

Amplitude detuning measurements

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- 1 Detuning with amplitude
- 2 Measurement with kicked beams
- 3 Measurement with driven oscillations
- 4 Alternative measurements
- 5 Conclusions

Many thanks to the Optics Measurement and Corrections team



■ Detuning with amplitude

→ dependence of tune on action ($J_{x,y}$) or CS-invariant ($\epsilon_{x,y} = 2J_{x,y}$)

$$\rightarrow N[\sigma_{\text{nominal}}] = \sqrt{\frac{2J}{\epsilon_{\text{nominal}}}}$$

$$Q_z(\epsilon_x, \epsilon_y) = Q_{z0} + \left(\frac{\partial Q_z}{\partial \epsilon_x} \epsilon_x + \frac{\partial Q_z}{\partial \epsilon_y} \epsilon_y \right) + \\ + \frac{1}{2!} \left(\frac{\partial^2 Q_z}{\partial \epsilon_x^2} \epsilon_x^2 + 2 \frac{\partial^2 Q_z}{\partial \epsilon_x \partial \epsilon_y} \epsilon_x \epsilon_y + \frac{\partial^2 Q_z}{\partial \epsilon_y^2} \epsilon_y^2 \right) + \dots$$

Order	Source (3 = sextupole)
$\frac{\partial Q}{\partial \epsilon}$	$(K_3)^2, K_4$
$\frac{\partial^2 Q}{\partial \epsilon^2}$	$(K_3)^4, (K_3)^2 K_4, (K_4)^2, K_3 K_5, K_6$

■ $\frac{\partial Q_x}{\partial \epsilon_x}$ → “Direct term”

■ $\frac{\partial Q_y}{\partial \epsilon_x}$ → “Cross term”

Amplitude detuning from an octupole

$$H_n = \frac{1}{B\rho} \operatorname{Re} \left[\frac{1}{n} [B_n(s) + iA_n(s)] (x + iy)^n \right]$$

$$\text{Normal octupole} \rightarrow H_4 = \frac{1}{4!} K_4 L (x^4 - 6x^2y^2 + y^4)$$

In action-angle coordinates $(x, y = \sqrt{2J_{x,y}\beta_{x,y}} \cos \phi_{x,y})$

$$H_4 = \frac{1}{4!} K_4 L (4J_x^2 \beta_x^2 \cos^4 \phi_x - 24J_x J_y \cos^2 \phi_x \cos^2 \phi_y + 4J_y^2 \beta_y^2 \cos^4 \phi_y)$$

$$Q_x = \frac{1}{2\pi} \frac{\partial \langle H \rangle}{\partial J_x} = \frac{1}{16\pi} K_4 L (J_x \beta_x^2 - 2J_y \beta_x \beta_y)$$

$$\frac{\partial Q_x}{\partial \epsilon_x} = \frac{1}{32\pi} \beta_x^2 K_4 L \quad \frac{\partial Q_x}{\partial \epsilon_y} = -\frac{1}{16\pi} \beta_x \beta_y K_4 L \quad \frac{\partial Q_y}{\partial \epsilon_y} = \frac{1}{32\pi} \beta_y^2 K_4 L$$

Equivalence of detuning cross terms

1st order detuning:

$$\frac{\partial Q_x}{\partial J_y} = \frac{1}{2\pi} \frac{\partial^2 \langle H \rangle}{\partial J_y \partial J_x} = \frac{\partial Q_y}{\partial J_x}$$

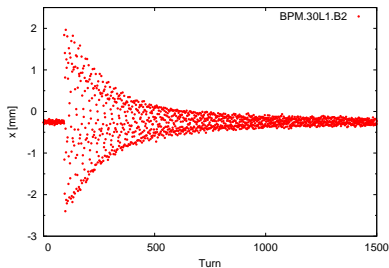
2nd order detuning:

$$\frac{\partial^2 Q_y}{\partial J_x^2} = \frac{1}{2\pi} \frac{\partial^3 \langle H \rangle}{\partial J_x^2 \partial J_y} = \frac{\partial^2 Q_x}{\partial J_x \partial J_y}$$

$$\frac{\partial^2 Q_x}{\partial J_y^2} = \frac{1}{2\pi} \frac{\partial^3 \langle H \rangle}{\partial J_x \partial J_y^2} = \frac{\partial^2 Q_y}{\partial J_x \partial J_y}$$

