

On the Trilinear Higgs Self-Couplings in the complex NMSSM

Dao Thi Nhung

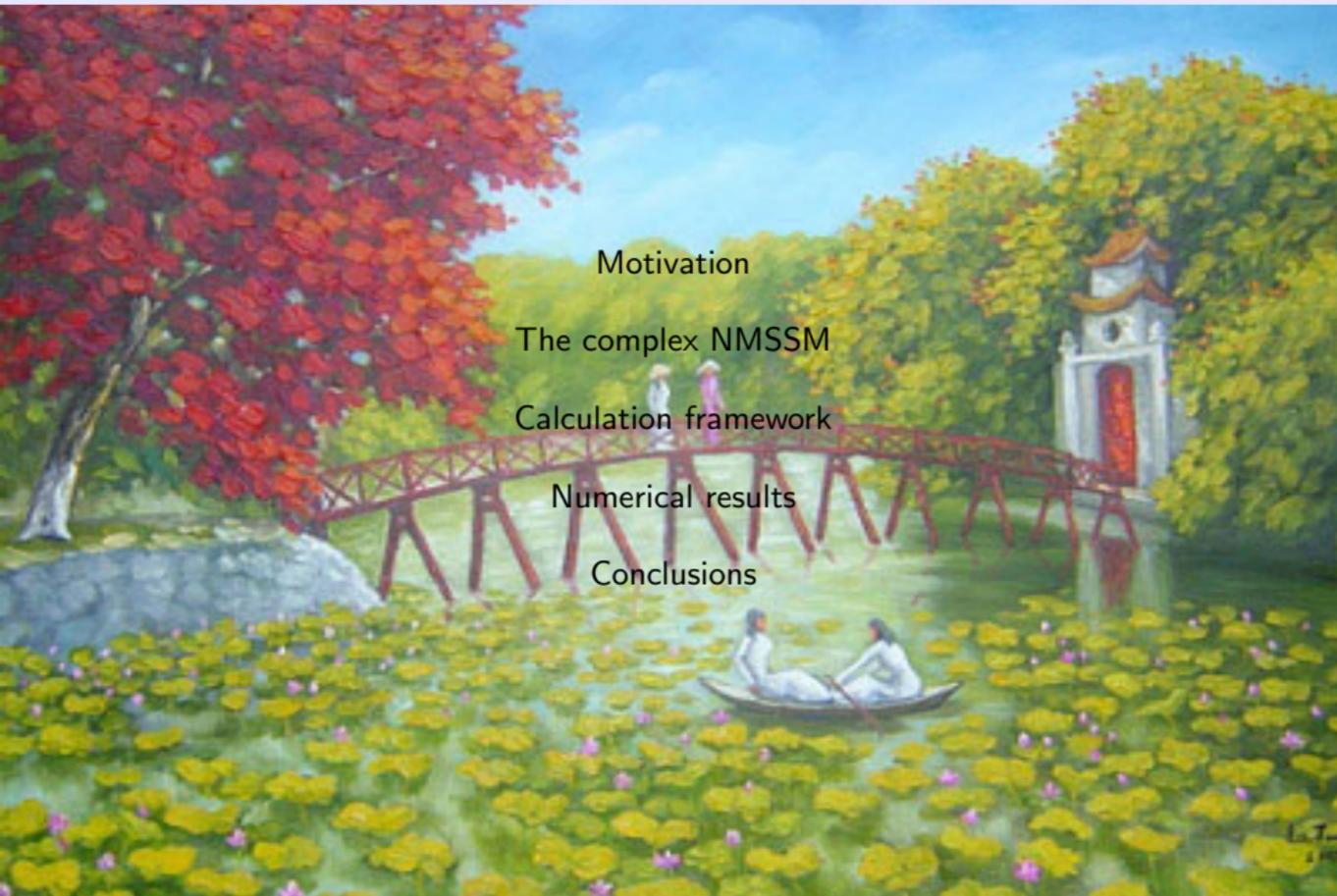
Institute for Interdisciplinary Research in Science and Education (IFIRSE),
ICISE, Quy Nhon, Vietnam

based on JHEP 1512 (2015) 034,
JHEP 1311 (2013) 181,

in coll. with Margarete Mühlleitner, Juraj Streicher, Kathrin Walz and Hanna Ziesche

XXXVI PIC, 13-17 Sept 2016, Quy Nhon, Viet Nam

Outline



Motivation

The complex NMSSM

Calculation framework

Numerical results

Conclusions

Why higher order trilinear Higgs couplings?

- Hints for the Higgs potential

- SM: $V_{SM} = -\mu^2 h^\dagger h + \lambda (h^\dagger h)^2$,

Higgs self couplings: $\lambda_{HHH} = \frac{3M_H^2}{v}$, $\lambda_{HHHH} = \frac{3M_H^2}{v^2}$

See also F. Brieuc's talk

- MSSM:

$$V_{MSSM} = (\mu^2 + m_{H_d}^2) H_{d,i}^* H_{d,i} + (\mu^2 + m_{H_u}^2) H_{u,i}^* H_{u,i} + \epsilon^{ij} (m_{12}^2 H_{u,i} H_{d,j} + \text{H.c.}) \\ + \frac{g_1^2 + g_2^2}{8} (H_{u,i}^* H_{u,i} - H_{d,i}^* H_{d,i})^2 + \frac{g_2^2}{2} |H_{u,i}^* H_{d,i}|^2$$

Decoupling limit: ($M_{H^\pm} \gg M_h$) then $\lambda_{hhh} \rightarrow \frac{3M_h^2}{v}$, h is the lightest Higgs boson, also at higher order

- NMSSM:

$$V_{NMSSM} = (|\lambda S|^2 + m_{H_d}^2) H_{d,i}^* H_{d,i} + (|\lambda S|^2 + m_{H_u}^2) H_{u,i}^* H_{u,i} + m_S^2 |S|^2 \\ + \frac{1}{8} (g_2^2 + g_1^2) (H_{d,i}^* H_{d,i} - H_{u,i}^* H_{u,i})^2 + \frac{1}{2} g_2^2 |H_{d,i}^* H_{u,i}|^2 \\ + |-\epsilon^{ij} \lambda H_{d,i} H_{u,j} + \kappa S^2|^2 + [-\epsilon^{ij} \lambda A_\lambda S H_{d,i} H_{u,j} + \frac{1}{3} \kappa A_\kappa S^3 + \text{H.c.}],$$

Decoupling limit: ($M_{H^\pm} \gg M_h$) then $\lambda_{hhh} \neq \frac{3M_h^2}{v}$, h is the SM-like Higgs boson depend on the mixing with singlet component.

