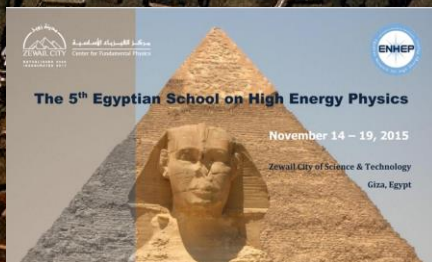


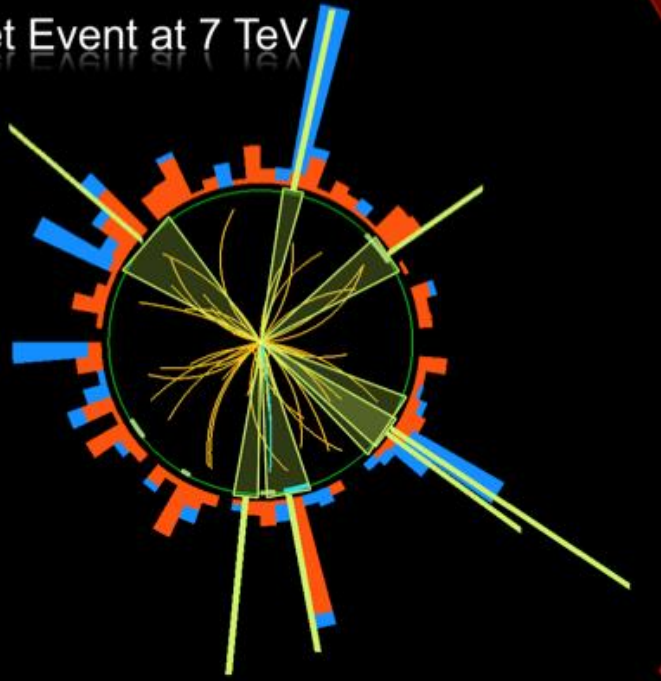
Dark Matter and Searches at the LHC

Albert De Roeck
CERN, Geneva, Switzerland
Antwerp University Belgium
UC-Davis California USA
IPPP, Durham UK
NTU, Singapore

November 14-19 Cairo, Egypt



Multi Jet Event at 7 TeV



Outline

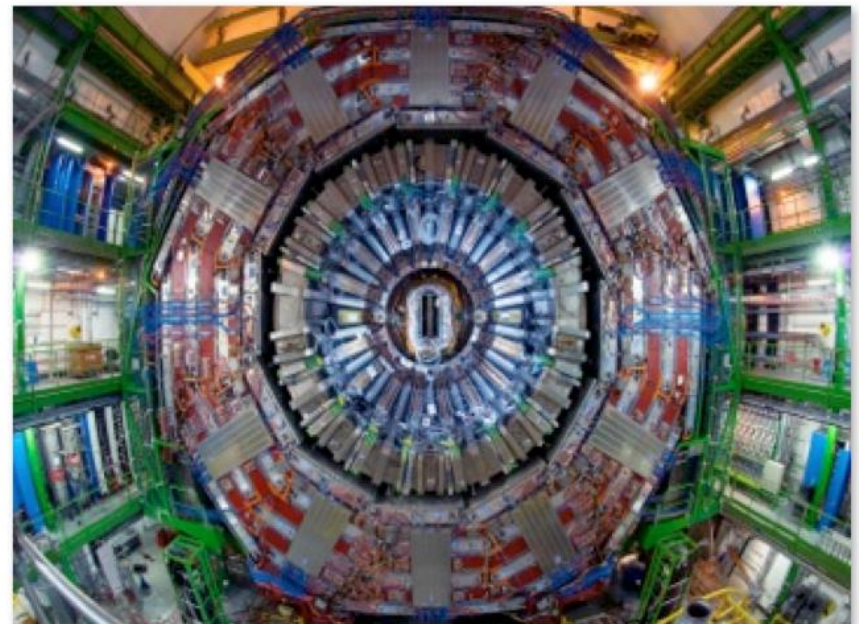
- Introduction: dark matter and the WIMP miracle
- The Higgs and dark matter
- Supersymmetry searches
- Generic searches via missing E_T , so called mono-objects
New: mono-Higgs production
- Summary

Dark Matter: Complementary Searches?

After the discovery of the Higgs particle @ the LHC:

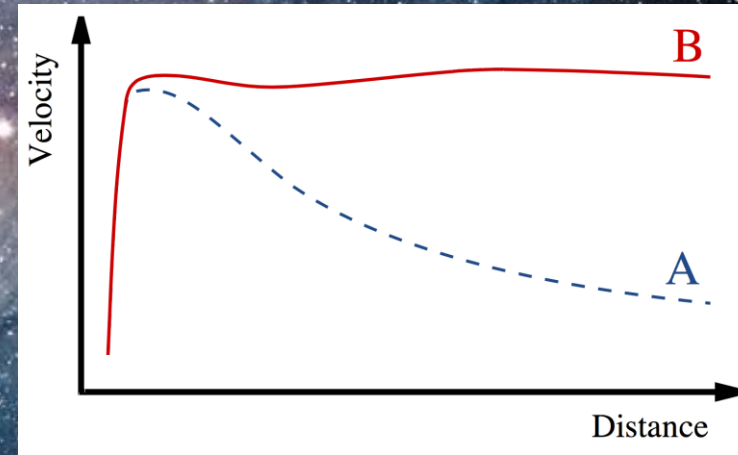
Dark matter is the next important physics problems to tackle for the LHC

The search is complementary to other experimental techniques used.



Dark Matter: The Next Challenge

Astronomers found that most of the matter in the Universe must be invisible Dark Matter



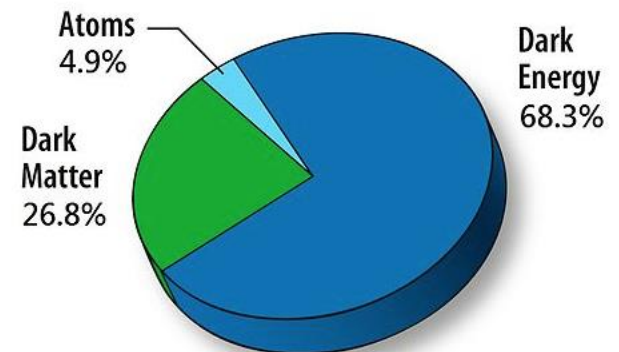
'Supersymmetric' particles ?



F. Zwicky 1898-1974



Vera Rubin ~ 1970

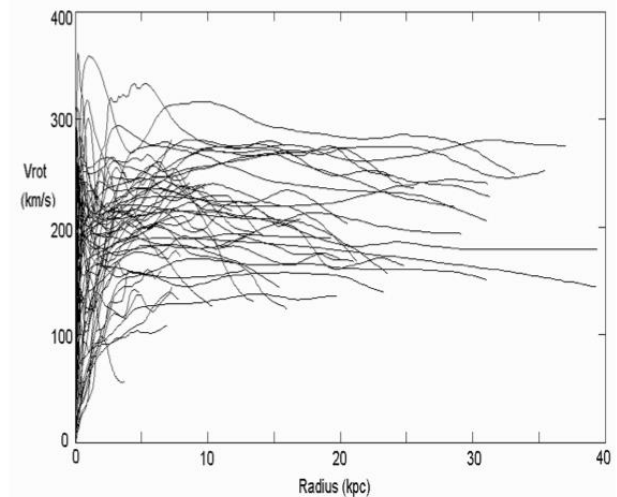
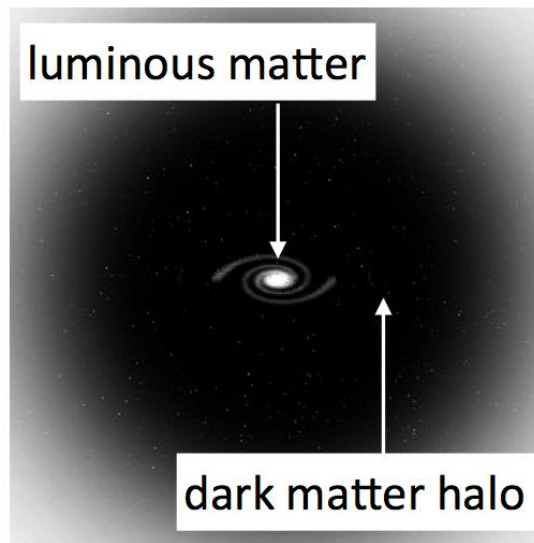
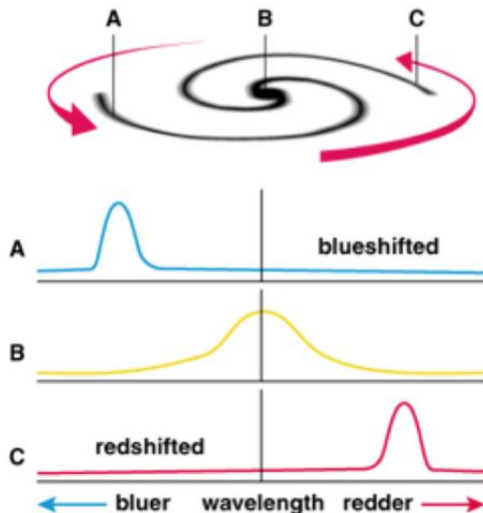
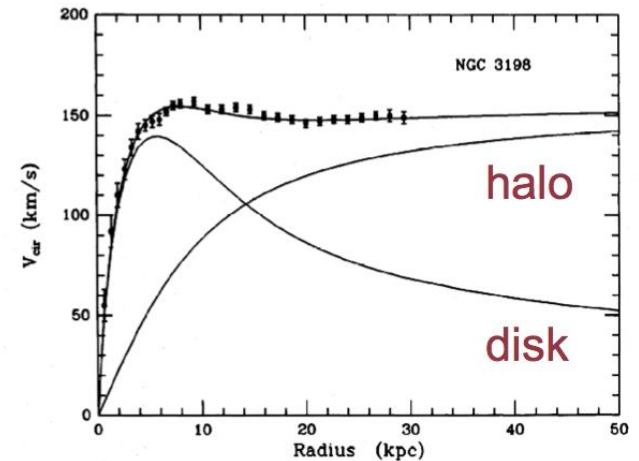


Galactic Rotation

- Starting in the 1970's, measured velocity vs. radius of edge-on spiral galaxies
- They found them to be flat, consistent with ~10x as much "dark" mass...

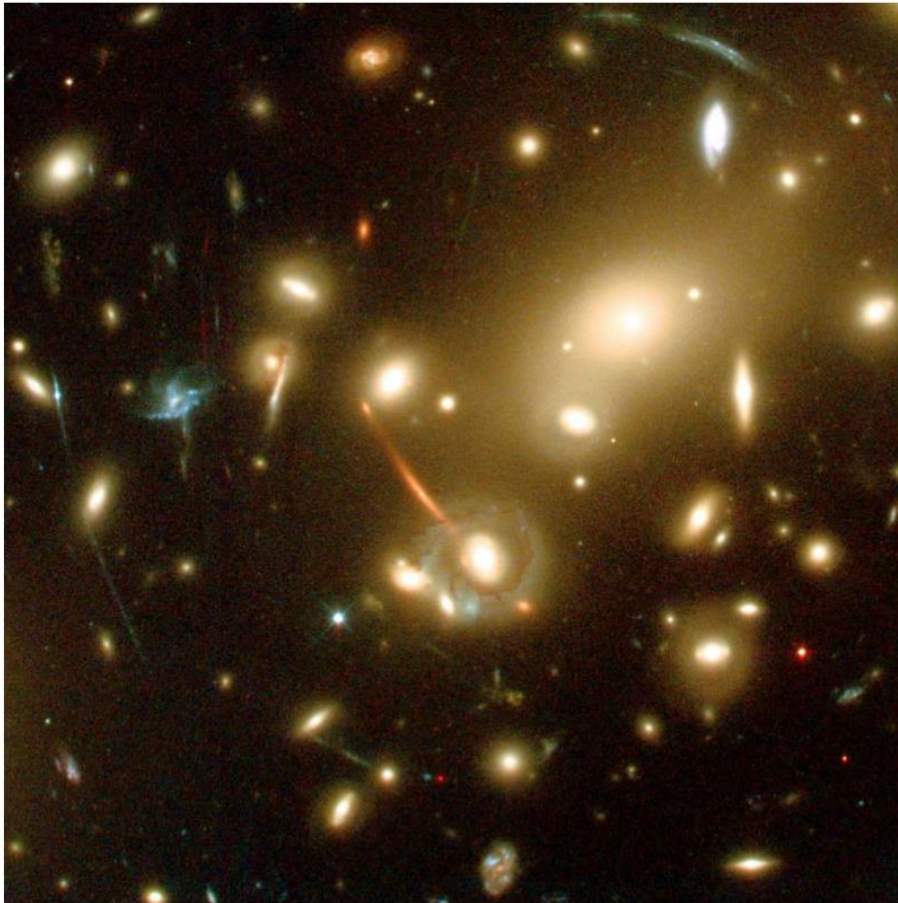
...and not just one galaxy

DISTRIBUTION OF DARK MATTER IN NGC 3198

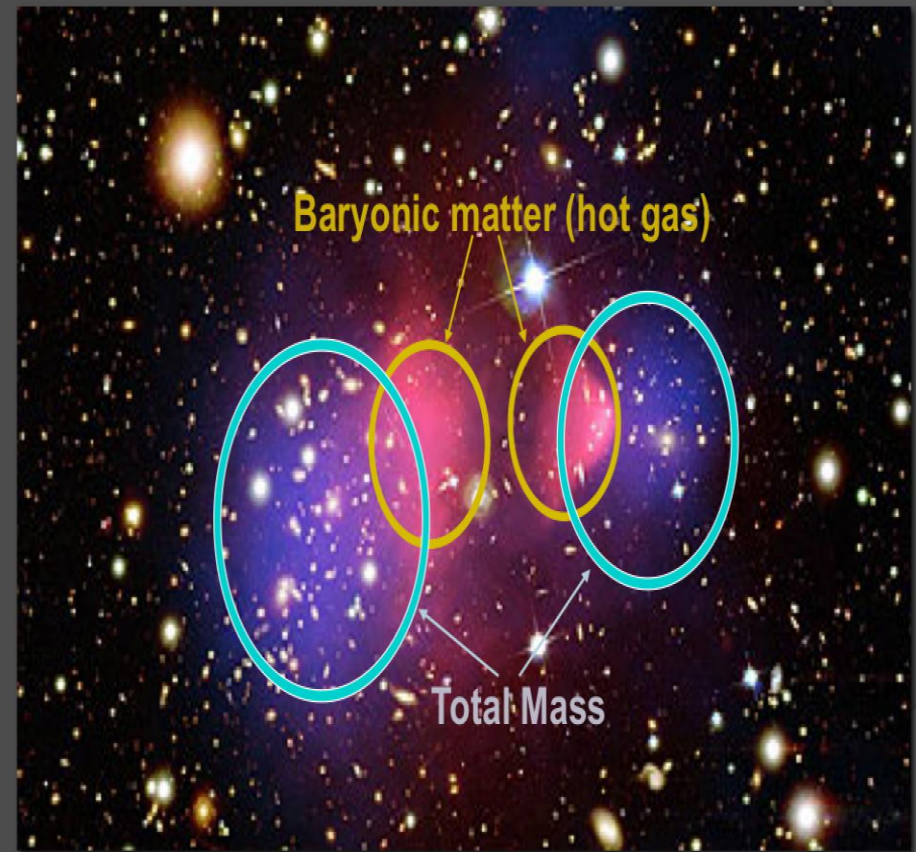


Evidence Piling-up

- Gravitational Lensing
 - much more lensing than can be explained by visible mass

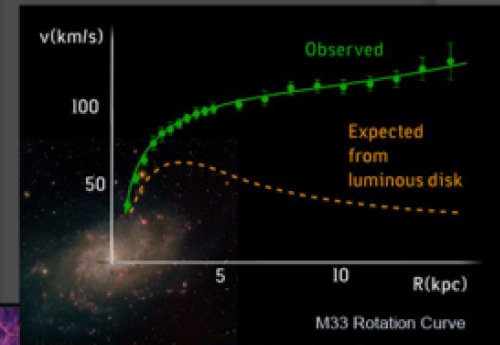
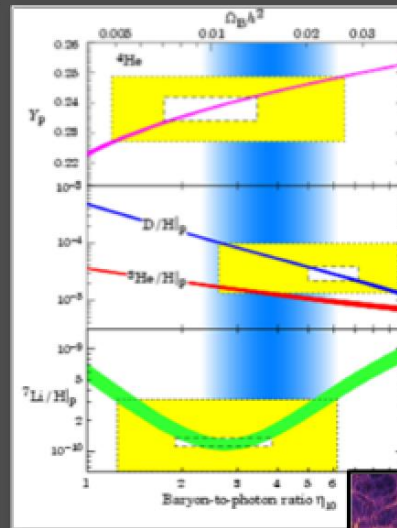
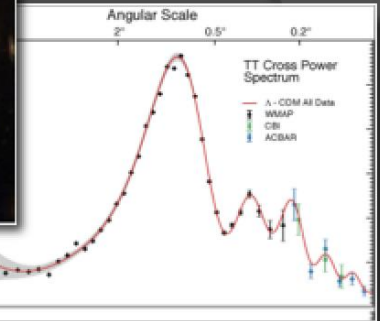


- Bullet Cluster; colliding galaxies
 - Composite x-ray, visible image, 10x DM

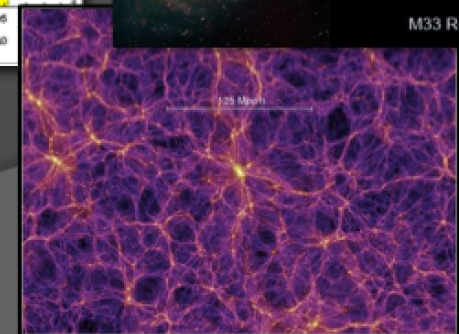


Evidence Piling-up

- There is a wide variety of evidence indicating that dark matter exists
- Each of these observations infer dark matter's presence uniquely through its gravitational influence
- To-date, no (non-controversial) observations have been made of dark matter's electroweak or other non-gravitational interactions



Instead of dark matter, might we not understand gravity?



Particle Dark Matter?

- We know only little about the nature of dark matter:
 - Cold (non-relativistic)
 - Stable
 - Dark and collisionless (no electric charge or QCD color)
- No particle contained in the Standard Model fulfills these criteria
- This leaves us with a vast range of possibilities from Planck/GUT scale “WIMPzillas” to ultra-light axions
- Dark matter candidates in the form of weakly interacting particles with masses in the GeV-TeV range (WIMPs) stand out for their
 - Testability
 - Theoretical motivation (solution to electroweak hierarchy problem)
 - The “WIMP Miracle”

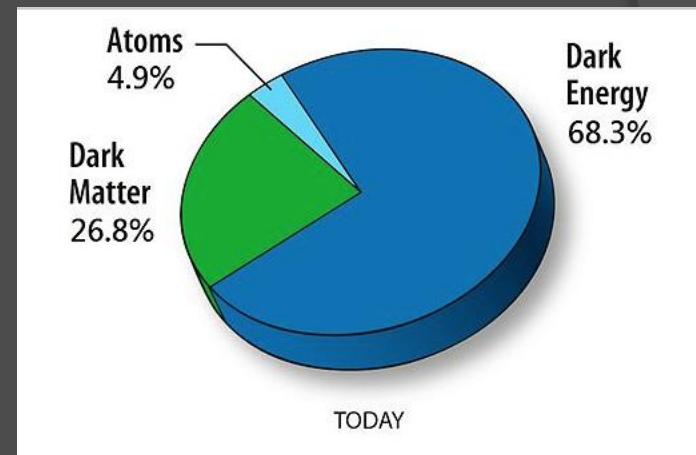
The observed density of dark matter is of the magnitude expected for a thermal relic weakly-interacting massive ($\sim 1-1000$ GeV) particle (WIMP).

Particle Dark Matter?