- Terms like $\frac{\partial^2}{\partial J_x \partial J_y}$ measured directly with diagonal kicks in H-V plane
- but from cross term equivalence actually determine all second order terms with pure H or V measurements
- Cross term equivalence gives good sanity check for data/fit quality

- **Traditional detuning measurement uses single kicks**
- Dipole kicker ramped up/down within single turn
- observe free betatron oscillations with turn-by-turn BPM data



- Oscillations **do not** decay
- Oscillations decohere
- **Measurement is destructive**
- Fresh beam for every kick

- BPM data post processed by **Singular Value Decomposition (SVD)**

R.Tomás & R.Calaga, *Statistical analysis of RHIC beam position monitors performance*,
Phys.Rev.ST.AB,7,042801

- Identifies malfunctioning BPMs
- Removes uncorrelated noise from BPM signals

Action determined from mean peak-to-peak TbT data over BPMs

$$2J_{x,y} = \frac{\sum_{BPMs} \left(\frac{1}{2} \frac{Peak-to-Peak}{\beta_{x,y}} \right)^2}{N_{BPMs}}$$

- Various sources of uncertainty:
beta-beat, coupling, BPM-scaling, BPM-nonlinearity, phase-space distortion from resonances

Tune determined via spectral analysis of TbT data

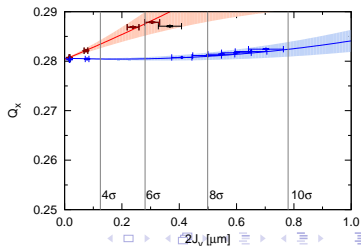
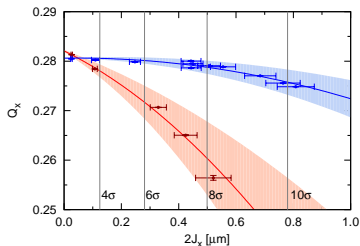
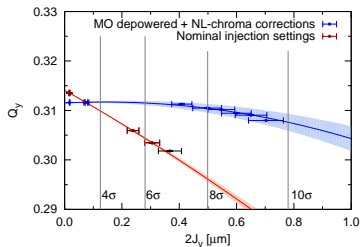
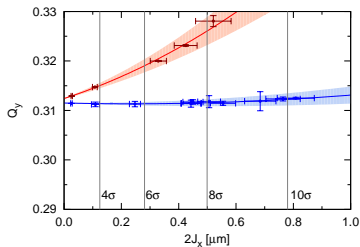
- Spectral analysis done via SUSSIX (interpolated FFT)
R.Bartolini & F.Schmidt, CERN SL/Note 98-017(AP),
'SUSSIX: a computer code for frequency analysis of non-linear betatron motion'
- Decoherence limits number of turns available for spectral analysis

- Beams kicked to varying amplitudes for several angles in H-V plane
(at least pure kicks in H and V)
- Limited by kicker strength, or machine / dynamic apertures

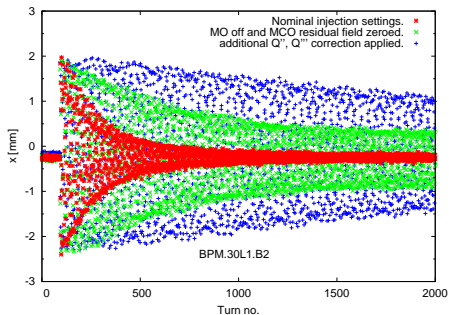
Traditional detuning measurements performed at injection in 2012

E.H.Maclean, R.Tomás, F.Schmidt, T.H.B.Persson. Phys.Rev.ST.AB,**18**,081002(2014)
Measurement of nonlinear observables in the Large Hadron Collider using kicked beams