- Match the accuracy of Higgs mass calculation
- Increase the accuracy of Higgs to Higgs decays

What have been done?

- **MSSM Trilinear Higgs Self-Couplings**
 - Full one-loop correction in Feynman diagram approach (complex MSSM)
[Williams, Rzehak, Weiglein]
 - Two-loop correction in effective potential approach (real MSSM)
[Brucherseifer, Gavin, Spira]
- **(complex) NMSSM Trilinear Higgs Self-Couplings**
 - Full one-loop correction in Feynman diagram approach
[Mühlleitner, DTN, Streicher, Walz]
 - Two-loop $\mathcal{O}(\alpha_t \alpha_s)$ correction in Feynman diagram approach
[Mühlleitner, DTN, Ziesche]

What is the complex NMSSM?

1 Superpotential

$$W_{NMSSM} = \epsilon_{ij} [y_e \hat{H}_d^i \hat{L}^j \hat{E}^c + y_d \hat{H}_d^i \hat{Q}^j \hat{D}^c - y_u \hat{H}_u^i \hat{Q}^j \hat{U}^c] - \epsilon_{ij} \lambda \hat{S} \hat{H}_d^i \hat{H}_u^j + \frac{1}{3} \kappa \hat{S}^3$$

2 Two complex Higgs doublets and one complex Higgs singlet

$$H_d = \begin{pmatrix} \frac{1}{\sqrt{2}}(v_d + h_d + ia_d) \\ h_d^- \end{pmatrix}, \quad H_u = e^{i\varphi_u} \begin{pmatrix} h_u^+ \\ \frac{1}{\sqrt{2}}(v_u + h_u + ia_u) \end{pmatrix}, \quad S = \frac{e^{i\varphi_s}}{\sqrt{2}}(v_s + h_s + ia_s).$$

3 Soft SUSY breaking terms

$$L_{soft} = \mathcal{L}_{soft, MSSM} - m_S^2 |S|^2 + (\epsilon_{ij} \lambda A_\lambda S H_d^i H_u^j - \frac{1}{3} \kappa A_\kappa S^3 + h.c.)$$

$\lambda, \kappa, A_\lambda, A_\kappa$ are in general complex.

4 dynamic μ term

$$\mu_{eff} H_d \cdot H_u \quad \text{with} \quad \mu_{eff} = \frac{\lambda v_s e^{i\varphi_s}}{\sqrt{2}}$$

5 CP-odd and CP-even Higgs bosons can already mix at tree-level.

$$(h_d, h_u, h_s, a_d, a_u, a_s) \rightarrow (h_1, h_2, h_3, h_4, h_5, G)$$

one CP-violating phase: $\phi_y = \phi_u - 2\phi_s - \phi_\kappa + \phi_\lambda$

The loop-corrected Higgs masses and trilinear couplings

- Loop-corrected Higgs mass matrices

$$M^2(p^2) = \begin{pmatrix} m_{h_1}^2 - \hat{\Sigma}_{h_1 h_1} & -\hat{\Sigma}_{h_1 h_2} & -\hat{\Sigma}_{h_1 h_3} & -\hat{\Sigma}_{h_1 h_4} & -\hat{\Sigma}_{h_1 h_5} \\ -\hat{\Sigma}_{h_2 h_1} & m_{h_2}^2 - \hat{\Sigma}_{h_2 h_2} & -\hat{\Sigma}_{h_2 h_3} & -\hat{\Sigma}_{h_2 h_4} & -\hat{\Sigma}_{h_2, h_5} \\ -\hat{\Sigma}_{h_3 h_1} & -\hat{\Sigma}_{h_3 h_2} & m_{h_3}^2 - \hat{\Sigma}_{h_3 h_3} & -\hat{\Sigma}_{h_3 h_4} & -\hat{\Sigma}_{h_3 h_5} \\ -\hat{\Sigma}_{h_4 h_1} & -\hat{\Sigma}_{h_4 h_2} & -\hat{\Sigma}_{h_4 h_3} & m_{h_4}^2 - \hat{\Sigma}_{h_4 h_4} & -\hat{\Sigma}_{h_4 h_5} \\ -\hat{\Sigma}_{h_5 h_1} & -\hat{\Sigma}_{h_5 h_2} & -\hat{\Sigma}_{h_5 h_3} & -\hat{\Sigma}_{h_5 h_4} & m_{h_5}^2 - \hat{\Sigma}_{h_5 h_5} \end{pmatrix},$$

$\hat{\Sigma}_{h_i h_j}(p^2)$ is renormalized self-energy of $h_i \rightarrow h_j$ transition

$$\hat{\Sigma}_{h_i h_j}(p^2) = \hat{\Sigma}_{h_i h_j}^{(\alpha)}(p^2) + \hat{\Sigma}_{h_i h_j}^{(\alpha_s \alpha_t)}(0)$$

loop-corrected Higgs mass eigenstates

$$(h_1, h_2, h_3, h_4, h_5) \rightarrow (H_1, H_2, H_3, H_4, H_5), \quad H_i = \mathbf{Z}_{ij} h_j$$

