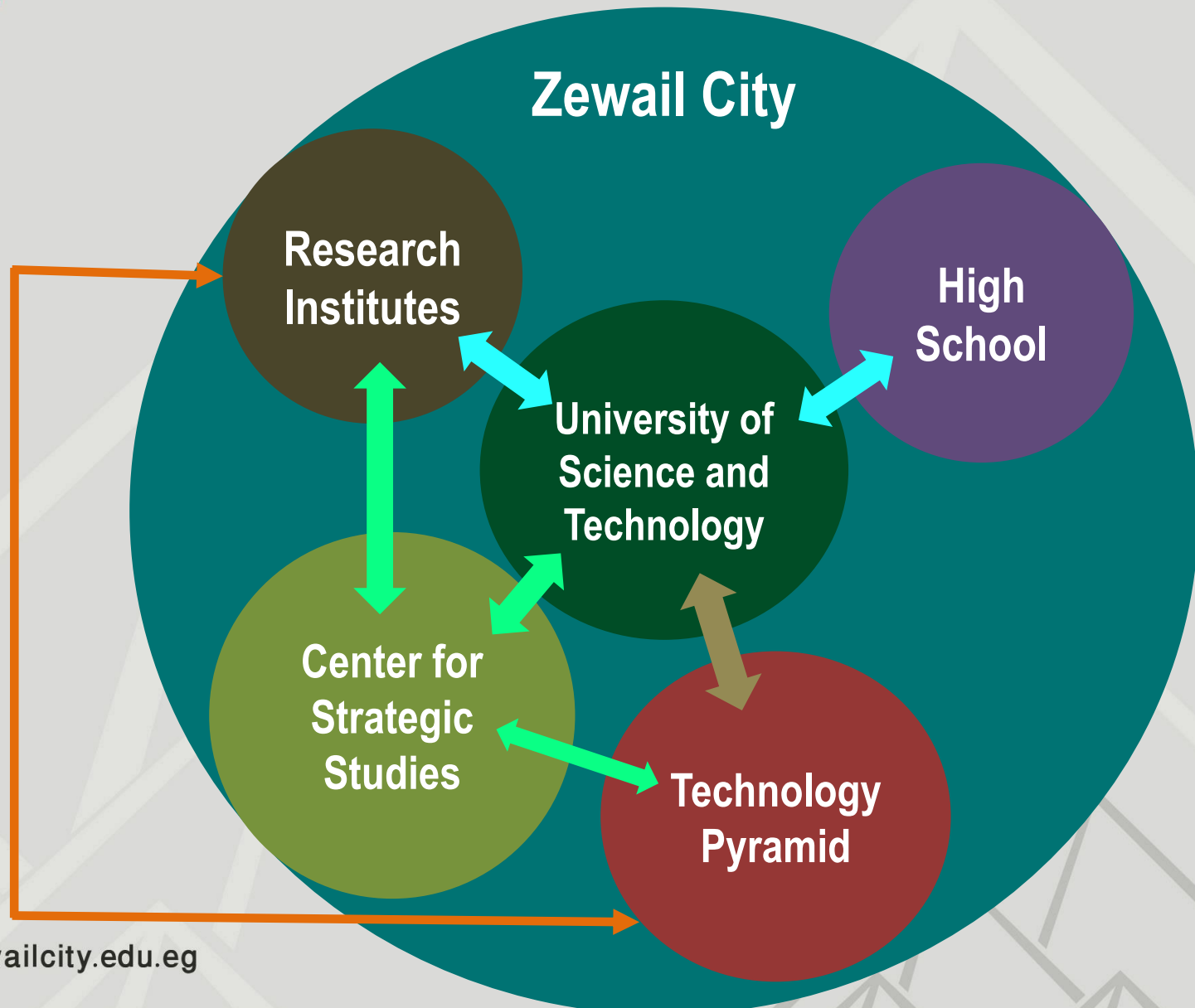


Zewail City Node



Structure of Zewail City





مركز الفيزياء الأساسية

Center for Fundamental Physics

ESTABLISHED 2000
INAUGURATED 2011

INAUGURATED 2011

Research at Zewail City

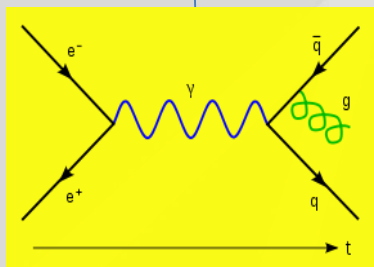
Research interest at Zewail City currently focuses on the following topics:

- Nano-science & Technology
- Bio-medical Science
- Energy, Environmental & Space
- **Basic Science (Physics)**

Center for Fundamental Physics (CFP)- HEP Theory group

- Faculty Members
 - S. Khalil
 - E. Lashin
- Postdoctoral Researchers
 - C. Un
 - A. Nassar
 - M. Abbass
- Graduate Students
 - W. Abdalla
 - A. Hammad
 - M. Amin
 - A. Moursy
 - A. Kassem
 - M. Abdelmoniem

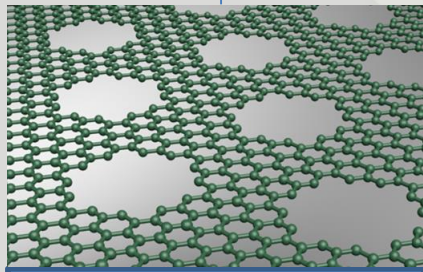
Research at Center for Fundamental Physics



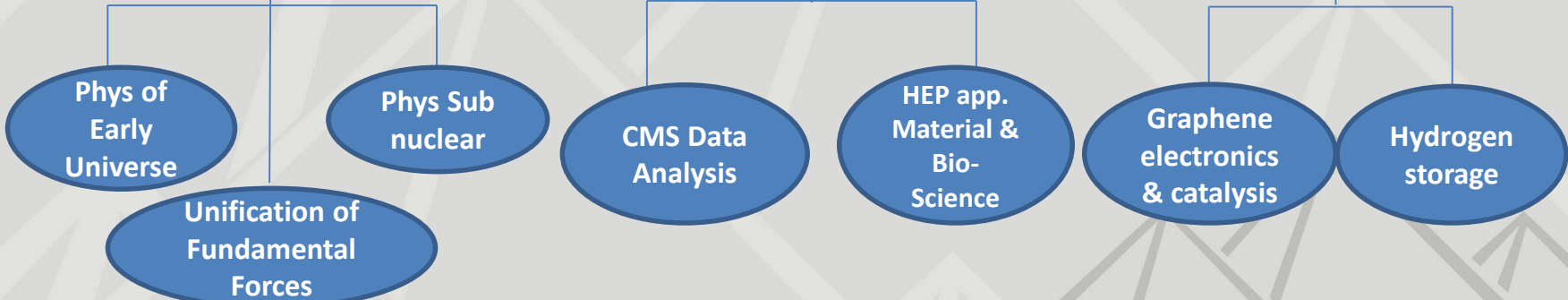
Theoretical High Energy Phys



Experimental High Energy Phys



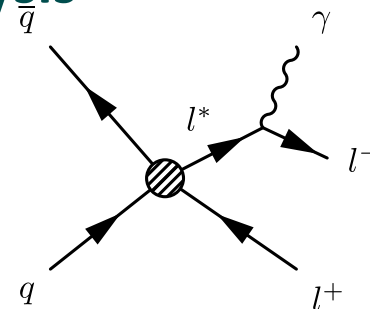
Condensed Matter Physics



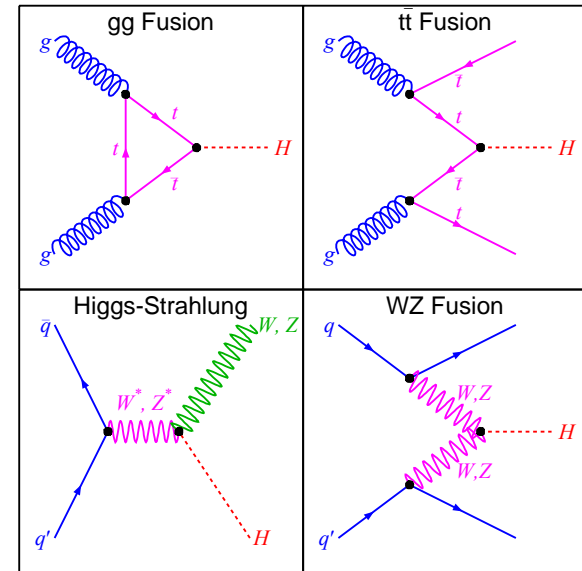
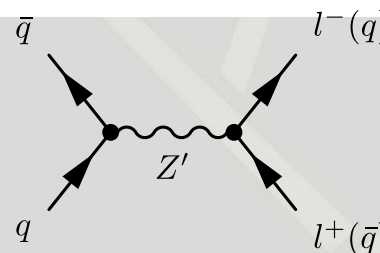
CMS: Physics analysis

- Higgs to 4 lepton analysis

- Excited Lepton analysis



- Heavy Neutral Gauge boson analysis



CMS group members:

- Dr. Amr Mohamed
- Dr. Ahmed Abdelalim
- Reham Mohamed
- Waleed Esmail
- Ashraf Kassem
- Ahmed Hammad
- Ahmed Mostafa

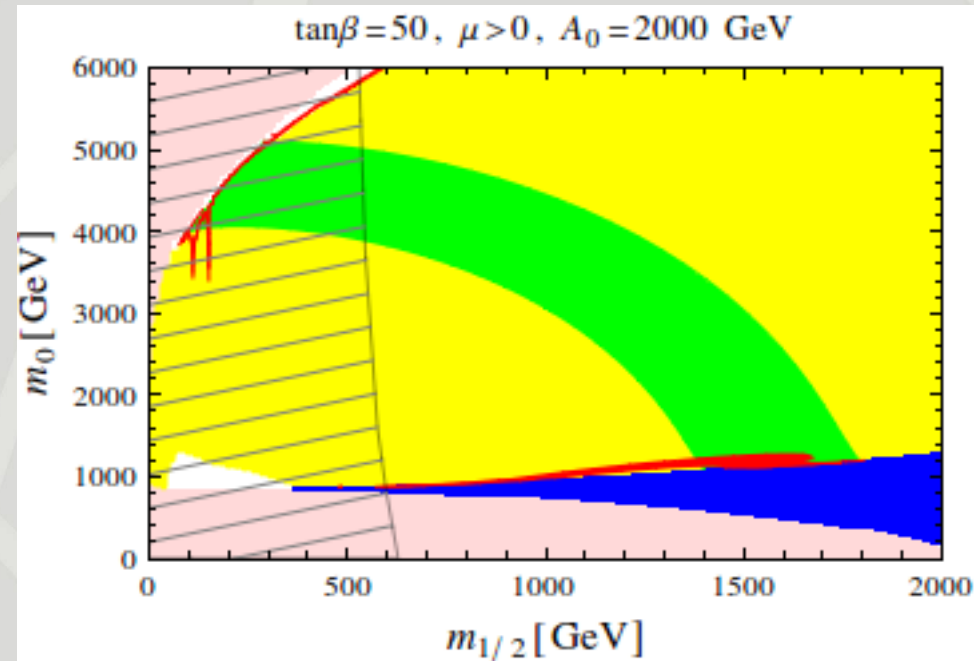
Recent and Current Activities

SUSY Dark Matter

- Combining the collider and astrophysics constraints almost rule out the MSSM.
- At large $\tan\beta$, a region with a possible resonance due to $MA = 2 m_{\tilde{\chi}}$ is allowed.

Combined LHC and relic abundance constraints rule out most of the MSSM parameter space except a very narrow region with very large $\tan\beta$ (~ 50).

- Non-universality in gaugino and/or scalar masses may remove this tension.
- Non-minimal supersymmetric extensions of the SM with a larger particle content or a higher symmetry can evade the problems of the MSSM.