- Nominal injection optics (Landau octupoles on)
- Landau octupoles off + beam-based correction of Q'' & Q'''

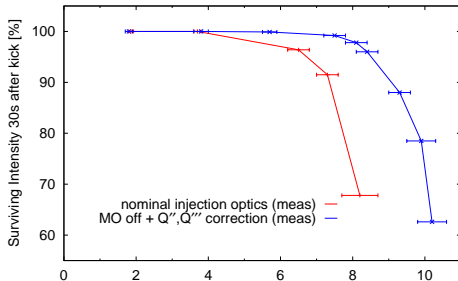


Beam-based correction also reduced decoherence and increased DA

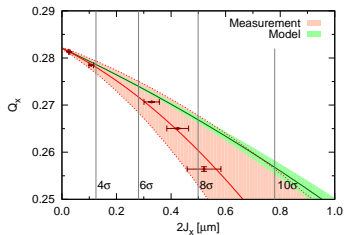
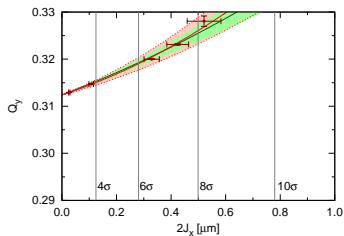


■ Decoherence used as online check b_4 corrections worked

- In general Q'' correction wont correct detuning & DA
- Implies local correction



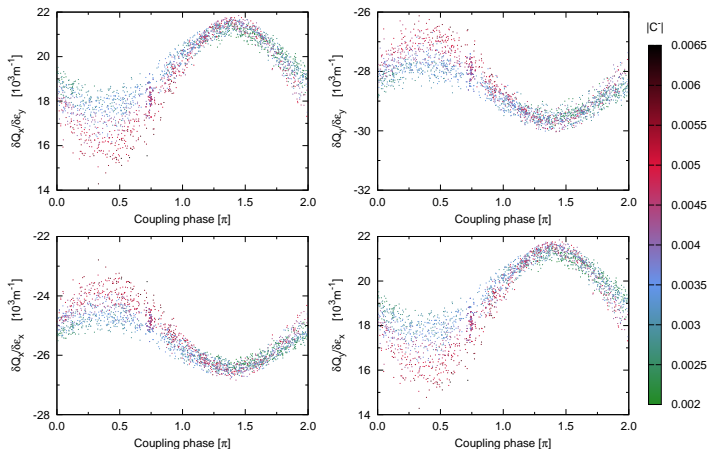
Comparison to LHC model



	[unit]	Meas' \pm err	Model \pm err
$\frac{\partial Q_x}{\partial \epsilon_x}$	$[10^3 \text{m}^{-1}]$	-29 7	-27.0 0.8
$\frac{\partial Q_y}{\partial \epsilon_x}$		19 3	21 2
$\frac{\partial Q_x}{\partial \epsilon_y}$		24 4	21 2
$\frac{\partial Q_y}{\partial \epsilon_y}$		-32.8 0.4	-30.5 0.9
$\frac{\partial^2 Q_x}{\partial \epsilon_x^2}$	$[10^9 \text{m}^{-2}]$	-60 30	-14 4
$\frac{\partial^2 Q_y}{\partial \epsilon_x^2}$		34 10	18 9

- Good agreement of 1st order detuning
- Qualitatively similar 2nd order

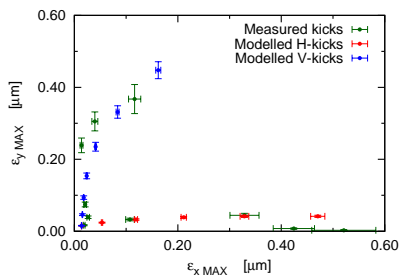
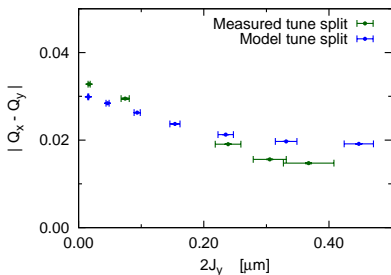
When comparing model and measurement must account for linear coupling