\mathbf{Z} : wave function renormalization factor

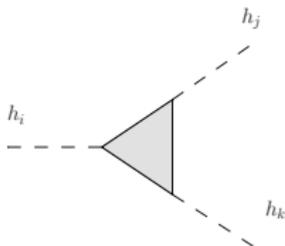
- Loop-corrected trilinear Higgs couplings

$$\begin{aligned} \mathcal{M}_{H_i \rightarrow H_j H_k} &= \sum_{i', j', k'=1}^5 \mathbf{z}_{ii'} \mathbf{z}_{jj'} \mathbf{z}_{kk'} (\lambda_{h_i' h_j' h_k'} + \Delta^{(\alpha)} \lambda_{h_i' h_j' h_k'} + \Delta^{(\alpha_s \alpha_t)} \lambda_{h_i' h_j' h_k'}) \\ &+ \Delta^{(\alpha)} M_{h_i \rightarrow h_j h_k}^{G, Z}. \end{aligned}$$

The full one-loop correction to the trilinear couplings

[Mühlleitner, DTN, Streicher, Walz]

$$\Delta^{(\alpha)} \lambda_{h_i h_j h_k}$$



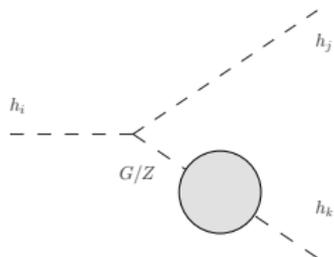
- $\mathcal{O}(\alpha)$ including full momentum dependence
- Renormalization scheme:

$$\underbrace{t_{h_d}, t_{h_u}, t_{h_s}, t_{a_d}, t_{a_s}, M_Z, M_W, M_{H^\pm}, e}_{\text{on-shell scheme}}, \underbrace{\tan \beta, \lambda, v_s, \kappa, A_{\kappa}}_{\overline{\text{DR}} \text{ scheme}}.$$

Higgs fields are renormalized in $\overline{\text{DR}}$ scheme
complex phases do not need to be renormalized

$$\Delta^{(\alpha)} M_{h_i \rightarrow h_j h_k}^{G,Z}$$

using tree-level masses to maintain gauge invariance

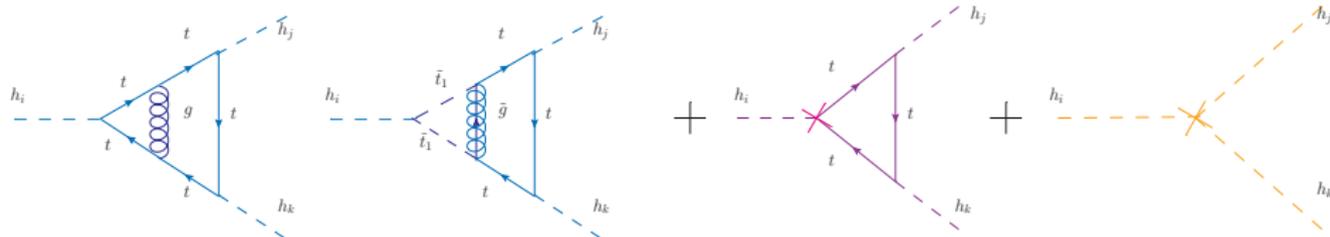


- Using FeynArt, FormCalc (D=4), LoopTools

The two-loop correction to the trilinear couplings

[Mühlleitner, DTN, Ziesche]

$$\Delta^{(\alpha_s \alpha_t)} \lambda_{h_i h_j h_k} = \Delta^{\Theta} \lambda_{h_i h_j h_k}^{\text{UR}} + \Delta^{\Theta} \lambda_{h_i h_j h_k}^{\text{CT1L}} + \Delta^{\Theta} \lambda_{h_i h_j h_k}^{\text{CT2L}}$$



- Zero external momentum approximation
- Dimensional reduction (DRED)
 - SUSY is preserved in two-loop Yukawa corrections to Higgs Masses
Stoekinger, Hollik' 2005
 - In practice: $\text{Tr}[1] = 4$, $[\gamma_5, \gamma^\mu] = 0$, $(g^4)_\nu^\mu (g^d)_\rho^\mu = (g^d)_\rho^\mu$
Loop momentum in d dimension \rightarrow tensor deduction in d dimension
- Tensor reduction: TARCER using Tarasov's Algorithm

$$c_1 A^d(m_i^2) A^d(m_j^2) + c_2 I(m_1^2, m_2^2, m_3^2)$$

A : one-loop tadpole integral, I : two-loop tadpole integral

- Using FeynArt, SARAH (model file), FeynCalc

Renormalization

- The following parameters need to be renormalized

$$\begin{aligned}t_\phi &\rightarrow t_\phi + \delta^{(0)}t_\phi + \delta^{(2)}t_\phi && \text{with } \phi = h_d, h_u, h_s, a_d, a_s, \\M_{H^\pm}^2 &\rightarrow M_{H^\pm}^2 + \delta^{(0)}M_{H^\pm}^2 + \delta^{(2)}M_{H^\pm}^2, \\v &\rightarrow v + \delta^{(0)}v + \delta^{(2)}v, \\\tan\beta &\rightarrow \tan\beta + \delta^{(0)}\tan\beta + \delta^{(2)}\tan\beta, \\|\lambda| &\rightarrow |\lambda| + \delta^{(0)}|\lambda| + \delta^{(2)}|\lambda|.\end{aligned}$$

- Chose the renormalization scheme as

$$\underbrace{t_{h_d}, t_{h_u}, t_{h_s}, t_{a_d}, t_{a_s}, M_{H^\pm}^2, v}_{\text{on-shell scheme}}, \underbrace{\tan\beta, |\lambda|}_{\overline{\text{DR}} \text{ scheme}}.$$

- Higgs wave function need also to be renormalized. Chose the $\overline{\text{DR}}$ scheme

$$H_u \rightarrow \left(1 + \frac{1}{2}\delta^{(0)}Z_{H_u} + \frac{1}{2}\delta^{(2)}Z_{H_u}\right) H_u, \quad \delta Z_{H_d} = \delta Z_{H_s} = 0$$

For $\overline{\text{DR}}$ scheme of top mass:

$$\delta^{(2)}Z_{H_u} = \frac{\alpha_s(m_t^2)^{\overline{\text{DR}}}}{8\pi^2 v^2 \sin^2\beta} \left(\frac{1}{\epsilon^2} - \frac{1}{\epsilon}\right), \quad \delta^{(2)}\tan\beta = \frac{1}{2}\tan\beta \delta^{(2)}Z_{H_u}, \quad \delta^{(2)}|\lambda| = \frac{-|\lambda|}{2}\delta^{(2)}Z_{H_u}$$