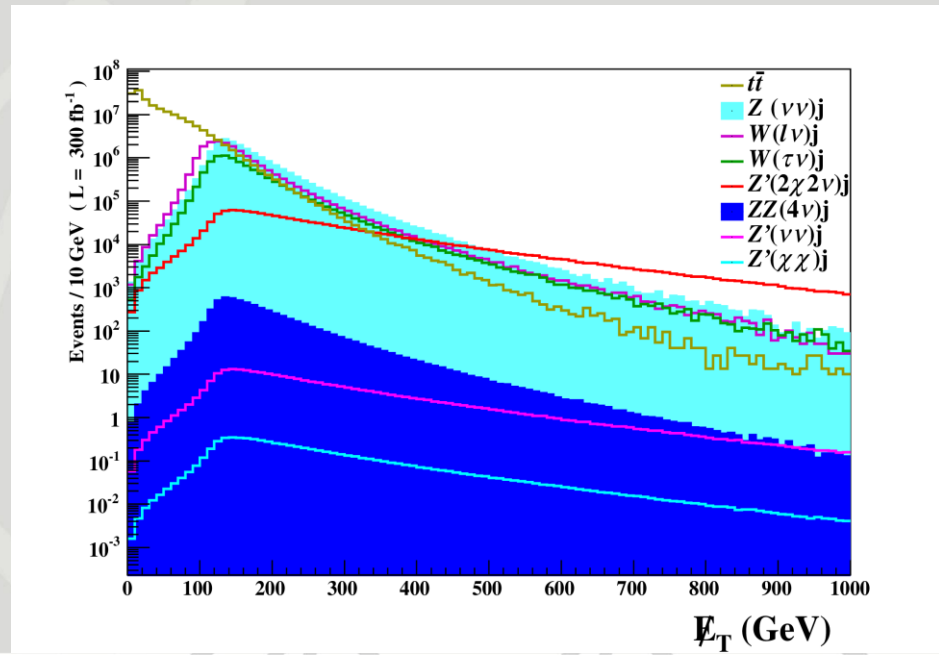
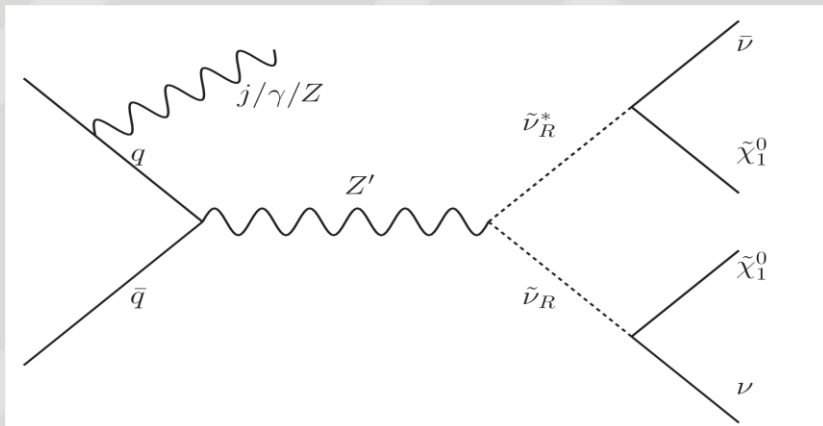
Monojet search in (B-L)SSM at LHC

The advantage of the signal is that the final state objects carry a very large (transverse) missing energy, since the Z' is naturally massive and constrained by direct searches to be at least in the TeV scale region.

Under these conditions the efficiency in accessing the invisible final state and rejecting the SM background is very high.

Another special feature of this invisible BLSSM signal is its composition, which is often dominated by RH sneutrino decays.

Sensitivity of the CERN machine can therefore help disentangling the BLSSM from more popular SUSY models.



Heavy Higgs in MSSM and beyond

- The coupling of second CP-even MSSM Higgs with SM particles is highly suppressed.

BLSSM \implies extending SM gauge groups by $U(1)_{B-L}$

BLSSM \implies Three generations of superfields right handed neutrinos $\hat{\nu}$.

BLSSM \implies Two bileptons $\chi_1 \chi_2$, needed to break $U(1)_{B-L}$.