- In simulation linear coupling significantly affects the detuning... even far from the coupling resonance
- Not only δQ_{min} that's important: also phase of RDT
- Best option is a good correction at start of measurement

Coupling effect may also impinge upon detuning measurements

- Detuning with J_y moved tunes together
- Tune separation saturates
- Kicks at large J_y couple significantly into H-plane
- Observed in real LHC and simulation



- Behaviour associated with transverse planes becoming strongly coupled
- **Still very far from measured** $|C^-| = 0.0036$
- Implies existence of amplitude dependent δQ_{min}
- **Coupling stopband will distort detuning...**

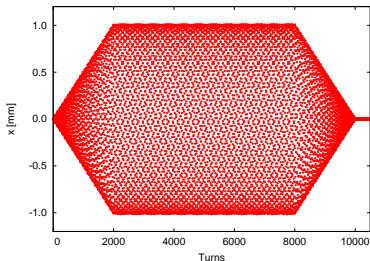
...but may have a nonlinear contribution

Single kick detuning measurements not possible at top energy in LHC

- Require new fill for every kick (destructive measurement)
- Machine protection

LHC is equipped with AC-dipole kickers

- Sinusoidally driven dipole kicker
- Driving frequency close (but not on!) natural tunes generates large response with little power, even at high energy
- If ramped up/down adiabatically kicks are non-destructive
- Routinely used for linear optics measurements throughout cycle



AC-dipole provides a tool to measure amplitude detuning at high energy

Actually rely on non-perfect adiabaticity to excite natural tune lines in spectrum

AC-dipole modifies solutions to equations of motion:

$$x(s) = \sqrt{2J_x\beta_x(s)} \cos \phi_x(s) \quad \rightarrow \quad x_D(s) = \sqrt{2J_x\beta_x(s)} \cos \phi_x(s) + \sqrt{2A_x\beta'_x(s)} \cos \phi_D(s)$$

- $A, \phi_D(s)$ are action angle variables of the driven oscillation
- β' is beta-function modified by the AC-dipole ($\beta' \approx \beta$)

Alters action-angle Hamiltonian & Q, eg octupole tune shift:

$$Q_x = \frac{1}{2\pi} \frac{\partial \langle H \rangle}{\partial J_x} = \frac{1}{16\pi} K_4 L (J_x \beta_x^2 - 2J_y \beta_x \beta_y)$$

$$\rightarrow \frac{1}{16\pi} K_4 L (J_x \beta_x^2 + 2A_x \beta'_x \beta_x - 2J_y \beta_x \beta_y)$$

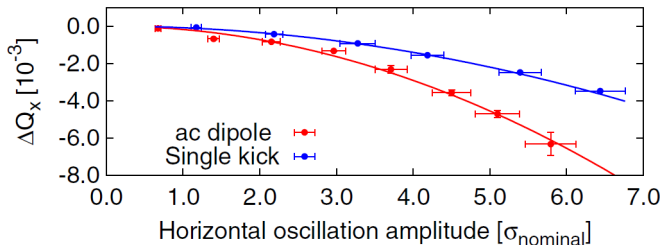
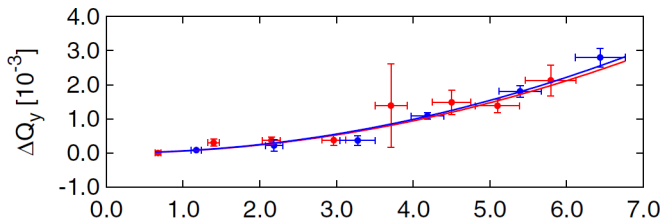
- $J_x \ll A_x \rightarrow$ direct detuning $2\times$ expectation for free oscillations
- Detuning cross term unaffected
- Similar result for Q_y

In general:

Direct detuning terms from n^{th} order are $\frac{n}{2}$ larger when measured with AC-dipole than free oscillations. Cross terms are unaffected.