Renormalization of (s)top sector

The parameters need to be renormalized at $\mathcal{O}(\alpha_s)$

$$m_t, m_{\tilde{Q}_3}, m_{\tilde{t}_R} \text{ and } A_t$$

- On-shell renormalization scheme.

$$\delta\mathcal{X}^{\text{OS}} = \frac{1}{\epsilon} \delta\mathcal{X}_{\text{pole}} + \delta\mathcal{X}_{\text{fin}},$$

Note: the terms which are proportional to ϵ are not taken into account

- $\overline{\text{DR}}$ renormalization scheme

$$\delta\mathcal{X}^{\overline{\text{DR}}} = \frac{1}{\epsilon} \delta\mathcal{X}_{\text{pole}}.$$

- Translation of the parameters from two schemes if needed
Rough treatment

$$\begin{aligned} A_t^{(\text{OS})} &= A_t^{(\overline{\text{DR}})} - \delta A_t^{\text{fin}}, \\ (m_{\tilde{Q}_L}^2)^{(\text{OS})} &= (m_{\tilde{Q}_L}^2)^{(\overline{\text{DR}})} - \delta(m_{\tilde{Q}_L}^2)^{\text{fin}}, \\ (m_{\tilde{t}_R}^2)^{(\text{OS})} &= (m_{\tilde{t}_R}^2)^{(\overline{\text{DR}})} - \delta(m_{\tilde{t}_R}^2)^{\text{fin}}. \end{aligned}$$

Delicate treatment for top mass

$$\begin{aligned} M_t &\rightarrow m_t^{\overline{\text{MS}}}(M_t) \rightarrow m_t^{\overline{\text{MS}}}(M_{\text{SUSY}}) \\ &\rightarrow m_t^{\overline{\text{DR}},\text{SM}}(M_{\text{SUSY}}) \rightarrow m_t^{\overline{\text{DR}},\text{NMSSM}} \end{aligned}$$

- Check UV finite
- Two independent calculations are in agreement
- Compared to the real MSSM, found a good agreement (Thanks Spira for this comparison)

- Using **NMSSMCALC** to compute effective couplings of the Higgs bosons, normalized to the corresponding SM values, as well as the masses, the widths and the branching ratios of the Higgs bosons.
- We chose the scenarios which are accordance with the LHC Higgs data by using the programs `HiggsBounds` and `HiggsSignals`
- The resulting supersymmetric particle spectrum is in accordance with present LHC searches for SUSY particles

inbox (77) - dtrhung... XXXVI PHYSICS IN C... 1602.00881.pdf NMSSMProg How to take screens... +

lenovo screenshot

Most Visited • email • Dict • News • HEP • Bank • musics • conferences • sites • Fellowships • VSOP2016 • manuals • Websites_OR • useful tips • Quy Nhon

NMSSMCALC

Calculator of One-Loop and $O(\alpha_t \alpha_s)$ Two-Loop Higgs Mass Corrections and of Higgs Decay Widths in the CP-conserving and the CP-violating NMSSM

Now with the computation of the EDMs in the complex NMSSM

The program package NMSSMCALC calculates the one-loop and $O(\alpha_t \alpha_s)$ corrected Higgs boson masses and the Higgs decay widths and branching ratios within the CP-conserving and the CP-violating NMSSM.
The decay calculator is based on an extension of the program HDECAY 6.10 **now**.

Released by: Julien Baglio, Ramona Gröber, Margarete Mühleleitner, Dao Thi Nhung, Heidi Rzehak, Michael Spira, Juraj Streicher and Kathrin Walz
Program: NMSSMCALC version 2.00 **NEW!** *Computation of the EDMs in the complex NMSSM*

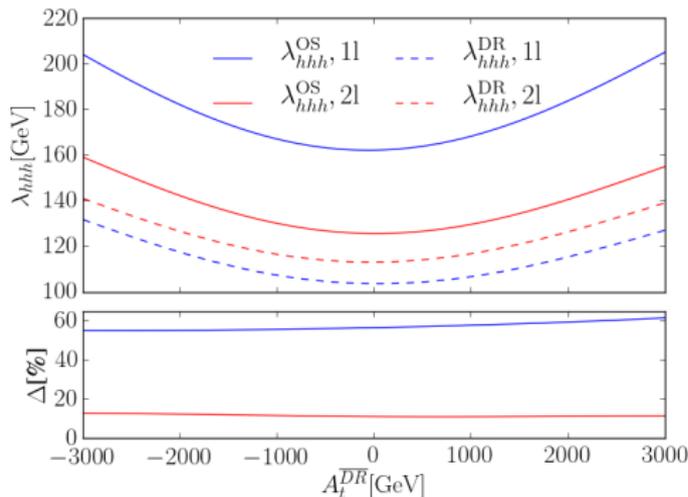
When you use this program, please cite the following references:

NMSSMCALC: [Julien Baglio, Ramona Gröber, Margarete Mühleleitner, Dao Thi Nhung, Heidi Rzehak, Michael Spira, Juraj Streicher and Kathrin Walz, in Comput. Phys. Commun. 185 \(2014\) 12](#)
One-Loop Masses: [K. Ender, T. Graf, M. Mühleleitner, H. Rzehak, in Phys. Rev. D85 \(2012\)075024](#)
[T. Graf, R. Gröber, M. Mühleleitner, H. Rzehak, K. Walz, in JHEP 1210 \(2012\) 122](#)
 $O(\alpha_t \alpha_s)$ Mass Corrections: [M. Mühleleitner, D.T. Nhung, H. Rzehak, K. Walz, in JHEP 1505 \(2015\) 128](#)
Computation of the EDMs in the cNMSSM: [S.F. King, M. Mühleleitner, R. Nevzorov, K. Walz, in arXiv:1508.03255](#)
HDECAY: [A. Djouadi, J. Kalinowski, M. Spira, Comput.Phys.Commun. 108 \(1998\) 56](#)
An update of HDECAY: [A. Djouadi, J. Kalinowski, Margarete Mühleleitner, M. Spira, in arXiv:1003.1643](#)

Informations on the Program:

- Short explanations on the program are given [here](#).
- To be advised about future updates or important modifications, send an E-mail to nmssmcalc@itp.kit.edu.