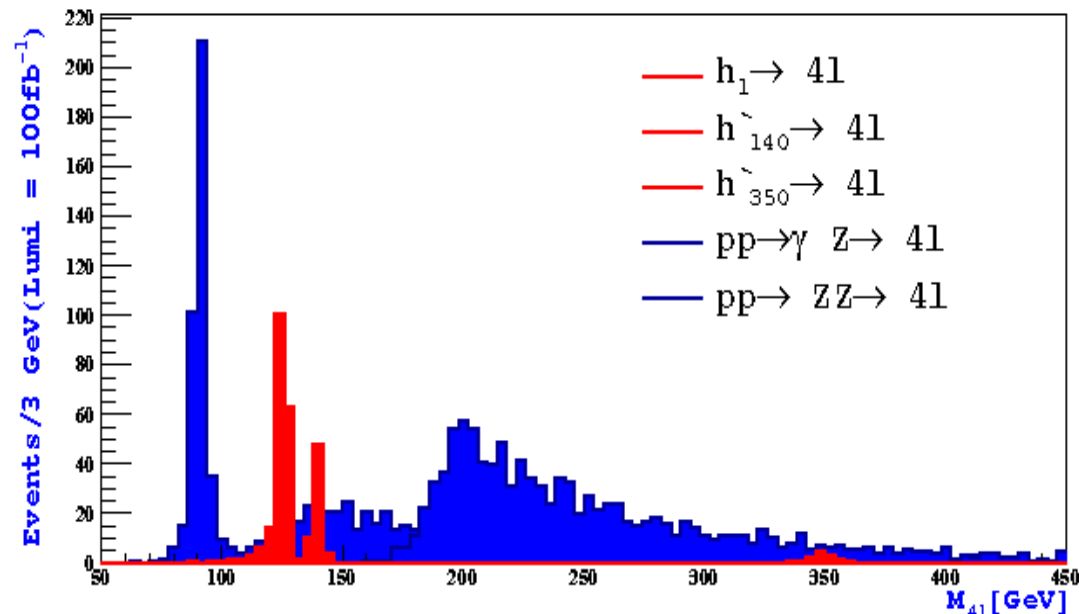
BLSSM \implies One gauge boson Z'