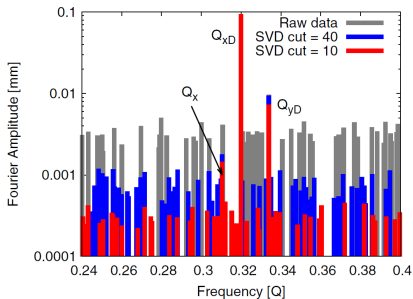
S.White, R.Tomás, E.H.Maclean, 'Direct amplitude detuning measurement with ac dipole', Phys.Rev.ST.AB,16,071002(2013)

Effect of AC-dipole on observed detuning was verified experimentally at injection



Measurement of natural tune variation with AC-dipole action is more challenging than with free oscillations

FFT of driven oscillations for different SVD cuts

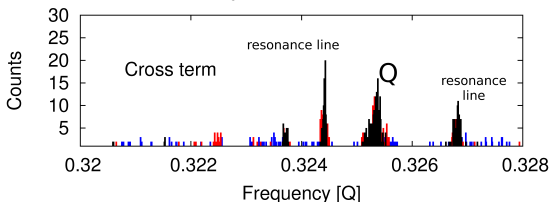


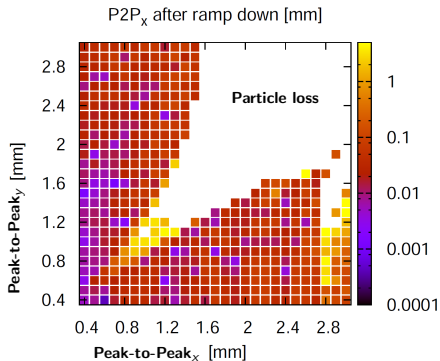
- Natural tune not a strong signal
- Need aggressive SVD cleaning
- Additional resonances

$$aQ_x + bQ_y = z$$

$$\rightarrow aQ_x + bQ_y + pQ_{ACx} + qQ_{ACy} = z$$

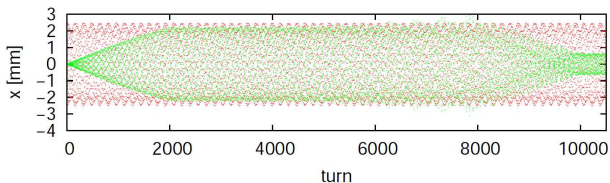
“Tune” identified by SUSSIX in ~ 500 LHC BPMs





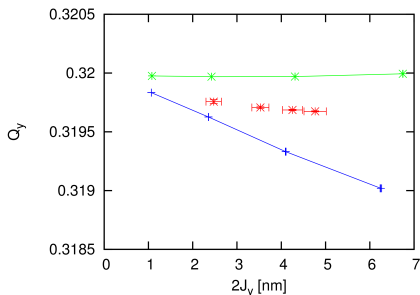
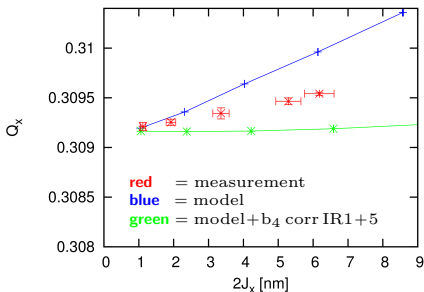
Watch out for dynamic aperture!

- In general DA smaller than un-driven motion
- Kicking to DA will cause blow up and particle loss
- DA depends on Q_{AC}



AC-dipole detuning measurements performed successfully at top energy

- Measured @ 6.5 TeV, $\beta^* = 0.4$ m during 2015 MD
- Comparison of to MAD-X tracking simulations, including AC-dipole



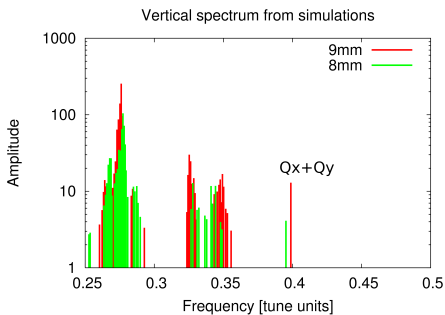
Amplitude detuning measurements by A.Langner, comparison to simulation by S.Monig

