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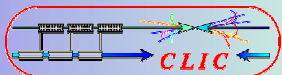
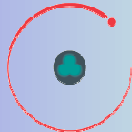
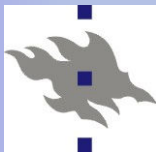
Progress in Breakdown Modelling – MD and PIC Breakdown Simulations

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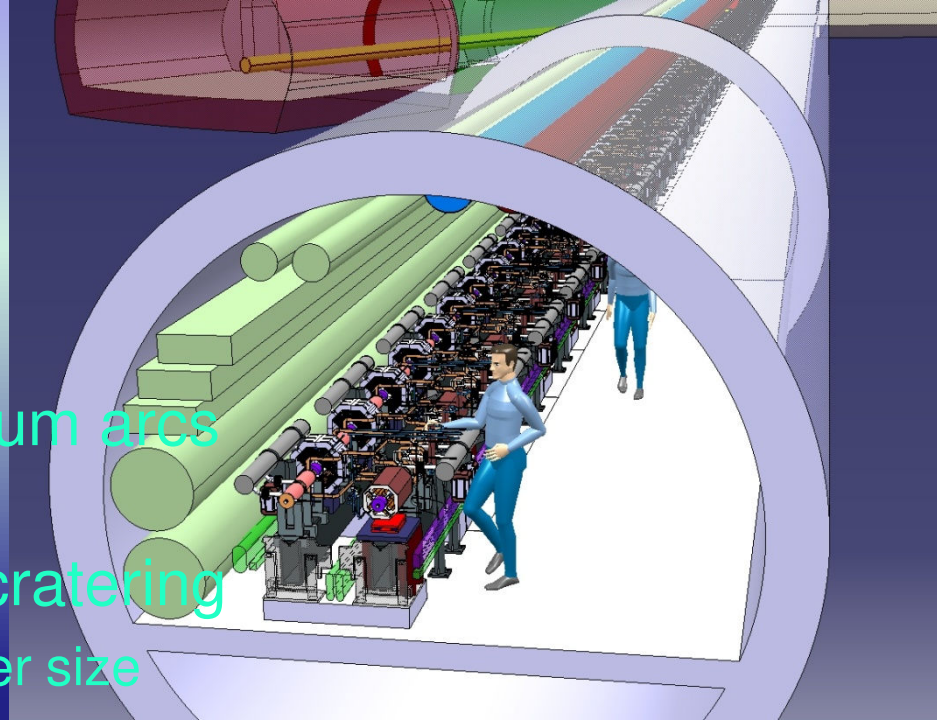
*Konstantin Matyash,
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Max-Planck Institut für Plasmaphysik

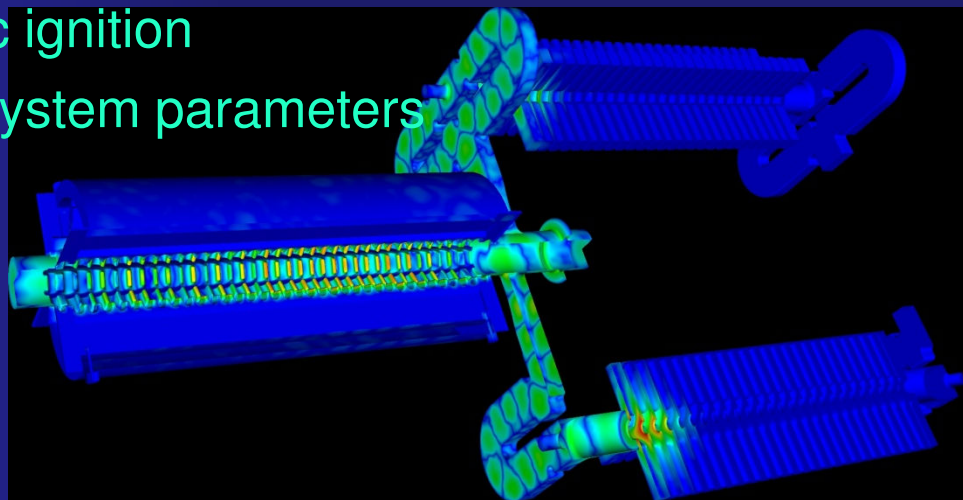
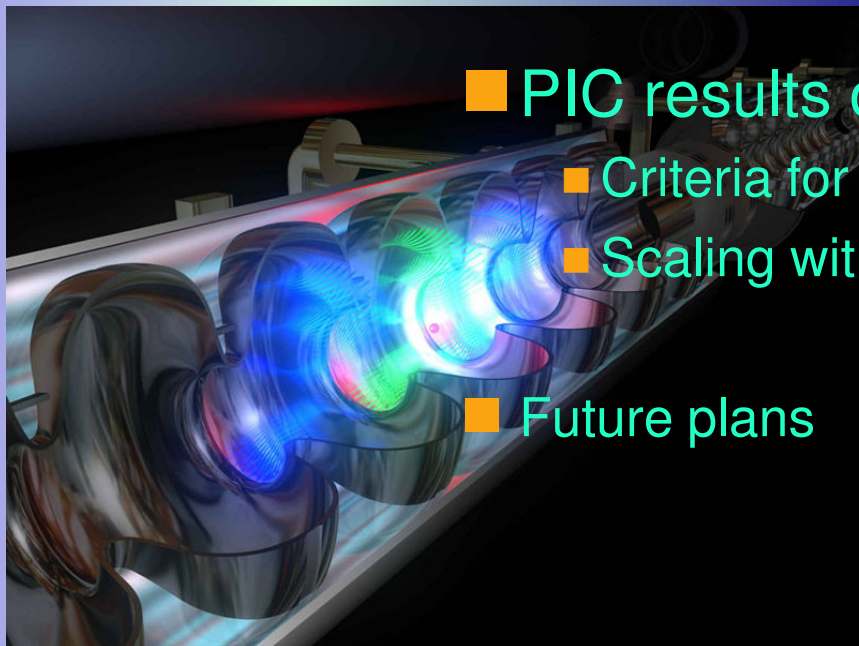


Outline

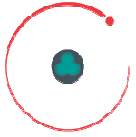
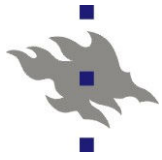
- Modelling vacuum arcs
- MD results on cratering
 - Scaling of crater size



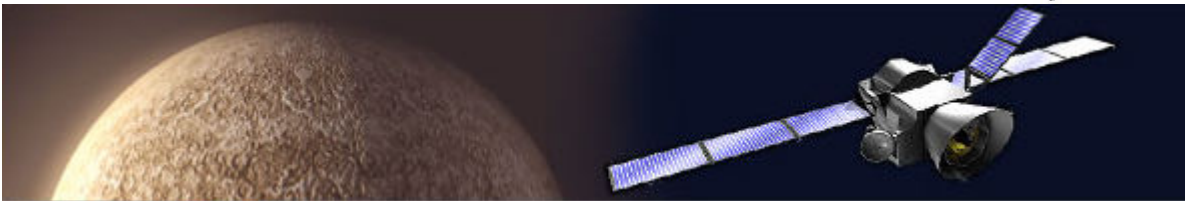
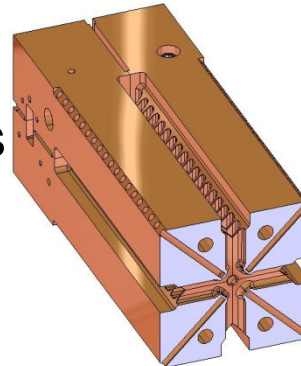
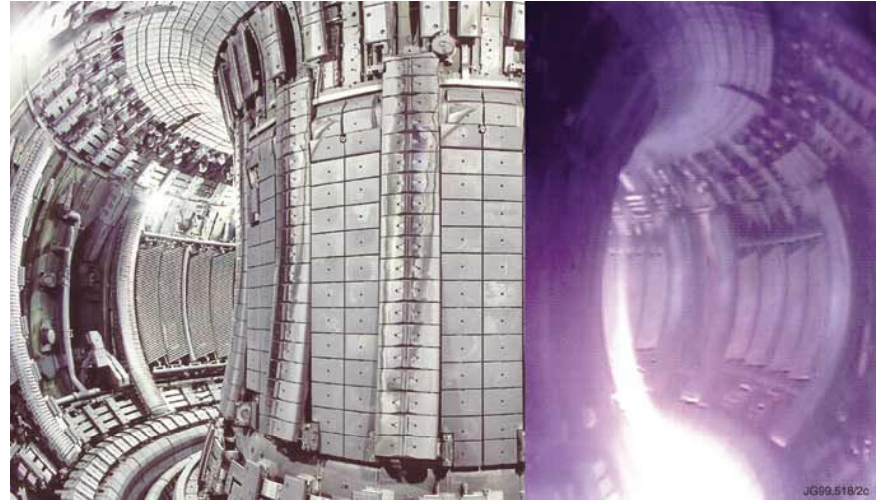
- PIC results on plasma build-up
 - Criteria for arc ignition
 - Scaling with system parameters
- Future plans