CP-even Higgs mass matrix can be diagonalized by unitary

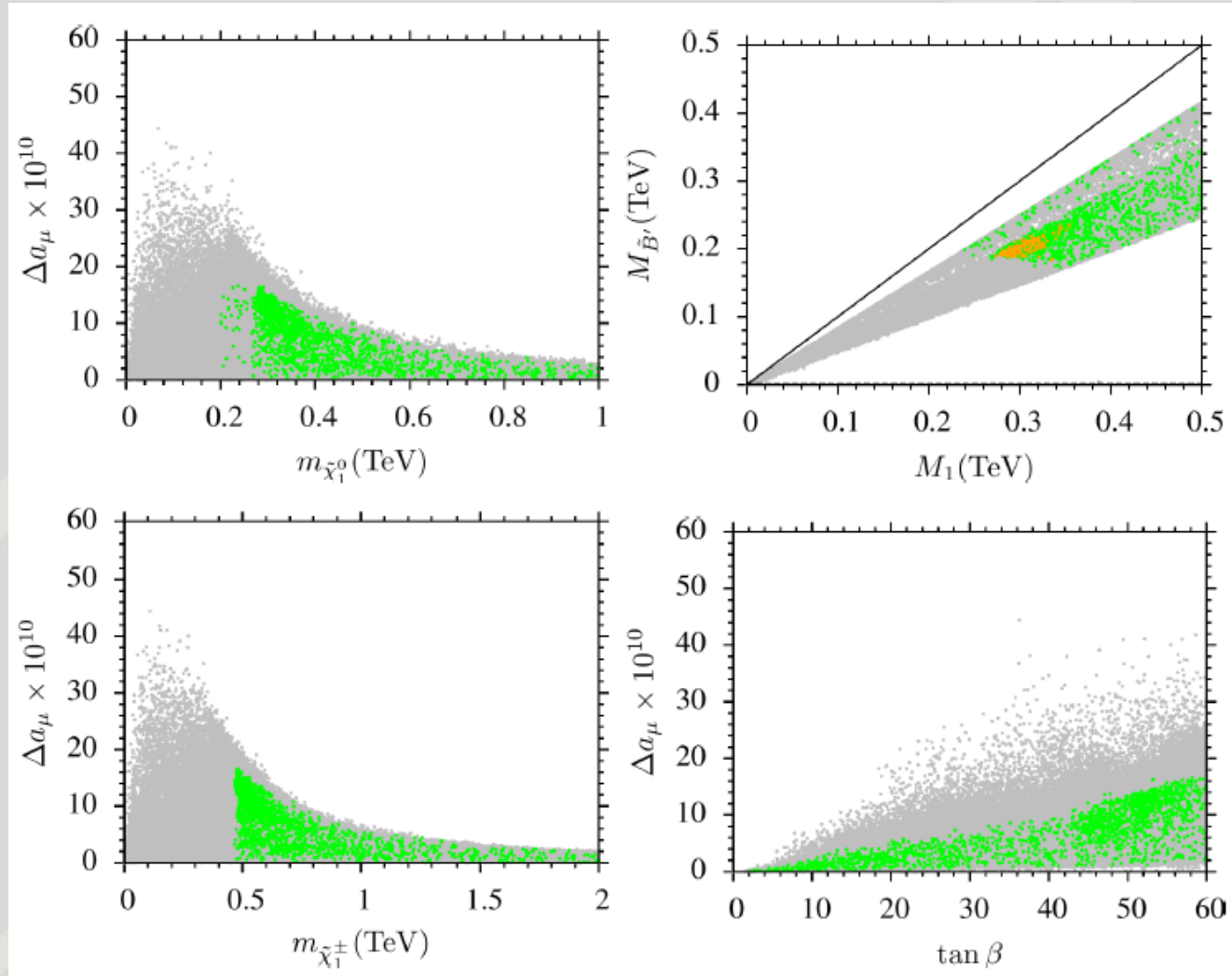
4×4 matrix

$$\Gamma M^2 \Gamma = \{ m_h^2, m_{h'}^2, m_{H'}^2, m_{H''}^2 \}$$

$h' \rightarrow ZZ \rightarrow 4l$ at Run-II



Muon $g-2$ in BLSSM