- For 0.4 m detuning dominated by b_4 errors in IR1+IR5
(negligible contribution of arcs, which dominate Q'')
- Implies $\sim \frac{1}{2}$ expected b_4 of IR1 + IR5

Amplitude detuning is not the only probe available for NL-errors

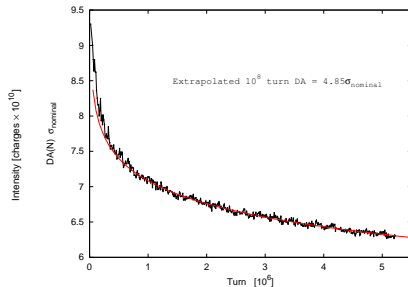
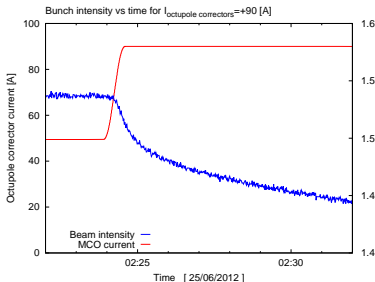
- With AC-dipole detuning measurements gain spectral info for free
- Studied for Wire Excitation experiments in SPS

U.Dorda et. al. *Wire excitation experiments in the CERN SPS*, EPAC'08



- Sextupole coupling line $Q_x + Q_y$
- Predict change in amplitude for change in beam-wire separation
- Qualitative agreement with observed spectrum

- **Traditional DA measurement with single kicks used at injection**
 - Not viable at top energy
- **Alternative method is blow up emittance with transverse damper**
 - study long term DA via intensity loss & scaling laws
 - Demonstrated at injection, viable at top energy
 - being considered for optimization of NL-correctors in IR



- possibilities for short-term DA measurement with AC-dipole

 Amp'
detuning

 Single
kicks

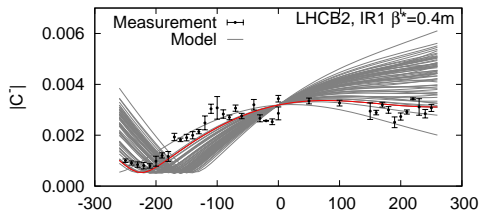
 AC-
dipole

 Other
method

Summary

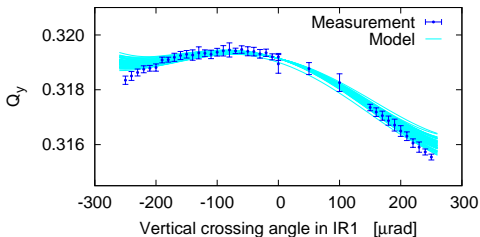
■ Currently study NL-errors in low- β^* IRs via feed-down

E.H.Maclean, R.Tomás, M.Giovanozzi, T.H.B.Persson. Accepted to Phys.Rev.ST.AB
 First measurement and correction of nonlinear errors in the experimental insertions of the CERN LHC



e.g. IR1 @ 0.4 m, 4 TeV

■ $b_3 + a_4$ feed-down to $|C^-|$



■ $a_3 + b_4$ feed-down to Q_y

■ Potentially quite useful in conjunction with other observables

Conclusions

- **Traditional detuning measurement** → **single kicks**
- 1st & 2nd order detuning measured @ 450 GeV with single kicks
- Traditional measurement not viable at LHC top energy

- **AC-dipole measurement possible at LHC top energy**
- Theory predicts driven oscillations have different detuning
- Verified experimentally @ 450 GeV
- AC-dipole measurement tougher than single-kick
- Demonstrated at top energy → now routine

- **Various additional methods also available**
- long-term DA (ADT), short-term DA (AC-dipole), feed-down, spectra

Amp'
detuningSingle
kicksAC-
dipoleOther
method

Summary