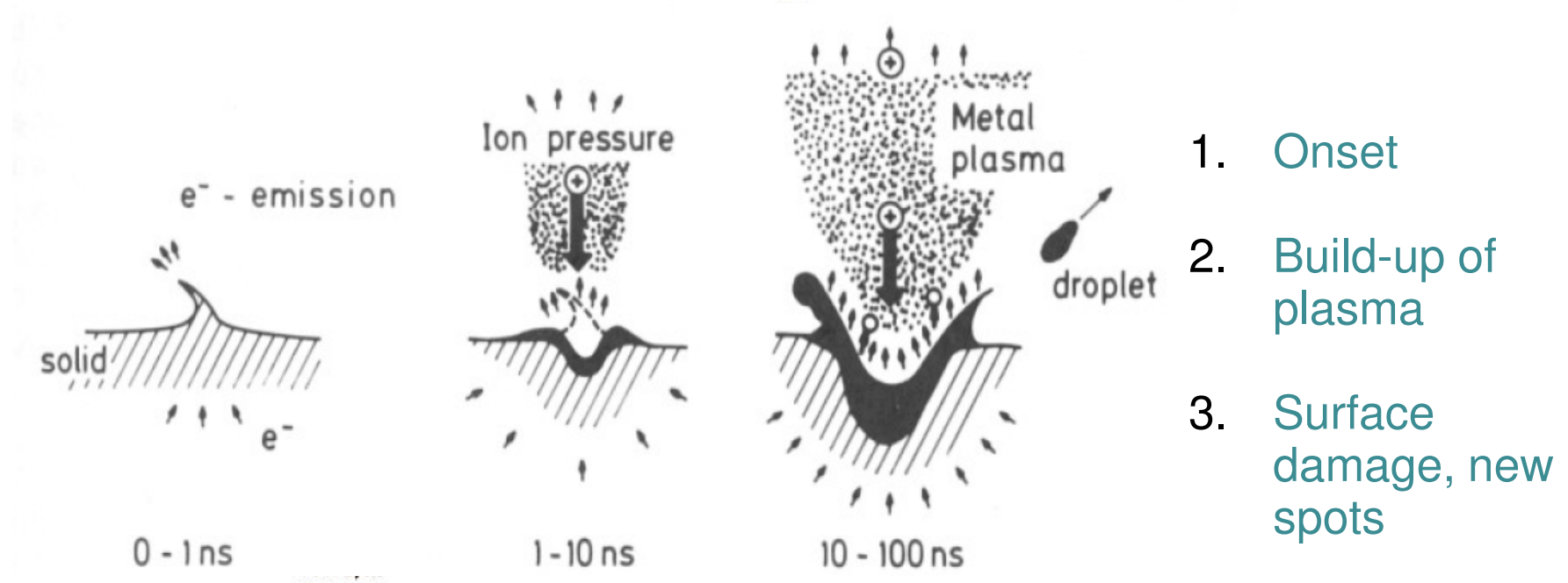
Breakdown studies have a broad application spectrum



- Fusion physics
- Satellite systems
- Industry
- Linear collider designs



Finally an electric field builds up between the plasma and the surface of the solid named the *Langmuir sheath potential*. Electrical arcs may ignite between the plasma representing the negative electrode and the vessel wall representing the positive electrode.



Schematic of the ignition of an electrical arc at a surface tip, as well as burning and movement of the cathode spot [21, 22].

Erosion yields by electrical arcs, in atoms/electron [21],[25]-[30]. The estimated total number of atoms removed for an arc current of 5 to 10 A and a burn time of about 10-100 μ s are also introduced.

Multiscale model

Our model corresponds to the above 3 phases.

Stage 0: Onset of tip growth; Dislocation mechanism

Method: MD, Molecular Statics...

~ sec/min

Stage 1: Charge distribution @ surface

Method: DFT with external electric field

~few fs

Stage 2: Atomic motion & evaporation

Method: Hybrid ED&MD model

Classical MD+Electron
Dynamics: Joule heating,
screening effect

Solution of Laplace
equation

~few ns

**Stage 3: Evolution of surface morphology due
to the given charge distribution**

Method: Kinetic Monte Carlo

~ sec/hours

=> Electron & ion & cluster emission ions

Stage 4: Plasma evolution, burning of arc

Method: Particle-in-Cell (PIC)

~10s ns

=> Energy & flux of bombarding ions

**Stage 5: Surface damage due to the
intense ion bombardment from plasma**

Method: Arc MD

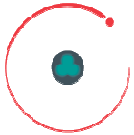
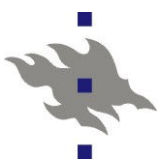
Onset of
plasma

Plasma
build-up

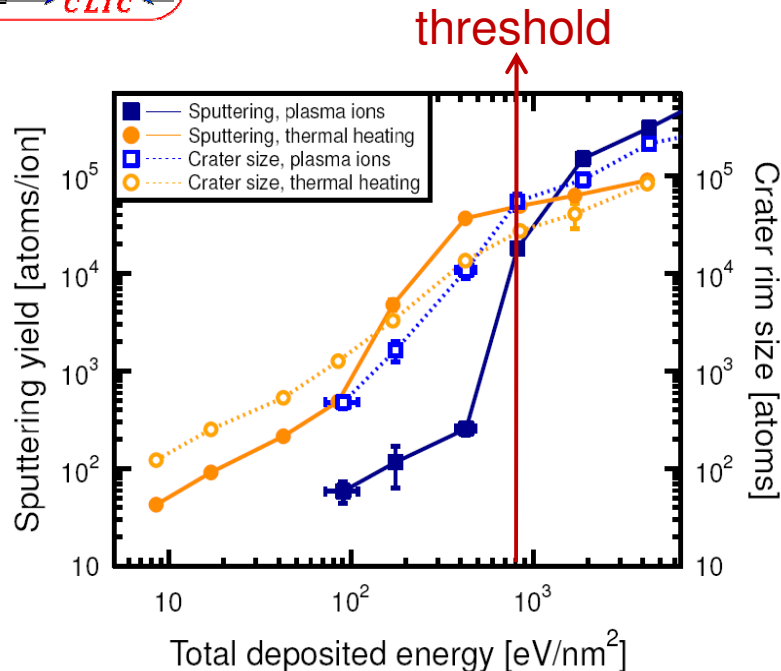
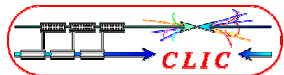
Surface
damage

~100s ns

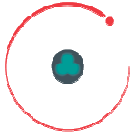
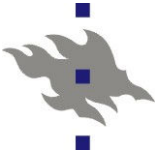
Achievements