Single field inflation In SUGRA

➤ The slow-roll parameters are defined as

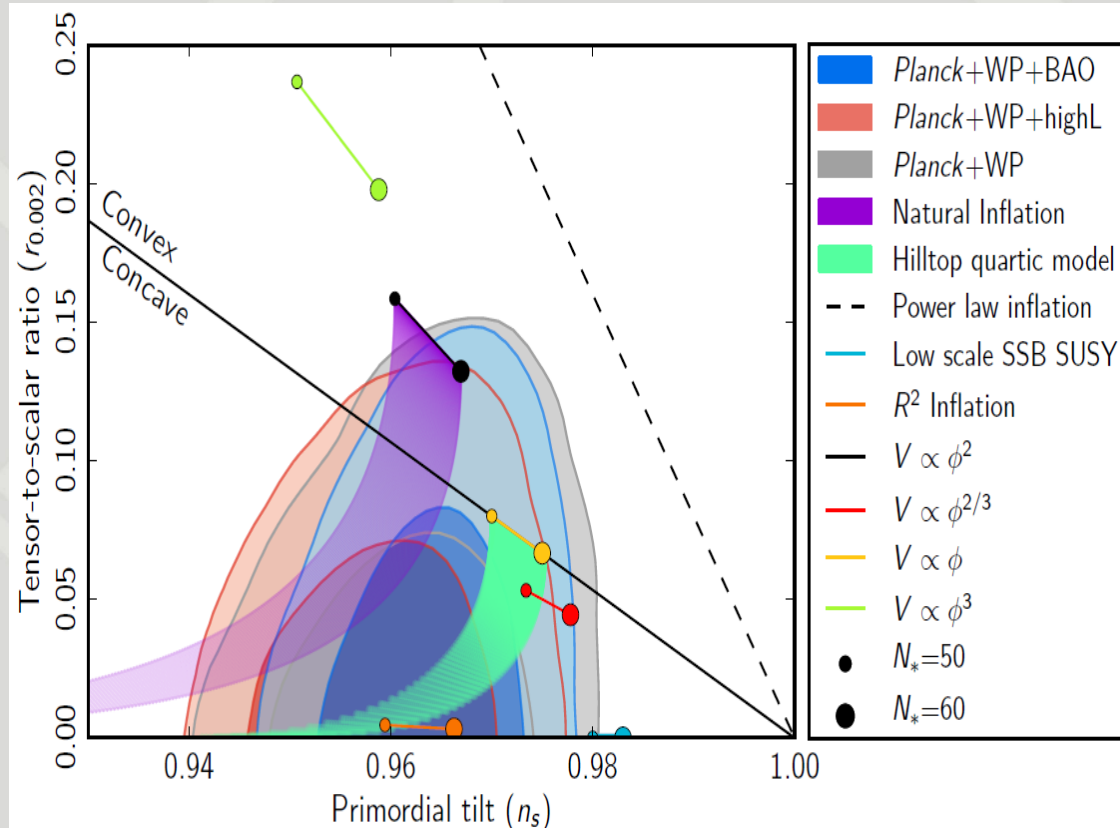
$$\varepsilon = \frac{m_{Pl}^2}{2} \left(\frac{V_\phi}{V} \right)^2, \quad \eta = m_{Pl}^2 \left(\frac{V_{\phi\phi}}{V} \right), \quad |\eta|, \varepsilon \ll 1$$