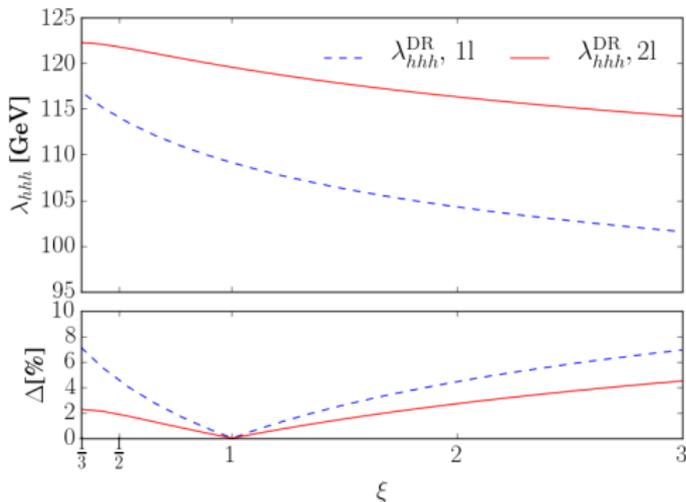
[Mühlleitner, DTN, Ziesche]



- h is dominated by h_u
- Including $\mathcal{O}(\alpha_t + \alpha_t \alpha_s)$, $p_2 = 0$
- In OS scheme: one-loop correction 140%, two-loop correction -24%
- In $\overline{\text{DR}}$ scheme: one-loop correction 74%, two-loop correction 9%
- Difference between $\mathcal{O}(\alpha_t)$ and $\mathcal{O}(\alpha)$ less than 4%
- $\Delta = \frac{|\lambda_{HHH}^{m_t(\overline{\text{DR}})} - \lambda_{HHH}^{m_t(\text{OS})}|}{\lambda_{HHH}^{m_t(\overline{\text{DR}})}}$
- Theoretical uncertainty decreases substantially

The SM-like trilinear coupling: $\overline{\text{DR}}$ scheme, scale uncertainty

[Mühlleitner, DTN, Ziesche]



- h is dominated by h_u
- Including $\mathcal{O}(\alpha_t + \alpha_t \alpha_s)$, $p_2 = 0$
- In $\overline{\text{DR}}$ scheme: one-loop correction 74%, two-loop correction 9%
- $\Delta = [\lambda_{hhh}(\mu_R) - \lambda_{hhh}(\mu_0)] / \lambda_{hhh}(\mu_0)$
 $\mu_0 = 2097 \text{ GeV}$
- $\overline{\text{DR}}$ parameters at different scales are estimated roughly by

$$p^{\text{OS}} + \delta p^{\text{OS}}(\mu) = p^{\overline{\text{DR}}}(\mu) + \delta p^{\overline{\text{DR}}}(\mu)$$

$$\begin{aligned} \tan \beta^{\text{pure}\overline{\text{DR}}}(\mu_1) - \tan \beta^{\text{pure}\overline{\text{DR}}}(\mu_2) &= a_1(m_t^{\text{OS}}) \ln \frac{\mu_1^2}{\mu_2^2} \\ &+ 2 a_2(m_t^{\text{OS}}, \alpha_s^{\overline{\text{DR}}}(\mu_1)) \ln \mu_1^2 - 2 a_2(m_t^{\text{OS}}, \alpha_s^{\overline{\text{DR}}}(\mu_2)) \ln \mu_2^2 \\ &+ 2 b_2(m_t^{\text{OS}}, \alpha_s^{\overline{\text{DR}}}(\mu_1)) \ln^2 \mu_1^2 - 2 b_2(m_t^{\text{OS}}, \alpha_s^{\overline{\text{DR}}}(\mu_2)) \ln^2 \mu_2^2. \end{aligned}$$

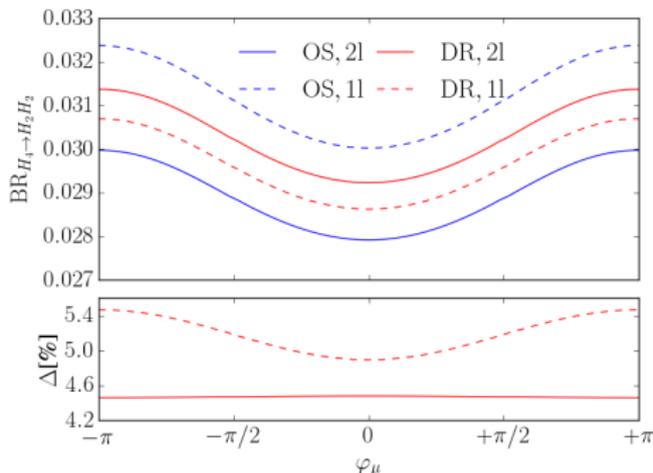
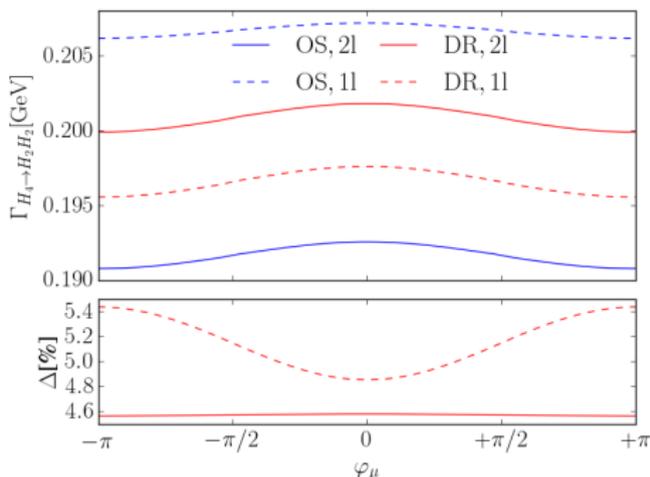
a_1 coef of 1-loop UV part. a_2, b_2 coef of double pole and single pole of 2-loop UV part