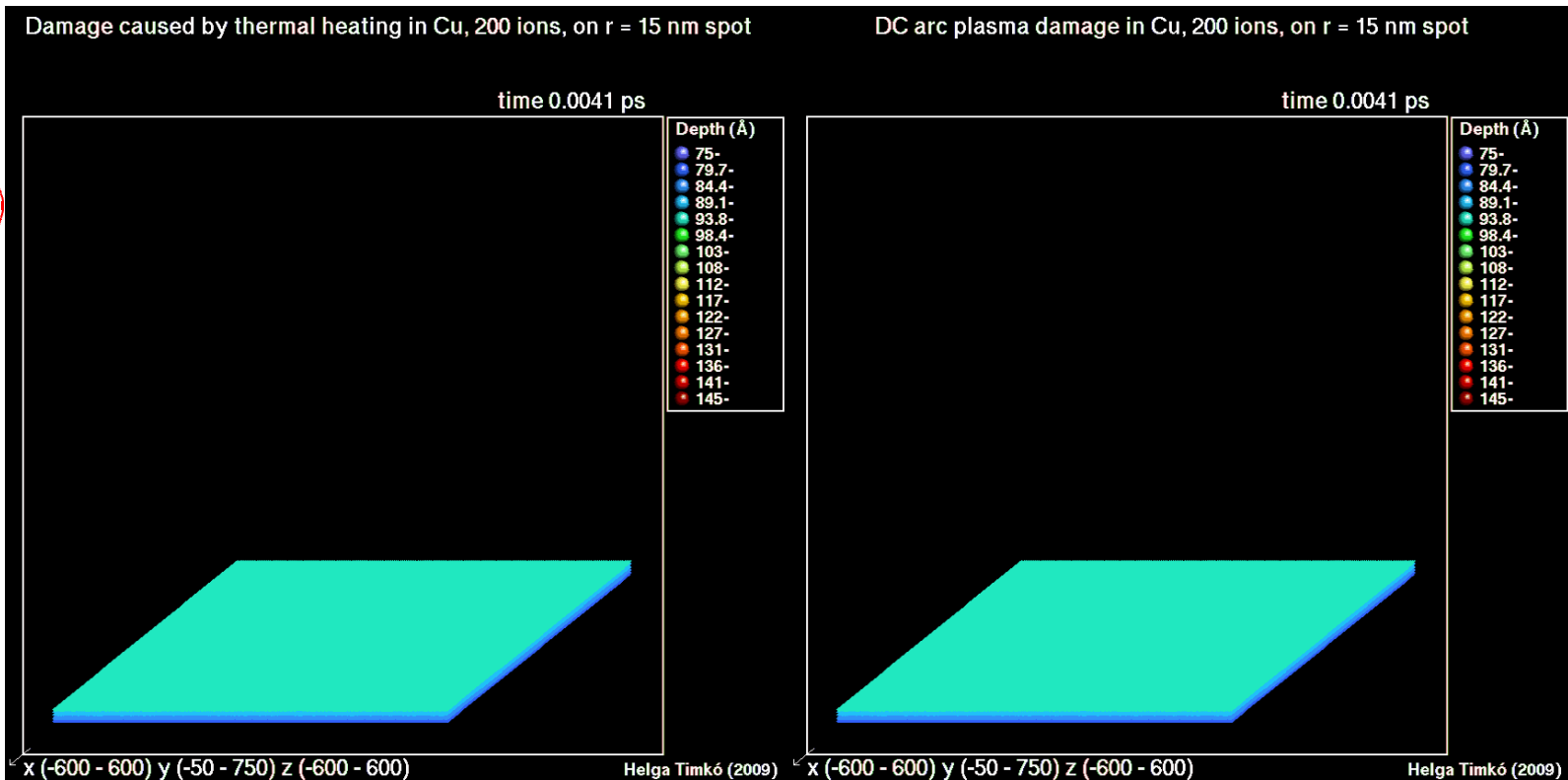
1. **Onset:** direct field evaporation from surfaces and tips
2. **Plasma build-up:** we have developed a one-dimensional PIC model and identified plasma build-up criteria
3. **Cratering:** knowing flux & energy distribution of incident ions, erosion and sputtering was simulated with MD



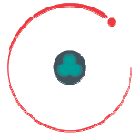
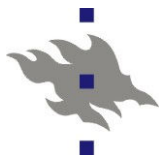
- Comparing arc plasma bombardment and thermal heating, we found that:
 - Enhanced sputtering yield above a threshold, corresp. to the melting point
 - Only for plasma bombardment:
 - (i) heat spike & cluster emission above the threshold
 - (ii) experimentally seen complex crater shapes can form



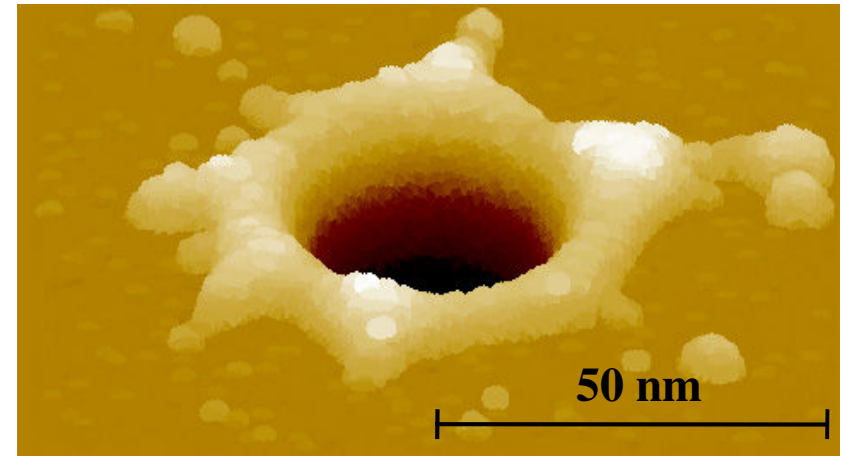
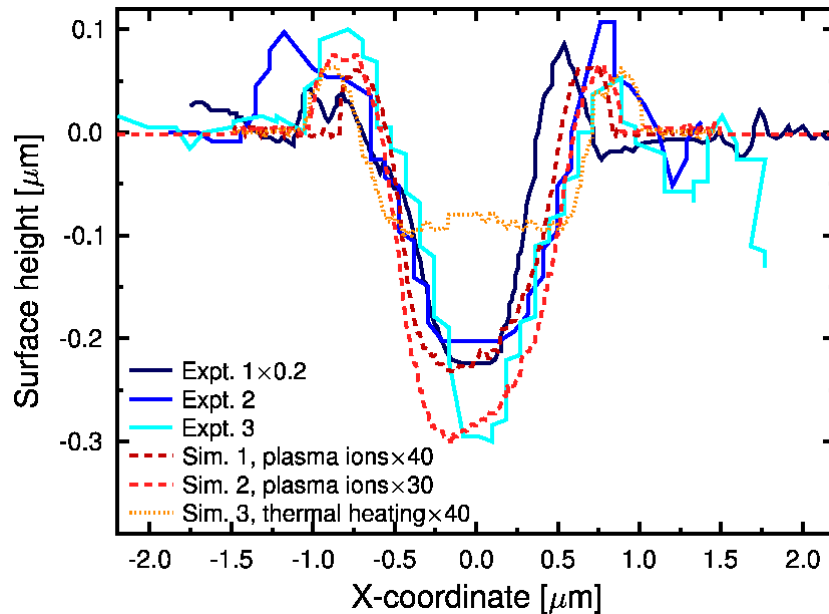
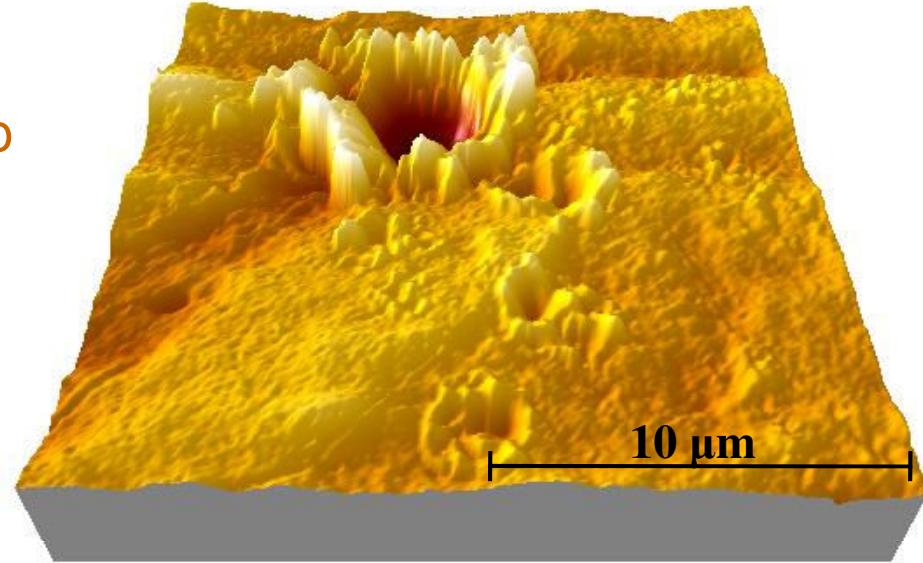
MD simulation of surface damaging