We define the inflation observables: The spectral index n_s and the tensor to scalar ratio r as following

$$n_s = 1 - 6\varepsilon + 2\eta,$$

$$r = 16\varepsilon,$$

➤ Planck Satellite constraints on inflation observables



Moduli Stabilization in String Theory

- Kachru, Kallosh, Linde and Trivedi (KKLT) could stabilize the volume modulus T using a non-perturbative superpotential and Kahler potential

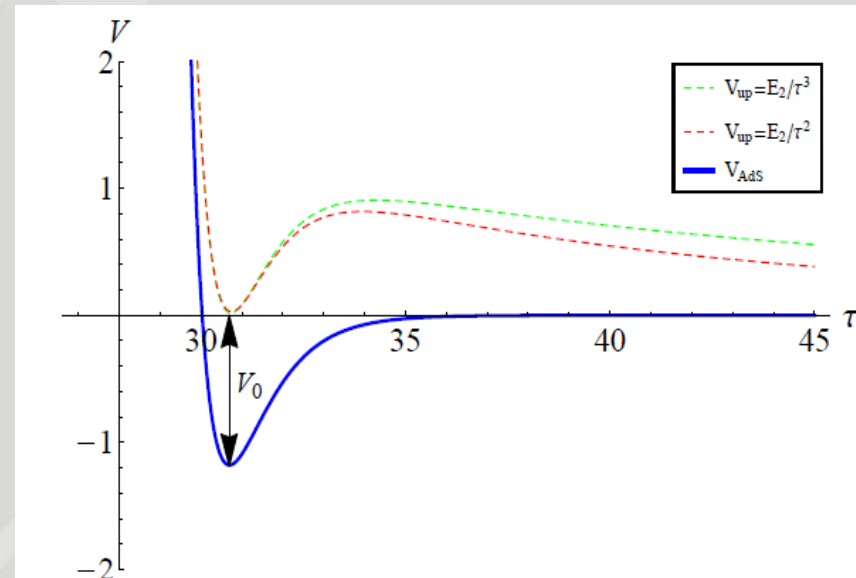
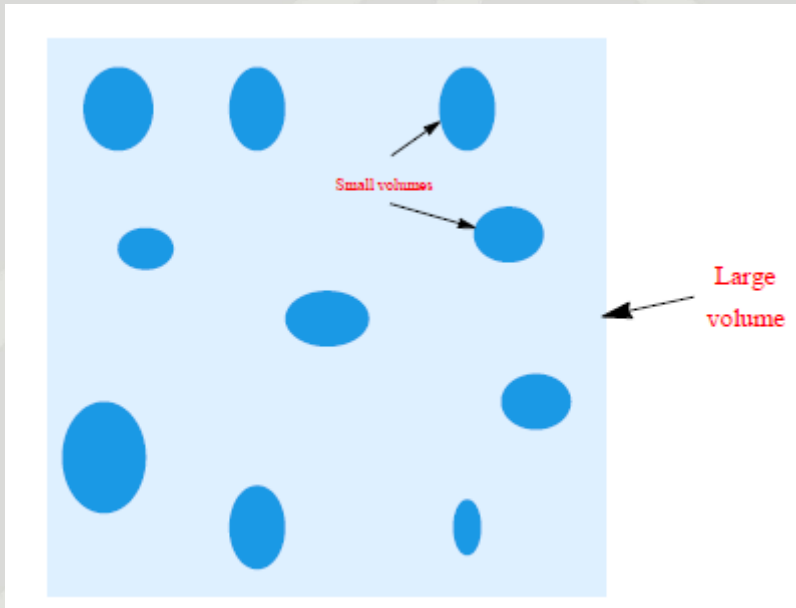
$$W = W_0 + Ae^{-aT}$$