The Dark Matter Candidate Zoo

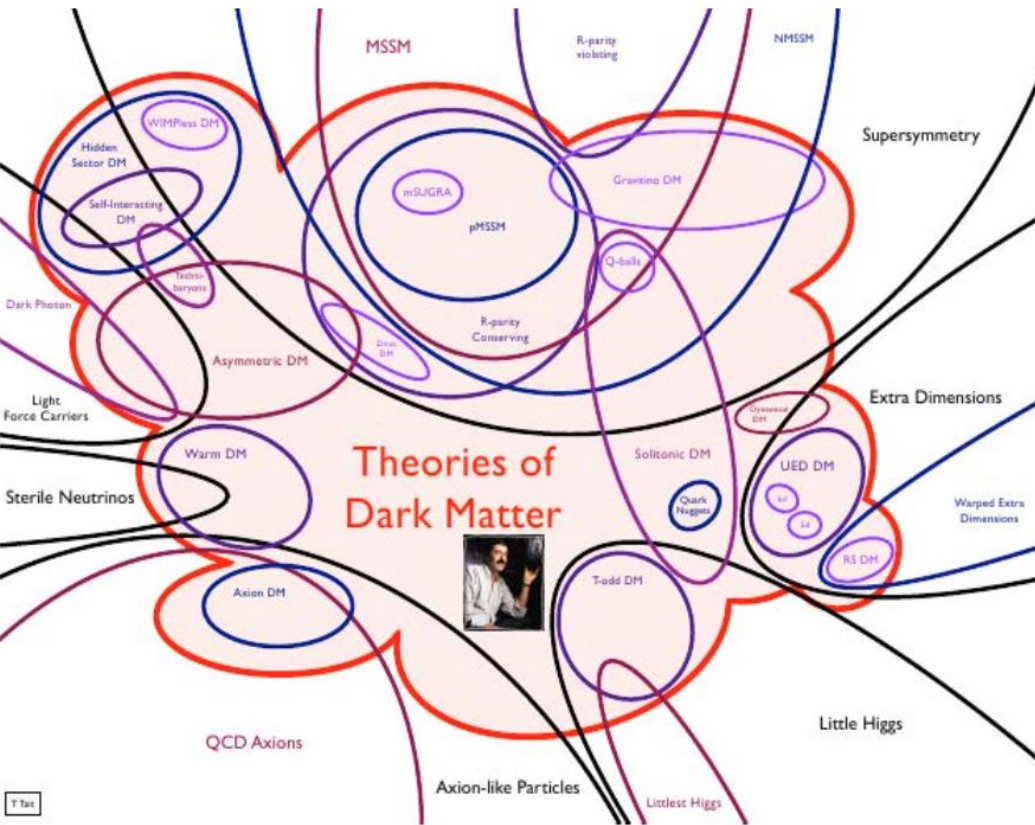
From D. Hooper

- Neutralinos (higgsino, bino, wino, singlino)
- Axinos
- Gravitinos
- Sneutrinos
- Axions
- Sterile neutrinos
- 4th generation neutrinos
- Kaluza-Klein photons
- Kaluza-Klein gravitons
- Brane world dark matter/D-matter
- Little higgs dark matter
- Light scalars
- Superheavy states (*ie.* “WIMPzillas”)
- Self-interacting dark matter
- Super-WIMPs
- Asymmetric dark matter
- Q-balls (and other topological states)
- CHAMPs (charged massive particles)
- Cryptons, mirror matter, and many, many, many others...

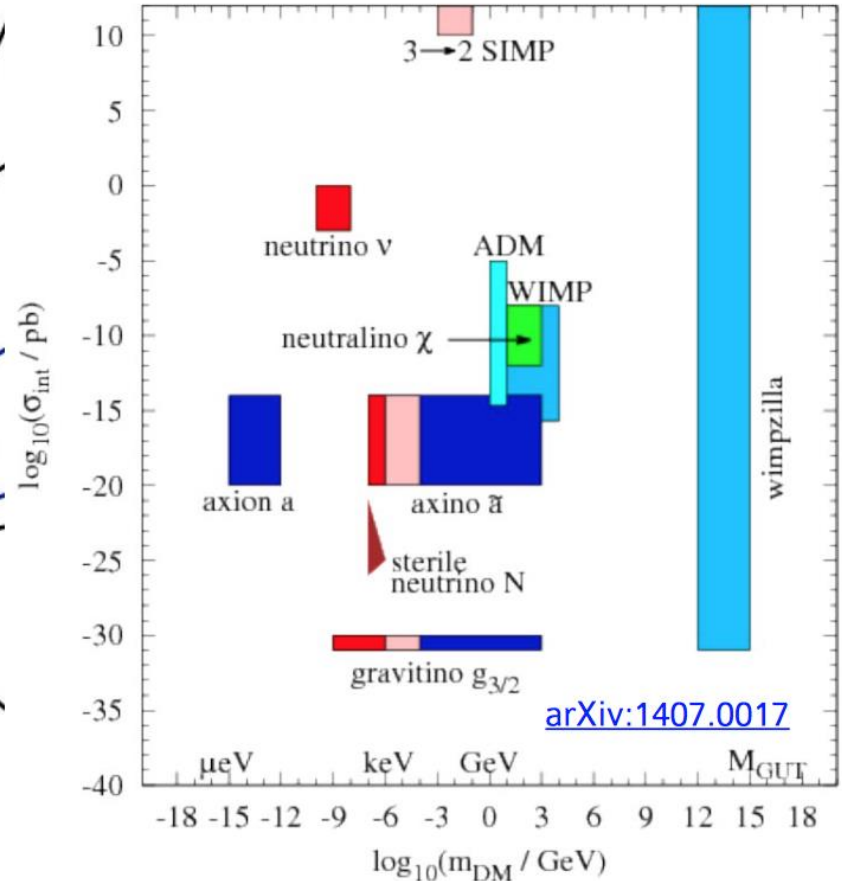


The ‘Standard Model’ of the Universe indicated by astrophysics and cosmology

Theories on Dark Matter



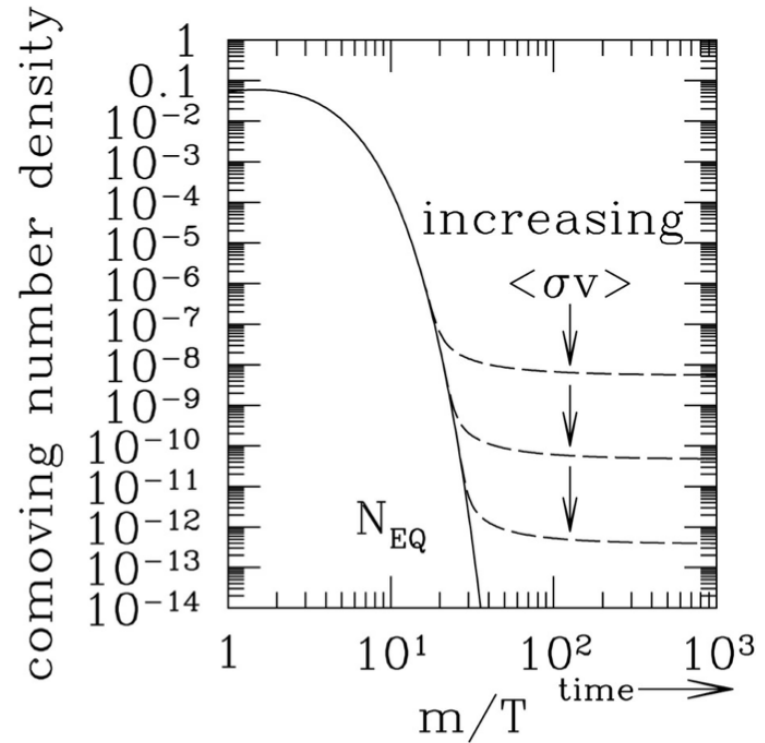
[Tim Tait, DM@LHC 2013](#)



(Our) preferred DM candidate matches cosmological observations (e.g. thermal relic density): dark, stable, cold, weakly interacting with SM particles, mass of up to a few TeV → a **WIMP**

WIMPs

- Perhaps Dark Matter is a particle with weak-scale mass?
 - *Weakly Interacting Massive Particles (WIMPs)*
 - Produced in the Big Bang, interact via $\chi + \chi \rightarrow q + q$
- As the universe expands and the temperature drops...
 - WIMPs become diluted, interact less often and ‘freeze out’.
 - Higher cross-section $\langle\sigma v\rangle$ yields lower relic density

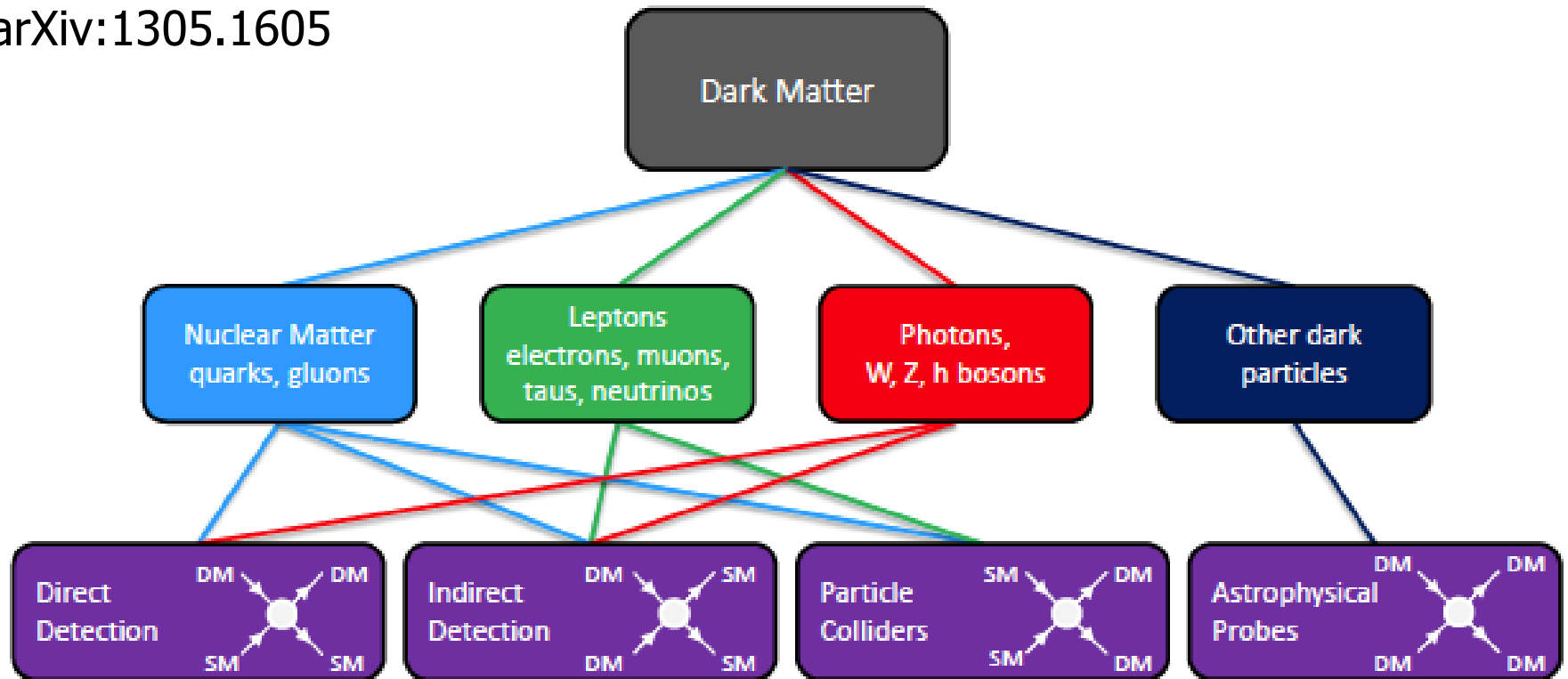


Weakly-interacting massive particles naturally provide the right relic abundance - “WIMP miracle”

Searches for Dark Matter

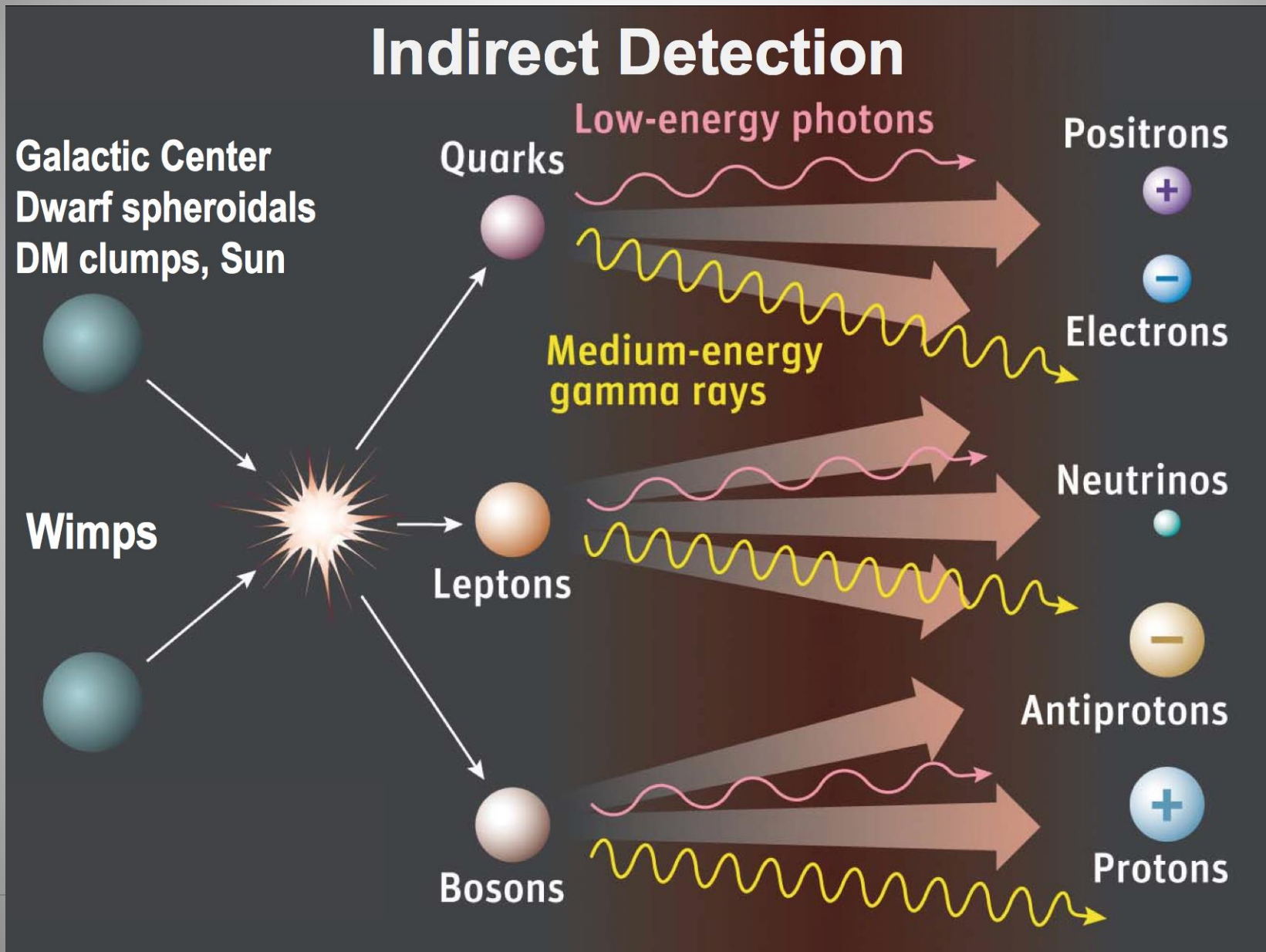
Search for WIMP candidates in events with Missing Transverse Momentum
EG: SUSY searches, monojet and mono-photon searches, W' searches...

arXiv:1305.1605

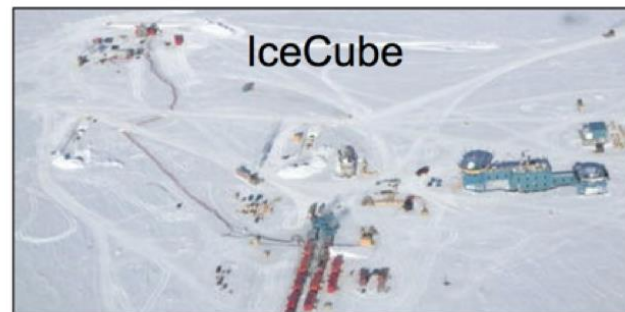
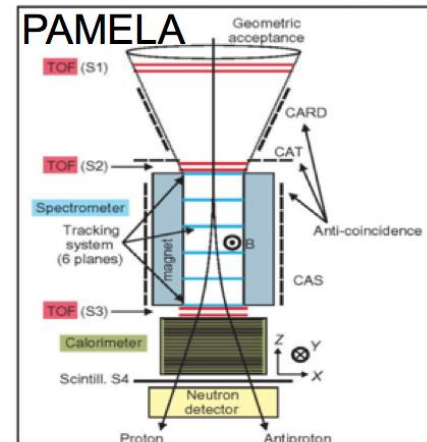
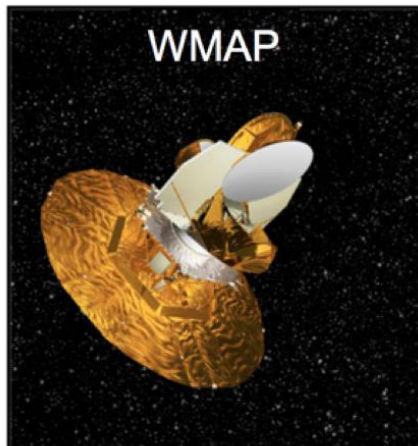


+ CAST experiment, searching for axion DM

Dark Matter: Indirect Detection



Indirect Detection Experiments



Indirect Detection

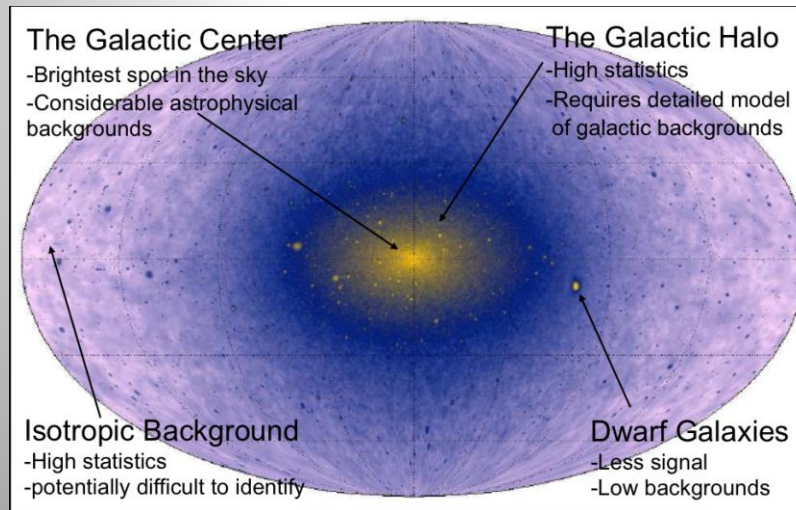
Some scientists are believers!!

arXiv:1402.6703v1

The Characterization of the Gamma-Ray Signal from the Central Milky Way: A Compelling Case for Annihilating Dark Matter

Tansu Daylan,¹ Douglas P. Finkbeiner,^{1,2} Dan Hooper,^{3,4} Tim Linden,⁵
Stephen K. N. Portillo,² Nicholas L. Rodd,⁶ and Tracy R. Slatyer^{6,7}

Using gamma-ray data from the FERMI satellite
DM ~ 40 GeV annihilation into b anti-b quarks?