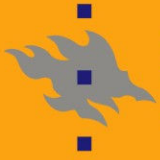


Comparison to experiment



■ Self-similarity:
Crater depth to width ratio
remains constant over
several orders of
magnitude, and is the
same for experiment
and simulation



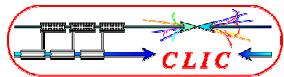
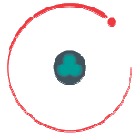
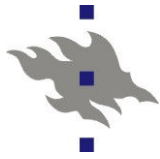


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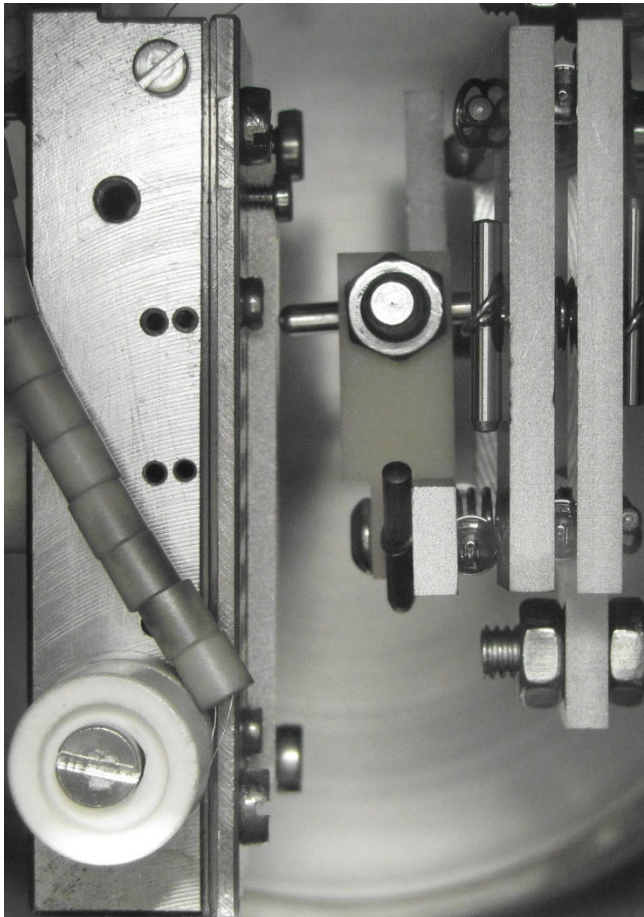


Now to the plasma part...





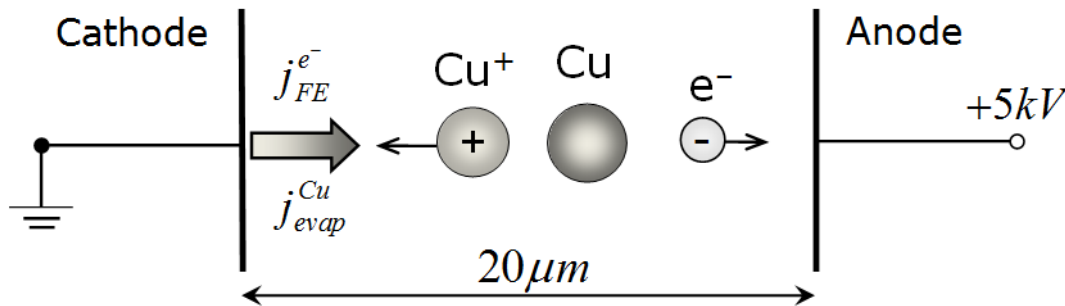
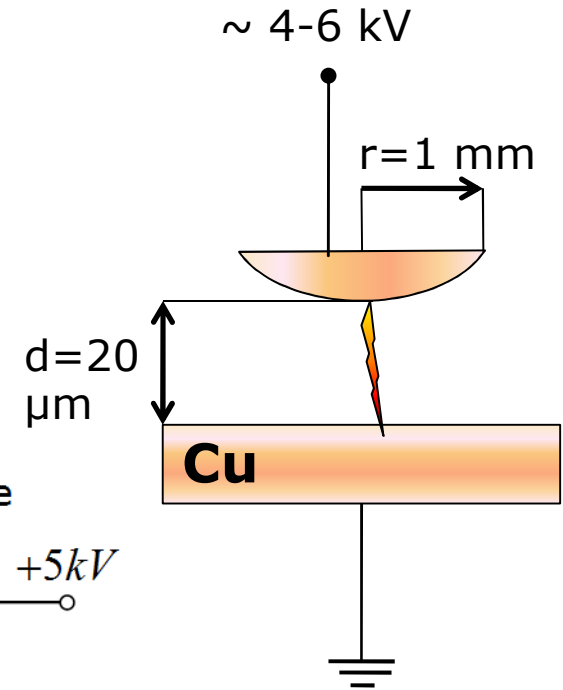
Modelling DC arcs



- First we have to understand breakdowns in DC, before we can generalise to RF
- To have a direct comparison with experiments, we adjusted simulation parameters to the DC setup at CERN
- However, the results we present here are completely general and not restricted to the DC setup!

Corresponding to experiment

- 1d3v electrostatic PIC-MCC code
 - Resolving the main stream of plasma
 - Areal densities of physical quantities

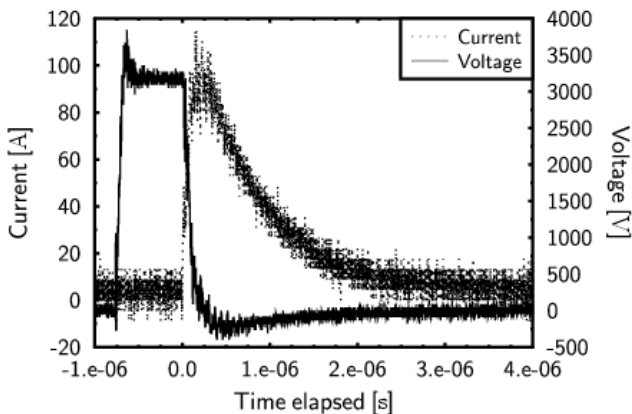
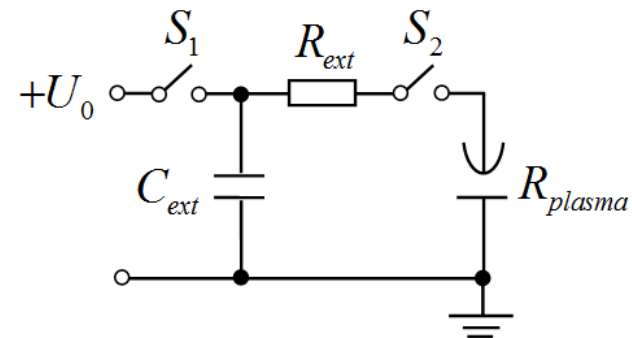


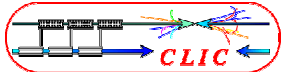
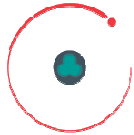
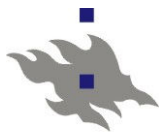
Exponential voltage drop mimiced

- Limited energy from the circuit

$$R_{ext} = 30\Omega$$

$$C_{ext} = 0.1 - 27.5nF$$





Phenomena taken into account

We started from a simple model with a code from IPP-MPG
(Collaborators: *R. Schneider, K. Matyash*)

