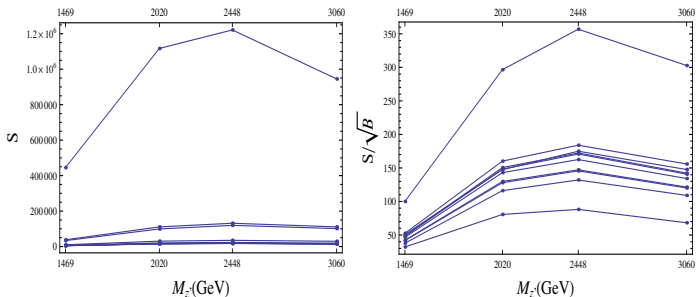
- This enhancement is also supplemented by requiring  $p_T(j_1) > 500$  GeV as the signals have harder  $p_T(j_1)$  spectra than the backgrounds.



Process	Backgrounds					Signals			
	$Z(\nu\bar{\nu})j$	$W(l\nu_l)j$	$W(\tau\nu_\tau)j$	$t\bar{t}$	$ZZj$	$Z'(2\bar{\chi}2\nu)j$	$Z'(\nu\bar{\nu})j$	$Z'(\bar{\chi}\bar{\chi})j$	
Before cuts	21573000	19248000	9390000	179058000	6621	1334400	278	7.54	
Cut	(1)	16823567 ± 1924	15817945 ± 1678	7719914 ± 1171	151390826 ± 4836	5732 ± 28	1219314 ± 324	255 ± 4.68	6.895 ± 0.77
	(2)	65275 ± 255	135191 ± 366	65423 ± 254	298430 ± 545	73 ± 8.5	130636 ± 343	27 ± 4.95	0.741 ± 0.82
	(3)	45530 ± 213	32569 ± 180	27102 ± 164	6836.8 ± 82.7	55.6 ± 7.43	118456 ± 328	25 ± 4.74	0.672 ± 0.78
	(4)	14283 ± 119	10566 ± 102	8668.5 ± 93.1	2808 ± 53	16.5 ± 4.06	35424 ± 185	7.4 ± 2.68	0.201 ± 0.44
	(5)	10831 ± 104	7395.3 ± 86	6088.7 ± 78	881.7 ± 29.7	12.2 ± 3.49	23330 ± 151	4.9 ± 2.18	0.132 ± 0.36
	(6)	8992.5 ± 94.8	6007.4 ± 77.5	4699.9 ± 68.5	379.8 ± 19.5	9.79 ± 3.13	18806 ± 136	3.9 ± 1.96	0.107 ± 0.33
	(7)	8969.8 ± 94.7	3343.1 ± 57.8	3929 ± 62.7	257.7 ± 16.1	9.78 ± 3.12	18786 ± 136	3.9 ± 1.96	0.107 ± 0.32
	(8)	8969.8 ± 94.7	871.2 ± 29.5	3207.4 ± 56.6	176.3 ± 13.3	9.77 ± 3.12	18782 ± 136	3.9 ± 1.96	0.107 ± 0.32
	(9)	8458.9 ± 92	790.2 ± 28.1	1378.8 ± 37.1	81.39 ± 9.02	9.21 ± 3.03	17878 ± 132	3.7 ± 1.92	0.102 ± 0.32
	(10)	8152.3 ± 90.3	769.9 ± 27.7	1334.4 ± 36.5	54.26 ± 7.37	8.8 ± 2.96	17357 ± 130	3.6 ± 1.89	0.098 ± 0.31

The cut flow on signal and background events after requiring the parton level cuts  $\cancel{E}_T > 100$  GeV and  $p_T(j_1) > 120$  GeV for  $M_{Z'} \simeq 2448$  GeV and  $g_{B-L} = 0.4$  at  $\sqrt{s} = 14$  TeV with  $\mathcal{L}dt = 300 \text{ fb}^{-1}$ :

- (1)  $n(\text{jets}) \geq 1$  with  $|\eta(j_1)| < 2$ ; (2)  $p_T(j_1) > 500$  GeV; (3)  $\cancel{E}_T > 500$  GeV;  
 (4)  $\Delta\phi(j_2, \cancel{E}_T) > 0.5$ ; (5) veto on  $p_T(j_2) > 100$  GeV,  $|\eta(j_2)| < 2$ ;  
 (6) veto on  $p_T(j_3) > 30$  GeV,  $|\eta(j_3)| < 4.5$ ; (7) veto on  $e$ ;  
 (8) veto on  $\mu$ ; (9) veto on  $\tau$ -jets; (10) veto on  $b$ -jets.



Number of events from the sum of all signals ( $S$ ) versus  $M_{Z'}$  and number of events from the sum of all signals divided by the square root of the total background ( $S/\sqrt{B}$ ) versus  $M_{Z'}$ . Rates are given at 14 TeV for an integrated luminosity of  $300 \text{ fb}^{-1}$ .

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- We were able, by BLSSM specific cuts, to establish signals with significances well above  $5\sigma$  discovery limit. Indeed, this has been possible due to the fact that the BLSSM mediator of such invisible signals is a very heavy  $Z'$ , thereby transferring to its decay products large transverse momenta that can be generically exploited in all cases for background reduction.

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- For all topologies considered, the dominant component of the signal is via RH sneutrinos (above neutrinos and neutralinos), this fact help in distinguishing the BLSSM model from the popular SUSY models.

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*Thank you*