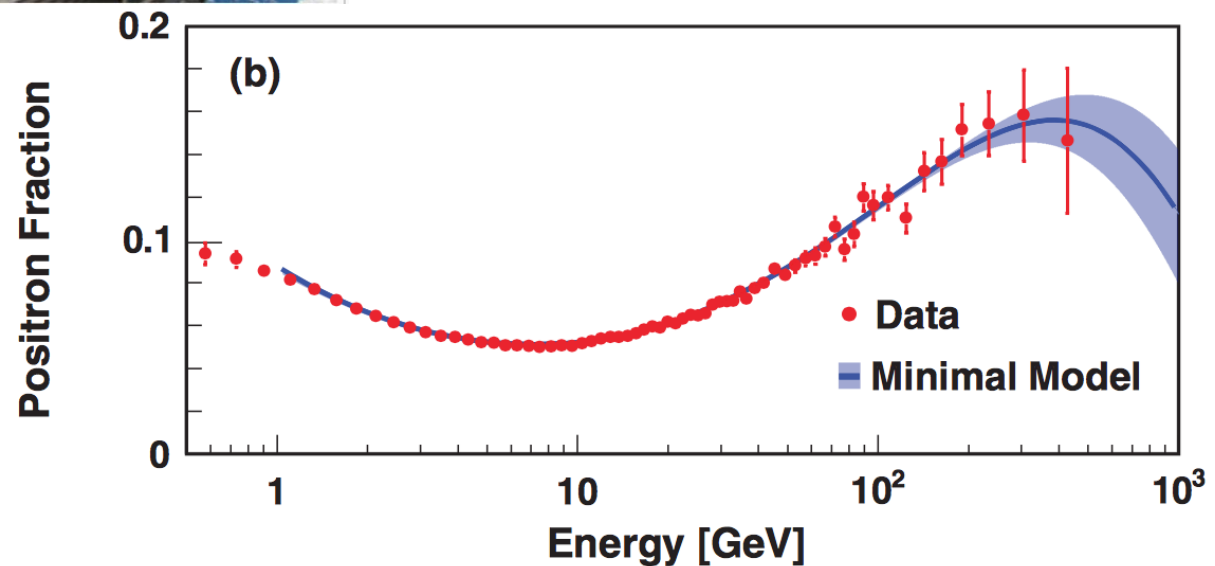
Also the 3.5 KeV line: light axion-like particle annihilation?

Dark Matter Searches in Space



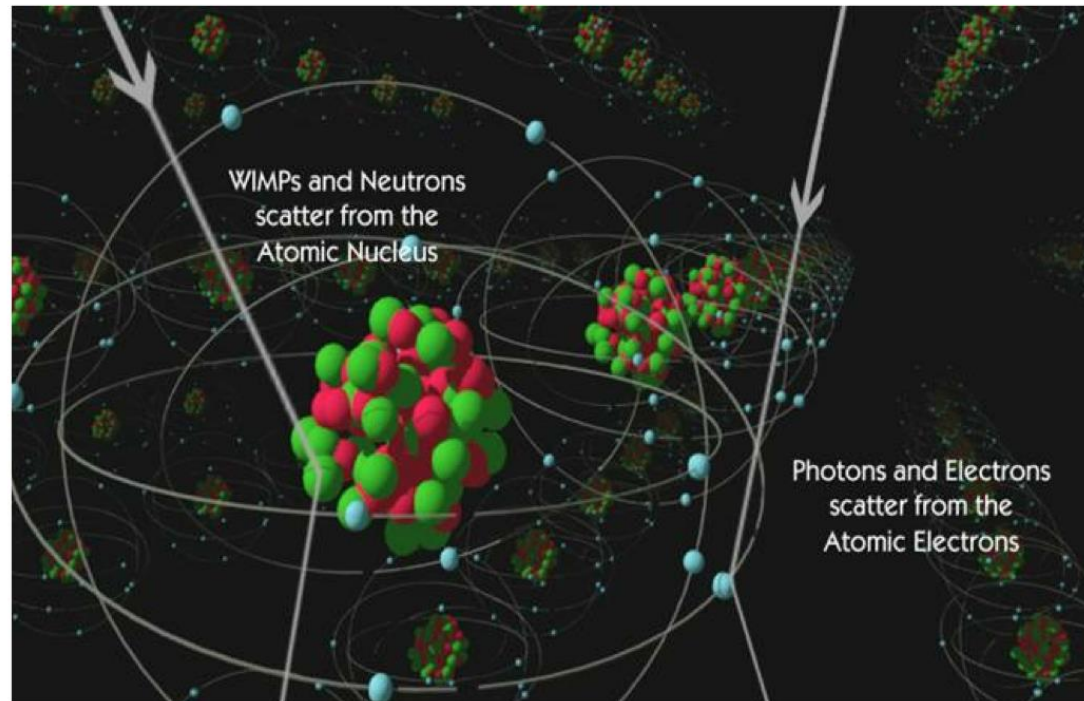
Searching for anti-matter in space.
A Dark Matter signal?

Difficulty: astrophysical sources...?



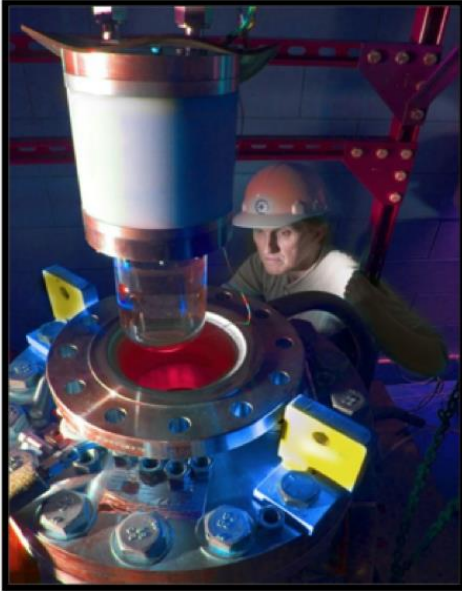
Dark Matter: Direct Detection

- Direct detection experiments: nuclear recoil from DM collision
 - Extremely sensitive, extremely difficult... extremely successful!
 - Excesses observed but not confirmed (10 GeV DM candidate?)
- Need for independent verification from non-astrophysical experiments
 - Low mass region not accessible to direct detection experiments
 - Limited by threshold effects, energy scale, bkgnds; spin-dependent couplings difficult...

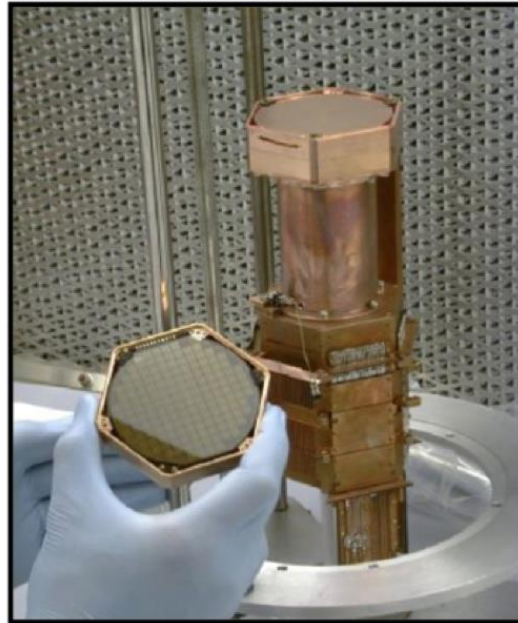


Direct Detection: Examples

COUPP



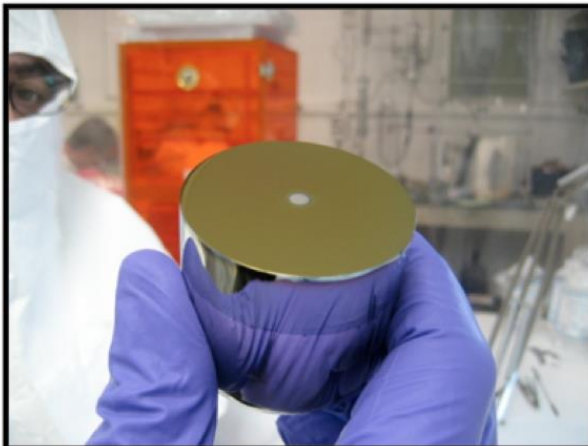
CDMS



CRESST



CoGeNT



Xenon

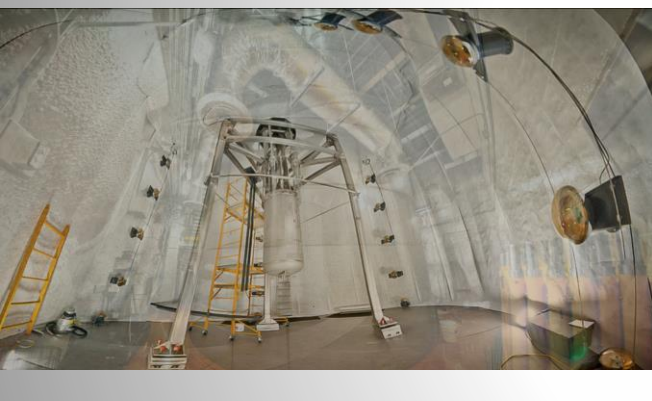


(+ EDELWEISS,
DAMA, EURECA,
ZEPLIN, DEAP, ArDM,
WARP, LUX, SIMPLE,
PICASSO, DMTPC,
DRIFT, KIMS,
DARKSIDE...)

Direct Searches for Dark Matter

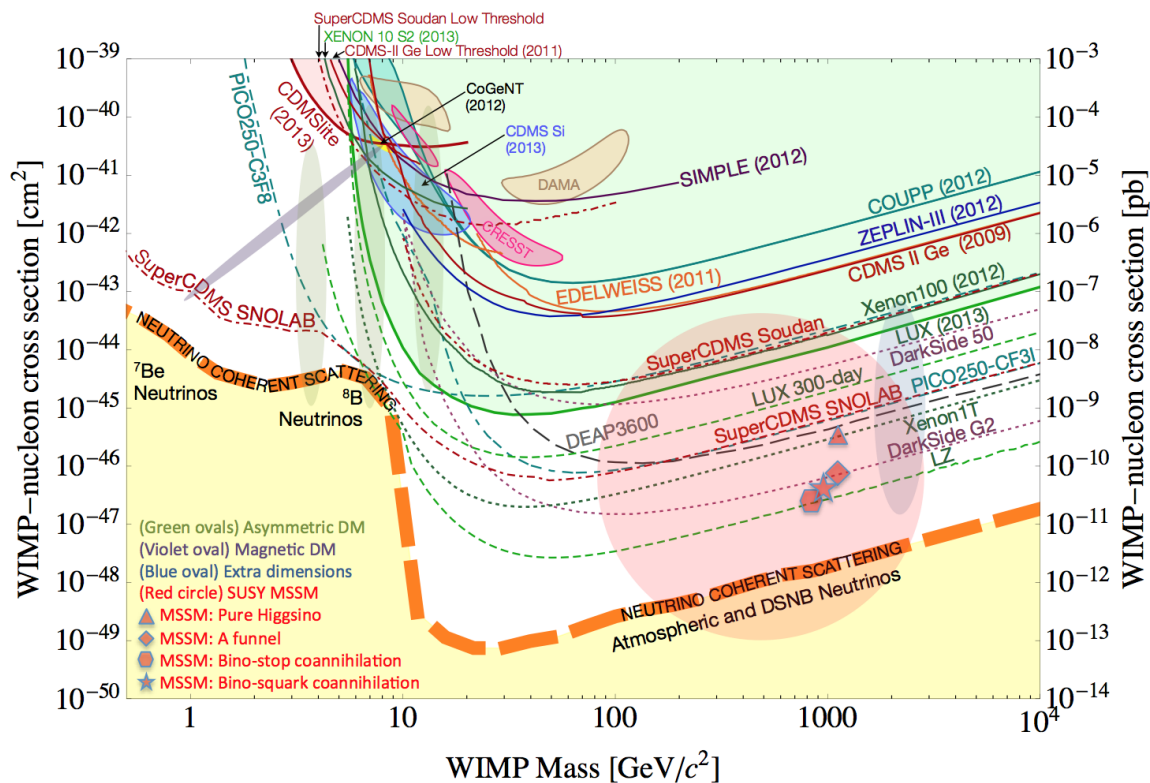


State of the art today:
Driven by the results of
the **LUX** experiment



Intensive campaign of
direct detection
experiments since more
than ~ 20 years

No (real) sign so far...

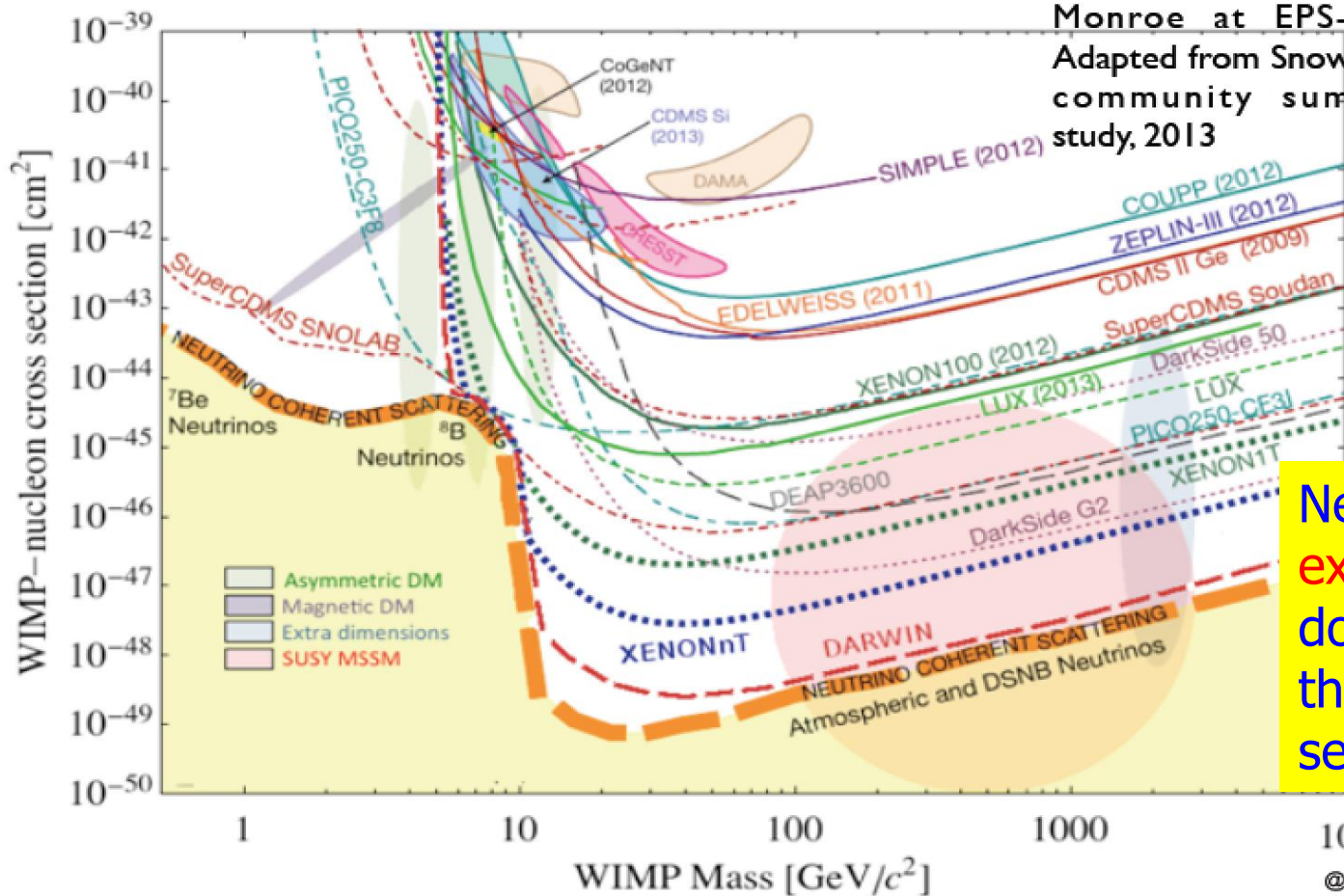


Direct Searches for Dark Matter

Underground low noise experiments

No non-ambiguous signal yet!!

There is a very large number of projects which are under construction or being planned for the future.

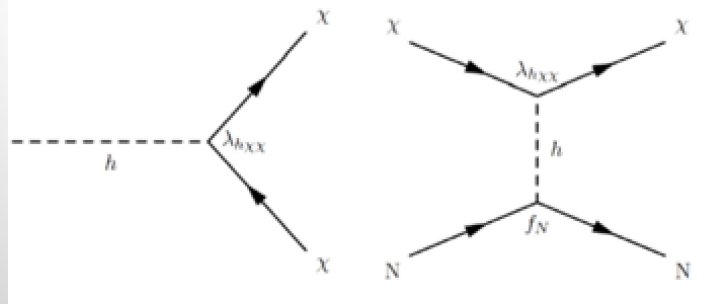


Monroe at EPS-HEP,
Adapted from Snowmass
community summer
study, 2013

New 8/9/15: XMASS
experiment (Japan)
does not confirm
the periodical effect
seen by DAMA/LIBRA

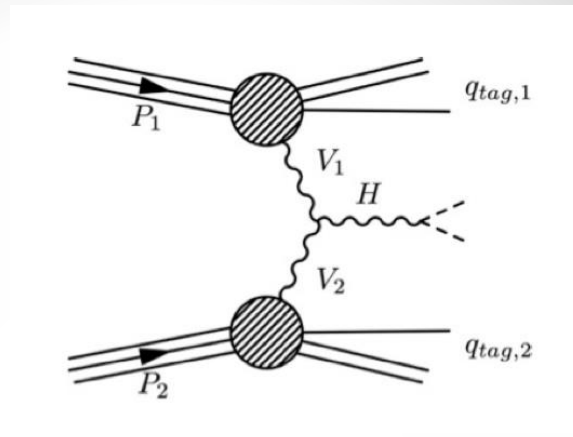
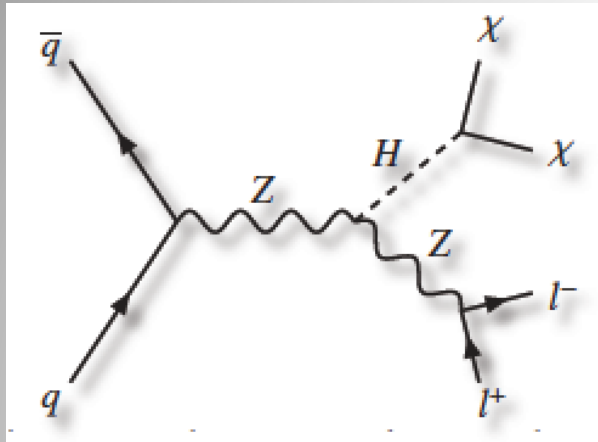
The LHC

Dark Matter and the Higgs



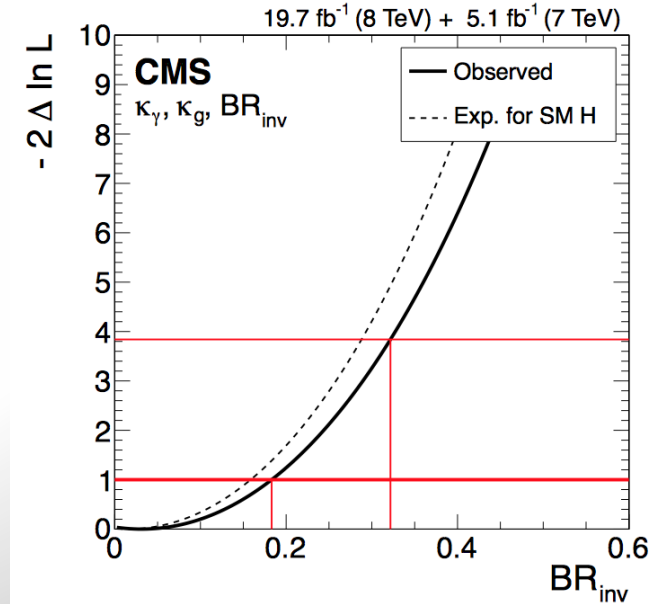
“higgs portal models”
Eg: arXiv:1205.3169