- Field emission of electrons, Fowler-Nordheim eq.:

$$j_{FE} = a_{FN} \frac{(eE_{LOC})^2}{\phi t(y)^2} e^{-b_{FN} \frac{\phi^{3/2} v(y)^2}{eE_{loc}}}, \quad \text{where } E_{loc} = \beta \cdot E$$

$$t(y) = 1, \quad v(y) = 0.956 - 1.062y^2 \quad \text{where } y = \sqrt{\frac{e^3 E_{LOC}}{4\pi\epsilon_0 \phi^2}}$$

- Evaporation of Cu neutrals

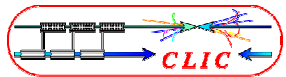
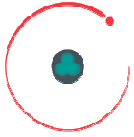
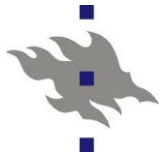
- Collisions, esp. ionisation collisions

- Sputtering of Cu neutrals at the wall
- Secondary electron yield due to ion bombardment

Start from these

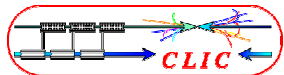
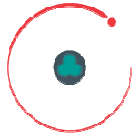
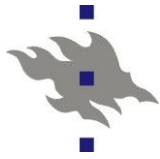
Produce ions

More e- & Cu



Plasma build-up from a field emitter tip

- We start from a field emitter tip → supply of electrons and neutrals → build-up of plasma
- The *field emitter* is assumed in terms of an initial field enhancement factor
 - Dynamic beta: the "erosion" and the "melting" of the tip was implemented
 - We define the "melting current" j_{melt} as the threshold of electron emission current, which, if exceeded, sets $\beta=1$
- *Neutral evaporation*: an estimate was needed
 - Define the neutral evaporation to electron FE ratio $r_{\text{Cu/e}} = r_{\text{Cu/e}}(E, t, \dots)$ and approximate it with a constant

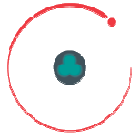
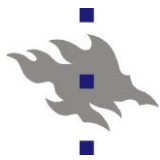


Under what conditions will an arc form?

Two conditions need to be fulfilled:

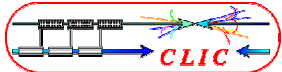
- High enough **initial local field** to have growing FE current
- Reaching the **critical neutral density** to induce an ionisation avalanche

- The sequence of events leading to plasma formation:
 - Due to high electric field: electron FE, neutral evaporation
 - Ionisation and acceleration of the charged particles
 - ⇒ e^- , Cu and Cu^+ densities build up
 - "Point of no return": $I_{mfp} < I_{sys}$ – corresponding to a critical neutral density $\sim 10^{18} \text{ 1/cm}^3$ in our case
 - Ions → sputtering neutrals → more ions ⇒ ionisation avalanche

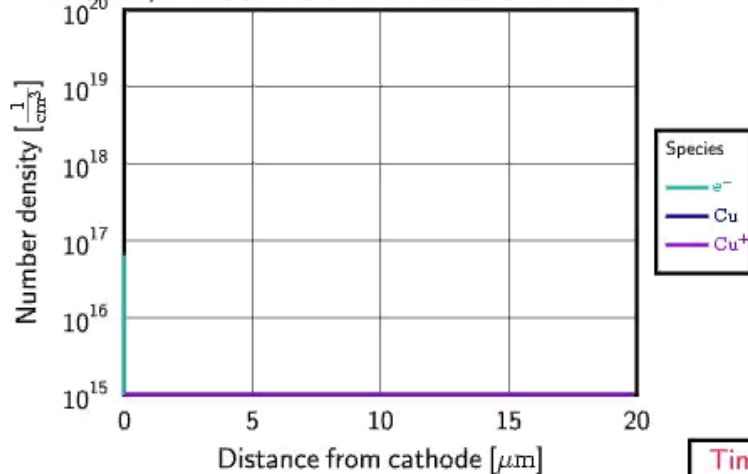


Plasma build-up

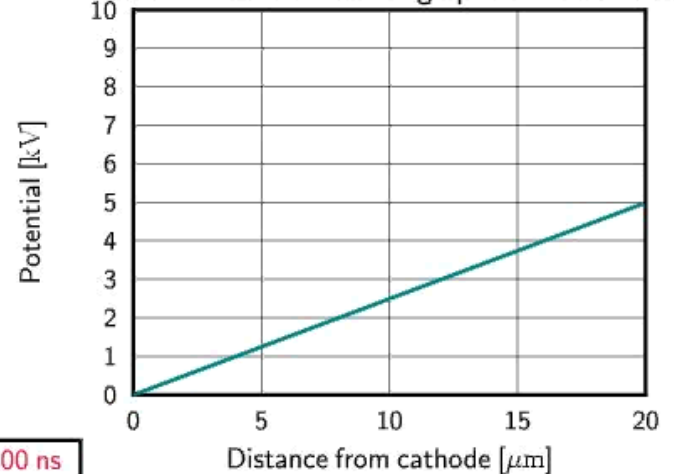
The only limiting parameter is what power can be supplied to the arc



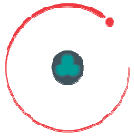
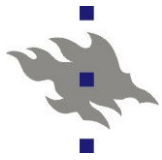
Electron, Neutral and Ion Densities in the Arc Plasma



Electrical Potential Building up in the Arc Plasma

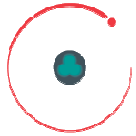
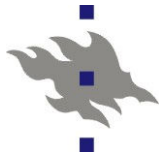


Parameter space investigated



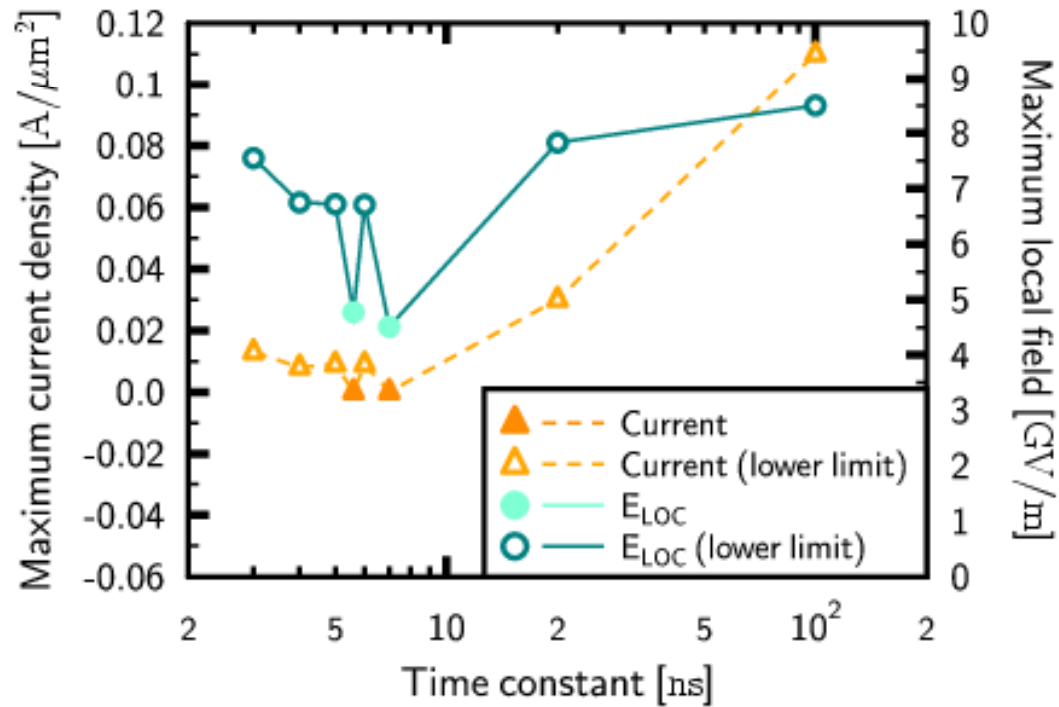
	$\tau =$ 2 ns	■ no plasma
	3 ns	realistic timescales
below critical Cu density		timescale below 1 ns
$r_{Cu/e} = 0.001, 0.005, 0.008,$	4 ns $0.01,$	$0.025, 0.05, 0.1$
		■ close to critical Cu density ■ separate current peaks for FE and plasma
$j_{melt} = 0.4,$	5 ns $0.5,$	$0.6, 0.8, 0.9, 1 \frac{A}{\mu m^2}$
qualitatively same behaviour	5.6 ns 6 ns 7 ns	■ close to critical Cu density ■ current peaks for FE and plasma grow together
	10 ns 20 ns 100 ns	■ ionisation avalanche ■ current peaks for FE and plasma grow together