The Higgs to Higgs decays

[Mühlleitner, DTN, Ziesche]

$$H_4 \rightarrow H_2 H_2$$

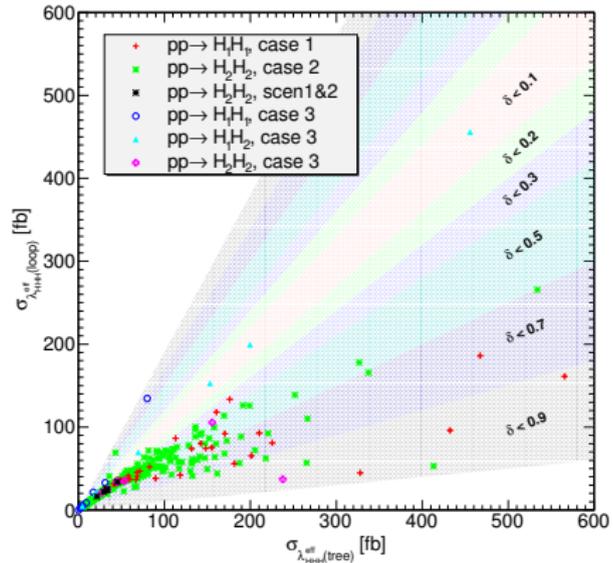
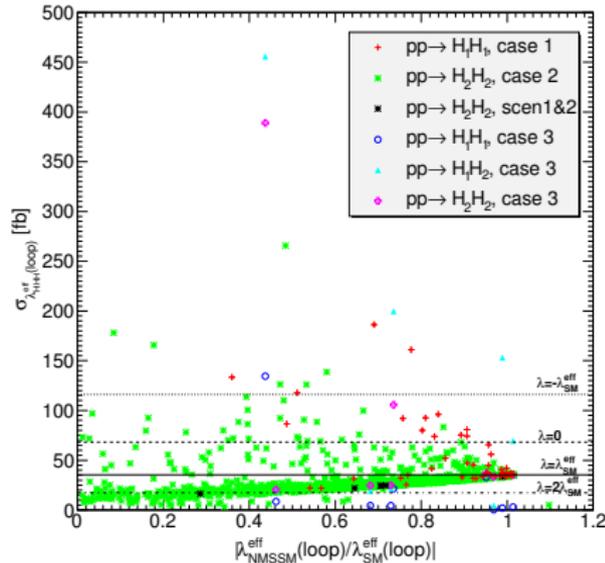
H_4 is mainly h_d , H_2 is dominated by h_u



- Include $\mathcal{O}(\alpha + \alpha_t \alpha_s)$, $\mathcal{O}(\alpha)$ with full momentum dependence
- tree-level CP violating phase $\phi_y = 0$
- OS: 21% at one-loop, -7% at two-loop
- $\overline{\text{DR}}$: 6.5% at one-loop, 2% at two-loop

Effect of trilinear Higgs self-coupling in Higgs boson pair production

[Mühlleitner, DTN, Streicher, Walz]



- Using HPAIR which include NLO QCD k-factor in the infinite top mass limit [Spira]
- Here Loop means using $\mathcal{O}(\alpha)$ effective Higgs Self-Coupling
- possible enhancement or detracton in the NMSSM

For the (c)NMSSM

- Full one-loop with full momentum dependence has been computed.
- Correction at one-loop level is of 75% in $\overline{\text{DR}}$ scheme, and 140% in OS scheme.
- The Leading two-loop correction of $\mathcal{O}(\alpha_t \alpha_s)$ in zero momentum approximation has been done
- The two-loop correction is of 9% in $\overline{\text{DR}}$ scheme and of -24% in OS scheme
- Theoretical uncertainty is significantly reduced for the SM-like trilinear couplings.
- Loop corrections to Higgs Self-coupling affect significantly the Higgs to Higgs bosons decay width and the total cross-section of Higgs pair production

For the (c)NMSSM

- Full one-loop with full momentum dependence has been computed.
- Correction at one-loop level is of 75% in $\overline{\text{DR}}$ scheme, and 140% in OS scheme.
- The Leading two-loop correction of $\mathcal{O}(\alpha_t\alpha_s)$ in zero momentum approximation has been done
- The two-loop correction is of 9% in $\overline{\text{DR}}$ scheme and of -24% in OS scheme
- Theoretical uncertainty is significantly reduced for the SM-like trilinear couplings.
- Loop corrections to Higgs Self-coupling affect significantly the Higgs to Higgs bosons decay width and the total cross-section of Higgs pair production

THANK YOU FOR YOUR ATTENTION

The one-loop top Yukawa ($\mathcal{O}(\alpha_t)$) correction to the trilinear couplings

$$\Delta^{(\alpha_t)} \lambda_{h_i, h_j, h_k}$$

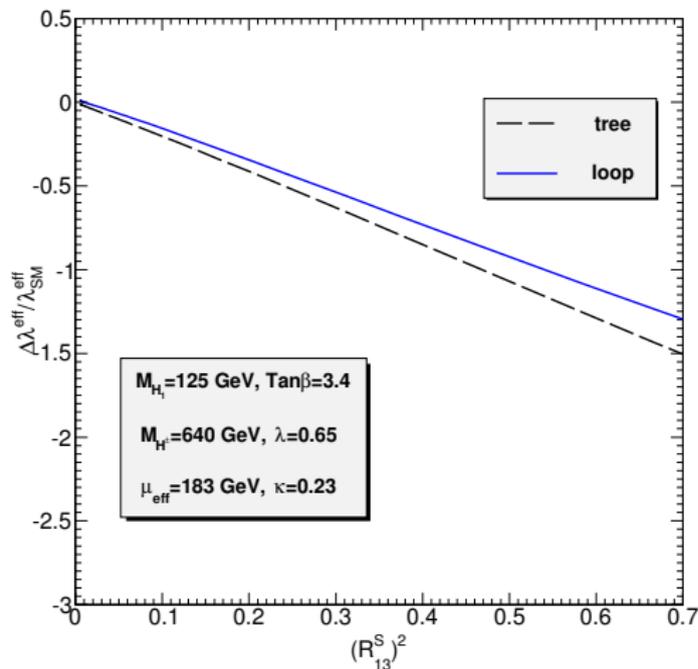
- Only one-loop diagrams with top/stop in loops
- Zero external momentum approximation is used
- This is the main contribution of the one-loop order, especially for SM-like trilinear coupling
- Useful computation to understand the two-loop $\mathcal{O}(\alpha_t \alpha_S)$ correction
- Renormalization of parameters (like the $\mathcal{O}(\alpha_t \alpha_S)$ case)