Invisible Higgs Decay Channel

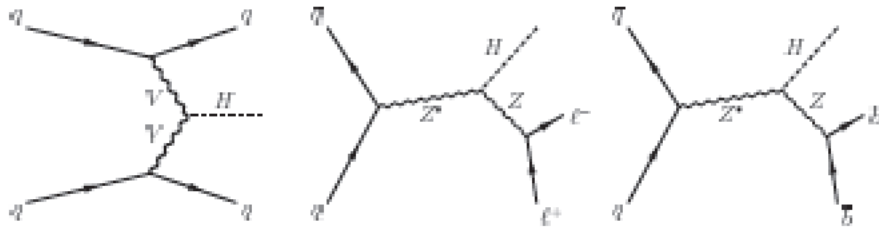


arXiv:1412.8662

- Possible decay of the Higgs in dark matter particles (if $M < M_H/2$)
- Different searches:
 - Direct search
 - Look for the invisible decay channels
 - Indirect search
 - Make a global fit of all production and decays (and some modest assumptions)

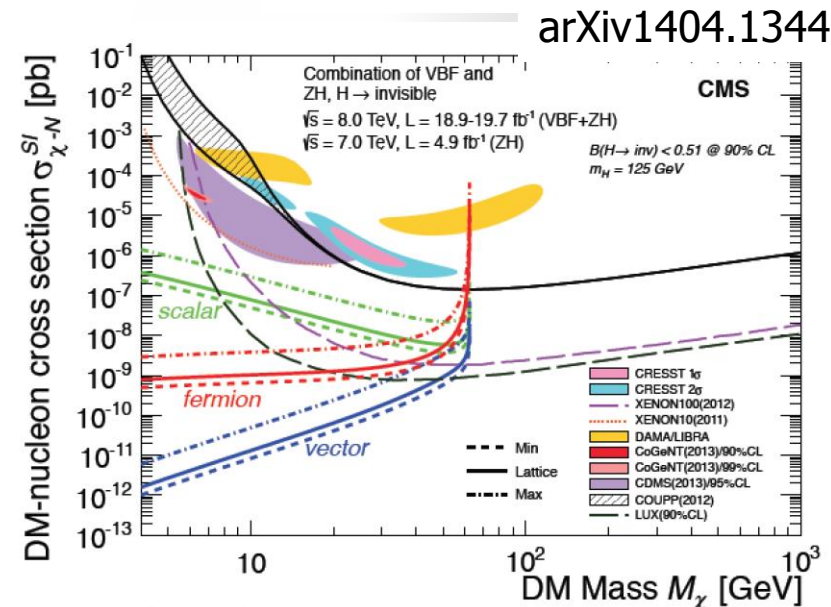
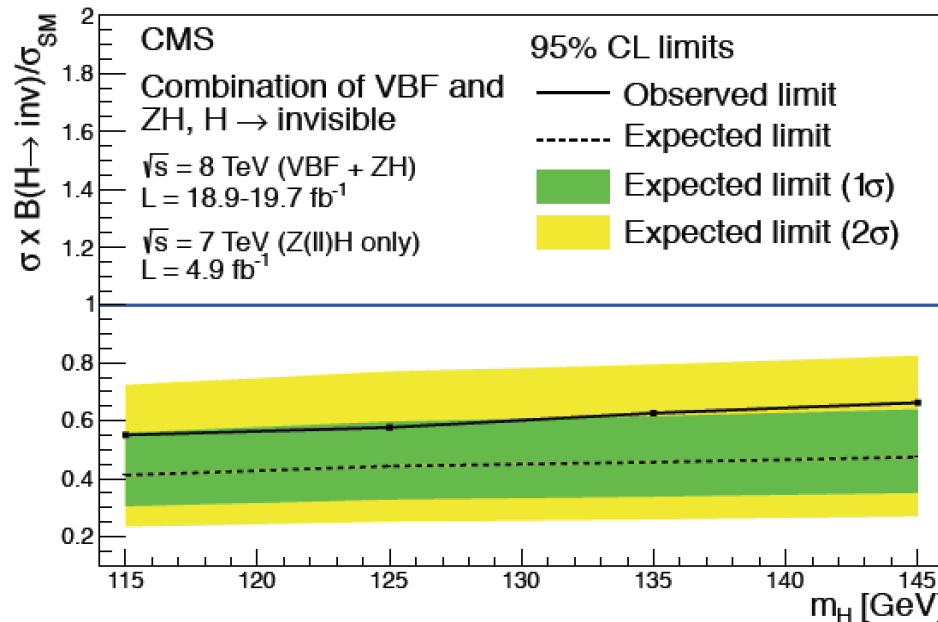
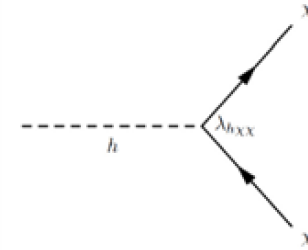


Invisible Higgs Decay Channel



Search for invisible Higgs decays using
 $Z+H \rightarrow 2 \text{ leptons} + \text{missing } E_T$
 $VBF H \rightarrow 2 \text{ jets} + \text{missing } E_T$
 Possible decay in Dark Matter particles
 (if $M < M_H/2$): Higgs Portal Models

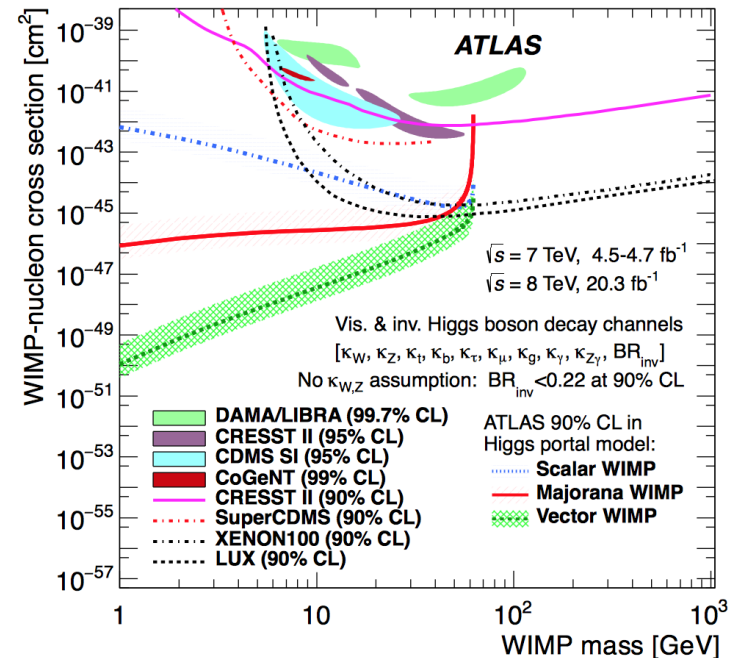
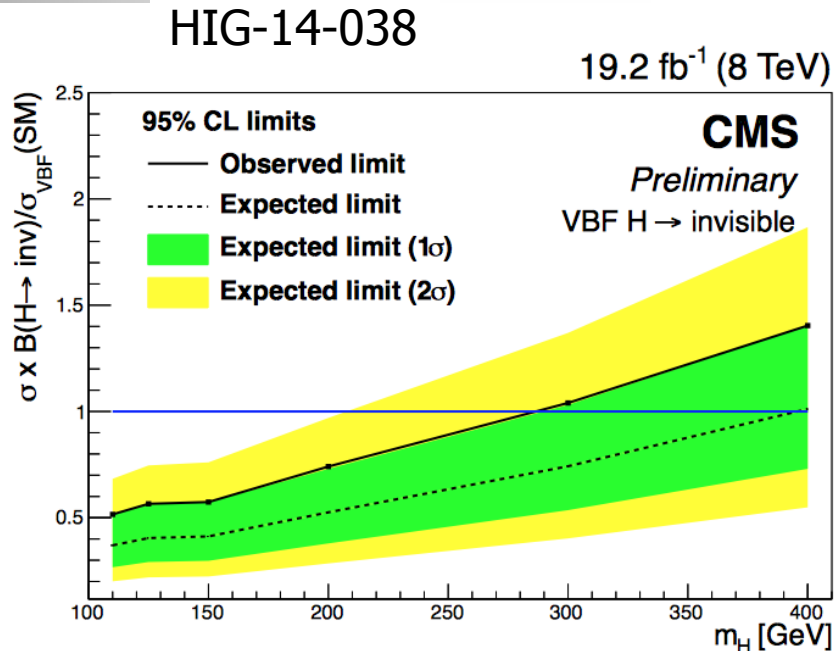
Combined result from the three channels
 $BR(H \rightarrow \text{invisible}) < 58\% (44\% \text{ exp})$ at 95% CL.
 for a Higgs with a mass of 125 GeV



New: Invisible Higgs in VBF

- CMS: VBF process with $p_T > 35(30)$ GeV, $M_{jj} > 700$ GeV and $\Delta\eta_{jj} > 3.5$ parked dataset with 11 fb^{-1} in run1. The $\text{BR}(H \rightarrow \text{inv}) = 0.57 (0.40)$ observed (expected) at 95% CL.
- ATLAS: First VBF analysis: 0.28 (0.31) observed (expected)
- CMS All channel combination 47% (35%) obs (exp).

arXiv:1509.00672

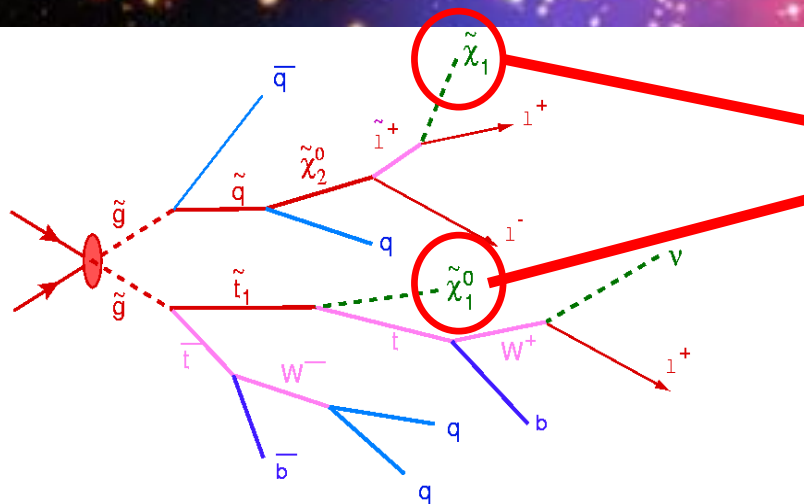
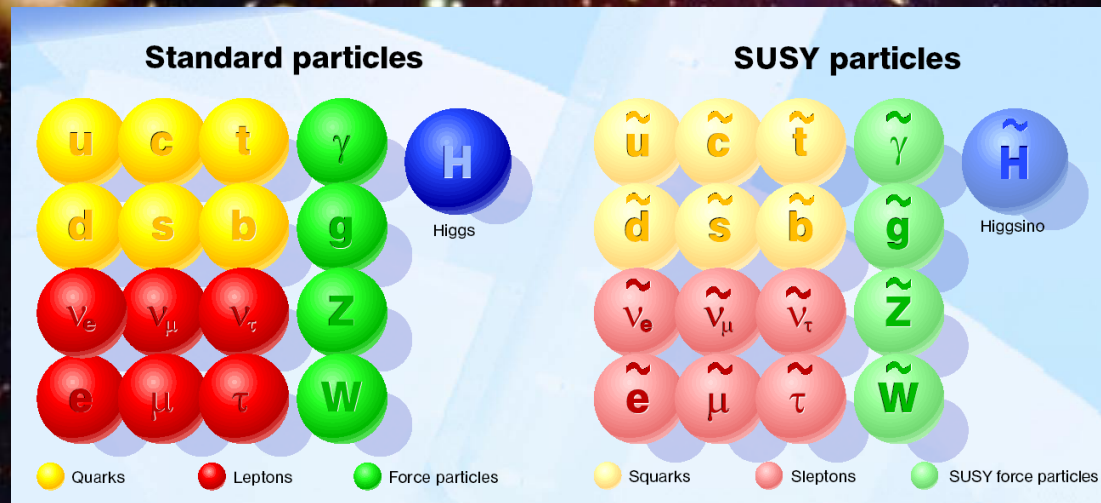


Searches for Supersymmetry

Supersymmetry was not “invented” to solve the dark matter problem, but can provide a great solution!

WIMP Dark Matter candidate comes for free ... (in many cases)

Supersymmetry: a new symmetry in Nature?



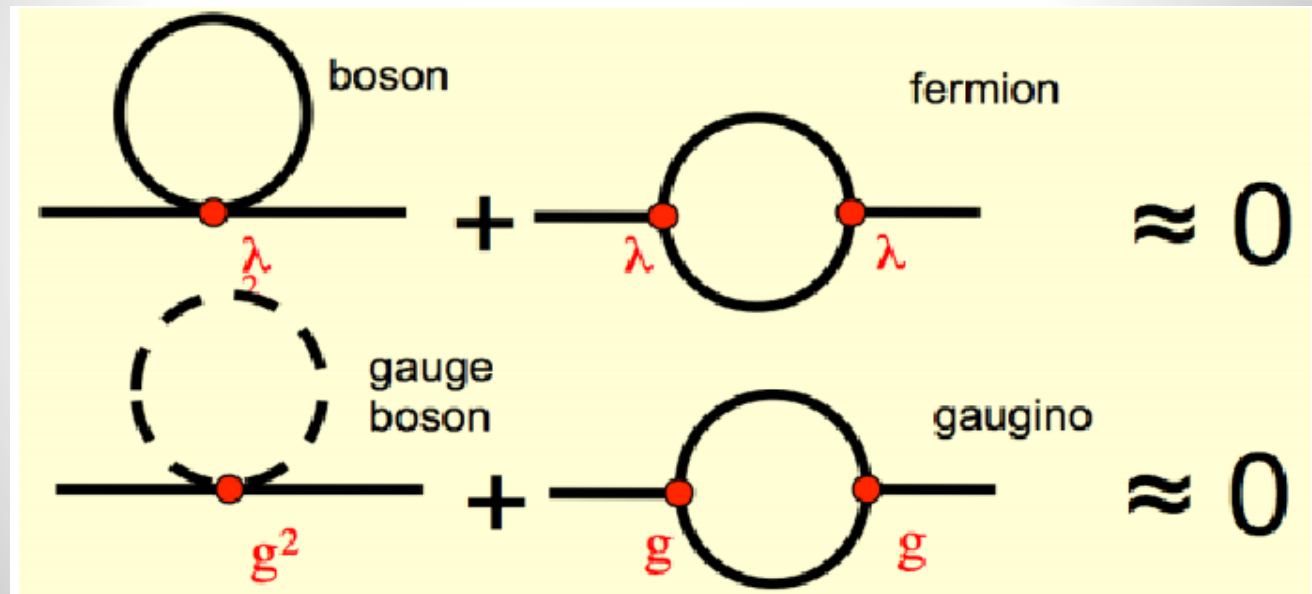
Candidate particles for Dark Matter
 \Rightarrow Produce Dark Matter in the lab

SUSY particle production at the LHC

Picture from Marusa Bradac

Supersymmetry

Supersymmetry (SUSY) → assumes a new hidden symmetry between the bosons (particles with integer spin) and fermions (particles with half integer spin). Stabilize the Higgs mass up to the Planck scale



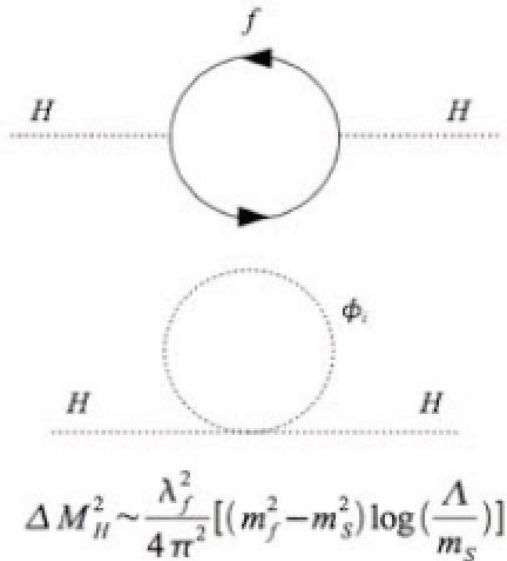
Fermion and boson loops cancel, provided $m_{\tilde{f}} \leq \text{TeV}$.

Why weak-scale SUSY ?

- ☞ stabilises the EW scale: $|m_F - m_B| < O(1 \text{ TeV})$
- ☞ predicts a light Higgs $m_h < 130 \text{ GeV}$
- ☞ accomodates gauge unification
- ☞ accomodates heavy top quark
- ☞ dark matter candidate: neutralino, sneutrino, gravitino, ...
- ☞ consistent with EW precision tests

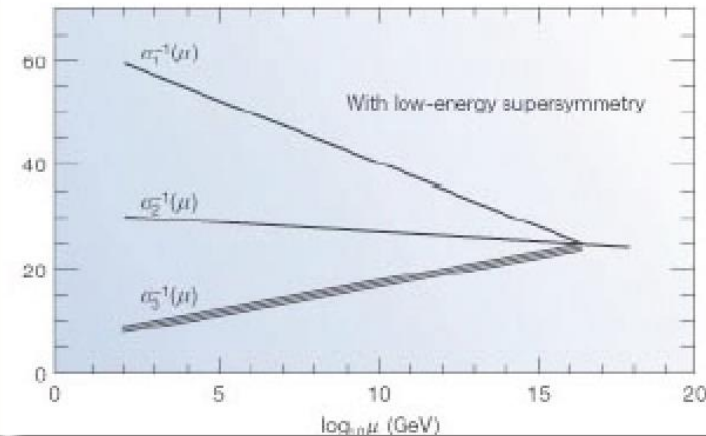
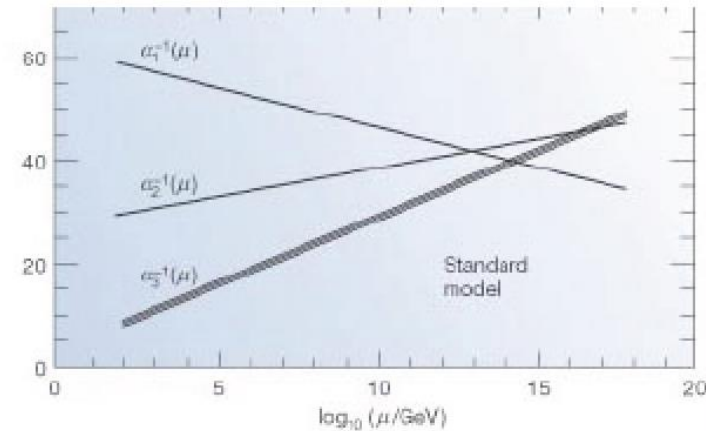
Discovering SUSY – A revolution in particle physics!!

Summary: Why SUSY is good for you!!