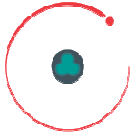
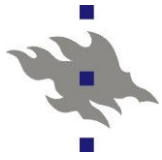
$$E_{LOC}^0 = 10 \frac{GV}{m}$$



Time constant

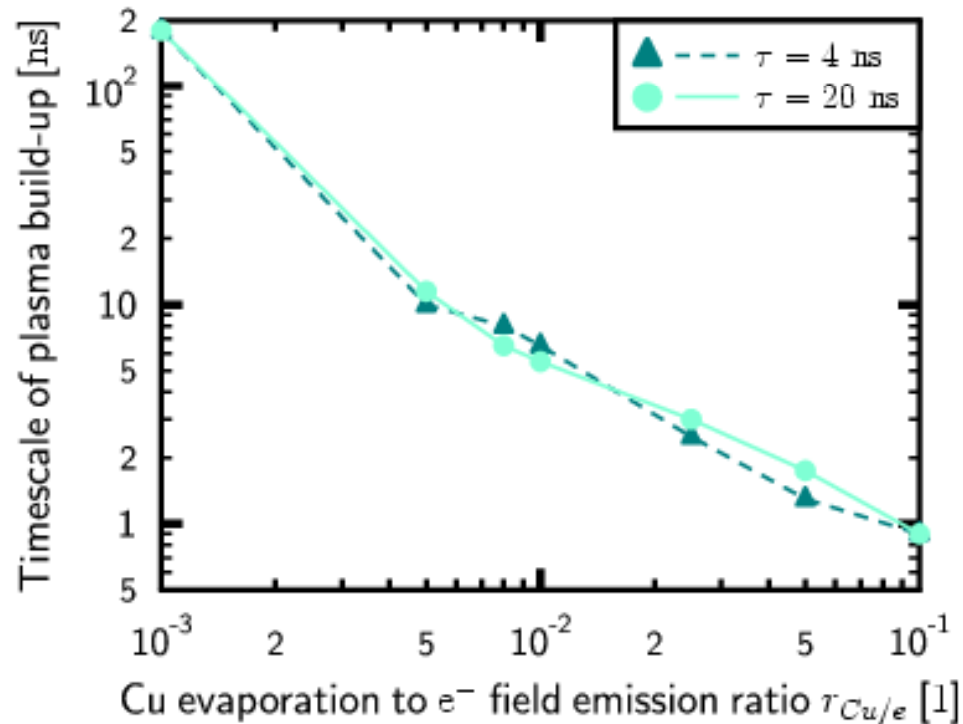
- Close to critical Cu density below ~ 10 ns
- Above ~ 10 ns, plasma formation is unavoidable

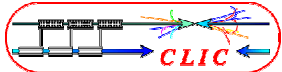
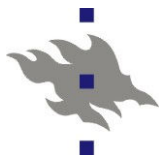




Neutral evaporation to electron FE ratio

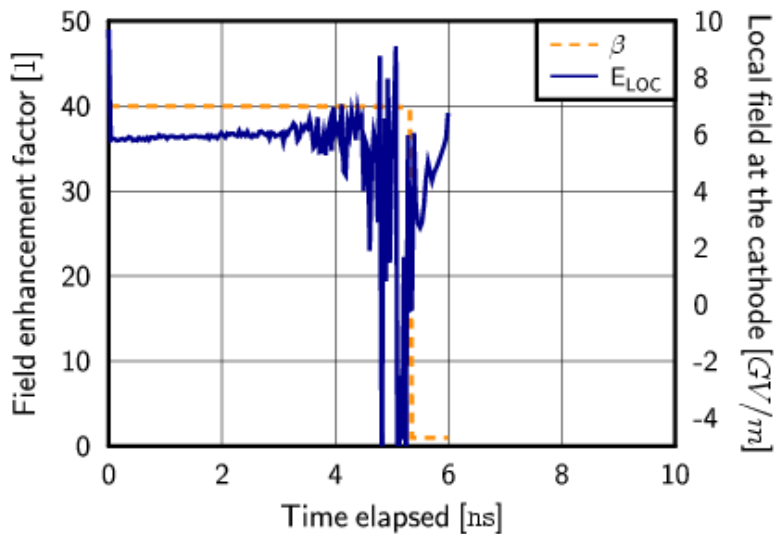
- 0,001 – 0,008: below critical Cu density
- 0,01 – 0,05 gives realistic timescales for plasma build-up



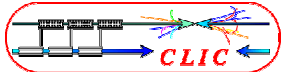
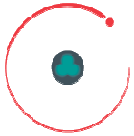
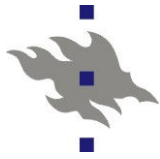


Initial local field required

- Up to now, 10 GV/m was assumed (measured value)
- Lowering E_{LOC} (either β or E) gave drastical changes
 - 8 MV/m: no ionisation avalanche any more
 - 7.5 MV/m and lower: no plasma at all

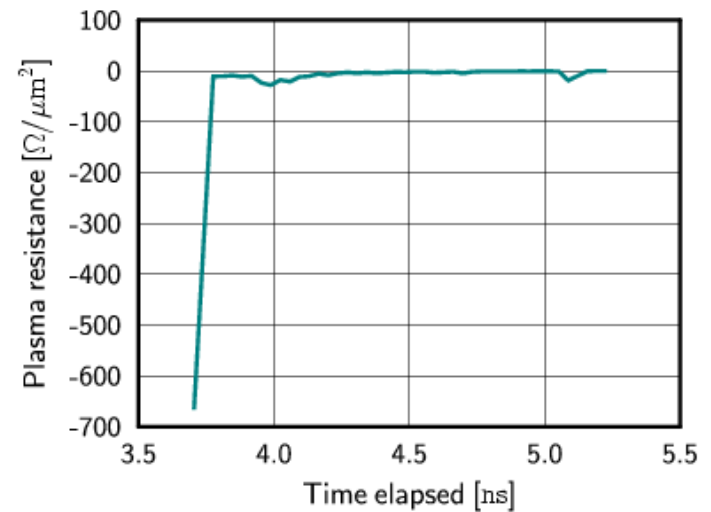
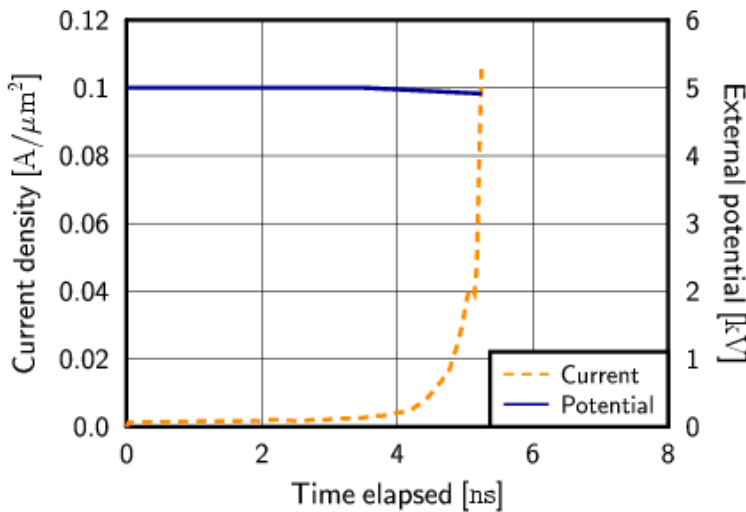
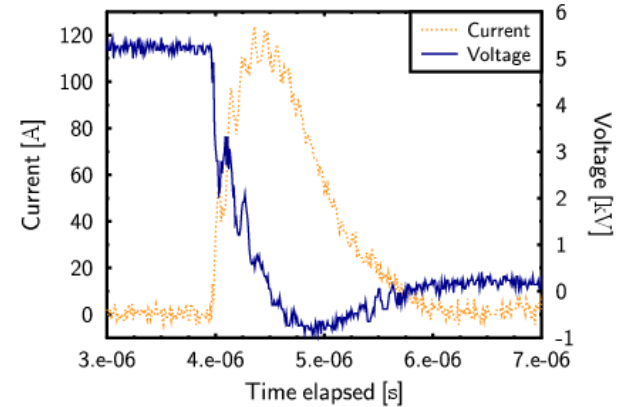