$$K = -3 \ln[T + \bar{T}]$$

- In Large Volume Scenario (LVS), potential possesses an AdS minimum at exponentially large volume

$$K = K_{cs} - 2 \log \left(\mathcal{V} + \frac{\xi}{2} \right),$$

$$W = W_0 + A_4 e^{-\frac{a_4}{g_s} T_4} + A_5 e^{-\frac{a_5}{g_s} T_5}$$



String Phenomenology

➤ SUSY is broken by the uplift mechanism

➤ In KKL_T, the soft terms are given by

➤ Unacceptable phenomenologically

➤ In LVS, the soft terms are given by

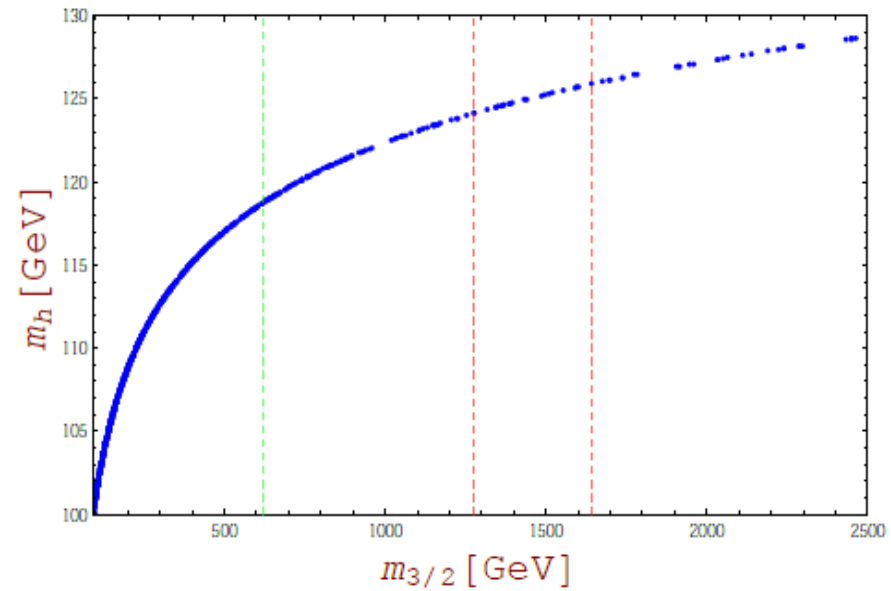
$$m_0^2 = \frac{|W|^2}{(2\tau)^3} \Big|_{dS} = m_{3/2}^2,$$

$$m_{1/2} = \frac{\sqrt{2\tau}}{6} D_T W(T) \frac{\partial}{\partial T} \ln(\text{Re} f^*) \Big|_{dS} \simeq \frac{m_{3/2}}{a\tau},$$

$$A_0 = -\frac{1}{\sqrt{2\tau}} \bar{D}_{\bar{T}} \bar{W}_h \Big|_{dS} = -\frac{3m_{3/2}}{a\tau},$$

$$m_0 \simeq m_{1/2} = m_{3/2}$$

$$A_0 \simeq -m_{3/2}.$$



Thank You