◆ Elegant solution to the hierarchy problem (i.e., why the Higgs mass is not at the Planck scale)

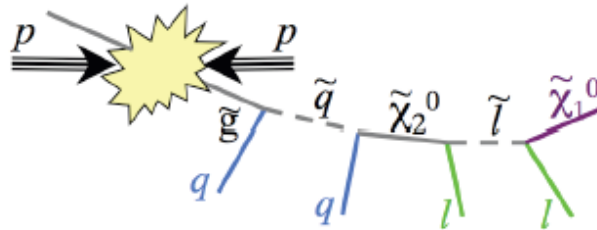
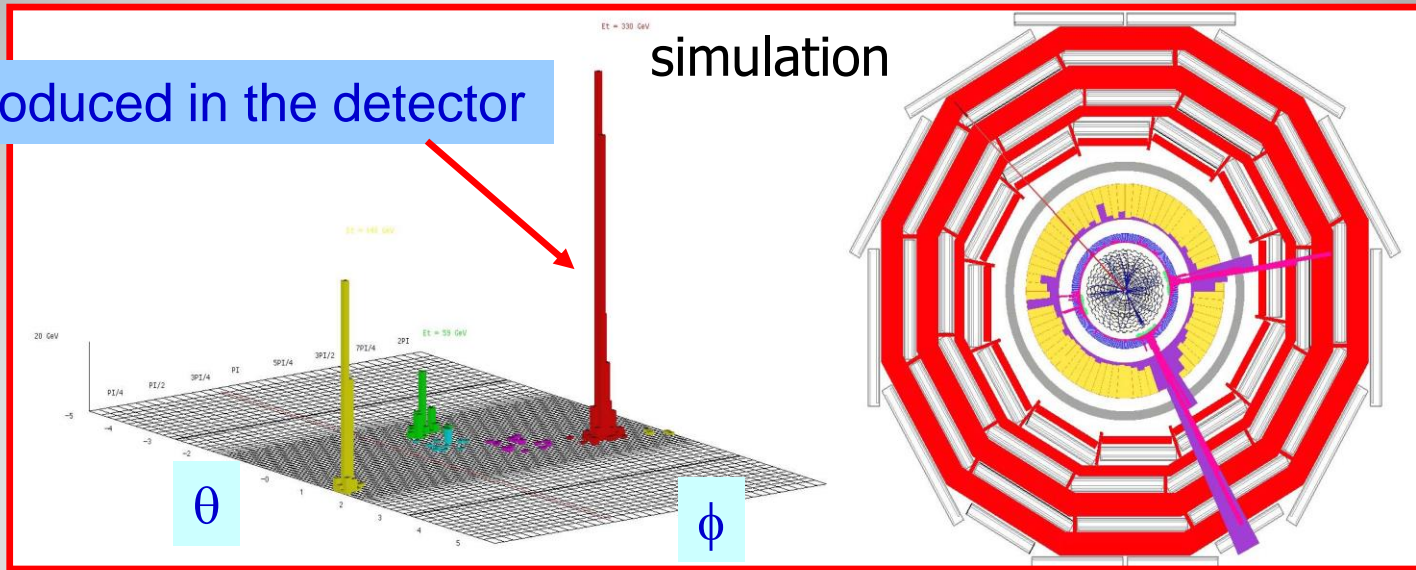
◆ Gauge unification



◆ Dark matter candidate with the right abundance

Detecting Supersymmetric Particles

Energy produced in the detector

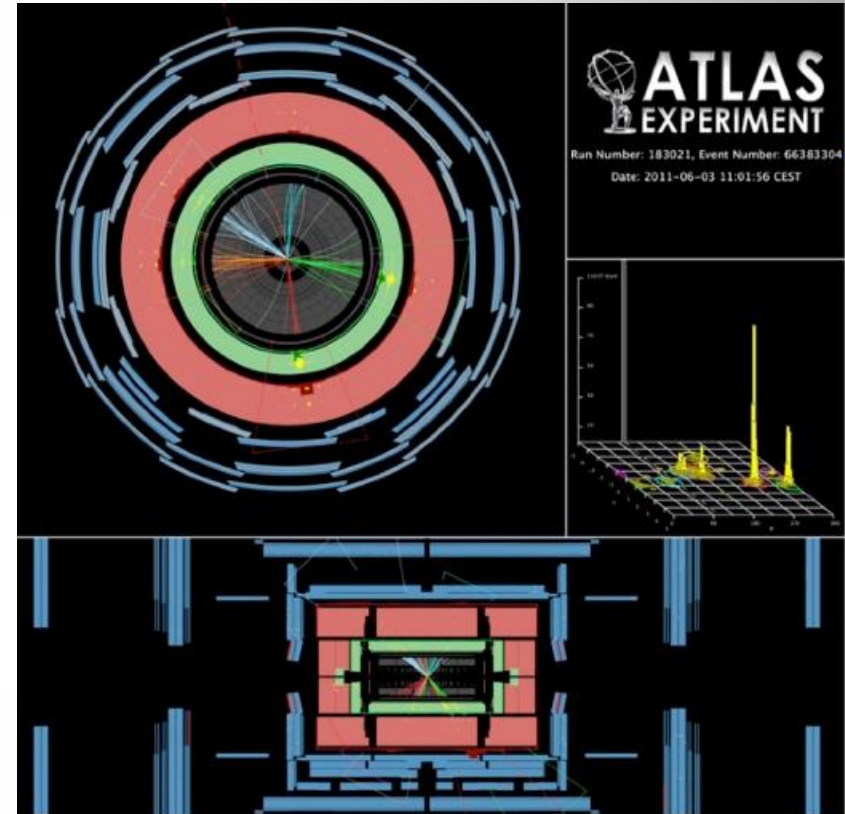
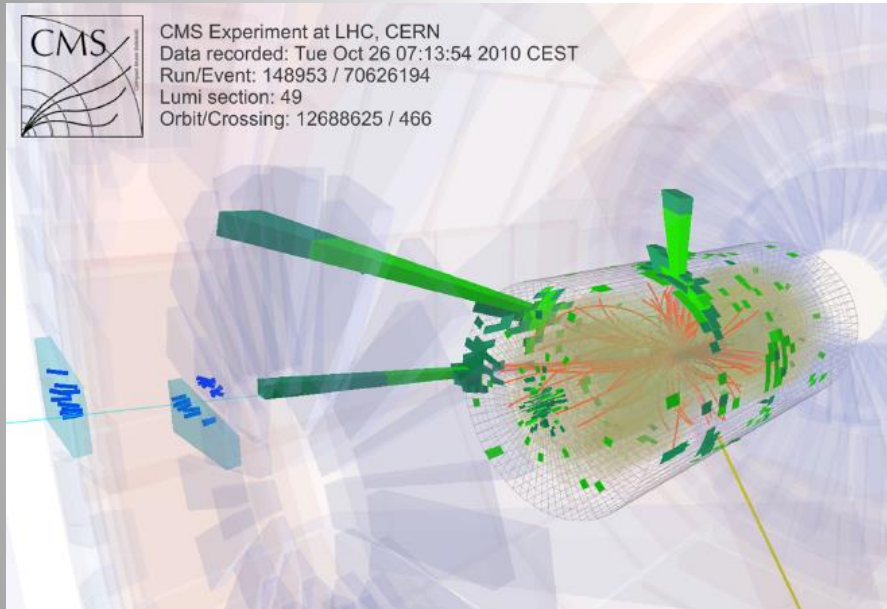


Supersymmetric particles decay and produce a cascade of jets, leptons and missing transverse energy (MET) due to escaping 'dark matter' particle candidates

Very prominent signatures in CMS and ATLAS

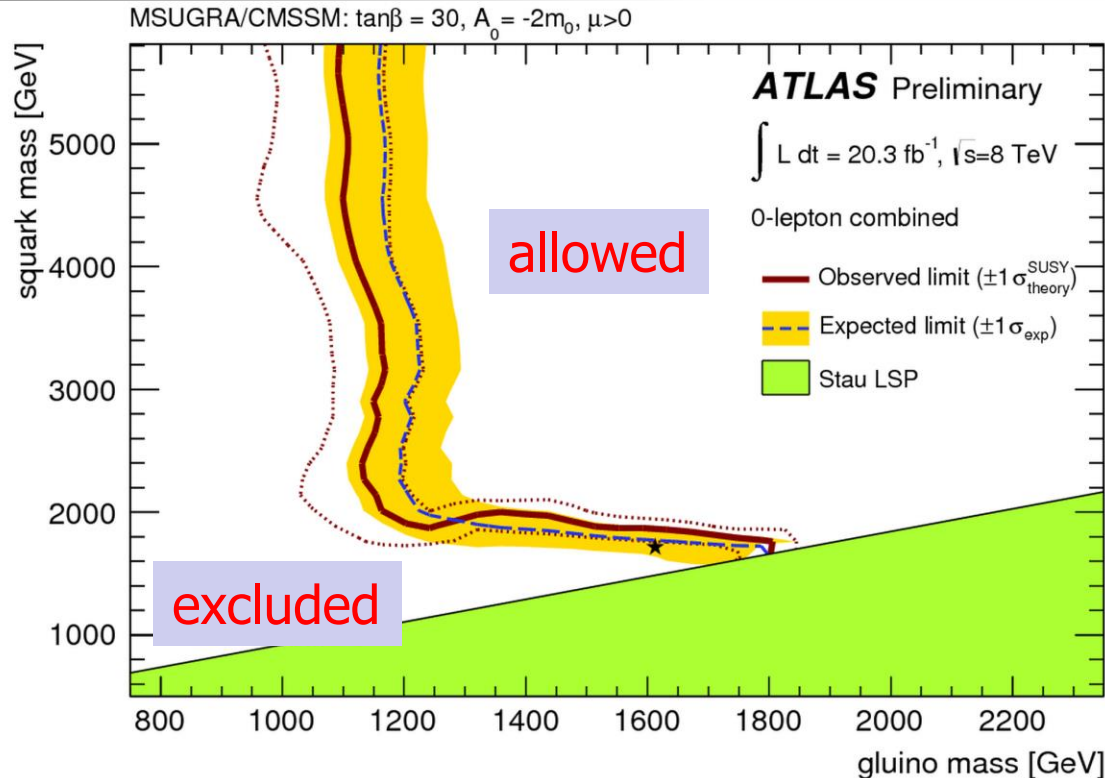
...Some Interesting Collisions...

...already in 2010...



- Events with five jets of particles **and large missing energy** which could come from a possible dark matter particle
- But a few events is not enough to prove we have something new
No visible excess has been building up with time...

SUSY Searches: No signal yet to date...



- So far **NO** clear signal of supersymmetric particles has been found

- We can exclude regions where the new particles could exist.

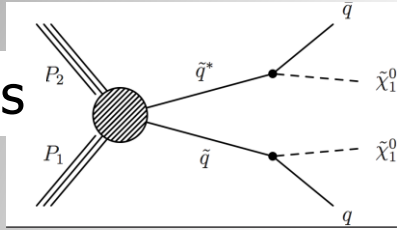
- Searches will continue for the **higher energy in 2015**

Plenty of searches ongoing: with jets, leptons, photons, W/Z, top, Higgs, with and without large missing transverse energy
Also special searches for more contrived model regions

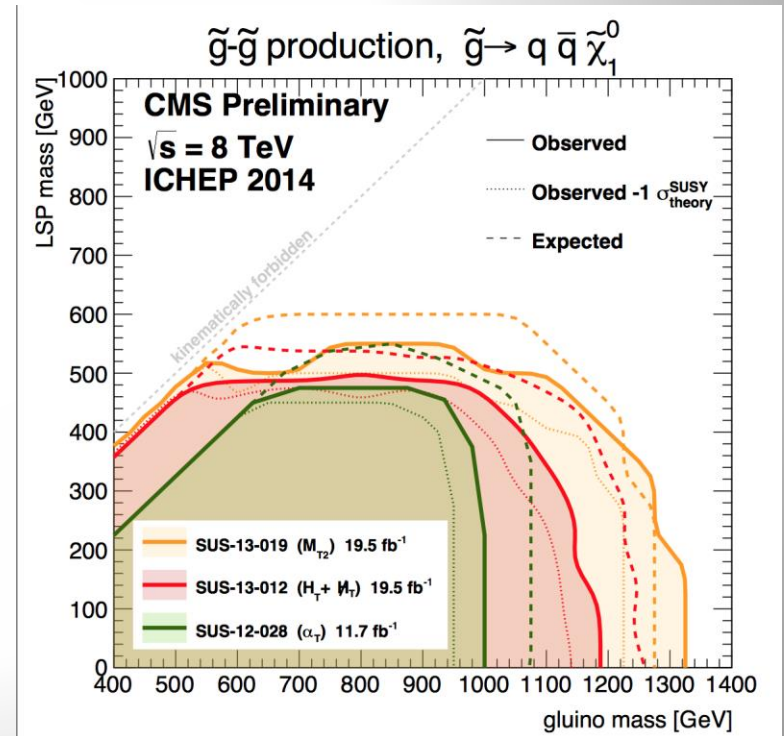
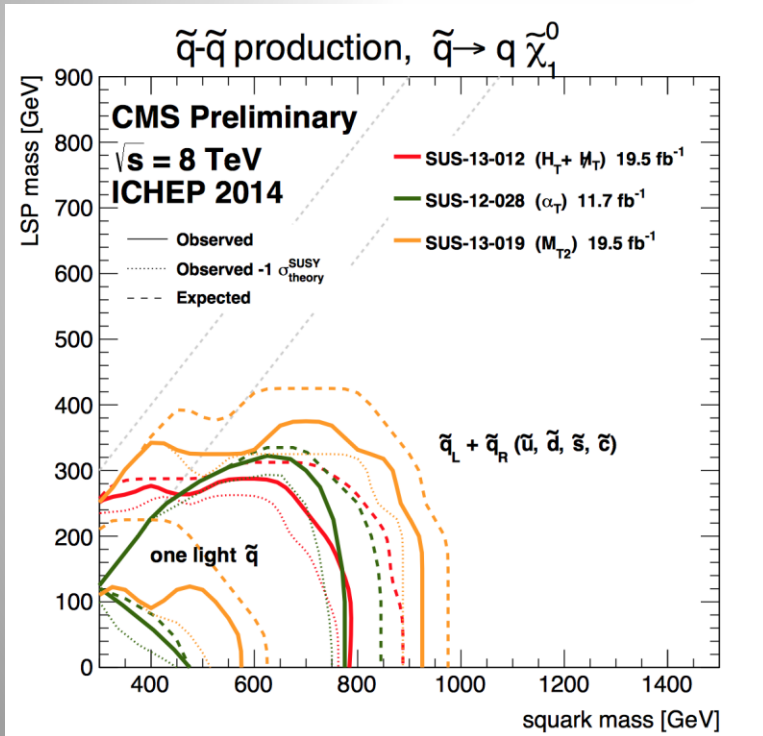
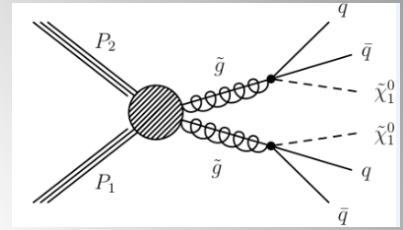
Limits on Squarks and Gluinos

Results depend on the topologies studies, assumed mass of the LSP etc.

Examples



Popular presentation of data:
Simplified ModelS (SMS)



Combined limits typically $> 1\text{-}1.3 \text{ TeV}$ on sparticle masses

What is really needed from SUSY?

End 2011: Revision!

N. Arkani-Ahmed
CERN Nov 2011

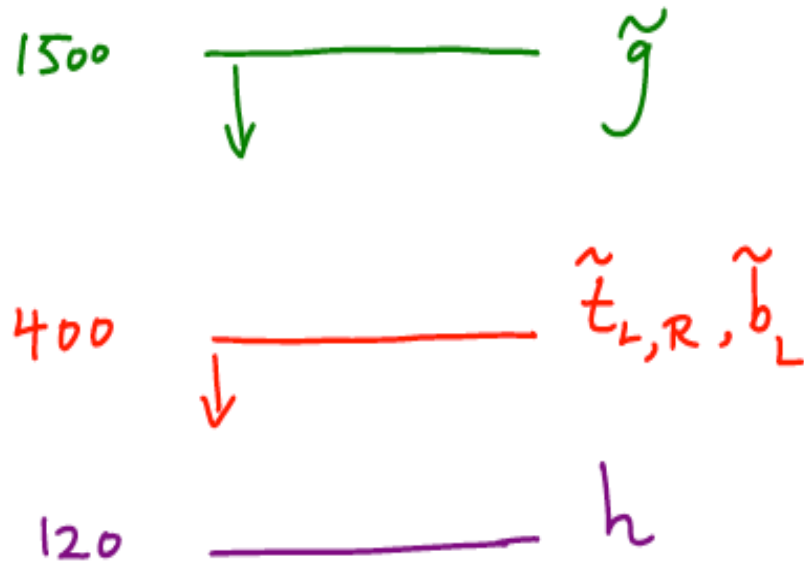
Papucci, Ruderman,
Weiler arXiv:1110.6926

LHC data end 2011
Stops > 200-300 GeV
Glino > 600-800 GeV

Moving away from
constrained SUSY models
to 'natural' models

Natural SUSY survived
LHC so far, but we
are getting close to
push it to its limits!

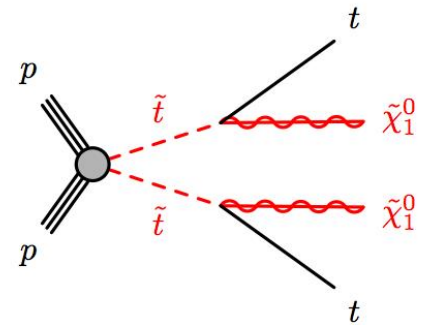
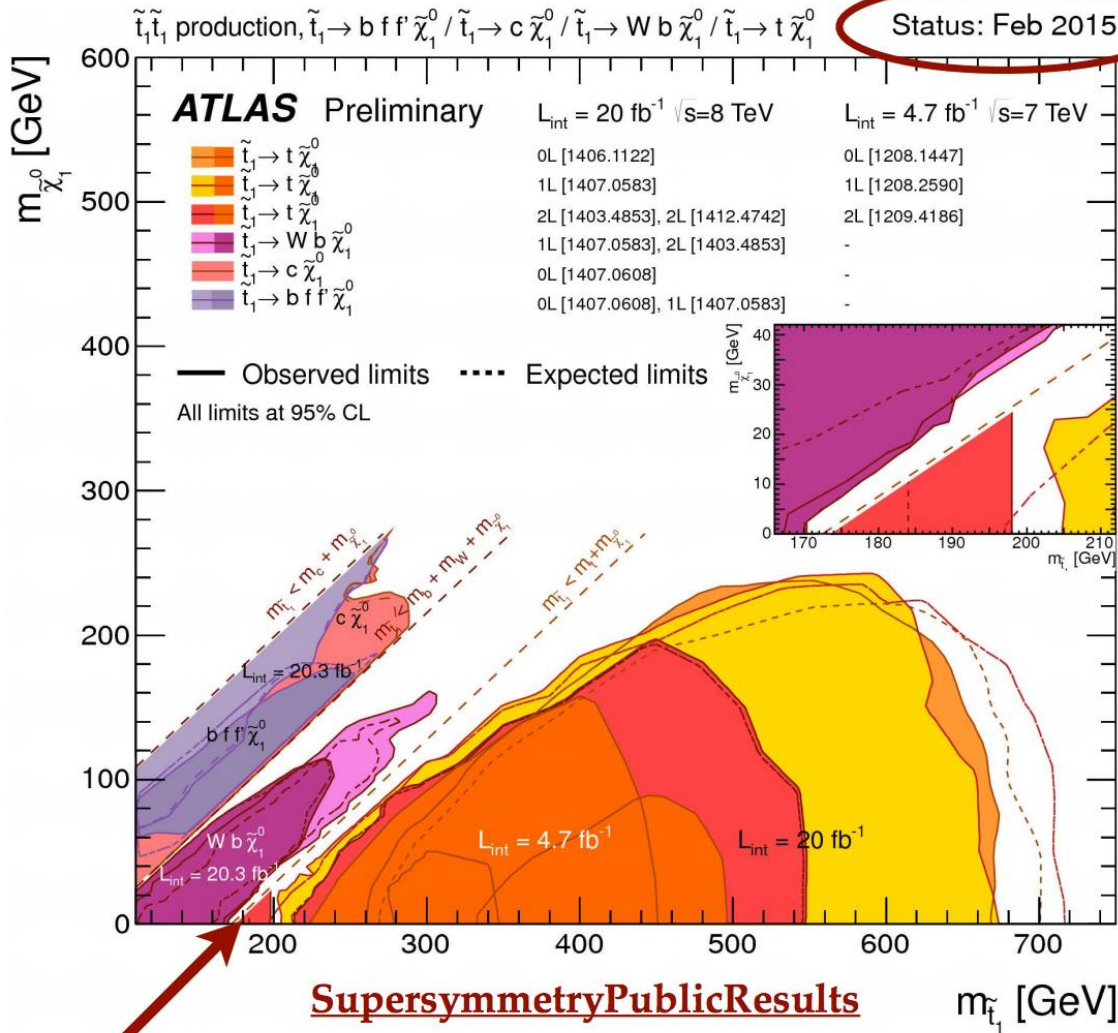
Compulsory Natural SUSY



Unavoidable tunings: $\left(\frac{400}{m_t}\right)^2$, $\left(\frac{4m_t}{M_{\tilde{g}}}\right)^2$

Natural SUSY?