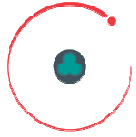
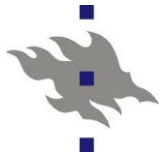
- The criterion seems to be:
 - to stabilise around ~ 6 GV/m to get growing FE current
- What happens if $E_{\text{LOC}} = 12$ GV/m?
 - It also stabilises to 6 GV/m only!
 - Note: BDR = 1 reached



Circuit characteristics

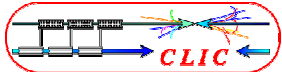
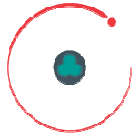
- Plasma has negative resistance
- The plasma seems to match the impedance of the external circuit to consume the available energy in the most effective way





Conclusions of the 1D model

- When the 2 required conditions (high enough initial local field, reaching the critical Cu density) are fulfilled, plasma formation is inevitable
- The 1D model is suitable to obtain information on fluxes, densities etc. in the main stream of the plasma
 - Restricted to the build-up phase of plasma
- Also RF can be simulated, requires only minimal modifications in the code

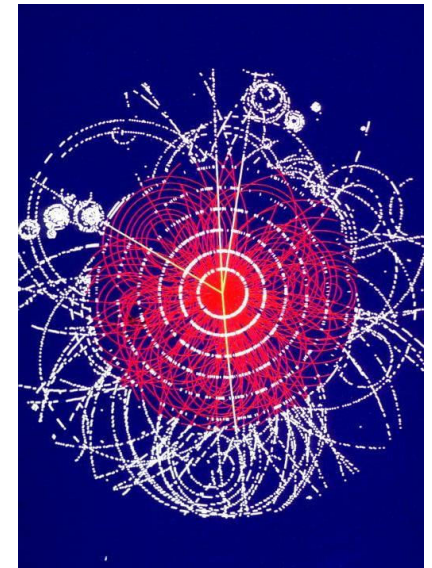


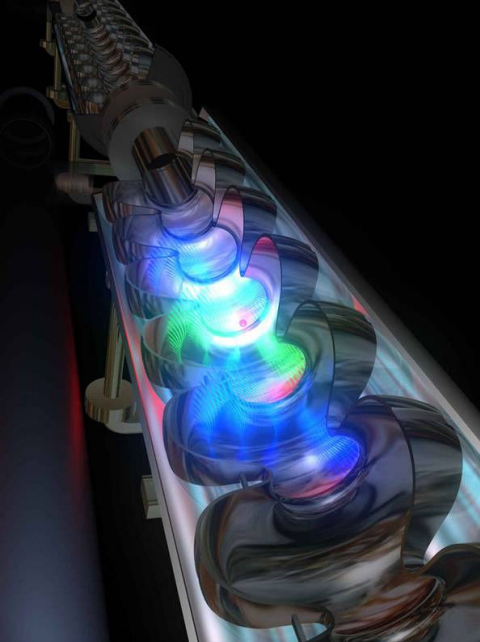
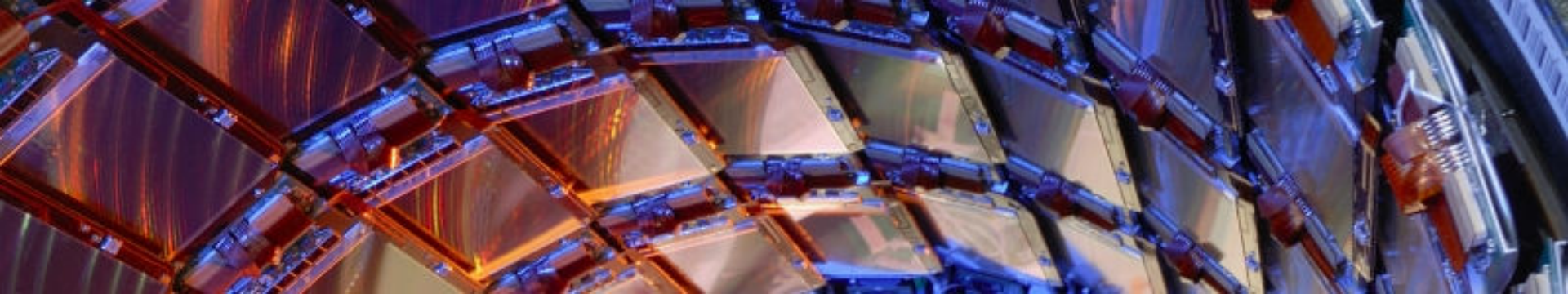
Future plans

- Extension to a 2D model; we gain:
 - Information on the radial distribution and diffusion of the plasma
 - Resolving area
 - Self-consistent PIC-MD coupling
 - Self-consistent coupling between the external circuit & discharge gap

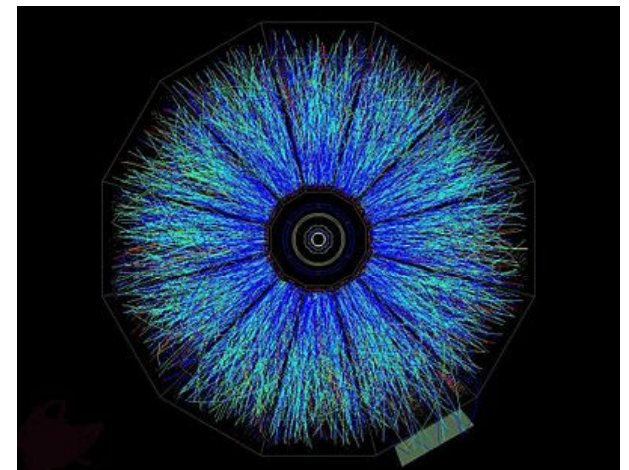
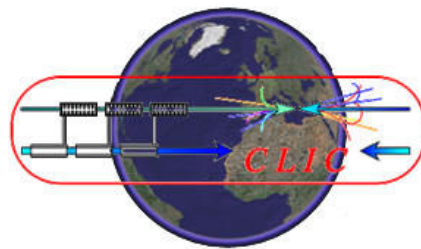
- Then we could build in also more easily
 - Thermionic emission
 - SEE

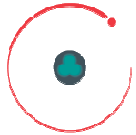
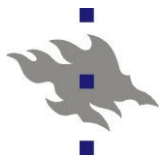
- Investigation of RF and other materials





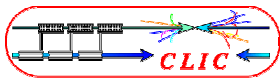
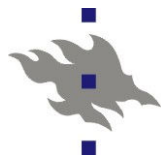
Thank you!