$$\underbrace{t_{h_d}, t_{h_u}, t_{h_s}, t_{a_d}, t_{a_s}, M_{H^\pm}^2, v}_{\text{on-shell scheme}}, \underbrace{\tan \beta, |\lambda|}_{\overline{\text{DR}} \text{ scheme}} .$$

- Do not include $\Delta^{(\alpha)} M_{h_i \rightarrow h_j h_k}^{G,Z}$, violate gauge invariance with zero momentum approximation

Effective coupling at decoupling Limit

[Mühlleitner, DTN, Streicher, Walz]



- Decoupling Limit: $M_{H^\pm} \gg M_h$
- Effective: $p^2 = 0$
- Loop: $\mathcal{O}(\alpha)$
- R_{13}^S : mixing with singlet component

Parameters - Scenario 1

The SM input parameters

$$\alpha(M_Z) = 1/128.962, \quad \alpha_s^{\overline{MS}}(M_Z) = 0.1184, \quad M_Z = 91.1876 \text{ GeV},$$
$$M_W = 80.385 \text{ GeV}, \quad m_t = 173.5 \text{ GeV}, \quad m_b^{\overline{MS}}(m_b^{\overline{MS}}) = 4.18 \text{ GeV}.$$

Using $\alpha_s^{\overline{DR}}$

Scenario 1:

$$m_{\tilde{u}_R, \tilde{c}_R} = m_{\tilde{d}_R, \tilde{s}_R} = m_{\tilde{Q}_{1,2}} = m_{\tilde{L}_{1,2}} = m_{\tilde{e}_R, \tilde{\mu}_R} = 3 \text{ TeV}, \quad m_{\tilde{t}_R} = 1909 \text{ GeV},$$
$$m_{\tilde{Q}_3} = 2764 \text{ GeV}, \quad m_{\tilde{b}_R} = 1108 \text{ GeV}, \quad m_{\tilde{L}_3} = 472 \text{ GeV}, \quad m_{\tilde{\tau}_R} = 1855 \text{ GeV},$$
$$|A_{u,c,t}| = 1283 \text{ GeV}, \quad |A_{d,s,b}| = 1020 \text{ GeV}, \quad |A_{e,\mu,\tau}| = 751 \text{ GeV},$$
$$|M_1| = 908 \text{ GeV}, \quad |M_2| = 237 \text{ GeV}, \quad |M_3| = 1966 \text{ GeV},$$
$$\varphi_{A_{d,s,b}} = \varphi_{A_{e,\mu,\tau}} = \varphi_{A_{u,c,t}} = \pi, \quad \varphi_{M_1} = \varphi_{M_2} = \varphi_{M_3} = 0.$$

$$|\lambda| = 0.374, \quad |\kappa| = 0.162, \quad |A_\kappa| = 178 \text{ GeV}, \quad |\mu_{\text{eff}}| = 184 \text{ GeV},$$
$$\varphi_\lambda = \varphi_\kappa = \varphi_{\mu_{\text{eff}}} = \varphi_u = 0, \quad \varphi_{A_\kappa} = \pi, \quad \tan \beta = 7.52, \quad M_{H^\pm} = 1491 \text{ GeV}.$$

Renormalization scale = SUSY scale

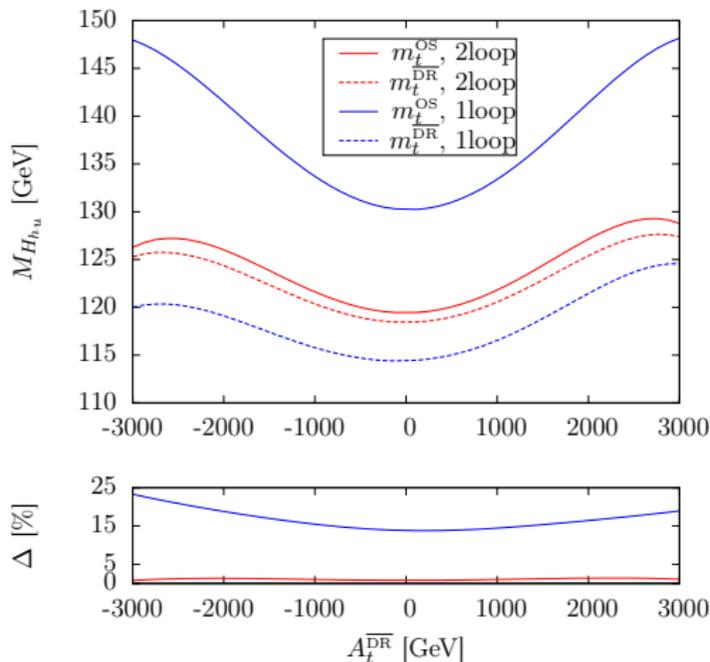
$$M_s = \sqrt{m_{\tilde{Q}_3} m_{\tilde{t}_R}}.$$