Direct Stop Searches



- Stealth stop region ($m_{\tilde{t}} \sim m_t$) nearly closed by precision $t\bar{t}$ measurements!

Summary of SUSY Searches

In short: no sign of SUSY with the data collected so far (similar for CMS)

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: Feb 2015

ATLAS Preliminary

$\sqrt{s} = 7, 8 \text{ TeV}$

Reference

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} d\mathcal{I} [\text{fb}^{-1}]$	Mass limit			
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g}	1.7 TeV	$m(\tilde{q})=m(\tilde{g})$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q}	850 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \tilde{q})=m(2^{\text{nd}} \text{ gen. } \tilde{q})$
	$\tilde{q}\tilde{q}\gamma, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed)	1 γ	0-1 jet	Yes	20.3	\tilde{q}	250 GeV	$m(\tilde{q})-m(\tilde{\chi}_1^0) = m(c)$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g}	1.33 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20	\tilde{g}	1.2 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}, m(\tilde{\chi}^{\pm}) = 0.5(m(\tilde{\chi}_1^0) + m(\tilde{g}))$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20	\tilde{g}	1.32 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	GMSB ($\tilde{\ell}$ NLSP)	1-2 $\tau + 0-1 \ell$	0-2 jets	Yes	20.3	\tilde{g}	1.6 TeV	$\tan\beta > 20$
	GGM (bino NLSP)	2 γ	-	Yes	20.3	\tilde{g}	1.28 TeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$
	GGM (wino NLSP)	1 $e, \mu + \gamma$	-	Yes	4.8	\tilde{g}	619 GeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$
	GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g}	900 GeV	$m(\tilde{\chi}_1^0) > 220 \text{ GeV}$
3 rd gen. squarks direct production	GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	\tilde{g}	690 GeV	$m(\text{NLSP}) > 200 \text{ GeV}$
	Gravitino LSP	0	mono-jet	Yes	20.3	$P^{1/2}$ scale	865 GeV	$m(\tilde{G}) > 1.8 \times 10^{-4} \text{ eV}, m(\tilde{g})=m(\tilde{q})=1.5 \text{ TeV}$
	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g}	1.25 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g}	1.1 TeV	$m(\tilde{\chi}_1^0) < 350 \text{ GeV}$
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	1.34 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$
	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^+$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	1.3 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1	100-620 GeV	$m(\tilde{\chi}_1^0) < 90 \text{ GeV}$
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^+$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{b}_1	275-440 GeV	$m(\tilde{\chi}_1^0) = 2 m(\tilde{\chi}_1^+)$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^+$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1	110-167 GeV	$m(\tilde{\chi}_1^0) = 2m(\tilde{\chi}_1^+), m(\tilde{\chi}_1^0) = 55 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	\tilde{t}_1	90-191 GeV	$m(\tilde{\chi}_1^0) = 1 \text{ GeV}$
EW direct	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	1-2 b	Yes	20	\tilde{t}_1	210-640 GeV	$m(\tilde{\chi}_1^0) = 1 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1	90-240 GeV	$m(\tilde{\chi}_1^0) = 1 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_1	150-580 GeV	$m(\tilde{\chi}_1^0) = 1 \text{ GeV}$
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_2	290-600 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$	90-325 GeV	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \tilde{\ell}\nu(\tilde{\ell}\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^+$	140-465 GeV	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}, m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^+) + m(\tilde{\chi}_1^-))$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \tilde{\tau}\nu(\tilde{\tau}\bar{\nu})$	2 τ	0	Yes	20.3	$\tilde{\chi}_1^+$	100-350 GeV	$m(\tilde{\chi}_1^0) = 0 \text{ GeV}, m(\tilde{\tau}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^+) + m(\tilde{\chi}_1^-))$
	$\tilde{\chi}_1^+\tilde{\chi}_1^0 \rightarrow \tilde{\ell}\nu\tilde{\chi}_1^0(\tilde{\ell}\bar{\nu}), \tilde{\ell}\tilde{\nu}\tilde{\chi}_1^0(\tilde{\ell}\bar{\nu})$	3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^+, \tilde{\chi}_1^0$	700 GeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_1^+), m(\tilde{\chi}_1^0) = 0, m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^+) + m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^+\tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$	2-3 e, μ	0-2 jets	Yes	20.3	$\tilde{\chi}_1^+, \tilde{\chi}_1^0$	420 GeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_1^+), m(\tilde{\chi}_1^0) = 0, \text{ sleptons decoupled}$
	$\tilde{\chi}_1^+\tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0, h \rightarrow b\bar{b}/W\tilde{\chi}_1^0/\tau\tau/\gamma\gamma$	e, μ, γ	0-2 b	Yes	20.3	$\tilde{\chi}_1^+, \tilde{\chi}_1^0$	250 GeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_1^+), m(\tilde{\chi}_1^0) = 0, \text{ sleptons decoupled}$
Long-lived particles	$\tilde{\chi}_1^0\tilde{\chi}_1^0, \tilde{\chi}_1^0\tilde{\chi}_1^0 \rightarrow \tilde{\ell}R\ell$	4 e, μ	0	Yes	20.3	$\tilde{\chi}_1^0, \tilde{\chi}_1^0$	620 GeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0) = 0, m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^0) + m(\tilde{\chi}_1^0))$
	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^+$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^+$	270 GeV	$m(\tilde{\chi}_1^+) - m(\tilde{\chi}_1^-) = 160 \text{ MeV}, \tau(\tilde{\chi}_1^+) = 0.2 \text{ ns}$
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	27.9	\tilde{g}	832 GeV	$m(\tilde{\chi}_1^0) = 100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$
	Stable \tilde{g} R-hadron	trk	-	-	19.1	\tilde{g}	1.27 TeV	
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{\ell}, \bar{\mu}) + \tau(e, \mu)$	1-2 μ	-	-	19.1	$\tilde{\chi}_1^0$	537 GeV	$10 < \tan\beta < 50$
	GMSB, $\tilde{\chi}_1^0 \rightarrow \text{RPV}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$	435 GeV	$2 < \tau(\tilde{\chi}_1^0) < 3 \text{ ns}, \text{SPS8 model}$
	$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow q\tilde{q}$ (RPV)	1 μ , displ. vtx	-	-	20.3	\tilde{q}	1.0 TeV	$1.5 < \tau < 156 \text{ mm}, \text{BR}(\mu) = 1, m(\tilde{\chi}_1^0) = 108 \text{ GeV}$
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$	1.61 TeV	$\lambda'_{311} = 0.10, \lambda_{132} = 0.05$
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 $e, \mu + \tau$	-	-	4.6	$\tilde{\nu}_\tau$	1.1 TeV	$\lambda'_{311} = 0.10, \lambda_{1233} = 0.05$
	RPV	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g}	1.35 TeV
$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_\tau$		4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^+, \tilde{\chi}_1^0$	750 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^+), \lambda_{121} \neq 0$
$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$		3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^+, \tilde{\chi}_1^0$	450 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^+), \lambda_{133} \neq 0$
$\tilde{g} \rightarrow q\tilde{q}q$		0	6-7 jets	-	20.3	\tilde{g}	916 GeV	$\text{BR}(\mu) = \text{BR}(b) = \text{BR}(c) = 0\%$
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow b\tilde{s}$		2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{g}	850 GeV	
Other		Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	2 c	Yes	20.3	\tilde{c}	490 GeV



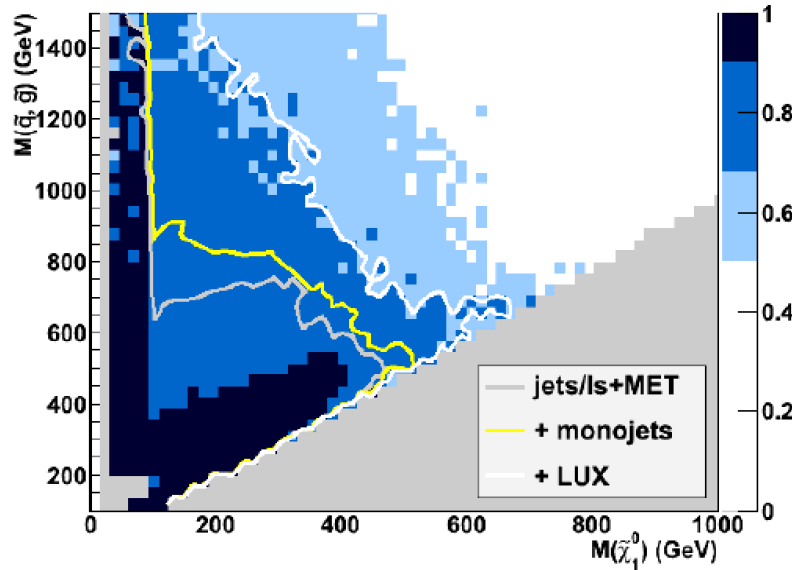
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$\sqrt{s} = 7 \text{ TeV}$ full data
 $\sqrt{s} = 8 \text{ TeV}$ partial data
 $\sqrt{s} = 8 \text{ TeV}$ full data

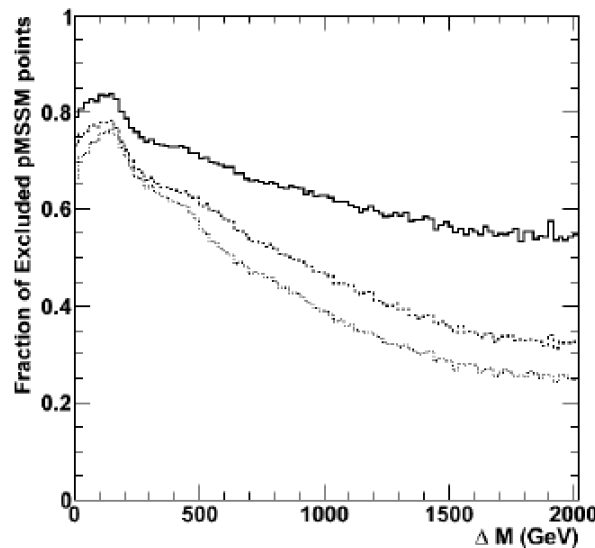
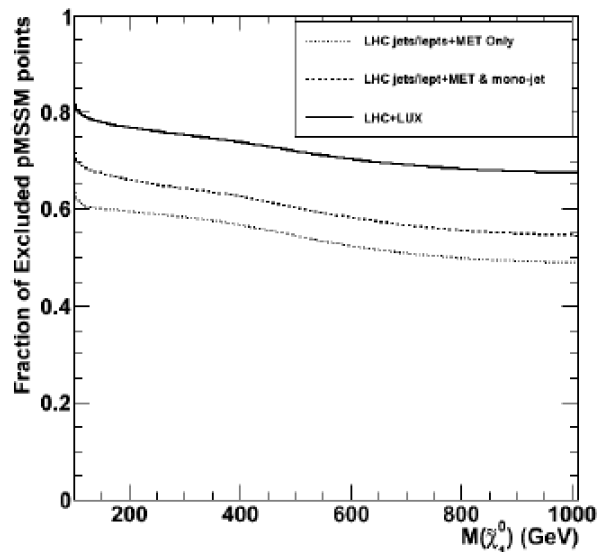
10⁻¹ 1 Mass scale [TeV]

Dark Matter SUSY Space Left?

arXiv:1311.7641



- Use the pMSSM SUSY model (19 parameters)
- Use all the ATLAS SUSY Data + mono-jet searches + LUX DM results
- Check what fraction of pMSSM solutions that is excluded.

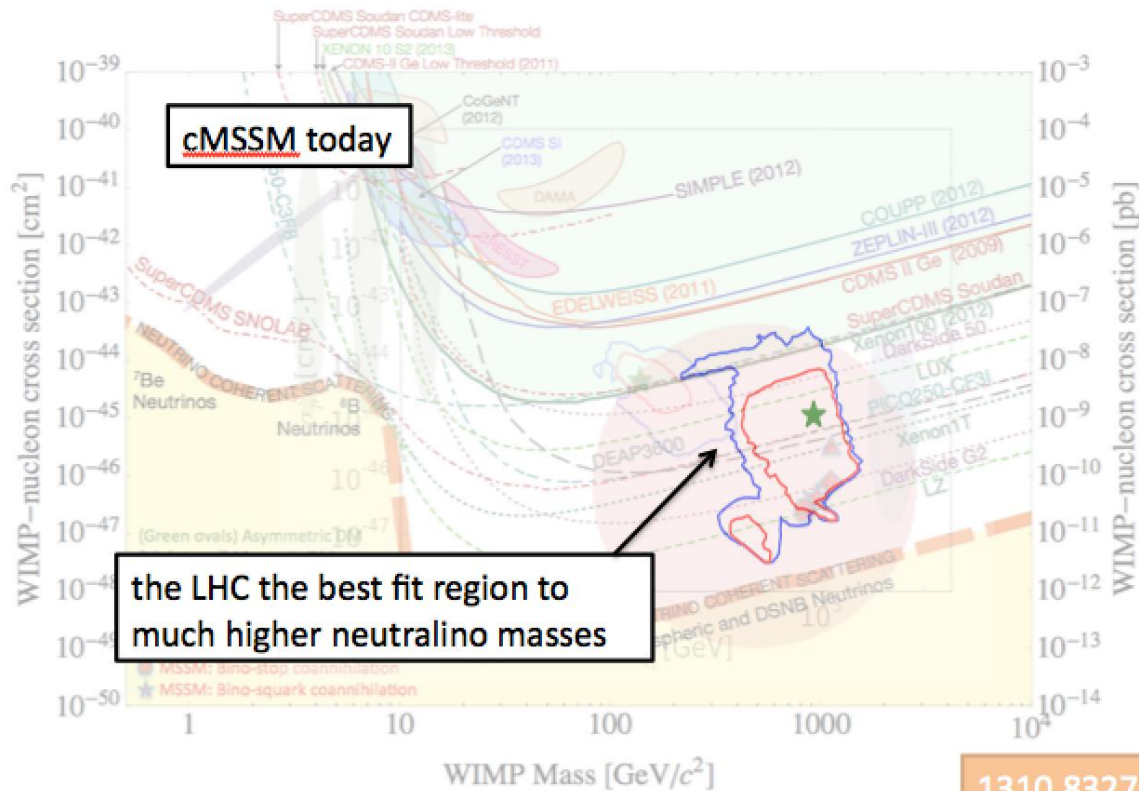


60-80% of the solutions excluded as function of the neutralino mass

Constrained SUSY Models: Now

O. Buchmuller, ... ADR, Ellis... et al: arXiv:13125250

Study of the allowed DM space in CMSSM model



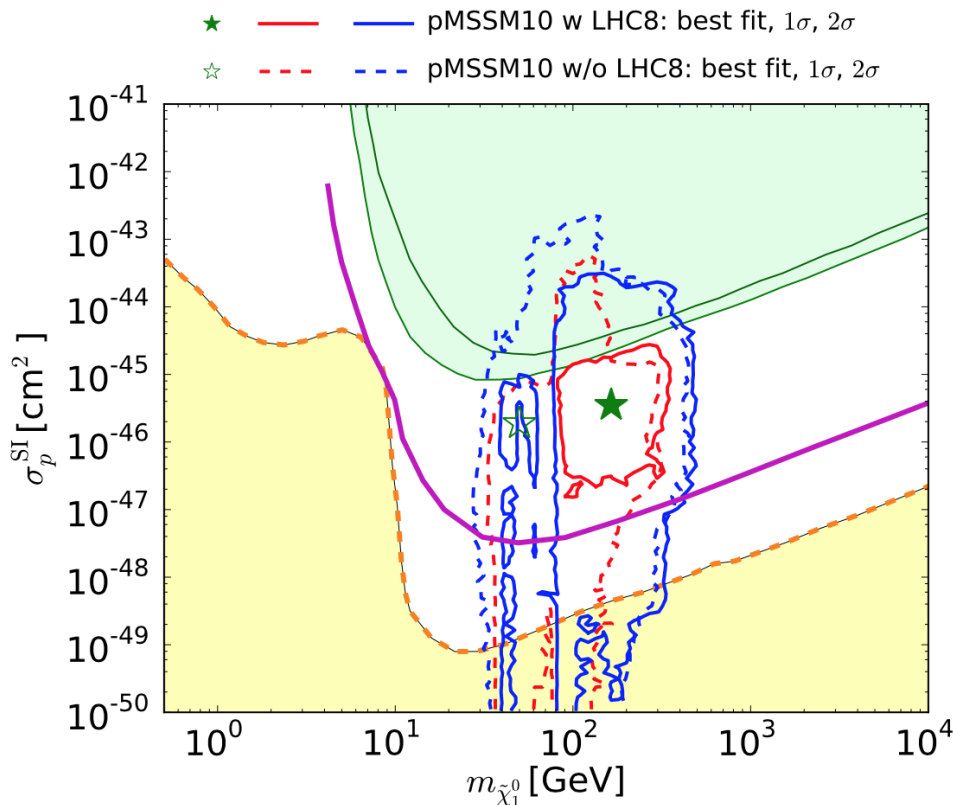
Region allowed in the CMSSM includes constraints of the Run-I LHC searches (SUSY) and precision data, $g-2$, cold dark matter constraints...

Available SUSY Space: Update

arXiv:1504.03260

The pMSSM10 after LHC Run 1

K.J. de Vries^a, E.A. Bagnaschi^b, O. Buchmueller^a, R. Cavanaugh^{c,d}, M. Citron^a,
 A. De Roeck^{e,f}, M.J. Dolan^g, J.R. Ellis^{h,e}, H. Flächerⁱ, S. Heinemeyer^j, G. Isidori^k,
 S. Malik^a, J. Marrouche^e, D. Martínez Santos^l, K.A. Olive^m, K. Sakurai^h, G. Weiglein^b



Region allowed in the pMSSM10
 (10 parameter SUSY space)
 includes constraints of the Run-I
 LHC searches (SUSY) and precision
 data, g-2, dark matter constraints...

$$\begin{aligned}
 & 3 \text{ gaugino masses : } M_{1,2,3}, \\
 & 2 \text{ squark masses : } m_{\bar{q}_1} = m_{\bar{q}_2} \neq m_{\bar{q}_3}, \\
 & 1 \text{ slepton mass : } m_{\bar{\ell}}, \\
 & 1 \text{ trilinear coupling : } A, \\
 & \text{Higgs mixing parameter : } \mu, \\
 & \text{Pseudoscalar Higgs mass : } M_A, \\
 & \text{Ratio of vevs : } \tan \beta.
 \end{aligned} \tag{1}$$

Once and a while it hits the blogosphere!

nature International weekly journal of science

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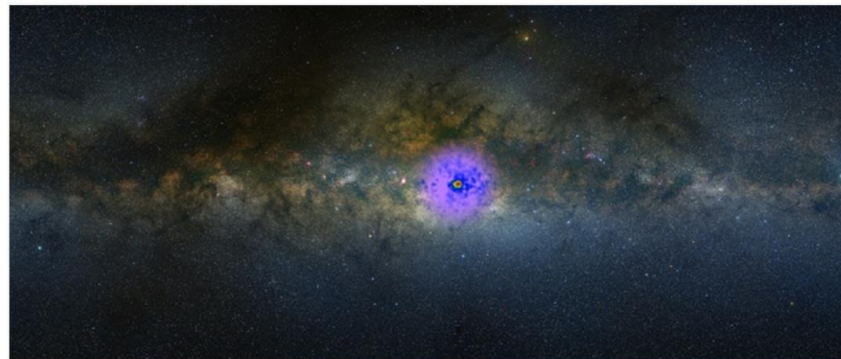
Mysterious galactic signal points LHC to dark matter

High-energy particles at centre of Milky Way now within scope of Large Hadron Collider.