Back-up slides

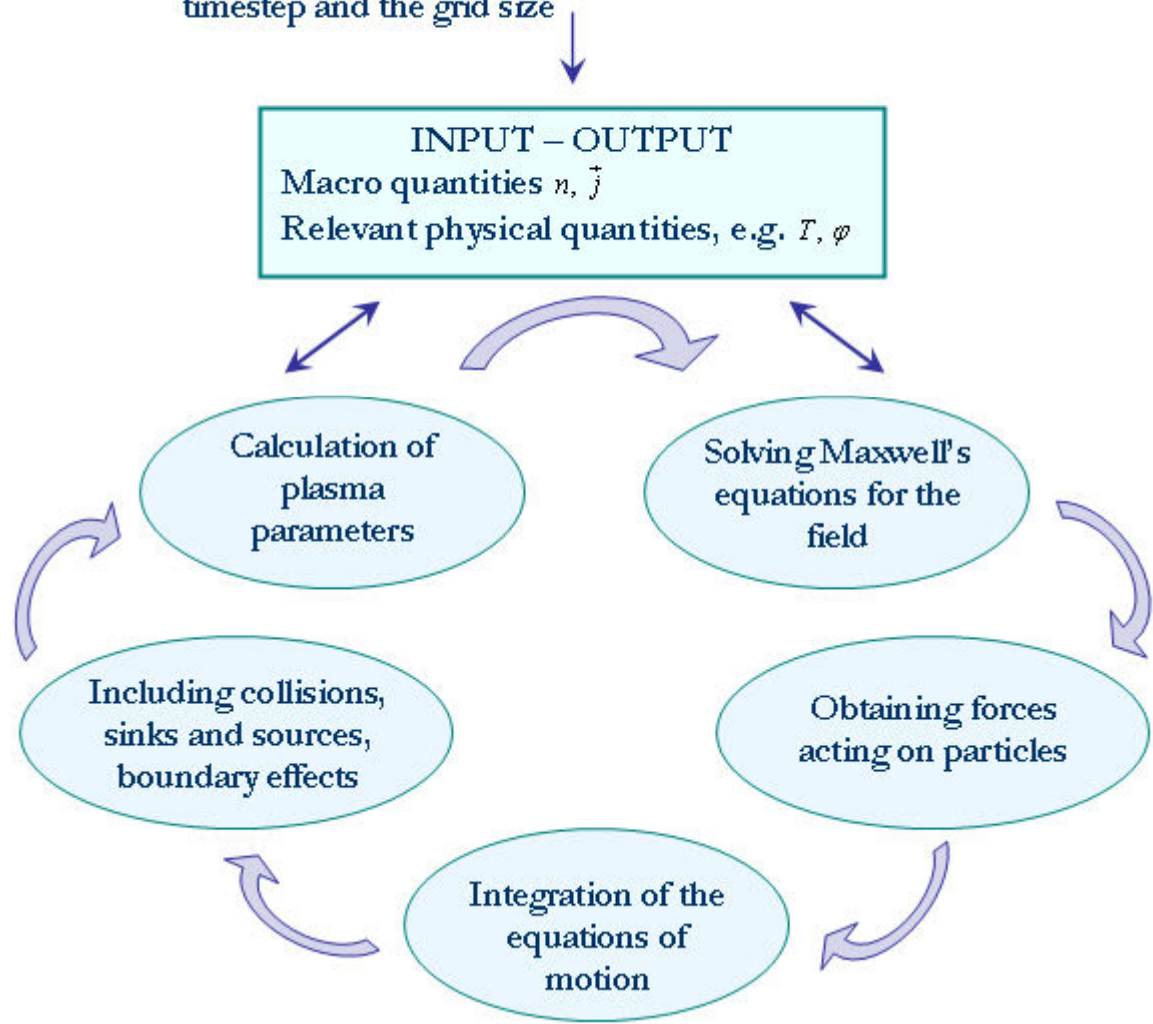


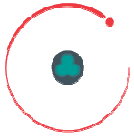
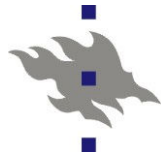


The Particle-in-Cell method

Setting up the simulation:

Defining superparticles, reference values, boundary conditions, timestep and the grid size

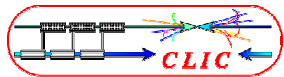
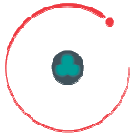
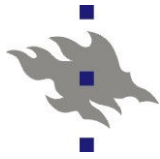




The Particle-in-Cell method

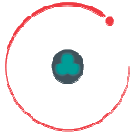
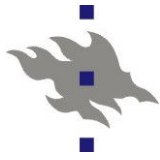
- Basic idea: simulate the time evolution of *macro quantities* instead of particle position and velocity (cf. MD method)
 - Fields and forces calculated on the grid
 - Superparticles
 - Restricted to certain regime of number density (ref. values)
 - Kinetic approach of plasma, but can be applied both for collisionless and collisional plasmas
- Application fields: solid state physics, quantum physics...
- Has become very popular in plasma physical applications
 - Esp. for modelling fusion reactor plasmas (sheath and edge)
- In our application:
 - 1D code (no side losses resolved)
 - Electrostatic: only Poisson's eq.

$$\frac{\varphi_{i+1} - 2\varphi_i + \varphi_{i-1}}{(\Delta x)^2} = -\frac{\rho}{\epsilon_0}.$$



I. Field emission, neutral evaporation

- Dynamic electron field emission current according to FN
 - Space charge corrections: Above $0.6 \text{ A}/\mu\text{m}^2$
 - PIC takes into account both the external+internal potential, so this is taken automatically care of
 - Thermionic emission can not be incorporated into 1D
- Beta expected to vary too; a dynamic beta, the "erosion" and the "melting" of the tip was implemented
 - We define the "melting current" j_{melt} as the threshold of electron emission current, which, if exceeded, sets $\beta=1$
- Neutral evaporation: An estimate was needed
 - Define the neutral evaporation to electron FE ratio $r_{\text{Cu/e}} = r_{\text{Cu/e}}(E, t, \dots)$ and approximate it with a constant
 - What is the possible range of $r_{\text{Cu/e}}$?



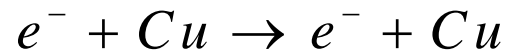
II. Collisions

With PIC we are limited in dynamic range/highest density that can be simulated; **restricts to plasma build-up phase**

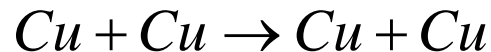
⇒ We treat only three species: e^- , Cu and Cu^+

- Coulomb collisions for all possible pairs,
 (e^-, e^-) , (e^-, Cu^+) and (Cu^+, Cu^+)

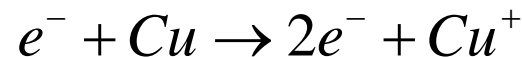
- Electron-neutral elastic collisions



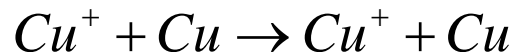
- Neutral-neutral elastic collisions

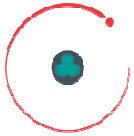
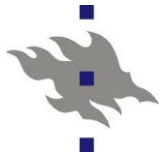


- Electron impact ionisation



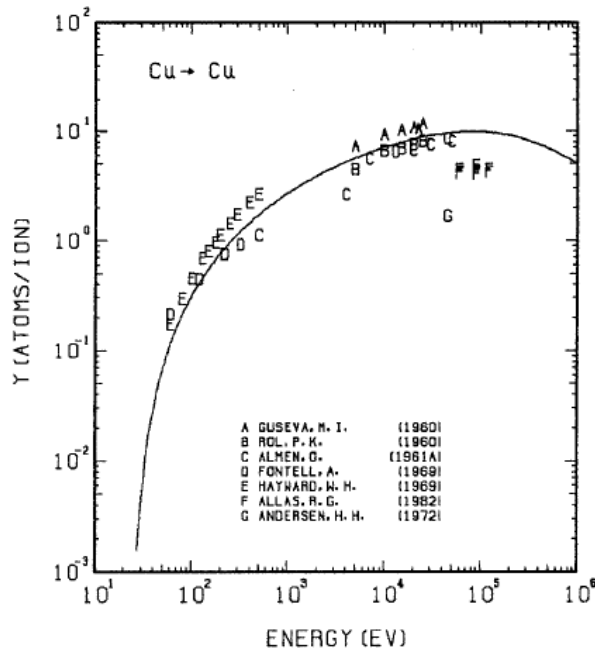
- Charge exchange and momentum transfer





III. Surface interaction model

- Neutrals and ions sputter neutrals at the walls
 - Energy dependent experimental yield
- Ions bombarding the cathode have in addition
 - An enhanced sputtering yield when the ion flux is above a given threshold (based on MD simulations)
 - A constant SEY = 0,5
 - SEE not yet included, but was implicitly parametrised through testing high SEY



[Yamamura & Tawara]