Senario 1: Mass spectrum

OS	H_1	H_2	H_3	H_4	H_5
mass tree [GeV] main component	71.14 h_u	117.49 h_s	211.12 a_s	1491 a	1492 h_d
mass one-loop [GeV] main component	98.65 h_s	139.17 h_u	217.27 a_s	1490 a	1491 h_d
mass two-loop [GeV] main component	94.68 h_s	125.06 h_u	217.32 a_s	1490 a	1491 h_d
DR	H_1	H_2	H_3	H_4	H_5
mass tree [GeV] main component	71.14 h_u	117.49 h_s	211.12 a_s	1491 a	1492 h_d
mass one-loop [GeV] main component	91.60 h_s	120.00 h_u	217.36 a_s	1491 a	1491 h_d
mass two-loop [GeV] main component	94.41 h_s	124.24 h_u	217.33 a_s	1490 a	1491 h_d

Higgs Masses: OS scheme vs $\overline{\text{DR}}$ scheme

[Mühlleitner, DTN, Rzehak, Walz]



- $\hat{\Sigma}_{h_i h_j}(p^2) = \hat{\Sigma}_{h_i h_j}^{(\alpha)}(p^2) + \hat{\Sigma}_{h_i h_j}^{(\alpha_S \alpha_t)}(0)$
- $\Delta = |M_{H_{h_u}}^{m_t(\overline{\text{DR}})} - M_{H_{h_u}}^{m_t(\text{OS})}| / M_{H_{h_u}}^{m_t(\overline{\text{DR}})}$

Scenario 2

$$\begin{aligned}
 m_{\tilde{t}_R} &= 1170 \text{ GeV}, m_{\tilde{Q}_3} = 1336 \text{ GeV}, m_{\tilde{b}_R} = 1029 \text{ GeV}, m_{\tilde{L}_3} = 2465 \text{ GeV}, m_{\tilde{\tau}_R} = 301 \text{ GeV} \\
 |A_{u,c,t}| &= 1824 \text{ GeV}, |A_{d,s,b}| = 1539 \text{ GeV}, |A_{e,\mu,\tau}| = 1503 \text{ GeV}, |M_1| = 862.4 \text{ GeV}, |\kappa| = 0.208 \\
 |M_2| &= 201.5 \text{ GeV}, |M_3| = 2285 \text{ GeV}, |\lambda| = 0.629, |A_\kappa| = 179.7 \text{ GeV}, |\mu_{\text{eff}}| = 173.7 \text{ GeV}, \\
 \tan \beta &= 4.02, M_{H^\pm} = 788 \text{ GeV}, \varphi_{A_{u,c,t}} = \varphi_{M_1} = \varphi_{M_2} = \varphi_{M_3} = \varphi_\lambda = \varphi_{\mu_{\text{eff}}} = \varphi_u = \varphi_{A_\kappa} = 0, \\
 \varphi_{A_{d,s,b}} &= \varphi_{A_{e,\mu,\tau}} = \varphi_\kappa = \pi.
 \end{aligned}$$

OS	H_1	H_2	H_3	H_4	H_5
mass tree [GeV]	79.15	103.55	146.78	796.62	803.86
main component	h_s	h_u	a_s	h_d	a
mass one-loop [GeV]	103.45	129.15	139.83	796.53	802.94
main component	h_s	a_s	h_u	h_d	a
mass two-loop [GeV]	102.99	126.09	128.94	796.45	803.07
main component	h_s	h_u	a_s	h_d	a
DR	H_1	H_2	H_3	H_4	H_5
mass tree [GeV]	79.15	103.55	146.78	796.62	803.86
main component	h_s	h_u	a_s	h_d	a
mass one-loop [GeV]	102.80	120.52	128.80	796.36	803.09
main component	h_s	h_u	a_s	h_d	a
mass two-loop [GeV]	103.09	124.55	128.91	796.36	803.03
main component	h_s	h_u	a_s	h_d	a

Scenario 3

$$\begin{aligned}
 m_{\tilde{t}_R} &= 1940 \text{ GeV}, \quad m_{\tilde{Q}_3} = 2480 \text{ GeV}, \quad m_{\tilde{b}_R} = 1979 \text{ GeV}, \quad m_{\tilde{L}_3} = 2667 \text{ GeV}, \quad m_{\tilde{\tau}_R} = 1689 \text{ GeV}, \\
 |A_{u,c,t}| &= 1192 \text{ GeV}, \quad |A_{d,s,b}| = 685 \text{ GeV}, \quad |A_{e,\mu,\tau}| = 778 \text{ GeV}, \quad |M_1| = 517 \text{ GeV}, \quad |M_2| = 239 \text{ GeV}, \\
 |M_3| &= 1544 \text{ GeV}, \quad |\lambda| = 0.267, \quad |\kappa| = 0.539, \quad |A_\kappa| = 810 \text{ GeV}, \quad |\mu_{\text{eff}}| = 104 \text{ GeV}, \quad \tan \beta = 8.97, \\
 M_{H^\pm} &= 613 \text{ GeV}, \quad \varphi_{A_{d,s,b}} = \varphi_{A_{e,\mu,\tau}} = \varphi_{M_1} = \varphi_{M_2} = \varphi_{M_3} = \varphi_\lambda = \varphi_\kappa = \varphi_{\mu_{\text{eff}}} = \varphi_u = 0, \\
 \varphi_{A_{u,c,t}} &= \varphi_{A_\kappa} = \pi.
 \end{aligned}$$

OS	H_1	H_2	H_3	H_4	H_5
mass tree [GeV]	49.17	99.83	609.21	611.77	715.92
main component	h_s	h_u	a	h_d	a_s
mass one-loop [GeV]	87.36	139.10	608.71	611.37	694.73
main component	h_s	h_u	a	h_d	a_s
mass two-loop [GeV]	83.66	124.95	608.73	611.37	694.76
main component	h_s	h_u	a	h_d	a_s
DR	H_1	H_2	H_3	H_4	H_5
mass tree [GeV]	49.17	99.83	609.21	611.77	715.92
main component	h_s	h_u	a	h_d	a_s
mass one-loop [GeV]	80.66	119.68	608.72	611.37	694.79
main component	h_s	h_u	a	h_d	a_s
mass two-loop [GeV]	83.03	124.34	608.71	611.36	694.78
main component	h_s	h_u	a	h_d	a_s