Davide Castelvecchi

05 May 2015

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A. Mellinger, CMU; T. Linden, Univ. of Chicago/NASA G

γ -rays (shown in false colour) emitted from the Galactic Centre are giving the LHC a firm target in its hunt for dark matter.

A description of the Galactic Center excess in the Minimal Supersymmetric Standard Model

Abraham Achterberg,^a Simone Amoroso,^e Sascha Caron,^{a,b} Luc Hendriks,^a Roberto Ruiz de Austri,^c Christoph Weniger^d

arXiv:1502.05703

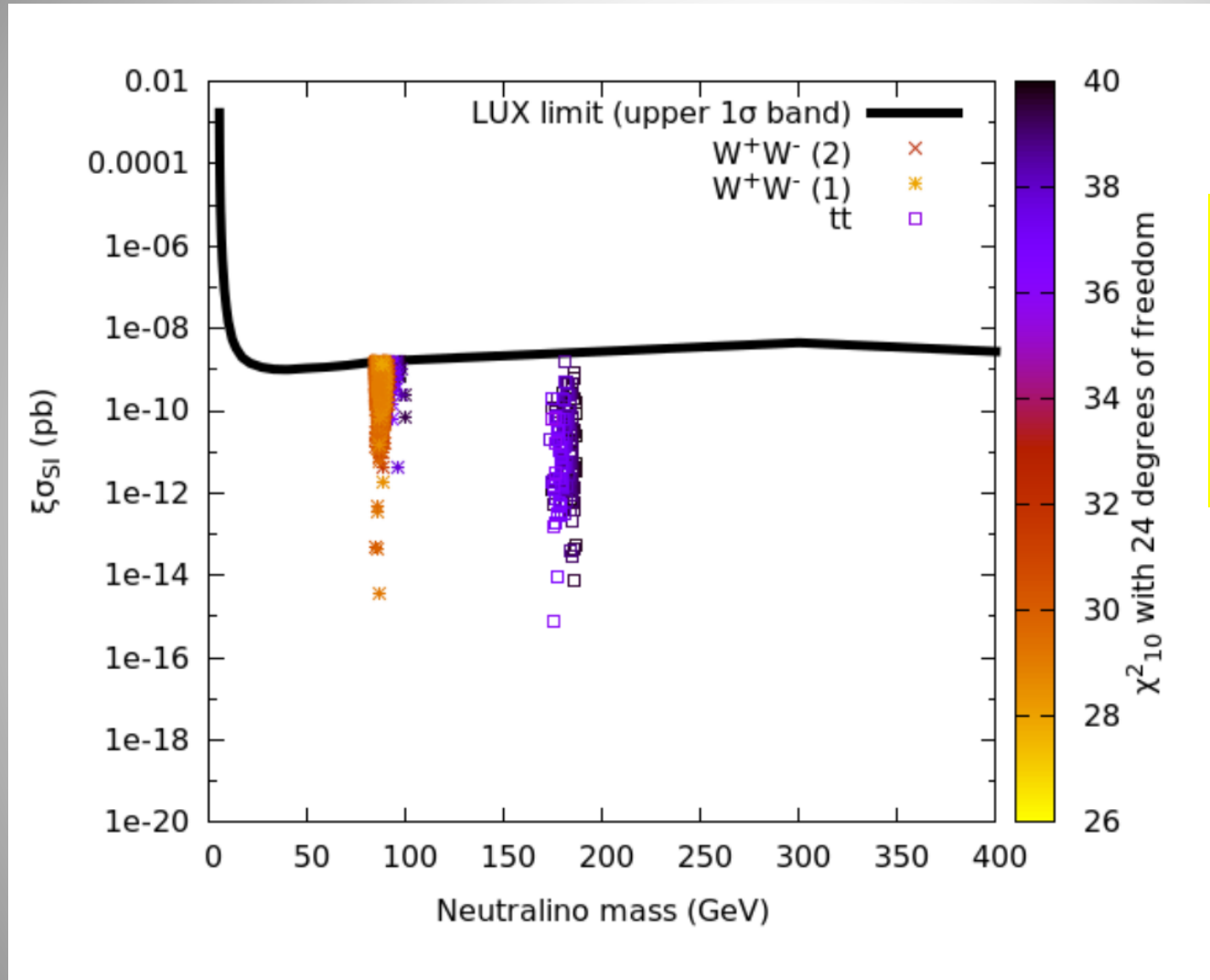
It is one of the most disputed observations in physics. But an explanation may be in sight for a mysterious excess of high-energy photons at the centre of the Milky Way. The latest analysis suggests that the signal could come from a dark-matter particle that has just the right mass to show up at the world's largest particle accelerator.

The Large Hadron Collider (LHC), housed at the CERN particle-physics laboratory near Geneva, Switzerland, is due to restart colliding protons this summer after a two-year hiatus (see '[LHC 2.0: A new view of the Universe](#)'). Physicists there have told *Nature* that they now plan to make the search for such a particle a top target for the collider's second run.

Resolution may be just around the corner. In addition to being produced at the LHC, the neutralino could also be within the shooting range of next-generation underground experiments that are trying to catch dark-matter particles that happen to fly through Earth, says physicist Albert De Roeck, who works on the CMS, one of the two LHC detectors that will hunt for dark matter. If such a particle is indeed the cause of the γ -rays, he says, "it seems that the dark-matter signals should be observed very soon now".

Search for Models compatible with the FERMI-LAT excess and LHC limits

Where Dark Matter Should Be Found...



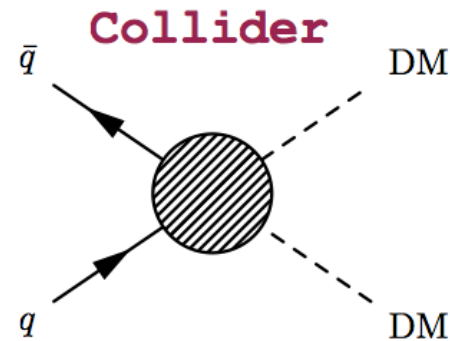
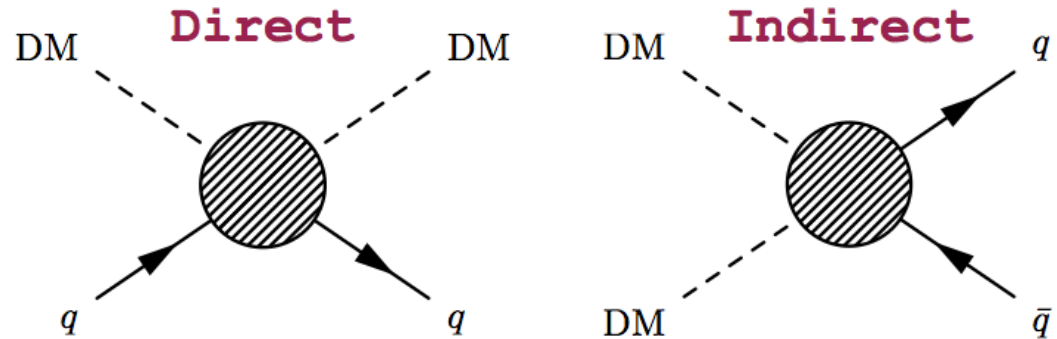
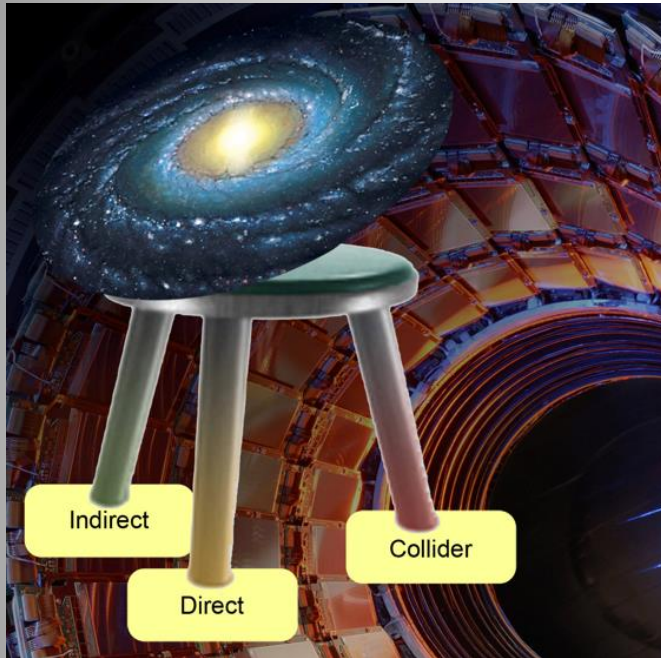
pMSSM solutions:
FERMI-LAT excess
Due to $DMDM \rightarrow WW$
or $DMDM \rightarrow t\text{top}$
annihilation

Who will get there first? The LHC? Direct Searches? To be continued...

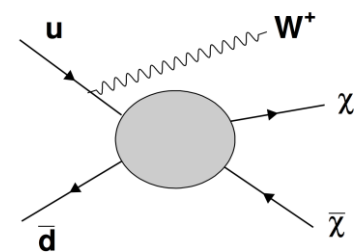
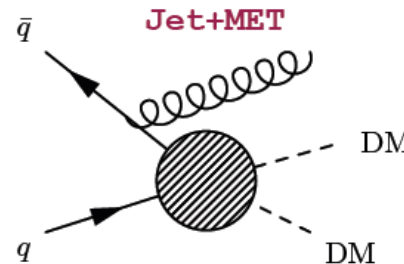
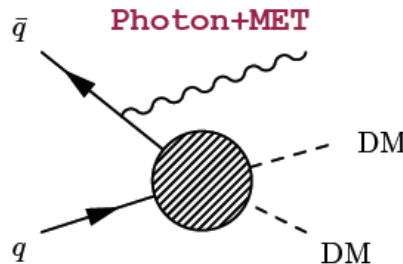
General Searches for Dark Matter

The Generic Dark Matter Connection

Searches for mono-jets and mono-photons can be used to search for Dark Matter (DM)



Use effective theory or better simplified models to relate measurements to Dark Matter studies



Dark Matter Interpretation

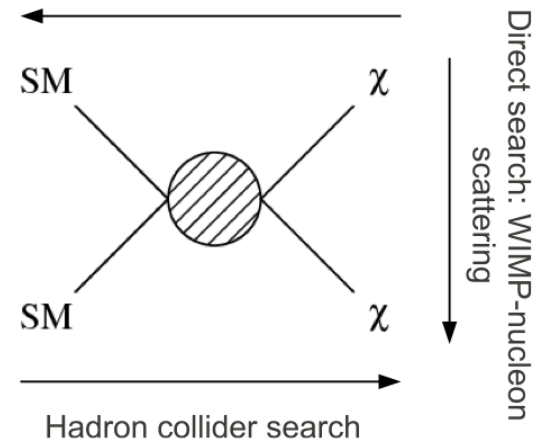
Convert experimental search results into limits on DM quantities

Two ways:

- Effective field theory (EFT):
 - Mediator too heavy to be generated directly
 - Contact interaction with suppression scale $M_\star \sim \frac{M}{\sqrt{g_\chi g_{SM}}}$, with g_χ and g_{SM} the couplings to Standard Model (SM) and DM, and M the mediator mass
- Simplified models: Popular in Supersymmetry
 - Specified massive mediator
 - UV-complete (no validity issue)

Types of interactions chosen in the studies @ LHC

Indirect search: WIMPs annihilation



Name	Initial state	Type	Operator
D1	qq	scalar	$\frac{m_q}{M_\star^3} \bar{\chi} \chi \bar{q} q$
D5	qq	vector	$\frac{1}{M_\star^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$
D8	qq	axial-vector	$\frac{1}{M_\star^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$
D9	qq	tensor	$\frac{1}{M_\star^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$
D11	gg	scalar	$\frac{1}{4M_\star^3} \bar{\chi} \chi \alpha_s (G_{\mu\nu}^a)^2$

Discussions on the region of validity for the EFTs...

Data Interpretation EFTs

- Pair-production of χ characterised by a contact interaction with operators

$$\mathcal{O}_V = \frac{(\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu q)}{\Lambda^2}$$

vector --> spin independent (SI)

$$\mathcal{O}_{AV} = \frac{(\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu\gamma_5q)}{\Lambda^2}$$

axial-vector --> spin-dependent (SD)

- Cross section depends on the mass (m_χ) and the scale Λ (for couplings g_χ, g_q)

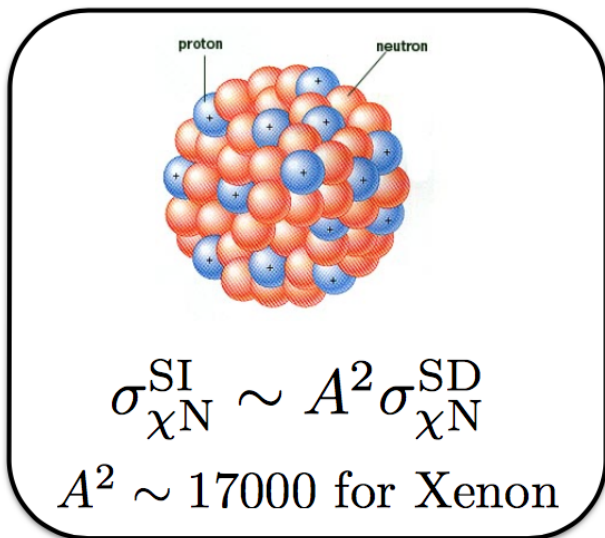
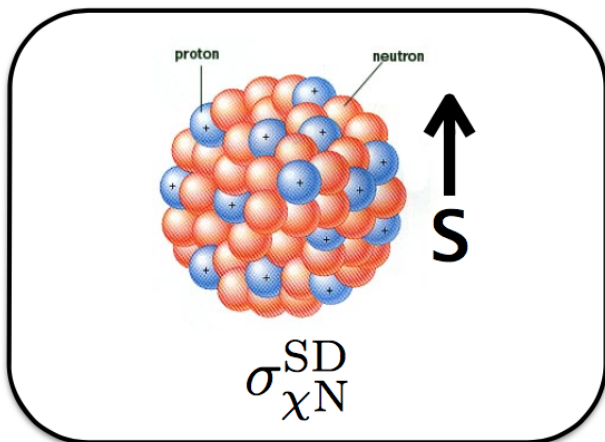
$$\sigma_{SI} = 9 \frac{\mu^2}{\pi\Lambda^4}$$
$$\sigma_{SD} = 0.33 \frac{\mu^2}{\pi\Lambda^4}$$

*spin-independent
and spin-dependent
cross sections*

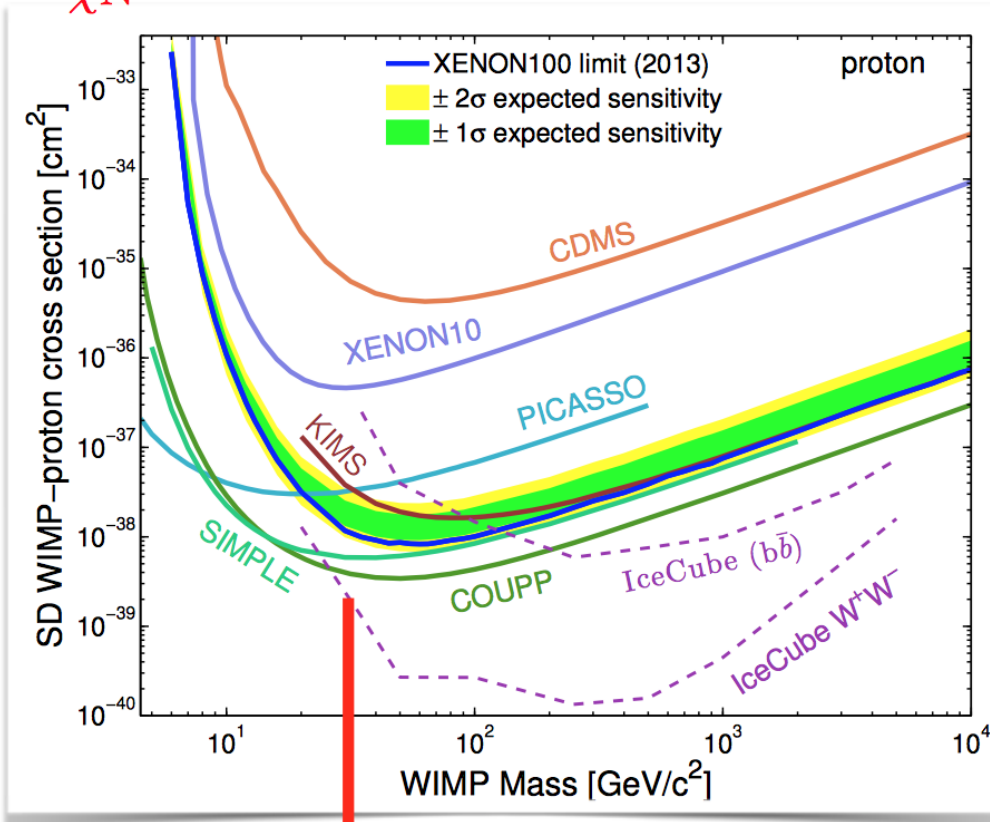
$$\Lambda = M/\sqrt{g_\chi g_q} \quad \mu = \frac{m_\chi m_p}{m_\chi + m_p}$$

[Bai, Fox and Harnik, JHEP 1012:048 (2010),
Goodman, Ibe, Rajaraman, Shepherd, Tait, Yu,
Phys.Rev.D82:116010 (2010), Beltran, Hooper,
Kolb, Krusberg, Tait, JHEP 1009:037 (2010)]

Spin Dependent and Independent σ



$\sigma_{\chi N}^{SD}$ limit

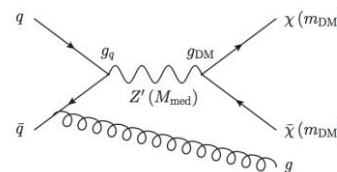


$\sigma_{\chi N}^{SI}$ limit is 10^5 times more stringent

Mono-object Searches in CMS

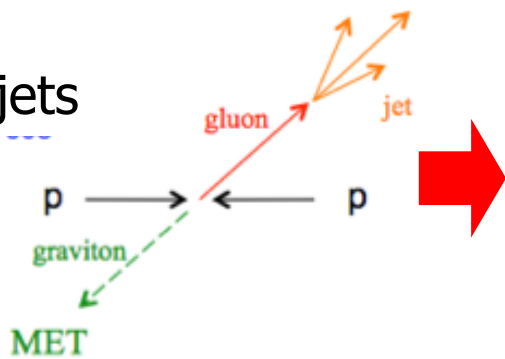
- **Mono-jets:** Generally the most powerful
- **Mono-photons:** First used for dark matter Searches
- **Mono-Ws:** Distinguish dark matter couplings to u- and d-type of quarks
- **Mono-Zs:** Clean signature
- **Mono-Tops:** Couplings to tops
- **Mono-Higgs:** Higgs-portals
- **Higgs Decays?**

Effective Field Theories for DM interpretation are under scrutiny!
 Alternatives such as SMS proposed

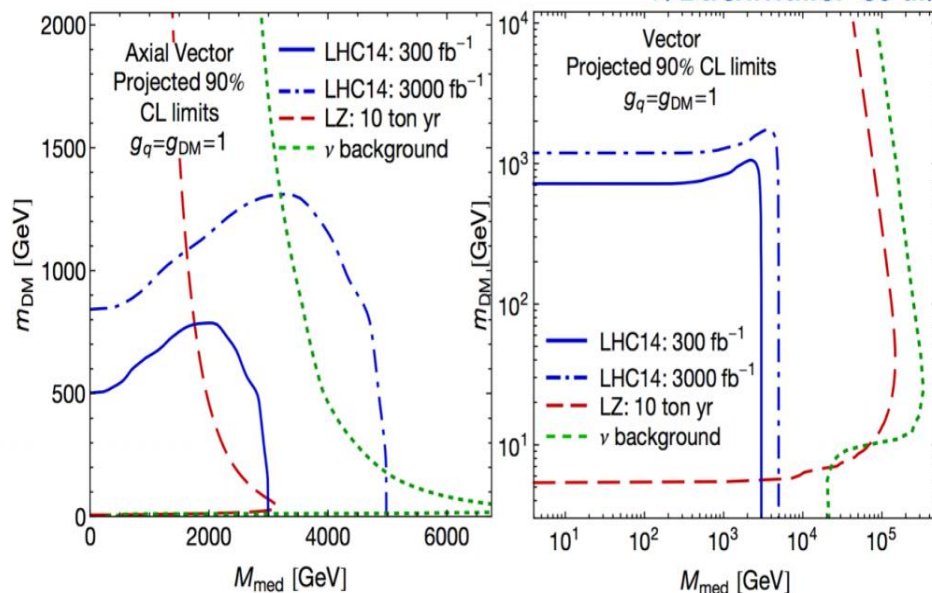


arXiv:1407.8257
 arXiv:1411.0535

Example Monojets



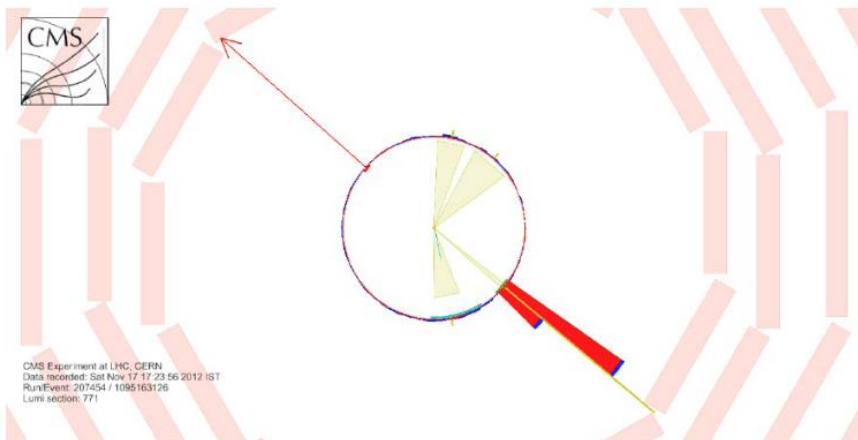
Dark Matter?



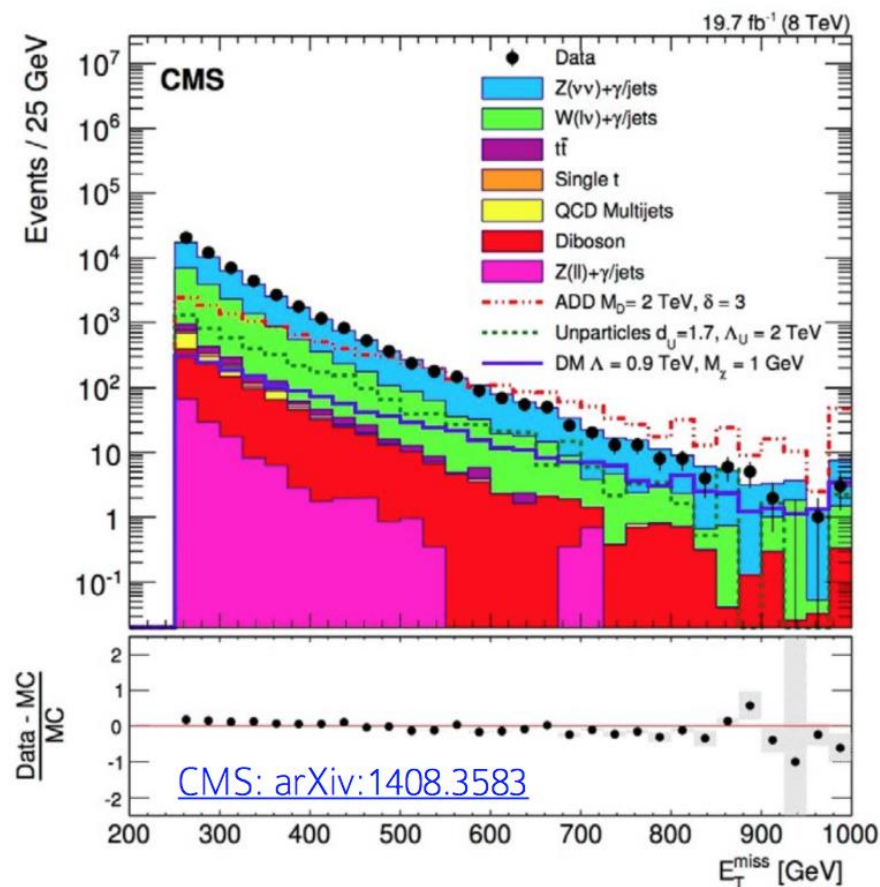
Mono-X Searches

Invisible DM particles **escape detection**:
LHC experiment strategy: tag events using recoiling object(s),
measure missing transverse momentum (Missing E_T)

[CMS: arXiv:1410.8812](https://arxiv.org/abs/1410.8812)

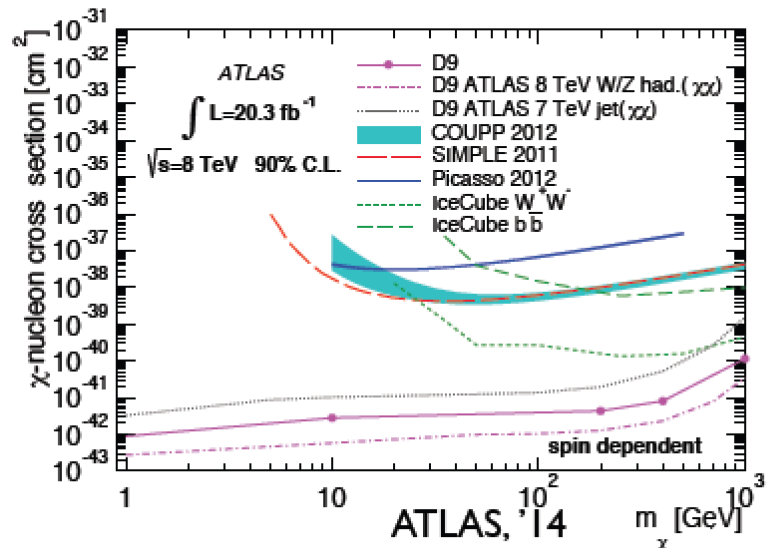
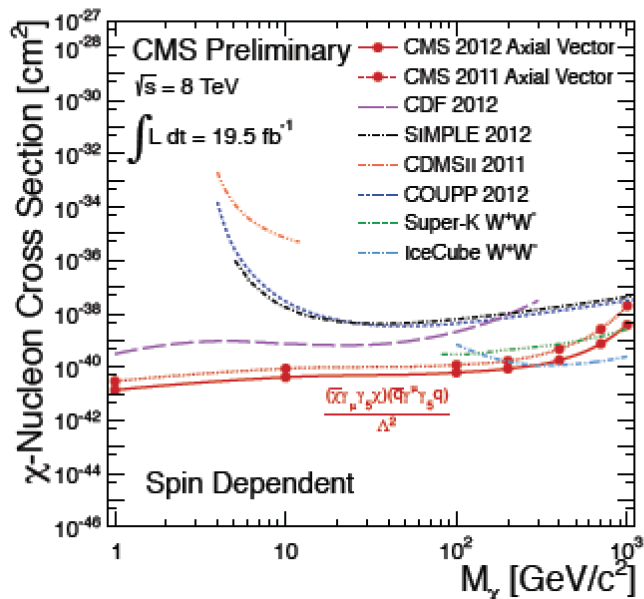
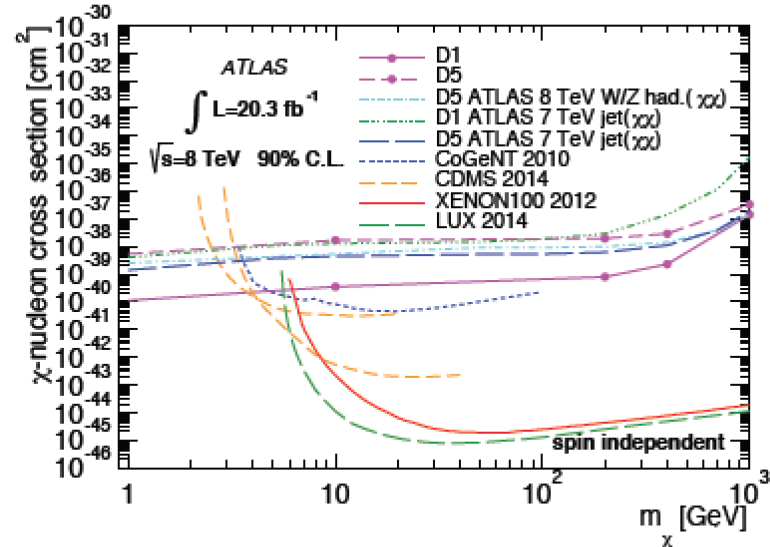
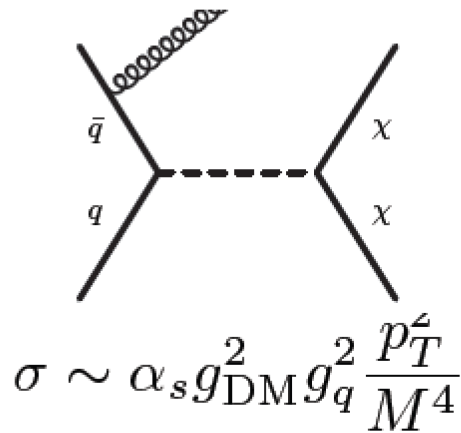


Dark Matter signature:
excess in tails of E_T distribution
(searches also sensitive to other models)



Some Recent Results (EFTs)

The cross section can be estimated to be



The Search for Mono-Higgs Next?

Mono-Higgs: a new collider probe of dark matter

Linda Carpenter,¹ Anthony DiFranzo,² Michael Mulhearn,³ Chase Shimmin,² Sean Tulin,⁴ and Daniel Whiteson²

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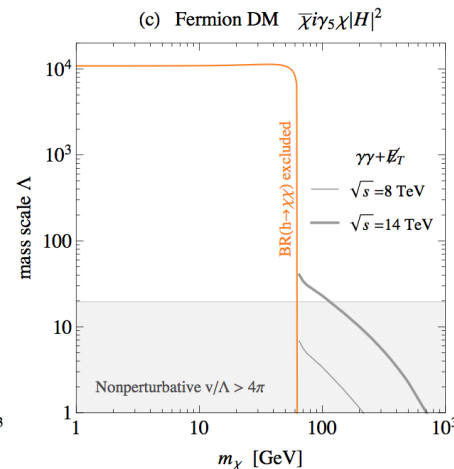
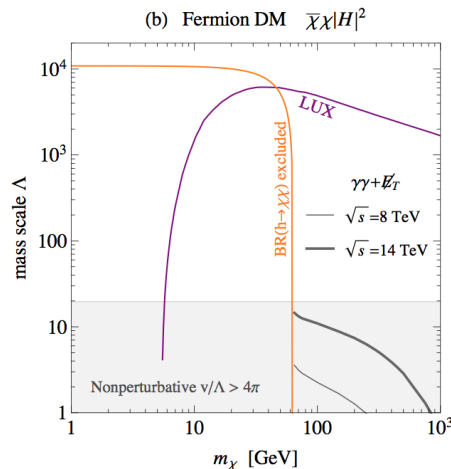
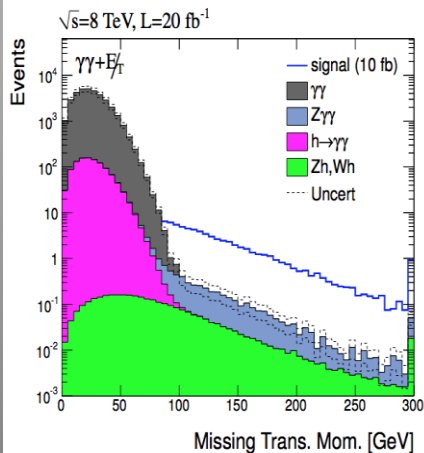
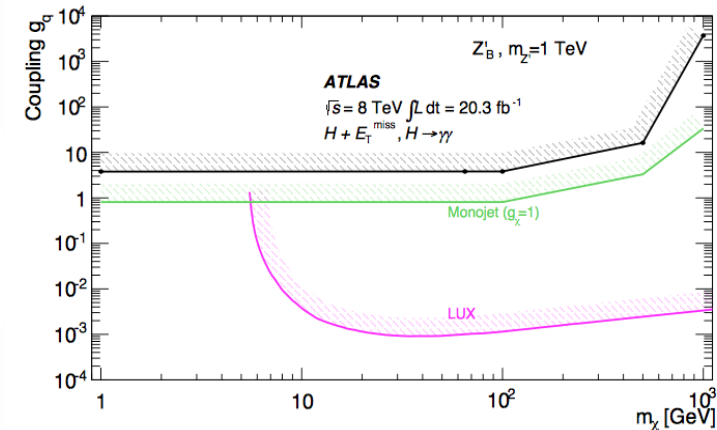
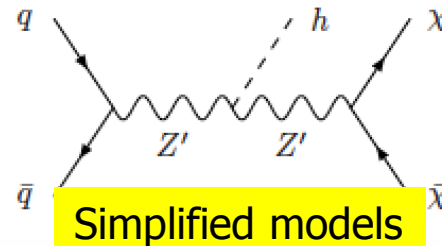
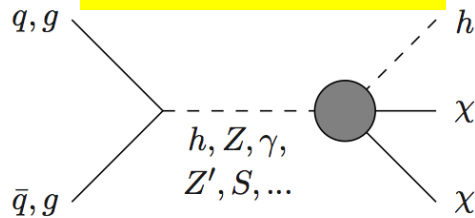
⁴Department of Physics and Astronomy, University of Michigan, MI

ATLAS

arXiv:1506.01081

The next mono-object to study?
No ISR, but coupling to the DM directly

Effective theories



Result:
Promising for
14 TeV?

2015: LHC-Wide Dark Matter Group

- **ATLAS-CMS Dark Matter Forum formed in Nov 2014**
- **Primary goal of forum : form consensus on the use of simplified models and EFTs to guide early Run-2 searches, by bringing together expertise on DM from experimental and theory communities.**

This forum will be followed up by a LPC with a new DM study group

LHC-Wide Dark Matter Group

- Final report of the DM forum, to be submitted to arXiv end of month
- 150 page document, detailing the recommended set of DM simplified models, presentation of EFT results, matrix element implementation etc
- Document reviewed by external reviewers
- ~150 signatories: CMS, ATLAS, theorists

Dark Matter Benchmark Models for Early LHC Run-2 Searches: Report of the ATLAS/CMS Dark Matter Forum

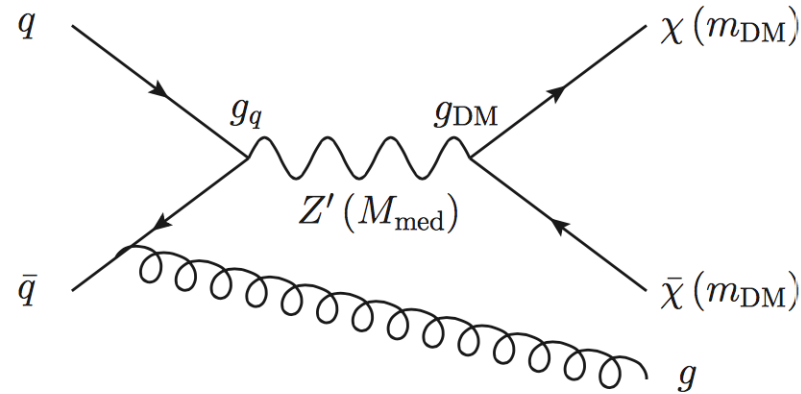
July 6, 2015

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Arely Cortes Gonzalez *IFAE Barcelona, Spain*

arXiv:1507.00966v1 [hep-ex] 3 Jul 2015

Using Simplified Models (SMS) as in SUSY

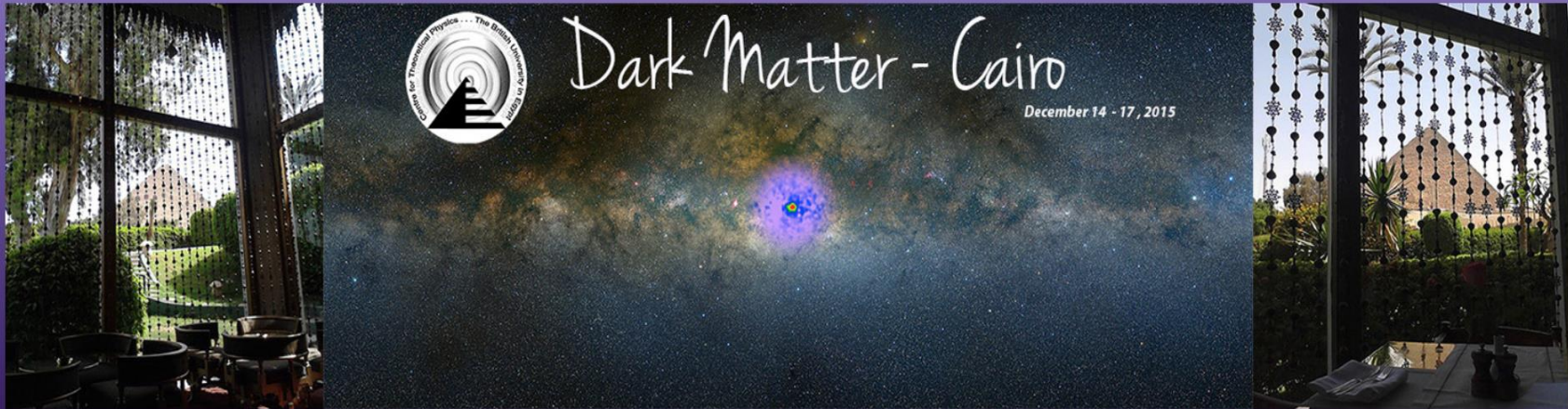
- ➔ Assume interaction exists between SM and DM
- ➔ DM is single particle, Dirac fermion
- ➔ Minimal Flavor Violation
- ➔ Assume minimal decay width for particle mediating interaction
- ➔ These assumptions only add a limited number of new particles and new interactions to the SM. Simplified models can be used as starting points to build more complete theories.



Factorised approach does not always lead to full theories but useful as a starting point giving a **distinct but complementary set of signatures to explore.**

More on Dark Matter in Egypt Soon

CTP@BUE



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CONTACT

ORGANISERS

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(Politecnico and INFN Bari)
- **Albert de Roeck**
(Antwerp/CERN)
- **Amr El-Zant**
(CTP-BUE, Cairo)
- **Joseph Silk**
(IAP/JHU/Oxford)

This meeting aims to bring together members of the astro-particle physics community - particularly those involved in dark matter dynamics, physics of the early universe and indirect detection signals -- with those working in dark matter searches at the LHC and in direct detection experiments. We aim this to be a forum for presentations, and informal interdisciplinary discussion, of the prospects for dark matter phenomenology, in light of the launch of Run 2 at the LHC and planned upgrades of direct search experiments, as well as recently planned experiments to explore the low energy hidden sector.

<http://dmc.ctp-bue.org>

Summary

- Dark Matter is an important open point in fundamental physics right now and the LHC data can contribute to the quest.
- The Higgs particle may couple to DM or may even decay into it. Invisible decays and deviations from SM Higgs couplings are explored
- Supersymmetry scenarios with R_p -conservation can have a natural DM candidate. Discovery of supersymmetry will have important impact for DM
- Generic searches for DM in analyses dealing with missing E_T : typically mono-object searches
- So far exclusion limits only, but maybe soon with the collisions at 13 TeV :



The Future?

